# Low Energy Analysis of $vN \rightarrow vN\gamma$

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- theoretical significance of neutrino-photonbaryon interactions
- chiral lagrangian and extrapolation
- some fits to MiniBooNE data

RJH: arXiv: 0905.0291

Harvey, Hill & Hill Phys.Rev.Lett..99, 261601 (2007); Phys.Rev. D77, 085017 (2008)

# Single photon emission in neutrino-nucleus collisions probe interesting physics

baryon anomaly. anomaly in baryon current in presence of  $SU(2)_L \times U(1)_Y \Rightarrow$  interactions like

$$\mathcal{L} \sim \epsilon^{\mu\nu\rho\sigma} \omega_{\mu} Z_{\nu} F_{\rho\sigma}$$

Harvey, Hill & Hill 2007

( like anomaly in axial current in the presence of  $U(I)_{EM} \Rightarrow \pi_0 \rightarrow \gamma \gamma$  )

skyrmion excitations. proximity of  $\Delta$  resonance leads to interesting effects: coherent-resonant phenomena, nuclear superradiance,...

Important background to  $v_e$  appearance experiments.

Applications beyond laboratory neutrino experiments

- parity violation (anapole moment at finite baryon density)
- astrophysics (neutron star cooling, supernova dynamics)
- axion interactions

Disentangling multiple effects requires systematic description ⇒ chiral lagrangian at low energy

 $U(x) = \exp\left[i\pi(x)/f_{\pi}\right]$ 

#### Some unfinished business in the baryon chiral lagrangian

"In order to avoid complications due to anomalies we disregard the isoscalar vector, axialvector and pseudoscalar currents." Gasser, Sainio and Svarc, 1988

Need full SU(2)<sub>L</sub>xU(1)<sub>Y</sub> in SU(2)<sub>L</sub>xSU(2)<sub>R</sub>xU(1)<sub>Y</sub>

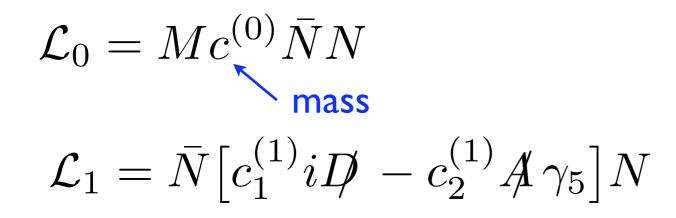
$$\begin{split} \psi_L &\to e^{i\epsilon_L \psi_L} \\ \psi_R &\to e^{i\epsilon_R \psi_R} \quad U(x) = \xi(x)^2 \quad \xi \to e^{i\epsilon_L} \xi e^{-i\epsilon'} = e^{i\epsilon'} \xi e^{-i\epsilon_R} \end{split}$$

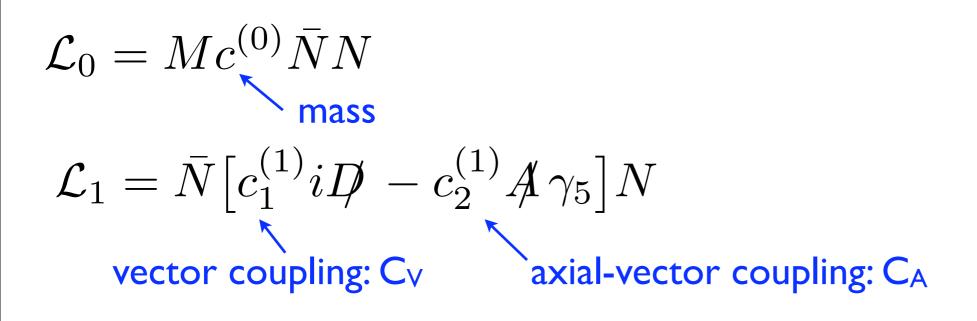
Represent nucleons as isodoublet spinor field, must have peculiar transformation under full  $SU(2)_L x SU(2)_R x U(1)_V$ 

$$N = \begin{pmatrix} p \\ n \end{pmatrix} \qquad \qquad N \to e^{i\epsilon'_{isovector} + 3i\epsilon'_{isoscalar}} N$$

 $\mathcal{L}_0 = M c^{(0)} \bar{N} N$ 

 $\mathcal{L}_0 = M c^{(0)} \bar{N} N$ 





 $\begin{aligned} \mathcal{L}_{0} &= M c_{\text{NN}}^{(0)} \bar{N} N \\ &\searrow \text{mass} \end{aligned}$   $\begin{aligned} \mathcal{L}_{1} &= \bar{N} \big[ c_{1}^{(1)} i D - c_{2}^{(1)} \mathcal{A} \gamma_{5} \big] N \\ &\texttt{vector coupling: } \mathbf{C}_{\mathsf{V}} \qquad \texttt{axial-vector coupling: } \mathbf{C}_{\mathsf{A}} \end{aligned}$   $\begin{aligned} \mathcal{L}_{2} &= \frac{1}{M} \bar{N} \big[ - c_{1}^{(2)} \frac{i}{2} \sigma^{\mu\nu} \mathrm{Tr}([iD_{\mu}, iD_{\nu}]) \big] N \end{aligned}$ 

 $\mathcal{L}_{0} = Mc_{\text{mass}}^{(0)} \bar{N}N$   $\mathcal{L}_{1} = \bar{N} \begin{bmatrix} c_{1}^{(1)}iD & -c_{2}^{(1)}A\gamma_{5} \end{bmatrix} N$ vector coupling: Cv
axial-vector coupling: CA  $\mathcal{L}_{2} = \frac{1}{M} \bar{N} \begin{bmatrix} -c_{1}^{(2)}\frac{i}{2}\sigma^{\mu\nu} \operatorname{Tr}([iD_{\mu}, iD_{\nu}]) \end{bmatrix} N$ 

anomalous magnetic moment: a<sub>N</sub>

$$\mathcal{L}_{0} = Mc_{\text{mass}}^{(0)} \bar{N}N$$

$$\max_{\text{mass}}$$

$$\mathcal{L}_{1} = \bar{N} \Big[ c_{1}^{(1)} i D - c_{2}^{(1)} \mathcal{A} \gamma_{5} \Big] N$$
vector coupling:  $C_{\text{V}}$  axial-vector coupling:  $C_{\text{A}}$ 

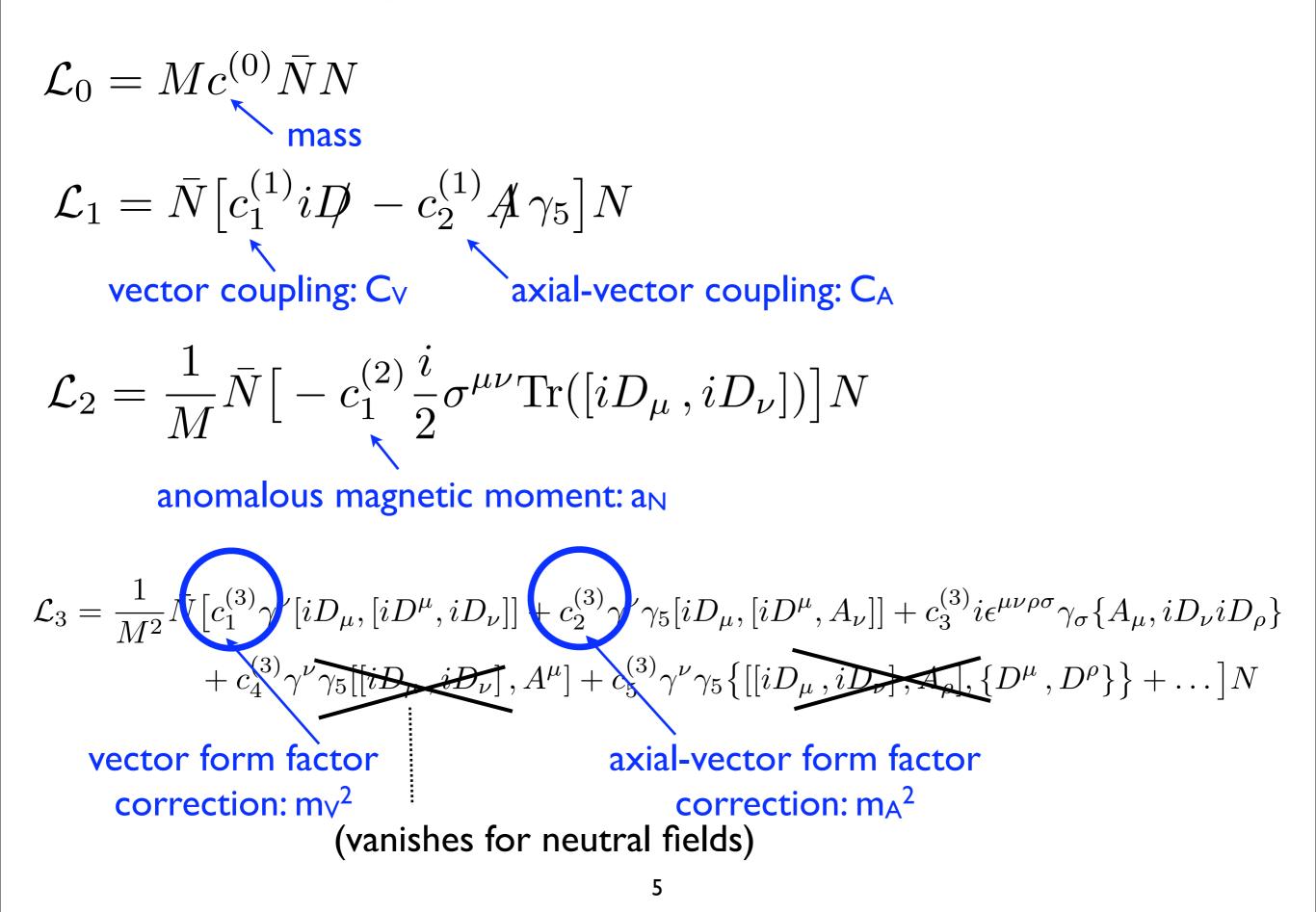
$$\mathcal{L}_{2} = \frac{1}{M} \bar{N} \Big[ -c_{1}^{(2)} \frac{i}{2} \sigma^{\mu\nu} \text{Tr}([iD_{\mu}, iD_{\nu}]) \Big] N$$

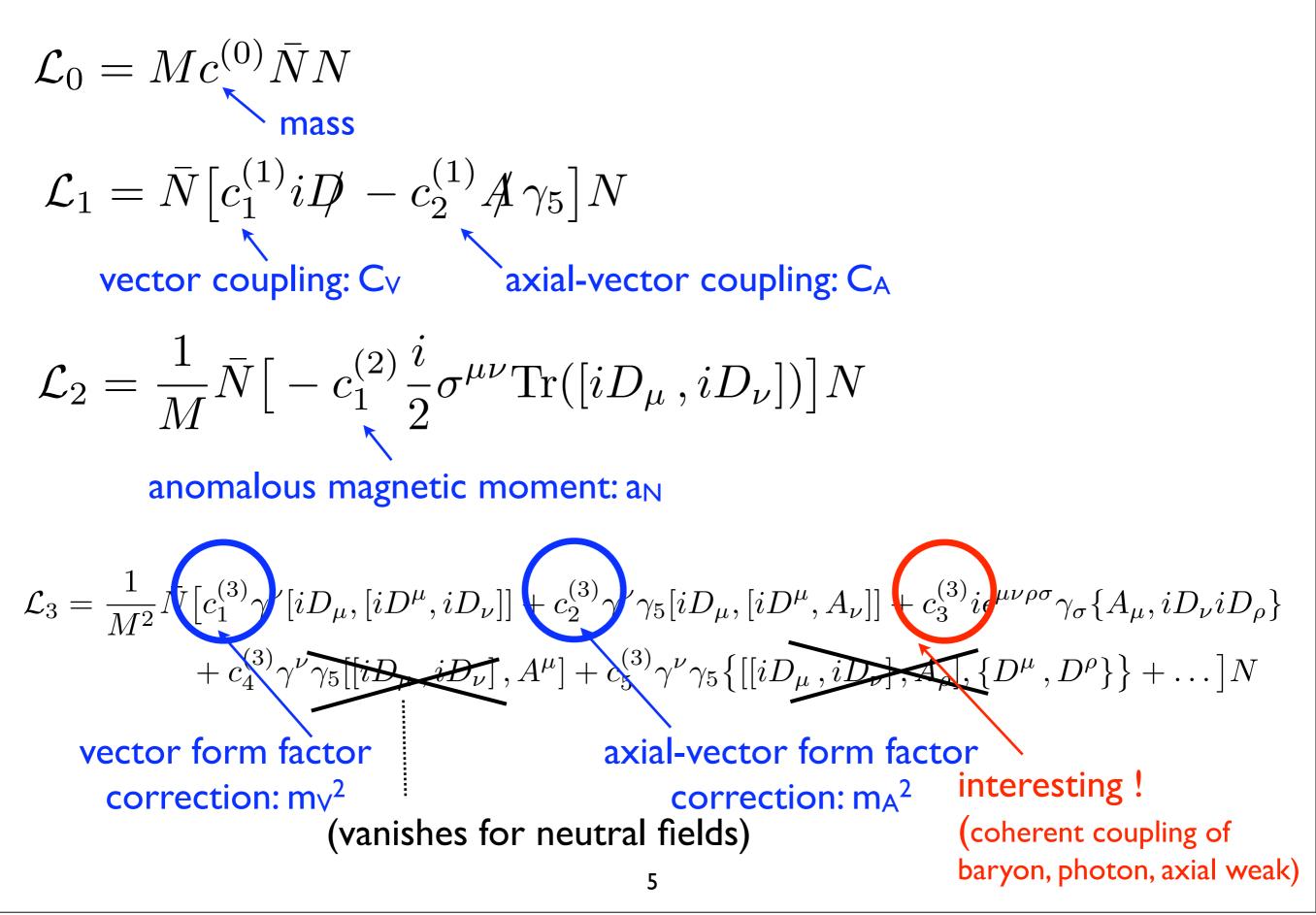
#### anomalous magnetic moment: a<sub>N</sub>

$$\mathcal{L}_{3} = \frac{1}{M^{2}} \bar{N} [c_{1}^{(3)} \gamma^{\nu} [iD_{\mu}, [iD^{\mu}, iD_{\nu}]] + c_{2}^{(3)} \gamma^{\nu} \gamma_{5} [iD_{\mu}, [iD^{\mu}, A_{\nu}]] + c_{3}^{(3)} i \epsilon^{\mu\nu\rho\sigma} \gamma_{\sigma} \{A_{\mu}, iD_{\nu} iD_{\rho}\} + c_{4}^{(3)} \gamma^{\nu} \gamma_{5} [[iD_{\mu}, iD_{\nu}], A^{\mu}] + c_{5}^{(3)} \gamma^{\nu} \gamma_{5} \{[iD_{\mu}, iD_{\nu}], A_{\rho}], \{D^{\mu}, D^{\rho}\}\} + \dots ]N$$

$$\begin{aligned} \mathcal{L}_{0} &= Mc_{0}^{(0)}\bar{N}N\\ &\text{mass}\\ \mathcal{L}_{1} &= \bar{N} \Big[ c_{1}^{(1)}iD - c_{2}^{(1)}\mathcal{A}\gamma_{5} \Big]N\\ &\text{vector coupling: } \mathbb{C}_{\vee} \quad \text{axial-vector coupling: } \mathbb{C}_{\mathsf{A}}\\ \mathcal{L}_{2} &= \frac{1}{M}\bar{N} \Big[ -c_{1}^{(2)}\frac{i}{2}\sigma^{\mu\nu}\mathrm{Tr}([iD_{\mu},iD_{\nu}]) \Big]N\\ &\text{anomalous magnetic moment: } \mathbf{a}_{\mathsf{N}}\\ \mathcal{L}_{3} &= \frac{1}{M^{2}}\bar{N} [c_{1}^{(3)}\gamma^{\nu}[iD_{\mu},[iD^{\mu},iD_{\nu}]] + c_{2}^{(3)}\gamma^{\nu}\gamma_{5}[iD_{\mu},[iD^{\mu},A_{\nu}]] + c_{3}^{(3)}i\epsilon^{\mu\nu\rho\sigma}\gamma_{\sigma}\{A_{\mu},iD_{\nu}iD_{\rho}\}\\ &+ c_{4}^{(3)}\gamma^{\nu}\gamma_{5}[[iD_{\mu},iD_{\nu}]] + c_{5}^{(3)}\gamma^{\nu}\gamma_{5}\{[[iD_{\mu},iD_{\nu},A_{\nu}]] + c_{5}^{(3)}\gamma^{\nu}\gamma_{5}\{[[iD_{\mu},iD_{\nu}]\} + \dots ]N \end{aligned}$$

(vanishes for neutral fields)





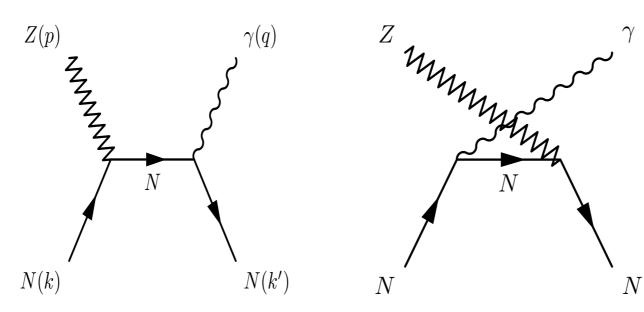
# Phenomenology

- unfortunately, convergence of the chiral lagrangian is poor for energies above a few hundred MeV
- $E_{v}$ -GeV not a great regime for precise calculation, but can perform phenomenological extrapolation to moderate energy
- include dominant resonances in each channel
- nuclear effects: Fermi motion, Pauli blocking not dominant effects, but should be systematically included; in-medium modifications - can compare to e.g. Compton scattering (won't discuss here)
- have to get our hands dirty to say something useful ...

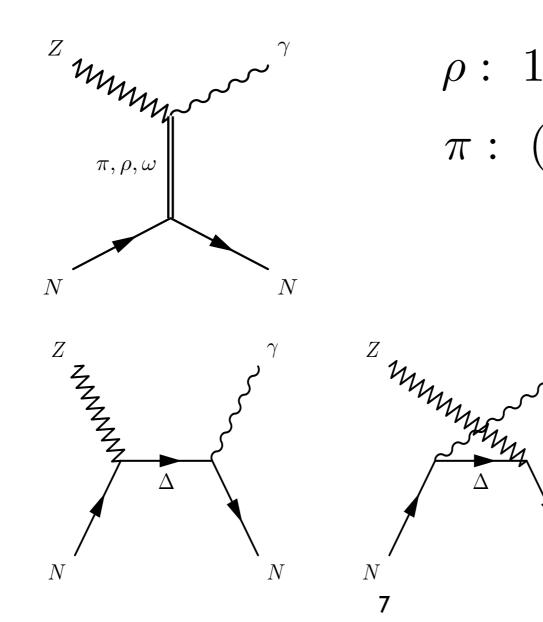
### Include ground state and leading resonances in each channel

Δ

N



at low E, match onto same (interesting) operator

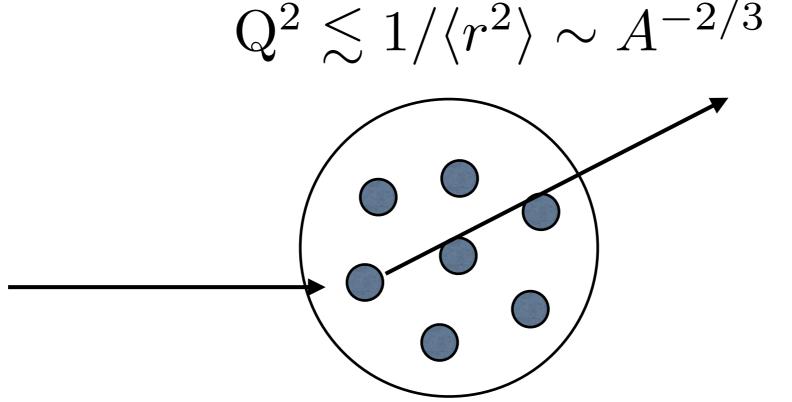


$$\rho: 1/N_c^2 \sim 1/9$$
  
 $\pi: (1 - 2s_W^2) - 2s_W^2 \ll 1$ 

Goldman and J. Jenkins 0906.0984

 $m_{\Delta} - m_N \sim 1/N_c$ 

Scattering on nucleus, can have both incoherent process (ejected nucleon) and coherent process (intact nucleus)



At small momentum transfer, amplitudes add,  $d\sigma \sim A^2$ 

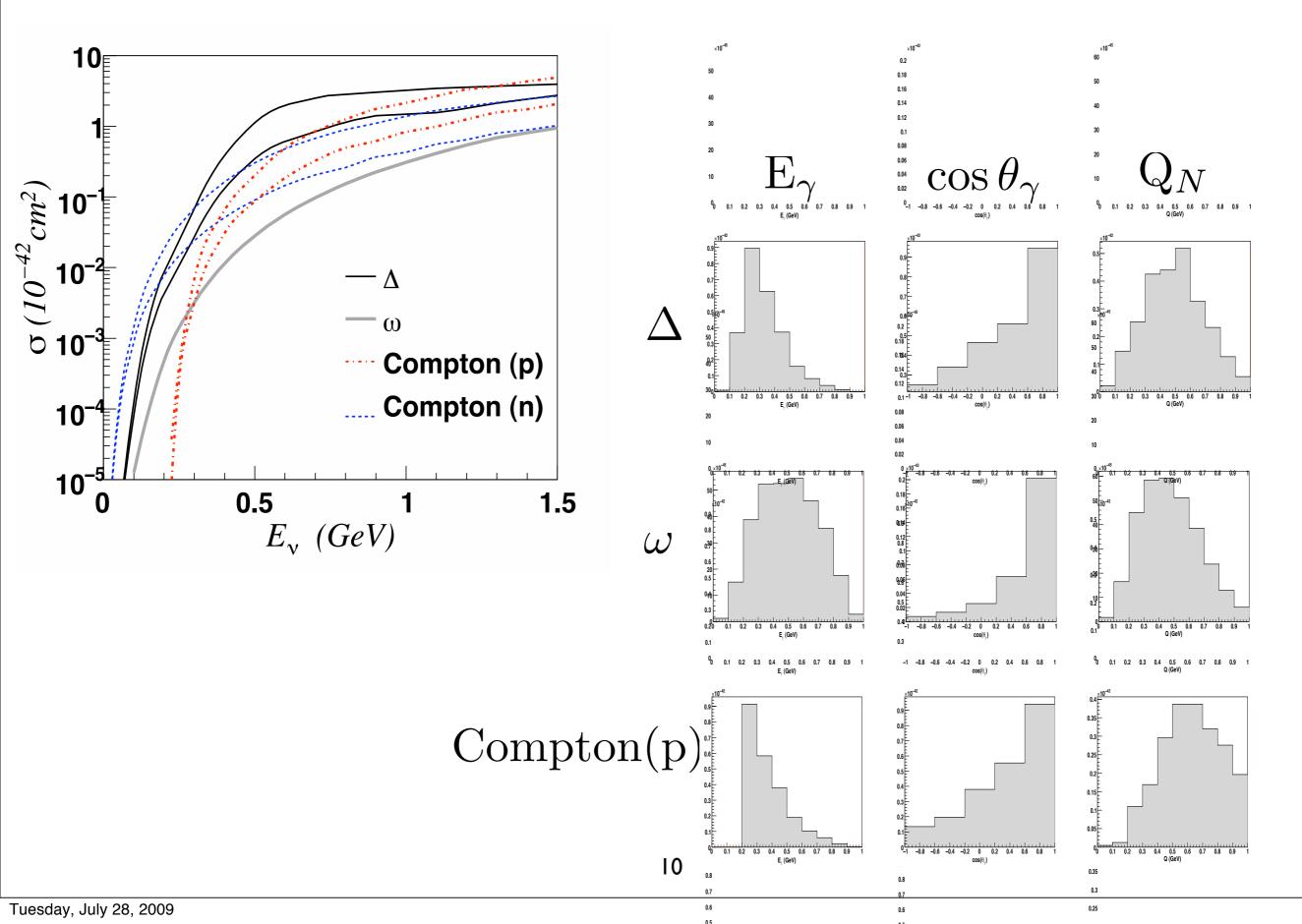
Nontrivial constraints on phase space, can analyze in limit of large nucleus:  $A^{\frac{1}{3}}E \rightarrow \infty$ 

$$\frac{\mathrm{d}\sigma_{\Delta}/\mathrm{d}\cos\theta}{\mathrm{d}\sigma_{\omega}/\mathrm{d}\cos\theta} \sim A^{4/3}(1-\cos^2\theta)$$
$$\frac{\mathrm{d}\sigma_{\omega}/\mathrm{d}\cos\theta}{\mathrm{d}\sigma_{\mathrm{Compton}}/\mathrm{d}\cos\theta} \sim A^{2/3}E^2\delta(\cos\theta)$$

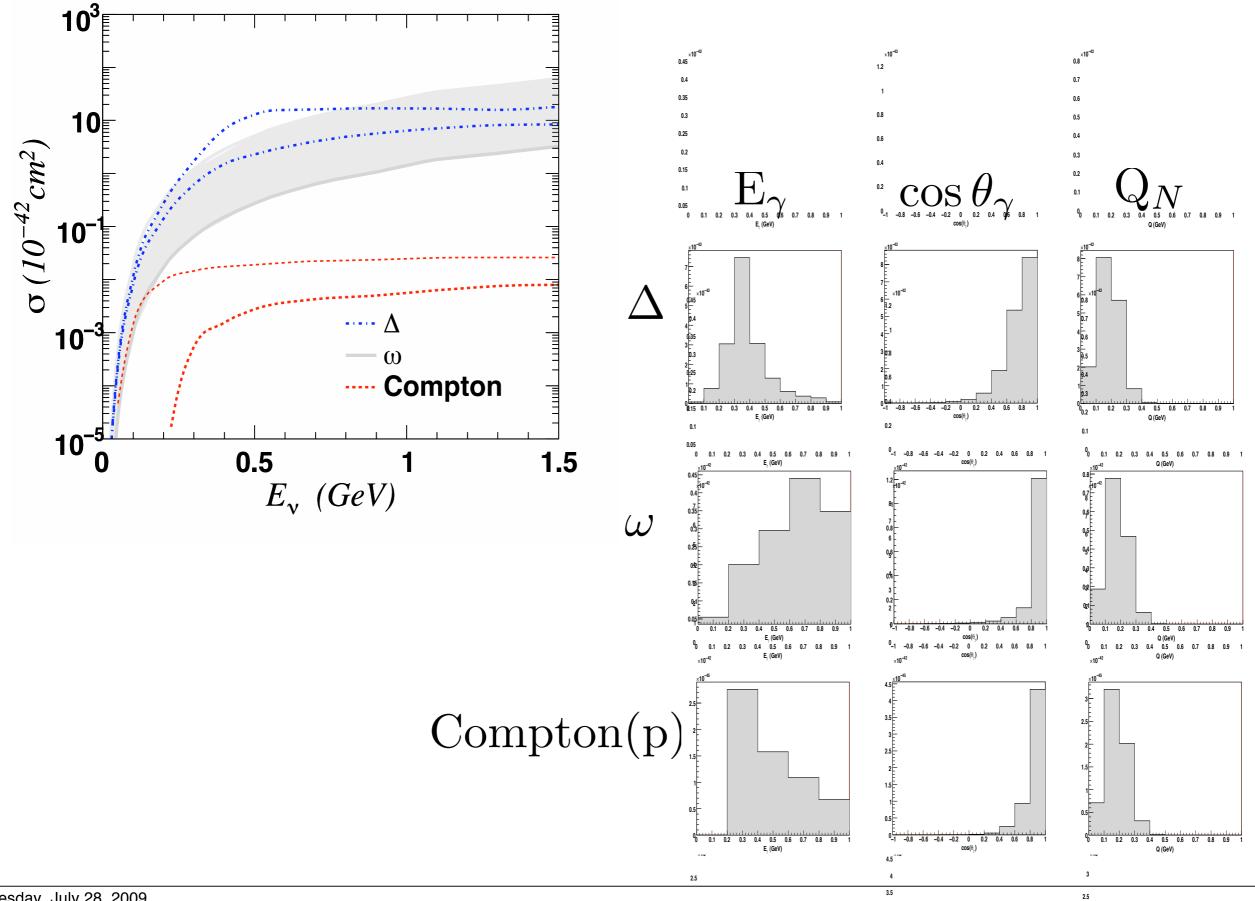
### Why is it so #!! hard to calculate?

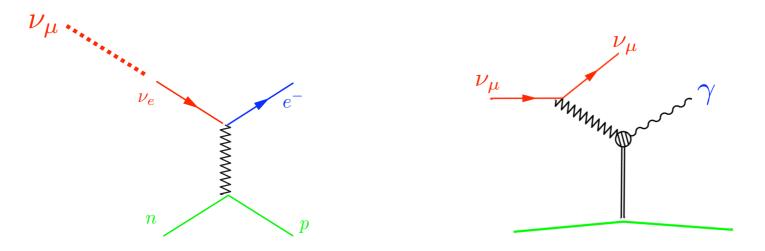
- what are the errors ?  $\approx$  what is the expansion ?
- need to get creative: I/N<sub>c</sub>, z(dispersive), I/A(nucleus), ...
- without support from data, errors to tree-level meson exchange are "I/N<sub>c</sub>" ~ 30% if all relevant states are considered (large energy ⇒ need more states)
- model independent approach: decompose into helicity amplitudes. but 12 of them, depending on multiple kinematic invariants - need dynamical model/small parameter expansion

### Single nucleon (incoherent) cross sections



### Coherent cross sections (<sup>12</sup>C)

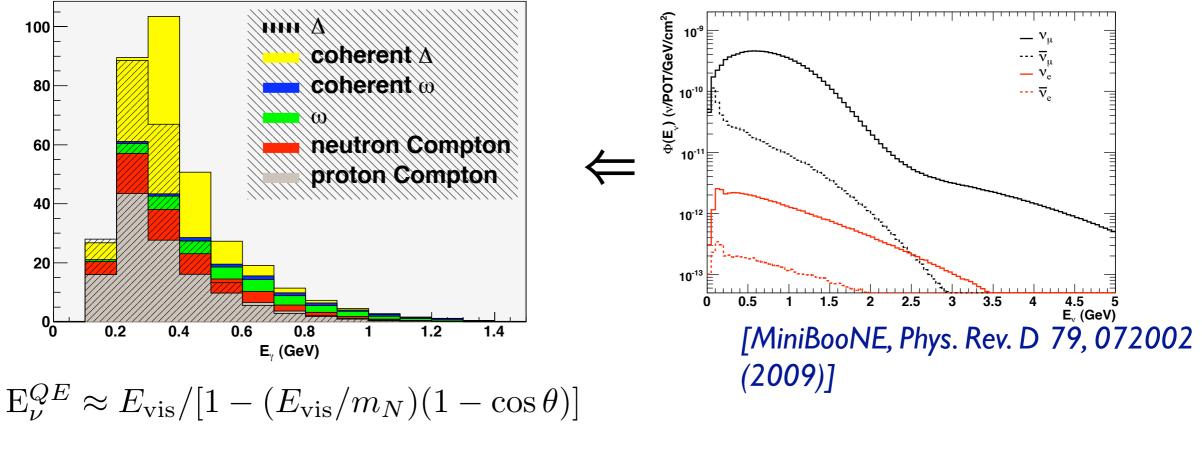




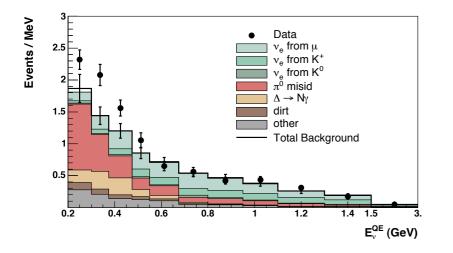
- MiniBooNE has presented a careful analysis of electron-like events in beams of primarily muon neutrinos / antineutrinos
- Interesting to compare this data, and apparent excess at low energy, with predictions of new single-photon events
- Potential background at other V<sub>e</sub> appearance experiments (higher energy not a focus of this talk - requires further theory)
- Preliminary results:

### flux averaging (MiniBooNE v mode)

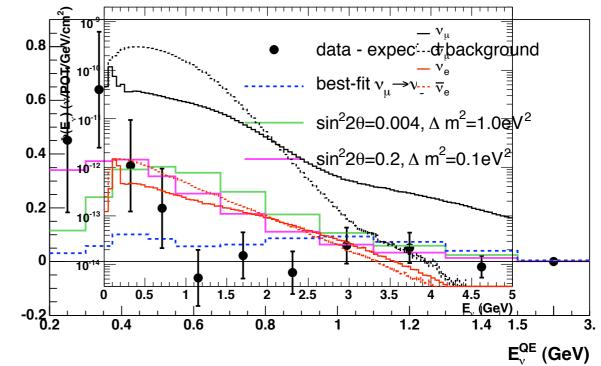
6.46e20 POT



Excess Events / MeV

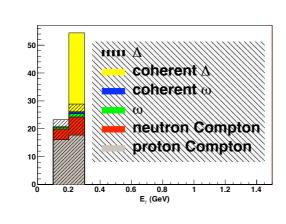


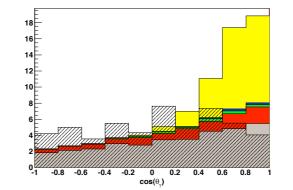
[MiniBooNE, PRL 102, 211801 (2009)]

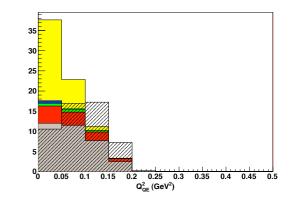


### flux averaged distributions (MiniBooNE v mode)

200-300 MeV

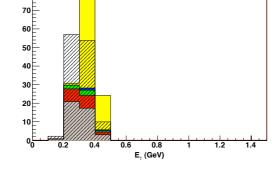








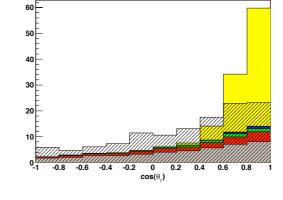


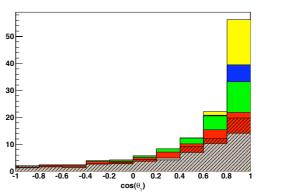


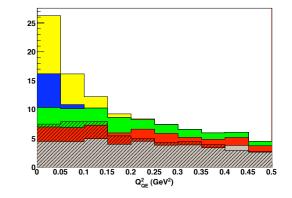
0.6

0.8

E, (GeV)



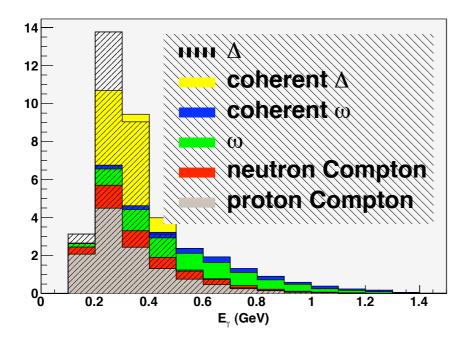


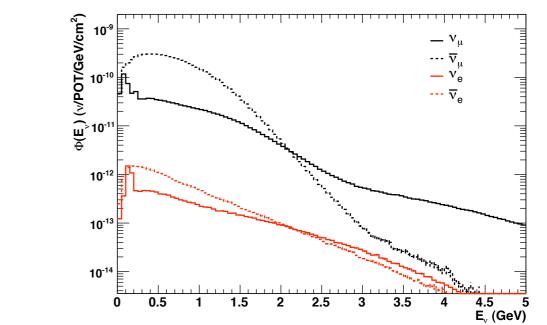


 $(\mathbf{Q}^2)^{QE} \approx 2E_{\nu}^{QE}E_{\mathrm{vis}}(1-\cos\theta)$ 

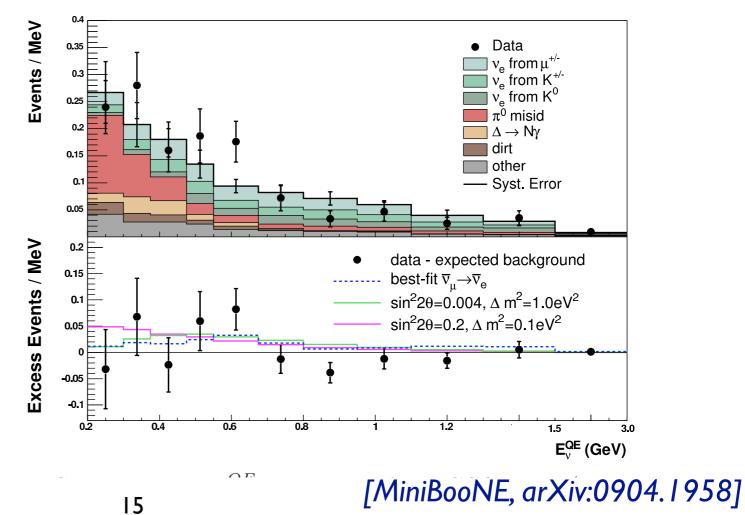
### flux averaging (MiniBooNE v-bar mode)

3.39e20 POT



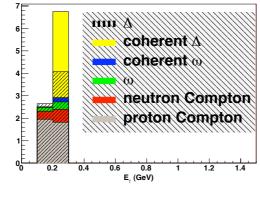


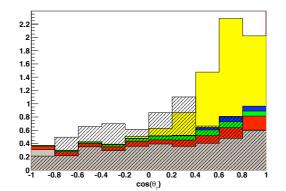
[MiniBooNE, Phys. Rev. D 79, 072002 (2009)]

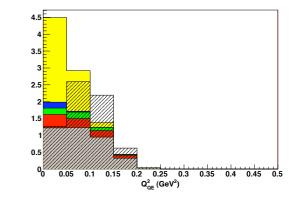


### flux averaged distributions (MiniBooNE v-bar mode)

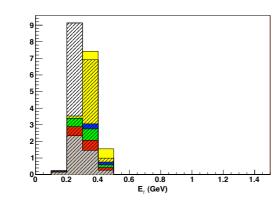
200-300 MeV

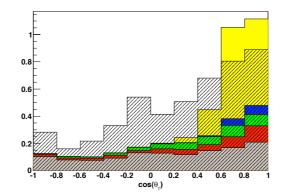


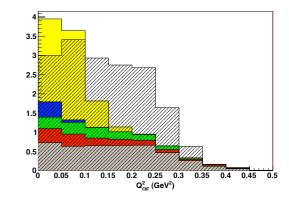




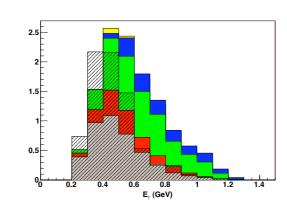
300-475

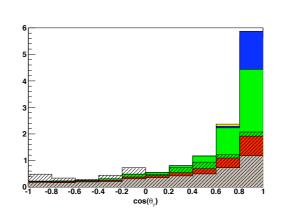


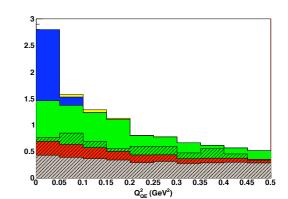




475-1250







	new events	excess	$\Delta$ -direct	Δ (MB)
200-300 MeV	75 [23]	45(26)	52 [16]	20
300-475	I 39 [42]	84(25)	I23 [37]	48
475-1250	119 [36]	22(36)	61 [18]	19

 $\chi^2 = 1.5/2$  d.o.f. (scale = 0.51),

 $\chi^2 = 3.8/3$  d.o.f. (scale = 0.3),

<u>v-bar</u>	200-300 MeV	9.3 [2.8]	05(117)	6.7 [2.0]	Ι.7
	300-475	<b>I3</b> [3.8]	-0.5(11.7)	17.3 [5.2]	4.9
	475-1250	12 [3.6]	3.2(10.0)	7.7 [2.3]	2.0

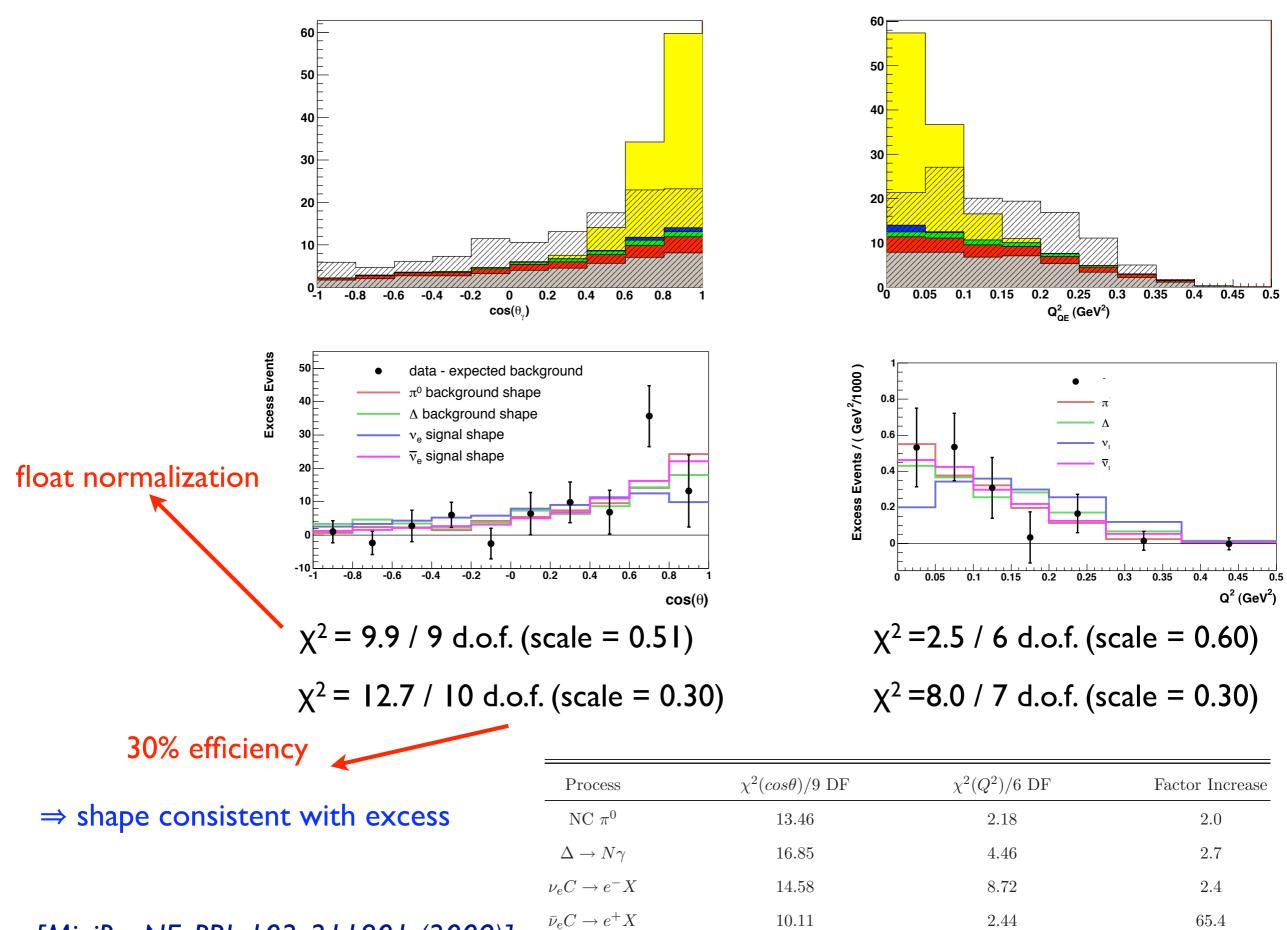
$$\chi^2 = 0.3/2$$
 d.o.f. (scale = 0.3),

 $\Rightarrow$  size consistent with data

- should do more complete efficiency analysis

17

<u>V</u>



18

[MiniBooNE, PRL 102, 211801 (2009)]

## Summary

- single photon events in neutrino-baryon scattering probe interesting (standard model) physics
- several mechanisms present, relevance varies with energy
- adding these events gives plausible explanation of MiniBooNE low-E excess
- potentially relevant background at T2K: mimics signal, varies with target material