

Discovery Potential for MSSM Higgs Bosons with the ATLAS Experiment at the LHC

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Introduction

Higgs Sector in the Minimal Supersymmetric Standard Model (MSSM)

- Two Higgs Doublets → Five physical bosons: A (CP-odd), h, H (CP-even) and H^\pm
- At tree-level, the masses of the 5 Higgs bosons are related:

$$m_{H,h}^2 = \frac{1}{2} [m_A^2 + m_Z^2 \pm \sqrt{(m_A^2 + m_Z^2)^2 - 4m_A^2 m_Z^2 \cos^2 2\beta}],$$

$$m_{H^\pm}^2 = m_W^2 + m_A^2,$$

- Branching ratios to down-type quarks and charged leptons are enhanced:

Φ	$g_{\Phi\bar{u}u}$		$g_{\Phi\bar{d}d}$		$g_{\Phi VV}$
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$	$\sin(\beta - \alpha)$
H	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$	$\cos(\beta - \alpha)$
A	$\cot \beta$	$\cot \beta$	$\cot \beta$	$\tan \beta$	0

A. Djouadi, hep-ph/0503173v2

Table 1.4: The neutral Higgs couplings to fermions and gauge bosons in 2HDMs of Type I and II compared to the SM Higgs couplings. The H^\pm couplings to fermions follow that of A.

$$\alpha = \frac{1}{2} \arctan \left(\tan 2\beta \frac{M_A^2 + M_Z^2}{M_A^2 - M_Z^2} \right), \quad -\frac{\pi}{2} \leq \alpha \leq 0$$

$$\tan \beta = \frac{v_2}{v_1} = \frac{(v \sin \beta)}{(v \cos \beta)}$$

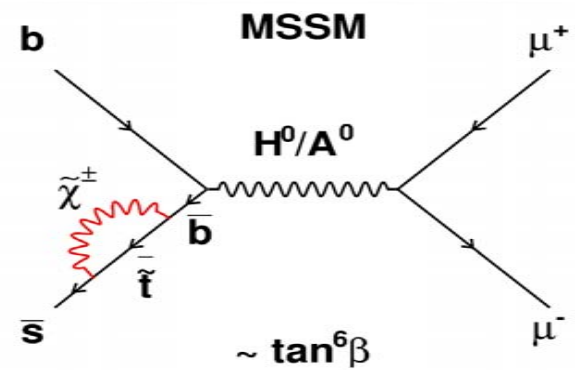
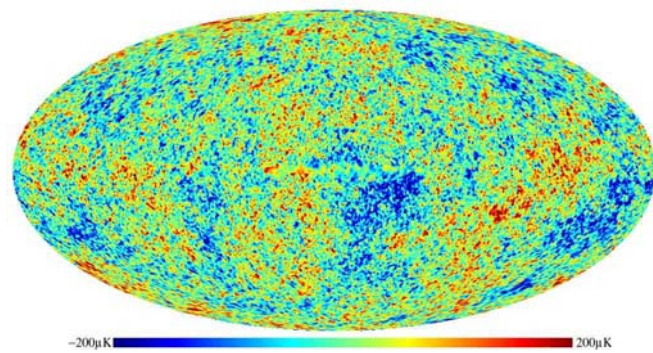
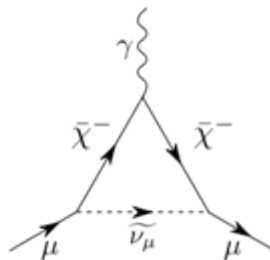
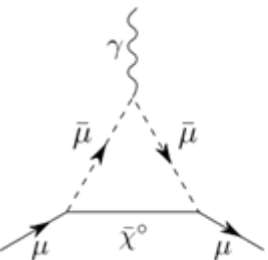
Large loop corrections to masses and couplings depend on SUSY parameters:

- Largely dependent on the top / stop sector “ m_h -max” scenario
- $X_t = 2$ TeV, $M_{\text{SUSY}} = 1$ TeV, $M_2 = 200$ GeV, $\mu = 200$ GeV and $M_{\text{gluino}} = 800$ GeV

Discovery Potential and Exclusion Bounds

- Scan the $m_A - \tan \beta$ plane [[CERN-OPEN-2008-020](#); [arXiv:0901.0512](#)]

Indirect Constraints from Experiment

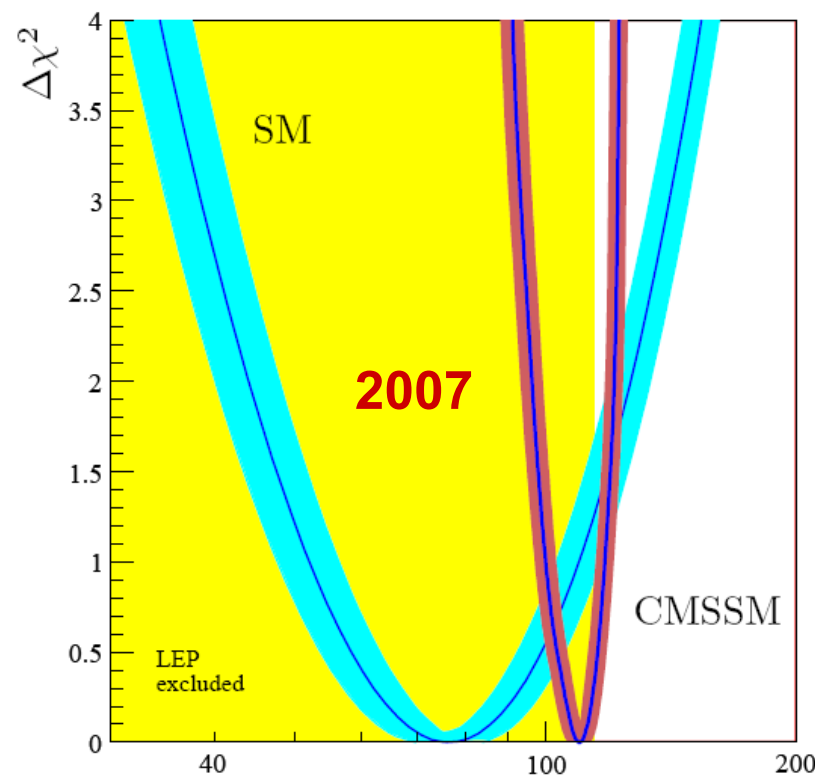


$(g - 2)_\mu$

Variable	Measurement	Fit	$ \sigma^{\text{meas}} - \sigma^{\text{fit}} / \sigma^{\text{meas}}$
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	0.02758 ± 0.00035	0.02774	0.1
m_Z [GeV]	91.1875 ± 0.0021	91.1873	0.02
Γ_Z [GeV]	2.4952 ± 0.0023	2.4952	0.0
σ_{had}^0 [nb]	41.540 ± 0.037	41.486	0.13
R_1	20.767 ± 0.025	20.744	0.11
$A_{\text{fb}}^{0,l}$	0.01714 ± 0.00095	0.01641	0.42
$A_1(P_\tau)$	0.1465 ± 0.0032	0.1479	0.09
R_b	0.21629 ± 0.00066	0.21613	0.07
R_c	0.1721 ± 0.0030	0.1722	0.00
$A_{\text{fb}}^{0,b}$	0.0992 ± 0.0016	0.1037	0.45
$A_{\text{fb}}^{0,c}$	0.0707 ± 0.0035	0.0741	0.48
A_b	0.923 ± 0.020	0.935	0.13
A_c	0.670 ± 0.027	0.668	0.03
$A_1(\text{SLD})$	0.1513 ± 0.0021	0.1479	0.22
$\sin^2\theta_{\text{eff}}^{\text{lept}}(Q_{\text{fb}})$	0.2324 ± 0.0012	0.2314	0.04
m_W [GeV]	80.398 ± 0.025	80.382	0.02
m_t [GeV]	170.9 ± 1.8	170.8	0.06
$R(b \rightarrow s\gamma)$	1.13 ± 0.12	1.12	0.09
$B_s \rightarrow \mu\mu$ [$\times 10^{-8}$]	< 8.00	0.33	N/A (upper limit)
Δa_μ [$\times 10^{-9}$]	2.95 ± 0.87	2.95	0.0
Ωh^2	0.113 ± 0.009	0.113	0.0

WMAP

BR($B_s \rightarrow \mu^+\mu^-$)



$m_h = 110 (+8)(-10) \pm 3$ (theo.) GeV

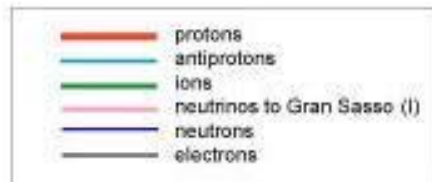
m_{Higgs} [GeV/ c^2] 3

O. Buchmueller et al., arXiv:0707.3447v2

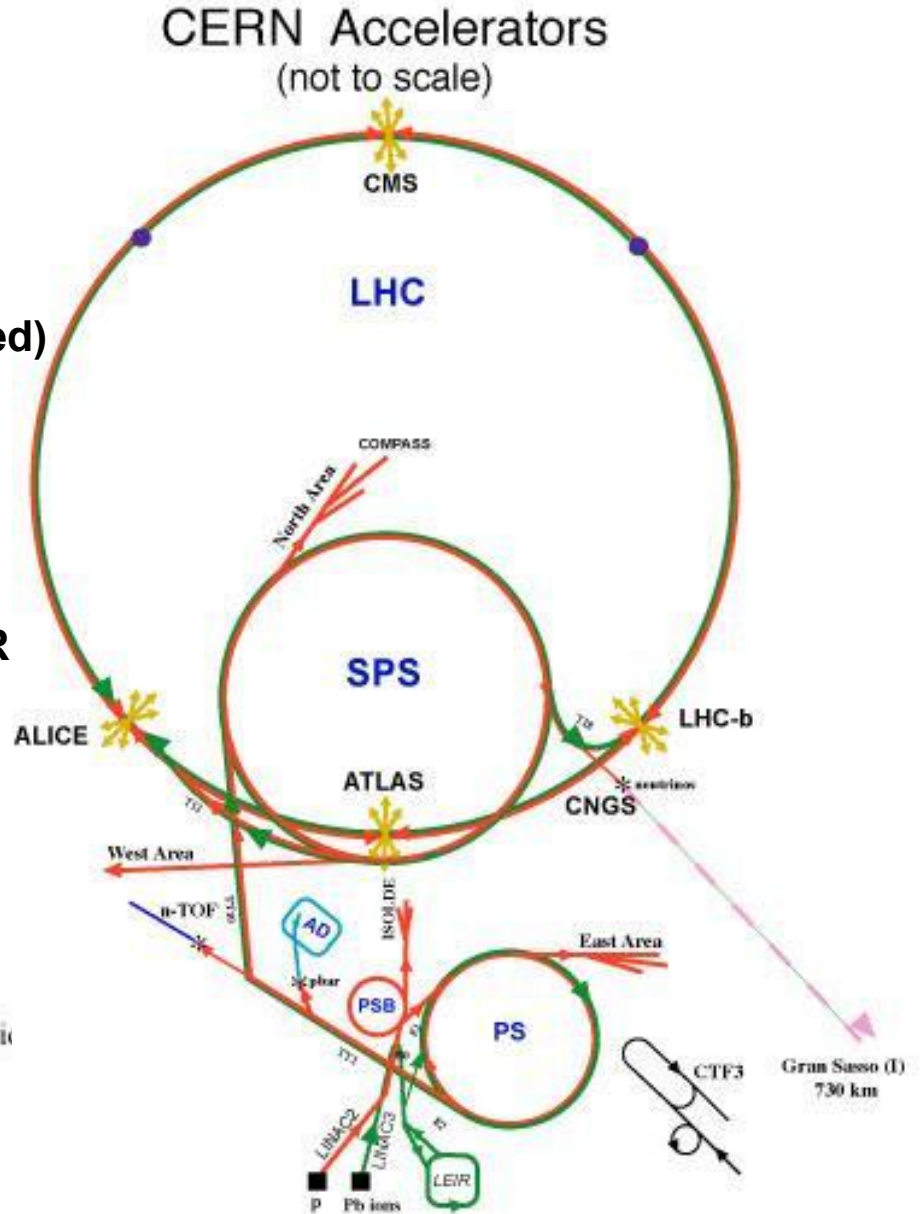
The Large Hadron Collider

Particle accelerator located at CERN (Geneva, Switzerland)

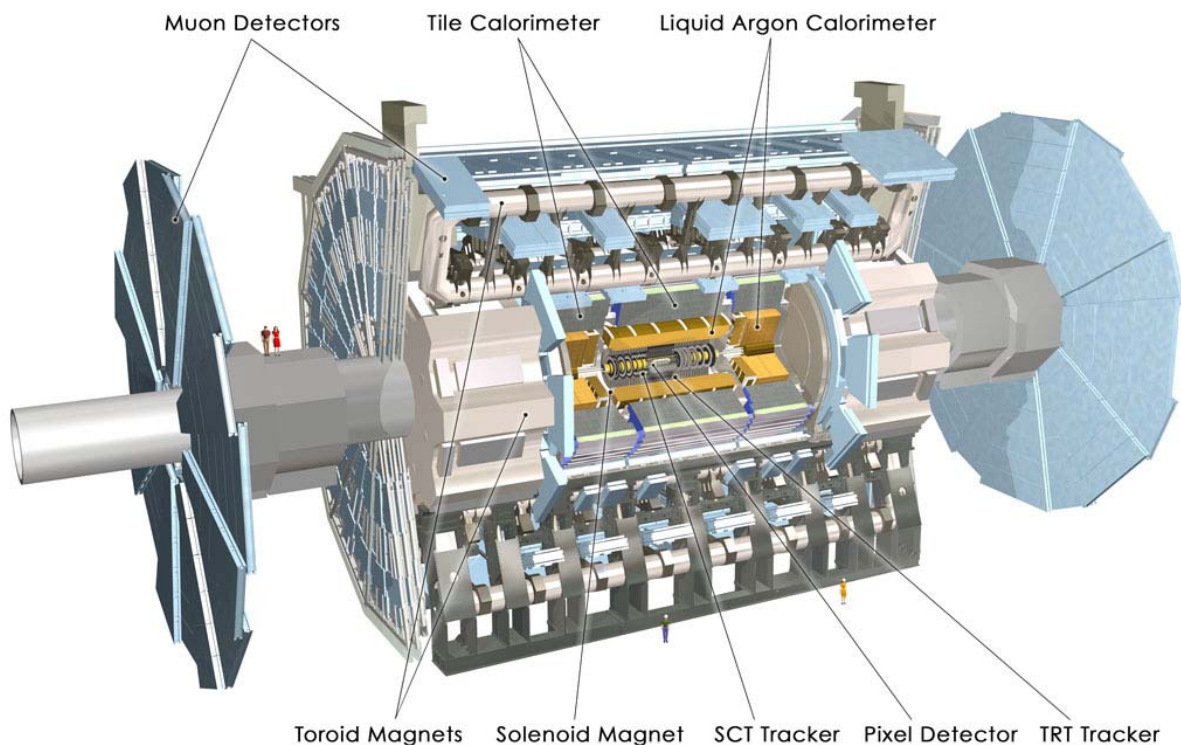
- Collide protons at a 14 TeV CME
- Housed in the former LEP tunnel
- Dipole field at 7 TeV is 8.33 T
- ~350 MJ per beam!
- Ultimately ~2800 bunches
- Vacuum 10^{-13} atm (~6500 m³ pumped)
- 1232 Dipoles (operate at 1.9 K)
- 858 Quadrupoles
- Typical store lasts ~10 hours
- Can be used for ion collisions (Pb)
- Final price tag estimated at 4G EUR



LHC: Large Hadron Collider
 SPS: Super Proton Synchrotron
 AD: Antiproton Decelerator
 ISOLDE: Isotope Separator OnLine DEvice
 PSB: Proton Synchrotron Booster
 PS: Proton Synchrotron
 LINAC: LINEar ACcelerator
 LEIR: Low Energy Ion Ring
 CNGS: Cern Neutrinos to Gran Sasso



The ATLAS Experiment



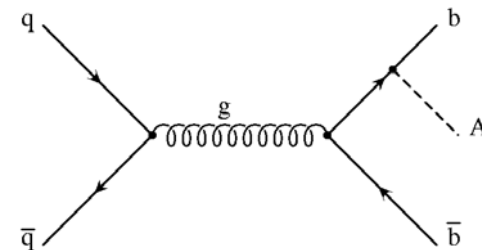
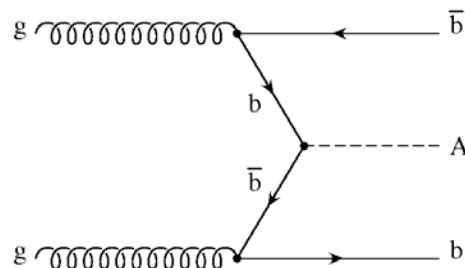
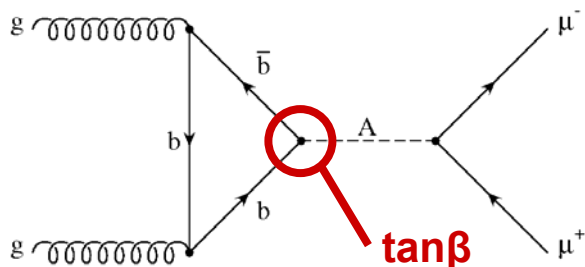
	ATLAS
Weight	7000 tons
Diameter	22m
Length	46m
Peak B Field	2T solenoid 3.9T (peak) BA toroid 4.1T (peak) EC toroids

PERFORMANCE		
Tracker	Si pixels, strips + TRT (pid)	$\sigma/p_T \approx 5 \times 10^{-4} p_T \oplus 0.01$
EM calorimeter	Pb + LAr	$\sigma/E \approx 10\%/\sqrt{E} \oplus 0.007$
Hadronic calorimeter	Fe+scintillator / Cu + LAr	$\sigma/E \approx 50\%/\sqrt{E} \oplus 0.03$
Combined Muons (ID+MS)	2%@50GeV to	10%@1TeV

Neutral MSSM Higgs Discovery Potential

Neutral MSSM Higgs

Direct and Associated Production of the h , H and A



Enhanced for large $\tan\beta$

Investigated the decay channels (14 TeV)

- $h/A/H \rightarrow \text{tau tau} \rightarrow 2l 4\nu$
- $h/A/H \rightarrow \text{mu mu}$
- Other final states (di-tau lepton-hadron and fully hadronic) are still under study
- Early running and low-luminosity scenarios for the above channels are also being considered (should have some preliminary results soon)

Neutral Higgs mass degeneracy

- For much of the parameter space the neutral Higgs masses are degenerate
- Cross-sections are summed

MSSM Higgs Di-Tau Analysis

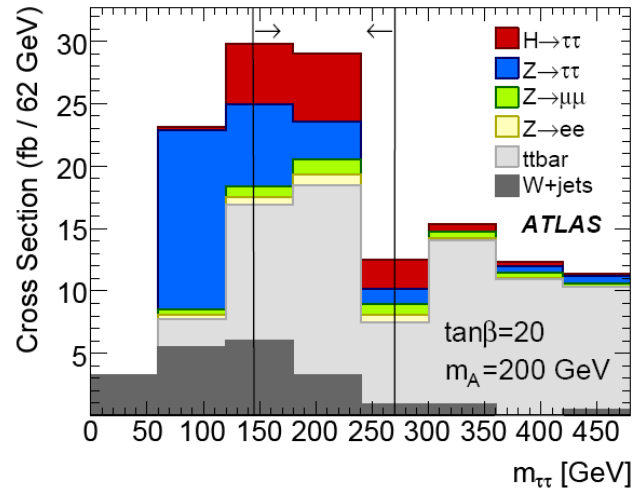
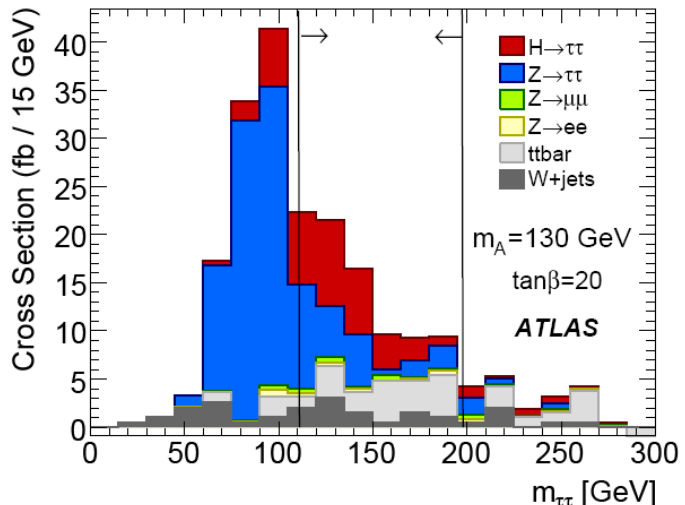
Branching-Ratio to taus is enhanced in the MSSM

- Investigated $h/A/H \rightarrow \tau\tau \rightarrow 2l\ 4\nu$ with associated b-jets
- High- p_T electron or muon triggers
- Imposed lepton kinematic requirements
- Required at least one b-jet to be present in the event
- Expect a large amount of missing transverse energy

Mass reconstruction is done via the collinear approximation

$$m_{\tau\tau} = \frac{m_{ll}}{\sqrt{x_1 x_2}}$$

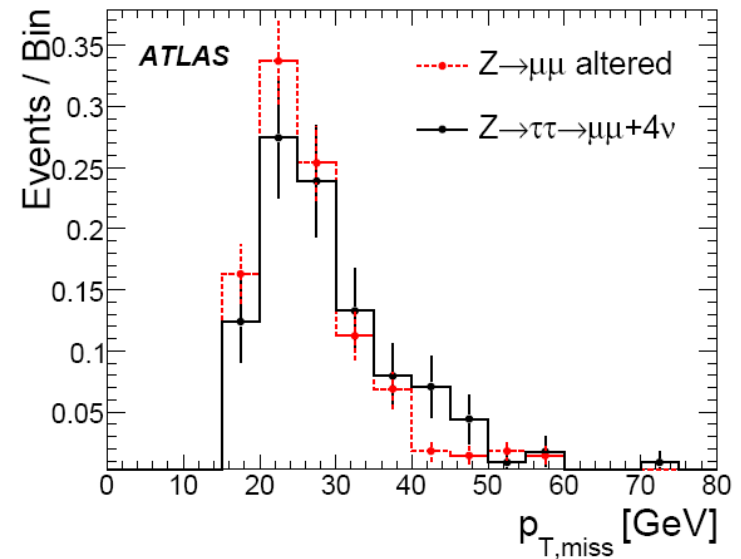
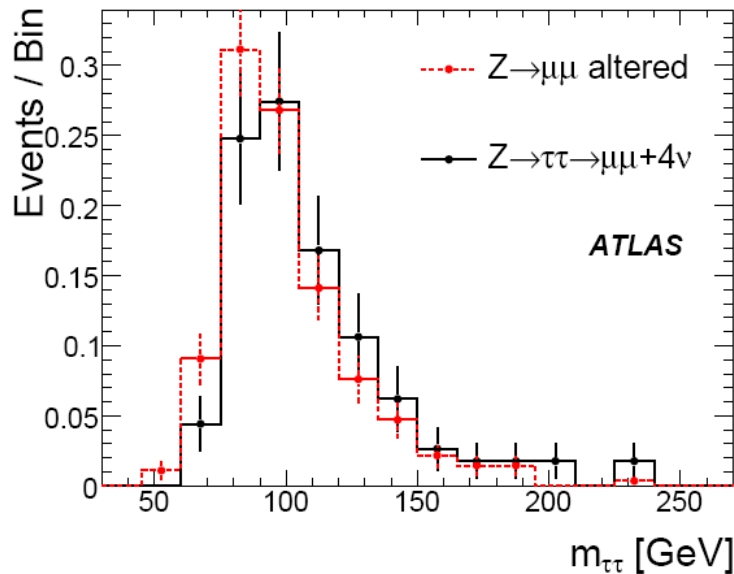
$$x_i = \frac{p_{T,li}}{p_{T,\tau i}}$$



Backgrounds to the Di-Tau Analysis

For $m_A < 200$ GeV, dominant background is $Z + \text{jets}$ with $Z \rightarrow \text{tau tau}$

- This is an irreducible background
- The shape and normalization can be taken from data-driven control samples
- Scale the energy of the $Z \rightarrow \mu\mu$ events collected in collision data to match that expected from $Z \rightarrow \text{tau tau}$



For $m_A \geq 200$ GeV, $t\bar{t}$ events become a significant background

- Can get a handle on this by cutting on the jet multiplicity ($N_{\text{jets}} \leq 2$)

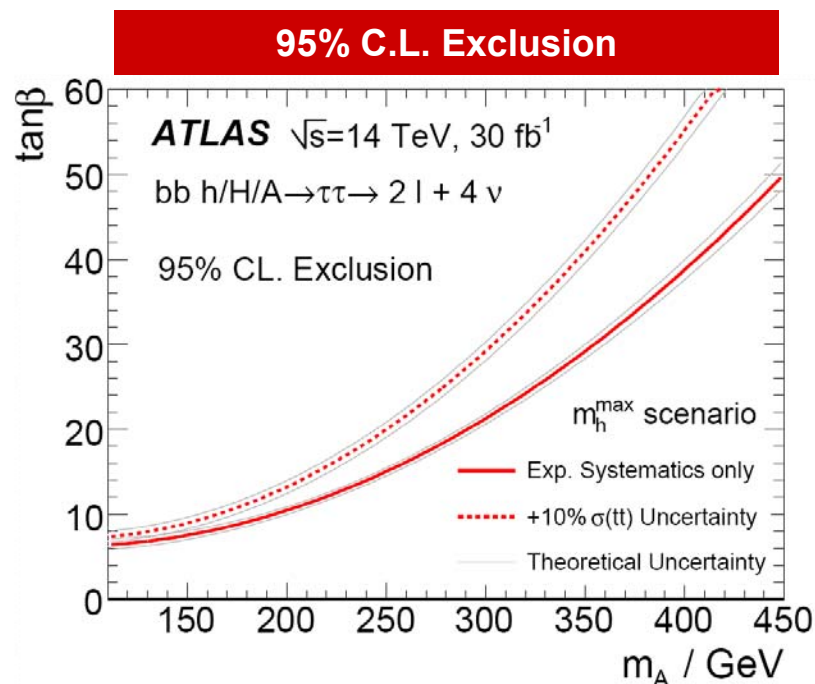
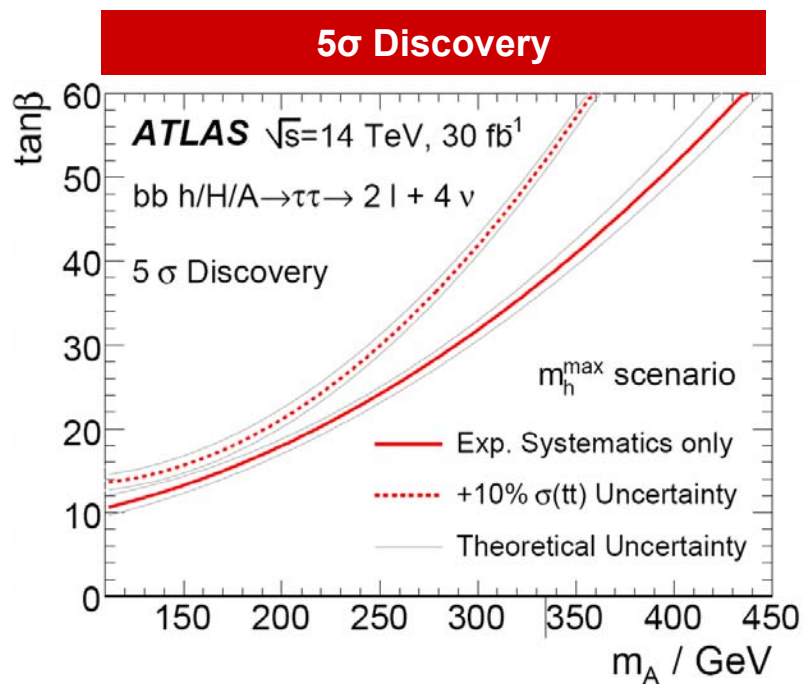
Di-Tau Analysis Potential

The high $\tan\beta$, low m_A region is well covered with 30 fb^{-1}

- Counting experiment with multiple mass windows

Dominant systematic uncertainties

- Jet resolution and energy scale
- b-jet identification



MSSM Higgs Di-Muon Analysis

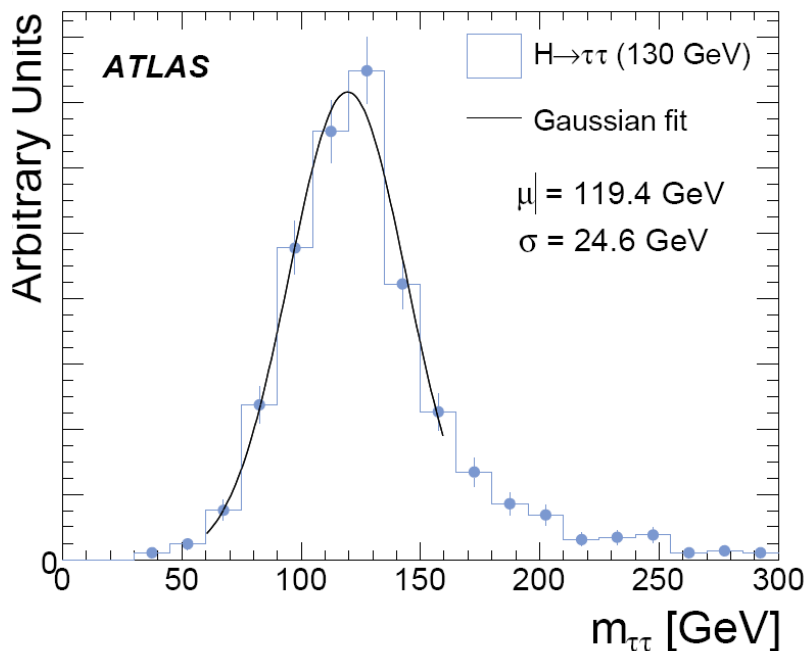
Some advantages

- Cleaner signal than the di-tau analysis
- Excellent mass resolution ($\sim 3\%$ versus $\sim 20\%$ for the di-tau)

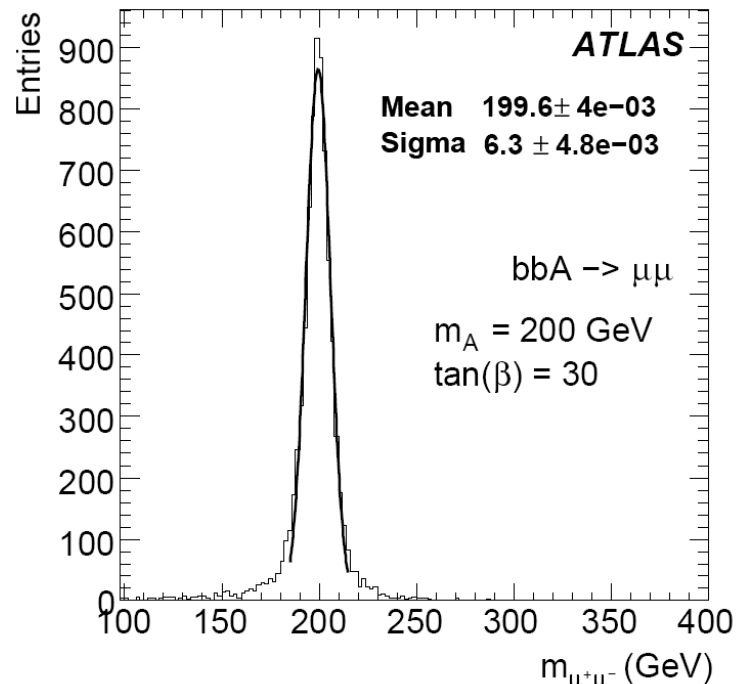
Disadvantage

- $h/A/H$ di-muon branching ratio is $\sim 300\%$ smaller than that of the di-tau

MSSM $h/A/H \rightarrow \tau\tau$ mass (collinear)



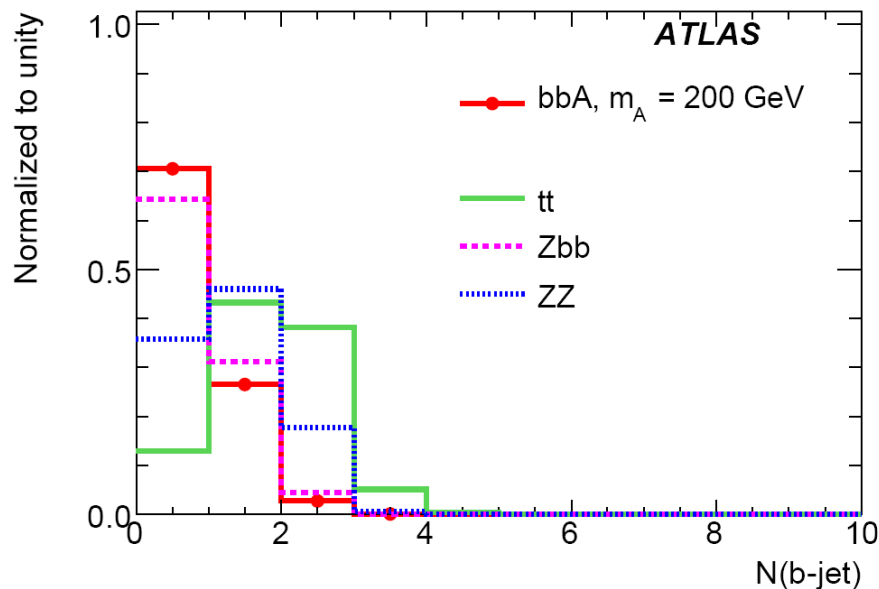
MSSM $h/A/H \rightarrow \mu\mu$ mass



MSSM Higgs Di-Muon Analysis

Divide the analysis into two uncorrelated channels

- 0 b-jets channel (to suppress the $t\bar{t}$ background)
- ≥ 1 b-jets channel (suppress the Z background; impose additional cuts to reduce $t\bar{t}$)



Data-driven background estimation

- For higher masses the tail of the Z resonance provides a large irreducible background, sensitive to detector systematic effects
- $\text{BR}(h/A/H \rightarrow ee) \sim 0$
- $\text{BR}(Z \rightarrow \mu\mu) = \text{BR}(Z \rightarrow ee)$, so use $Z \rightarrow ee$ events from data as a control sample

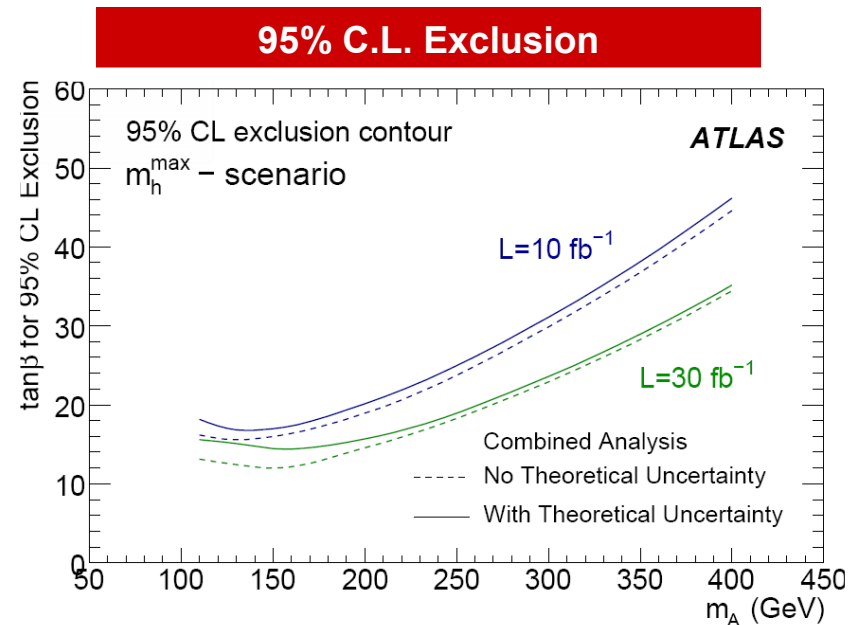
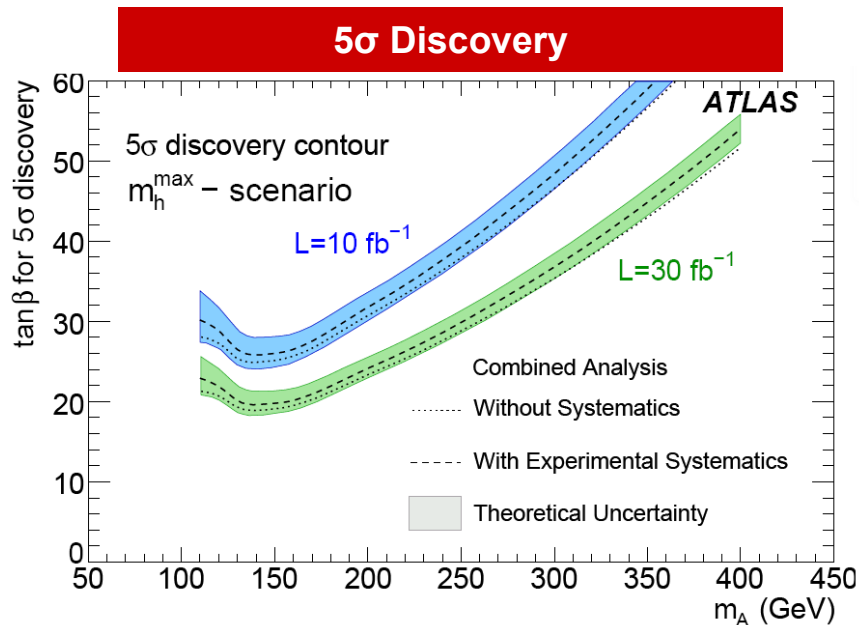
Di-Muon Analysis Potential

Less coverage than the di-tau analysis

- But the two analyses could be combined to increase the sensitivity

Systematic uncertainties

- Around 5 – 10% for the signal processes
- Predominantly from the jet energy scale and b-jet identification
- Systematic uncertainties degrade the signal significance by up to 20% at large values of $\tan\beta$



Charged MSSM Higgs Discovery Potential

Charged MSSM Higgs

Production mode greatly depends on m_{H^\pm}

Three different analyses for a low mass ($m_{H^\pm} < m_{\text{top}}$)

- $t\bar{t} \rightarrow bH^\pm bW \rightarrow b\tau_H \nu bqq$
- $t\bar{t} \rightarrow bH^\pm bW \rightarrow b\tau_L \nu bqq$
- $t\bar{t} \rightarrow bH^\pm bW \rightarrow b\tau_H \nu bl\nu$

Two analyses considered for a high mass ($m_{H^\pm} > m_{\text{top}}$)

- Production via: $gg \rightarrow H^\pm tb$ and $gb \rightarrow H^\pm t$
- Decay modes:

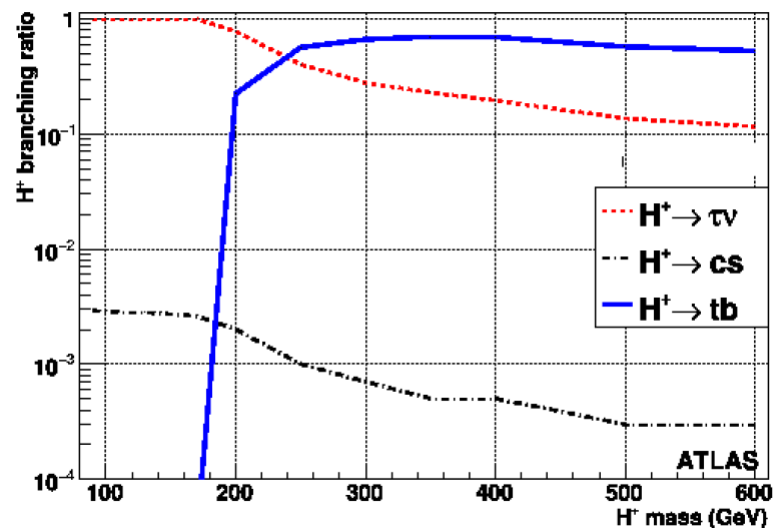
$$H^\pm t \rightarrow \nu\tau_H bqq$$

$$H^\pm t \rightarrow tbt \rightarrow bWbbW \rightarrow bqqbb\nu$$

Dominant Backgrounds

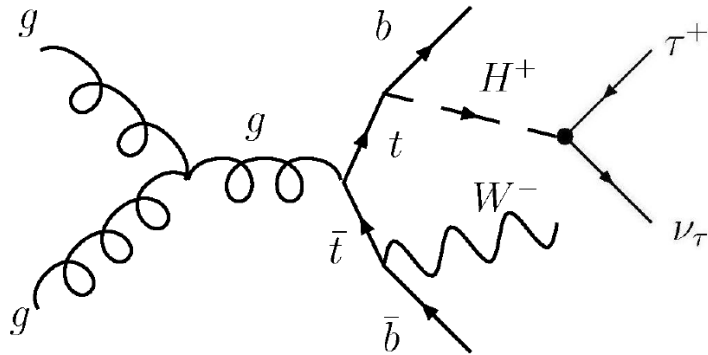
- $t\bar{t}$ (primary)
- QCD di-jets
- W+jets
- Single top

“ m_h -max” scenario with $\tan\beta = 35$



Data-driven Background Estimation

Signal Final State

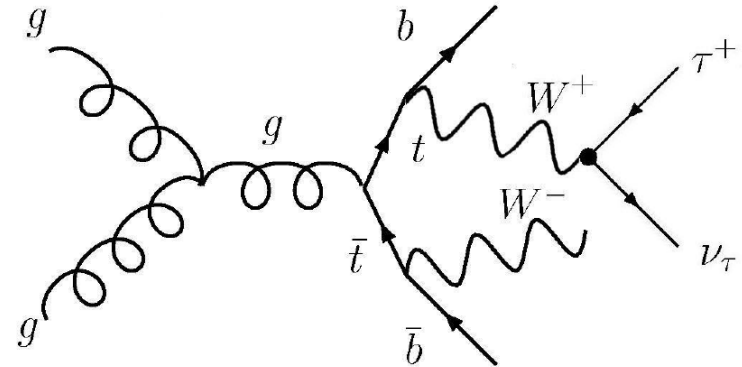


$$H^+ \rightarrow \tau_H \nu; W \rightarrow qq$$

$$H^+ \rightarrow \tau_L \nu; W \rightarrow qq$$

$$H^+ \rightarrow \tau_H \nu; W \rightarrow l\nu$$

Dominant Background



$$W \rightarrow \tau_H \nu; W \rightarrow qq$$

$$W \rightarrow \tau_L \nu; W \rightarrow qq$$

$$W \rightarrow \tau_H \nu; W \rightarrow l\nu$$

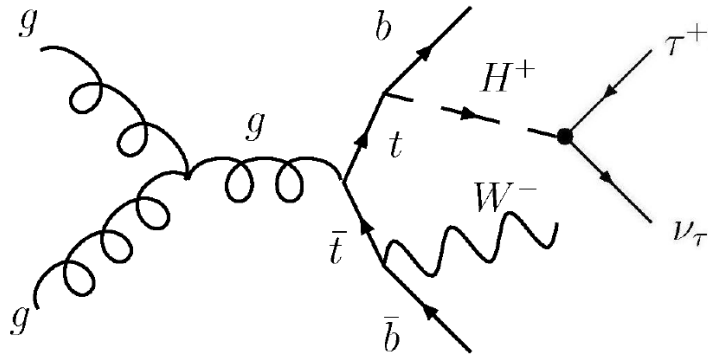
Do not trust Tevatron extrapolations

Difficult to obtain clean samples from data

Unknowns related to analysis-specific variables exist

Data-driven Background Estimation

Signal Final State



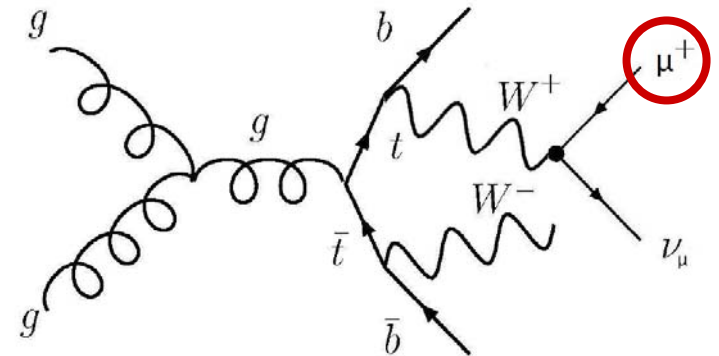
$$H^+ \rightarrow \tau_H \nu; W \rightarrow qq$$

$$H^+ \rightarrow \tau_L \nu; W \rightarrow qq$$

$$H^+ \rightarrow \tau_H \nu; W \rightarrow l\nu$$

Background Control Sample

Change muons into taus
using the TAUOLA package



Leptonically- and
hadronically-decaying
taus can be emulated

Does not rely on the Tevatron

Clean samples can be obtained from data

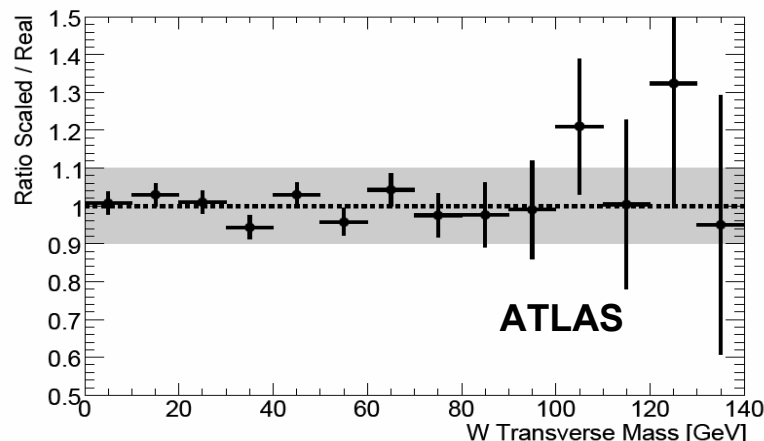
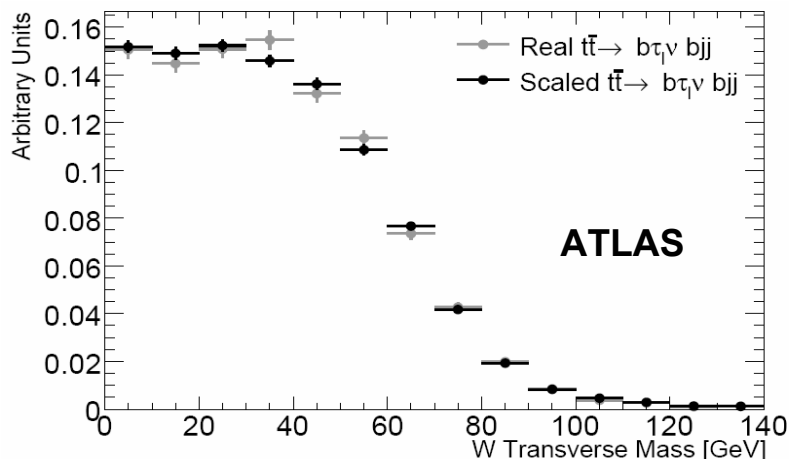
Unknowns related to analysis-specific variables included

Data-Driven Background Estimation

W Transverse Mass (complex quantity; relevant correlations preserved)

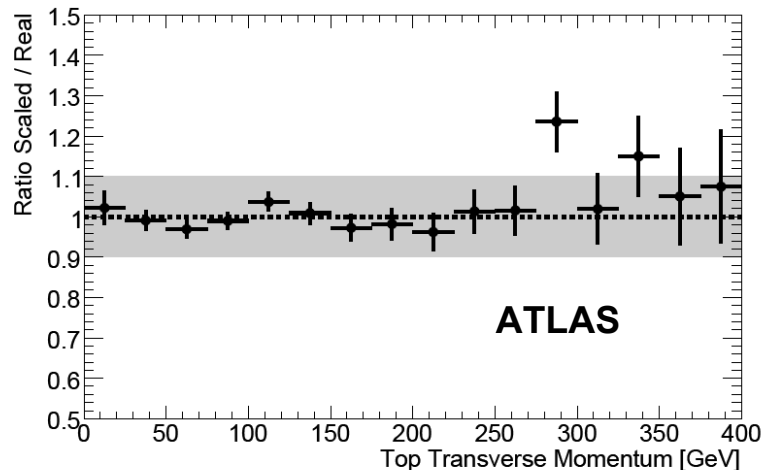
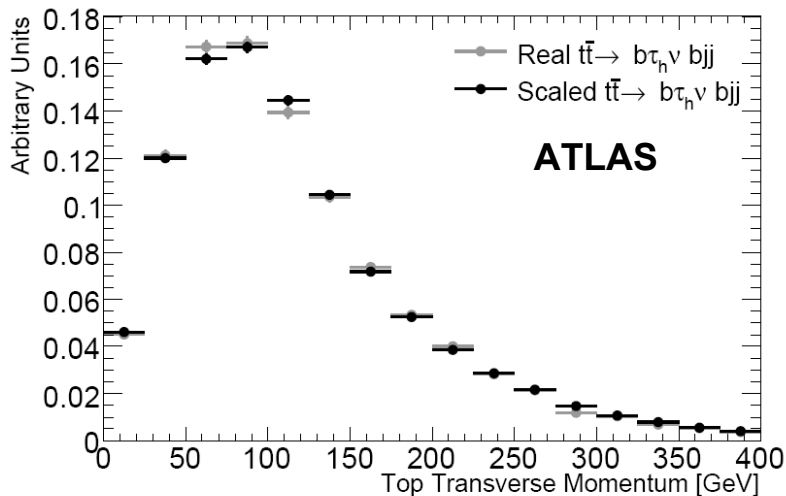
- Leptonically decaying tau ($t\bar{t} \rightarrow b\tau_L\nu bqq$)

$$m_T = \sqrt{2p_T^l p_T^{miss} (1 - \cos(\Delta\phi))}$$



Top Quark Transverse Momentum (complex quantity)

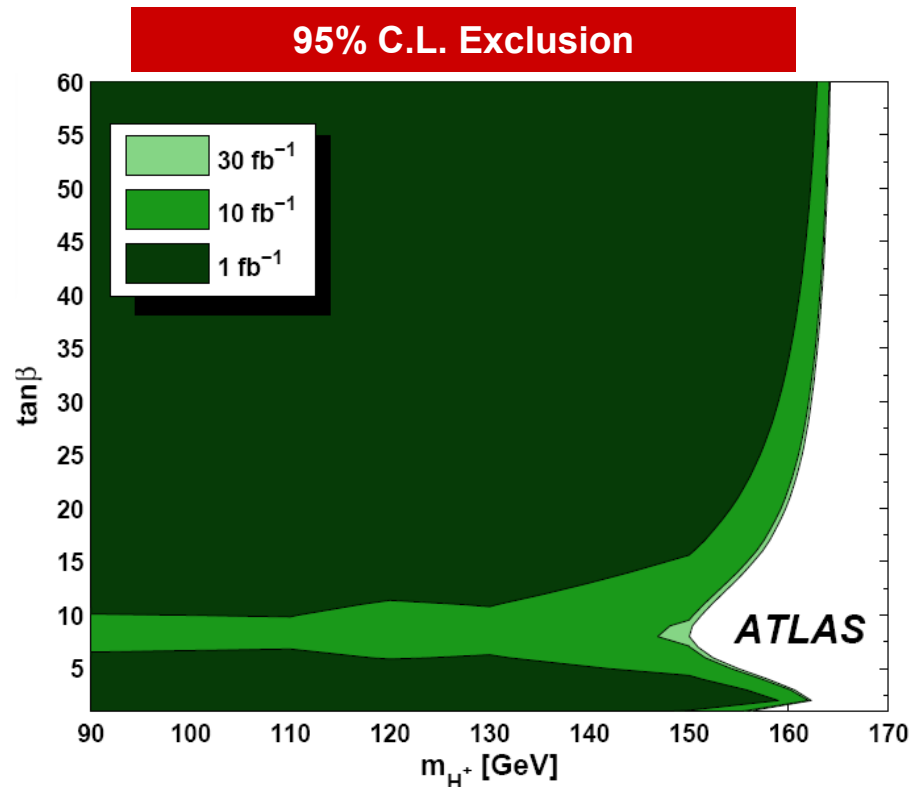
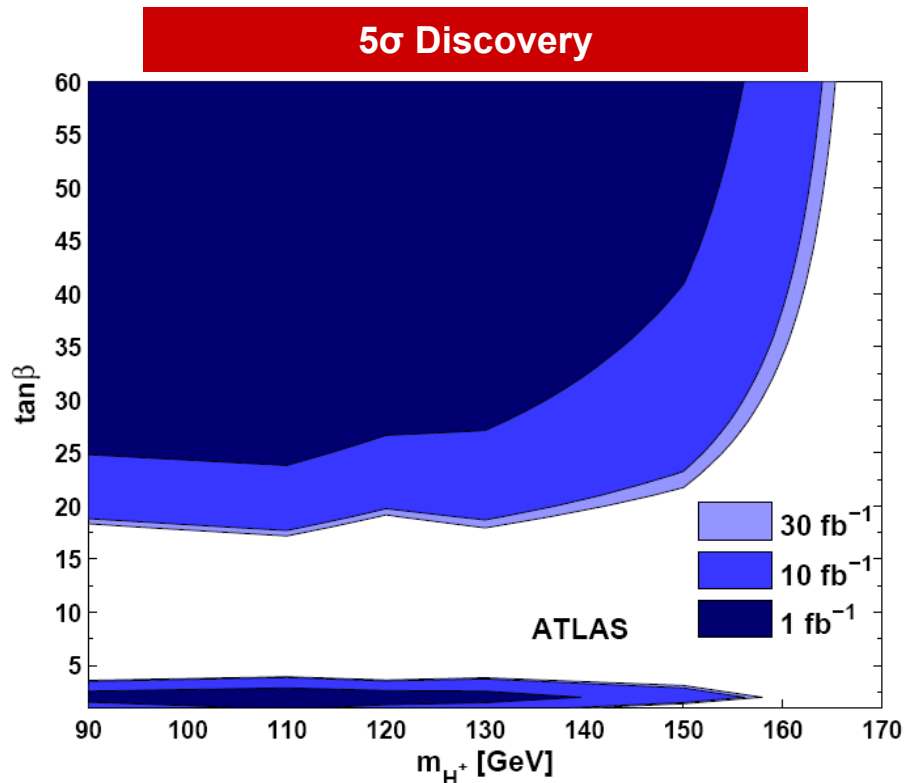
- Hadronically decaying tau ($t\bar{t} \rightarrow b\tau_H\nu bqq$)



Light H^\pm Discovery Potential

Individual analysis cuts vary depending on the final state

- Most promising is $t\bar{t} \rightarrow bH^\pm bW \rightarrow b\tau_H\nu bqq$ due to the large branching fractions into this final state; also challenging due to the high hadronic activity and lack of leptons

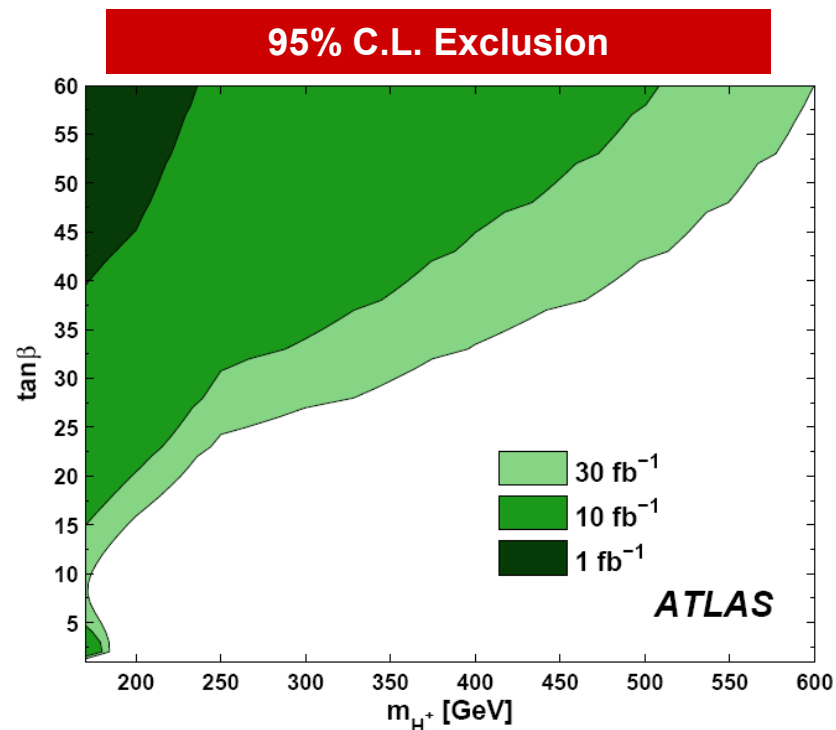
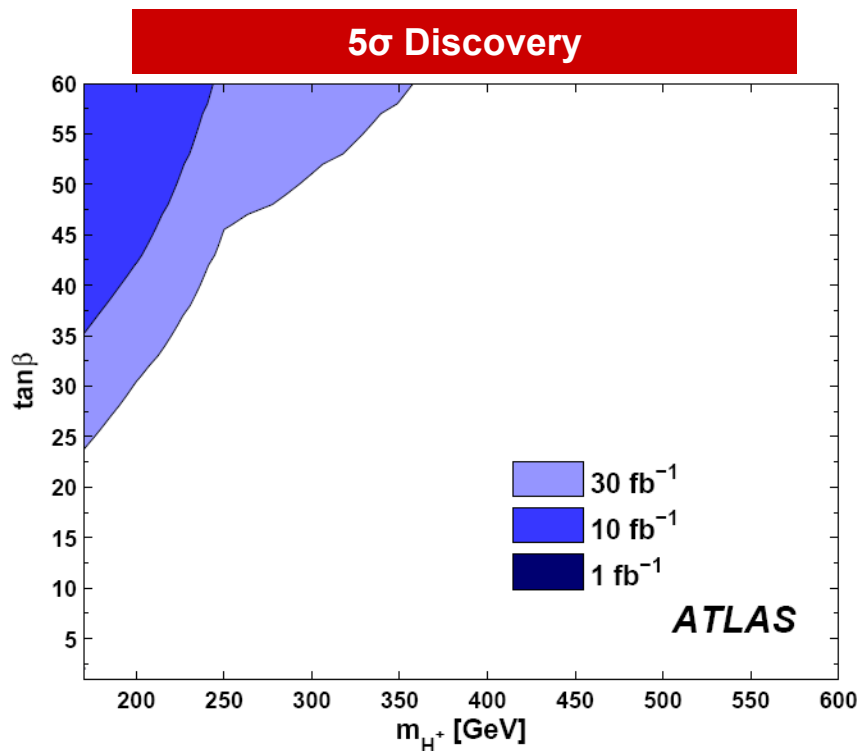


- The other final states contain one charged lepton; exploit signal and background kinematics to get the upper hand on the backgrounds

Heavy H^\pm Discovery Potential

Individual analysis cuts vary depending on the decay of top and H^\pm

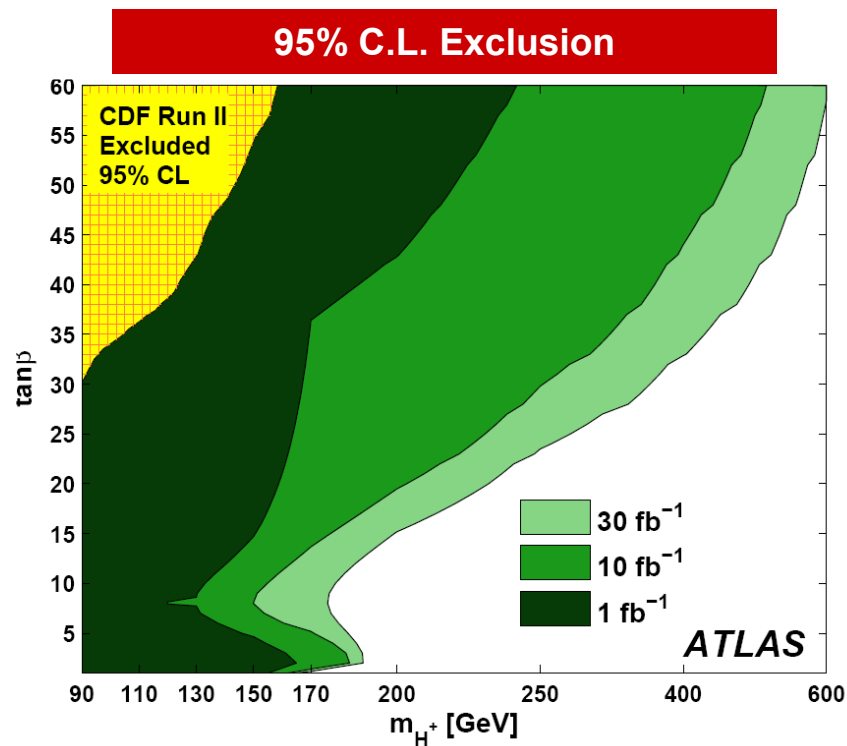
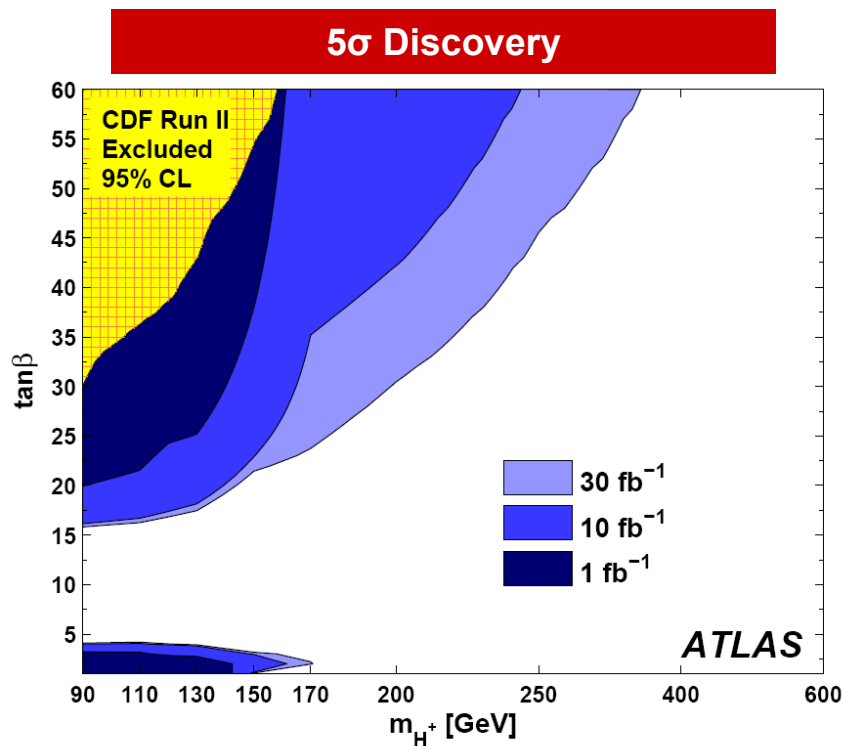
- For $H^\pm t \rightarrow \nu \tau_H b q q$ cut on the quality of the reconstructed top and W boson; use likelihood background discrimination based on the hadronic tau and MET



- For $H^\pm t \rightarrow t b t \rightarrow b W b b W \rightarrow b q q b b l \nu$ jet assignment combinatorics make this channel difficult; reduce the background by reconstructing the W and top quark; a combinatorial likelihood analysis is used

Combined H^\pm Discovery Potential

Good sensitivity for high $\tan\beta$ and low m_{H^\pm} even with 1 fb^{-1} of data



H^\pm is invisible in the so-called “wedge region” of intermediate $\tan\beta$ where the charged Higgs cross-section is at a minimum

Conclusions

Neutral Higgs discovery potential

- With 10 - 30 fb⁻¹ we have good discovery potential for high $\tan\beta$ and low m_A
- Of the results shown here, the di-tau analysis has the best sensitivity
- Discovery potential in other final states are currently being investigated

Charged Higgs discovery potential

- Decays of the H^\pm to a tau and a neutrino offer the best sensitivity for light and heavy charged Higgs bosons in ATLAS; good discovery potential for 1 – 30 fb⁻¹
- Other final states are being investigated here as well (e.g., via decay to a chargino and neutralino)

Data-driven background estimation

- Each analysis shown here contains data-driven methods for estimating dominant and irreducible backgrounds
- Further refinements of these studies are currently underway

Results shown here are for 14 TeV with 1 – 30 fb⁻¹ of data

- Studies currently underway in ATLAS to evaluate non-Standard Model Higgs boson exclusion and discovery with both a reduced center-of-mass energy and less integrated luminosity

Backup Slides

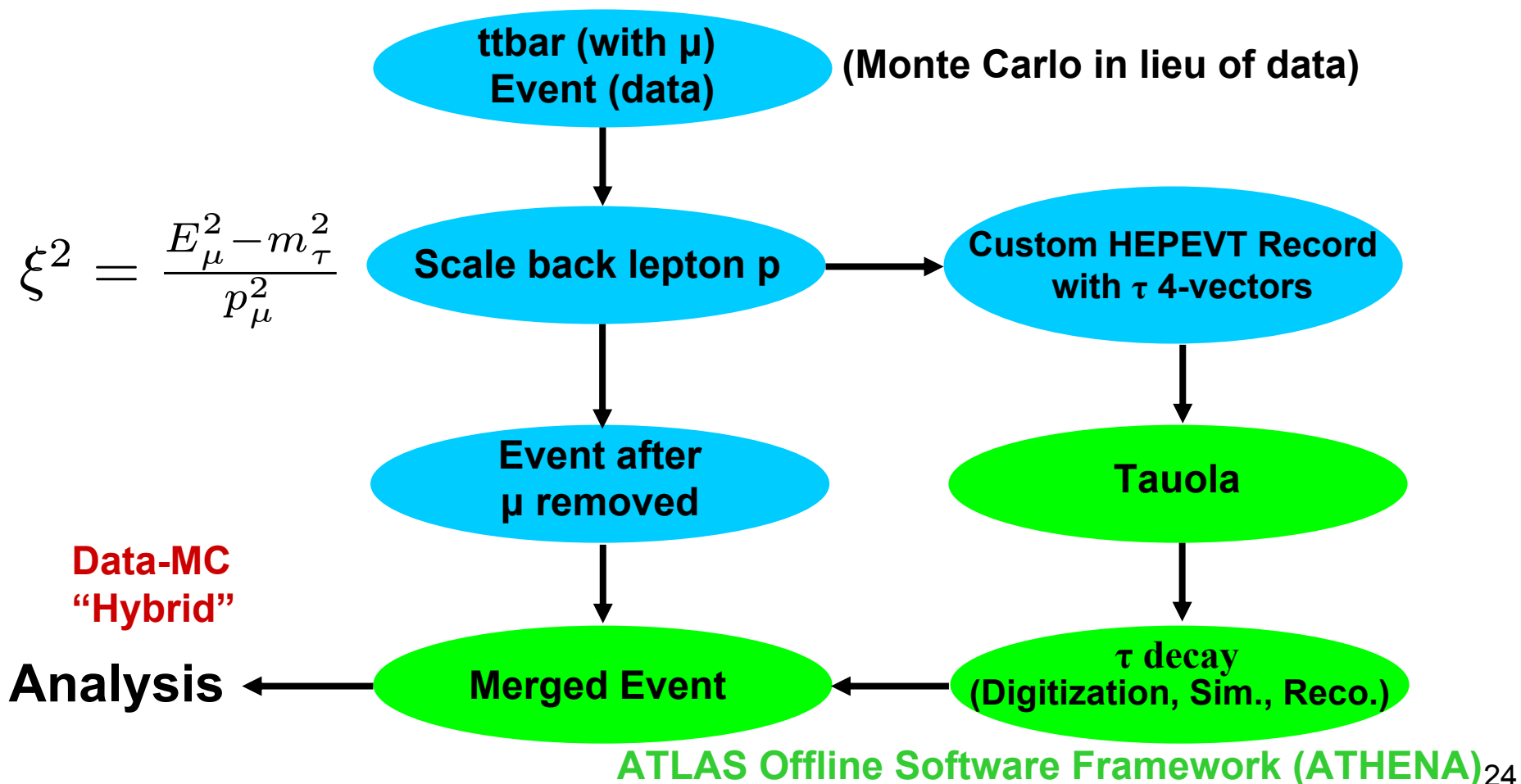
Data-driven Background Method

Based on a method used in ATLAS for SM and MSSM neutral Higgs searches

- Generate control samples for the Z+jets backgrounds

Original implementation (ATLAS CSC studies)

- Done at the ntuple-level and used the full ATLAS detector simulation
- Applicable to many different final states



“Best-fit” Supersymmetric Spectra

The Constrained MSSM (CMSSM) predicts A/H masses ~ 425 GeV

In the single-parameter Non-Universal Higgs Model (NUHM1) A/H ~ 300 GeV

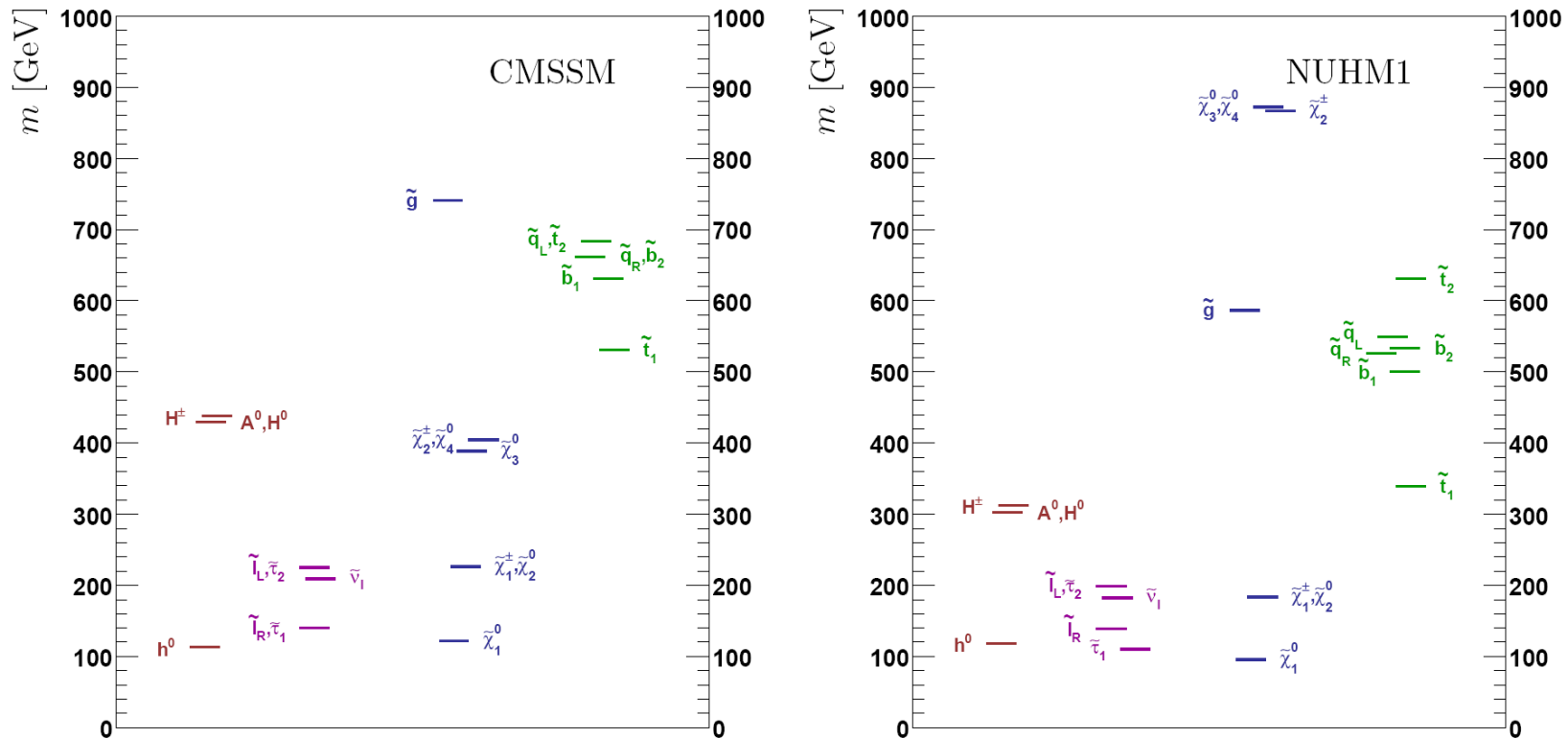


Figure 2. The spectra at the best-fit points: left — in the CMSSM with $m_0 = 60$ GeV, $m_{1/2} = 310$ GeV, $A_0 = 240$ GeV, $\tan\beta = 11$, and right — in the NUHM1 with $m_0 = 100$ GeV, $m_{1/2} = 240$ GeV, $A_0 = -930$ GeV, $\tan\beta = 7$, $m_H^2 = -6.9 \times 10^5$ GeV² and $\mu = 870$ GeV.

Figure taken from O. Buchmueller et al., [arXiv:0707.3447v2](https://arxiv.org/abs/0707.3447v2) [hep-ph]

mh-max Scenario

Evolved out of the LEP2 era

- Extremely common in the literature (e.g., PDG review)
- Using this scenario will allow for easy / direct comparison with previously published results

$$m_t = 174.3 \text{ GeV}, \quad M_{SUSY} = 1 \text{ TeV}, \quad \mu = 200 \text{ GeV}, \quad M_2 = 200 \text{ GeV},$$
$$X_t^{\text{OS}} = 2 M_{SUSY} \text{ (FD calculation)}, \quad X_t^{\overline{\text{MS}}} = \sqrt{6} M_{SUSY} \text{ (RG calculation)}$$
$$A_b = A_t, \quad m_{\tilde{g}} = 0.8 M_{SUSY} .$$

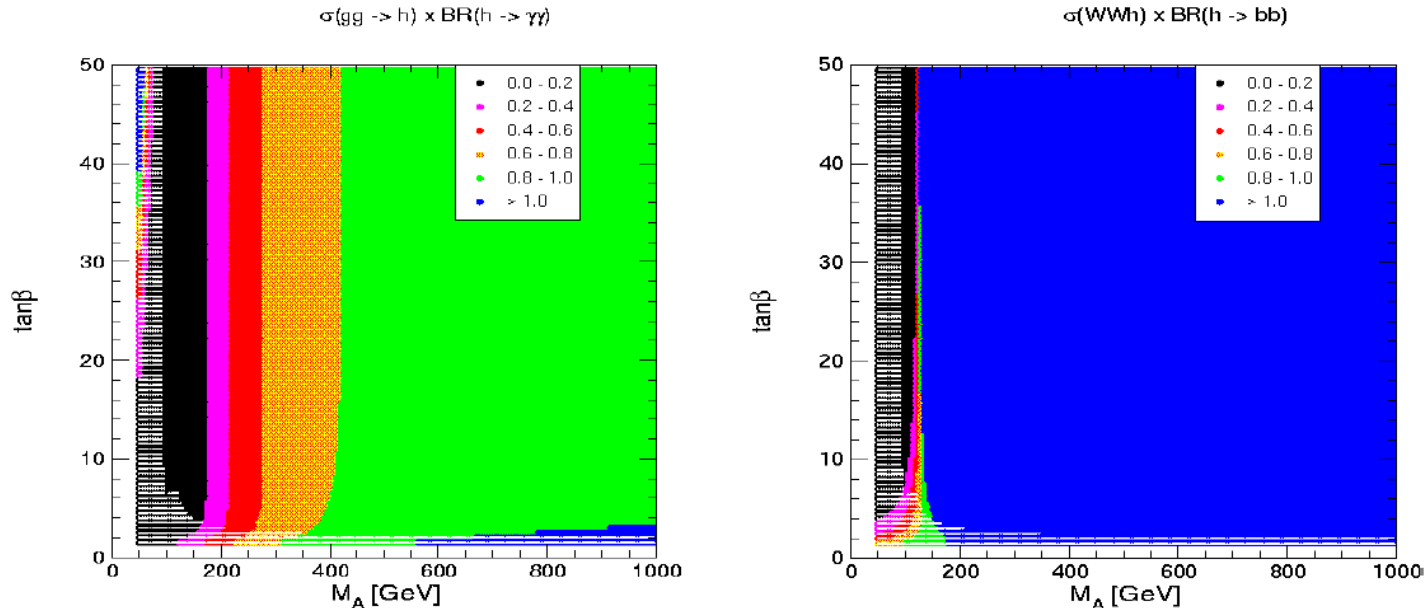


Figure 1: $[\sigma \times \text{BR}]_{\text{MSSM}} / [\sigma \times \text{BR}]_{\text{SM}}$ is shown for the channels $gg \rightarrow h \rightarrow \gamma\gamma$ (left plot) and $W^* \rightarrow Wh \rightarrow Wb\bar{b}$ (right plot) in the $M_A - \tan\beta$ -plane for the m_h^{max} scenario. The white-dotted area is excluded by LEP Higgs searches.

No-Mixing Scenario

Evolved out of the LEP2 era

$$m_t = 174.3 \text{ GeV}, \quad M_{SUSY} = 2 \text{ TeV}, \quad \mu = 200 \text{ GeV}, \quad M_2 = 200 \text{ GeV}, \\ X_t = 0 \text{ (FD/RG calculation)}, \quad A_b = A_t, \quad m_{\tilde{g}} = 0.8 M_{SUSY}.$$

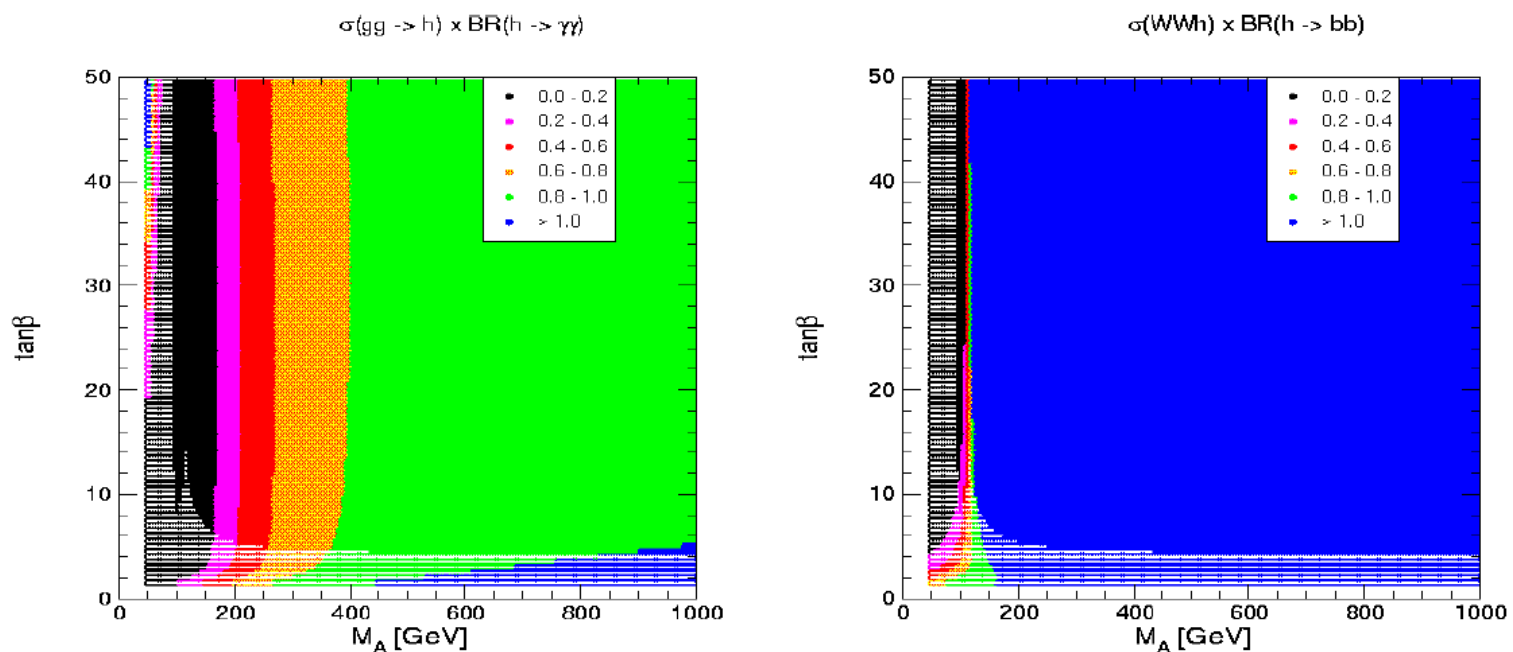


Figure 2: $[\sigma \times BR]_{MSSM}/[\sigma \times BR]_{SM}$ is shown for the channels $gg \rightarrow h \rightarrow \gamma\gamma$ (left plot) and $W^* \rightarrow Wh \rightarrow Wb\bar{b}$ (right plot) in the $M_A - \tan\beta$ -plane for the no-mixing scenario. The white-dotted area is excluded by LEP Higgs searches.

Gluophobic Higgs

Hadron collider

$$m_t = 174.3 \text{ GeV}, \quad M_{SUSY} = 350 \text{ GeV}, \quad \mu = 300 \text{ GeV}, \quad M_2 = 300 \text{ GeV},$$

$$X_t^{\text{OS}} = -750 \text{ GeV} \text{ (FD calculation)}, \quad X_t^{\overline{\text{MS}}} = -770 \text{ GeV} \text{ (RG calculation)}$$

$$A_b = A_t, \quad m_{\tilde{g}} = 500 \text{ GeV}.$$

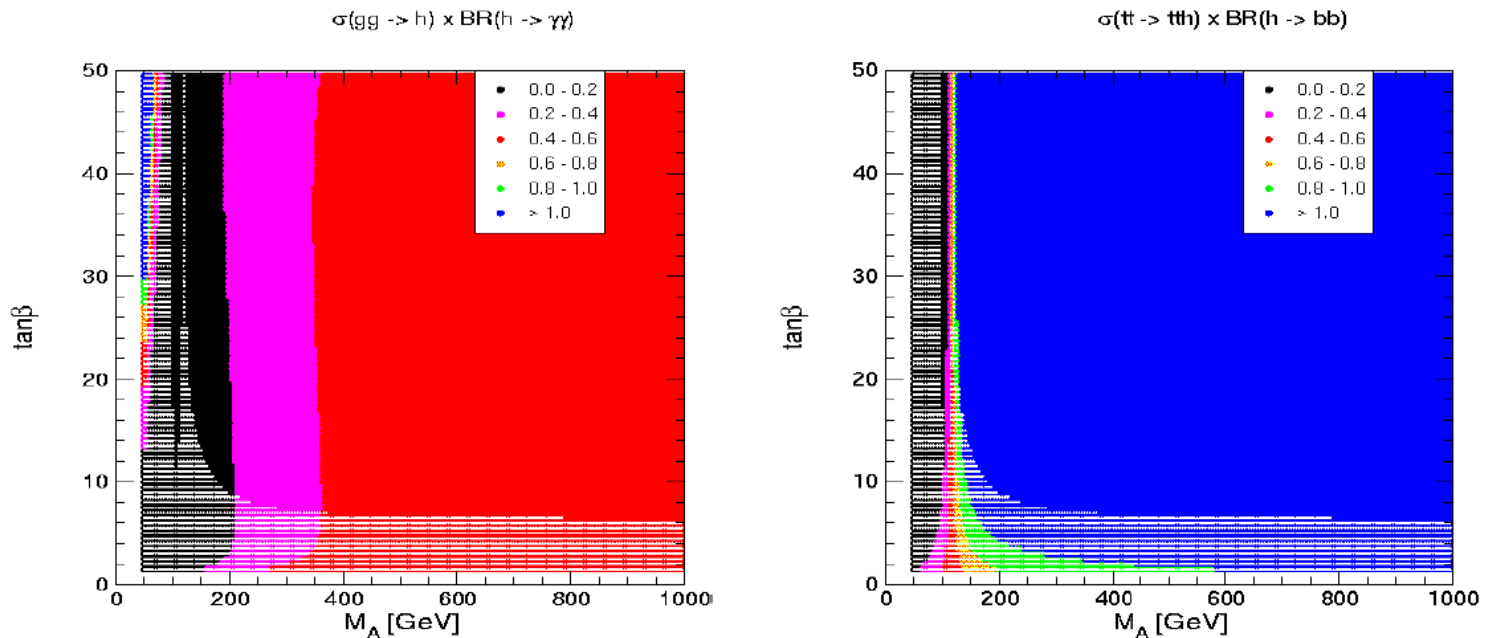


Figure 3: $[\sigma \times \text{BR}]_{\text{MSSM}}/[\sigma \times \text{BR}]_{\text{SM}}$ is shown for the channels $gg \rightarrow h \rightarrow \gamma\gamma$ (left plot) and $t\bar{t} \rightarrow t\bar{t}h \rightarrow t\bar{t}b\bar{b}$ (right plot) in the $M_A - \tan\beta$ -plane for the gluophobic Higgs scenario. The white-dotted area is excluded by LEP Higgs searches.

Small α_{eff}

Hadron collider

$$m_t = 174.3 \text{ GeV}, \quad M_{SUSY} = 800 \text{ GeV}, \quad \mu = 2.5 M_{SUSY}, \quad M_2 = 500 \text{ GeV},$$
$$X_t^{\text{OS}} = -1100 \text{ GeV} \text{ (FD calculation)}, \quad X_t^{\overline{\text{MS}}} = -1200 \text{ GeV} \text{ (RG calculation)}$$
$$A_b = A_t, \quad m_{\tilde{g}} = 500 \text{ GeV}.$$

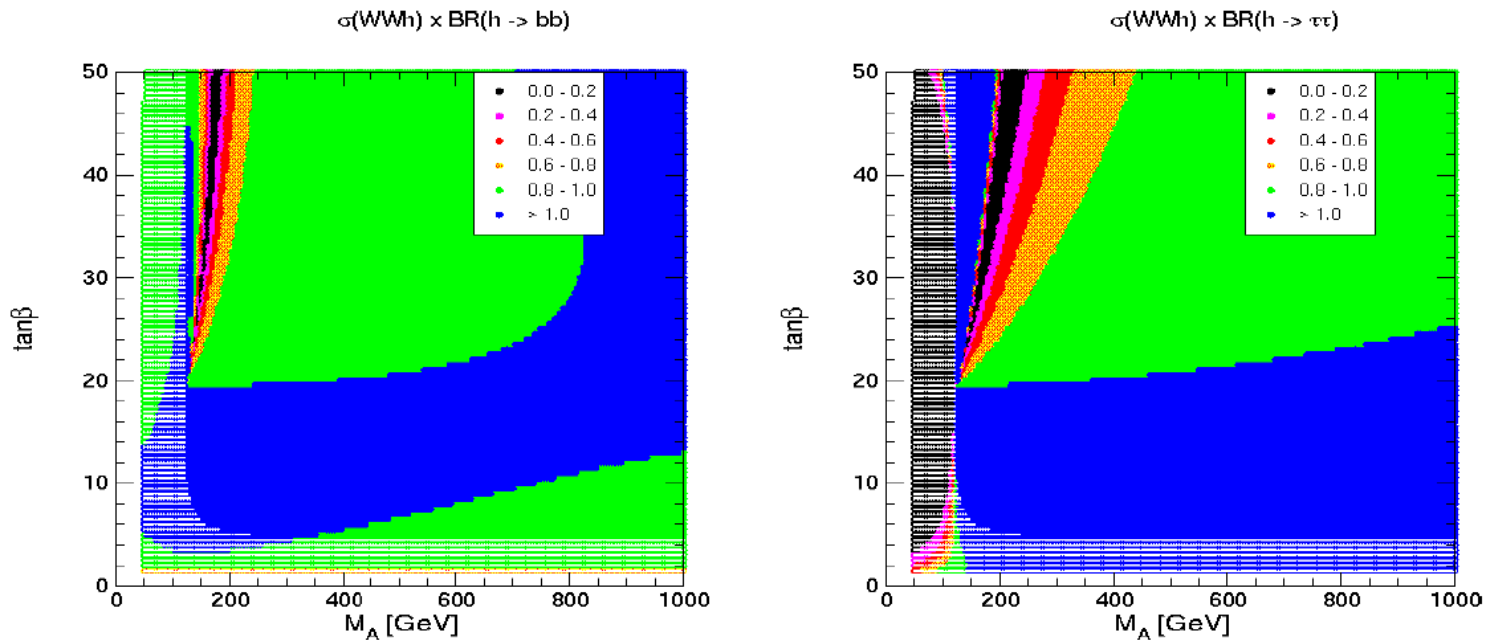


Figure 4: $[\sigma \times \text{BR}]_{\text{MSSM}} / [\sigma \times \text{BR}]_{\text{SM}}$ is shown for the channels $W^* \rightarrow Wh \rightarrow Wb\bar{b}$ (left plot) and $W^* \rightarrow Wh \rightarrow W\tau^+\tau^-$ (right plot) in the $M_A - \tan\beta$ -plane for the small α_{eff} scenario. The white-dotted area is excluded by LEP Higgs searches.