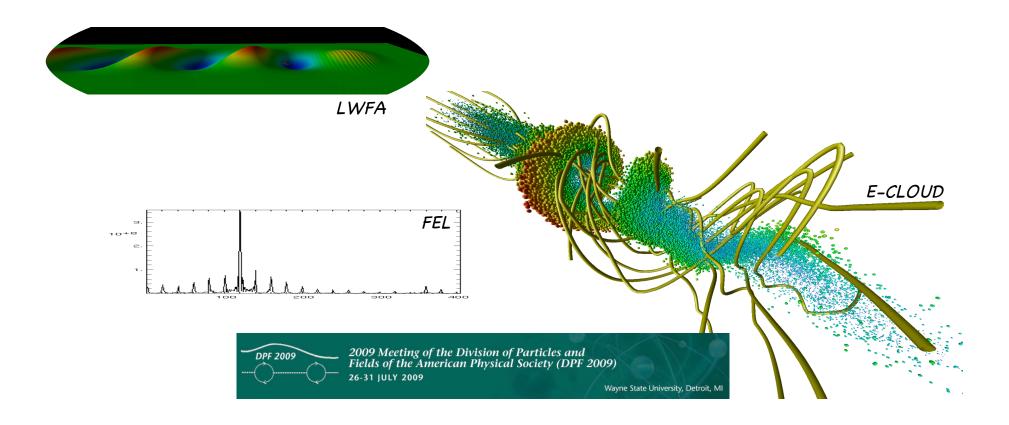
Speeding up simulations of relativistic systems using an optimal boosted frame

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- Concept
- Difficulties
- Examples of application
  - laser wakefield acceleration
  - electron cloud effects
  - free electron laser
- Conclusion



## Special relativity

### Lorentz transformation (LT) for v along x

$$t' = \gamma (t - vx/c^2) \qquad \gamma = (1 - v^2/c^2)^{-1/2}$$
  

$$x' = \gamma (x - vt)$$
  

$$y' = y$$
  

$$z' = z$$

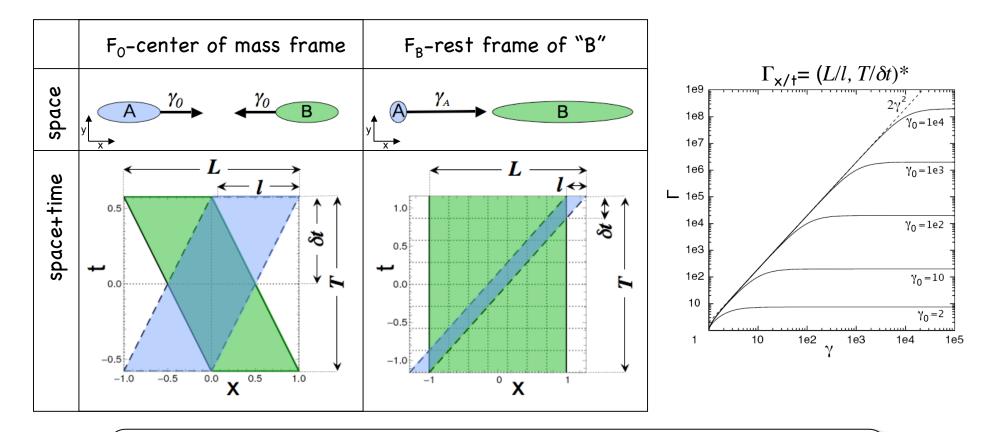
### Time dilation/space contraction

at rest: 
$$\Delta t$$
,  $\Delta x=0 \rightarrow$  in motion:  $\Delta t'=\gamma \Delta t$   
 $\Delta x$ ,  $\Delta t=0$   $\Delta x'=\Delta x/\gamma$ 

Lorentz invariant (invariant to change of reference frame)  $\Delta s^{2} = \Delta x^{2} + \Delta y^{2} + \Delta z^{2} - c^{2} \Delta t^{2} = \Delta x'^{2} + \Delta y'^{2} + \Delta z'^{2} - c^{2} \Delta t'^{2}$ 



### Range of space and time scales spawned by two identical beams crossing each other



- $\Gamma$  is not invariant under the Lorentz transformation:  $\Gamma_{\text{x/t}} \propto \gamma^2.$
- There exists an "optimum" frame which minimizes it.
- Result is general and applies to light beams too.

\*J.-L. Vay, Phys. Rev. Lett. 98, 130405 (2007)

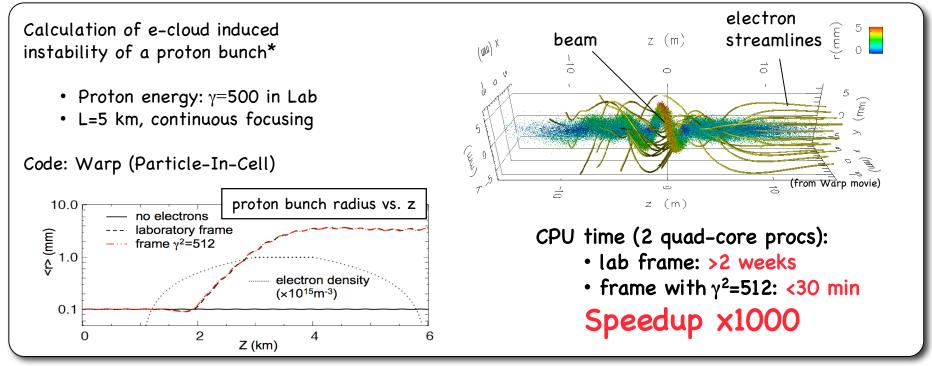


## Consequence for computer simulations

# of computational steps grows with the full range of space and time scales involved



Choosing optimum frame of reference to minimize range can lead to dramatic speed-up for relativistic matter-matter or light-matter interactions.



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# Seems simple but <u>1</u>. Algorithms which work in one frame may break in another. Example: the Boris particle pusher.

- Boris pusher ubiquitous
  - In first attempt of e-cloud calculation using the Boris pusher, the beam was lost in a few betatron periods!
  - Position push:  $X^{n+1/2} = X^{n-1/2} + V^n \Delta t$  -- no issue
  - Velocity push:  $\gamma^{n+1}\mathbf{V}^{n+1} = \gamma^{n}\mathbf{V}^{n} + \frac{q\Delta t}{m} (\mathbf{E}^{n+1/2} + \frac{\gamma^{n+1}\mathbf{V}^{n+1} + \gamma^{n}\mathbf{V}^{n}}{2\gamma^{n+1/2}} \times \mathbf{B}^{n+1/2})$

issue:  $E+v \times B=0$  implies  $E=B=0 \Rightarrow large errors$  when  $E+v \times B\approx 0$  (e.g. relativistic beams).

• Solution

- Velocity push: 
$$\gamma^{n+1}\mathbf{V}^{n+1} = \gamma^{n}\mathbf{V}^{n} + \frac{q\Delta^{\dagger}}{m} (\mathbf{E}^{n+1/2} + \frac{\mathbf{V}^{n+1} + \mathbf{V}^{n}}{2} \times \mathbf{B}^{n+1/2})$$

• Not used before because of implicitness. We solved it analytically\*

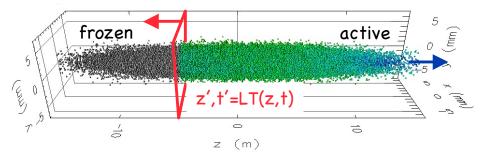
$$\begin{cases} \gamma^{i+1} = \sqrt{\frac{\sigma + \sqrt{\sigma^2 + 4(\tau^2 + u^{*2})}}{2}} & \text{(with } \mathbf{u} = \gamma \mathbf{v}, \quad \mathbf{u}' = \mathbf{u}^{\mathbf{i}} + \frac{q\Delta t}{m} \left( \mathbf{E}^{i+1/2} + \frac{\mathbf{v}^i}{2} \times \mathbf{B}^{i+1/2} \right), \quad \tau = (q\Delta t/2m) \mathbf{B}^{i+1/2}, \\ \mathbf{u}^{i+1} = [\mathbf{u}' + (\mathbf{u}' \cdot \mathbf{t})\mathbf{t} + \mathbf{u}' \times \mathbf{t}]/(1+t^2) & u^* = \mathbf{u}' \cdot \tau/c, \quad \sigma = \gamma'^2 - \tau^2, \quad \gamma' = \sqrt{1 + u'^2/c^2}, \quad \mathbf{t} = \tau/\gamma^{i+1}). \end{cases}$$

\*J.-L. Vay, Phys. Plasmas 15, 056701 (2008)

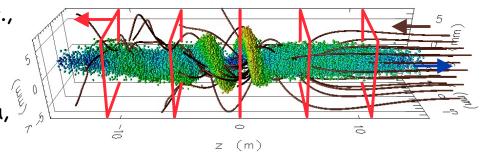


### Other possible complication: inputs/outputs

- Often, initial conditions known and output desired in laboratory frame
  - relativity of simultaneity => inject/collect at plane(s)  $\perp$  to direction of boost.
- Injection through a moving plane in boosted frame (fix in lab frame)
  - fields include frozen particles,
  - same for laser in EM calculations.



- Diagnostics: collect data at a collection of planes
  - fixed in lab fr., moving in boosted fr.,
  - interpolation in space and/or time,
  - already done routinely with Warp for comparison with experimental data, often known at given stations in lab.





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# Several areas in which simulations in a boosted may be beneficial were identified

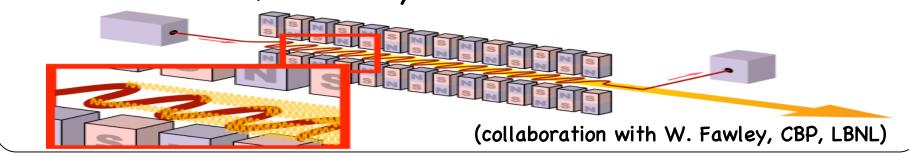
# Laser-plasma wakefield accelerators Plasma wake e- beam Laser pulse

(collaboration with LBNL's LOASIS group, lead by Wim Leemans)

#### Electron cloud driven beam instabilities



Free electron lasers/coherent synchrotron radiation

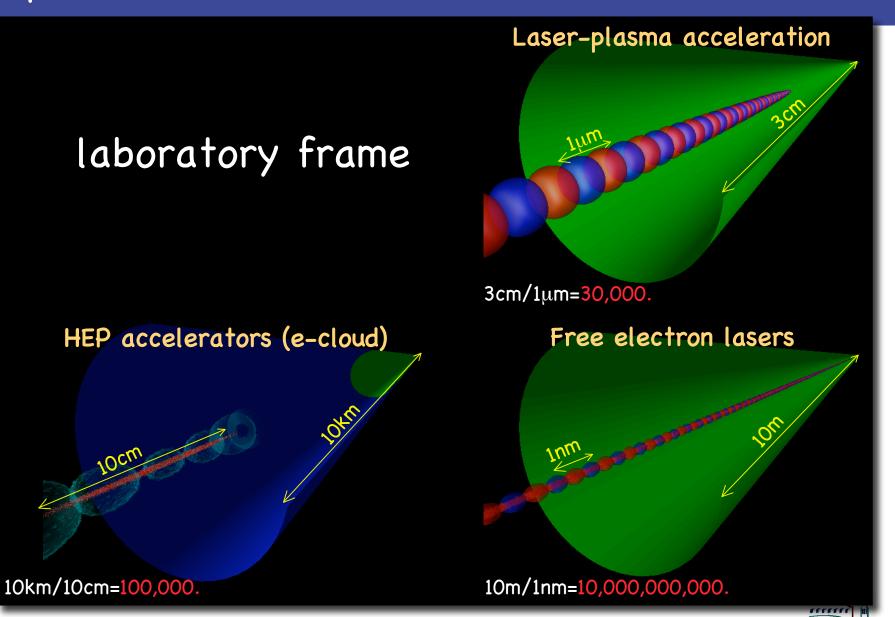




E<sub>//</sub> (GV/m)

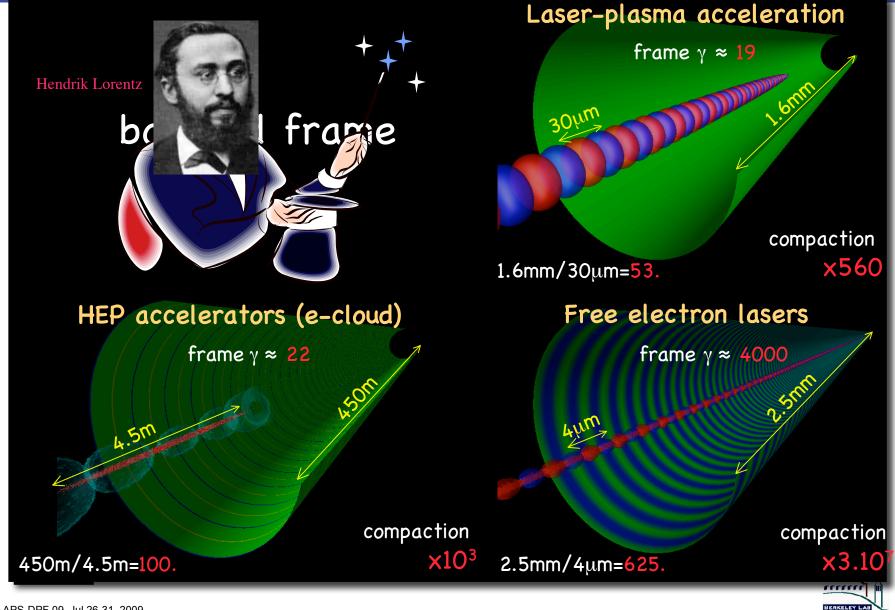
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# Large scale range renders simulation difficult, if not impractical, in lab frame



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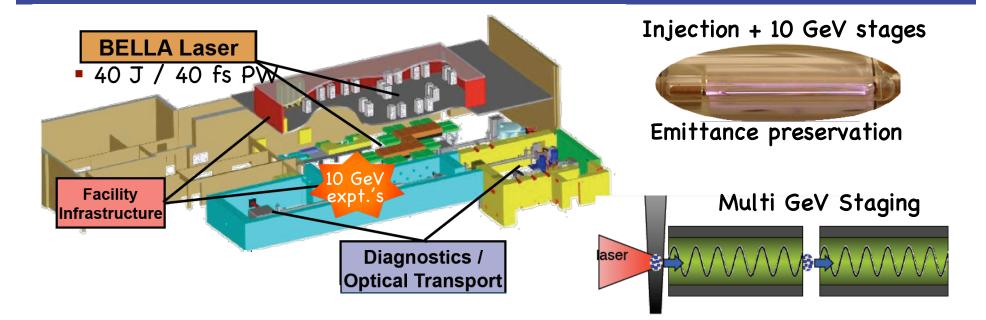
#### Lorentz transformation => large level of compaction of scales range



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#### BELLA 40 J PW Laser – Components for a Laser Plasma Collider



Simulating 10 GeV stages explicitly (PIC) in lab frame needs ~1G CPU•hours  $\Rightarrow$  impractical\*

Predictions have relied on theory, reduced models (fluid, envelope, quasistatic), scaling:

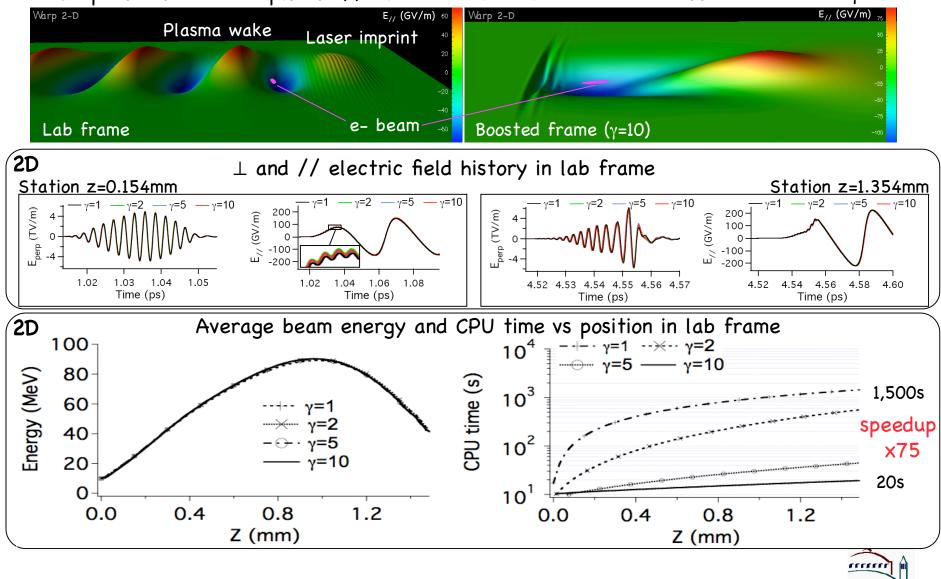
- Energy gain  $\propto n^{-1}$  : 10 GeV at 10<sup>17</sup>/cc  $\Rightarrow$  100 MeV at 10<sup>19</sup>/cc
- Length  $\propto n^{-3/2}$  : 1m at 10<sup>17</sup>/cc  $\Rightarrow$  1mm at 10<sup>19</sup>/cc
- Gradient  $\propto n^{1/2}$  : 10 GV/m at 10<sup>17</sup>/cc  $\Rightarrow$  100 GV/m at 10<sup>19</sup>/cc

Can simulations of full scale 10 GeV stages be practical using a Lorentz boosted ref. frame?

- difficulty: backward emitted radiation frequency upshifted in boosted frame.

# Scaled simulations of a 10 GeV LWFA stage ( $\lambda$ =0.8µm, $a_0$ =1, $k_p$ L=2, $L_p$ =1.5mm in lab)

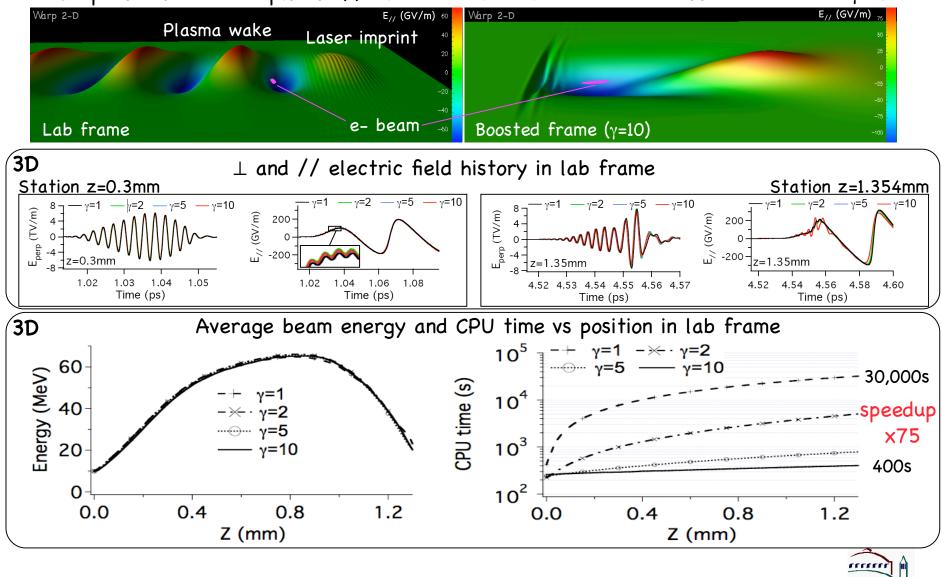
Snapshots of surface plot of // electric field in lab frame and boosted frame at  $\gamma$ =10



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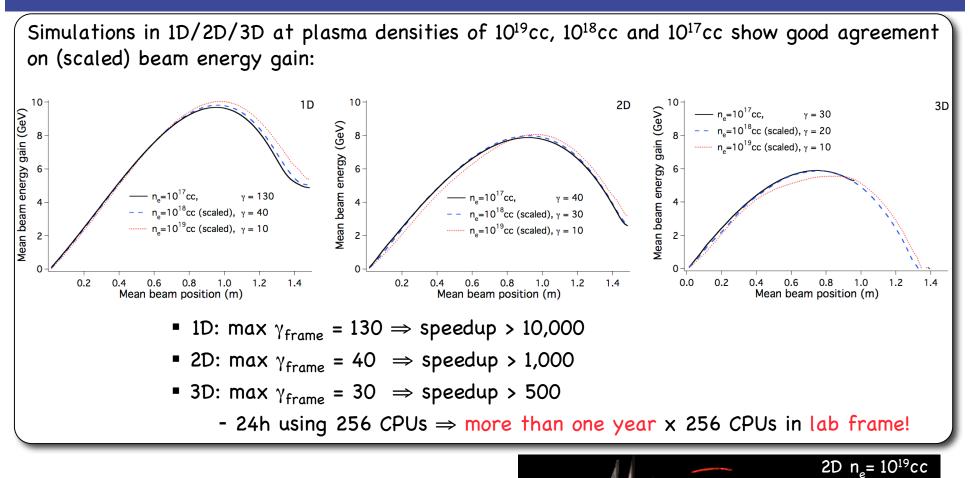
Snapshots of surface plot of // electric field in lab frame and boosted frame at  $\gamma$ =10



Vay, APS-DPF 09, Jul 26-31, 2009

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### Full scale simulations of a 10 GeV LWFA stage



Max  $\gamma_{\text{frame}}$  achieved in 2D and 3D limited by instability developing at front of plasma

origin and cures are being studied...

\*similar work by Bruhwiler et al (Tech X), Martins et al (UCLA/IST)

 $\gamma_{\text{frame}} = 13$ 

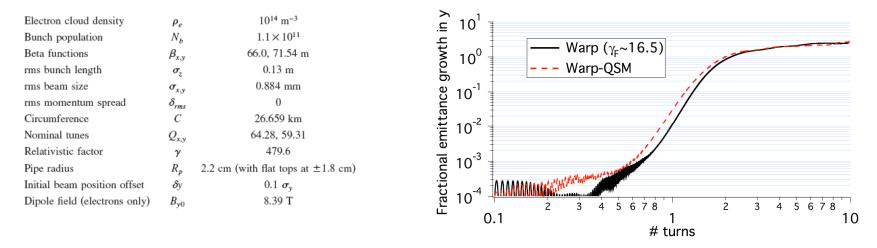
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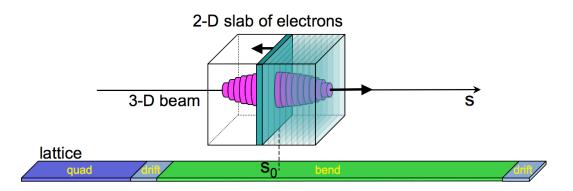


### E-cloud: benchmarking against quasistatic model for LHC scenario

Excellent agreement on emittance growth between boosted frame full PIC and "quasistatic" for e-cloud driven transverse instability in continuous focusing model of LHC



The "quasistatic" approximation uses the separation of time scales for pushing beam and ecloud macro-particles with different "time steps"





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### FEL in Boosted-Frame E&M Code

Physics ignored by Eikonal codes but accessible to boosted frame approach: **Backward wave emission** Wide-angle emission (generally highly red-shifted) CSE for all undulator, e-beam configurations Emission from very short beams Emission from beams with rapidly-varying envelope properties Emission from beams bunched with "multiple colors" Properties of \*very\* high gain systems ( $L_G/\lambda_u < 5$ ) FEL emission from beams in multiple harmonic undulators Biharmonic (or triharmonic undulators) Effects of adiabatic match sections FEL emission in waveguides where  $v_{group}$  strongly varying with  $\omega$  (normally relevant to microwave FEL's operating near cutoff) Overall computational speed impressive compared to full E&M but much slower than standard eikonal method: Not likely to become dominant paradigm for short wavelength FEL's but *might* be useful for very high gain microwave/far-IR devices or situations with wideband spectral output

### Conclusion and outlook

• The range of scales of a system is not a Lorentz invariant ( $\propto \gamma^2$ ), and there exists an optimum frame minimizing it => orders of magnitude speedup predicted for some simulations.

- Calculating in a boosted frame more demanding, eventually:
  - developed new particle pusher for e-cloud problems,
  - added capabilities for injection/diagnostics in boosted frame.
- Orders of magnitude speedup demonstrated for a class of firstprinciple simulations of multiscale problems: laser-plasma acceleration, e-cloud in HEP accelerators, free electron lasers.
- Explore other applications: CSR, astrophysics,...
- Can we develop methods which costs do not depend on frame?

