



天 地

CP Violation for the Heaven and the Earth

— Sighting the 4th Generation ?

George W.S. Hou (侯維恕)

National Taiwan University

February 13, 2009, WG4/LHC2FC @ CERN



臺灣大學



National Taiwan University





Can *all this* be understood from my vantage?



Capital Reef National Park
(c) Wally Pacholka



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CP Violation for the Heaven and the Earth

— Sighting the 4th Generation ?

George W.S. Hou (侯維恕)

National Taiwan University

February 6, 2009, Wine & Cheese Seminar, Fermilab

see: <http://theory.fnal.gov/jetp/>



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Outline



I. Intro

II. $\Delta A_{K\pi}$ & Nondecoupling: t' in Z Penguin

1st Prediction for $\sin 2\Phi_{B_s}$

III. Δm_{B_s} & $\mathcal{B}(B \rightarrow X_s \ell^+ \ell^-)$: 2nd Prediction for $\sin 2\Phi_{B_s}$

[CKM Consistency & Implications]

IV. A $10^{13}+$ Enhancement of CPV for BAU?

[V. (In)Direct Sighting: Tevatron vs LHC]

VI. Discussion/Conclusion: Know in 3-5 Years

WSH, Nagashima, Soddu,
PRL'05; PRD'05; PRD'07; PRD'08
Belle, Nature, 452, 20 (2008)
WSH, arXiv:0803.1234 [hep/ph]



I. Intro



4th Generation Still?



- N_ν counting? 4th “neutral lepton” heavy
Massive neutrinos call for new Physics

- Disfavored by **EW Precision** (see e.g. J. Erler hep-ph/0604035; PDG06)

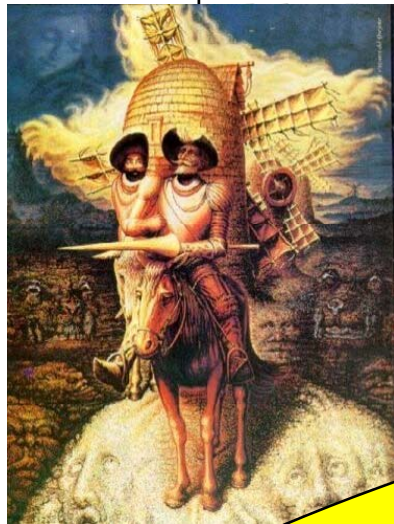
An extra generation of ordinary fermions is excluded at the 99.999% CL on the basis of the S parameter alone, corresponding to $N_F = 2.81 \pm 0.24$ for the number of fermion families. This result assumes that there are no new contributions to T or U and that the fermion families are degenerate. In principle this restriction can be relaxed.

July 14, 2006 10:28

10. Electroweak model constraints on new physics 37

As well, since $T > 0$, a non-degenerate extra family. However, currently favor strengthening the exclusion limits. A more detailed analysis of the S parameter (or the extra down-type quark) is close to zero. This can drive S to small or even negative values but at the same time large contributions to T . These results are in agreement with a fit to the S parameter of light neutrinos, $N_\nu = 2.986 \pm 0.007$ (which favors a larger value for N_ν) and $S = 0.1231 \pm 0.0020$ mainly from R_ℓ and τ_τ). However, the S parameter fits are still valid even for a very heavy fourth family neutrino.

Friday the 13th (before Valentines!!) Punishment!





4th Generation Still?



- N_ν counting? 4th “neutral lepton” heavy
Massive neutrinos call for new Physics

~ Tao Han’s talk

- Disfavored by **EW Precision** (see e.g. J. Erler hep-ph/0604035; PDG06)

An extra generation of ordinary fermions is excluded at the 99.999% CL on the basis of the S parameter alone, corresponding to $N_F = 2.81 \pm 0.24$ for the number of families. This assumes that there are no new contributions to T or U and therefore that the extra families are degenerate. In principle this restriction can be relaxed by allowing

July 14, 2006 10:37

10. Electroweak model and constraints on new physics 37

As well, since $T > 0$ is expected from a non-degenerate extra family. However, current data currently favor $T < 0$, thus strengthening the exclusion limits. A more detailed analysis is required if the extra neutrino (or the extra down-type quark) is close to the Dirac mass limit [208]. This can drive S to small or even negative values but at the expense of too-large contributions to T . These results are in agreement with a fit to the number of light neutrinos, $N_\nu = 2.986 \pm 0.007$ (which favors a larger value for $\alpha_s(M_Z) = 0.1231 \pm 0.0020$ mainly from R_ℓ and τ_τ). However, the S parameter fits are valid even for a very heavy fourth family neutrino.

- 4th generation **not** in such great conflict with EWPrT
Kribs, Plehn, Spannowsky, Tait, PRD’07

Emi Kou’s talk



In fact, well presented !



General: Emi Kou (plenary)
[Heavy Lepton: Tao Han (plenary); could be related

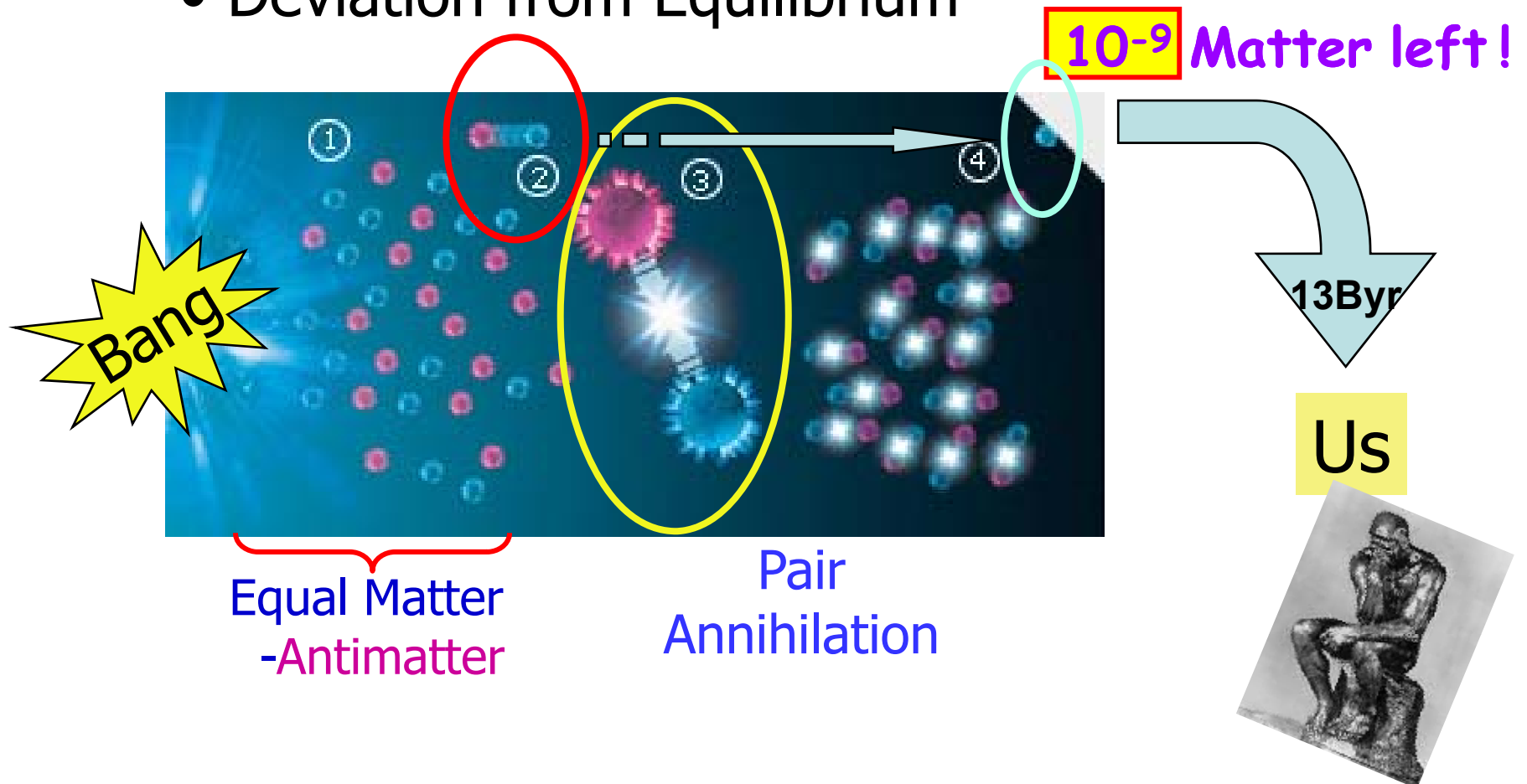
ATLAS: Saleh Sultansoy (also on Higgs in WG1)
CMS: Kai-Feng Chen

CPV: this talk

(1967)

CPV & BAU (& U): The Sakharov View

- *Baryon Number Violation*
- *CP Violation*
- Deviation from Equilibrium



Sakharov Stimulated by ... Discovery of CP Violation

- Phys. Rev. Lett. 13, 138 (1964)



27 JULY 1964

VOLUME 13, NUMBER 4

PHYSICAL REVIEW LETTERS

EVIDENCE FOR THE 2π DECAY OF THE K_2^0 MESON*†

J. H. Christenson, J. W. Cronin,† V. L. Fitch,‡ and R. Turlay§
Princeton University, Princeton, New Jersey

(Received 10 July 1964)

1980 Nobel

This Letter reports the results of experimental studies designed to search for the 2π decay of the K_2^0 meson. Several previous experiments have served^{1,2} to set an upper limit of 1/300 for the fraction of K_2^0 's which decay into two charged pions. The present experiment, using spark chamber techniques, proposed to extend this limit.

In this measurement, K_2^0 mesons were produced at the Brookhaven AGS in an internal Be target bombarded by 30-BeV protons. A neutral beam was defined at 30 degrees relative to the circulating protons by a $1\frac{1}{2}$ -in. \times $1\frac{1}{2}$ -in. \times 48-in. collimator at an average distance of 14.5 ft. from

The analysis program computed the vector momentum of each charged particle observed in the decay and the invariant mass, m^* , assuming each charged particle had the mass of the charged pion. In this detector the K_{e3} decay leads to a distribution in m^* ranging from 280 MeV to ~536 MeV; the $K_{\mu 3}$, from 280 to ~516; and the $K_{\pi 3}$, from 280 to 363 MeV. We emphasize that m^* equal to the K^0 mass is not a preferred result when the three-body decays are analyzed in this way. In addition, the vector sum of the two momenta and the angle, θ , between it and the direction of the K_2^0 beam were determined. This



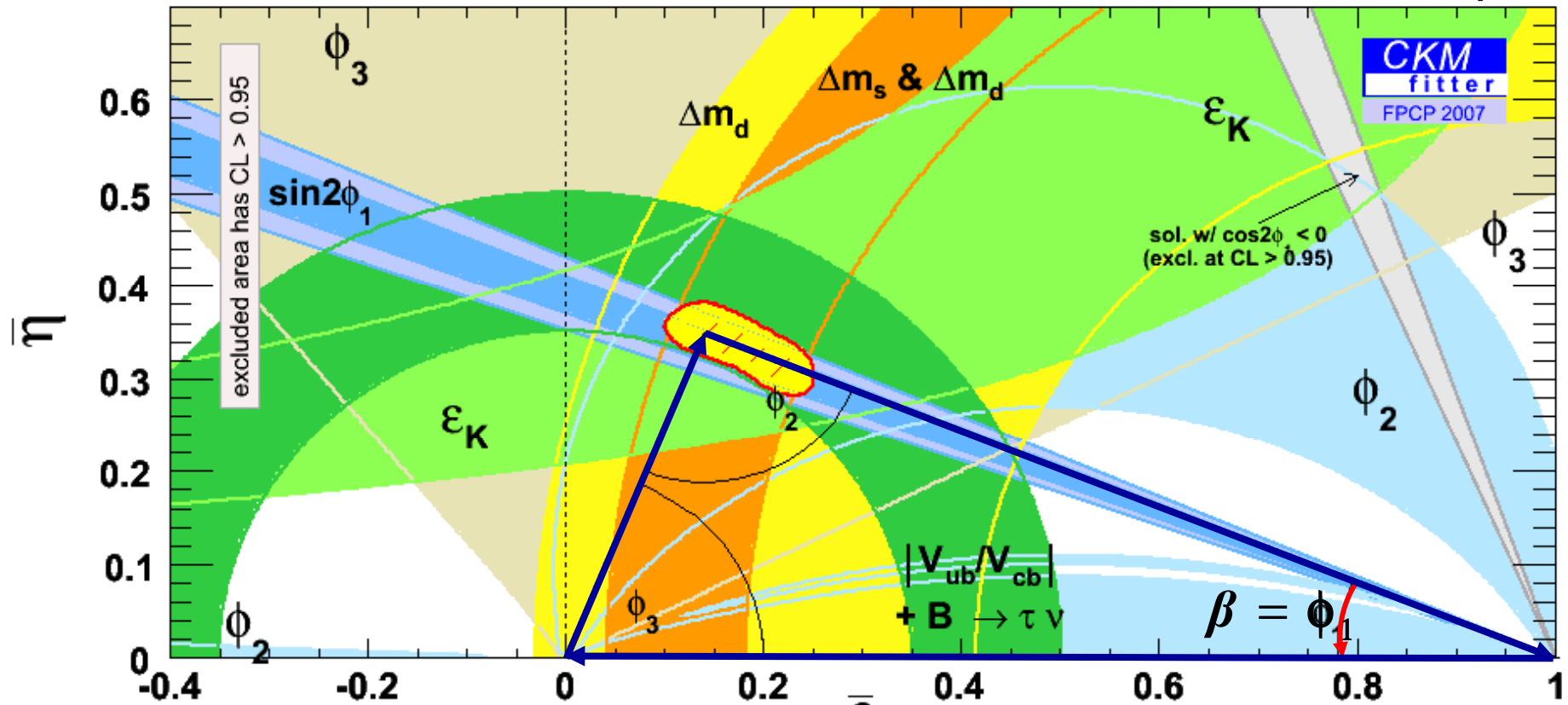
2×10^{-3} : Too Small for Sakharov !



KM CPV Confirmed ~ 2001



the MOMA plot



“Nontrivial”

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



The Nobel Prize in Physics 2008



"for the discovery of the mechanism of spontaneous broken symmetry in subatomic physics"



Photo: University of Chicago

Yoichiro Nambu

🏆 1/2 of the prize

USA

Enrico Fermi Institute,
University of Chicago
Chicago, IL, USA

b. 1921
(in Tokyo, Japan)

"for the **discovery** of the **origin** of the broken symmetry which **predicts** the existence of at least three families of quarks in nature"



Photo: KEK

Makoto Kobayashi

🏆 1/4 of the prize

Japan

High Energy Accelerator
Research Organization
(KEK)
Tsukuba, Japan

b. 1944



Photo: Kyoto University

Toshihide Maskawa

🏆 1/4 of the prize

Japan

Kyoto Sangyo University;
Yukawa Institute for
Theoretical Physics (YITP),
Kyoto University
Kyoto, Japan

b. 1940

CP Violation in SM



naturenews
7 October 2008



The Belle detector in Japan helped to confirm the symmetry breaking effects predicted by theoretical physicists.

KEK

B Factories (BaBar & Belle)



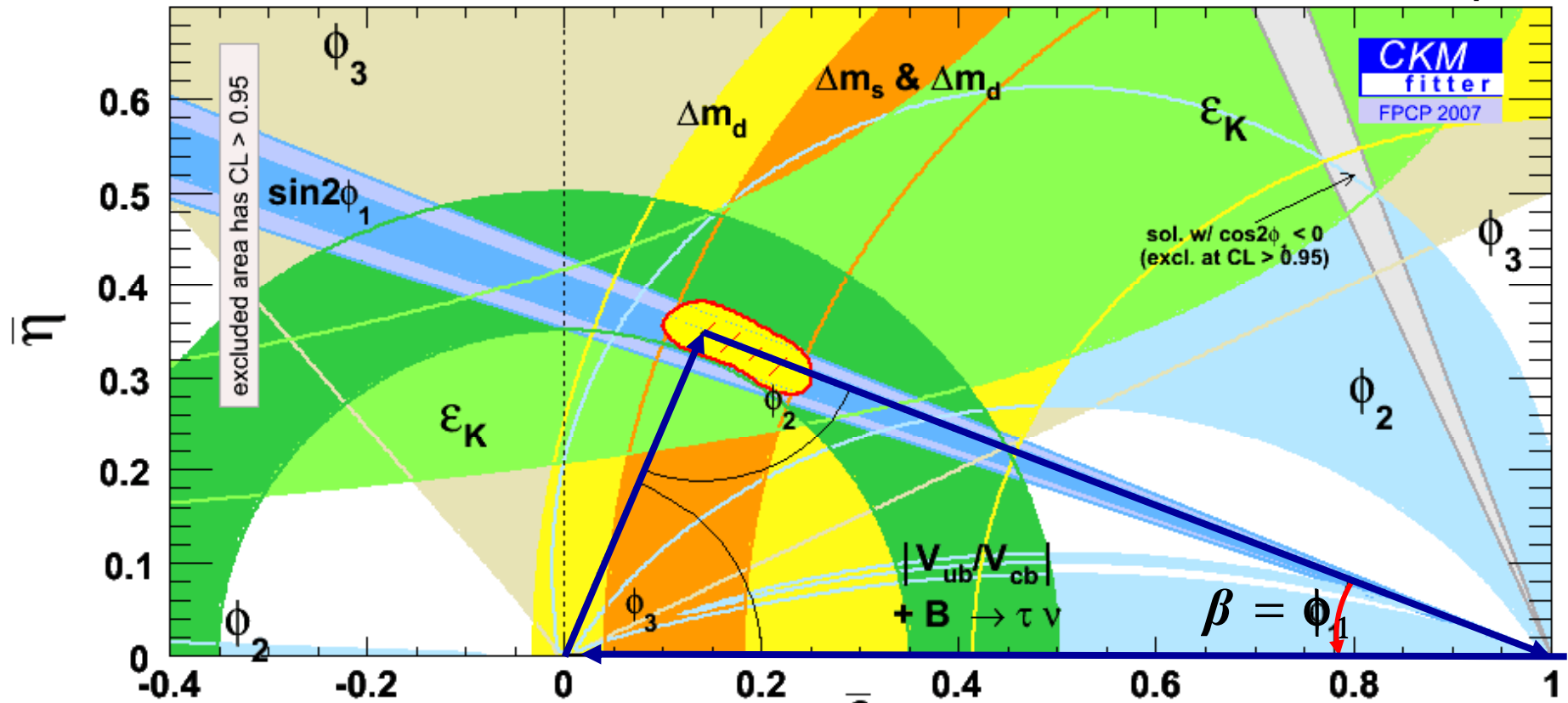


$b \rightarrow d$ transitions consistent with SM



$b \rightarrow s$: the Current Frontier

the MOMA plot



“Nontrivial”

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



A Real Hint , ... or Not !?



II. $\Delta A_{K\pi}$ & Nondecoupling: t' in Z Penguin

1st Prediction for $\sin 2\Phi_{B_s}$



Belle 2008 Nature: Simple Bean Count

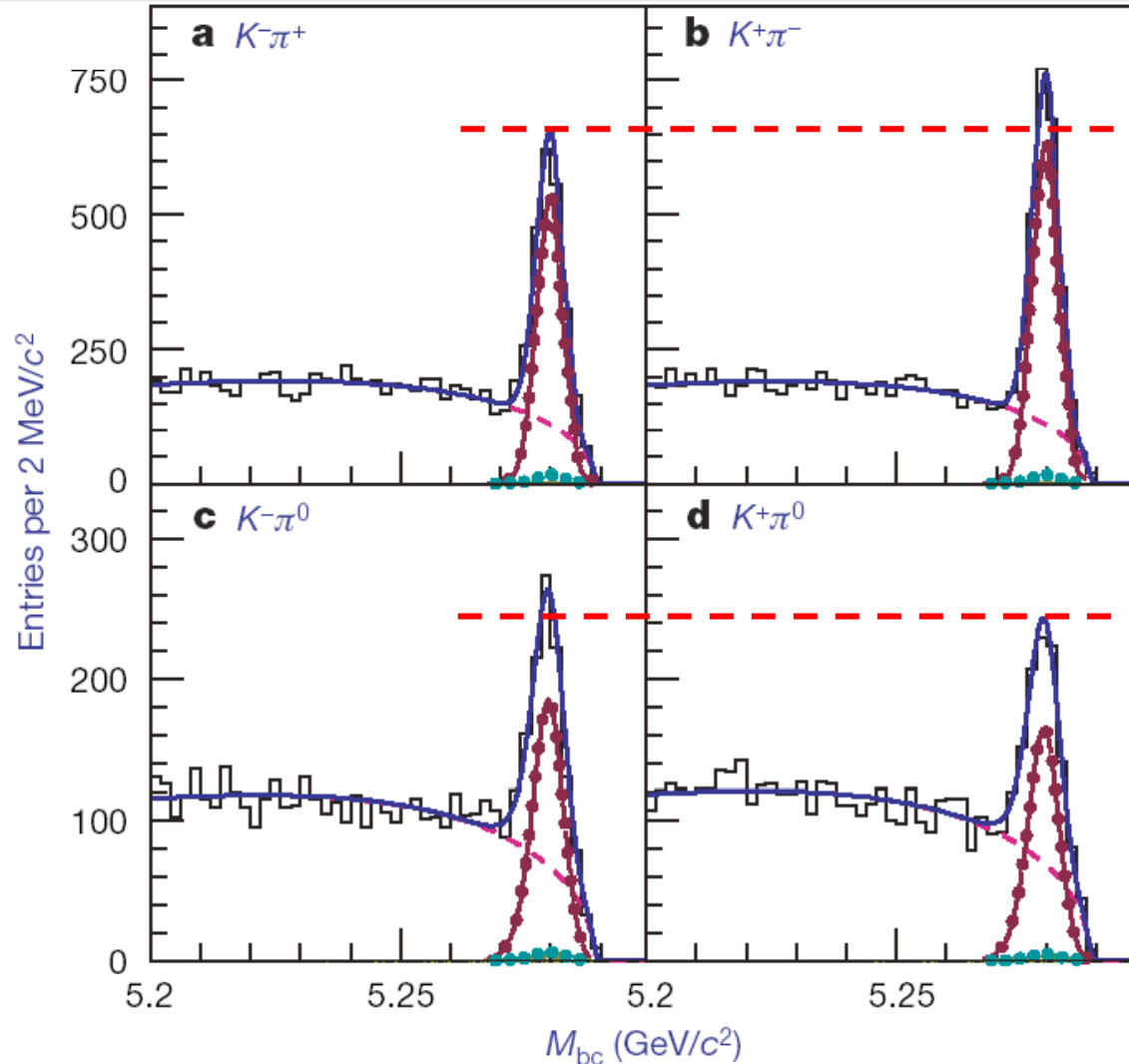


$$\Delta A_{K\pi} = A_{K^+\pi^0} - A_{K^+\pi^-} = +0.164 \pm 0.037 \quad 4.4\sigma$$

$$+0.07 \pm 0.03 \quad \text{vs} \quad -0.094 \pm 0.020$$

NATURE | Vol 452 | 20 March 2008

LETTERS



$b \rightarrow s$ CPV

Not Predicted !

Difference
Is
Large !

And Established
Belle + BaBar (+ CDF)

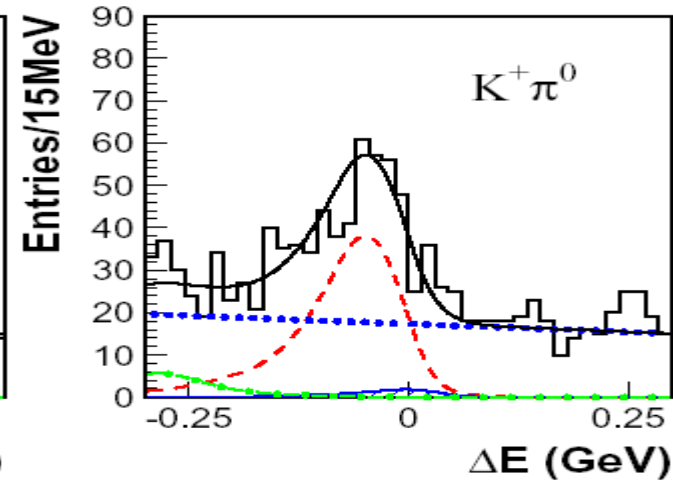
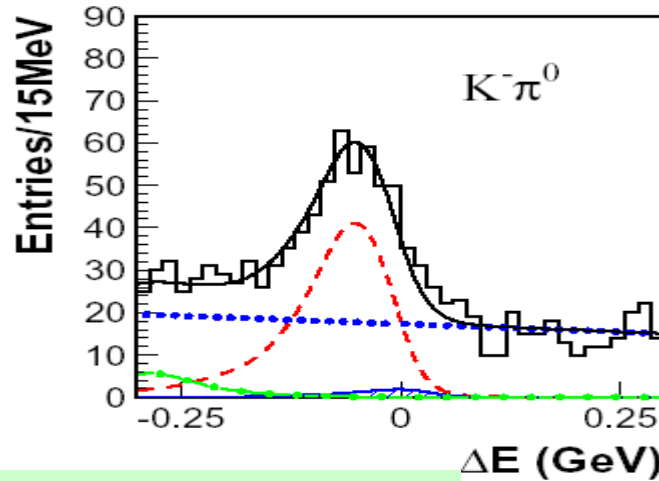


$A_{CP}(B \rightarrow K^+ \pi^0)$

Sakai

B 275M $B\bar{B}$
BELLE
New

$K^\pm \pi^0: 728 \pm 53$

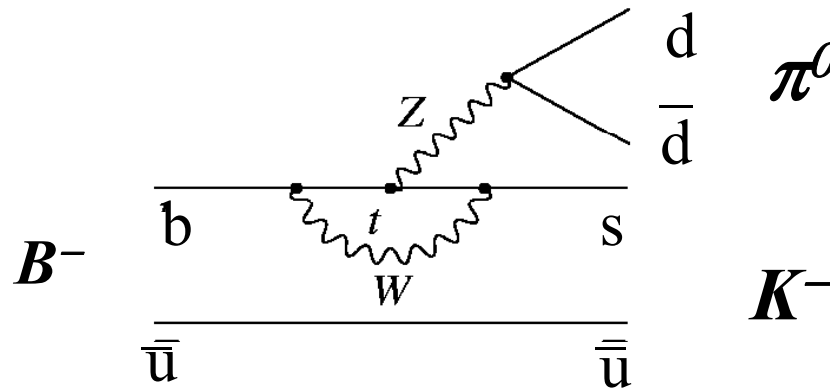


$A_{CP}(K^\pm \pi^0) = 0.04 \pm 0.05 \pm 0.02$

hint that $A_{CP}(K^+ \pi^-) \neq A_{CP}(K^\pm \pi^0)$? (2.4σ)

[also seen by BaBar]

Large EW penguin (Z^0) ?
New Physics ?





It would seem that we are well on the way to understanding the basis of particle–antiparticle asymmetry in the early Universe.

In fact, we are not. The KM predictions depend crucially on the masses of the intermediate-mass *s* and *c* quarks. But the high temperature of the Universe just after the Big Bang makes these masses irrelevant in calculations of the cosmic-matter excess. The degree of asymmetry predicted by the KM model is ten orders of magnitude too small.

reveal exotic
the Universe.

of quark were known: strange (*s*). But in the 1970s more were discovered: the heavy bottom (*b*) quark. This astounding success came from specific experiments that produced *c*–antiquark pairings. In 1973, it was a *b* quark or *b* antiquark (the Kobayashi–Maskawa mechanism). The idea, proposed by Makoto Kobayashi and Toshihide Maskawa, was that experiments could be done using two beams of different energies (one of positrons and one of electrons), motivated the accelerators at KEK and SLAC and Belle⁶ reported a KM asymmetry in a

the experimentalist



elementary particles of matter — has an antimatter counterpart with exactly the same mass, and exactly the opposite electric charge. Over the past 20 years, the theories of the weak and strong nuclear forces that have been built up on this basis have passed numerous rigorous experimental tests. The mathematical form of these theories allows little space for interactions that treat particles and antiparticles differently.

And yet the Universe, as far out as we can see, is made of matter, not of antimatter. We see no signals of the matter–antimatter annihilation that would happen on the edge of our local region if only this region were dominated by matter. So did the initial conditions of the Big Bang perhaps contain more matter than antimatter? It is possible. But in inflationary cosmology, the model that has successfully

process (shown here from left to right): a, in a standard 'box' diagram of weak quark-mixing interactions, quarks change type by exchanging a pair of particles, for example a heavy top (*t*) quark and a *W* boson, the intermediary of the weak force. Here, a \bar{B}^0 meson (quark content *db*) converts into a B^0 (*bd*). b, In a penguin process, the change of quark type occurs via a particle loop, which connects via a boson (wavy line; a gluon, *g*, gives a 'strong penguin'; a *Z*⁰ or a photon, *γ*, gives an 'electroweak penguin'); *γ* is a photon) to a further particle. Here, for example, a \bar{B}^0 or B^0 could be decaying into a K^0 (*ns*) or \bar{K}^0 (*d̄s*), plus an additional *u* or *d* quark that combines with the *u* or *d* antiquark in the *B* meson. The other end product is a π^0 particle, which can have quark content *uū* or *d̄d̄*. In both penguin and box processes, the particles represented by the heavy lines (square in a, circle in b) could be as-yet-undiscovered exotic particles. Recent results from the Belle and BaBar⁶ collaborations indicate

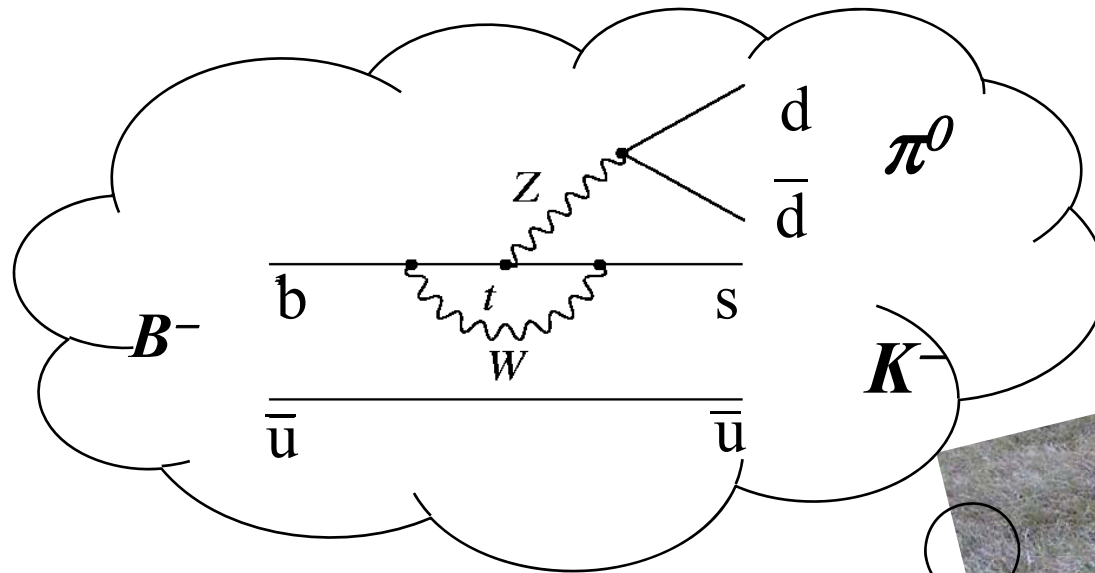
Since then, evidence accumulated by BaBar and Belle, in a data set of more than 1.2 million *B*-meson decays, has been used to fix the two crucial parameters of the KM theory to an accuracy of about 5%. Complementary measurements from other processes involving *B* mesons^{10–12} have confirmed these parameters to accuracies of between 10% and 20%.

It would seem that we are well on the way to understanding the basis of particle–antiparticle asymmetry in the early Universe.

In fact, we are not. The KM predictions depend crucially on the masses of the intermediate-mass *s* and *c* quarks. But the high temperature of the Universe just after the Big Bang makes these masses irrelevant in calculations of the cosmic-matter excess. The degree of asymmetry predicted by the KM model is ten orders of magnitude too small.



The Crawl' of one Ant



Going Up a Hill ...





My first B paper



WSH, Willey, Soni

VOLUME 58, NUMBER 16

PHYSICAL REVIEW LETTERS

20 APRIL 1987

an by Inami and Lim,⁹ and we follow their notation. The effective Lagrangean arising from Fig. 1 is

$$\mathcal{L}_{\text{eff}}^{b\bar{s} \rightarrow l^+ l^-} = 2\sqrt{2}G_F \chi v_i \{ \bar{C}_i (\bar{s} \gamma_\mu L b) (\bar{l} \gamma_\mu L l) - s_W^2 (F_1^i + 2\bar{C}_i^Z) (\bar{s} \gamma_\mu L b) (\bar{l} \gamma_\mu l) - s_W^4 F_2^i [\bar{s} i \sigma_{\mu\nu} (q_\nu / q^2) (m_s L + m_b R) b] (\bar{l} \gamma_\mu l) \}, \quad (1)$$

$$\mathcal{L}_{\text{eff}}^{b\bar{s} \rightarrow \nu \bar{\nu}} = -2\sqrt{2}G_F \chi v_i \bar{D}_i (\bar{s} \gamma_\mu L b) (\bar{\nu} \gamma_\mu L \nu), \quad (2)$$

where $\chi = g^2/16\pi^2$, $v_i \equiv V_{is}^* V_{ib}$, i is summed from 2 to n (where n is the number of generations),¹⁰ s_W is the sine of the Weinberg angle, and we exhibit¹¹

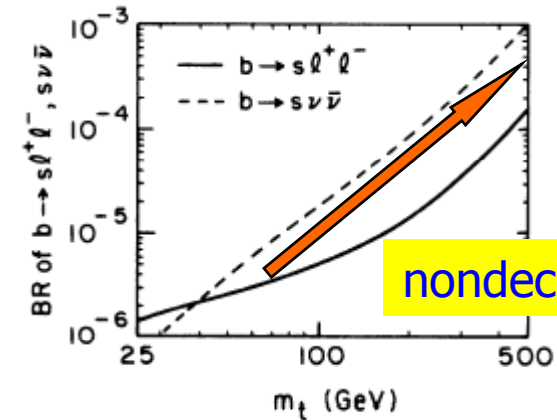
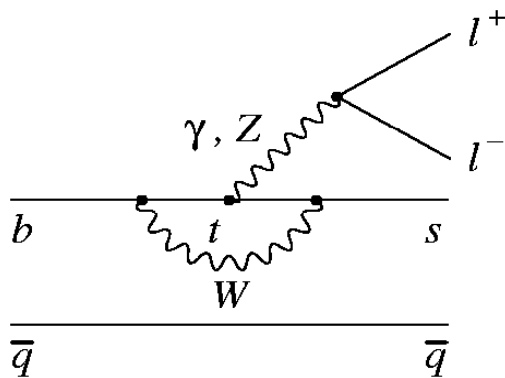
$$\bar{C}_i \equiv \bar{C}_i^Z + \bar{C}_i^{\text{box}} = \frac{1}{4} x_i + \frac{3}{4} \left(\frac{x_i}{x_i - 1} \right)^2 \ln x_i - \frac{3}{4} \frac{x_i}{x_i - 1},$$

$$\bar{D}_i \equiv \bar{D}_i^Z + \bar{D}_i^{\text{box}} = \frac{1}{4} x_i + \frac{3}{4} \frac{x_i (x_i - 2)}{(x_i - 1)^2} \ln x_i + \frac{3}{4} \frac{x_i}{x_i - 1},$$

dimensions

γ	Z	(3)
$\alpha G_F < G_F^2 m_t^2$		

where $x_i = m_i^2/M_W^2$, and m_i is the internal quark mass. The important feature of Eqs. (3) and (4) is the term $x_i/4$,⁸





Nondecoupling



Decoupling Thm: Heavy **Masses** are decoupled in QED/QCD
∴ Appear in Propagator

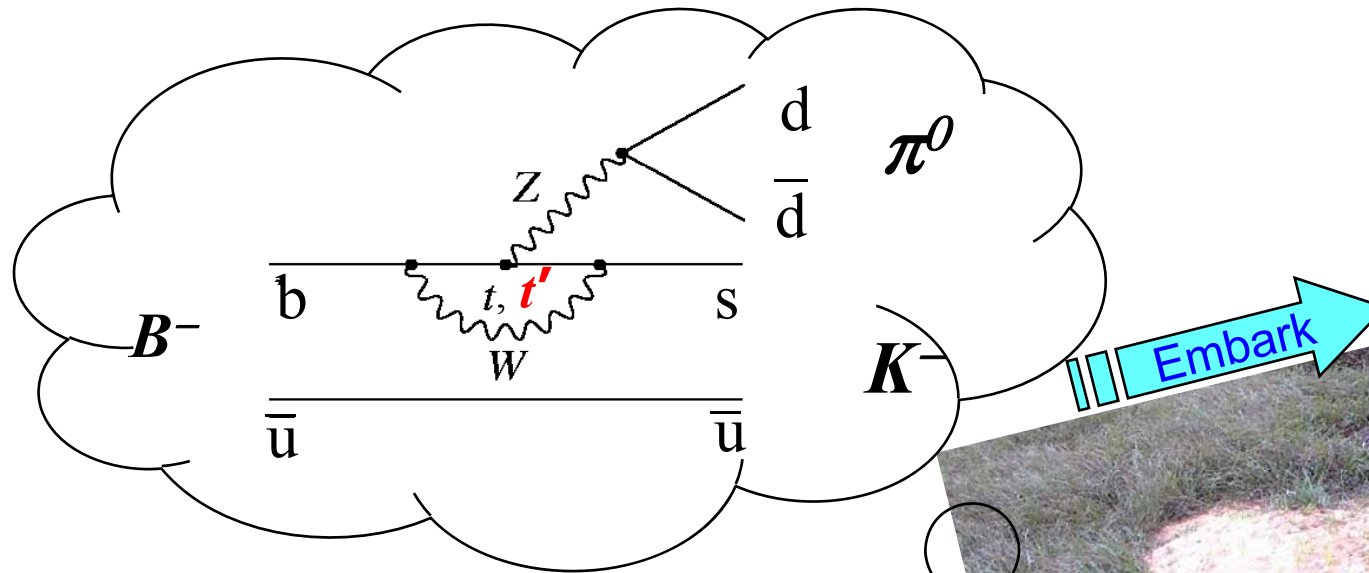
Nondecoupling: Yukawa Couplings λ_Q Appear in Numerator

Subtlety of Spont. Broken Gauge Theory

dynamical



The Crawl' of one Ant



Going Up a Hill* ...



* I am aware of hadronic (color-suppressed C) complications



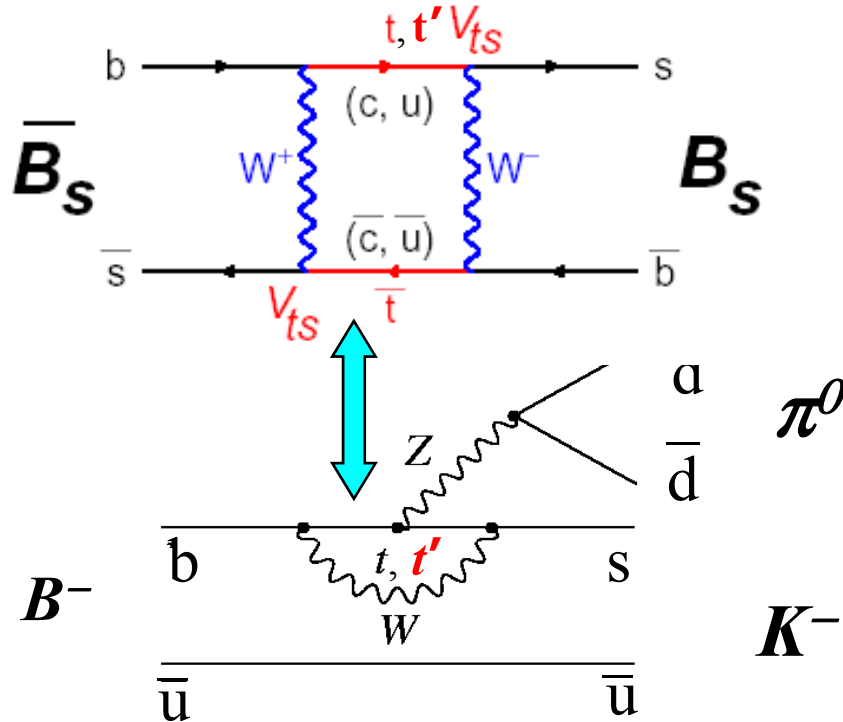
This is Still the Standard Model



$$\lambda_{t'} \equiv V_{t's}^* V_{t'b} \equiv r_{sb} e^{i\phi_{sb}}$$



Arhrib and WSH, EPJC'03

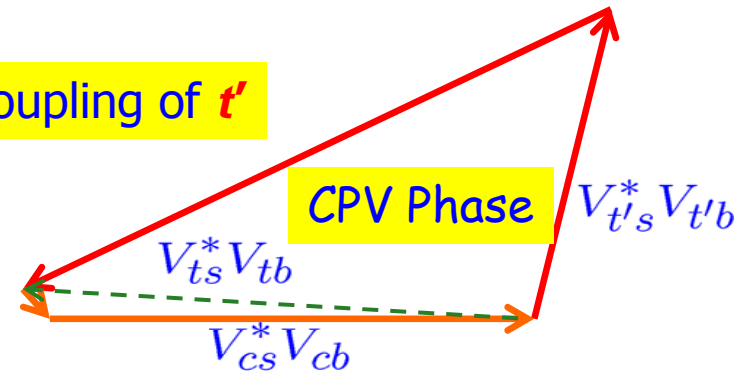


$$t \Rightarrow t, t'$$

$$\lambda_u + \lambda_c + \lambda_t + \lambda_{t'} = 0$$

$$\lambda_t \approx -\lambda_c - \lambda_{t'}$$

Nondecoupling of t'



$$M_{12} \propto f_{B_s}^2 B_{B_s} \left\{ \lambda_c^2 S_0(t, t) + 2\lambda_c \lambda_{t'} [S_0(t, t) - S_0(t, t')] + \lambda_{t'}^2 [S_0(t, t) - 2S_0(t, t') + S_0(t', t')] \right\}$$

GIM Respecting

$$H_{\text{eff}}^4 = \frac{G_F}{\sqrt{2}} \left[\lambda_u (C_1 O_1 + C_2 O_2) + \sum_{i=3}^{10} (\lambda_c C_i^t - \lambda_{t'} (C_i^{t'} - C_i^t)) O_i \right]$$



EWP/Box Sensitivity to 4th Gen.

γ, g less sensitive

(No New Operators)

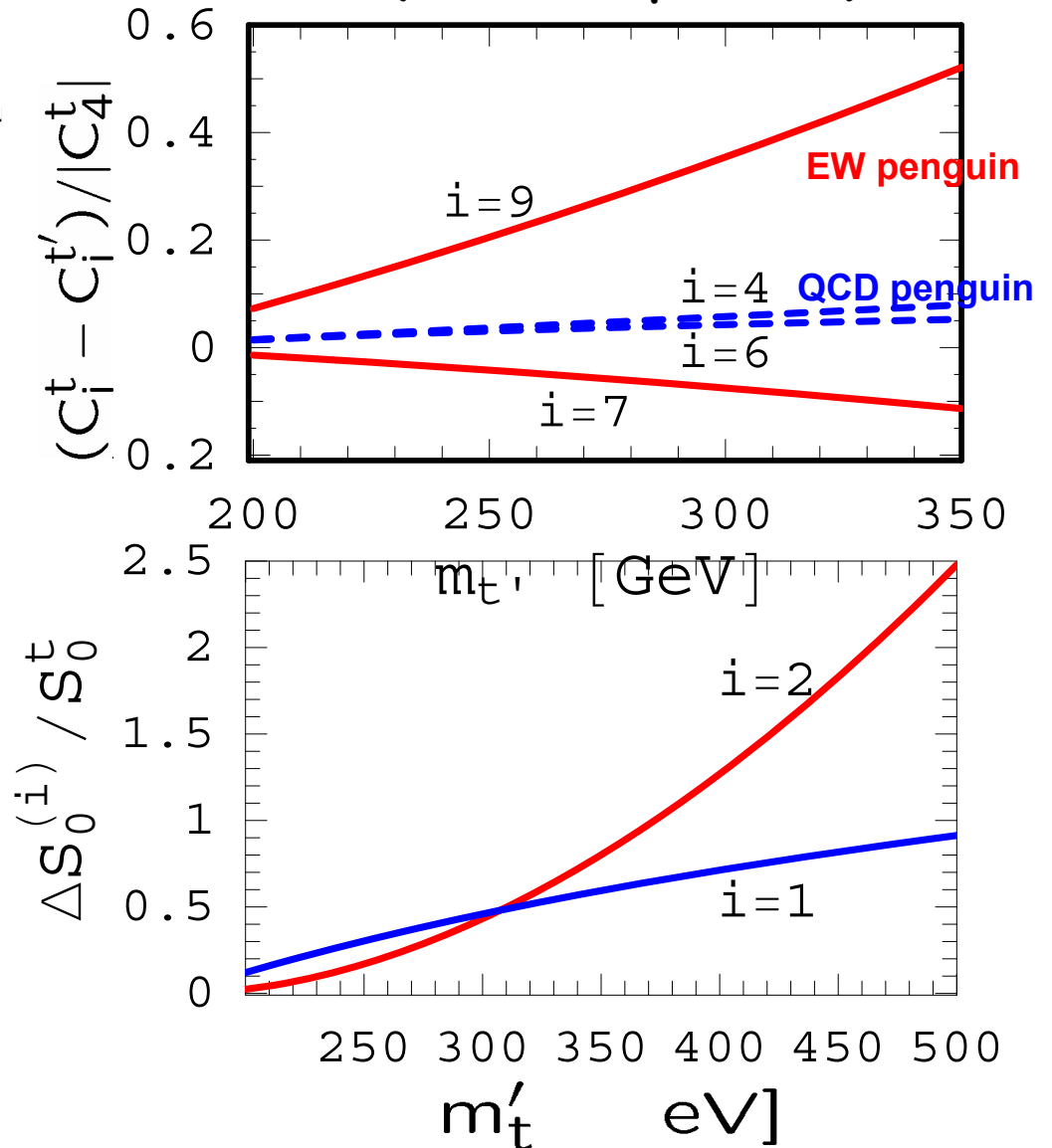


$$C_9^t - C_9^{t'} \propto x_t - x_{t'}$$

nondecoupling

$$\Delta S_0^{(1)} = S_0(t, t') - S_0(t, t)$$

$$\Delta S_0^{(2)} = S_0(t', t') + S_0(t, t) - 2S_0(t, t')$$

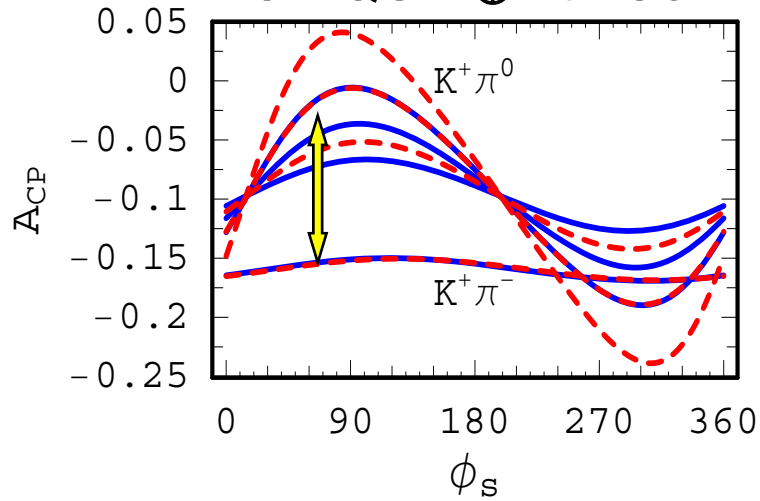




$$\Delta A = A_{K^+\pi^0} - A_{K^+\pi^-} \sim 15\% \text{ and } P_{EW}^{b \rightarrow s}$$



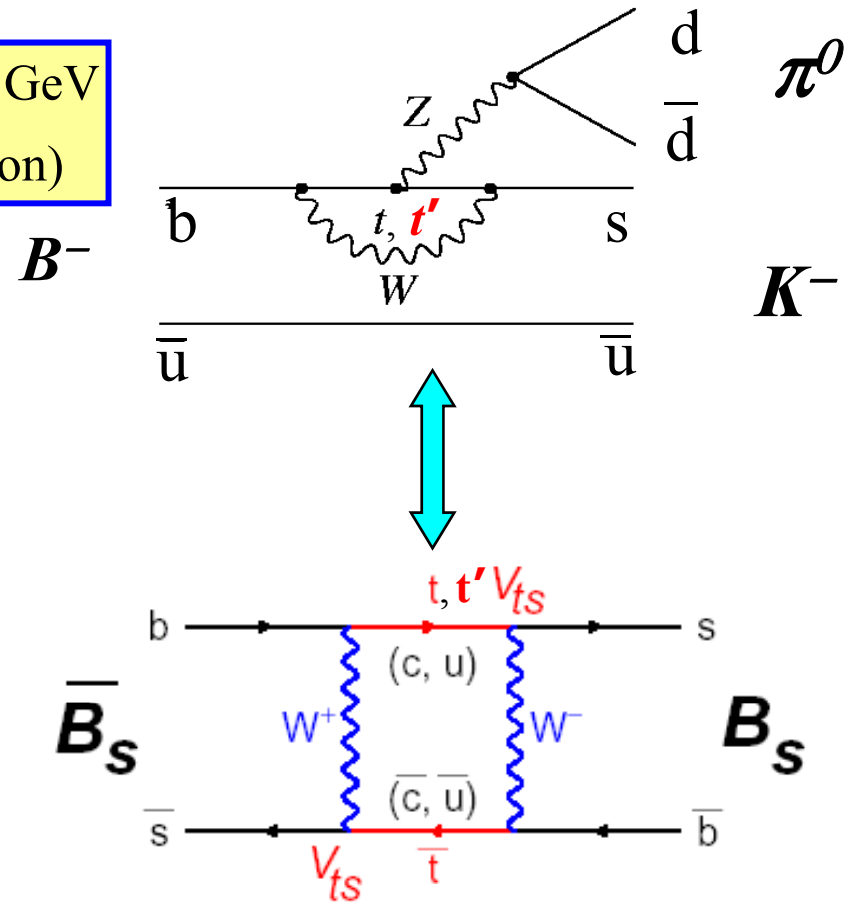
LO PQCD \oplus 4th Gen.



WSH, Nagashima, Soddu, PRL'05

$$\Delta A \approx 12\% \text{ vs } 15\% \text{ (data)}$$

$m_{t'} = 300 \text{ GeV}$
(illustration)





Difference in B^+ and B^0 Direct CP Asymmetry as an Effect of a Fourth Generation

Wei-Shu Hou, Makiko Nagashima, and Andrea Soddu

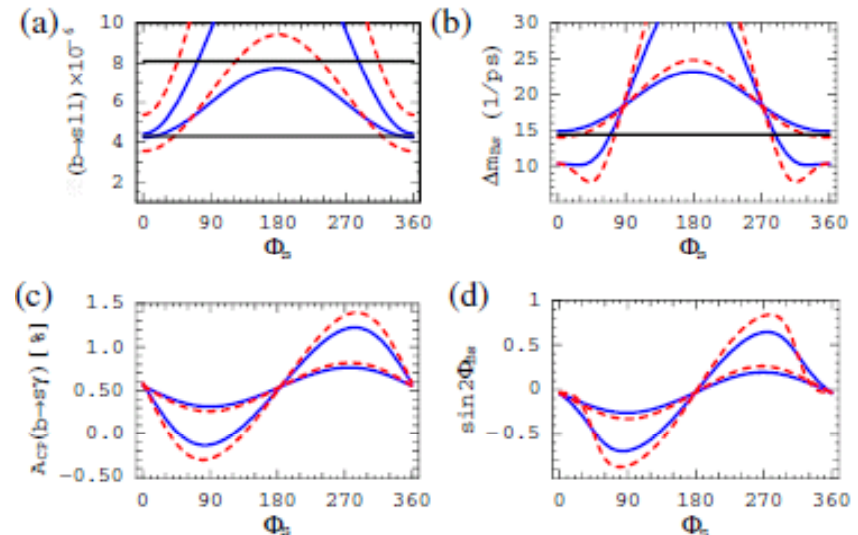
Department of Physics, National Taiwan University, Taipei, Taiwan 106, Republic of China

(Received 8 March 2005; revised manuscript received 20 June 2005; published 30 September 2005)

Direct CP violation in $B^0 \rightarrow K^+ \pi^-$ decay has emerged at the $\sim 10\%$ level, but the asymmetry in $B^+ \rightarrow K^+ \pi^0$ mode is consistent with zero. This difference points towards possible new physics in the electroweak penguin operator. We point out that a sequential fourth generation, with sizable $V_{t's}^* V_{t'b}$ and near maximal phase, could be a natural cause. We use the perturbative QCD factorization approach for $B \rightarrow K\pi$ amplitudes. While the $B^0 \rightarrow K^+ \pi^-$ mode is insensitive to t' , we critically compare t' effects on direct CP violation in $B^+ \rightarrow K^+ \pi^0$ with $b \rightarrow s\ell^+\ell^-$ and B_s mixing. If the $K^+ \pi^0 - K^+ \pi^-$ asymmetry difference persists, we predict $\sin 2\Phi_{B_s}$ to be negative

As prediction, we find $\sin 2\Phi_{B_s} < 0$ for CPV in B_s mixing, which is plotted versus ϕ_s in Fig. 3(d). We find $\sin 2\Phi_{B_s}$ in the range of -0.2 to -0.7 and correlating with $\mathcal{A}_{K^+\pi^0} - \mathcal{A}_{K^+\pi^-}$. Three generation SM predicts zero.

Note that refined measurements of $\mathcal{B}(b \rightarrow s\ell\ell)$ and future measurements of Δm_{R_s} and $\sin 2\Phi_{R_s}$, together with theory improvements, can pinpoint $m_{t'}$, r_s , and ϕ_s . We note further that [6] $14.4 \text{ ps}^{-1} < \Delta m_{B_s} < 21.8 \text{ ps}^{-1}$ cannot yet be excluded because data are compatible with a signal in this region. We eagerly await B_s mixing and associated CPV measurement in the near future.





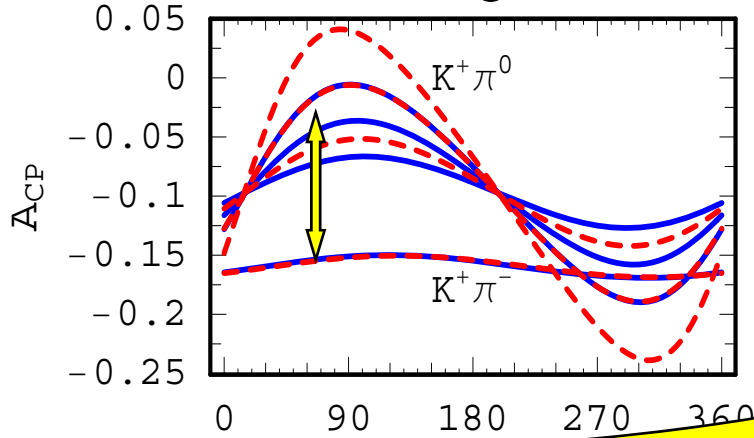
$$\Delta A = A_{K^+\pi^0} - A_{K^+\pi^-} \sim 15\% \text{ and } P_{EW}^{b \rightarrow s}$$

$$P_{EW}^{b \rightarrow s}$$



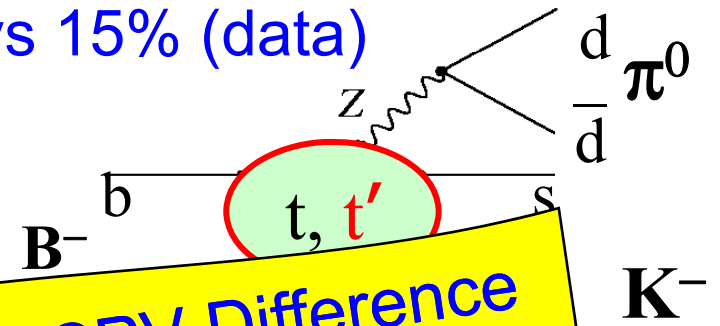
LO PQCD \oplus 4th Gen.

WSH, Nagashima, Soddu, PRL'05



$\Delta A \approx 12\%$ vs 15% (data)

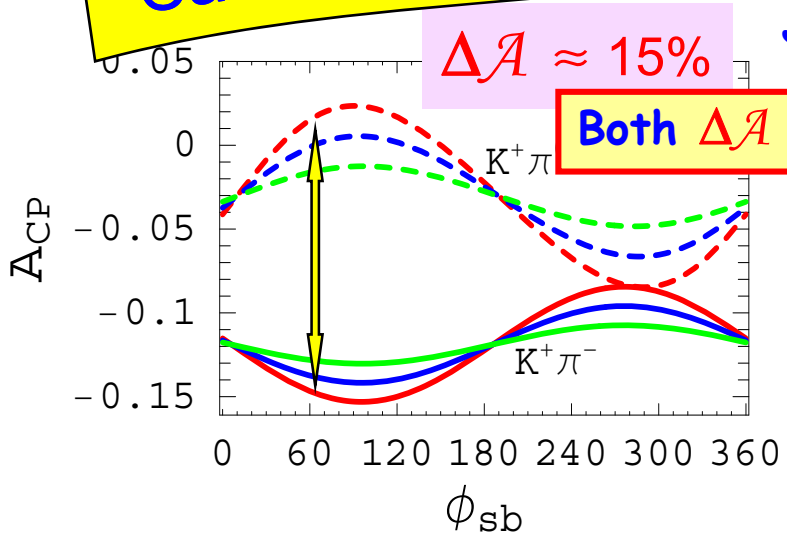
$m_{t'} = 300 \text{ GeV}$
(illustration)



Can Account for Belle/BaBar Direct CPV Difference

Li, Mishima, Nagashima, PRL'07

$r_{sb} = 0.03$: red, dash
0.02: blue, solid
0.01: green, dot-dash



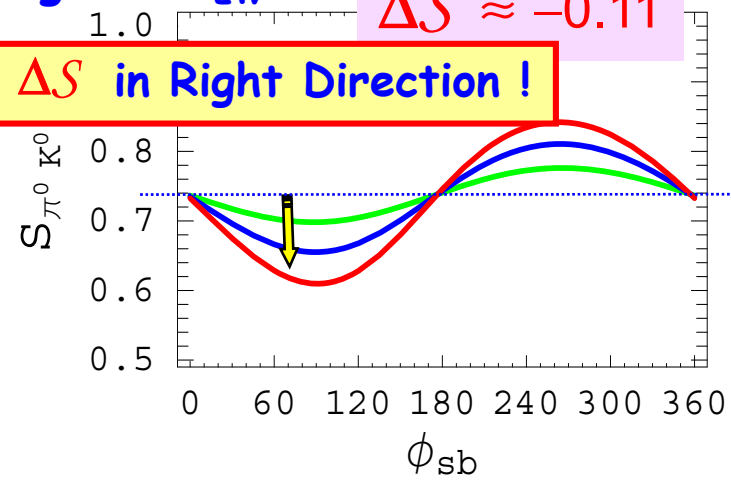
$\Delta A \approx 15\%$

Joining C & P_{EW}

Both ΔA and ΔS in Right Direction !

$\Delta S \approx -0.11$

consistent with data



SM3 input

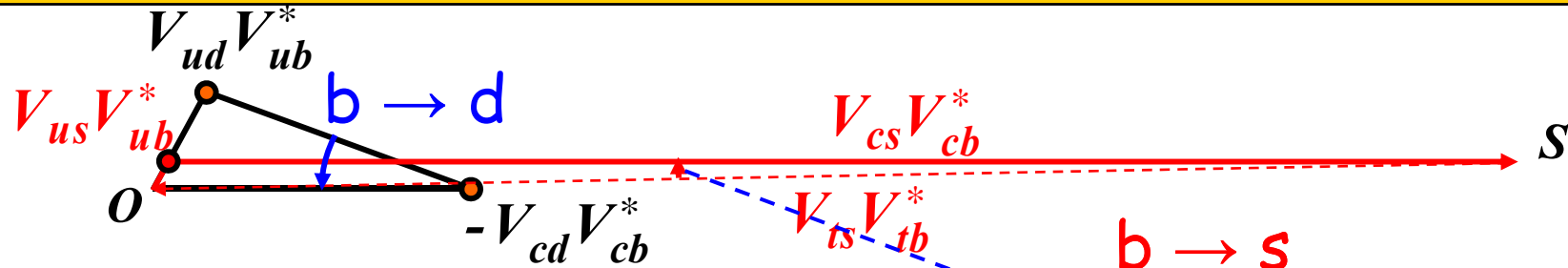


$b \leftrightarrow s$ CPV

III. Δm_{B_s} & $\mathcal{B}(B \rightarrow X_s \ell^+ \ell^-)$: 2nd Prediction for $\sin 2\Phi_{B_s}$



Mixing-dep. CPV in B_d and B_s in SM



$$\sin 2\phi_1 = \sin 2\beta$$

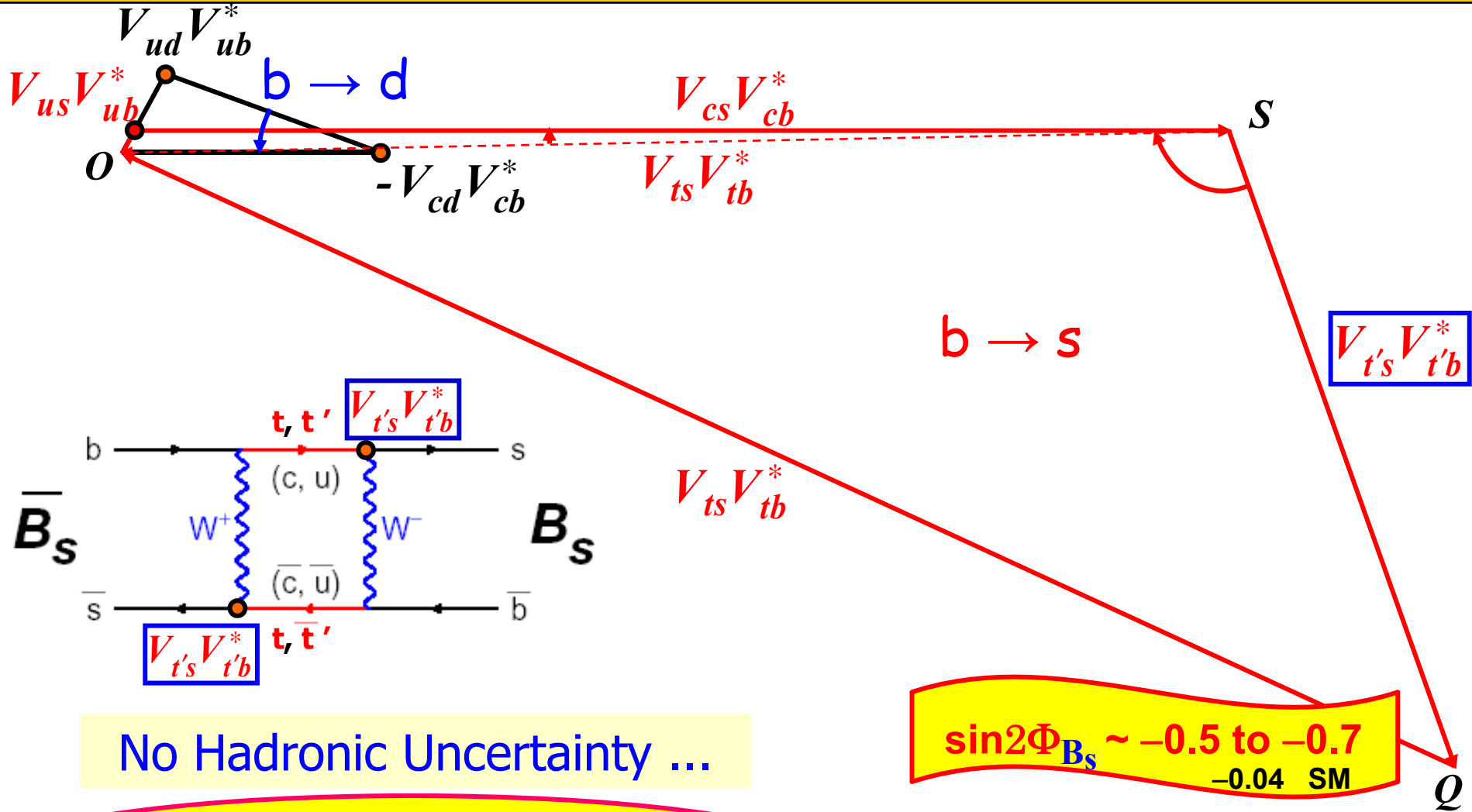
Measured by Belle/BaBar
in $B_d \rightarrow J/\psi K_S$

$$\sin 2\Phi_{B_s} \approx -0.04 \quad \text{in SM3}$$

Measure in $B_s \rightarrow J/\psi \phi$



Prediction: Large CPV in B_s Mixing



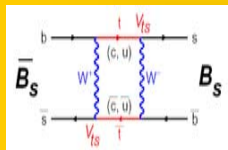
No Hadronic Uncertainty ...

Strength and Size of $\sin 2\Phi_{B_s}$

$\sin 2\Phi_{B_s} \sim -0.5 \text{ to } -0.7$
 -0.04 SM

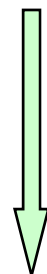
Despite Δm_{B_s} , $\mathcal{B}(b \rightarrow s \ell \ell)$ SM-like

WSH, Nagashima, Soddu, PRD'07
 PRL'05

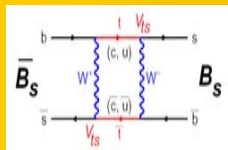


B_s Mixing vs $B \rightarrow X_s \ell^+ \ell^-$

different nondecoupl. functions



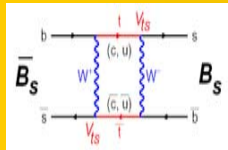
Large CPV in B_s Mixing



Use nominal $m_{t'}$ = 300 GeV
Change $m_{t'}$, Change parameter range
Effect the Same.

(Similar)

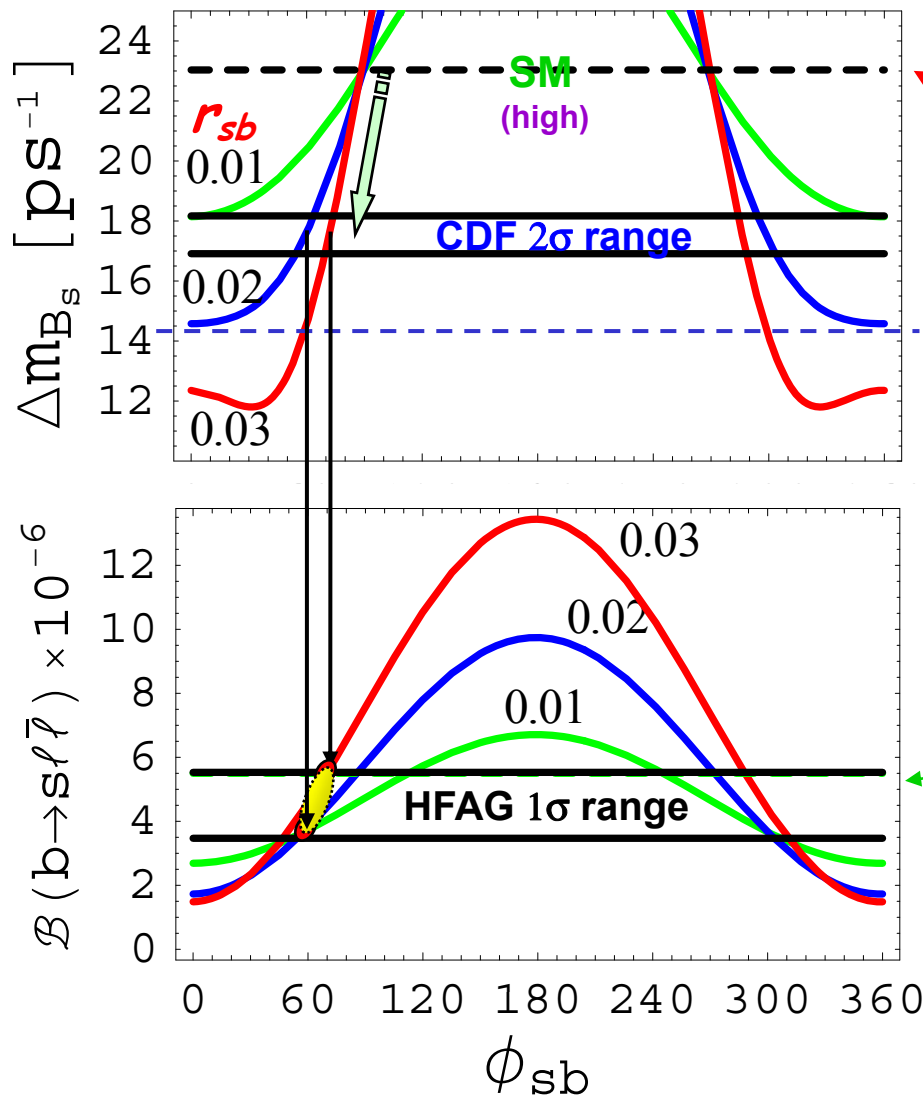
N.B. $\Delta A_{K\pi}$ suggest parameter space



$$\lambda_{t'} \equiv V_{t's}^* V_{t'b} \equiv r_{sb} e^{i\phi_{sb}}$$



WSH, Nagashima, Soddu, hep-ph/0610385 (PRD'07)



$$f_{B_s} \sqrt{B_{B_s}} = 295 \pm 32 \text{ MeV}$$

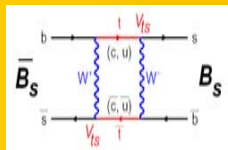
- Fixed $r_{sb} \Rightarrow$ Narrow ϕ_{sb} Range destructive with top
- For $r_{sb} \sim 0.02 - 0.03$, $[V_{cb} \sim 0.04$

ϕ_{sb} Range $\sim 60^\circ - 70^\circ$
Finite CPV Phase

Consistent w/ $\mathcal{B}(b \rightarrow s \ell \bar{\ell})$
 SM-like!

Large CPV Possible !

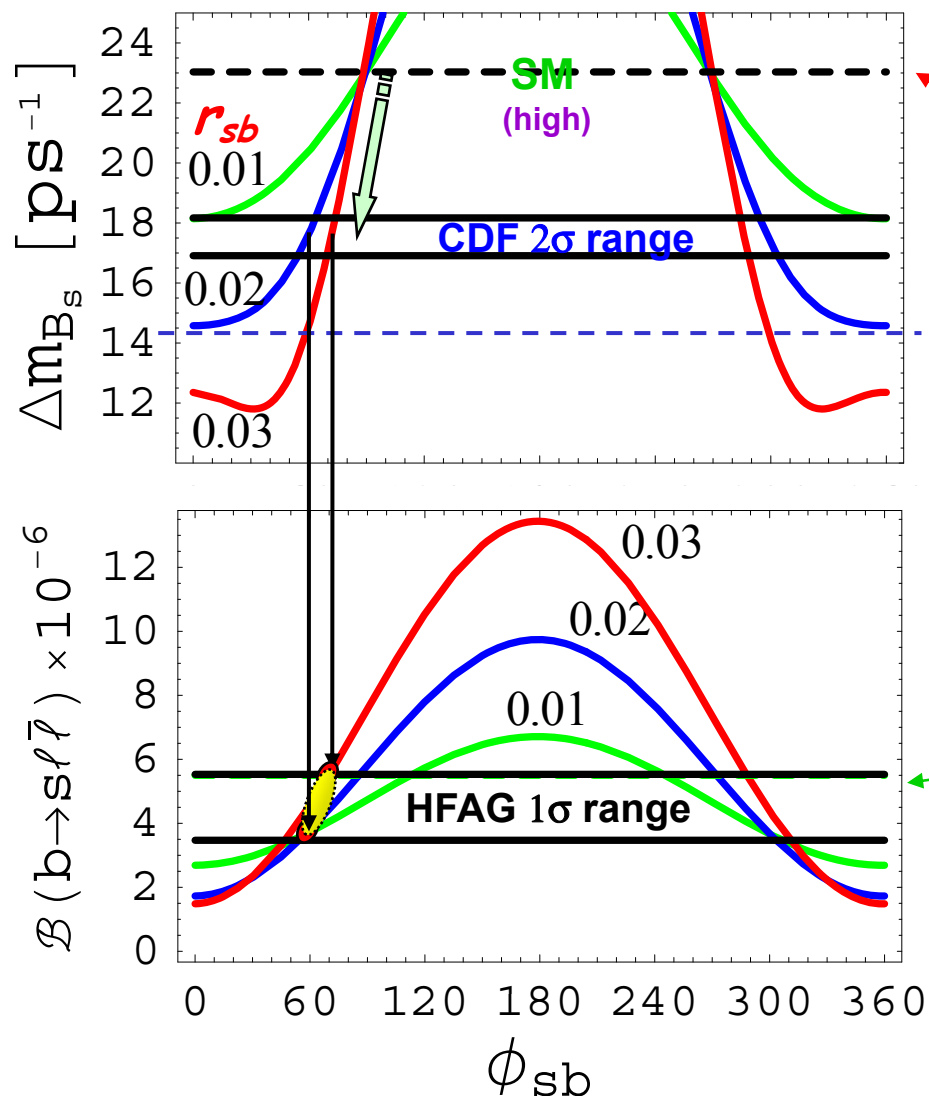
Despite Δm_{B_s} , $\mathcal{B}(b \rightarrow s \ell \bar{\ell})$ SM-like



$$\lambda_{t'} \equiv V_{t's}^* V_{t'b} \equiv r_{sb} e^{i\phi_{sb}}$$



WSH, Nagashima, Soddu, hep-ph/0610385 (PRD'07)



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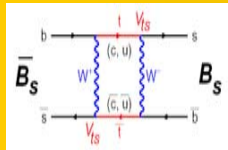
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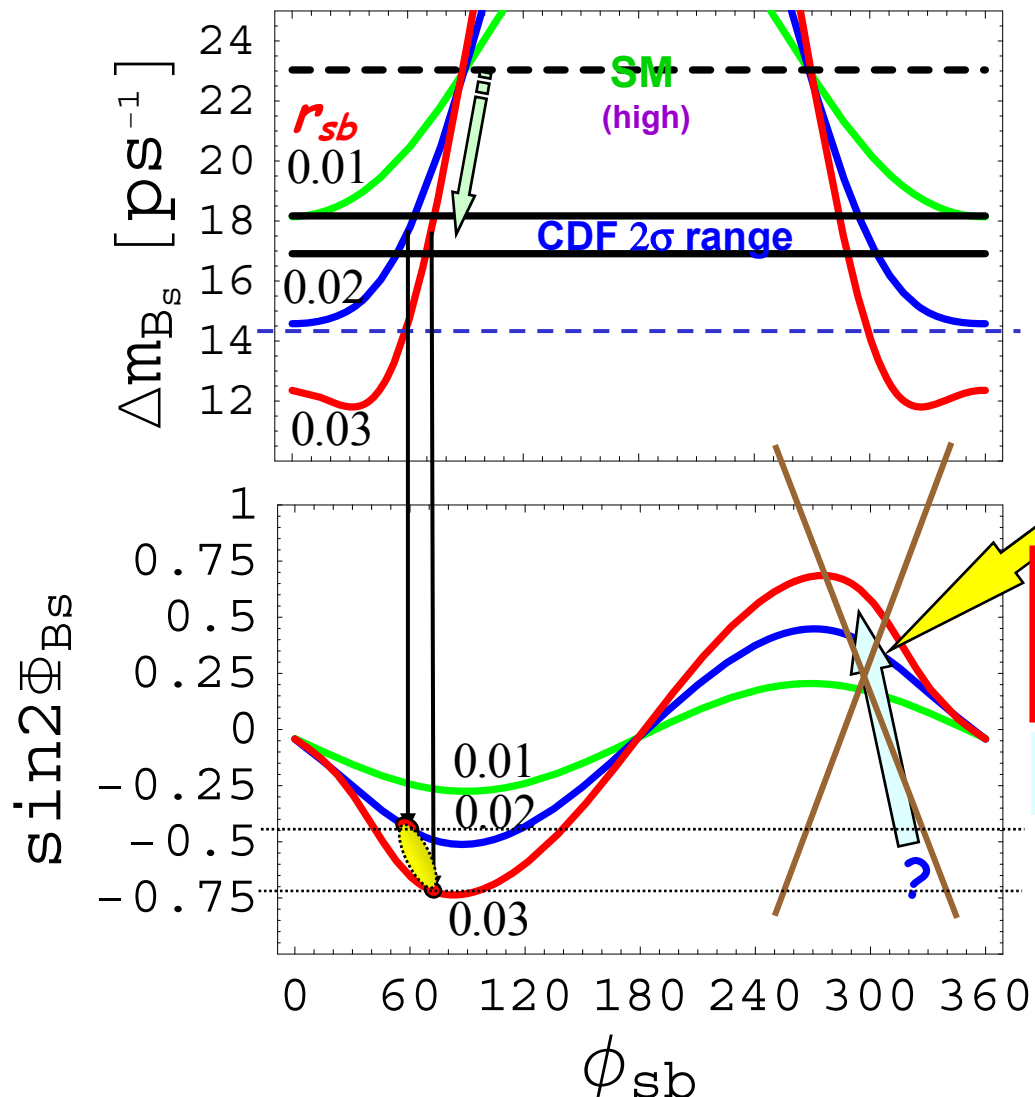
Despite Δm_{B_s} , $\mathcal{B}(b \rightarrow s \ell \bar{\ell})$ SM-like



Large CPV in B_s Mixing



WSH, Nagashima, Soddu, hep-ph/0610385



$\Delta \mathcal{A}_{K\pi}, \Delta S$

Can Large CPV in B_s Mixing
Be Measured @ Tevatron ?

Sign Predicted ! Sure thing by
LHCb ca. 2008

$\sin 2\Phi_{B_s} \sim \pm 0.5 - \pm 0.7$

Despite $\Delta m_{B_s}, \mathcal{B}(b \rightarrow sll)$ SM-like



Prediction: Large CPV in B_s Mixing



WSH, Nagashima, Soddu, PRD'07

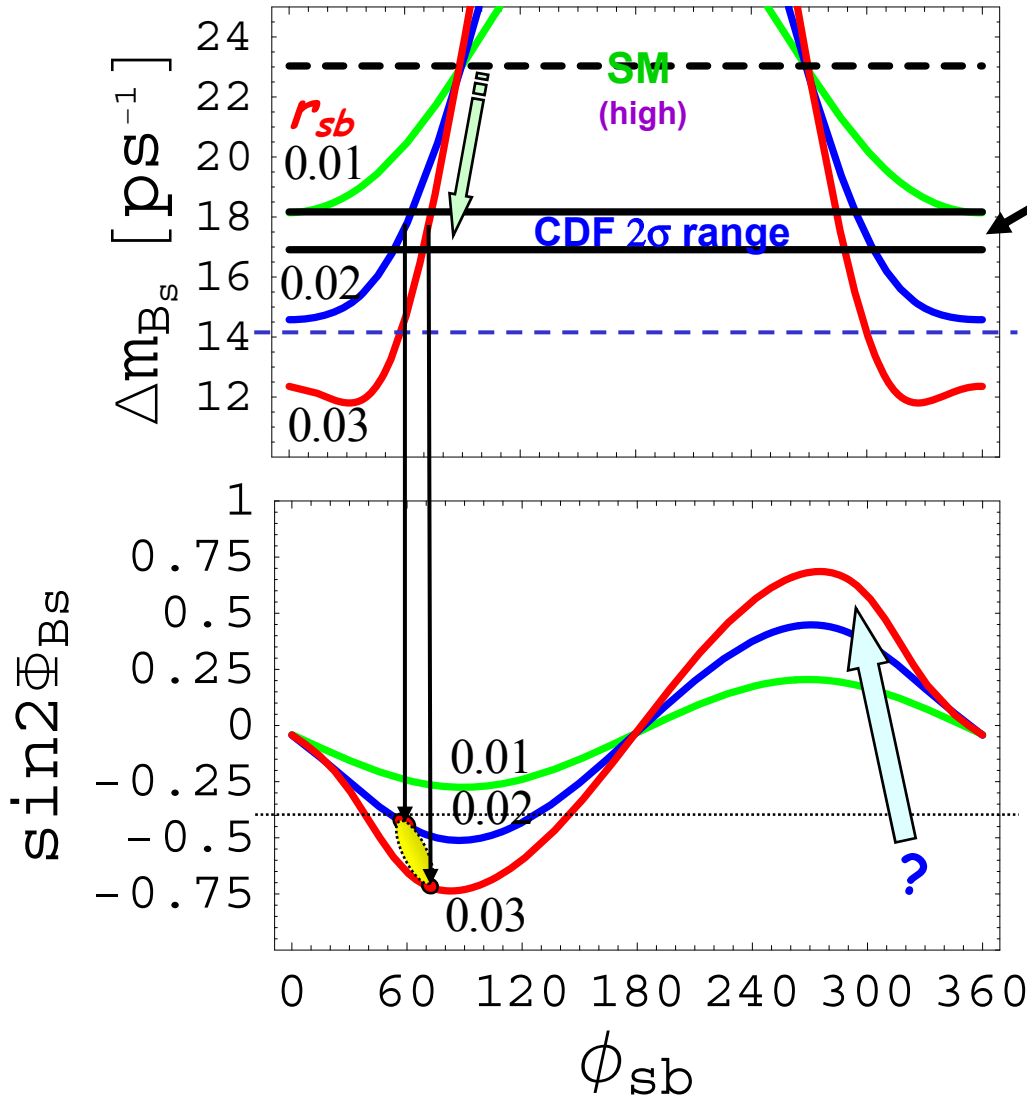
$$f_{B_s} \sqrt{B_{B_s}} = 295 \pm 32 \text{ MeV}$$

B_s Mixing Measured @ Tevatron in 4/2006

• For $r_{sb} \sim 0.02 - 0.03$, $[V_{cb} \sim 0.04$

ϕ_{sb} Range $\sim 60^\circ - 70^\circ$

Finite CPV Phase



$\sin 2\Phi_{B_s} \sim -0.5 - -0.7$
 -0.04 SM

Despite Δm_{B_s} , $\mathcal{B}(b \rightarrow s \ell \ell)$ SM-like

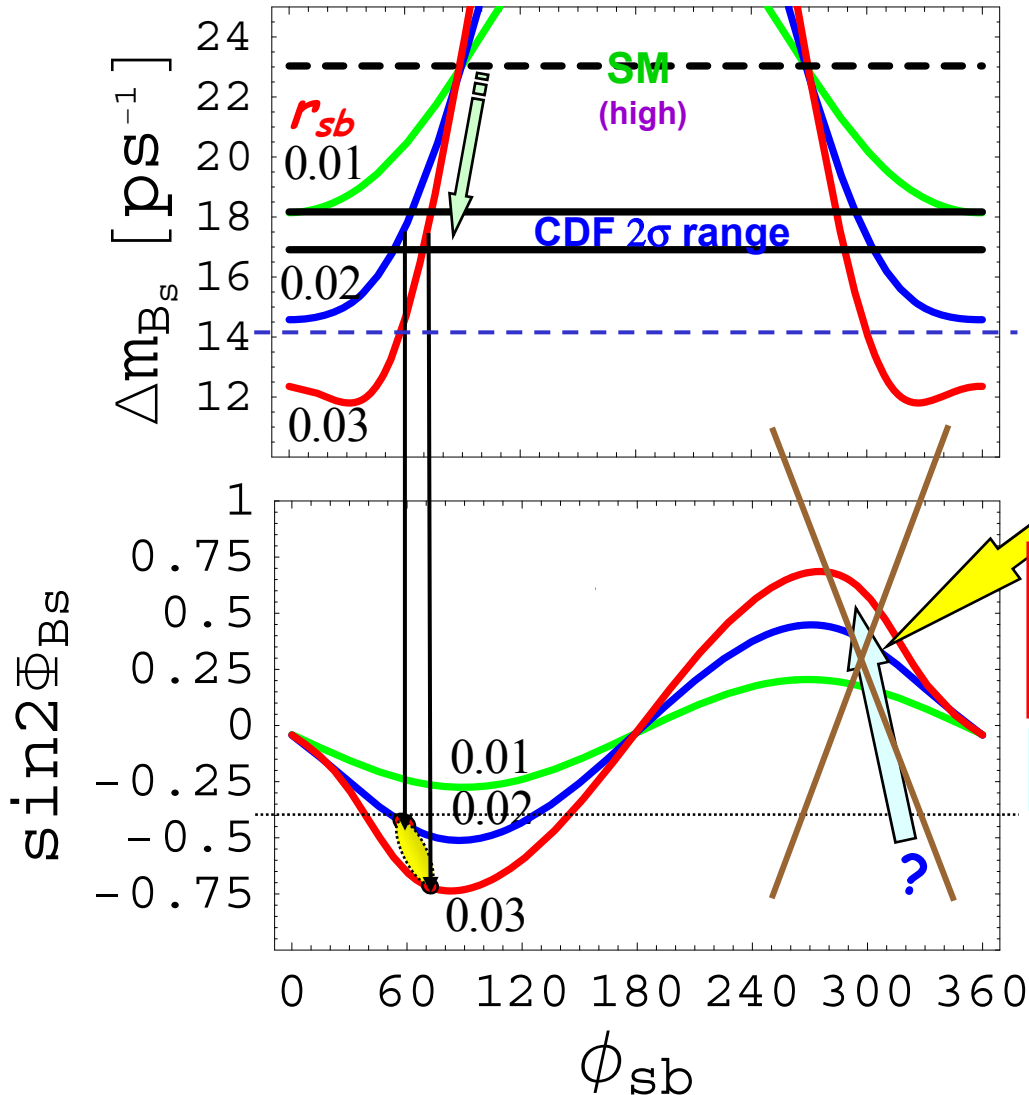
WSH, Nagashima, Soddu, PRL'05



Prediction: Large CPV in B_s Mixing



WSH, Nagashima, Soddu, PRD'07



$\Delta A_{K\pi}, \Delta S$

Can Large CPV in B_s Mixing
Be Measured @ Tevatron ?

Sign Predicted !

Sure thing by
LHCb ca. 2010 (?)

$\sin 2\Phi_{B_s} \sim -0.5 - -0.7$
-0.04 SM

Despite $\Delta m_{B_s}, \mathcal{B}(b \rightarrow sll)$ SM-like

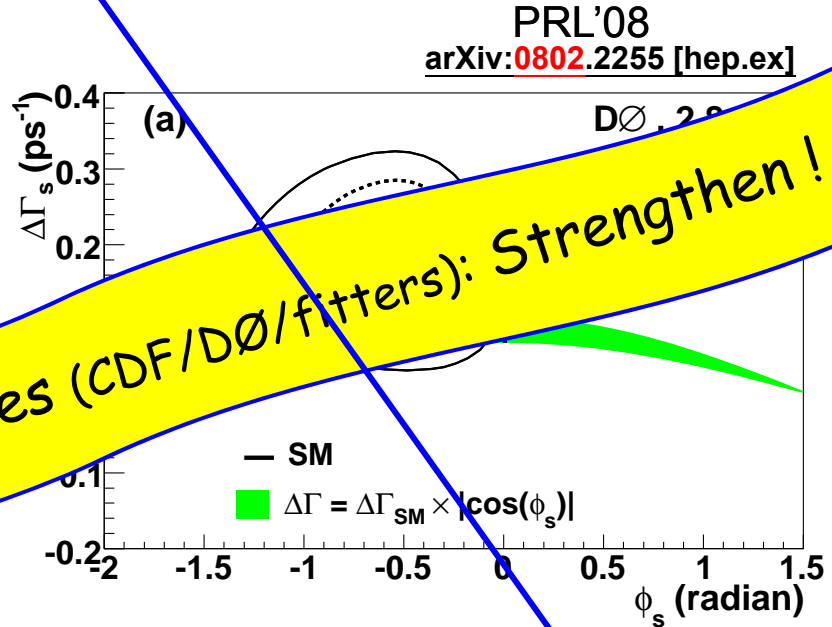
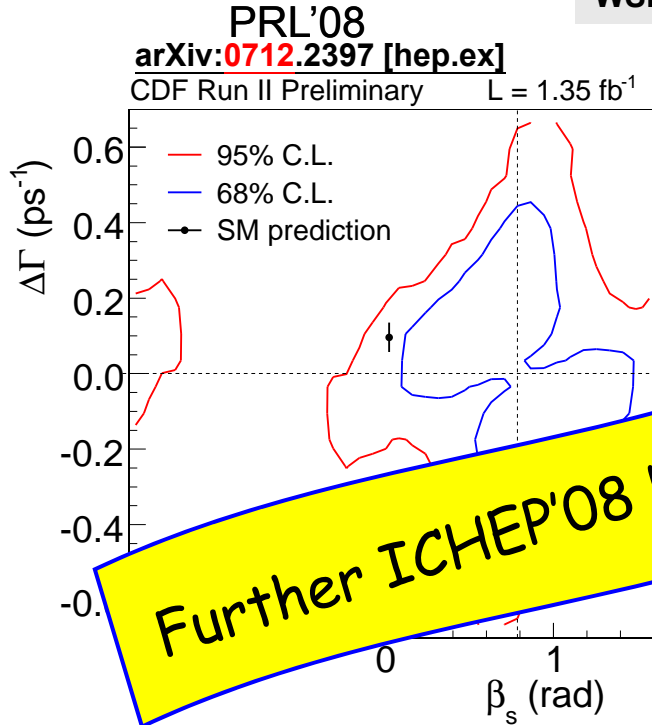
WSH, Nagashima, Soddu, PRL'05



$$\sin 2\Phi_{B_s} \sim -0.5 \text{ - } -0.7$$



WSH, Nagashima, Soddu, PRD'07 (already in 05)



Further ICHEP'08 Updates (CDF/DØ/fitters): Strengthen!

Observable	68% Prob.	95% Prob.
$\phi_{B_s} [^\circ]$	-19.9 ± 5.6	$[-36.45, 9.29]$
	-68.2 ± 4.9	$[-78.45, -58.2]$

UTfit

arXiv:0803.0659 [hep.ph]

$$\sin 2\Phi_{B_s} = -0.64 \pm ? \quad \sim 2.8\sigma$$

Incredible !!!



An Updated Measurement of the CP Violating Phase $\beta_s^{J/\psi\phi}$



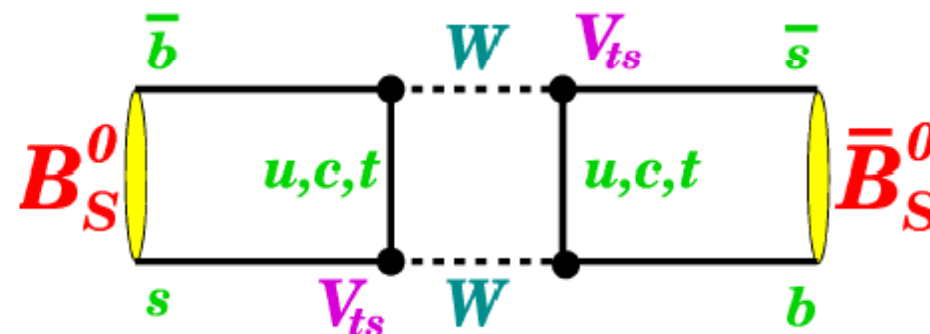
The CDF Collaboration¹

CDF/ANAL/BOTTOM/PUBLIC/9458

Version 1.0

August 7, 2008

It is interesting to note that the Belle and *BABAR* collaborations have observed an asymmetry between direct CP asymmetries of charged and neutral $B \rightarrow K\pi$ decays with 5σ significance [5, 6]. In the absence of an under-estimation of the contribution from color-suppressed tree decays, it is difficult to explain this discrepancy without some source of new physics contributing to the electroweak penguin which governs the $b \rightarrow s$ transition. In the standard model, this isospin-violating diagram should be highly suppressed, but if a new source of physics is indeed present in these transitions it may be enough to cause the different CP asymmetries that have been observed. In the $B_s^0 \rightarrow J/\psi\phi$ decay, the $b \rightarrow s$ transition occurs through the mixing box diagram shown in Fig. 1. It is possible that new particles could enter this transition through the $b \rightarrow s$ quark transition. While there are surely a number of possible sources of new physics that might give rise to such discrepancies, George Hou predicted the presence of a t' quark with mass between ~ 300 and $1,000 \text{ GeV}/c^2$ in order to explain the Belle result and predicted *a priori* the observation of a large CP -violating phase in $B_s^0 \rightarrow J/\psi\phi$ decays [7, 8]. Another result of interest in the context of these measurements is the excess observed at $\sim 350 \text{ GeV}/c^2$ in the recent t' search at CDF using 2.3 fb^{-1} of data [9]. In this direct search for a fourth generation up-type quark, a significance of less than 2σ is obtained for the discrepancy between the data and the predicted backgrounds, so that the effect, while intriguing, is presently consistent with a statistical fluctuation. A updated search with more data would also clearly be of interest, particularly if a large value of $\beta_s^{J/\psi\phi}$ persists with the addition of more data.



$$\sin 2\Phi_{B_s} = -\sin\beta_s = \sin\phi_s$$

(Conservative) outlook

% of CDF 'clones' that would observe a 5σ -effect, as a function of β_s

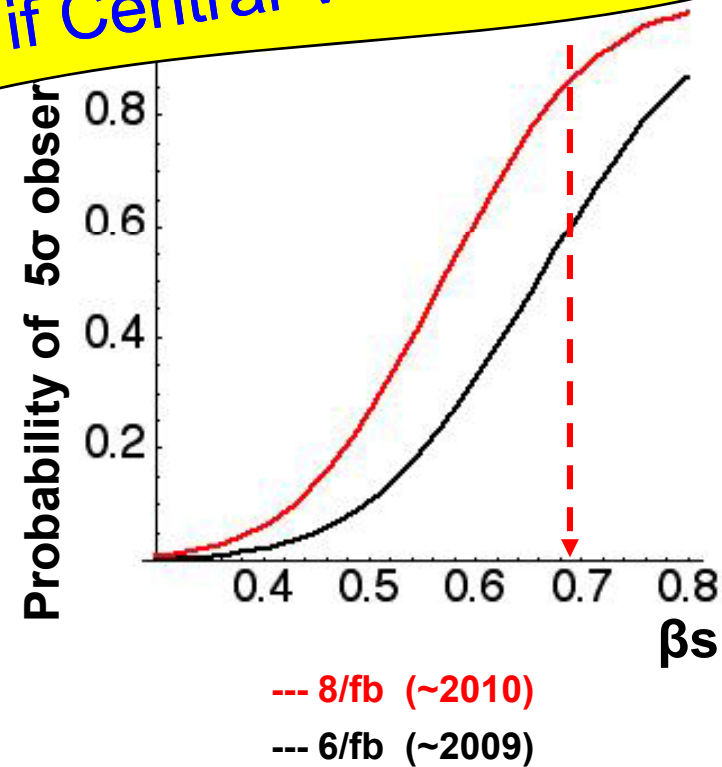
Assumptions

- ✓ $\Delta\Gamma_c = 0.1 \text{ ps}^{-1}$
- ✓ CDF β_s measuring efficiency
- ✓ No analysis improvements.
- ✓ No external constraints (A_{SL} , lifetimes) used.

Observation w/ 2010 Data if Central Value Stays!

CDF future will probably be better than that.

And $D\bar{D}$ will contribute too.





4 x 4 Unitarity \Rightarrow Z/K Constraints

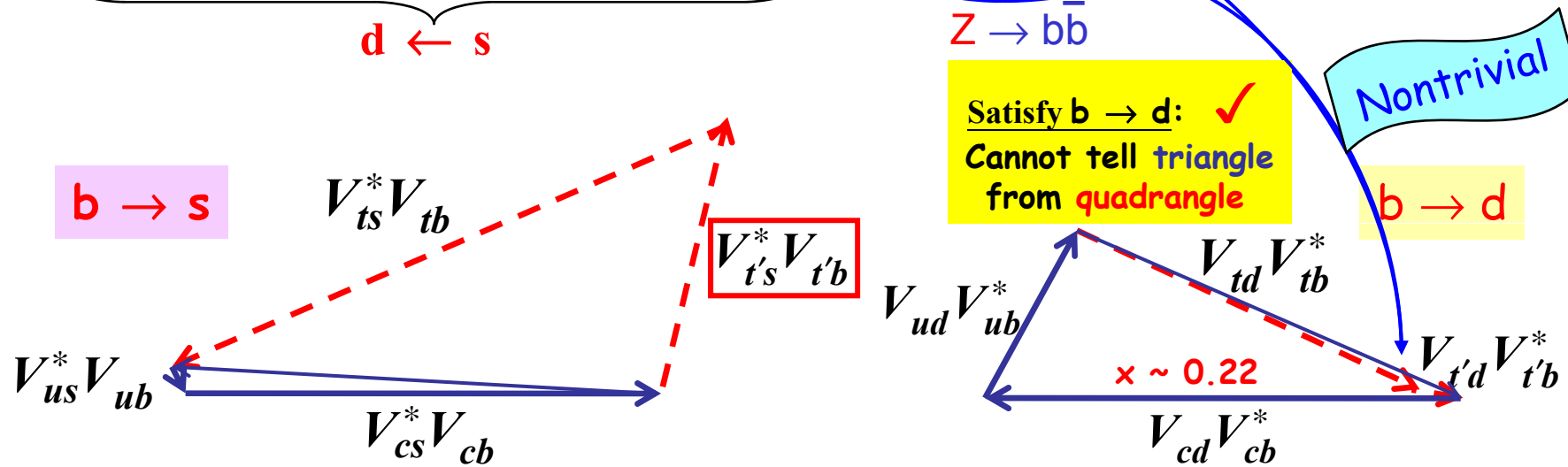
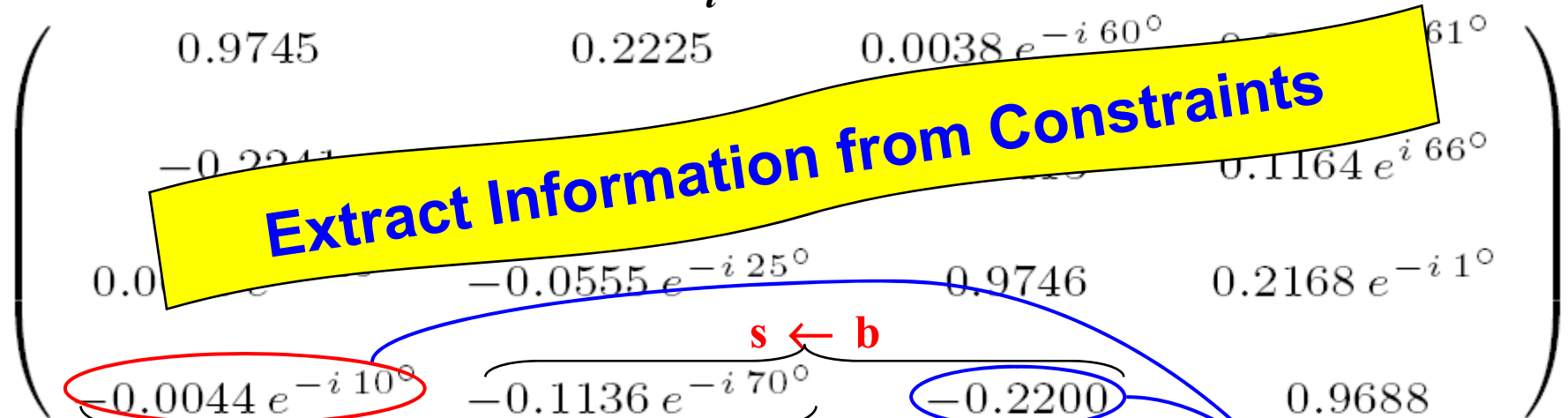


$$V_{CKM}^4 =$$

“Typical” CKM Matrix

WSH, Nagashima, Soddu, PRD'05

$$m_{t'} = 300 \text{ GeV}$$





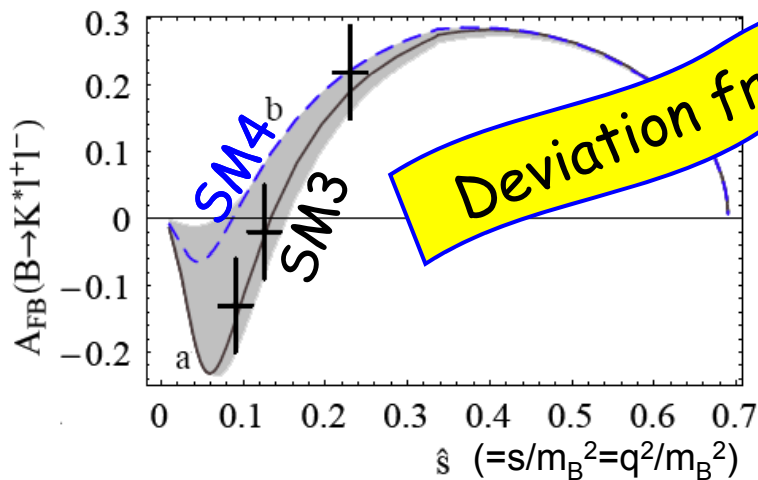
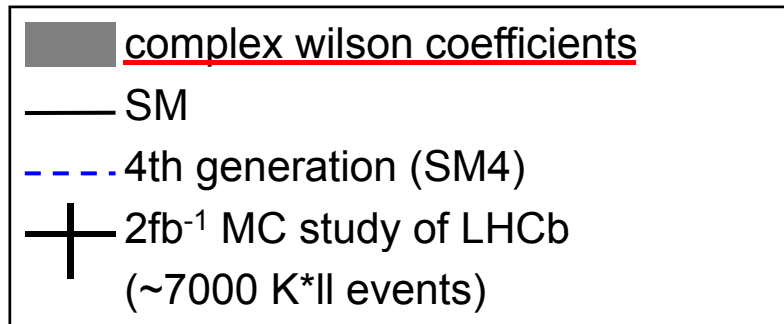
$A_{FB}(B \rightarrow K^*l^+l^-)$ and Other Predictions

sent to Backup

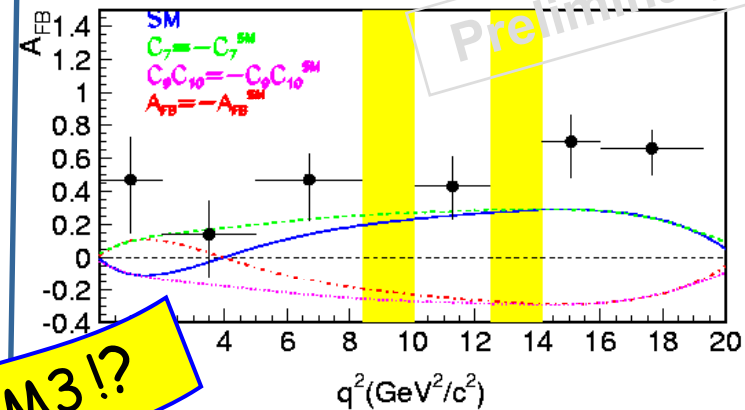
Instead flipped $C_7 \dots$

$$\frac{dA_{FB}}{d\hat{s}} \propto - \left\{ \text{Re}(C_9^{eff} C_{10}) V A_1 + \frac{\hat{m}_b}{\hat{s}} \text{Re}(C_7^{eff} C_{10}) [V T_2 (1 - \hat{m}_{K^*}) + A_1 T_1 (1 + \hat{m}_{K^*})] \right\}$$

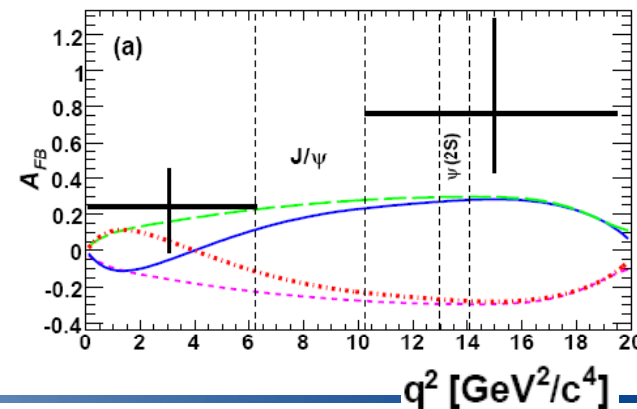
W.-S. Hou, A. Hovhannisyanyan, and N. Mahajan, PRD 77, 014016 (2008)



Belle 657M arXiv:0810.0335



BABAR, arXiv:0804.4412 386M





IV. A $10^{13}+$ Enhancement of CPV for BAU?

If ... KM4



The Abyss: CPV in KM and B.A.U.



The Lore

$$\frac{n_{\bar{B}}}{n_{\gamma}} \cong 0$$

$$\frac{n_B}{n_{\gamma}} = (6.2 \pm 0.2) \times 10^{-10}$$

WMAP

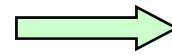
$$\text{KM} \sim 10^{-20}$$

Too Small in SM

Jarlskog Invariant in SM3 (need 3 generation in KM)

$$J = (m_t^2 - m_u^2)(m_t^2 - m_c^2)(m_c^2 - m_u^2)(m_b^2 - m_d^2)(m_b^2 - m_s^2)(m_s^2 - m_d^2) A$$

Normalize by $T \sim 100 \text{ GeV}$



$$J/T^{12} \sim 10^{-20}$$

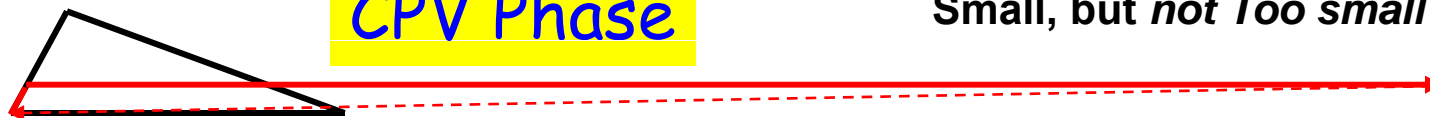
EW Phase Transition Temperature
 $\sim \text{v.e.v.}$

Masses too Small!

$A \sim 3 \times 10^{-5}$ is common (unique) area of triangle ^{in SM}

CPV Phase

Small, but *not Too small*





B.A.U. from CPV in KM



$$\frac{n_{\bar{B}}}{n_{\gamma}} \cong 0$$

$$\frac{n_B}{n_{\gamma}} = (6.2 \pm 0.2) \times 10^{-10}$$

WMAP

KM ~

Enough CPV?

~~Too Small in SM~~

If shift by One Generation in SM4 (need 3 generation in KM)

$$J = (m_t^2 - m_u^2)(m_t^2 - m_c^2)(m_c^2 - m_u^2)(m_b^2 - m_d^2)(m_b^2 - m_s^2)(m_s^2 - m_d^2) A$$

Providence

WSH, arXiv:0803.1234 [hep/ph]

$$J_{(2,3,4)}^{sb} \simeq (m_{t'}^2 - m_c^2)(m_{t'}^2 - m_t^2)(m_t^2 - m_c^2)(m_{b'}^2 - m_s^2)(m_{b'}^2 - m_b^2)(m_b^2 - m_s^2) A_{234}^{sb}$$
$$\sim \frac{m_{t'}^2}{m_c^2} \left(\frac{m_{t'}^2}{m_t^2} - 1 \right) \frac{m_{b'}^4}{m_b^2 m_s^2} \left(\frac{A_{234}^{sb}}{A} \right) J \sim 10^{+15} \text{ Gain}$$

Order 1 ~ 30

Gain mostly in Large Yukawa Couplings!

Nature would likely use this !?

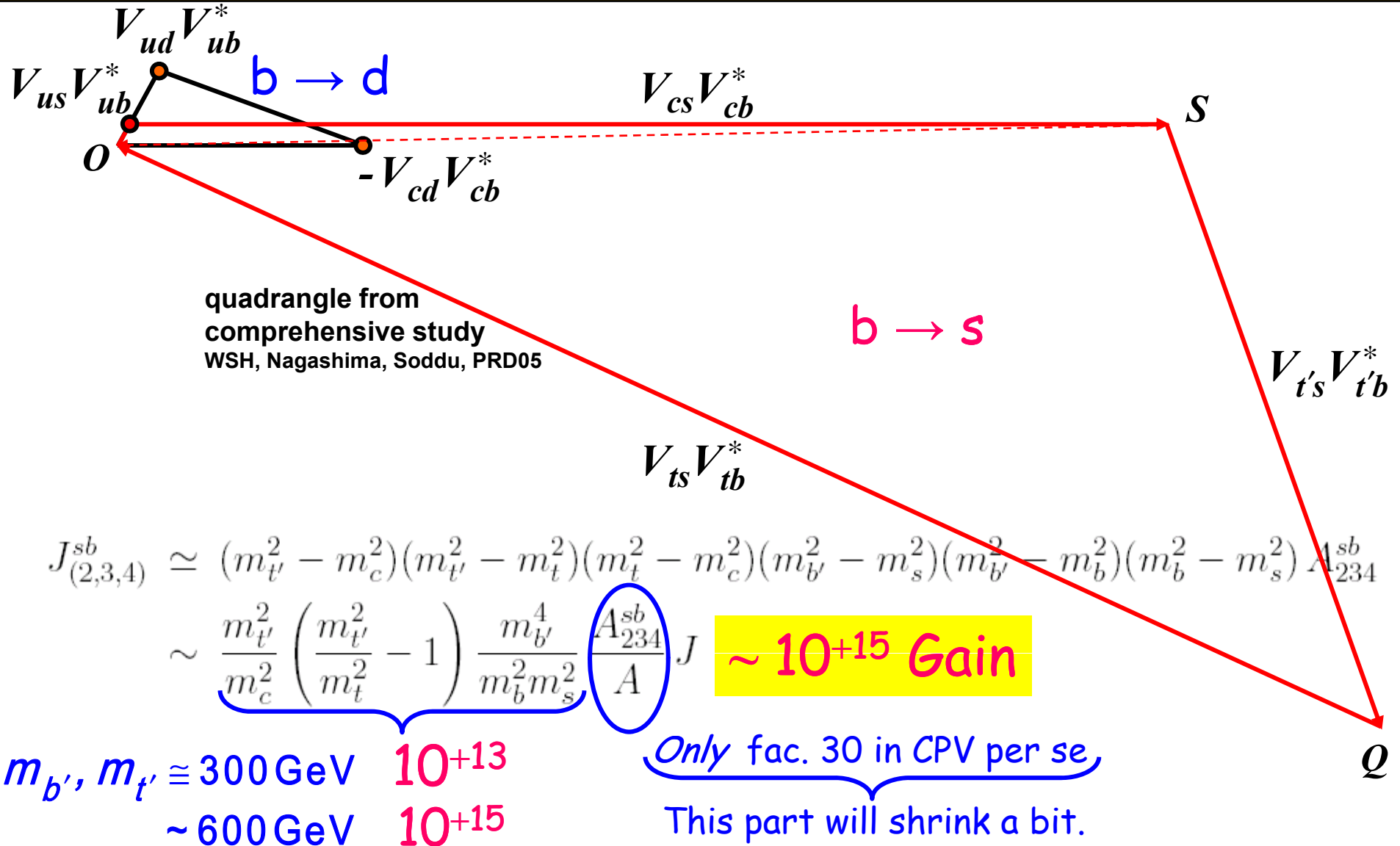


The **Abyss** between CPV in SM3 vs BAU
bridged in SM4 by *Heaviness of t' and b'*

Why wasn't this clearly
pointed out in past 20 years?



Gain mostly in Large Yukawa Couplings !





V. Direct Sighting @ Tevatron vs LHC

the Experimentalist





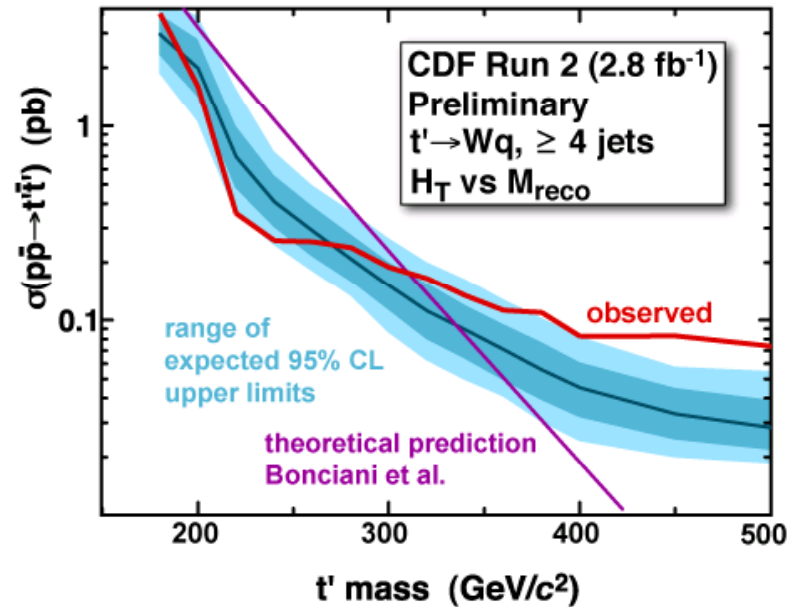
Tevatron/LHC Verification



Tevatron *Unequivocal BSM*
... if true

• $\sin 2\Phi_{B_s}$ “Evidence” by 2009 ?
“Observe” by 2010 ?

• t' Search Ongoing:
 $m_{t'} > 311 \text{ GeV @ 95\% CL}$



LHC

• $\sin 2\Phi_{B_s}$ “Confirmation” — “Easy” for LHCb

But when ?

• b', t' Discovery — Straightforward/full terrain

Glimpse of agenda at ATLAS/CMS

previous 2 talks



VI. Discussion/Conclusion: Know in 3-5 Years



CPV for BAU: 2-3-4 Dominance



Jarlskog'85, 3 generations

$$\text{Im det} \begin{bmatrix} m_u m_u^\dagger & \\ & m_d m_d^\dagger \end{bmatrix}$$

$S \quad S'$



Jarlskog'87, n generations

$$\text{Im tr}[S, S']^3$$

“3 cycles”

also Gronau, Kfir, Loewy '87

4 generations: 3 indep. phases

long and short



d - s degenerate

(on v.e.v. scale)

2-3-4 generation only !

Effectively 3 generations

$$J_{(2,3,4)}^{sb} \simeq (m_{t'}^2 - m_c^2)(m_{t'}^2 - m_t^2)(m_t^2 - m_c^2)(m_{b'}^2 - m_s^2)(m_{b'}^2 - m_b^2)(m_b^2 - m_s^2) A_{234}^{sb}$$

$$\sim \frac{m_{t'}^2}{m_c^2} \left(\frac{m_{t'}^2}{m_t^2} - 1 \right) \frac{m_{b'}^4}{m_b^2 m_s^2} \frac{A_{234}^{sb}}{A} J$$

$J(1,2,3)$ very small

suppressed by m_s, m_c



1st Order EW Phase Trans. for BAU ?



0803.1234 will appear in Chin. J. Phys.

Ran out of time, and knowledge ...

(perturbative)

- Fok & Kribs: Not possible in 4th generation PRD'08
arXiv:0803.4207 [hep-ph]
- Conjecture: Could Strong Yukawa's do it ?

Beyond Unitarity Limit

arXiv:0901.1962v1 [hep-ph]

The strongly coupled fourth family and a first-order electroweak phase transition

(I) quark sector

Not quite conclusive (?)

Yoshio Kikukawa,^{1,*} Masaya Kohda,^{2,†} and Junichiro Yasuda^{3,‡}

¹*Institute of Physics, University of Tokyo Tokyo 153-8092, Japan*

²*Department of Physics, Nagoya University Nagoya 464-8602, Japan*

³*Center for the Studies of Higher Education, Nagoya University Nagoya 464-8601, Japan*

(Dated: January 14, 2009)

In models of dynamical electroweak symmetry breaking due to strongly coupled fourth-family quarks and leptons, their low-energy effective descriptions may involve multiple composite Higgs fields, leading to a possibility that the electroweak phase transition at finite temperature is first order due to the Coleman-Weinberg mechanism. We examine the behavior of the electroweak phase transition based on the effective renormalizable Yukawa theory which consists of the fourth-family quarks and two SU(2)-doublet Higgs fields corresponding to the bilinear operators of the fourth-family quarks with/without imposing the compositeness condition. The strength of the first order



Thoughts on the other 1/2 Nobel Prize



SSB

"for the discovery of the mechanism of spontaneous broken symmetry in subatomic physics"



Photo: University of Chicago

Yoichiro Nambu

🏆 1/2 of the prize

USA

Enrico Fermi Institute,
University of Chicago
Chicago, IL, USA

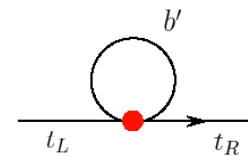
b. 1921
(in Tokyo, Japan)

$\langle \bar{Q}Q \rangle$ can Condense
by Large Yukawa !

Could EWSB be
due to b' and t'
above unitarity bound $\sim 500\text{-}600$ GeV ?

Bob Holdom: [Bardeen, Hill, Lindner

N-J-L



Gustavo Burdman: "Holographic" 4th gen.



VI. Conclusion: Know in 3-5 Years



$$J_{(2,3,4)}^{sb} \simeq (m_{t'}^2 - m_c^2)(m_{t'}^2 - m_t^2)(m_t^2 - m_c^2)(m_{b'}^2 - m_s^2)(m_{b'}^2 - m_b^2)(m_b^2 - m_s^2) A_{234}^{sb}$$

$$\sim \underbrace{\frac{m_{t'}^2}{m_c^2} \left(\frac{m_{t'}^2}{m_t^2} - 1 \right) \frac{m_{b'}^4}{m_b^2 m_s^2}}_{\text{Even if } O(1)} \left(\frac{A_{234}^{sb}}{A} \right) J \quad \sim 10^{+15} \text{ Gain}$$

$$m_{b'}, m_{t'} \cong 300 \text{ GeV} \quad 10^{+13}$$

$$\sim 600 \text{ GeV} \quad 10^{+15}$$

Enough CPV
for B.A.U.

Maybe there is a 4th Generation !

$\sin 2\Phi_{B_s}$
@ Tevatron
by 2010(1)

Will Really Know in $\sim 3-5$ years !

@ LHC



Heaven on Earth?



Backup



$$\Delta A_{K\pi} = A_{B \rightarrow K^+ \pi^0}^{+0.050 \pm 0.025} - A_{B \rightarrow K^+ \pi^-}^{-0.097 \pm 0.012} \neq 0$$

World



$$= +0.147 \pm 0.028 > 5\sigma$$

Experiment is Firm

Why a Puzzle ?

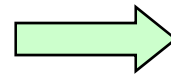
$\Delta A_{K\pi} \sim 0$ expected

$$\mathcal{M}(B^0 \rightarrow K^+ \pi^-) \propto (T + P) = r e^{i\phi_3} + e^{i\delta}$$

$$\sqrt{2} \mathcal{M}_{K^+ \pi^0} - \mathcal{M}_{K^+ \pi^-} \propto (P_{EW} + C) ?$$

$$r = \frac{\text{[Image of a tree]}}{\text{[Image of a penguin]}}$$

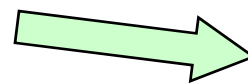
Large C ?



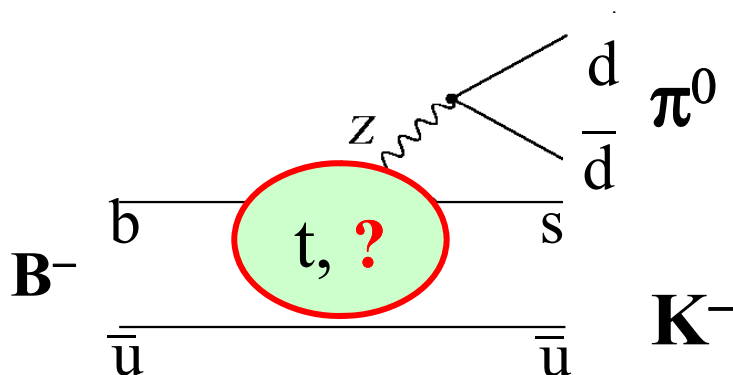
A lot of (hadronic) finesse

Baek, London, PLB653, 249 (2007)

Large EW Penguin ?



Need NP CPV Phase



P_{EW} has practically no weak phase in SM



$$\Delta A_{K\pi} = A_{B \rightarrow K^+ \pi^0}^{+0.050 \pm 0.025} - A_{B \rightarrow K^+ \pi^-}^{-0.097 \pm 0.012} \neq 0$$

World



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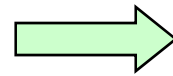
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$$\sqrt{2}\mathcal{M}_{K^+ \pi^0} - \mathcal{M}_{K^+ \pi^-} \propto (P_{EW} + C) ?$$

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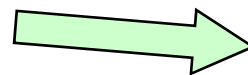
Large C ?



A lot of (hadronic) finesse

Baek, London, PLB653, 249 (2007)

Large EWPenguin ?

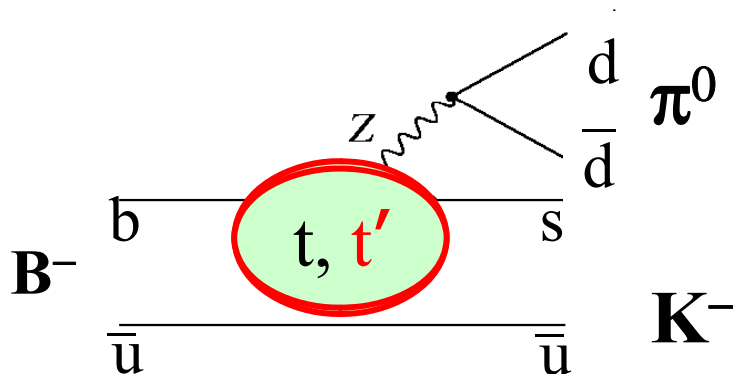


Need NP CPV Phase

P_{EW} has practically no weak phase in SM

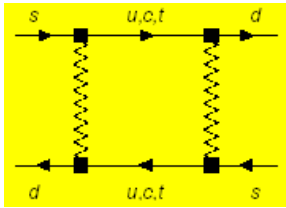
4th Gen. in EWP Natural

nondecouplin

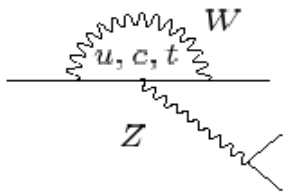




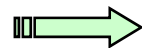
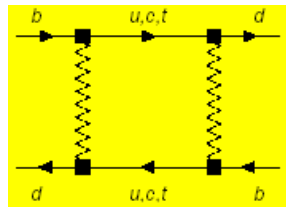
On Boxes and Z Penguins



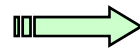
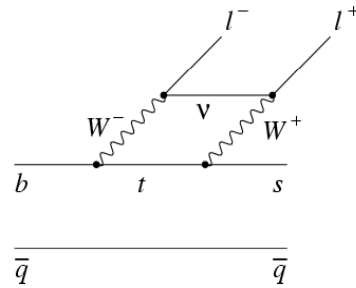
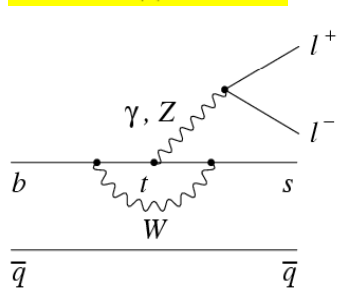
GIM, charm, ϵ_K



small ϵ'/ϵ , $K \rightarrow \pi\nu\nu$ (still waiting)



heavy top, $\sin 2\phi_1/\beta$



Z dominance for heavy top

1986 \rightarrow 2002

Most Flavor/CPV learned from these diagrams/processes

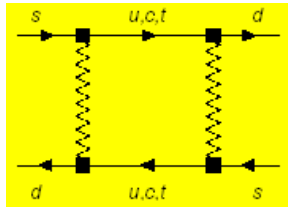


On Boxes and Z Penguins

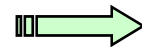
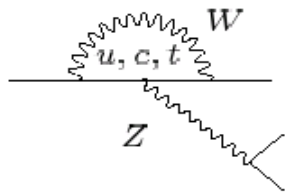


Nondecoupling

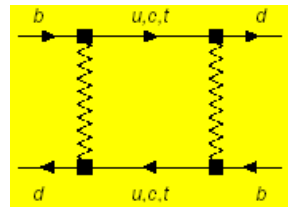
∴ Large Yukawa!



GIM, charm, ϵ_K

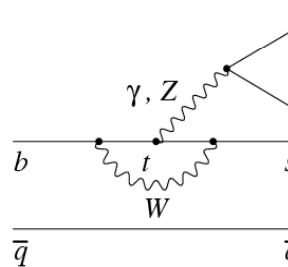


small ϵ'/ϵ , $K \rightarrow \pi\nu\nu$ (still waiting)

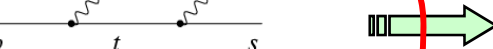


heavy top, $\sin 2\phi_1/\beta$

B_s



A_{FB}



Z dominance for heavy top

1986 → 2002

All w/ 3-generations,
Just wait if there's a 4th

D!

b' , t' @ LHC



4 x 4 Unitarity \Rightarrow Constraints



	d	s	b	b'
u	$c_{12} c_{13} c_{14}$ $-c_{13} s_{12} s_{14} s_{24} \exp[-i(\phi_{db} - \phi_{sb})]$ $-c_{24} s_{13} s_{14} s_{34} \exp[-i(\phi_{db} + \phi_{ub})]$	$c_{13} c_{24} s_{12}$ $-s_{13} s_{24} s_{34} \exp[-i(\phi_{sb} + \phi_{ub})]$	$c_{34} s_{13} \exp[-i\phi_{ub}]$	$c_{12} c_{13} s_{14} \exp[i\phi_{db}]$ $+c_{13} c_{14} s_{12} s_{24} \exp[i\phi_{sb}]$ $+c_{14} c_{24} s_{13} s_{34} \exp[-i\phi_{ub}]$
c	$-c_{14} c_{23} s_{12}$ $-c_{12} c_{14} s_{13} s_{23} \exp[i\phi_{ub}]$ $-c_{12} c_{23} s_{14} s_{24} \exp[-i(\phi_{db} - \phi_{sb})]$ $+s_{12} s_{13} s_{14} s_{23} s_{24} \exp[-i(\phi_{db} - \phi_{sb} - i\phi_{ub})]$ $-c_{13} c_{24} s_{14} s_{23} s_{34} \exp[-i\phi_{db}]$	$c_{12} c_{23} c_{24}$ $-c_{24} s_{12} s_{13} s_{23} \exp[i\phi_{ub}]$ $-c_{13} s_{23} s_{24} s_{34} \exp[-i\phi_{sb}]$	$c_{13} c_{34} s_{23}$	$-c_{23} s_{12} s_{14} \exp[i\phi_{db}]$ $-c_{12} s_{13} s_{14} s_{23} \exp[i(\phi_{db} + \phi_{ub})]$ $+c_{12} c_{14} c_{23} s_{24} \exp[i\phi_{sb}]$ $-c_{14} s_{12} s_{13} s_{23} s_{24} \exp[i(\phi_{sb} + \phi_{ub})]$ $+c_{13} c_{14} c_{24} s_{23} s_{34}$
t	$-c_{12} c_{14} c_{23} s_{13} \exp[i\phi_{ub}]$ $+c_{14} s_{12} s_{23}$ $+c_{23} s_{12} s_{13} s_{14} s_{24} \exp[-i(\phi_{db} - \phi_{sb} - i\phi_{ub})]$ $+c_{12} s_{14} s_{23} s_{24} \exp[-i(\phi_{db} - \phi_{sb})]$ $-c_{13} c_{23} c_{24} s_{14} s_{34} \exp[-i\phi_{db}]$	$-c_{23} c_{24} s_{12} s_{13} \exp[i\phi_{ub}]$ $-c_{12} c_{24} s_{23}$ $-c_{13} c_{23} s_{24} s_{34} \exp[i\phi_{sb}]$	$c_{13} c_{23} c_{34}$	$-c_{12} c_{23} s_{13} s_{14} \exp[i(\phi_{db} + \phi_{ub})]$ $+s_{12} s_{14} s_{23} \exp[i\phi_{db}]$ $-c_{14} c_{23} s_{12} s_{13} s_{24} \exp[i(\phi_{sb} + \phi_{ub})]$ $-c_{12} c_{14} s_{23} s_{24} \exp[i\phi_{sb}]$ $+c_{13} c_{14} c_{23} c_{24} s_{34}$
t'	$-c_{24} c_{34} s_{14} \exp[-i\phi_{db}]$	$-c_{34} s_{24} \exp[-i\phi_{sb}]$	$-s_{34}$	$c_{14} c_{24} c_{34}$

SM3

We need to deal with mixing matrix in detail to keep **Unitarity**

$$V_{t's}^* V_{t'd} = c_{24} c_{34}^2 s_{14} s_{24} e^{i(\phi_{sb} - \phi_{db})}$$

Kaon $\equiv r_{ds} e^{i\phi_{ds}}$

$$V_{t's}^* V_{t'b} = c_{34} s_{24} s_{34} e^{i\phi_{sb}}$$

$b \rightarrow s \equiv r_{sb}$

$$V_{t'd}^* V_{t'b} = c_{24} c_{34} s_{14} s_{34} e^{i\phi_{db}} = \frac{r_{ds} s_{34}^2}{r_{sb}} e^{i\phi_{db}}$$

$b \rightarrow d$

Cross Check!

$\Gamma(Z \rightarrow \text{hadrons})$

impose $s_{34} = 0.22 \simeq V_{us}$

$$|V_{tb}|^2 + 3.4|V_{t'b}|^2 < 1.14 \text{ for } m_{t'} = 300 \text{ GeV} \Rightarrow s_{34} < 0.25$$

From $b \rightarrow s$ study

$$r_{sb} e^{i\phi_{sb}} \simeq 0.025 e^{i70^\circ}$$



Constrain $s \leftrightarrow d$ from K Physics



$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (14.7^{+13.0}_{-8.9}) \cdot 10^{-11}$$

(shaded)

$$BR(K_L \rightarrow \mu^+ \mu^-)_{SD} < 3.75 \cdot 10^{-9}$$

$$\epsilon_K = (2.284 \pm 2 \times 0.014) \cdot 10^{-3}$$

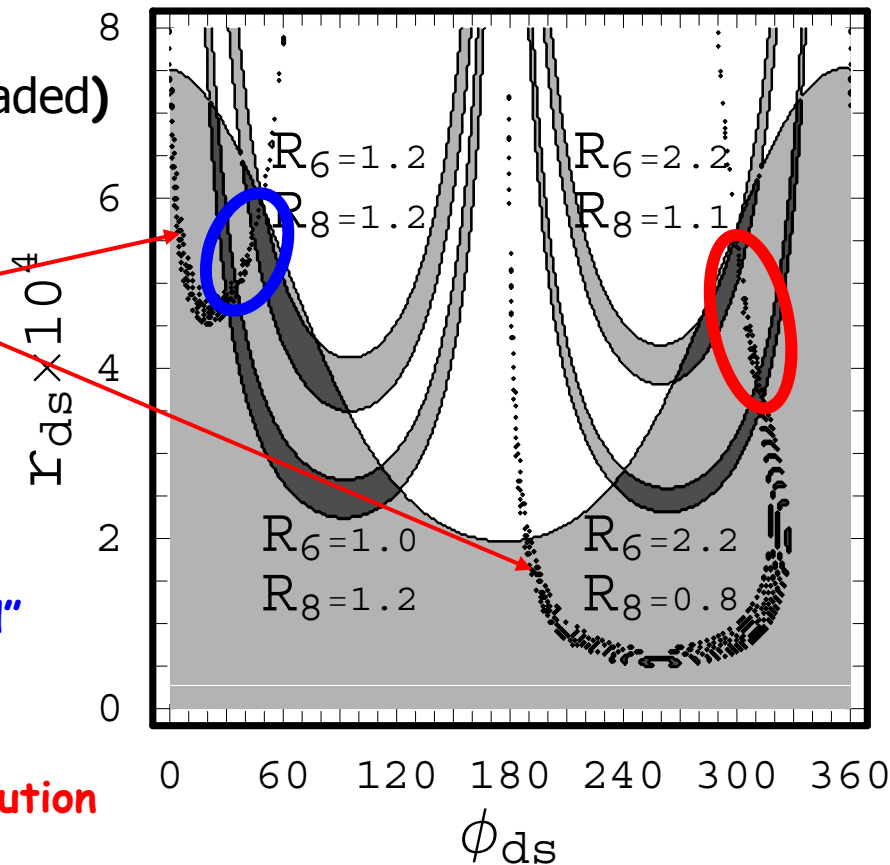
$$\frac{\epsilon'}{\epsilon} = (16.6 \pm 2 \times 1.6) \cdot 10^{-4}$$

$$R_6 = 1.2 \quad (\text{E. Pallante et al.})$$

$$R_8 = 0.7 - 1.3 \quad \text{"Standard"}$$

$$R_6 = 2.2 \quad (\text{J. Bijnens et al.})$$

$$R_8 = 0.8 - 1.4 \quad \text{No SM3 solution}$$



Therefore....

$$r_{ds} \sim 5 \times 10^{-4}, \quad \phi_{ds} \sim -60^\circ \text{ or } +35^\circ$$

well-satisfy Δm_{B_d} and $\sin 2\phi_1$!



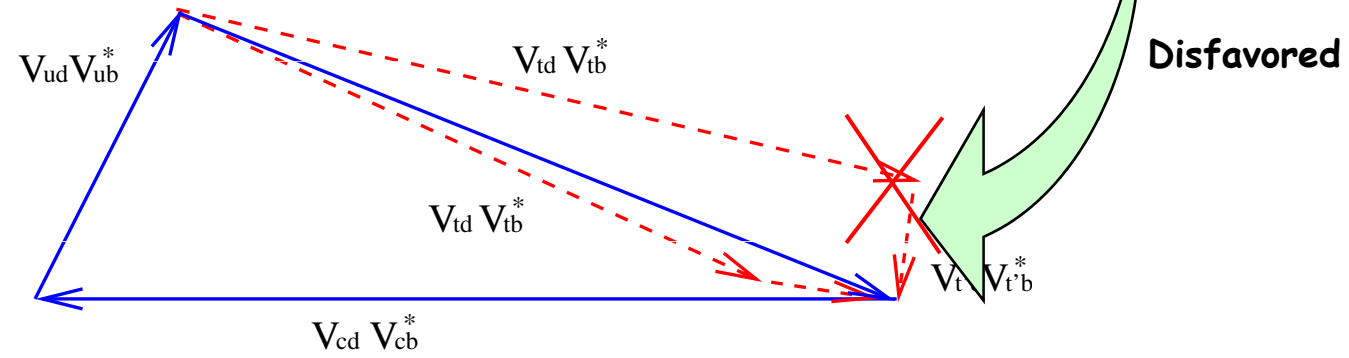
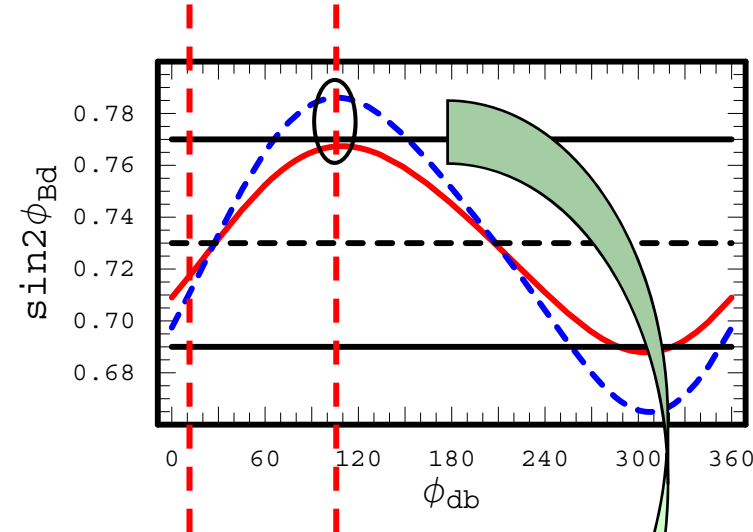
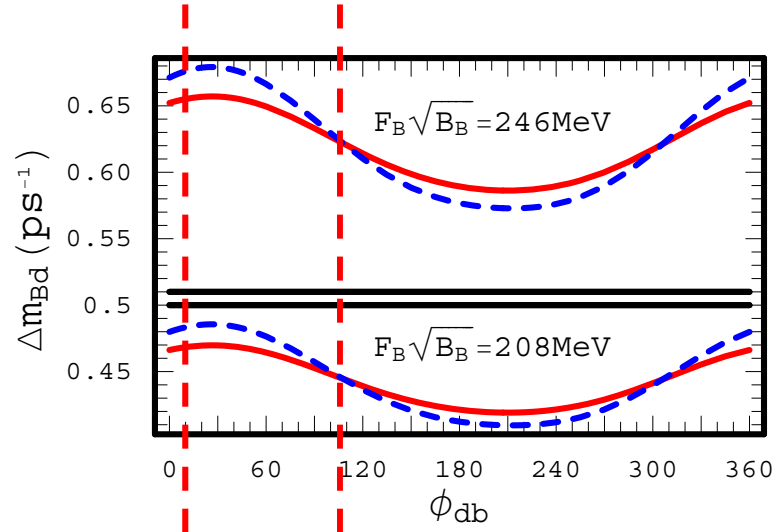
$$r_{ds} \sim 5 \times 10^{-4}, \quad \phi_{ds} \sim -60^\circ \text{ or } +35^\circ$$

$$r_{db} \sim 1 \times 10^{-3}, \quad \phi_{db} \sim 10^\circ (105^\circ)$$



well-satisfy Δm_{B_d} and $\sin 2\phi_1$

vs $V_{ub} \sim 0.01 e^{-i\gamma}$



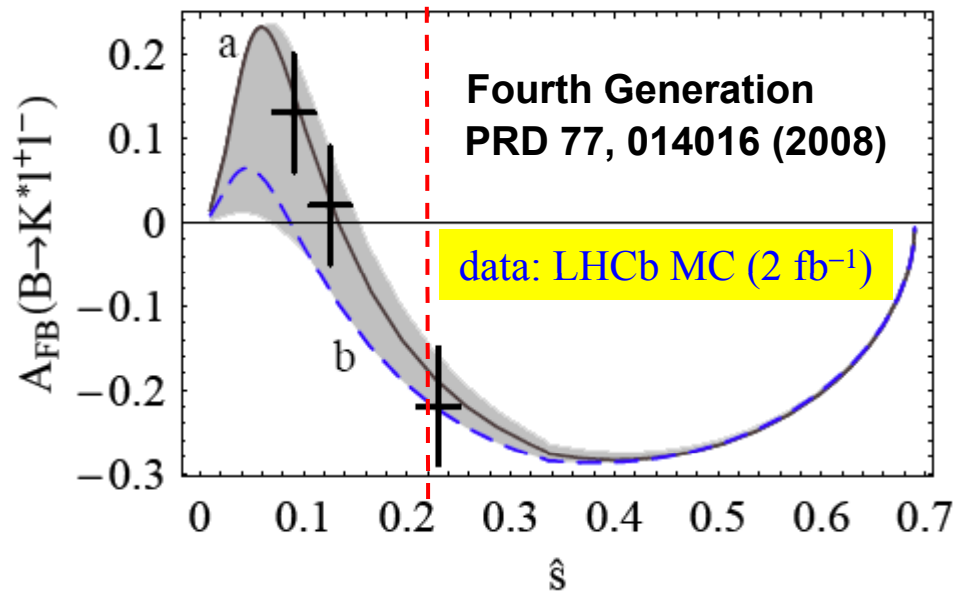
Hard to tell apart (**non-trivial**) with present precision
 \therefore stringent $s \rightarrow d$



$A_{FB}(B \rightarrow K^*l^+l^-)$ and Other Predictions

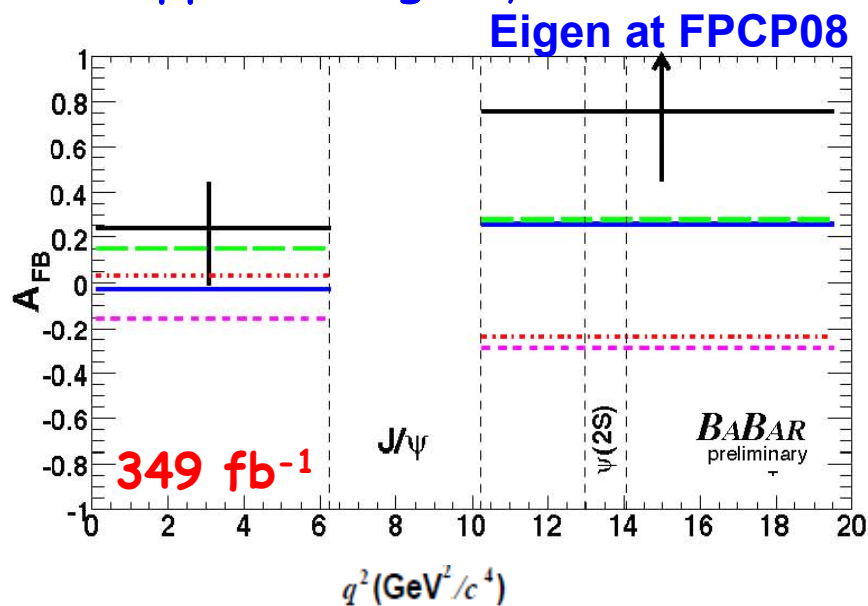
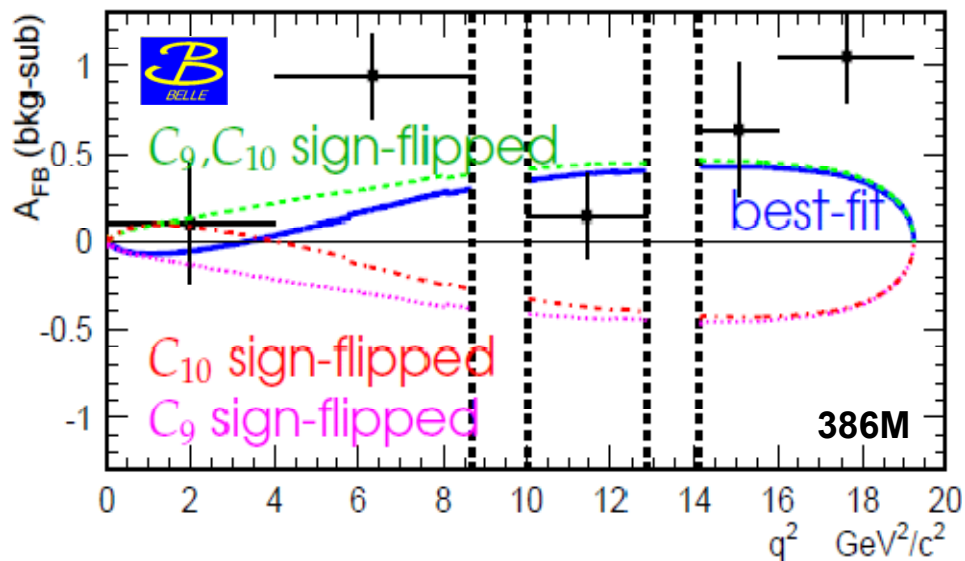
sent to Backup

Quoted by **Tsybychev at FPCP08**



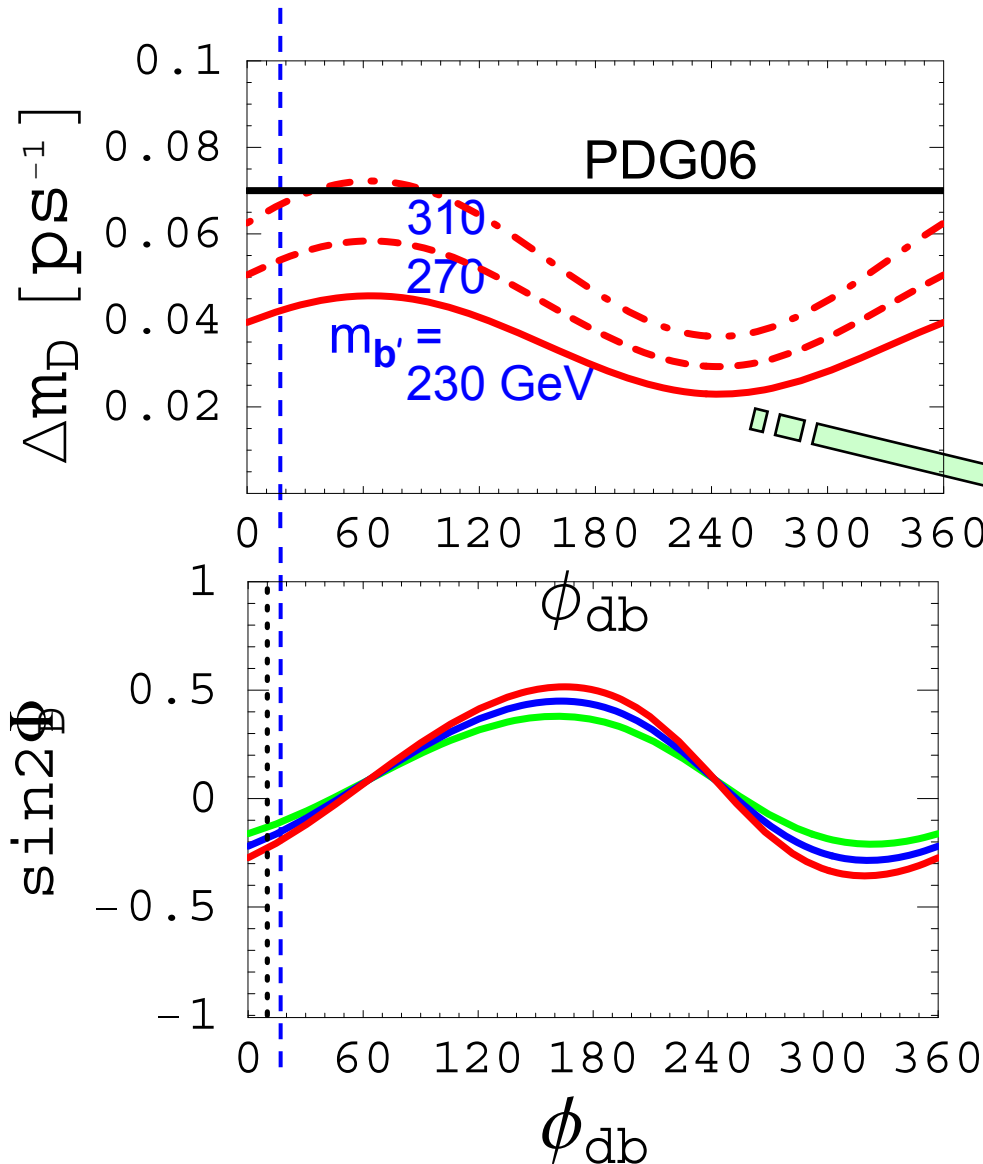
a: SM; b: **4 Gen.**
better

● (F_L and) A_{FB} (and A_I) favor the "opposite-sign C_7 model"





D Mixing (Short-distance Only)



$$f_D \sqrt{B_D} = 200 \text{ MeV}$$

$$V_{t'd}^* V_{t'b} \equiv r_{db} e^{i\phi_{db}}$$

From 4 x 4 Unitarity

$$V_{ub'} V_{cb'}^*$$

$x = \Delta m/\Gamma \sim 1 - 3$ plausible

w/ Sizable (but not huge)
CPV in Mixing $\sim -15\%$

N.B. SM LD could generate

$y \sim 1\%$, $x \approx y$

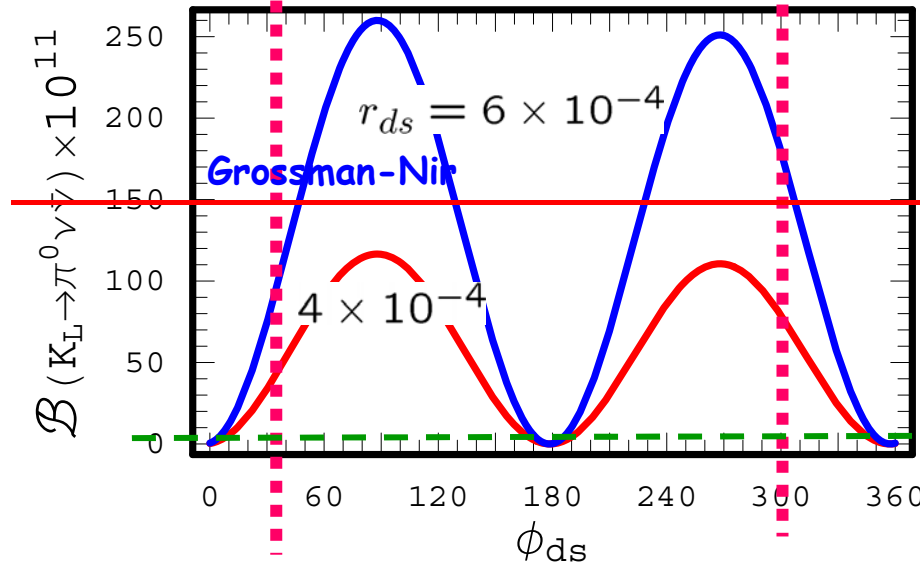
[Falk, Grossman, Ligeti, (Nir,) Petrov]



Implication for $\mathcal{B}(K_L \rightarrow \pi^0 \nu \bar{\nu})$



Nontrivial (phase) $V_{t'd}^* V_{t's}$



Current E391A U.L.

2.86×10^{-7} (90% c.l.)

Very hard to measure

$\mathcal{B}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \simeq 3 \times 10^{-11}$

SM3

Rate could be enhanced by up to almost two orders !!

$K_L \rightarrow \pi^0 \nu \bar{\nu}$ enhanced to 5×10^{-10} or even higher !!

In general larger than $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ ($2-3 \times 10^{-10}$)

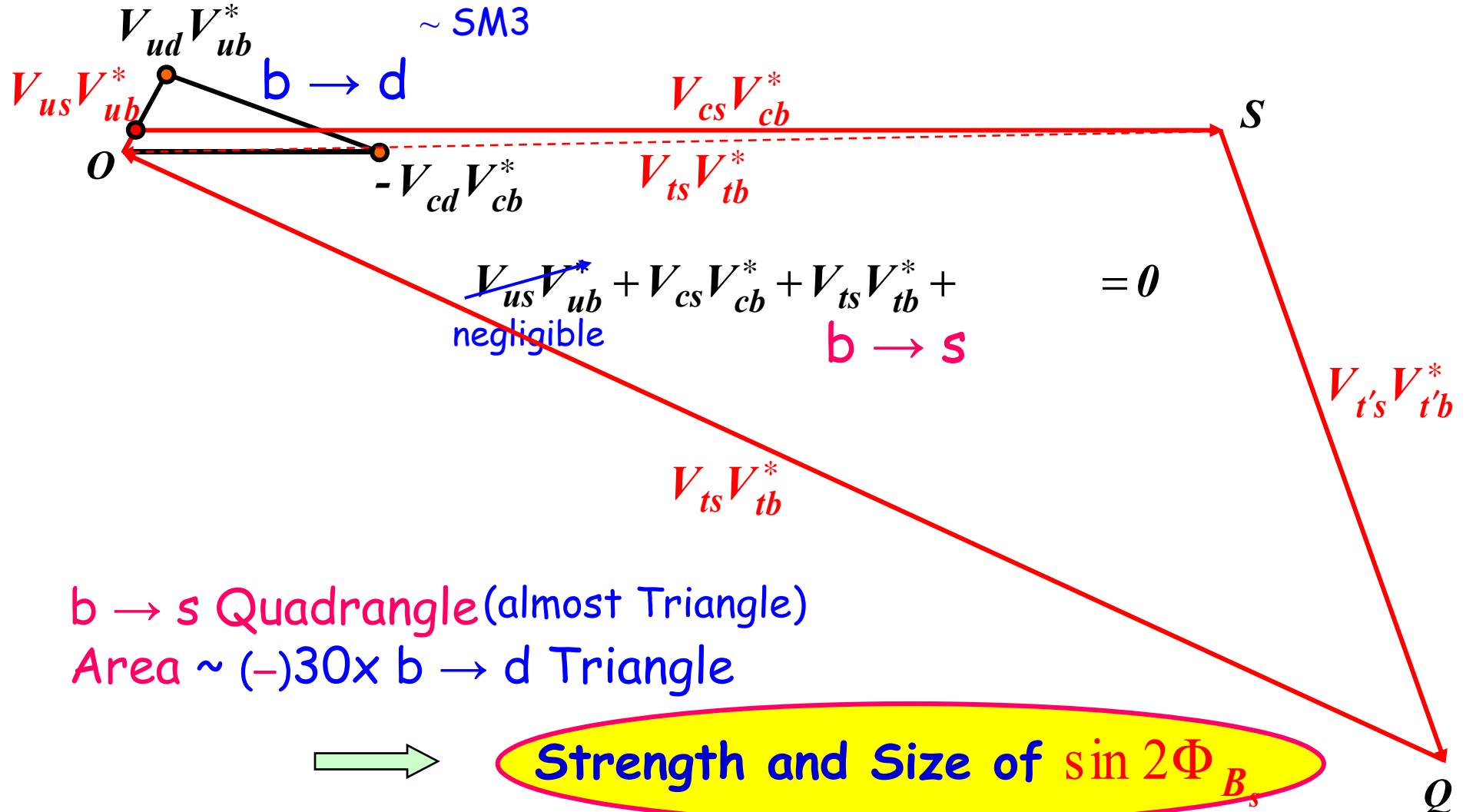
\therefore Large CPV Phase



$b \rightarrow d$ "Triangle" and $b \rightarrow s$ Quadrangle

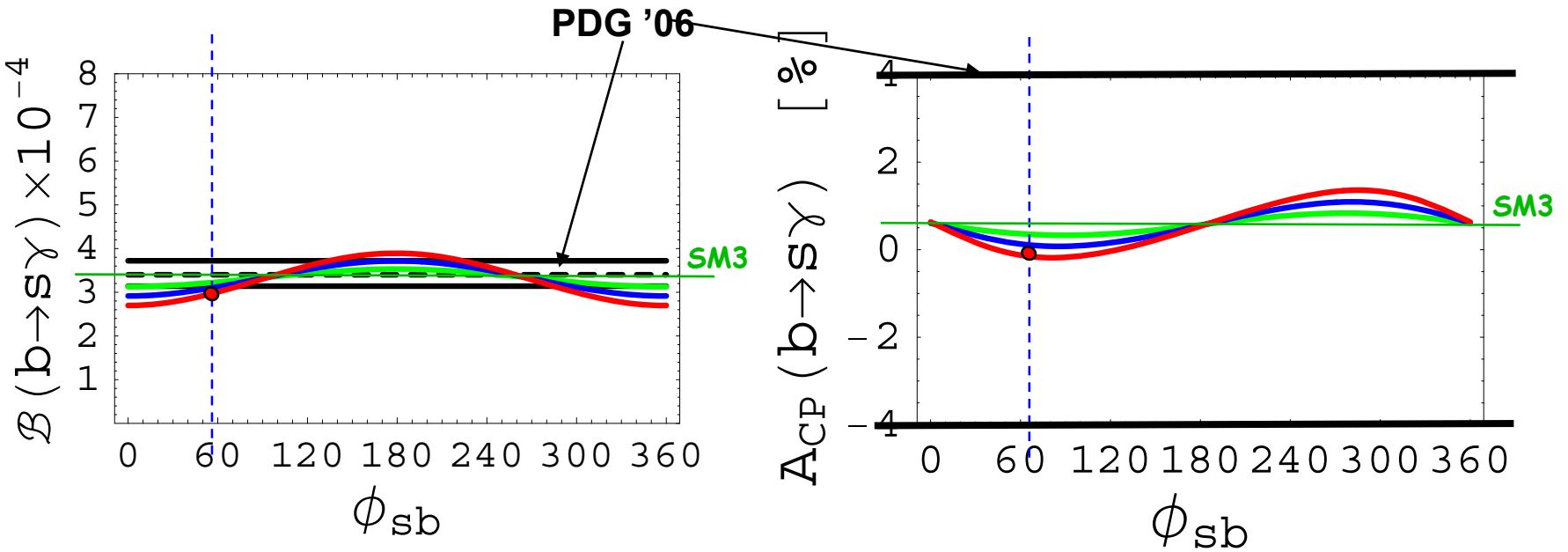


$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + \underbrace{V_{td}V_{tb}^* + V_{t'd}V_{t'b}^*}_{\sim \text{SM3}} = 0$$





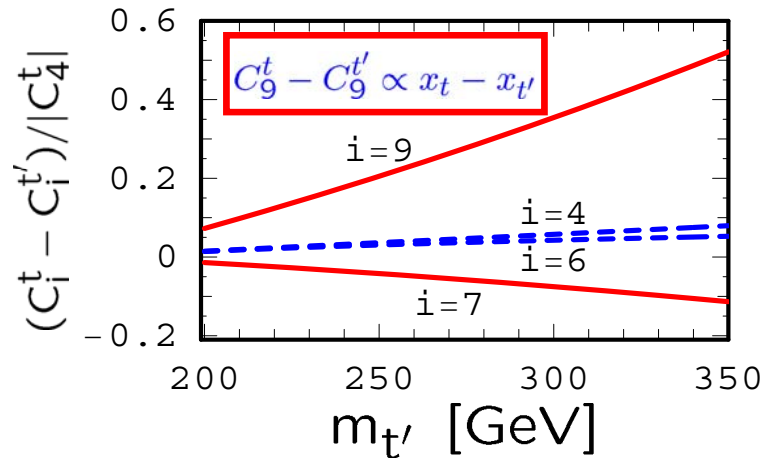
Consistency and $b \rightarrow s\gamma$ Predictions



BR OK

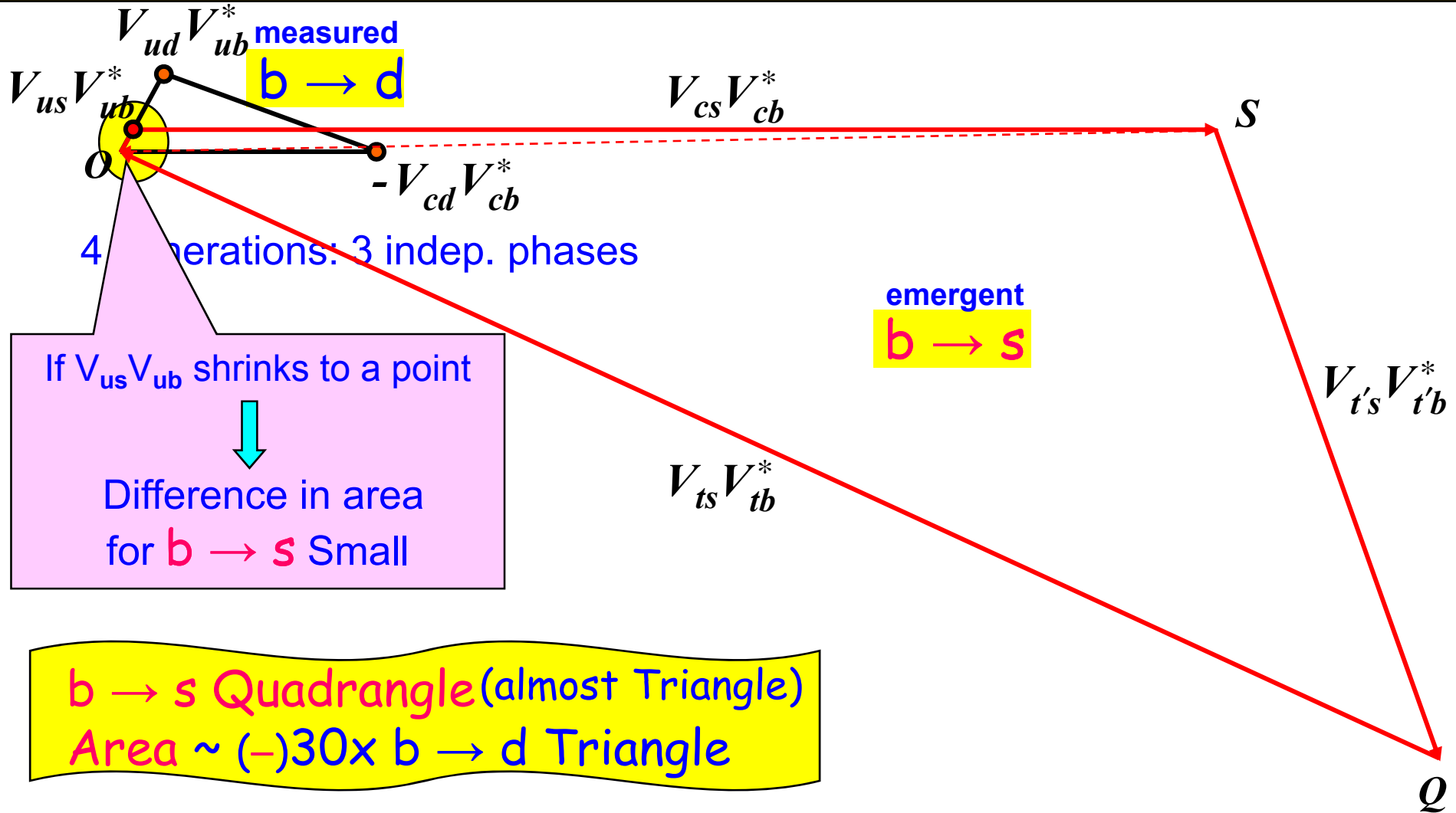
**$A_{CP} \sim 0$ far away
beyond SuperB**

Heavy t' effect
decoupled
for $b \rightarrow s\gamma$





4 generations: 3 indep. phases



2nd argument that $J_{(2,3,4)}^{sb}$ is predominant CPV



Sighting

Vision ~ Early '06



4th generation? — The jury is out ...

In era of LHC, can **Directly Search for b' , t'**
Once and For All !

Find b' , t' , or Rule Out @ LHC

It's a Duty.

Strategy Considerations (漢中策略)

- **Well shielded** training ground — **All Tools**
 - ☞ Move on to Greener Pastures ~ in 2 years
- **Publish early** — Large Cross Section
 - If “Limits”, then easy to publish
 - If “Signal”, Lucked Out!



b' Signatures



For $m_{b'} < m_t + M_W = 255 \text{ GeV}$

$b' \rightarrow cW$ dominance for sizable $V_{cb'}$
 $b' \rightarrow tW^*$ dominance for suppressed $V_{cb'}$



Kinematic suppressed for $m_{b'} \lesssim 230 \text{ GeV}$

Initial discovery should consider

$b' \rightarrow cW \sim b' \rightarrow bZ, bH \sim b' \rightarrow tW^*$

Rich Signature

$cc(\text{bar})WW; cWbZ; cWbH;$
 $tc(\text{bar})WW^*;$
 $tt(\text{bar})W^*W^*; tW^*bZ; tW^*bH;$

Bonus !!

For $m_{b'} > m_t + M_W = 255 \text{ GeV}$

$b' \rightarrow tW$ dominance; FCNC searchable

$tt(\text{bar})WW \rightarrow bb(\text{bar})W^+W^-W^+W^-$

Heavy Q related To EWSB ?

4 W's + 2b's



Available on the CMS information server

CMS PAS EXO-08-09

CMS Physics Analysis Summary

2008/08/29

Search for Heavy Bottom-like **Fourth Generation Quark**
Pair at CMS in pp Collisions at $\sqrt{s} = 14$ TeV

The CMS collaboration



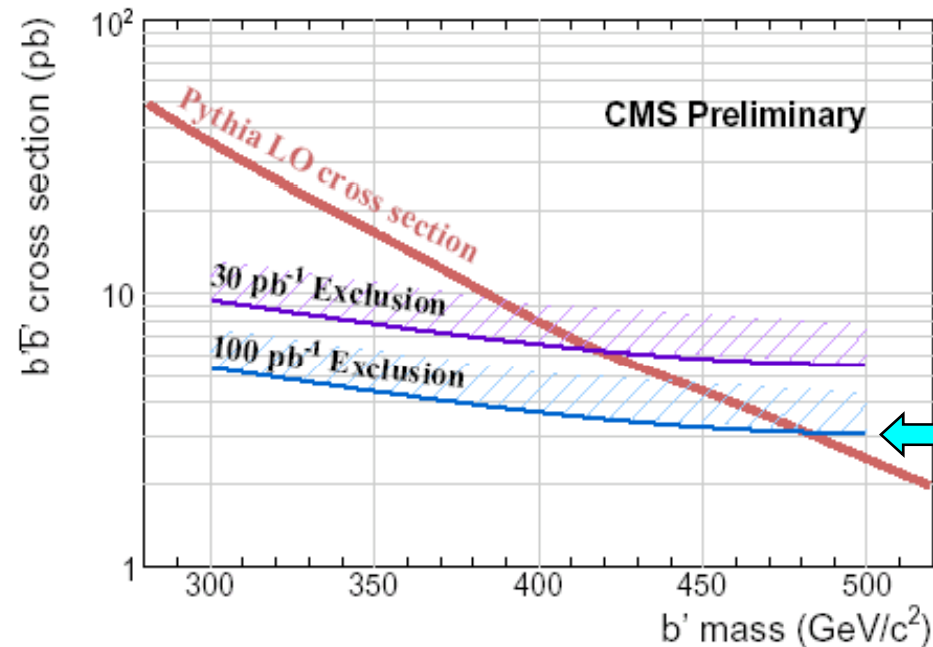
$$pp \rightarrow b'\bar{b}' \rightarrow t\bar{t}W^+W^-$$

100 pb⁻¹



same-sign dilepton and trilepton

b' Mass	300 GeV/ c^2	400 GeV/ c^2	500 GeV/ c^2
$b'\bar{b}'$ LO cross section	34.9 pb	8.05 pb	2.45 pb
Expected signal yield	68.2	22.2	8.0
Expected background yield		7.3 ^{+10.5} _{-4.8}	
S_{12}	7.5 σ	2.0 σ	0.0 σ
S_{CP}	N/A	2.1 σ	0.0 σ



Limit to 480 GeV
w/ 100 pb⁻¹