

Theoretical Review of Ultra-Relativistic Heavy Ion Physics

Larry McLerran BNL and Riken-BNL Center

Topics:

Color Glass Condensate

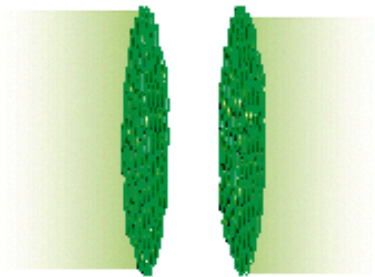
Glasma

Quark Gluon Plasma

Quarkyonic Matter

Topological Charge Fluctuations

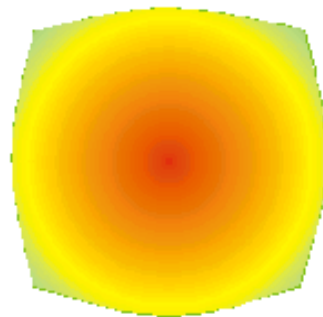
Art due to S. Bass



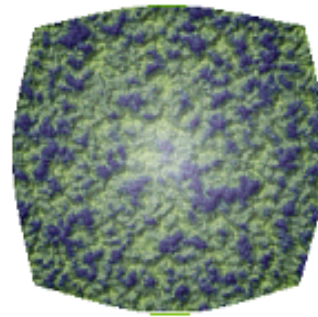
CGC



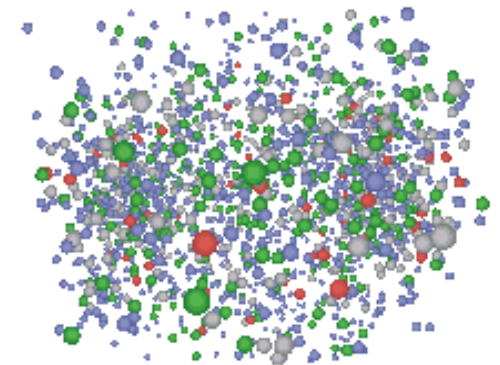
Initial
Singularity



Glasma



sQGP

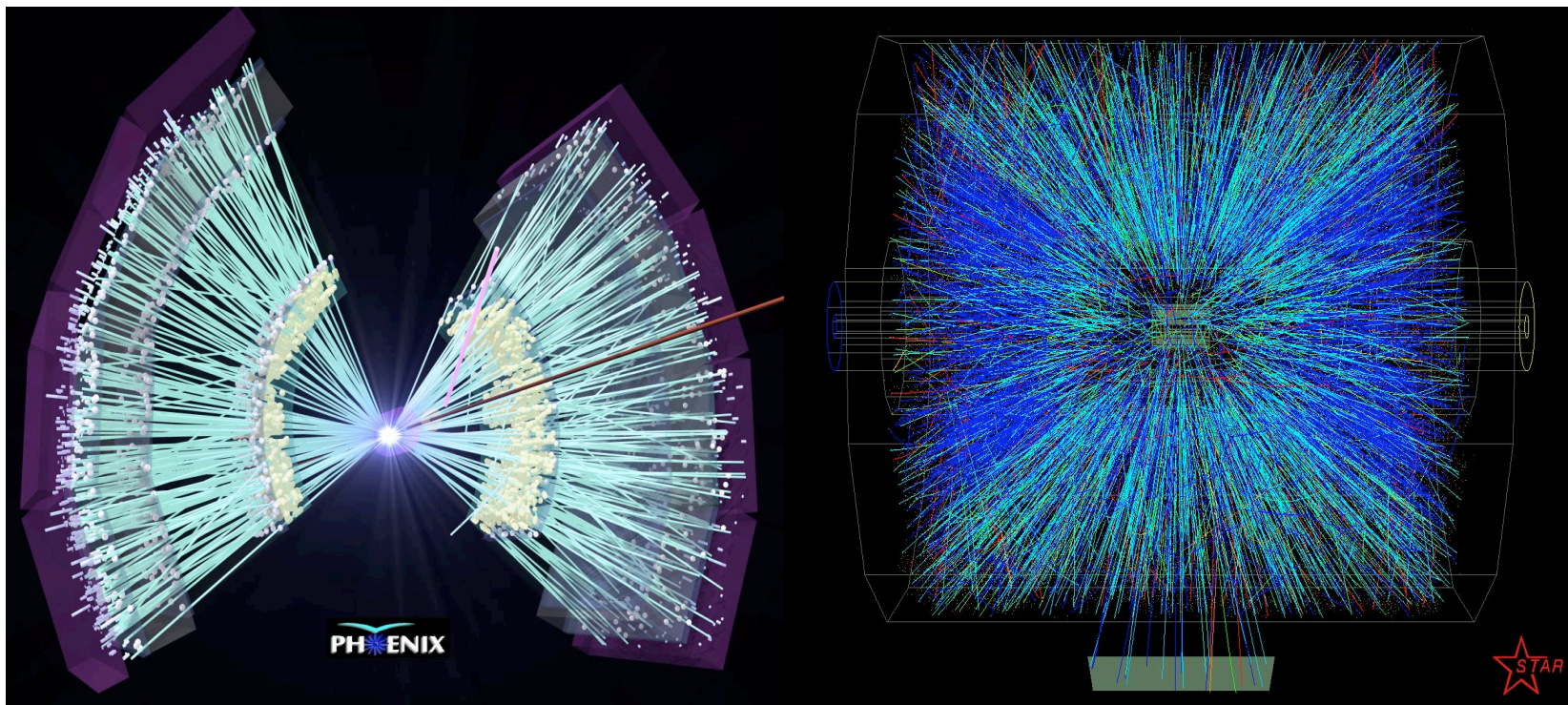
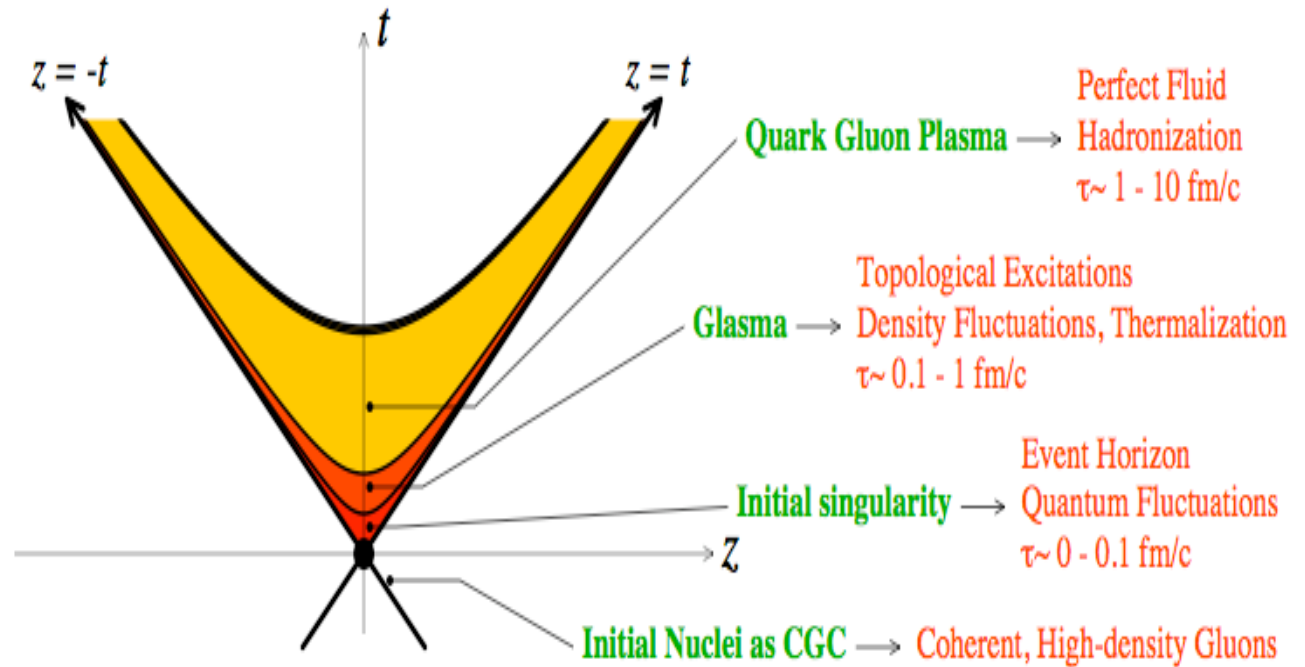


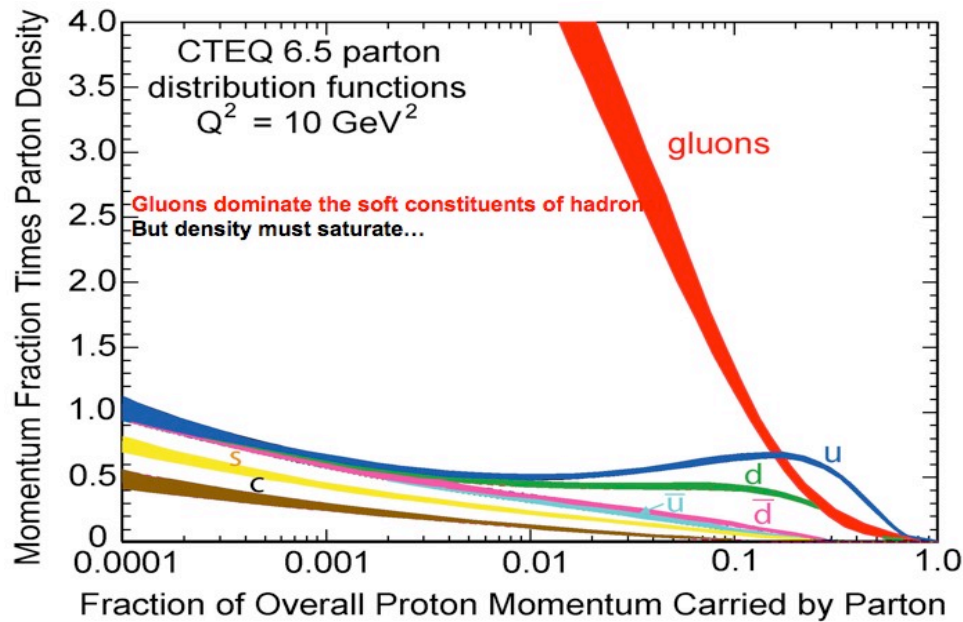
Hadron Gas

What is the high energy limit of QCD?

What are the possible form of high energy density matter?

How do quarks and gluons originate in strongly interacting particles?





Cross sections for hadrons rise very slowly with energy

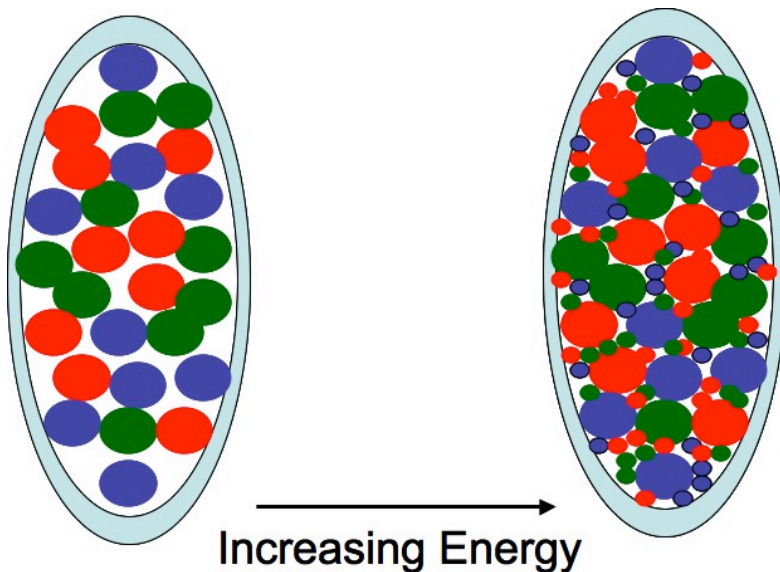
$$\sigma_{tot} \sim \ln^2(E/\Lambda_{QCD})$$

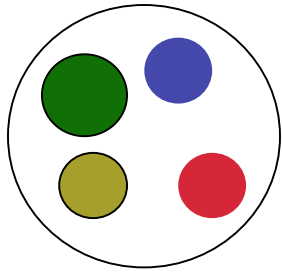
$$\Lambda_{QCD} \sim 200 \text{ MeV}$$

But the gluon density rises much more rapidly!

The high energy limit is the high gluon density limit.

Surely the density must saturate for fixed sizes of gluons at high energy.





Baryon:

3 quarks

3 quarks 1
gluon

.....

3 quarks and
lots of gluons



Color Glass Condensate

Color: Gluons

Glass:

The partons which make
the CGC fields are moving
fast => Lorentz time
dilation => fields evolve
slowly compared to natural
times scales

Condensate:

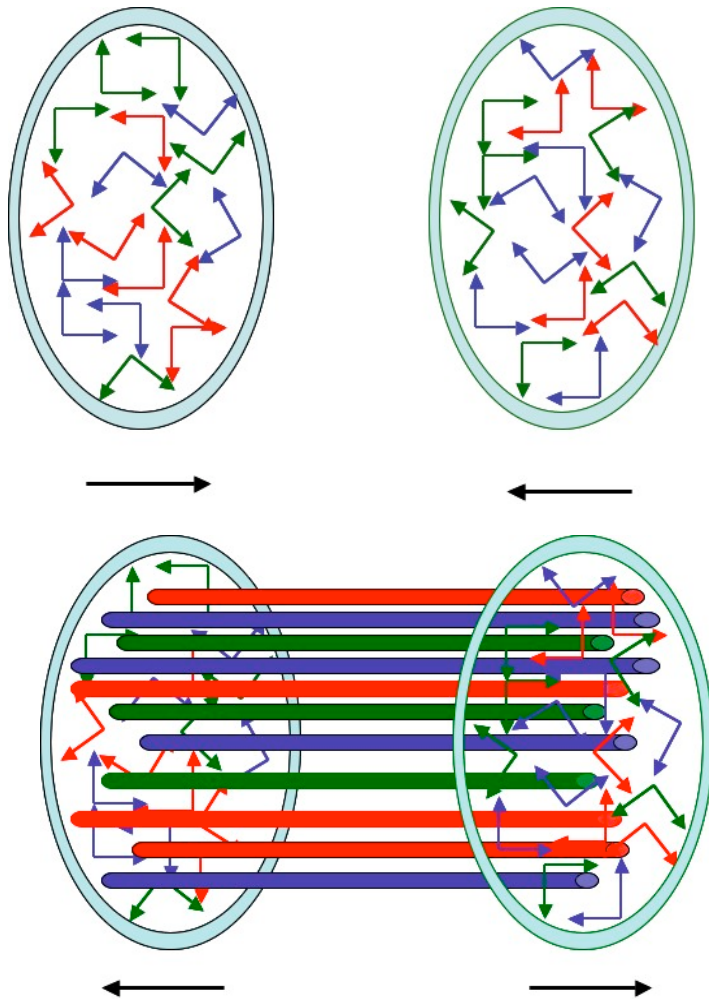
$$\frac{dN}{dyd^2p_Td^2x_T} = \rho \quad \text{Phase space density}$$

$$E = -\kappa\rho + \kappa'\alpha_S\rho^2 \Rightarrow \rho \sim 1/\alpha_S$$

Coupling weak because density is high

Fields are coherent and classical

CGC Gives Initial Conditions for QGP in Heavy Ion Collisions

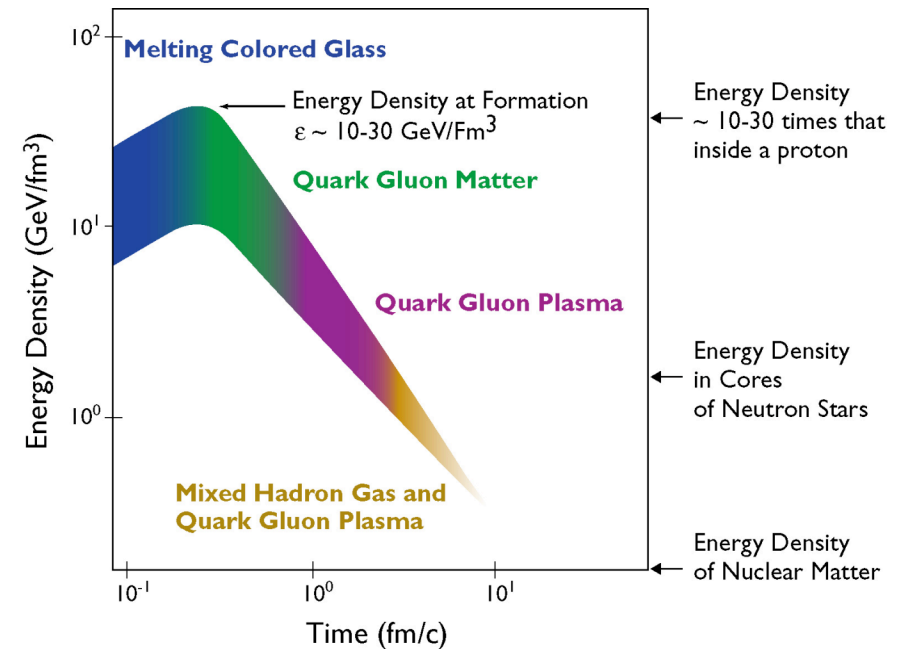


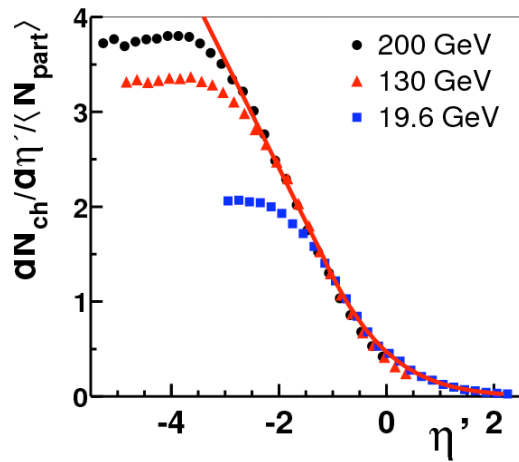
“Instantaneously” develop longitudinal color E and B fields

Two sheets of colored glass collide

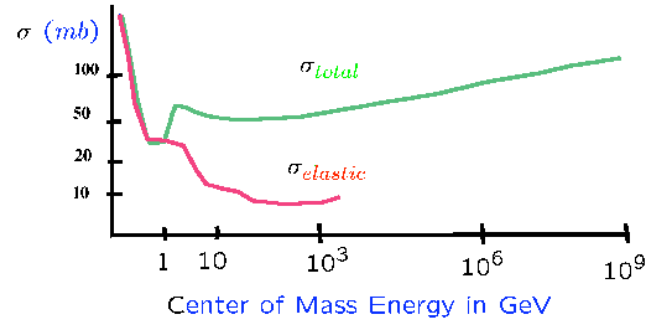
Glass melts into gluons and thermalize

QGP is made which expands into a mixed phase of QGP and hadrons





The total hadronic cross section:



CGC Provides a Successful Phenomenology of:

Limiting Fragmentation

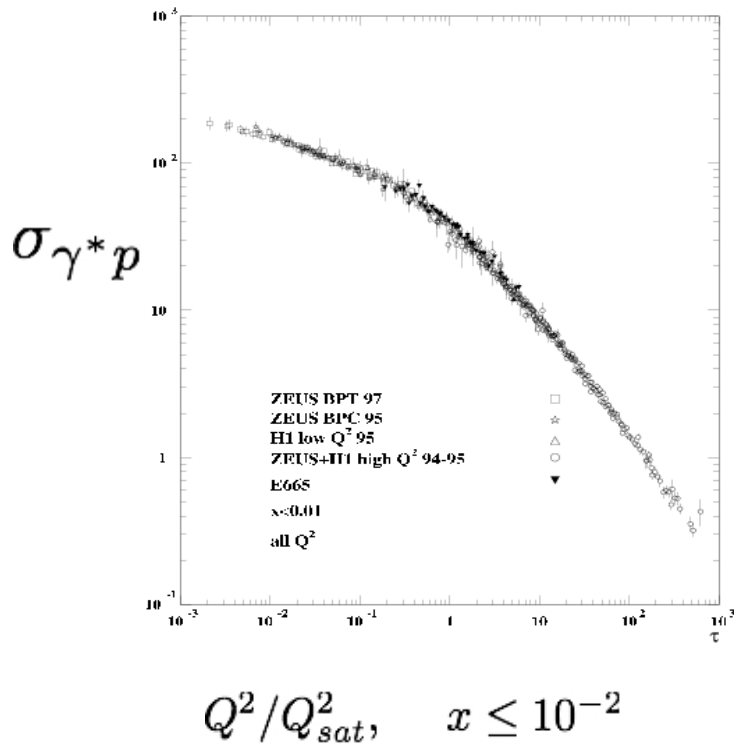
Growth and Scaling Properties of Gluon Distributions

Diffractive DIS and DIS Structure Functions

Multiplicities in Heavy Ion Collisions

Shadowing

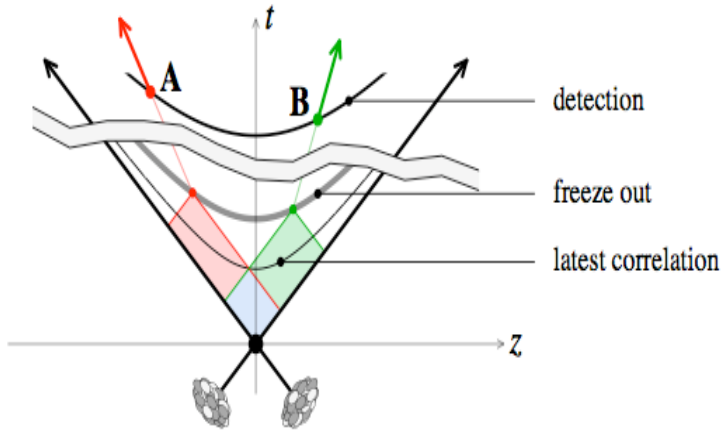
Sophisticated theoretical formalism from first principles in QCD



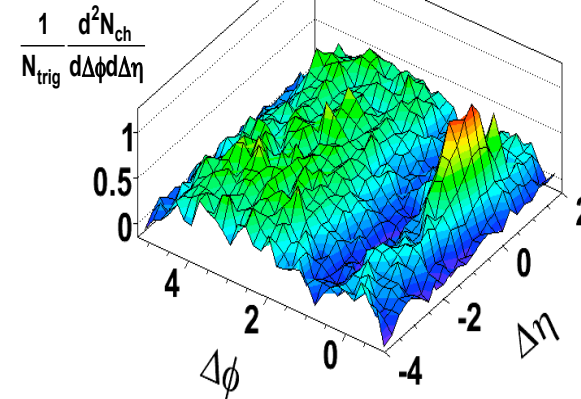
The Glasma, Long Range Correlations and the Ridge

Causality dictates:

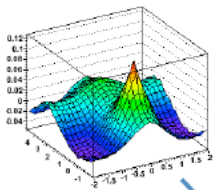
$$\tau \leq \tau_{\text{freeze-out}} \exp\left(-\frac{1}{2}|y_A - y_B|\right)$$



Au+Au 200 GeV, 0 - 30% PHOBOS preliminary



proton-proton

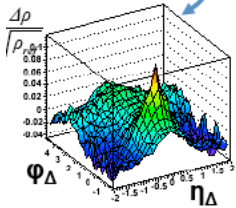


200 GeV Data

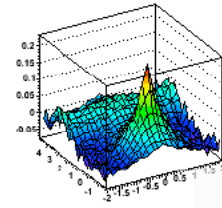
Analyzed 1.2M minbias 200 GeV Au+Au events, and 13M 62 GeV minbias events (not shown) Included all tracks with $p_T > 0.15$ GeV/c, $|\eta| < 1$, full ϕ

note: 38-46% not shown

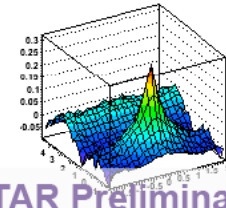
84-93%



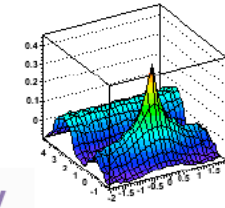
75-84%



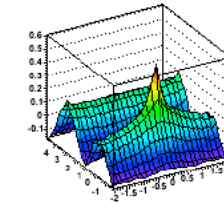
65-75%



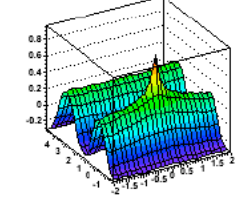
55-65%



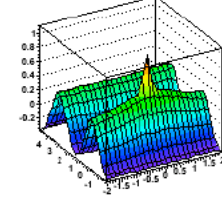
46-55%



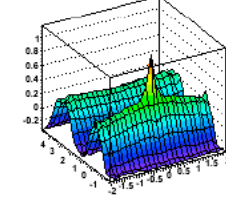
28-38%



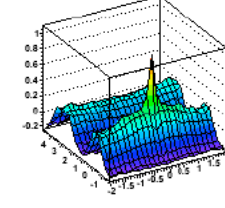
19-28%



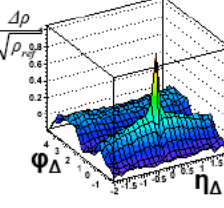
9-19%



5-9%



0-5%



STAR Preliminary

Shuryak

Dumitru, Gavin,
Gelis, Venugopalan
Moschelli

“Glittering” Glasmas

n-particle correlation can be expressed as

Gelis,Lappi,McLerran

$$\left\langle \frac{d^n N}{dy_1 d^2 \mathbf{p}_{\perp 1} \cdots dy_n d^2 \mathbf{p}_{\perp n}} \right\rangle = \frac{(n-1)!}{k^{n-1}} \left\langle \frac{dN}{dy_1 d^2 \mathbf{p}_{\perp 1}} \right\rangle \cdots \left\langle \frac{dN}{dy_n d^2 \mathbf{p}_{\perp n}} \right\rangle$$

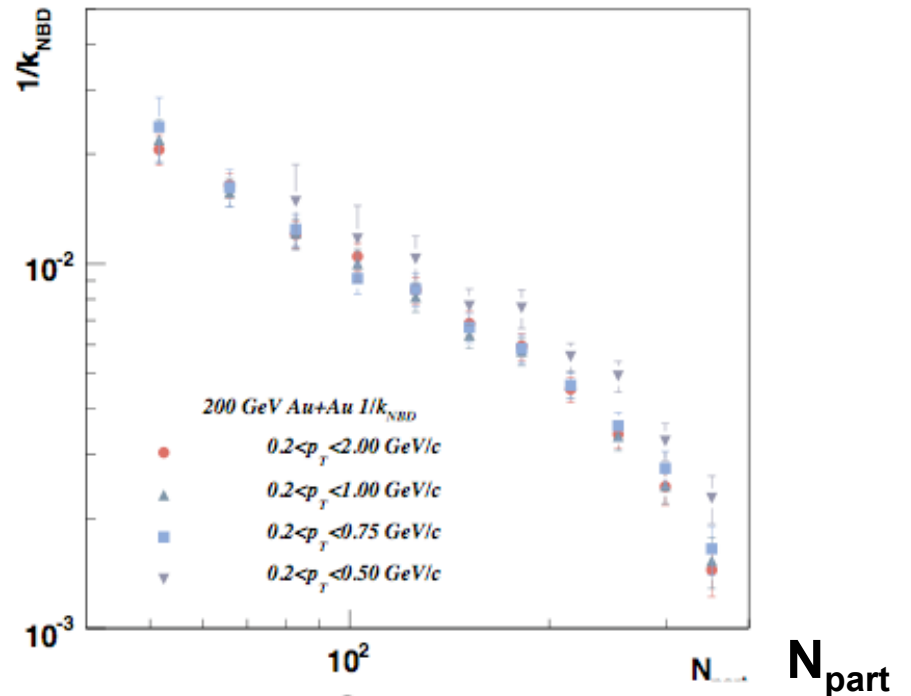
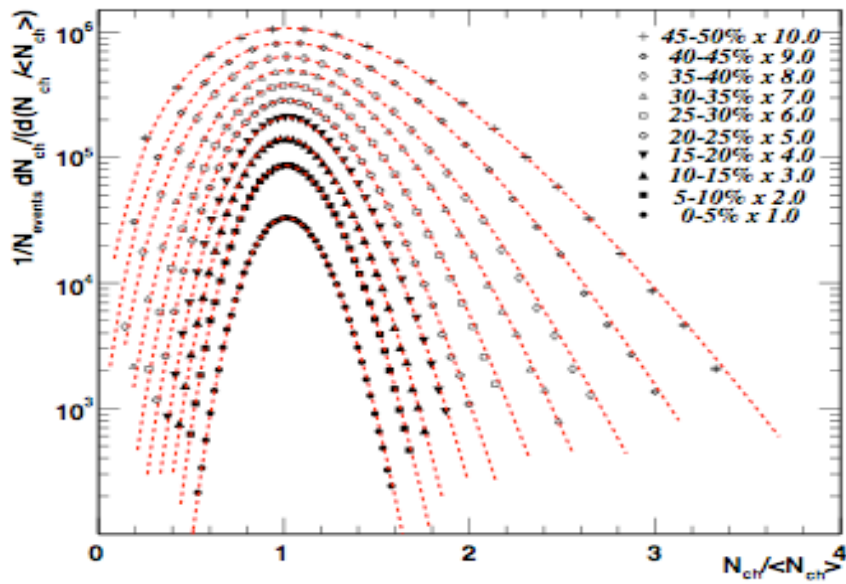
with $k = \zeta_n \frac{(N_c^2 - 1) Q_S^2 S_{\perp}}{2\pi}$

For $k = 1$, Bose-Einstein dist. For $k = \infty$, Poisson Dist.

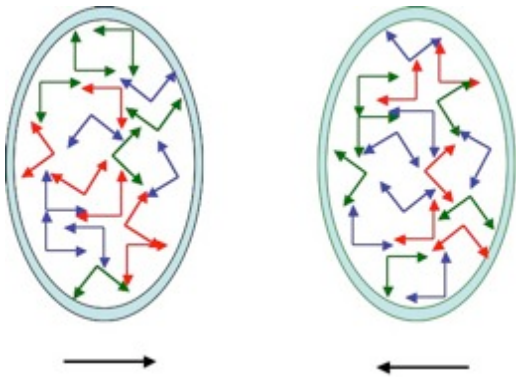
This is a negative binomial distribution which is known to describe well multiplicity distributions in hadronic and nuclear collisions

PHENIX
arXiv:0805.1521

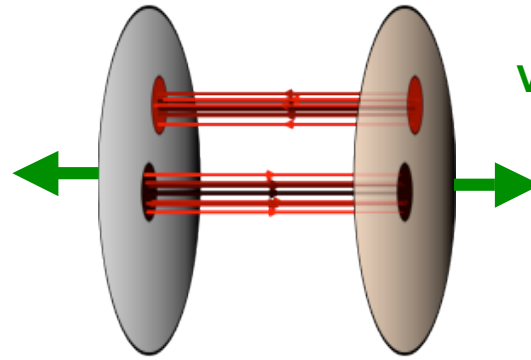
$$\sigma^2 = \bar{n} + \frac{\bar{n}^2}{k}$$



Glasma Flux Tubes and Chern-Simons Number



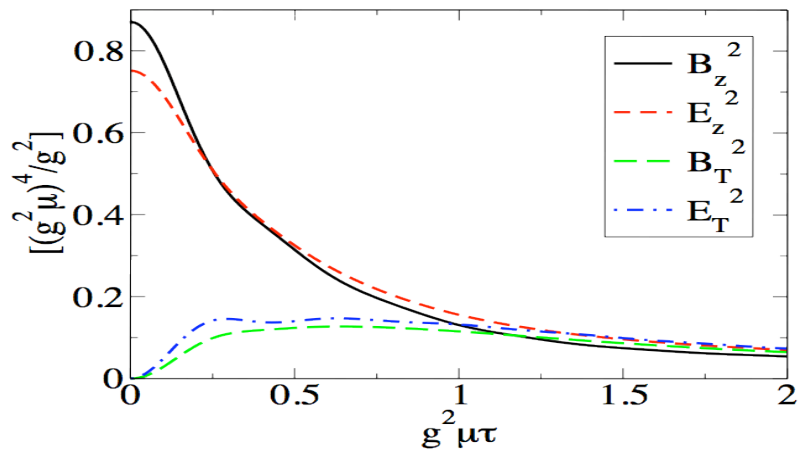
Before: transverse E & B
"Weizsacker-Williams fields"



After: **boost invariant** flux tubes of size $1/Q_s$

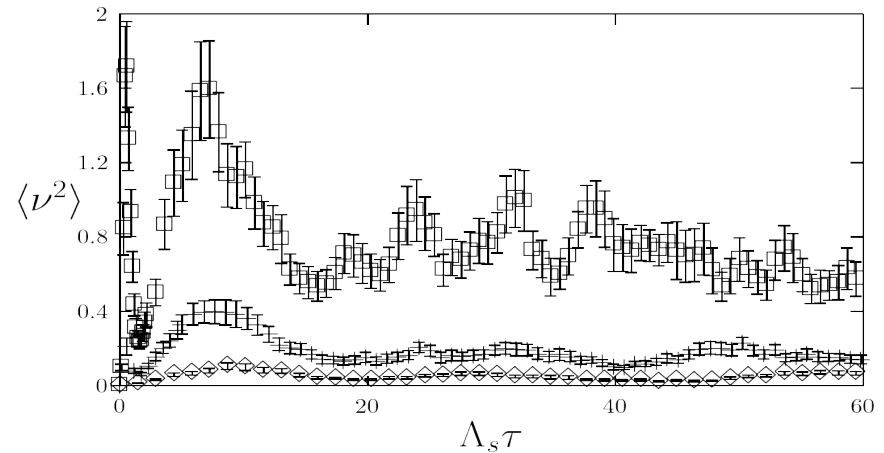
Krasnitz, Nara,
Venugopalan; Lappi

Parallel color E & B fields



Lappi, McLerran

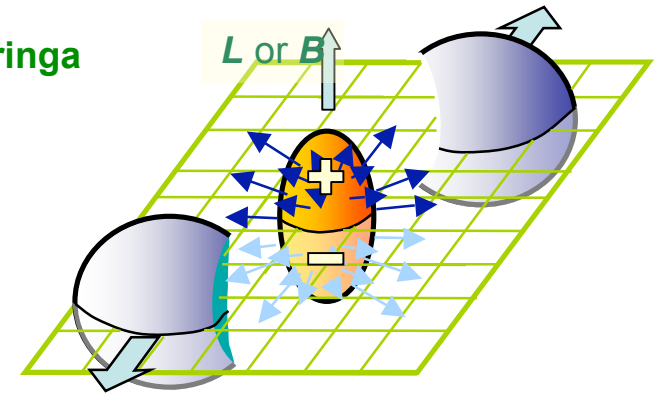
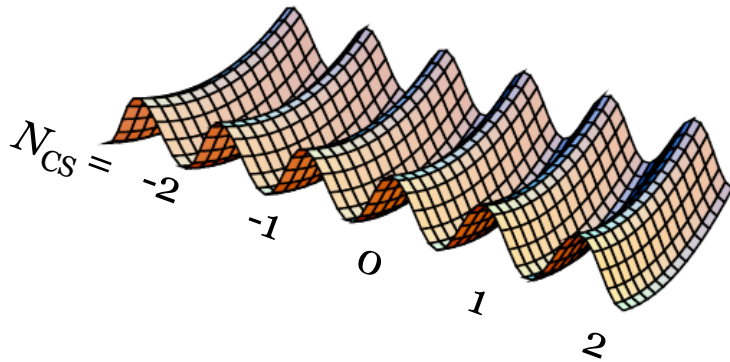
Chern-Simons topological charge



Kharzeev, Krasnitz, Venugopalan,

P and CP violation: Chiral Magnetic Effect

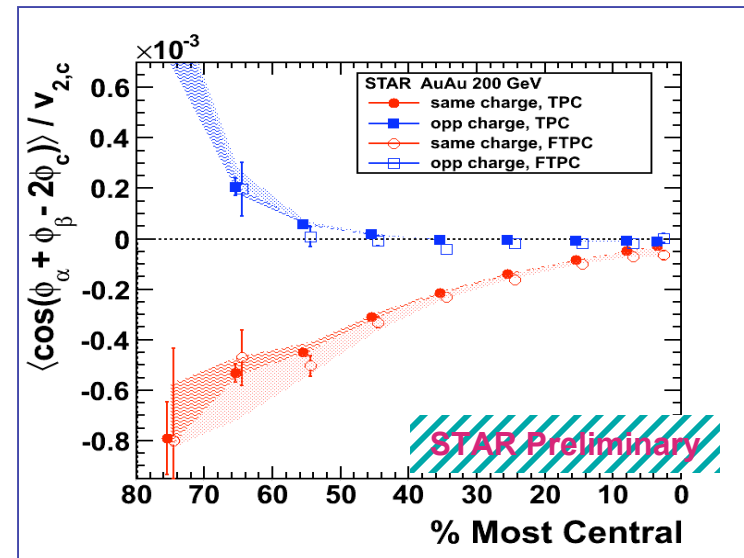
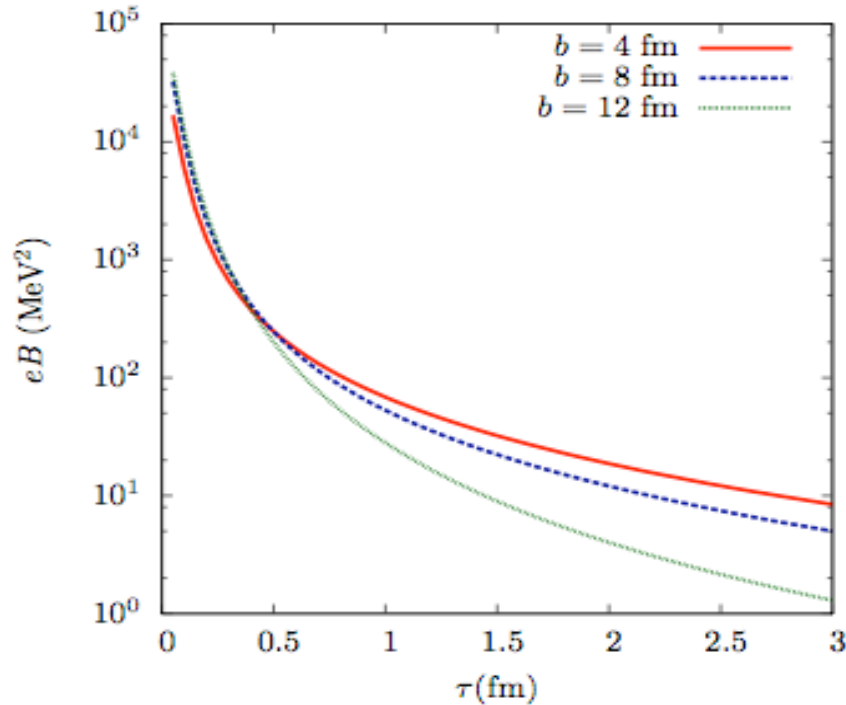
Kharzeev, McLerran, Warringa



+

Topological fluctuations
-sphaleron transitions in Glasma

External (QED) magnetic field



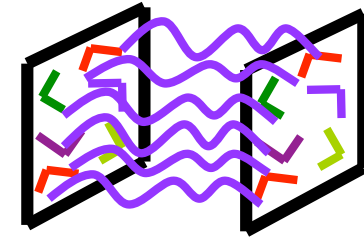
Possible experimental signal of
charge separation (Voloshin)

The Unstable Glasma

Romatschke, Venugopalan

- Small rapidity dependent **quantum fluctuations** of the LO Yang-Mills fields grow rapidly as

$$\sim e\sqrt{Q_s\tau}$$

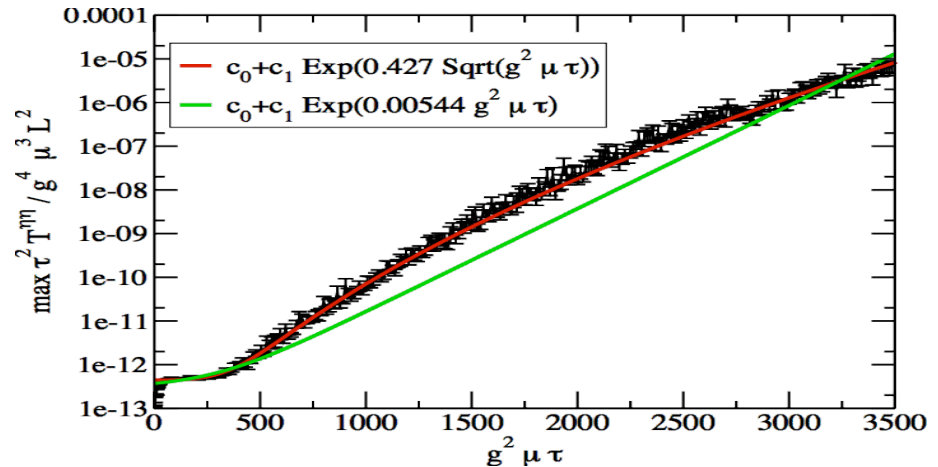


- E_{\perp} and B_{\perp} fields as large as E_{\parallel} and B_{\parallel} at time

$$\tau \sim \frac{1}{Q_s} \ln^2 \left(\frac{1}{\alpha_s} \right)$$

Possible mechanism for **rapid isotropization**

Problem: collisions can't 'catch up'



- Turbulent “thermalization” may lead to “anomalously” low viscosities

Asakawa, Bass, Muller; Dumitru, Nara, Schenke, Strickland

- Significant energy loss in Glasma because of synchrotron like radiation?

Shuryak, Zahed; Zakharov; Kharzeev

Thermalization:

Growth of Instabilities and Turbulance?

Strong Coupling Dynamics?

System thermalizes and makes a very good fluid: Hydrodynamic simulations of flow induced by the collisions, and of particle distributions with transverse momenta $< 1\text{-}2$ GeV work very well. Estimates of viscosity give very small values

$$\eta/S > 1/4\pi$$

$$\lambda T \sim 1/2\pi$$

Klebanov

Strongly interacting QGP paradigm.

AdSCFT Modeling for qualitative understanding of the strong coupling limit

AdSCFT has allowed insight about the form of viscous corrections to relativistic hydrodynamics that allow for a consistent theoretical formulation and numerical implementation

Danielowicz, Gyulassy
Son, Starinets, Gubser

Romatschke and Romatschke,
Muronga, Rischke

AdSCFT and Jet Quenching:

Theoretical problem: No jets in strong coupling N=4:

Electron positron annihilation

Deep inelastic Jet production

Cascade to low momentum modes occurs very quickly

A single hadron at high energy in strong coupling N=4:

Distribution of partons $\frac{dN}{dx} \sim \frac{1}{x^2}$

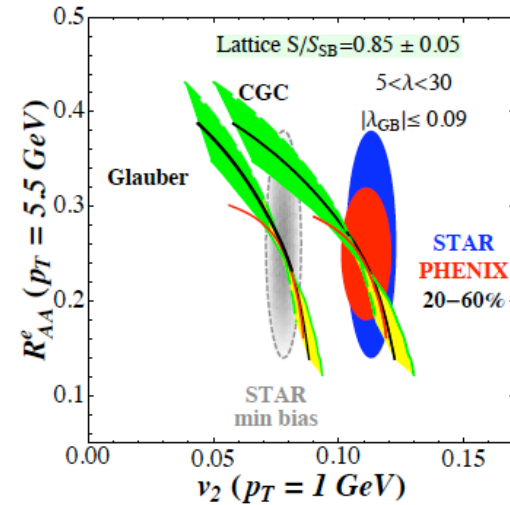
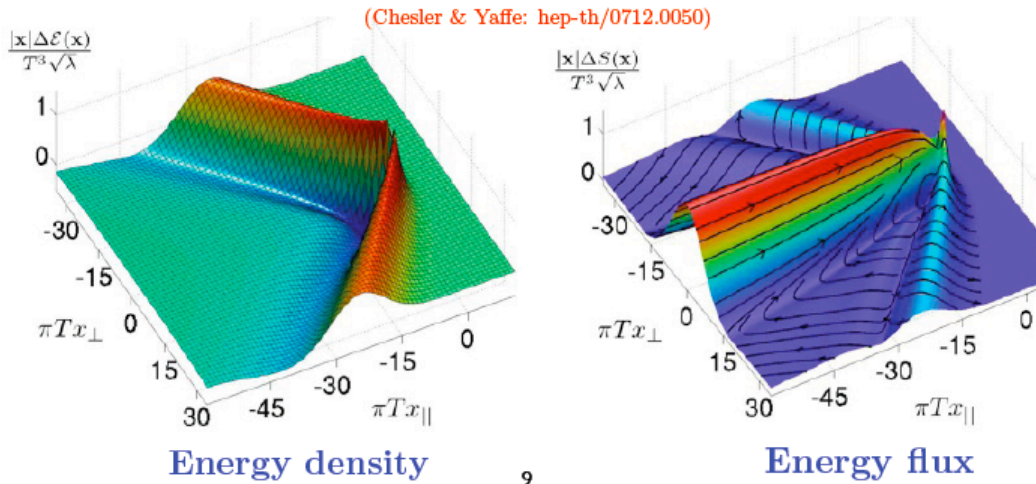
Saturated gluon density with phase
space occupation ~ 1 , not $\frac{1}{\alpha}$

Leading twist contribution to deep
inelastic scattering vanishes

Gylassy, Vitev
Rajagopal, Wiedmann; Yaffe; Mueller, Iancu and Hatta;
Maldacena; Kovchegov

Nevertheless:

Data on jet energy loss and heavy quark flow and energy loss is so stunning:

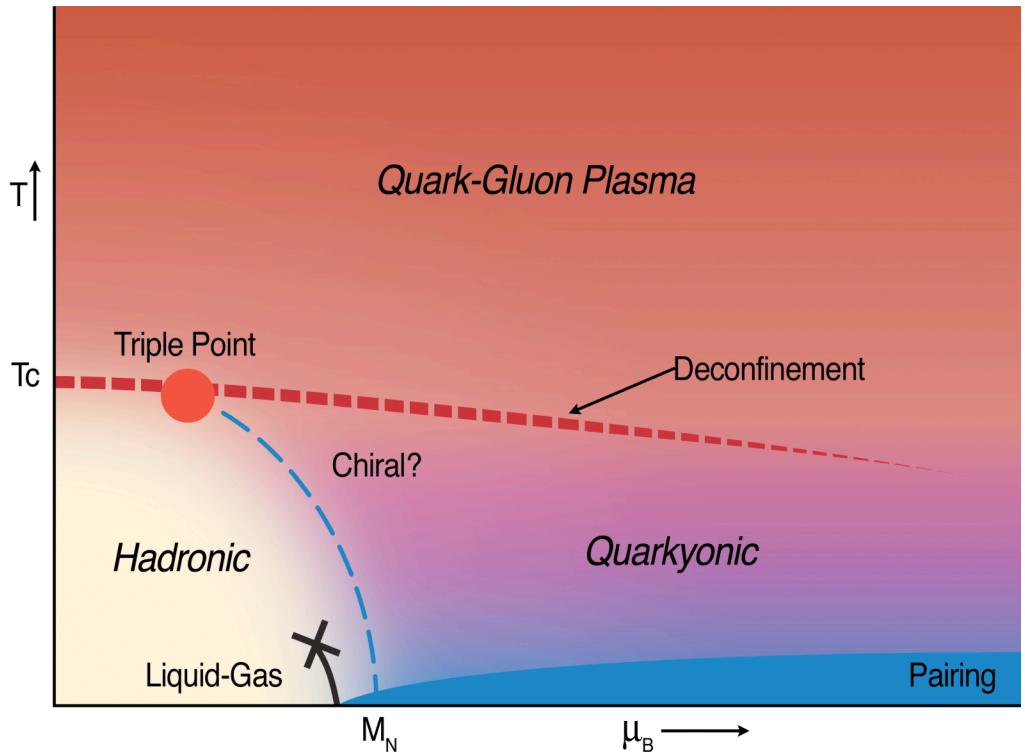


Heavy quark energy loss computations, Mach cones

Heavy quark energy loss vs v_2

Chasler and Yaffe;

Betz, Gyulassy, Noronha, Torrieri, Rischke



Hadronic: mesons, glueballs and no baryons

QGP quark and gluons

Quarkyonic: net quarks, mesons and glueballs

Large N_c :

Quark Loops do not affect the confinement potential



$$g^2 \mu_Q^2 \sim \alpha_N \mu_Q^2 / N_c$$

$$e^{(\mu_B - M_B)/T} = 0 \text{ if } \mu_B < M_B$$

Baryon mass $\sim N_c$,
no baryons in low density phase