Computation of the string tension in Yang-Mills theory using large N reduction

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### Results on

3D string tension (T=0) (arXiv:0807.1315)
4D string tension (T=0) (Real soon now)
3D string tension (T>T<sub>c</sub>) (arXiv:0906.3015)

# Outline

Quick results
Introduction
Details
Conclusion

#### Quick results: T=0 3 Dimensions 4 Dimensions $\odot$ 6<sup>4</sup> lattice $\odot$ 5<sup>3</sup> lattice ⊘ N = 47 ⊘ N = 47 and 59 ${ m O} b = rac{1}{q^2 N}$ = 0.6 to 0.8 0.3500 Wilson loops 1x1 to Wilson loops 1x1 to 7x7 9x9 $\sqrt{\sigma}b = 0.1964 \pm 0.0009$ $\sigma a^2 = 0.099 \pm 0.016$ (continuum extrapolation) at b=0.3480 N=47

# Quick result: T>Tc

 L<sup>3</sup> lattice L=4, 5, 6, 7
 ⊘ N = 59  $\bullet b = \frac{1}{a^2 N}$  = 0.9 to 1.75 Wilson loops 1x1 to 10x10  $\sigma a^2 \approx rac{1}{4 h L}$  (High T dimensional reduction)  $\bullet \sigma b^2 \approx \frac{T}{4b}$ 

## Introduction

Large N
Large N reduction
Phase structure
Project description

# Large N

Second Expansion parameters  $\alpha(Q^2)$  or 1/N  $\odot$   $N \to \infty$  simplifications Planar graphs Factorization Non-interacting mesons Ø OZI rule  $1/3 \approx 1/\infty$ 

# Large N reduction

Reduction to a one point 1<sup>d</sup> lattice (Eguchi-Kawai)

 $Z_N^d$  center symmetry

But broken at weak coupling

### Work-arounds

Quenched E-K But Bringoltz and Sharpe Twisted E-K But Teper and Vairinhos Continuum or partial reduction i.e. reduction to finite physical size  $l > 1/T_{c}$ 

# Center symmetry breaking at physical scale $Z_N^d \to Z_N^{d-1} \to Z_N^{d-2} \to \dots$



Kiskis, Narayanan, and Neuberger

## Details

- Wilson gauge field action with bare coupling g
- Smearing
- $b = \frac{1}{g^2 N}$  Tadpole improved to  $b_I = e(b)b$ with e(b) the average plaquette
- 3d: Space-like and time-like separations K, T in lattice units.  $k = K/b_I$   $t = T/b_I$



 $U' = P_{SU(N)}[(1 - f)U + \frac{f}{4}S_{12}^{+} + \frac{f}{4}S_{12}^{-} + \frac{f}{4}S_{13}^{+} + \frac{f}{4}S_{13}^{-}]$ Iterate n times  $\tau = fn$ 4d: f = 0.45 n = 5  $\tau = 2.25$   $\sqrt{\tau} = 1.5$ 3d (4->2): f = 0.1 n = 25  $\tau = 2.5$   $\sqrt{\tau} = 1.6$   $\odot$  5<sup>3</sup> lattice ⊘ N = 47 Smear space-like links with staples in the same time slice 

Fit to get quark-antiquark potential and string tension

 $\odot$  3-dimensions, T>T<sub>c</sub>  $\odot$  L<sup>3</sup> lattice L=4, 5, 6, 7 Ø N = 59 Smear space-like links with staples in the same time slice Wilson loops 1x1 to 10x10 (folded loops)

Fit to get quark-antiquark potential and string tension

4-dimensions, T=0  $\bullet$   $6^4$  lattice N = 37, 47, 59 b = 0.3450, 0.3480, 0.3500 Smear space-like links with staples in the same time slice Wilson loops 1x1 to 9x9 (folded loops) Fit to get quark-antiquark potential and string tension

#### Compute all Wilson loops 1x1 to 7x7 Fit to $W(k,t) = e^{-a-m(k)t}$



# Fit m(k) to $m(k) = \sigma b_I^2 k + c_0 b_I + \frac{c_1}{k}$



#### **Extrapolate:** $b_I \rightarrow \infty \quad \sqrt{\sigma} b_I \rightarrow 0.1964 \pm 0.0009$



3d

#### Are N and L large enough?





#### Are the results sensitive to smearing?



### Compute all Wilson loops 1x1 to 9x9 Fit 2x2 through 9x9 to $W(k,t) = e^{-a-m(k)t}$



4d

# Fit m(k) to $m(k) = \sigma a^2 k + c_0 + \frac{c_1}{k}$

b=0.348, L=6, N=47





4d

#### Scaling



**4**d



3d T



3d T



3d T

### Comparison with large L results <u>3-dimensions, T=0</u>

Karabali, Kim, and Nair  $\sqrt{\sigma b} = \frac{1}{\sqrt{8\pi}} \approx 0.1995$ Bringoltz and Teper  $\sqrt{\sigma b} = 0.1975 \pm 0.0002 - 0.0005$ J.K. & R.N. arXiv:0807.1315  $\sqrt{\sigma b} = 0.1964 \pm 0.0009$ 

### Comparison with large L results <u>4-dimensions, T=0</u>

• Lucini, Teper, and Wenger  $\frac{T_c}{\sqrt{\sigma}} = 0.597 \pm 0.004$ • J.K. & R.N.  $\frac{T_c}{\sqrt{\sigma}} = 0.61 \pm 0.05$ • At the same b<sub>I</sub>=0.182 Lucini, Teper, and Wenger N=8  $\sigma a^2 = 0.116 \pm 0.001$ J.K. & R.N.  $\sigma a^2 = 0.099 \pm 0.016$ 

# Conclusion

 In both 3 and 4 dimensions, continuum reduction gives good results for quantities based on the space-time dependence of large Wilson loops, e.g. the heavy quark potential and the string tension.