

Analysis of $B_s^0 \rightarrow \phi\phi$ decay mode

Benedetto Di Ruzza for the CDF Collaboration



D.P.F. MEETING 2009

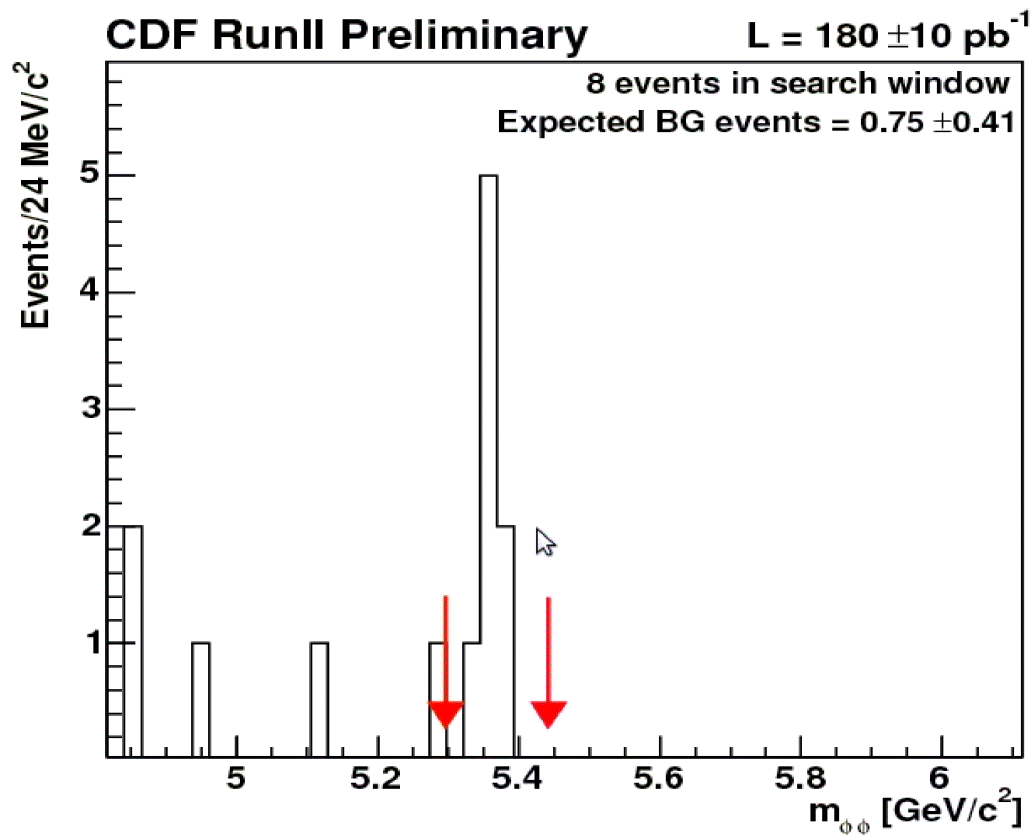
Detroit July 28 2009



INTRODUCTION

The $B_s^0 \rightarrow \phi \phi$ decay was observed for the first time by CDFII in 2005:
D. Acosta et al. (CDF Collaboration), Phys. Rev. Lett. 95,031801 (2005).

Using 180 pb⁻¹ of data available at that time 8 events were observed.





INTRODUCTION

Now this measurement was done using 2.9 fb⁻¹ of data available at CDFII (2002-up to April 2008).

In this talk I will describe the motivations, the strategy, the result and the perspectives of this measurement.



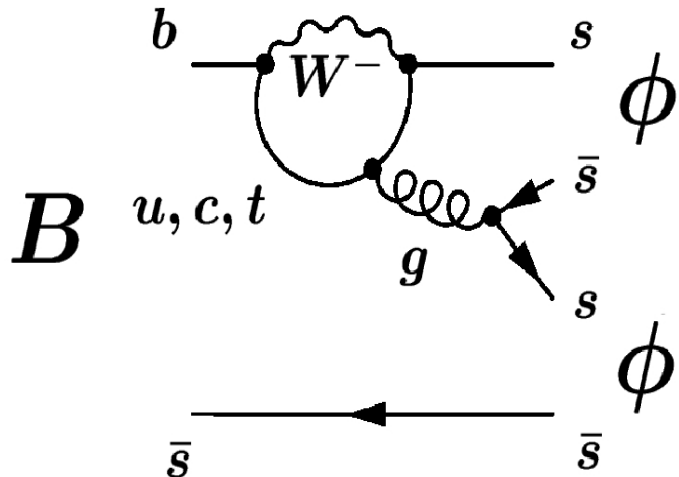
OUTLINE

- 1) Motivation.
- 2) Theoretical predictions.
- 3) Analysis strategy.
- 4) Signal selection.
- 5) Background evaluation.
- 6) Branching ratio measurement.
- 7) Conclusions and perspectives.



MOTIVATION

The $B_s^0 \rightarrow \phi \phi$ is a vector vector decay mode



The dominant decay element is the **penguin** $b \rightarrow s \bar{s} s$

The final state is self - conjugate.
 It can be used to measure $\Delta \Gamma_s$, CKM studies,
 and test of decay polarization prediction.



MOTIVATION

Considering a B meson with four-momentum p_B decaying into light vector mesons $V_1(p_1, \eta^*)$, $V_2(p_2, \epsilon^*)$, with masses $m_{1,2}$ of order Λ_{QCD} , the decay amplitude can be decomposed into three scalar amplitude $S_{1,2,3}$ according to

$$\mathcal{A}_{B \rightarrow V_1 V_2} = i\eta^{*\mu}\epsilon^{*\nu} \left(S_1 g_{\mu\nu} - S_2 \frac{p_{B\mu} p_{B\nu}}{m_B^2} + S_3 i\epsilon_{\mu\nu\rho\sigma} \frac{p_1^\rho p_2^\sigma}{p_1 \cdot p_2} \right),$$

or alternatively an helicity basis can be defined as :

$$\begin{aligned} \mathcal{A}_0 &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_0^*) V_2(p_2, \epsilon_0^*)) = \frac{im_B^2}{2m_1 m_2} \left(S_1 - \frac{S_2}{2} \right), \\ \mathcal{A}_\pm &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_\pm^*) V_2(p_2, \epsilon_\pm^*)) = i(S_1 \mp S_2), \end{aligned}$$

or a transversity amplitude basis can be defined using $\mathcal{A}_L = \mathcal{A}_0$, and replacing \mathcal{A}_\pm with

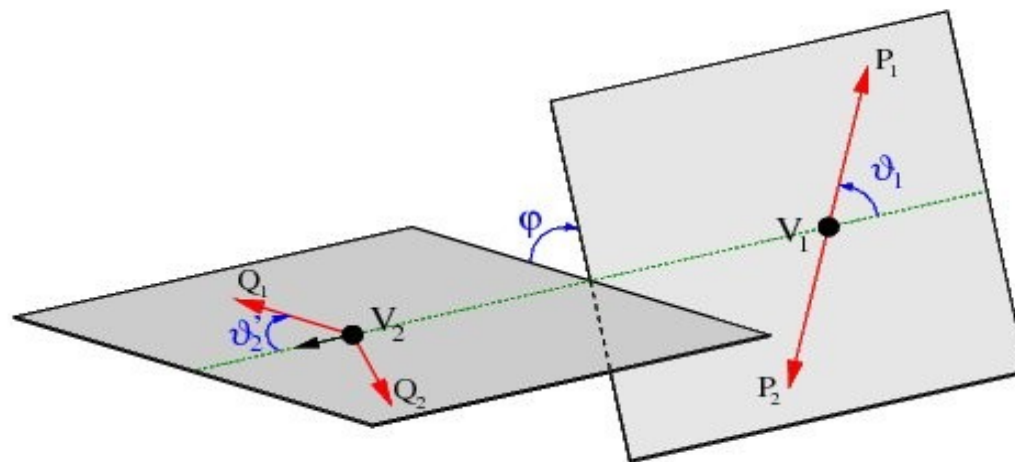
$$\begin{aligned} \mathcal{A}_\parallel &= \frac{(\mathcal{A}_+ + \mathcal{A}_-)}{\sqrt{2}} \\ \mathcal{A}_\perp &= \frac{(\mathcal{A}_+ - \mathcal{A}_-)}{\sqrt{2}} \end{aligned}$$



MOTIVATION

The relation between **differential cross-section** and the observable **angular distribution** in the B meson rest frame using a helicity basis is:

$$\begin{aligned} \frac{d\Gamma_{B \rightarrow V_1 V_2 \rightarrow \dots}}{d\cos\vartheta_1 d\cos\vartheta_2 d\varphi} &\propto |\mathcal{A}_0|^2 \cos^2\vartheta_1 \cos^2\vartheta_2 + \frac{1}{4} \sin^2\vartheta_1 \sin^2\vartheta_2 (|\mathcal{A}_+|^2 + |\mathcal{A}_-|^2) \\ &\quad - \cos\vartheta_1 \sin\vartheta_1 \cos\vartheta_2 \sin\vartheta_2 [\operatorname{Re}(e^{-i\varphi} \mathcal{A}_0 \mathcal{A}_+^*) + \operatorname{Re}(e^{+i\varphi} \mathcal{A}_0 \mathcal{A}_-^*)] \\ &\quad + \frac{1}{2} \sin^2\vartheta_1 \sin^2\vartheta_2 \operatorname{Re}(e^{2i\varphi} \mathcal{A}_+ \mathcal{A}_-^*), \end{aligned}$$





MOTIVATION

In Experimental analyses, observables are preferably defined in terms of the transversity as they have definite CP transformation properties. A typical set of observable consists of the branching fraction, two out of the three polarization fractions $f_L, f_{\parallel}, f_{\perp}$, and two phases $\phi_{\parallel}, \phi_{\perp}$, where

$$f_{L,\parallel,\perp}^B = \frac{|\mathcal{A}_{0,\parallel,\perp}|^2}{|\mathcal{A}_0|^2 + |\mathcal{A}_{\parallel}|^2 + |\mathcal{A}_{\perp}|^2}; \quad \phi_{\parallel,\perp}^B = \arg \frac{\mathcal{A}_{\parallel,\perp}}{\mathcal{A}_0}$$

and naively, within the Standard Model, one expects

$$\frac{f_T}{f_L} \propto \frac{m_V}{m_B}$$

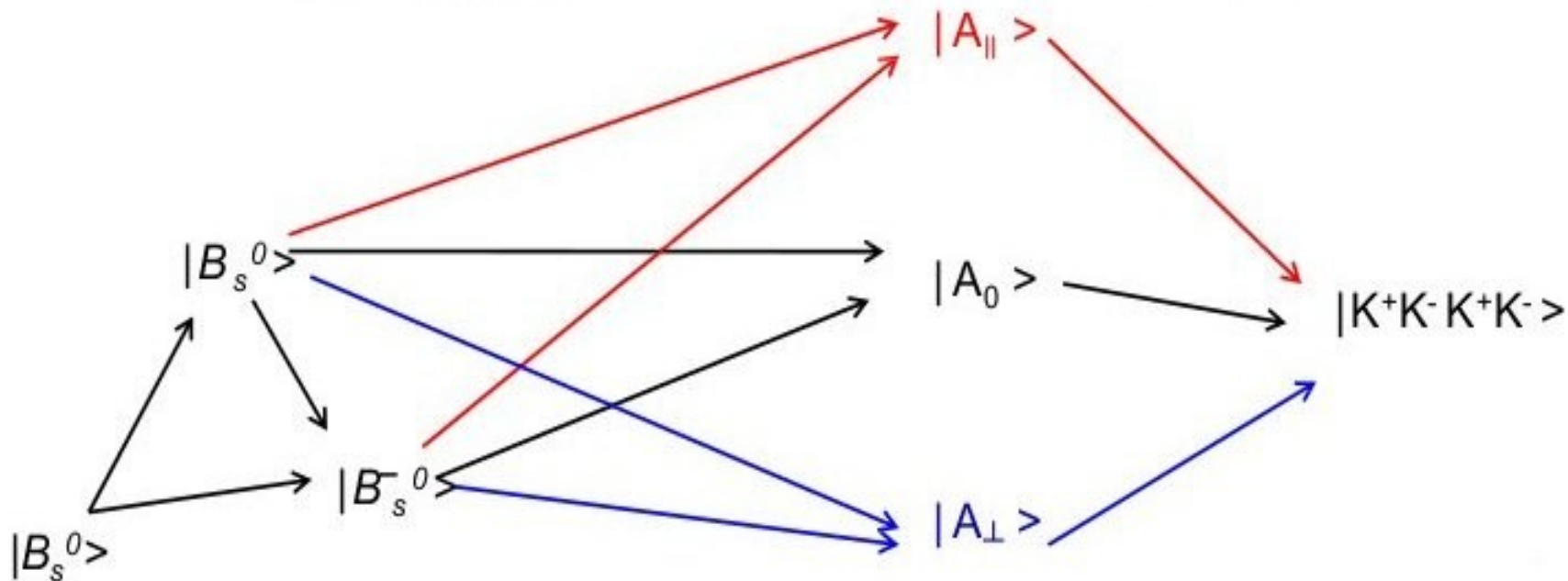


MOTIVATION

In the transversity basis the amplitudes have defined CP value:

Corresponding decay amplitudes: $A_0, A_{\parallel}, A_{\perp}$

- transverse (\perp perpendicular to each other) \rightarrow CP odd
- transverse (\parallel parallel to each other) \rightarrow CP even
- longitudinal (0) \rightarrow CP even





Polarization puzzle

Within the standard model one expect:

$$\frac{f_T}{f_L} \propto \frac{m_V}{m_B}$$

And for the $B_s^0 \rightarrow \phi \phi$ decay: $\frac{f_T}{f_L} \propto \frac{1020}{5369}$

However it was observed in some decays, for example

$$B \rightarrow \varphi K^* \quad \text{and} \quad B \rightarrow \rho K^*$$

that **this relation** is not valid

Babar collaboration : hep-hex: 0705.1798 :

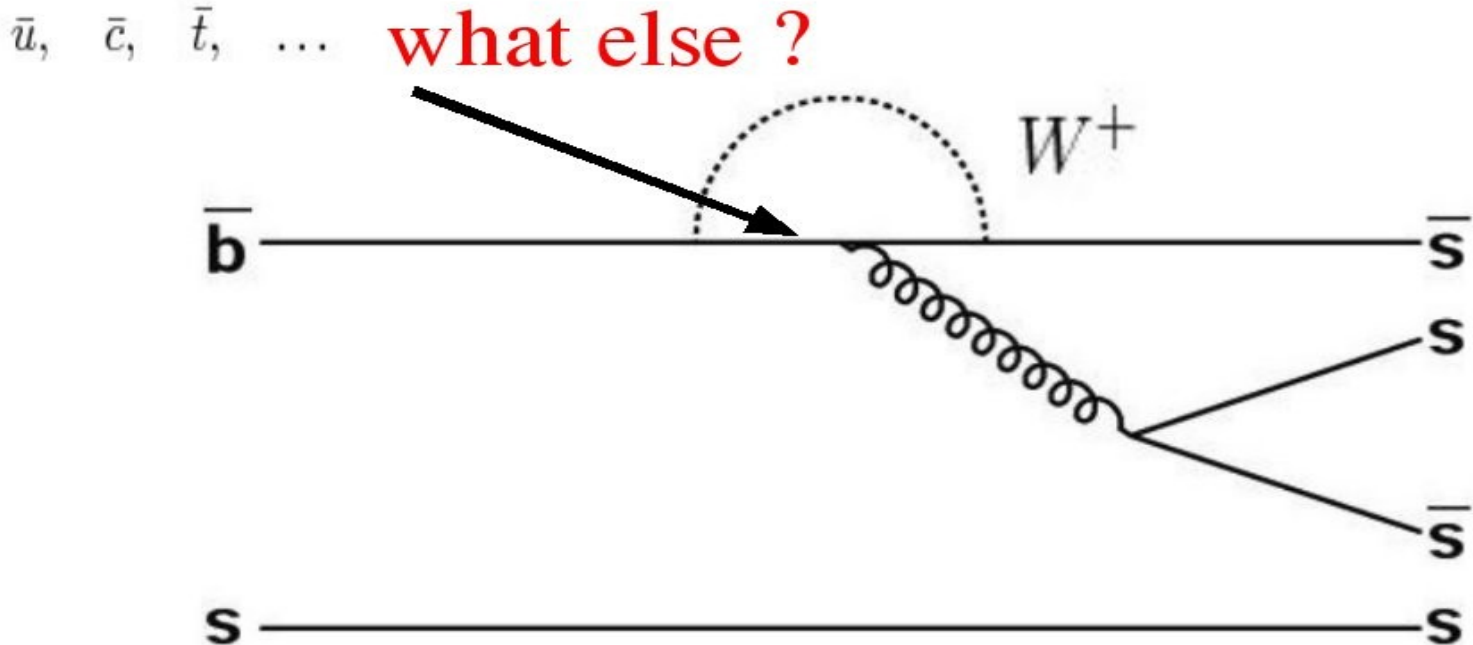
$$f_T = 0.21_{-0.05}^{+0.02} \quad f_L = 0.49_{-0.05}^{+0.03}$$



Polarization puzzle

This result, known also as “Polarization Puzzle”, is more evident in the B meson decays involving penguin diagram decays. Different solutions were proposed for this “puzzle”, beyond and inside the Standard Model.

One solution beyond the SM could be New Physics inside the penguin loop:
(Grossman hep-ph/03110229; Yang et al. hep-ph/0411211)





Polarization puzzle

Others explain this puzzle inside the SM, for example

Ali & Kramer hep-ph/070316, on recent BR and polarization prediction in the perturbative QCD approach.

Fleisher & Gronau in hep-ph/07094013

Some of them do a prevision on the Branching ratio too:

	BR[10 ⁻⁶]
QCD Factorisation	21.8 ^{+1.1+30.4} _{-1.1-17.0}
	19.5 ^{+1.0+13.1} _{-1.0-8.0}
QCD Factorisation	13.1
Naive Factorisation	9.05

M. Beneke et al., hep-hex/0612290.

X. Li et al. P.R.L. D.68, 114015(2003);
D71 019902(2005);
hep-hex/030936.



MEASUREMENT STRATEGY

The measurement consists in the determination of this ratio:

$$\frac{BR(B_s \mapsto \phi\phi)}{BR(B_s \mapsto J/\psi \phi)} = \frac{N_{\phi\phi}}{N_{J/\psi\phi}} \cdot \frac{\epsilon_{(B_s \mapsto J/\psi\phi)}}{\epsilon_{(B_s \mapsto \phi\phi)}} \cdot \frac{BR(J/\psi \mapsto \mu\mu)}{BR(\phi \mapsto KK)} \cdot \epsilon_{\mu}$$

From MC From PDG From data

$$\phi \rightarrow K^+ K^-$$

Where

$$J/\psi \rightarrow \mu\mu$$

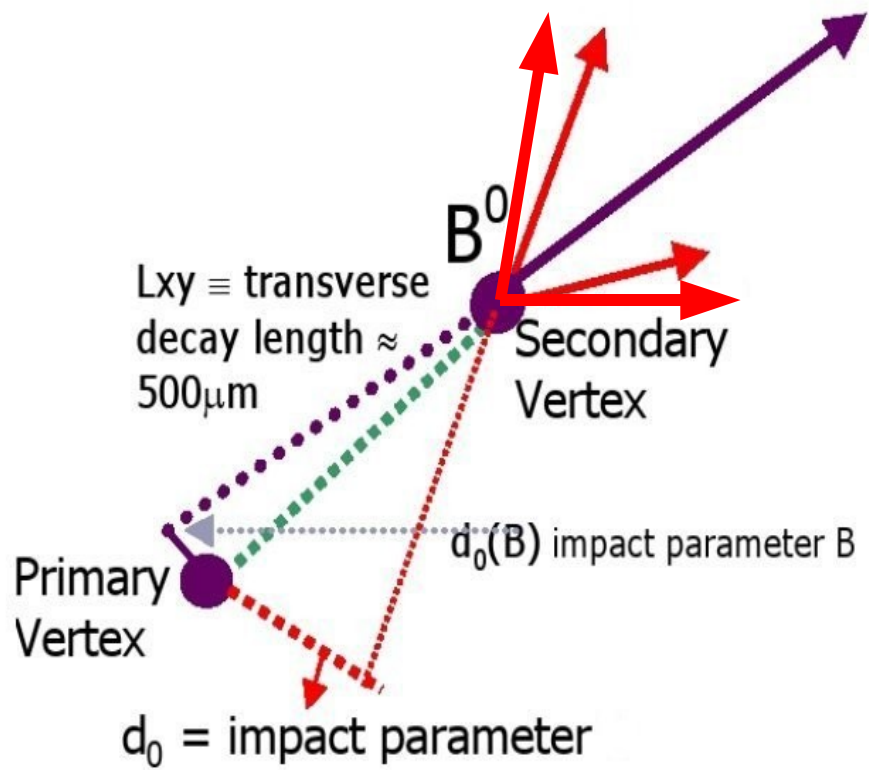
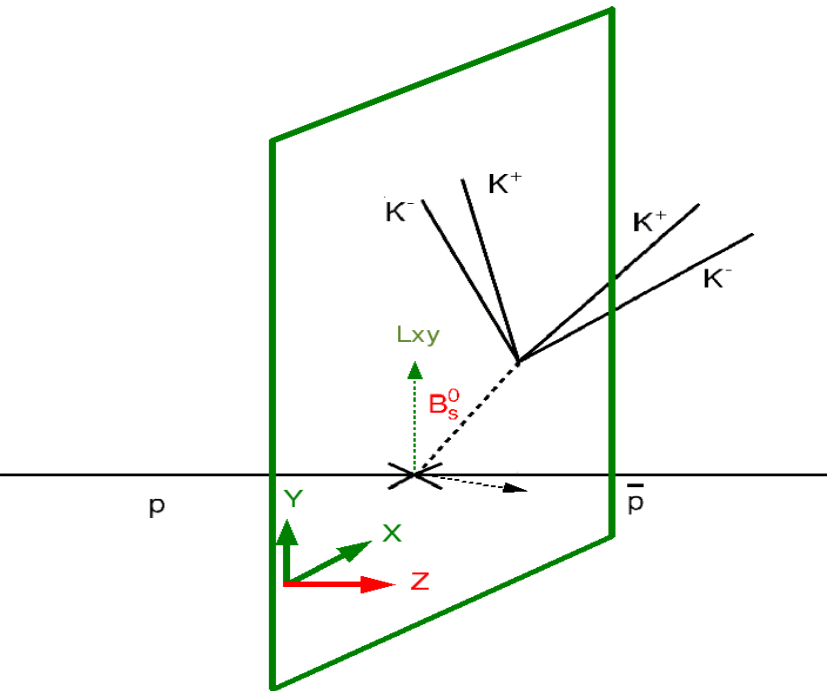
We have performed a normalized branching ratio and not absolute because:

- 1) The two decays are topologically very similar.
- 2) In the ratio we can simplify a lot of sistematics.



Event signature

The events are characterized by 4 charged particle in the final state





Optimization Selection procedure

Event selection is performed using the following variables

- transverse momentum of the B meson: p_T^B ;
- transverse decay length of the B meson: L_{xy} ;
- impact parameter of the B meson: d_0^B ;
- impact parameter of the more energetic ϕ : d_{0max}^ϕ ;
- impact parameter of the less energetic ϕ : d_{0min}^ϕ ;
- transverse momentum of the J/ψ : $p_T^{J/\psi}$;
- transverse momentum of the less energetic kaon from ϕ decay: p_T^K ;
- the bi-dimensional χ^2 of the primary vertex fit: χ_{xy}^2 .



Optimization Selection procedure

In order to choose the variables to be used in the event selection and select the best interval values of these variables, we followed the procedure of the “maximization of the score function”.

maximizes the score function $\frac{S}{\sqrt{S+B}}$ where S is signal events evaluated with MonteCarlo simulation and B is background evaluated from data sideband



Optimization results:

$B_s \longrightarrow \phi\phi$	
Variable	cut
L_{xy}	$> 330\mu m$
$P_T^{K \min}$	$> 0.7 \text{ GeV}/c$
χ_{xy}^2	< 17
$d0(B)$	$< 65\mu m$
$d0_{\max}^{\phi}$	$> 85\mu m$

$B_s \longrightarrow J/\psi\phi$	
Variable	cut
L_{xy}	$> 290\mu m$
P_T^{ϕ}	$> 1.4 \text{ GeV}/c$
χ_{xy}^2	< 15
$d0(B)$	$< 80\mu m$
$P_T^{J/\psi}$	$> 2.0 \text{ GeV}/c$



Backgrounds

Expected background are:

- ▶ combinatorics (expected smooth function);
- ▶ decays not correctly reconstructed (reflections) peaking under the signal:
- ▶ for $B_s \rightarrow J/\psi\phi$:
 - ▶ $B_d \rightarrow J/\psi K^*$

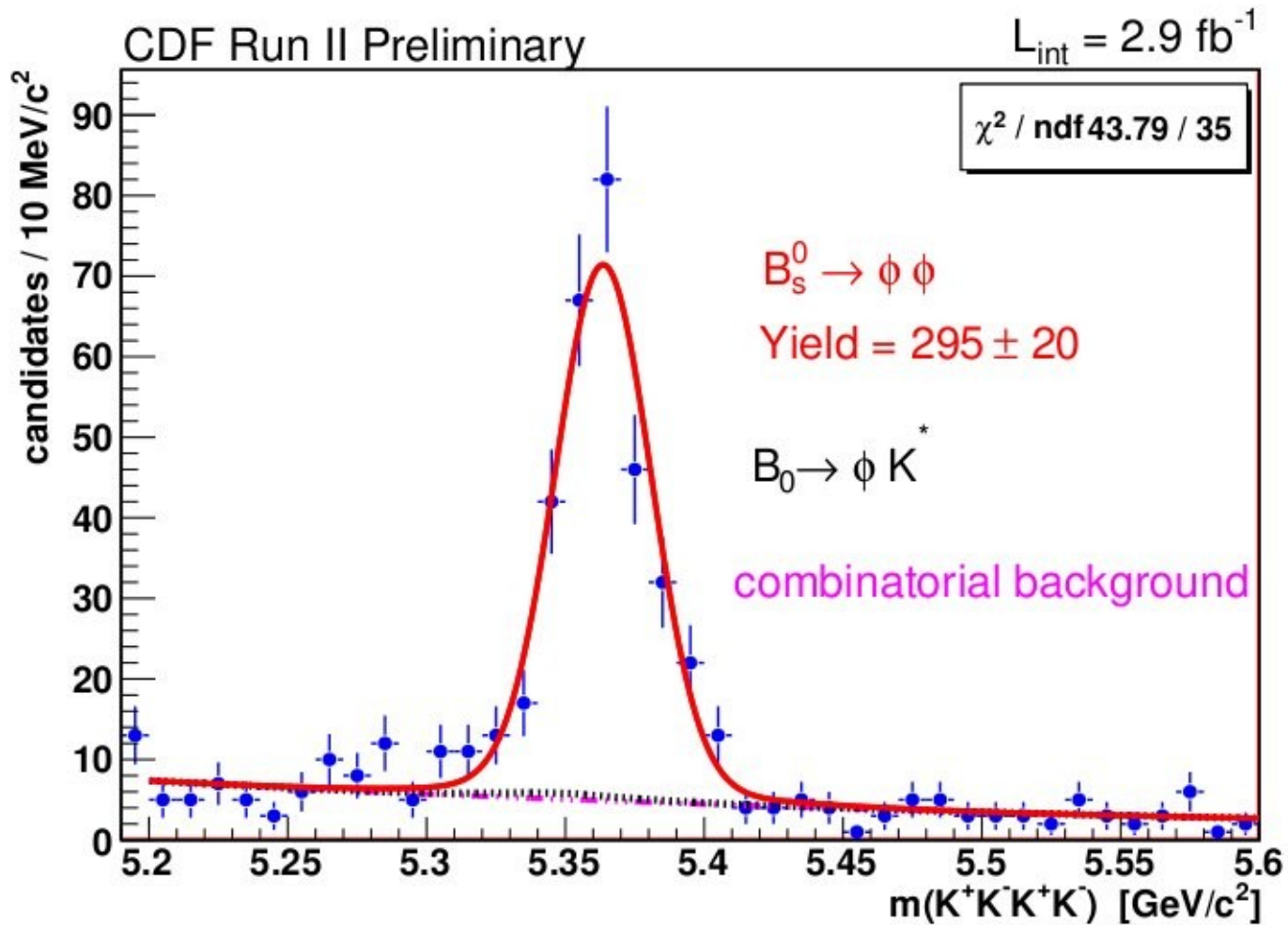
- ▶ for $B_s \rightarrow \phi\phi$:
 - ▶ $B_d \rightarrow \phi K^*$
 - ▶ $B_s \rightarrow \bar{K}^* K^*$

The reflections were evaluated using Monte Carlo sample.

The $B_s \rightarrow \bar{K}^* K^*$ reflection was negligible.

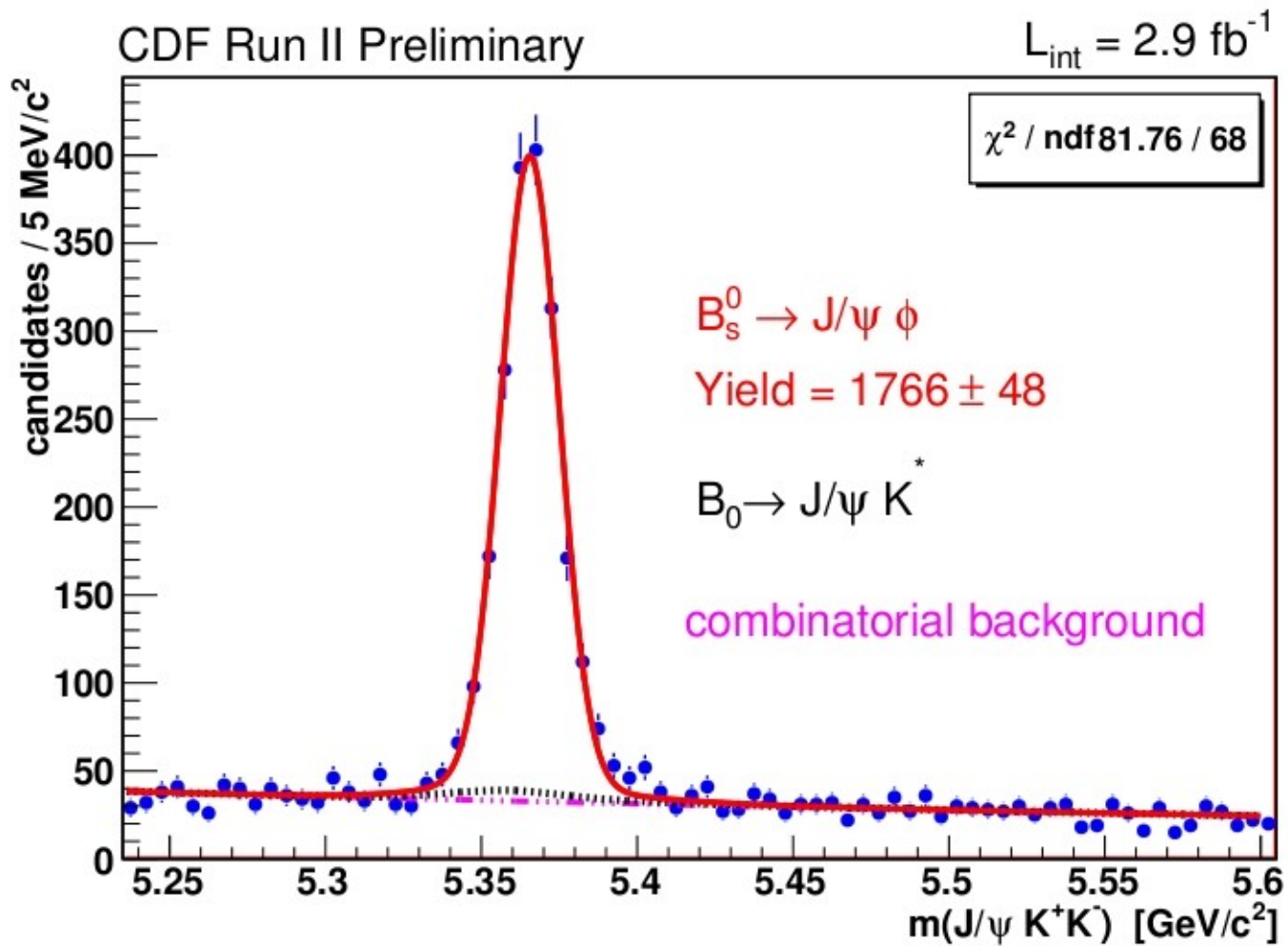


Selected events





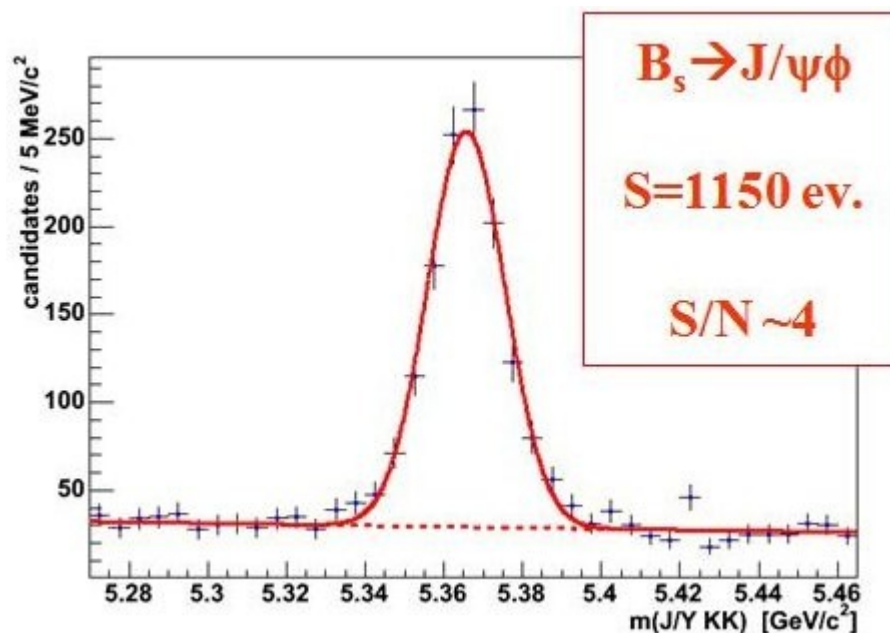
Selected events





Selected events

Adding more trigger restriction:

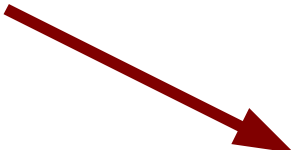


These events are only in the trigger sample we used in this analysis (**TwoTrackTrigger**), not in the samples used for other **J/psi phi** measurements (**DI_MUON Trigger**)



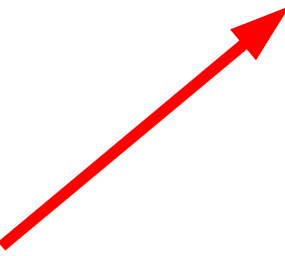
Efficiencies

Evaluated using MC



$$\frac{BR(B_s \mapsto \phi\phi)}{BR(B_s \mapsto J/\psi \phi)} = \frac{N_{\phi\phi}}{N_{J/\psi\phi}} \cdot \frac{\epsilon_{(B_s \mapsto J/\psi\phi)}}{\epsilon_{(B_s \mapsto \phi\phi)}} \cdot \frac{BR(J/\psi \mapsto \mu\mu)}{BR(\phi \mapsto KK)} \cdot \epsilon_{\mu}$$

Evaluated on data





Systematics Evaluation

We evaluated systematic uncertainty:

- ▶ due to different parametrization:
 - ▶ change in fit range \Rightarrow take into account possible structures below the mass peak due to unknown partially reconstructed
 - ▶ using one gaussian instead of two for the signal
- ▶ change in background subtraction \Rightarrow driven by $B_s \rightarrow J/\psi\phi$
- ▶ due to effect not simulated in MonteCarlo used for example:
 - ▶ the polarization



BR result

The result is:

$$\frac{BR(B_s \rightarrow \phi\phi)}{BR(B_s \rightarrow J/\psi\phi)} = [1.78 \pm 0.14^{stat} \pm 0.20^{syst}] \cdot 10^{-2}$$

$$BR(B_s^0 \rightarrow \phi\phi) = [24.0 \pm 2.1(stat) \pm 2.7(syst) \pm 8.2(BR)]10^{-6}$$



BR result

$$\text{BR}(B_s^0 \rightarrow \phi\phi) = [24.0 \pm 2.1(\text{stat}) \pm 2.7(\text{syst}) \pm 8.2(\text{BR})]10^{-6}$$

	BR[10 ⁻⁶]
Experiment	14 ₋₅ ⁺⁶ (stat.) ± 6(syst.)
QCD Factorisation	21.8 _{-1.1-17.0} ^{+1.1+30.4} 19.5 _{-1.0-8.0} ^{+1.0+13.1}
QCD Factorisation	13.1
Naive Factorisation	9.05



Conclusion and perspectives

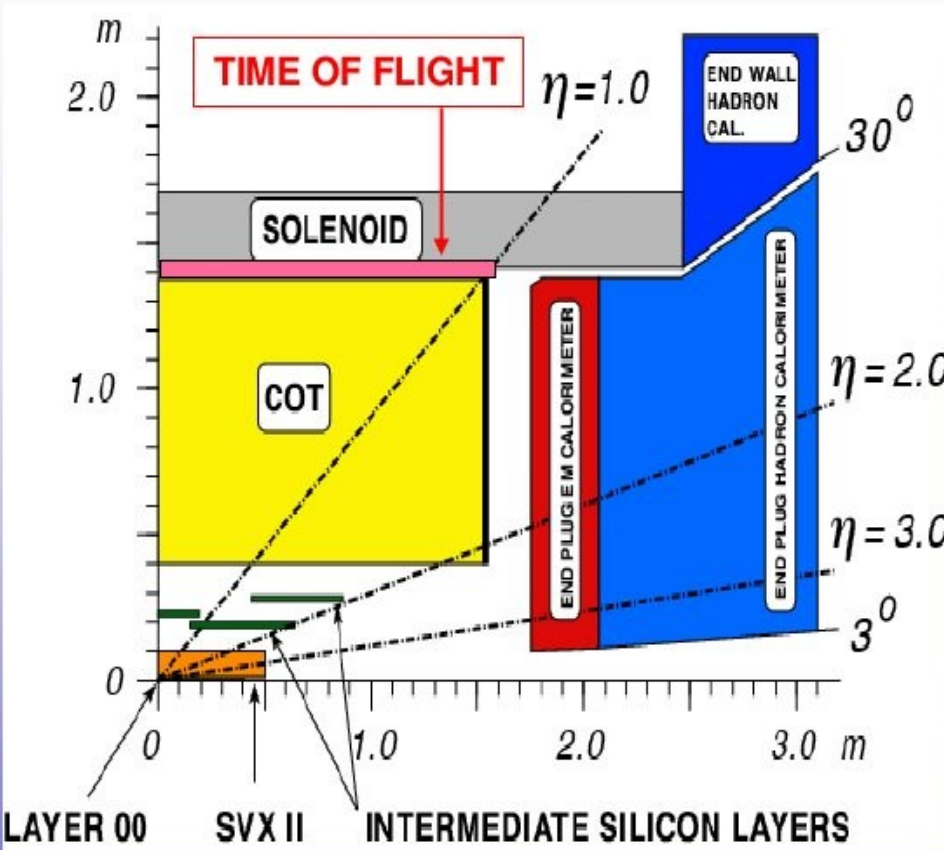
- 1) The final result is in agreement with SM prediction.
- 2) Error is dominated by J_{ψ}/Φ Branching ratio error
- 3) The sample selected allow to polarization studies



Backup slides



Tracking system



TOF: 100ps resolution, 2 sigma K/π separation for tracks below 1.6 GeV/c (significant improvement of B_s flavor tag effectiveness)

COT: large radius (1.4 m) Drift C.

- 96 layers, 200ns drift time
- Precise P_T above 400 MeV/c
- Precise 3D tracking in $|\eta| < 1$

$\sigma(1/P_T) \sim 0.1\% \text{GeV}^{-1}$; $\sigma(\text{hit}) \sim 150 \mu\text{m}$

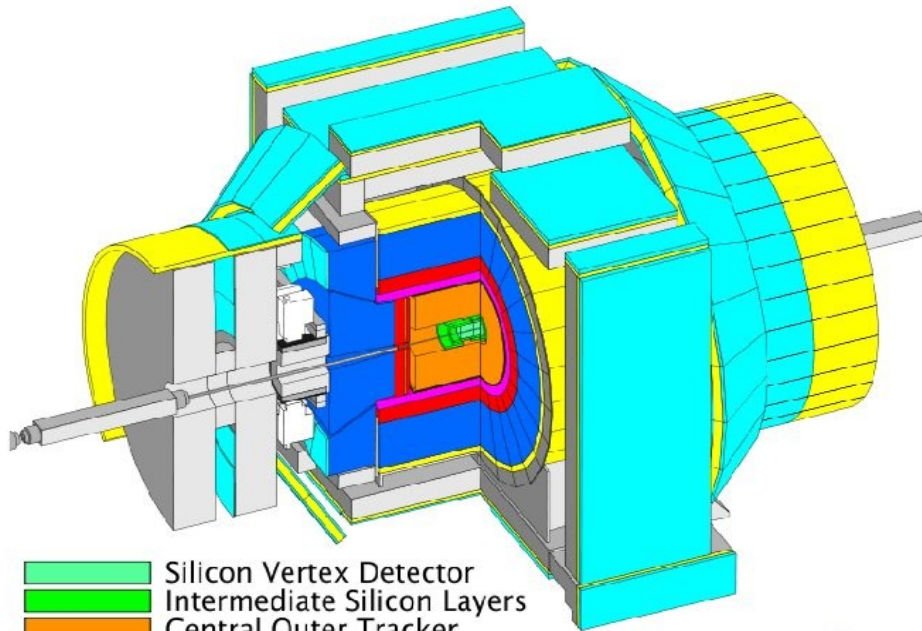
- dE/dx info provides >1.3 sigma K/π separation above 2 GeV

SVX-II + ISL: 6 (7) layers of double-side silicon ($3\text{cm} < R < 30\text{cm}$)

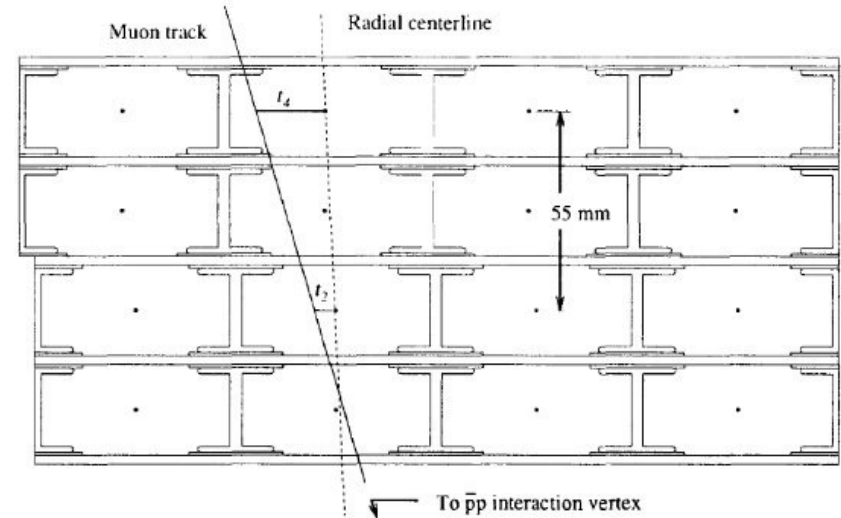
- Standalone 3D tracking up to $|\eta| = 2$
- Very good I.P. resolution: $\sim 30 \mu\text{m}$ ($\sim 20 \mu\text{m}$ with Layer 00)



Muons detectors



- Silicon Vertex Detector
- Intermediate Silicon Layers
- Central Outer Tracker
- Time Of Flight
- 1.4 T Superconducting Solenoid
- EM Calorimeter
- Hadron Calorimeter
- Muon Counters/Chambers

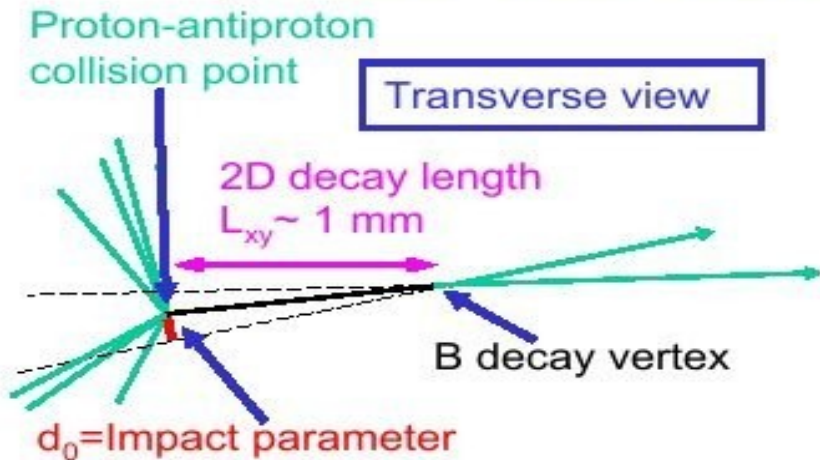




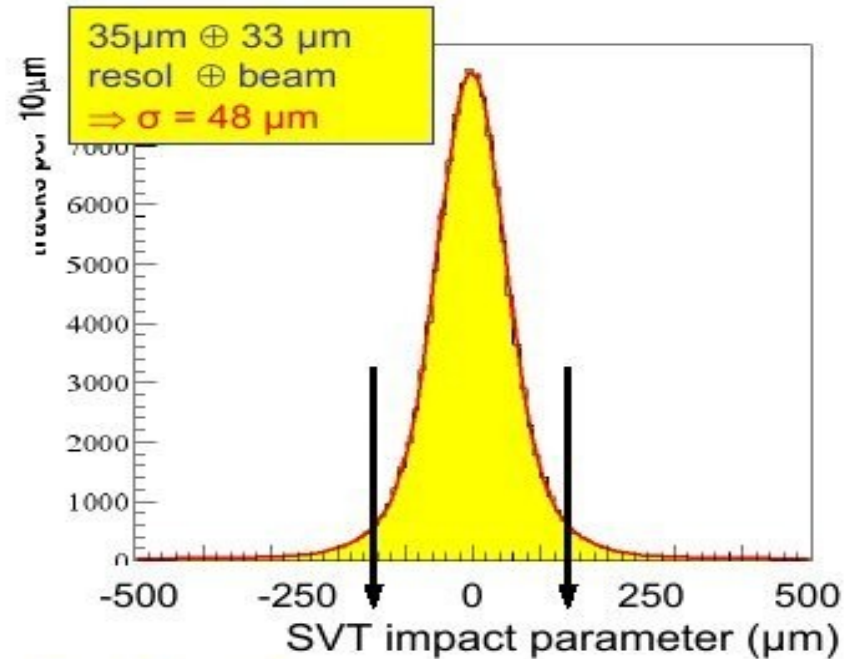
Relevant detector elements for this analysis:



CDF Level 2 Silicon Vertex Trigger



Exploit long b, c lifetimes in Trigger!
 L1 track + Si hits = Impact parameter @L2
 A first at a hadron collider!
 CDF is a charm/ B Factory!



Lepton (e, μ) + displaced track trigger

Lepton: $p_T > 4 \text{ GeV}$

Track: $p_T > 2 \text{ GeV}$, $d_0 > 120 \mu\text{m}$

Semi-leptonic B decays ($B \rightarrow \ell \nu X$)

Displaced two track trigger

Tracks: $p_T > 2 \text{ GeV}$, $d_0 > 120 \mu\text{m}$

$\Sigma p_T > 5.5 \text{ GeV}$

Fully hadronic B decays ($B \rightarrow hh'$, $B_s \rightarrow D_s \pi$, $D \rightarrow K \pi \dots$)

systematics

	$B_s^0 \rightarrow \phi\phi$	$B_s^0 \rightarrow J/\psi\phi$
	$\Delta N_{\phi\phi}/N_{\phi\phi}$	$\Delta N_{J/\psi\phi}/N_{J/\psi\phi}$
fit range	3%	-
signal parametrization	3%	2%
background subtraction: error on BRs	1%	1%
	$\Delta \varepsilon_{\phi\phi}/\varepsilon_{\phi\phi}$	$\Delta \varepsilon_{J/\psi\phi}/\varepsilon_{J/\psi\phi}$
polarization in MC	7%	6%
	$\Delta \varepsilon_{\phi\phi}/\varepsilon_{J/\psi\phi}$	
XFT particle dep.	4%	
p_T reweight	0.9%	
	$\Delta \varepsilon_{\mu}/\varepsilon_{\mu}$	
η parametrization & correlation	0.9%	

Table 16: *Contributions to the total relative uncertainty from the systematic uncertainty sources considered.*