(INTRODUCTION TO) PRIMORDIAL COSMOLOGY

JEAN-PHILIPPE UZAN

Institut d'Astrophysique de Paris





UK PKUGKAM

1- Introduction:

What is the size of the universe? What is cosmology? Models and hypothesis

- 2- The standard hot big-bang model

 Friedmann equations. Dynamics. Successes and problems.
- 3- Gravitational dynamics and large scale structures

 Newtonian regime. General properties. Gravity waves.
- 4- Observing the large scale structures

 Galaxy catalogs. CMB. Lensing. Summary of the observational status.
- 5- Inflation: a scenario for the origin of structures
- 6- The input of high energy physics

 dark energy. dark matter. Extra-dimension and string inspired cosmology
- 7- Conclusion

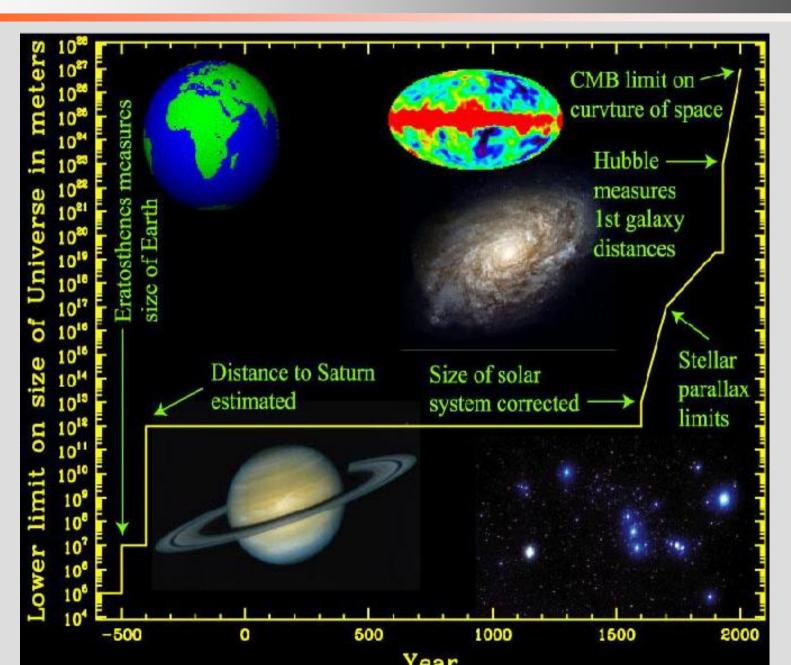
Open issues and questions.- exotica

PART 1: COSMOLOGY

Main topics

- What is the size of the universe?
- What is cosmology?
- Models and hypothesis

AUVI 12 IUF 217F OL IUF ONIAFK2F:



MAI IS COSMOLOGY?

Cosmology is a **description** of what we think is the universe It requires to define the framework of its study:

What is the universe?

It is developed within a precise **geometrical theory** (Euclid, Descartes, Riemann,...)

which serves as a test bank for **new physical theories** (Aristotle, Newton, Einstein, ...)

Observations allow to distinguish, draw new questions,... (planetary orbits, cosmic microwave background...)

AKTICULAKITIES

Cosmology is however peculiar:

- we observe only a *small* part of the universe from a *fixed point* of spacetime,
- only *one* universe is observed by us.

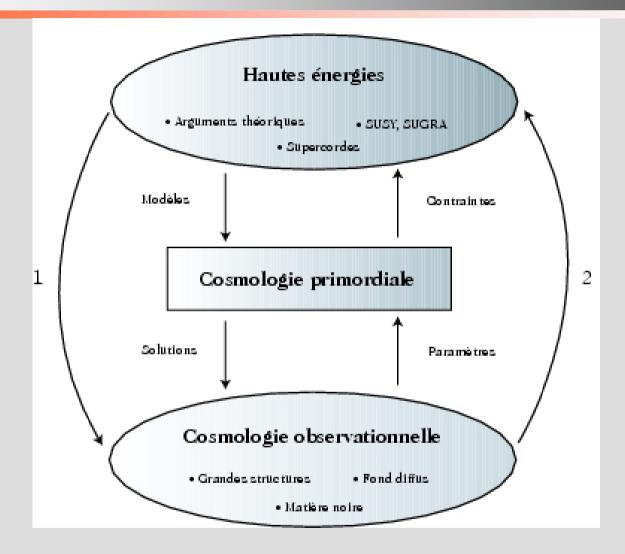
t implies that *some* hypothesis *cannot be tested*.

We start from the standard fundamental theories (as tested in laboratory and accelerator).

This theories are extrapolated to higher energies in the early energies.

- we can constrain various extension of this fundamental setting.
- it is important to diagnostic the failures of these extrapolations in order to infer the more fruitful extensions

KIMOKPIAL COSMOLOGY



High energy theories provide *models* of the primordial universe

Cosmology allows to set constraints on various extensions of the standard fundamental theories

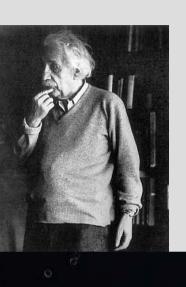
Each models must provide unambiguous observational predictions

PART II: THE STANDARD BIG-BANG MODEL

Main topics

- Friedmann-Lemaître universes
- dynamics and properties
- Successes
 - •Hubble diagram
 - •Age of the universe
 - •Relics and BBN
 - •Cosmic microwave background
- Problems

ELATIVISTIC COSMOLOGY



Physics : Einstein

Geometry: Riemann

Formulation of General Relativity (1915)

First cosmological models (1917-1927) static or in evolution?

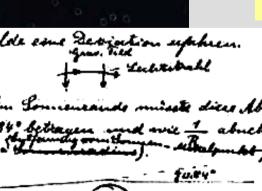


Development of physics in expanding universe (>1948)

Hot big bang model

Links with high-energy physics (>1981)

Inflation, string cosmology











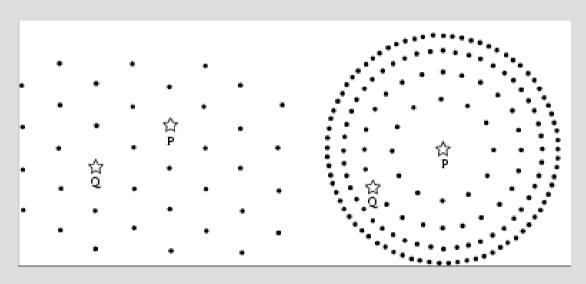
- ¤ gravitation is described by general relativity,
- ¤ physical laws are universal
- ¤ we are not seating at the particular place in the universe:
 cosmological principle Copernician principle
- **matter** contains
 - * radiation
 - * dust (fluid of galaxies without pressure)

Conservative hypothesis!

OSMOLOGICAL PRINCIPLE

Cosmological principle: the universe is homogeneous and isotropic

Copernician principle: we do not occupy a particular place in the universe

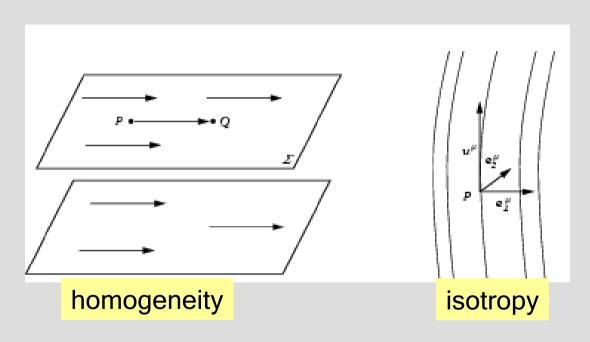


Distribution statistically isotropic Around each point P and Q are equivalent Distribution statistically isotropic around P alone P and Q are not equivalent

Copernician + isotropy → Cosmological principle

OMOGENETT AND ISOTROPT

Homogeneity: at any time, any point of space is equivalent to any other **Isotropy**: the universe is seen isotropic from any point

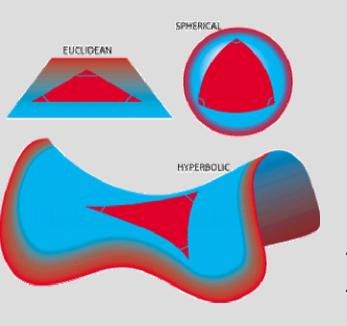


pace can be foliated by hypersurfaces Σ ich that on each hypersurface there exists isometry bringing P to Q

It exists a geodesic flow with tangent Vector u^{μ} such that at any P, there is an isometry letting P and u^{μ} invariant which makes a rotation of the directions of observation

KIEDWANN-FEWALIKE ONIVEKSE

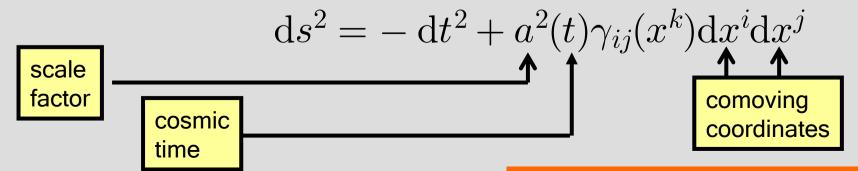
Homogeneity + isotropy → constant curvature space



There are only three 3D constant curvature spaces

$$\mathrm{d}s_3^2 = \gamma_{ij}(x^k)\mathrm{d}x^i\mathrm{d}x^j$$

The most general 4D metric satisfying these hypotheses is



Friedmann Lamaîtra enacetim

IK31 PKEDICTION: MUBBLE LAW

Space is expanding

The relative distance between 2 comoving observers

$$\mathbf{r}_{12} = a(t)(\mathbf{x}_1 - \mathbf{x}_2)$$

evolves as

$$\dot{\mathbf{r}}_{12} = H\mathbf{r}_{12} \quad H \equiv \dot{a}/a$$

This is the Hubble law (Lemaître, 1924!)

- Purely kinematic consequence of the expansion
- The function H(a) will be central to describe the dynamics

ECOND PREDICTION: REDSHIFT

From the geodesic equation for photon, one can deduce that the energy of the photons satisfy

$$\dot{E}/E = -H \Rightarrow E \propto 1/a$$

<u>Interpretation:</u> the wavelength of the photon is stretched by the expansion

redshift

$$1 + z \equiv \frac{E(t_{obs})}{E(t_{em})} = \frac{a(t_{em})}{a(t_{obs})}$$



Consequence: achromatic stretching of atomic spectra

KIEDMANN EQUATIONS

Start from the Einstein equations

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

for the Friedmann-Lemaître metric

The Cosmological Principle implies that $T_{\mu\nu} = \mathrm{diag}(-\rho, P, P, P)$

So that the Friedmann equations take the form

$$3\left(H^2 + \frac{K}{a^2}\right) = 8\pi G\rho$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

The conservation equation gives

$$\dot{\rho} + 3H(\rho + P) = 0$$

KIEDWANN EQUATIONS: MATTER DESCRIPTION

Only 2 equations are independent

To describe the matter, we need to specify an equation of state

$$P(\rho) = w\rho$$

At the cosmological level

- pressureless matter: P=0
- radiation $P=\rho/3$

3 variables (a,P,ρ) -1 relation - 2 equations = OK

The conservation equation implies

$$\rho = \rho_0(\frac{a}{a_0})^{-3(1+w)} = (1+z)^{3(1+w)}$$

KIEDWANN EQUATIONS: ADDING A COSMOLOGICAL CONSTANT

One can add a constant term to the einstein action

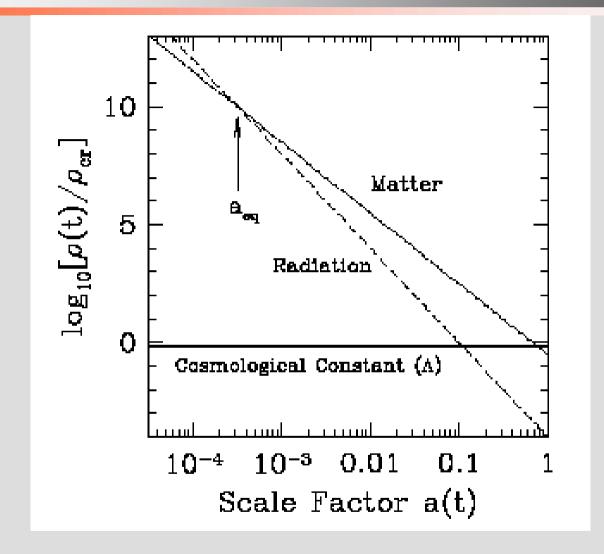
$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

So that the Friedmann equations become

$$3\left(H^2 + \frac{K}{a^2}\right) = 8\pi G\rho + \Lambda$$
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3}$$

- Interpretation of this term to be discussed later
- Equivalent to a fluid with w=-1

KIEDWANN EQUATIONS: CONSEQUENCE



The universe was dominated by radiation in the past HOT BIG BANG MODEL

KIEDMANN EQUATIONS: KEWKITING

t is common to rewrite the Friedmann equations with adimensional quantities

Density parameters

$$\Omega = \frac{8\pi G\rho}{3H^2}$$
 $\Omega_{\Lambda} = \frac{\Lambda}{3H^2}$ $\Omega_{K} = -\frac{K}{a^2H^2}$

Friedmann equation

$$E^2(z) \equiv \left(\frac{H}{H_0}\right)^2$$

$$= \sum \Omega_0 (1+z)^{3(1+w)} + \Omega_{\Lambda 0} + \Omega_{K 0} (1+z)^2$$

Constraint

$$\sum \Omega_0 + \Omega_{\Lambda 0} + \Omega_{K0} = 1$$

KIEDWANN EQUATIONS: EXEKCISE

Find the evolution of the scale factor when K=0 and w≠-1

$$\rho \propto a^{-3(1+w)} \longrightarrow H^2 \propto a^{-3(1+w)} \longrightarrow \dot{a}/a \propto a^{-3(1+w)/2}$$

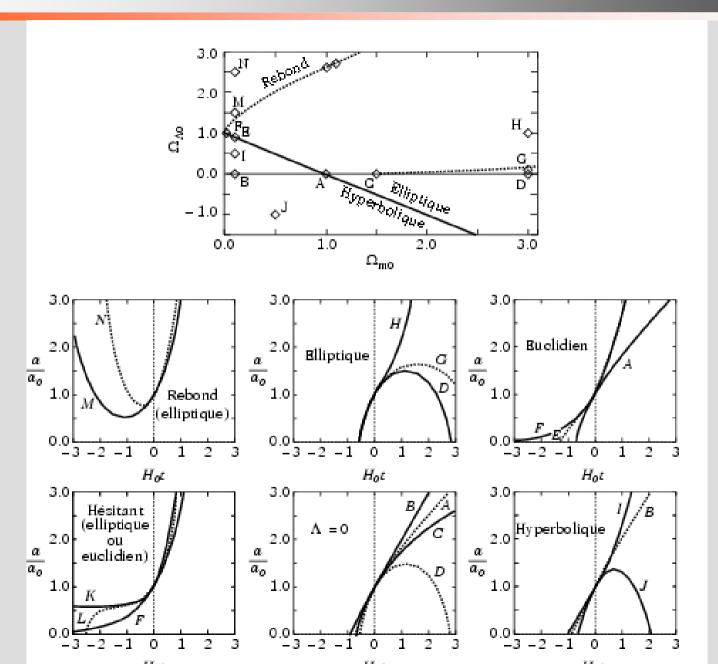
$$a(t) \propto t^{2/3(1+w)}$$

Find the evolution of the scale factor when K=0 and w=-1

$$ho \propto a^0 \longrightarrow H^2 \propto a^0 \longrightarrow \dot{a} \propto Ha$$

$$a(t) \propto \exp Ht$$

KIEDWANN EQUATIONS: DINAMICS



IWE AND DISTANCE SCALES

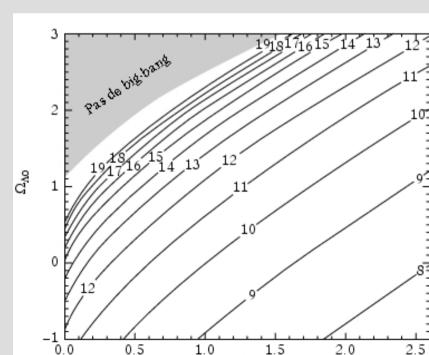
H₀ sets a typical time and distance scales for the universe

$$H_0 = 100 h \text{ km/s/Mpc}$$

 $D_{H_0} = c/H_0 = 3000 h^{-1} \text{Mpc} = 2.9 h^{-1} \times 10^{25} \text{ m}$
 $t_{H_0} = 1/H_0 = 9.78 h^{-1} \times 10^9 \text{yr}$

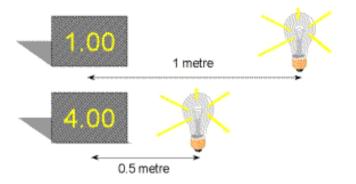
Dynamical age of the universe

$$\mathrm{d}t = \frac{\mathrm{d}a}{\dot{a}} = t_{H_0} \frac{\mathrm{d}a}{aE(a)}$$



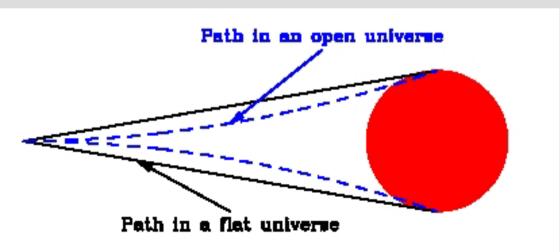
12 I ANCES

Measuring Distances with Standard Light Bulbs



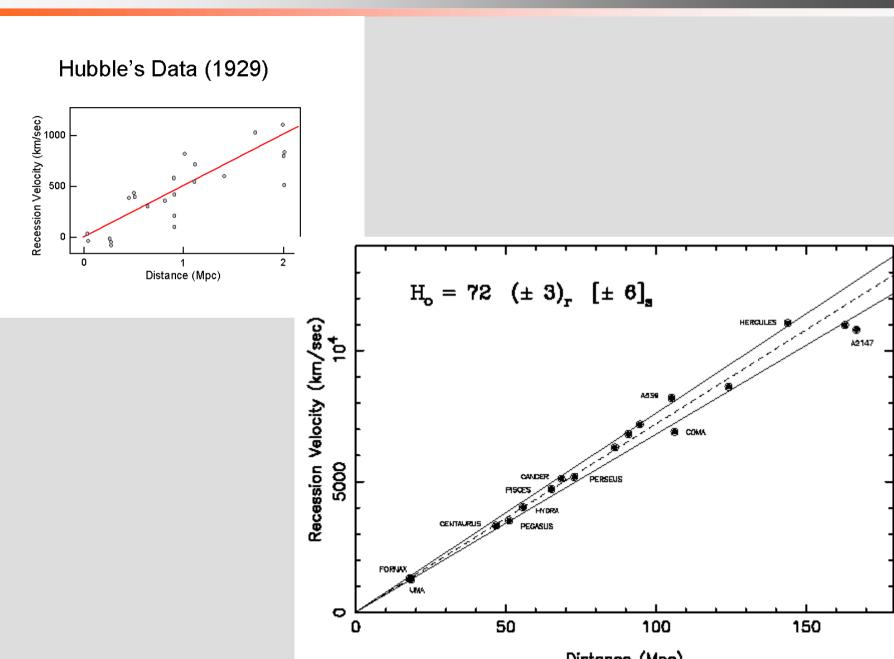
An Object becomes fainter by the square of its distance

$$\phi_{
m obs} = rac{L_{
m source}}{4\pi D_L^2}$$



$$D_A^2 = rac{\mathrm{d}S_{\mathrm{source}}^{\mathrm{phys}}}{\mathrm{d}\Omega_{\mathrm{obs}}}$$

OCCESSES: EXPANSION



OCCE22E2: I LEKWAL LIZIOKI

Consider a particle of mass M.

In radiation era $H \propto T^2$

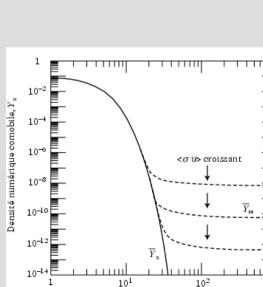
T>M It is in thermal equilibrium with its anti-particle

Competition between dilution and annihilation
$$\Gamma \sim \eta$$

If $\Gamma \propto T^{n+3}$ (n=2 for weak interaction)

There ALWAYS a temperature below which an interaction is frozen

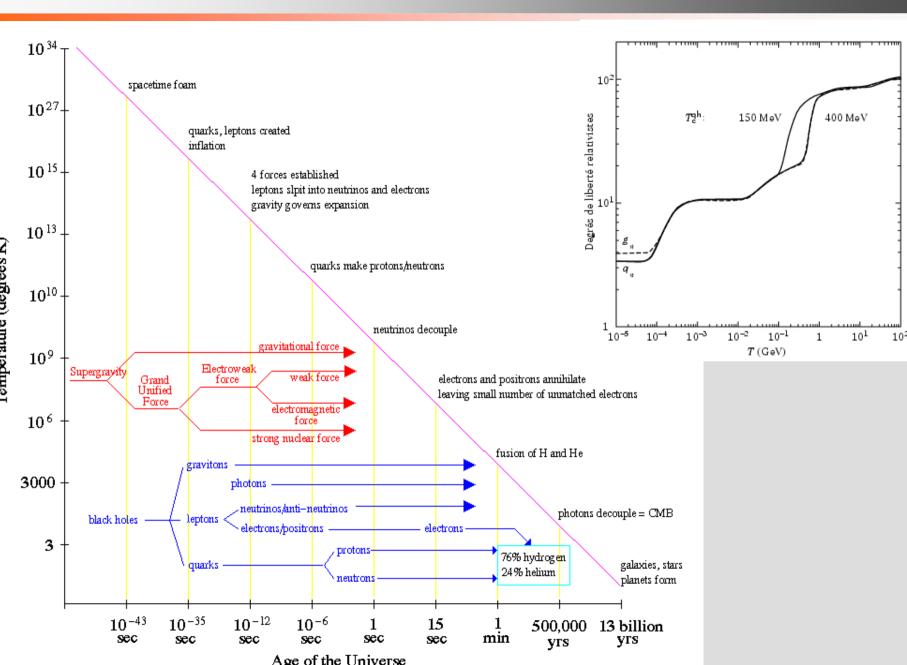
The universe has a THERMAL HISTORY



 $H > \Gamma$ Interaction froze

Interaction effect

OCCE22E2: I LEKWAL LI2 LOKA



OCCE22E2: BIG-RANG MOCLEO21M1HE2I2

T>> 100 Mev Electron, positron, neutrinos and photons: UR / proton, neutron: NR

$$(n/p)_{eq} \sim \exp(-Q/T) \sim 1, \quad Q = m_n - m_p \sim 1.3 \,\text{MeV}$$

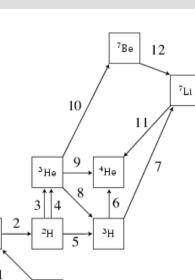
$$(n/p)_f \sim \exp(-\,Q/T_f) \sim 1/5$$
 Free neutrons decay in **887 sec**

 $\Gamma_{\rm weak} \sim H \Leftrightarrow T_f \sim 0.8 \,{\rm MeV}$

Light nuclei are formed by a series of nuclear reactions

$$n + p \rightarrow D + \gamma$$

D can be produced only when T<0.066MeV is low enough so that photo-dissociation



 $^{1}\mathrm{H}$

1. $p \leftrightarrow n$ 2. $p(n, \gamma)d$

3. $d(p, \gamma)^3$ He 4. $d(d, n)^3$ He

d (d, p)t
 d (d, n)⁴He

7. t (α, γ)⁷Li

8. ³He (n, p)t 9. ³He (d, p)⁴h 10. ³He (α, γ)

11. ⁷Li (p, α)⁴

Helium-4

T=0.7-0.05 MeV

negligible is negligible



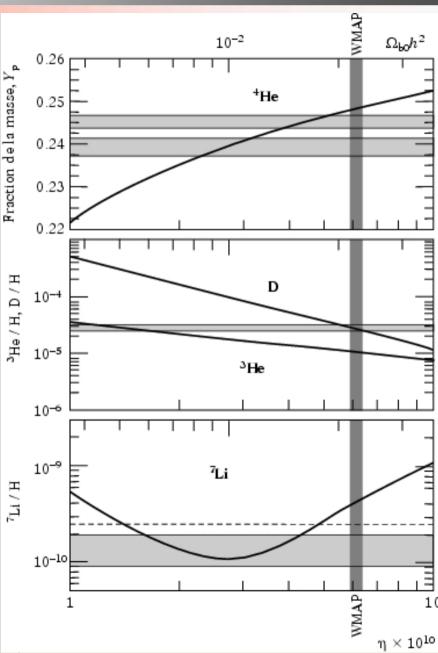
OCCE22E2: RIG-RANG NOCLEO21N1 HE212

Parameters:

- -Number of relativistic particles - lifetime of neutron
- η = n_{baryon}/n_{γ} $G, G_F, \alpha,...$

Allows to test extensions of the standard model of particle physics





OCCE22E2: CO2WIC WICKOMANE RACKROOND

The universe cools during expansion

Around **T~4000 K**, protons and électrons Combine to form hydrogene.

The universe becomes transparent

<u>Samow argument</u>

Knowing the baryonic density today

$$n_{b0} \sim 10^{-7} \, \mathrm{cm}^{-3}$$

and at BBN

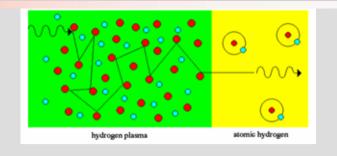
$$n_b \sim 10^{18} \, {\rm cm}^{-3}$$

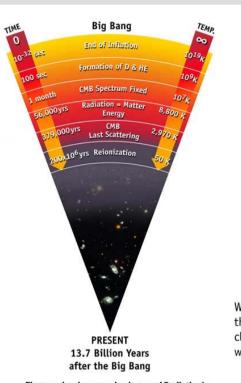
he inferred the redshift at BBN

$$1 + z_{BBN} \sim (n_b/n_{b0})^{1/3} \sim 2 \times 10^8$$

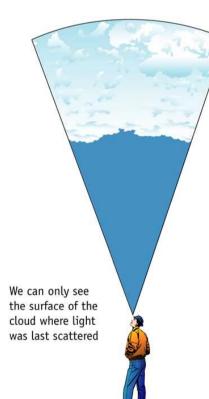
and the temperature of the photon bath today

$$T_{\gamma 0} = rac{T_{BBN}}{1+z_{BBN}} \sim 5K$$





The cosmic microwave background Radiation's "surface of last scatter" is analogous to the light coming through the clouds to our eye on a cloudy day.



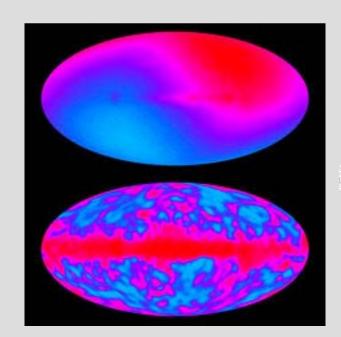
OCCE22E2: CO2WIC WICKOMANE RACKOKOONN

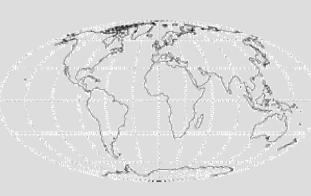
Emission d'un fond de photons avec un spectre de corps noir à une température de 2.725K aujourd'hui.

COBE observation

Dipole after Monopole substraction

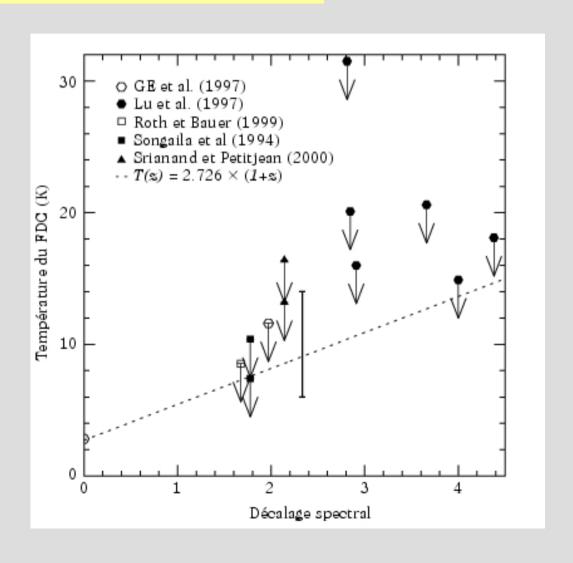
After dipole Substraction: Fluctuation of order μK





OCCE22E2: CO2WIC WICKOMAAF RACKROOND

Temperature of CMB scales as 1/(1+z)



AKAMETERS OF THE WODEL

4 numbers to describe the dynamics of the universe

$$\Omega_m$$
, Ω_r , Ω_K , Ω_Λ , H_0

They start to be accurately measured

- we shall see how later

- ``precision cosmology"
- This is a very successful model
- The universe cools down from a hot thermal equilibrium state

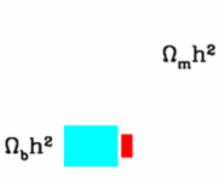
<u>Successes:</u>

- -expansion observed (Hubble law and redshift)
- light nuclei abundances (RG and weak interaction)
- CMB (RG and electromagnetism)



■ 10 Years ago

Now





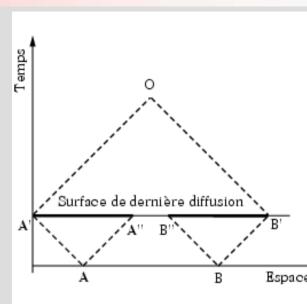
KORLEW2 AND GOESTIONS

Flatness

$$|\Omega_{K0}| < 0.1 \Rightarrow |\Omega_K(z_{pl})| < 10^{-60}$$

Horizon

CMB isotropic but corresponds to 10⁸⁷ causal zones. How do they reach thermal equilibrium?



Origin of structures

The universe is obviously not smooth. Where do the structure come from?

Dark sector $\Omega_r:\Omega_b:\Omega_m:\Omega_\Lambda\sim 10^{-3}:1:5:14$

Good description up to approx. 10¹⁶ GeV

- Effect of the GUT unification scheme on the particle content
- Topological defects...
- Close to 10¹⁹ GeV, we expect to have effect of quantum gravity
- We have a window on energies not accessible in accelerator!

PART III: GRAVITATIONAL DYNAMICS AND LARGE SCALE STRUCTURES

Main topics

- Newtonian regime
- General cosmological perturbation
- Description and parameters

KAVITATIONAL DYNAMICS IN A STATIC SPACETIME

In a static spacetime, hydrodynamics takes its standard form

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = 0$$
$$\partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{1}{\rho} \nabla P - \nabla \Phi$$

If we expand quantities as $P = \bar{P} + \delta P, \dots$ we obtain the wave equation

$$\partial_t^2 \delta \rho - \Delta \delta P = \bar{\rho} \Delta \Phi$$

Use the Poisson equation

$$\Delta\Phi = 4\pi G \delta\rho$$

and the definition of the sound speed

$$c_s^2 = (\delta P/\delta \rho)_s$$

You get the propagation equation

$$\partial_t^2 \delta \rho - c_s^2 \Delta \delta \rho = 4\pi G \bar{\rho} \delta \rho$$

EANS LENGIN

If we decompose in plane waves exp[i(ωt-k.x)], we get the dispersion relation

$$\omega^2 = \frac{4\pi c_s^2}{\lambda_J^2} \left(\frac{\lambda_J^2}{\lambda^2} - 1 \right)$$

$$\lambda_J \equiv c_s \, \sqrt{rac{\pi}{Gar
ho}}$$
 is the Jean length

$$\lambda > \lambda_J$$
 Gravity dominates, density perturbations grow

$$\lambda < \lambda_J$$
 Gravity negligible, sound waves

$$t_{\rm sound} \sim \lambda c_s, \quad t_{\rm collapse} \sim \sqrt{G\bar{\rho}}$$

 $t_{
m sound} \ll t_{
m collapse}$ Pressure can compensate gravitational collapse

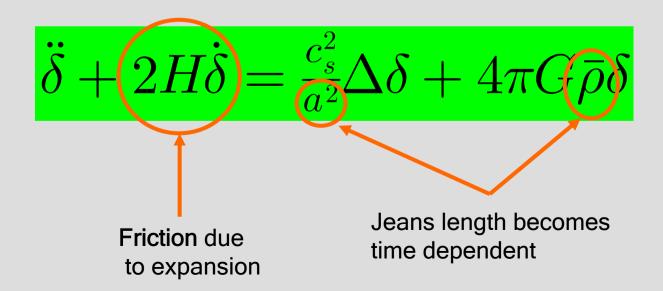
EWIONIAN DYNAMICS IN A EXPANDING UNIVERSE

In an expanding universe, one can follow the same route for wavelengths smaller than the Hubble radius, H^{1} ,

Physical coordinates are time dependent $\mathbf{r} = a(t)\mathbf{x}$

It follows that
$$\begin{cases} \nabla_{\mathbf{r}} \to \frac{1}{a} \nabla_{\mathbf{x}} \\ \partial_t \rho(\mathbf{r},t) \to \partial_t \rho(\mathbf{x},t) - H\mathbf{x}. \nabla_{\mathbf{x}} \rho \end{cases}$$

The wave equation becomes

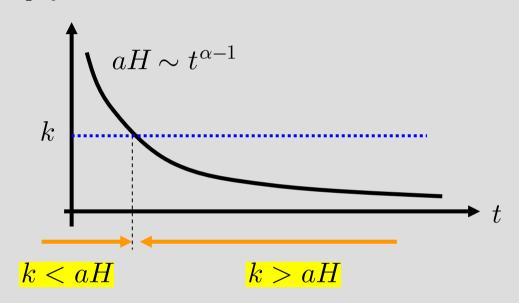


ONG WAVELENGINS

All wavelength are redshifted \rightarrow

$$\lambda_{\rm phys} \propto a$$

Wavenumbers $k_{\rm phys} = k/a$



k < aH

Super-Hubble mode. Pressure negligible.

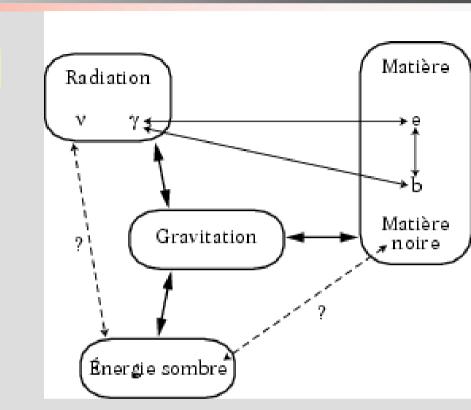
k > aH

Sub-Hubble mode. Expansion negligible.

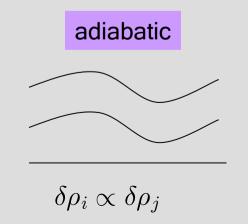
All modes was SUPER-HUBBLE in the past

LE 2121FW IO NE2CKIRE

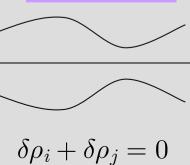
System of coupled differential equations



Nature of the initial conditions



isocurvature



ALLIAL PLECIKOW AND ORPEKAED PLECIKOW

Harrison-Zel'dovich initial power spectrum:

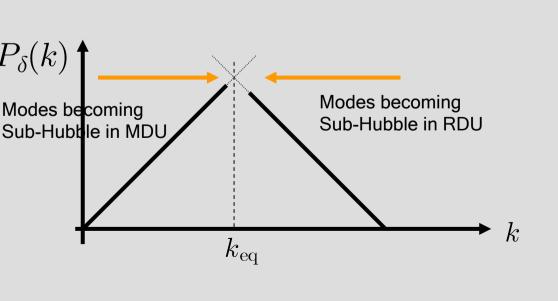
Initial power spectrum of fluctuations on super-Hubble scales

$$P_{\Phi}(k) \propto k^{-3} \iff P_{\delta}(k) \propto k$$

No justification at this level. Origin of the spectrum?

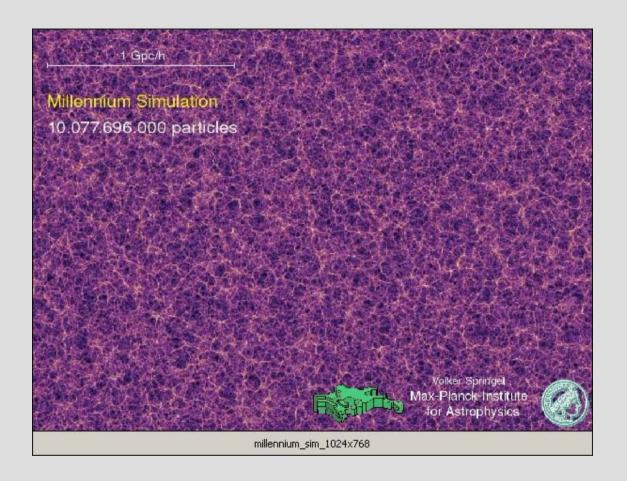
Observed spectrum

Modes that became sub-Hubble during radiation era remained constant until equalit There amplitude is thus suppressed by a factor $(k_{eq}/k)^2$.



 $P_{\delta}(k) \propto k^{-1}$ $(k < k_{\rm eq})$

OWEKICAL SIMOLATION



AKAWEIEK2 OF THE WONER

5 numbers to describe the dynamics of the background universe

$$\Omega_m$$
, Ω_b , Ω_r , Ω_K , Ω_Λ , H_0

1 function to describe the scalar perturbation

$$P(k): A, n_S, \alpha_S$$
 (?)

1 function to describe the gravity waves

$$P_T(k): T/S, \quad n_T, \quad \alpha_T \ (?)$$

Reionisation parameter (needed for CMB)

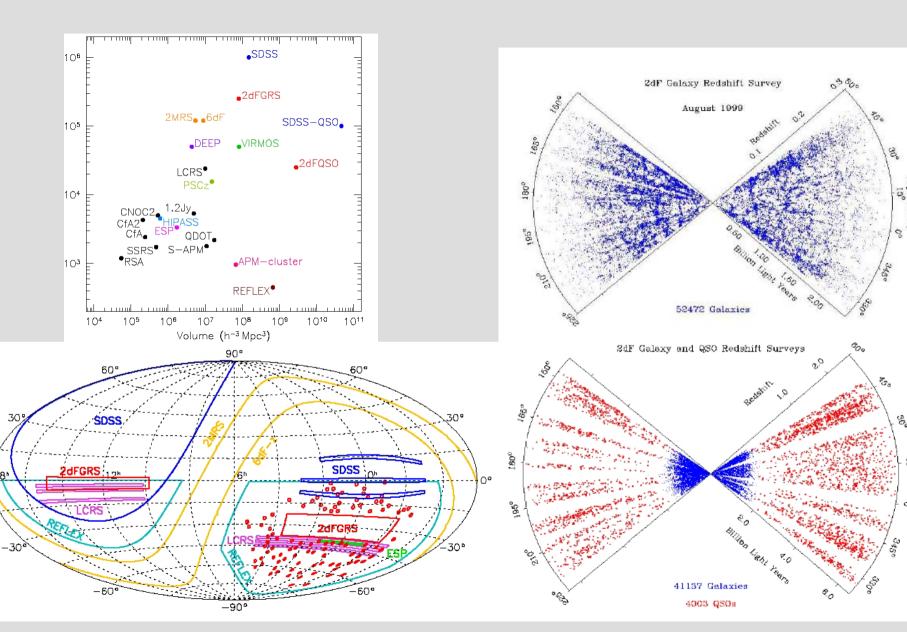
 ${\mathcal T}$

PART IV: OBSERVING THE LARGE SCALE STRUCTURE

Main topics

- Galaxy catalogs
- Cosmic microwave background
- Weak lensing
- •Observational status

ALAXY CATALOGS

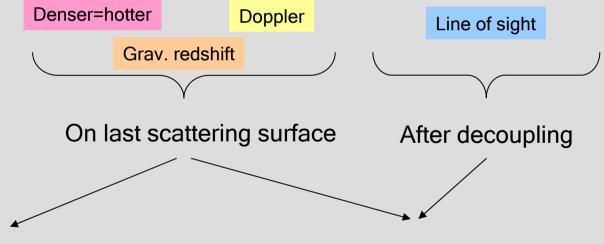


MD ANISUTKUPIES: UKIGIN

COBE has detected small temperature fluctuation in the CMB of order 10⁻⁵.

Sachs-Wolfe relation

$$\frac{\delta T}{T}(\mathbf{e}) = \frac{1}{4}\delta_{\gamma} + \Phi + \mathbf{e}.\mathbf{v}_b + \int (\Phi' + \Psi')d\eta$$

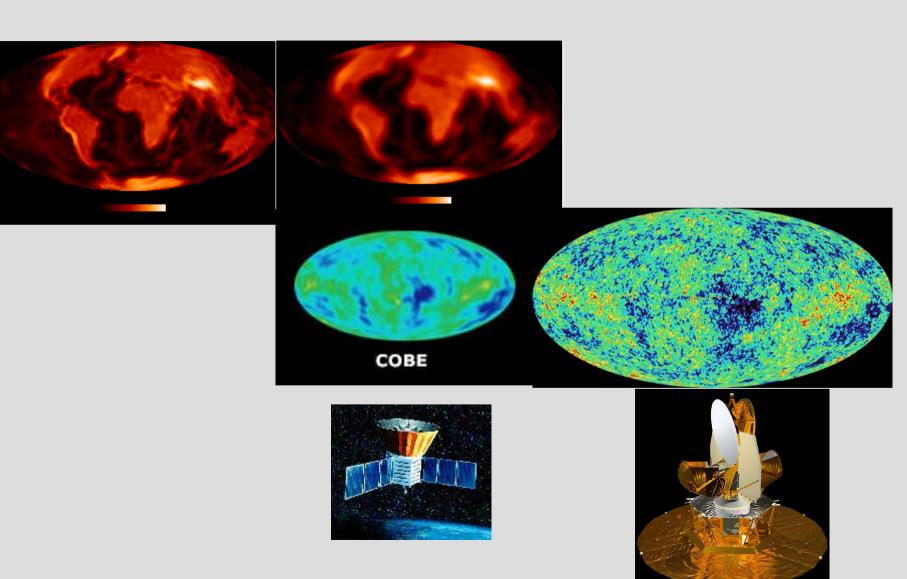


On large scales: observed modes were super-hubble at the time of decoupling. Initial conditions!

Evolution effects

WID: ORPEKAUTION

This temperature fluctuations have been observed with increasing resolution



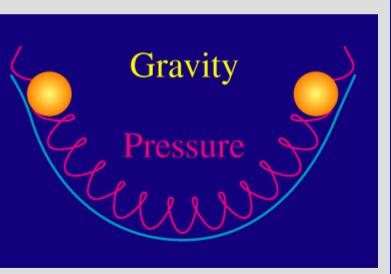
MD: ANGULAK PUWEK SPECIKUM

Angular power spectrum

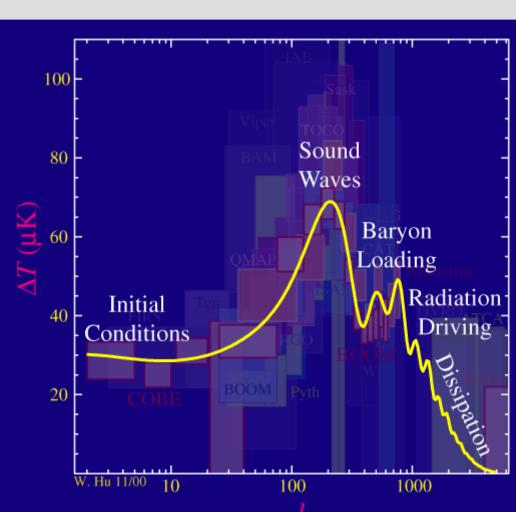
$$\langle \delta T(\mathbf{e}_1) \delta T(\mathbf{e}_2) \rangle = \sum \frac{2\ell+1}{4\pi} P_{\ell}(\cos \theta)$$

Before decoupling

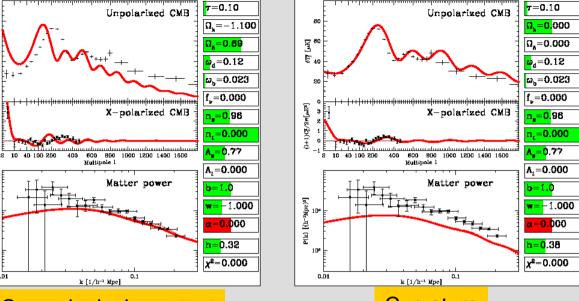
Oscillation in the photon-baryon plasma

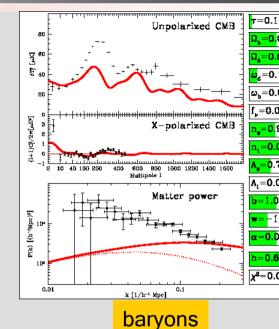


Characteristic scale



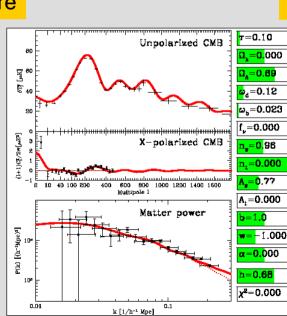
MD: EFFECT OF COSMOLOGICAL PAKAMETERS



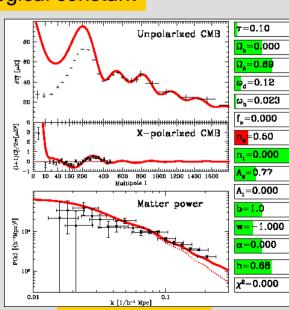


Cosmological constant

Curvature



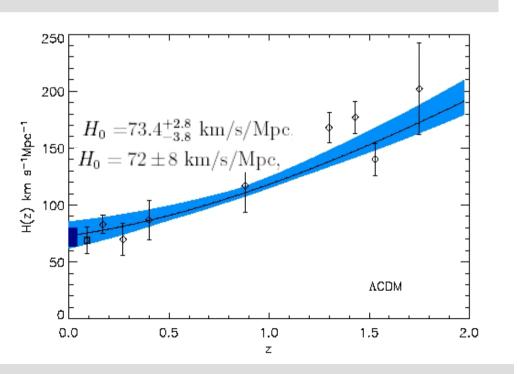
nautrings

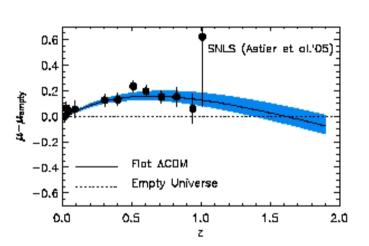


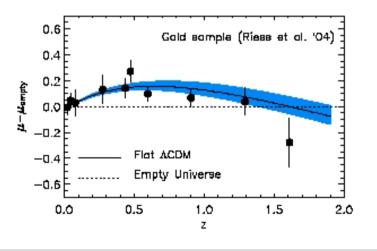
Spectral index

Analysis of a Λ CDM model with 6 parameters (h, $\Omega_{\rm m}$, $\Omega_{\rm b}$, τ , n_s, $\sigma_{\rm 8}$)

$$t_0 = 13.73^{+0.13}_{-0.17} \text{ Gyr}$$

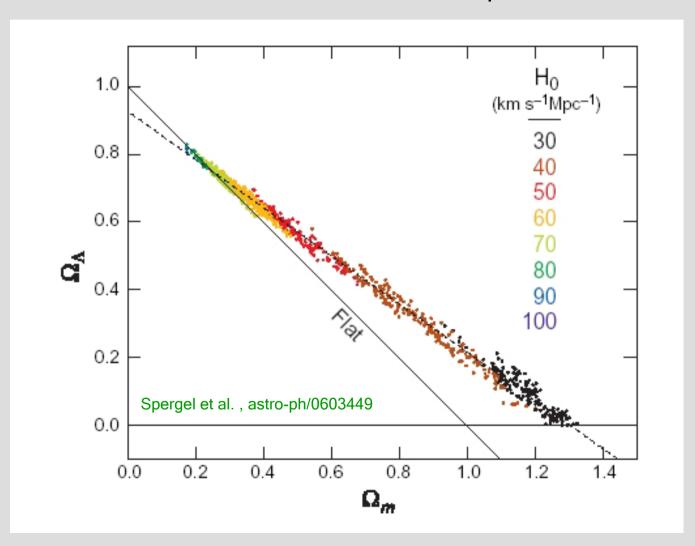






MMAP - CUKVATUKE OF SPACE

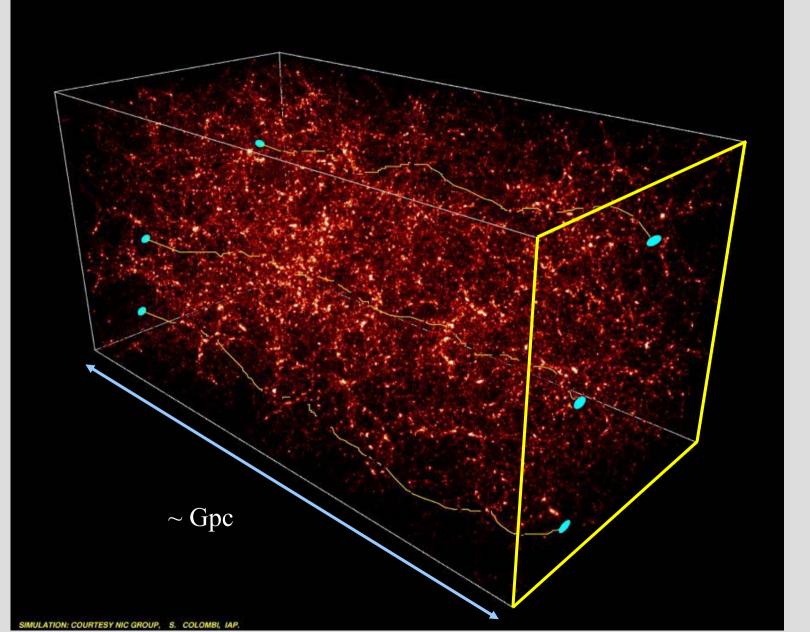
Non-flat ACDM models are compatible with WMAP



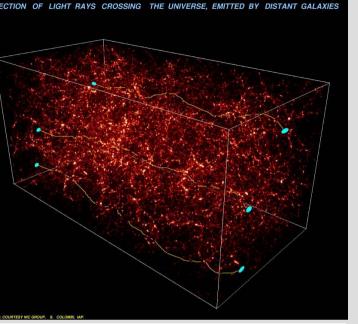
KAVITATIONAL LENSING: PRINCIPLE

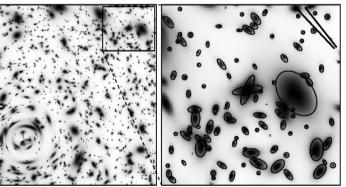


DEFLECTION OF LIGHT RAYS CROSSING THE UNIVERSE, EMITTED BY DISTANT GALAXIES



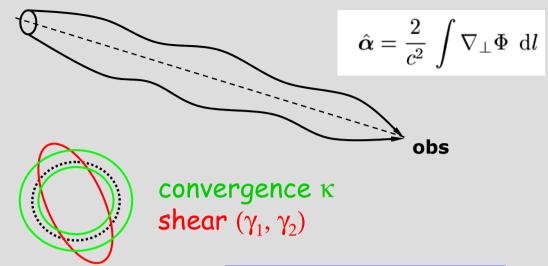
VEAK LENSING





$$M_{ij} = \frac{\int I(\boldsymbol{\theta}) \,\theta_i \,\theta_j \,d^2 \theta}{\int I(\boldsymbol{\theta}) \,d^2 \theta} \qquad \frac{a^2 - b^2}{a^2 + b^2}$$

$$\frac{a^2 - b^2}{a^2 + b^2}$$

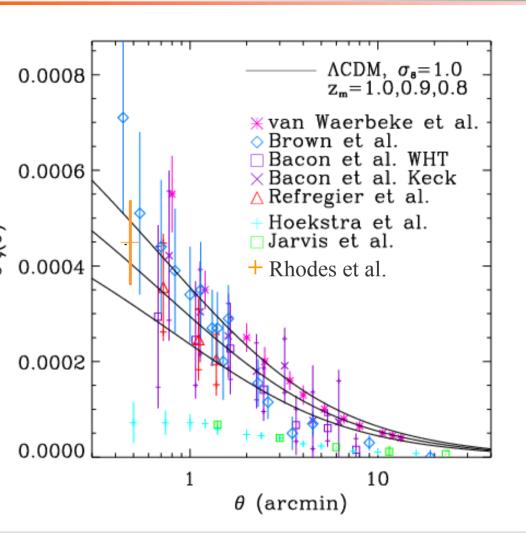


$$A = \begin{pmatrix} 1 - \kappa - \gamma_1 & \gamma_2 \\ \gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix}$$

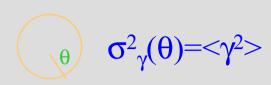
$$\delta = \frac{2\gamma (1 - \kappa)}{(1 - \kappa)^2 + |\gamma|^2} = \frac{a^2 - b^2}{a^2 + b^2}$$

$$\delta = 2\gamma$$

LEAK FENDING ORDEKANITOND



Shear variance

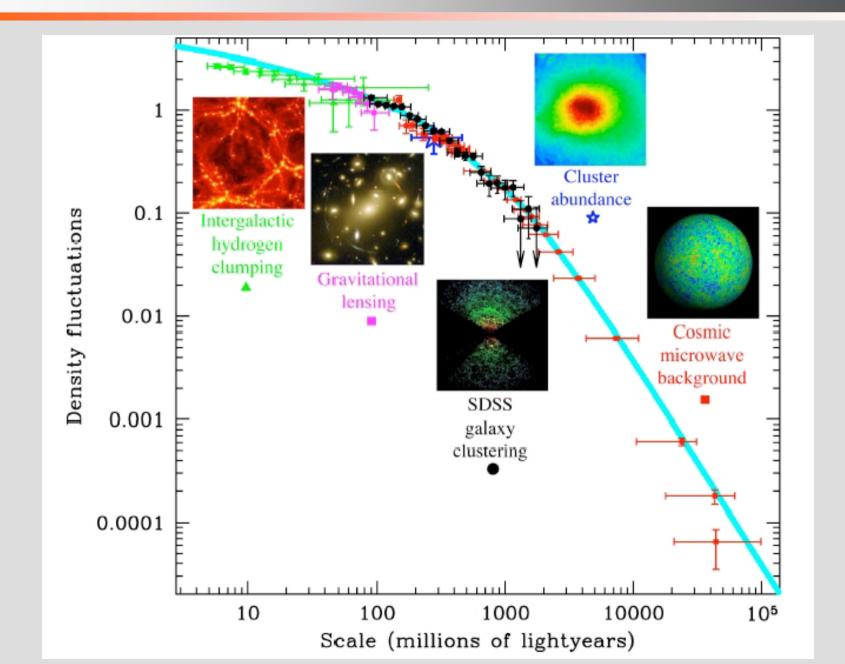


Bacon, Refregier & Ellis 2000* Bacon, Massey, Refregier, Ellis 2001 Kaiser et al. 2000* Maoli et al. 2000* Rhodes, Refregier & Groth 2001* Refregier, Rhodes & Groth 2002 van Waerbeke et al. 2000* van Waerbeke et al. 2001, 2005 Wittman et al. 2000* Hammerle et al. 2001* Hoekstra et al. 2002 * Brown et al. 2003 * not shown Hamana et al. 2003 * Jarvis et al. 2003 Casertano et al 2003* Rhodes et al 2004 Massey et al. 2004

Heymans et al 2004*

In agreement with a ACDM and the gravitational Instability paradigm

OUCTO210N



PART V: INFLATION

Main topics

- Principle
- Resolution of the BB problems
- Origin of perturbations
- Eternal inflation
- Observational constraints

ALLAHON: RASICS

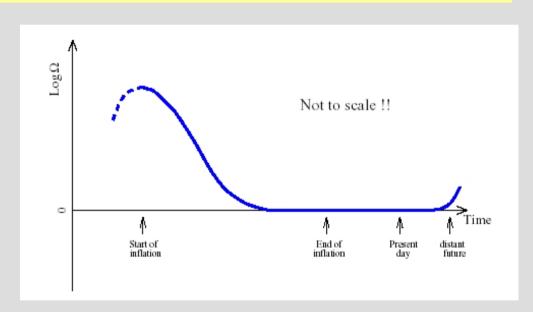
The origin of the flatness problem is clear: $\Omega_K = -\frac{K}{a^2H^2}$

$$\Omega_K = -rac{K}{a^2H^2}$$

During the cosmological evolution *aH* decreases

Assume there is a primordial phase during which aH increases

$$\Omega_K \to 0$$



Inflation = primordial phase of accelerated expansion

$$aH \nearrow \Leftrightarrow \ddot{a} > 0 \Leftrightarrow \rho + 3P < 0$$

TAINE 22 LKORTEW

If the inflation phase is long enough then WK is brought very close to 0, explaining the flatness of our universe.

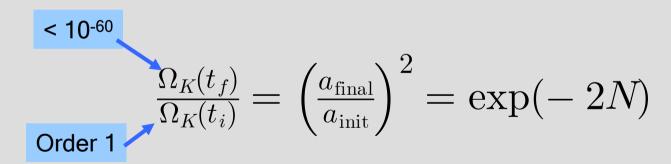
E-fold

We define the number of e-fold of inflation by

$$N = \ln \left(\frac{a_{\text{final}}}{a_{\text{init}}} \right)$$

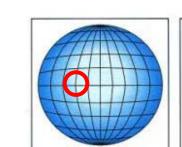
Solving the problem

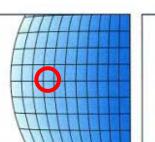
If H is almost constant during this phase then

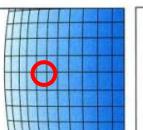


Thus, we need at least











OKIZON PROBLEM Infinite redshift Vithout Inflation Horizon at decoupling N~10⁸⁷ Surface de dernière diffusion Last scattering surface Espace With Inflation

MPLEMENIAIION

We need matter satisfying ρ + 3P <0

Cosmological constant: de Sitter phase

exponential expansion

Scalar field:

$$\rho = \frac{1}{2}\dot{\phi}^2 + V(\phi), \quad P = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$

so that

$$\rho + 3P = -2V(\phi) \times \left(1 - \frac{\dot{\phi}^2}{V}\right)$$

The field must be in slow-roll

TOM-KOLL CONDITIONS

The 2 slow-roll conditions

$$\dot{\phi}^2 \ll V$$

$$\ddot{\phi}^2 \ll 3H\dot{\phi}$$

Evolution equations

obtained from the Friedmann and Klain-Gordon equations

$$H^2 \simeq \frac{8\pi G}{3}V$$

$$3H\dot{\phi} \simeq -V'$$

The expansion will be quasi-exponential

Validity conditions

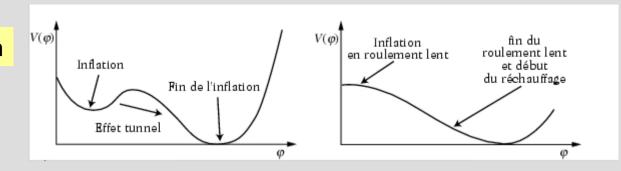
The slow-roll conditions are fulfilled if

$$(V'/V)^2 \ll 24\pi G$$
, $V''/V \ll 24\pi G$

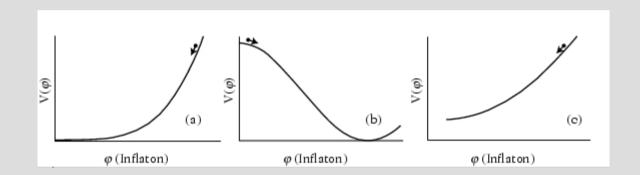
The potential must be flat

HE TOO OF LOIENIIST

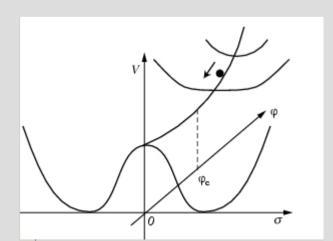
Old and new inflation



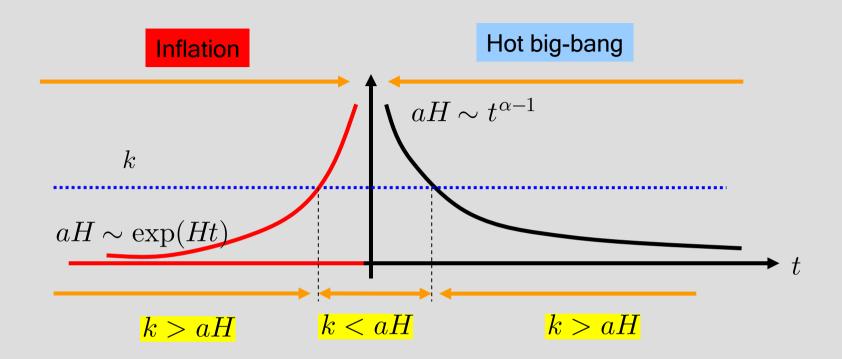
Large/small field



Hybrid inflation



KIGIN OF FLUCTUATION: MODE EVOLUTION



The modes

- * start sub-Hubble
- * exit the Hubble radius during inflation
- * enters the Hubble radius during the matter or radiation era

F21 LIFTA IN AF 211 IFK

To grasp the origin of perturbation, consider a test massless scalar field in de Sitter

Its equation of evolution is the Klein-Gordon equation

$$\ddot{\chi} + 3H\dot{\chi} + \frac{k^2}{a^2}\chi = 0$$

k > aH

Friction negligible - harmonic oscillator.

k < aH

Gradient negligible - constant mode

Setting $v=a\chi$, and using $\eta=-exp(-Ht)/H$ the Klein-Gordon equation becomes

$$v'' + \left(k^2 - \frac{2}{\eta^2}\right)\chi = 0$$

Its solution is

$$v = A(k) \left(1 + \frac{1}{ik\eta}\right) \exp(-ik\eta) + B(k)c.c.$$

MITIAL CONDITION

 $k\eta\gg 1$ The curvature is negligible

We can quantify as in Minkowski space

Initial conditions

$$v \to \frac{\exp(-ik\eta)}{\sqrt{2k}} \quad k\eta \to -\infty$$

Solution

$$\chi = \frac{H\eta}{\sqrt{2k}} \left(1 + \frac{1}{ik\eta} \right) \exp(-ik\eta)$$

Frozen super-Hubble fluctuations

 $\chi_k \sim \frac{H}{\sqrt{2k^3}}$

scales $\chi_k \xrightarrow[l_m]{} \frac{\mathrm{e}^{-ikx}}{\sqrt{2k}}$

Plane wave on small

10⁻³³ cm

 10^{-28} cm amplitude ~ 10^{-5} M₄

OLUIZ I ICA I IONZ

The inflaton also fluctuates

But its fluctuations are coupled to gravity

What is the variable to quantify?

Heuristically

Because of the initial fluctuations in the scalar field, inflation last more or less longer from one Hubble patch to another.

The curvature perturbation from one patch to the other is

$$\delta \mathcal{R} \sim H \delta t \sim H \delta \phi / \dot{\phi}$$

The fluctuation in the inflaton are of order

$$\delta\phi/\dot{\phi}\sim H/2\pi$$



Gravity waves

One produces GW with an amplitude of order

$$h \sim \frac{H}{2\pi M_p}$$

CALE OF INFLATION

On super-Hubble scales, the CMB anisotropy is of order

$$\frac{\delta T}{T} \sim \frac{1}{3}\Phi$$

The amplitude of the CMB temperature anisotropies allow to calibrate the spectrum

$$\frac{\delta T}{T} \sim 2 \times 10^{-6} \ K$$
 $\sim 5.7 \times 10^{16} \ \mathrm{GeV}$

From WMAP data, we can infer the bound

$$\varepsilon < 0.032 \Leftrightarrow \frac{H_{inf}}{M_p} < 1.4 \times 10^{-5}$$

$$\varepsilon \equiv (V/V)^2/16\pi G$$

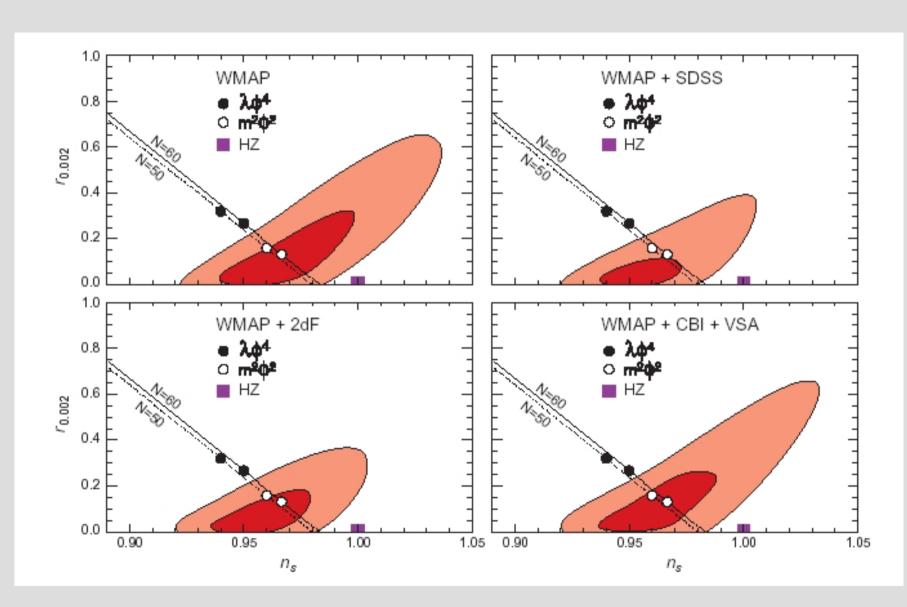
ENERIC PREDICTIONS

- 1- universe is *flat*: Ω =1
- 2- classical inhomogeneities are erased Justification of the *cosmological principle*
- 3- metric perturbations are Gaussian, almost scale invariant origin of the Harrison-Zel'dovich spectrum
- 4- no vector perturbation
- 5- tensor perturbations are almost scale invariant
- 6- there exists a *consistency relation* between T/S, n_t et n_s

extensions

- multi-field non-gaussianity, isocurvature modes

MAT CONSTRAINTS



IEKNAL INFLATION

Consider a potential $V = \frac{1}{2}m^2\phi^2$

In slow-roll, during a time step $\delta t \sim H^{-1}$

Classical motion

$$\Delta\phi_{cl}\sim\dot{\phi}\delta t\sim\dot{\phi}/H\sim-M_p^2/4\pi\phi$$

Quantum fluctuation

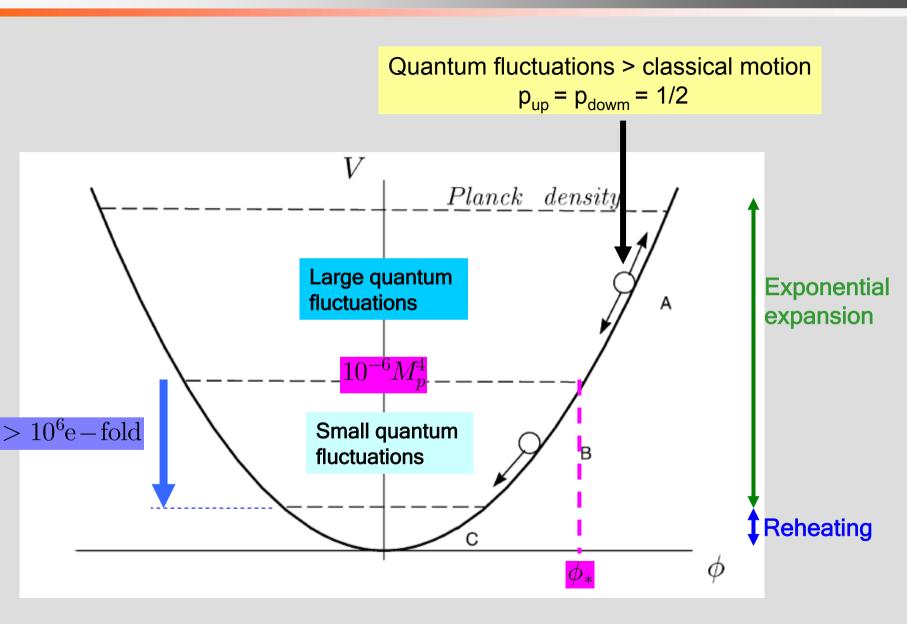
$$\delta\phi_q \sim H/2\pi$$

The fluctuations are of same amplitude when

$$\phi \sim \phi_* = \frac{M_p}{2} \sqrt{\frac{M_p}{m}}$$

Note that $V(\phi_*) \sim 10^{-6} M_p^4$

IEKNAL INFLATION



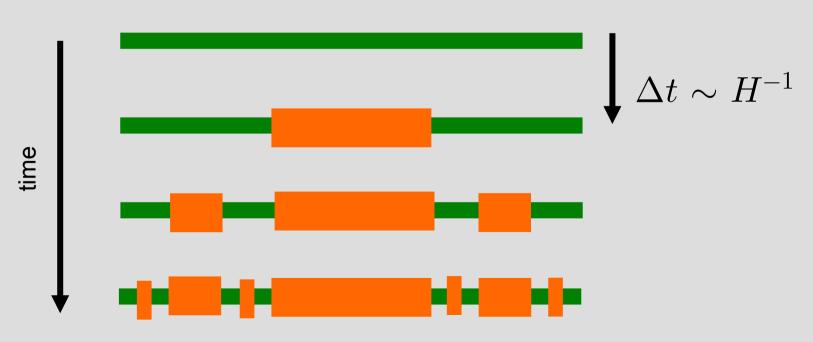
ELF KEPKODUCING UNIVERSE

De Sitter inflation : $a(t) = \exp(Ht)$

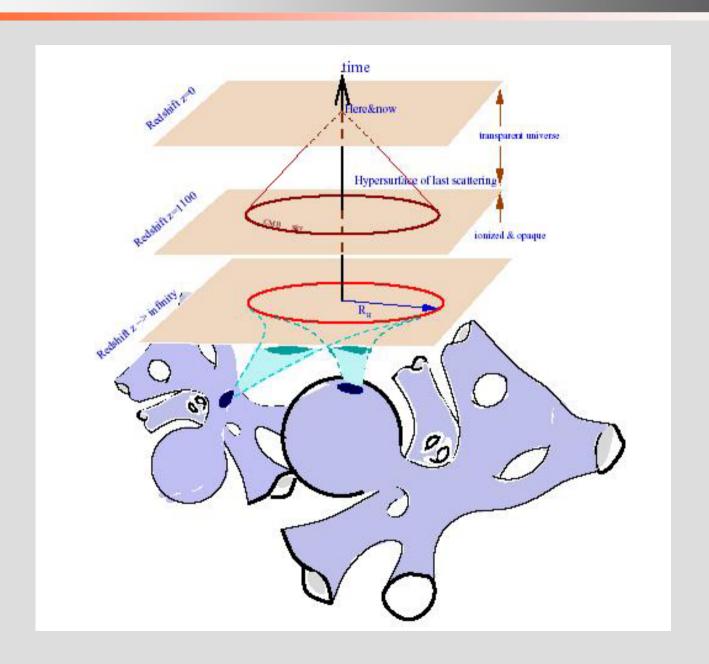
Consider a patch of size H-1

During a time step of H-1

$$V_{t+\Delta t}(\phi > \phi_*) \sim \frac{1}{2} \left(e^{H\Delta t}\right)^3 V_t(\phi > \phi_*) \sim \frac{1}{2} \times 20 \times V_t(\phi > \phi_*)$$



MEM LICIORE OF THE ANIAEK2E



PART VI: THE INPUT FROM HIGH ENERGY PHYSICS

Main topics

- Dark matter
- Dark energy
- •String inspired cosmology
 - Extra-dimensions
 - •branes

AKK MATTEK: GALAXT KUTATION CUKVE

Keplerian motion

$$\frac{v^2(r)}{r} = \frac{GM(\langle r)|}{r^2}$$

Spiral galaxy

$$v \to v_{\infty}$$

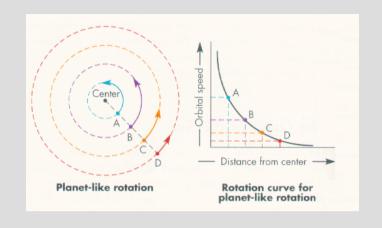
which implies

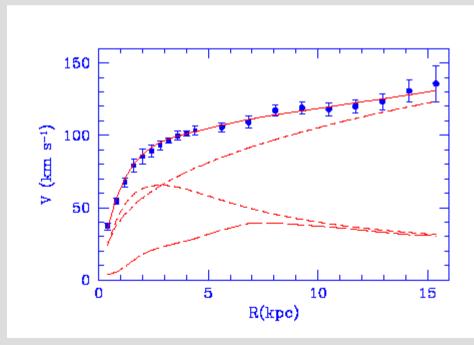
$$M(< r) \propto r, \qquad \rho \propto 1/r^2$$

Tully-Fisher

There is a scaling law

$$L \propto v_{\infty}^4$$





Most of the mass of the galaxy (halo) is dark

SAKTUN USCILLATIUNS

Radiation pressure



Oscillations in the photon fluid

Thomson scattering

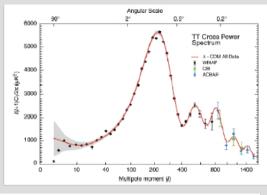


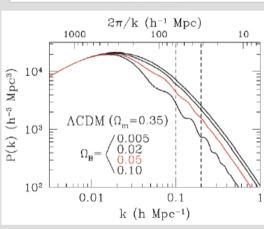
Baryons oscillate with photon Not the dark matter

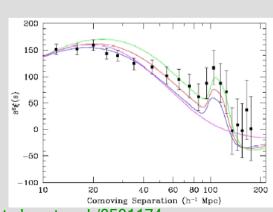
Characteristic scales of the oscillations sonic horizon

Amplitude is attenuated by a factor $\Omega_{\rm b}$

Detected in jan. 2005 (SDSS+2dF)

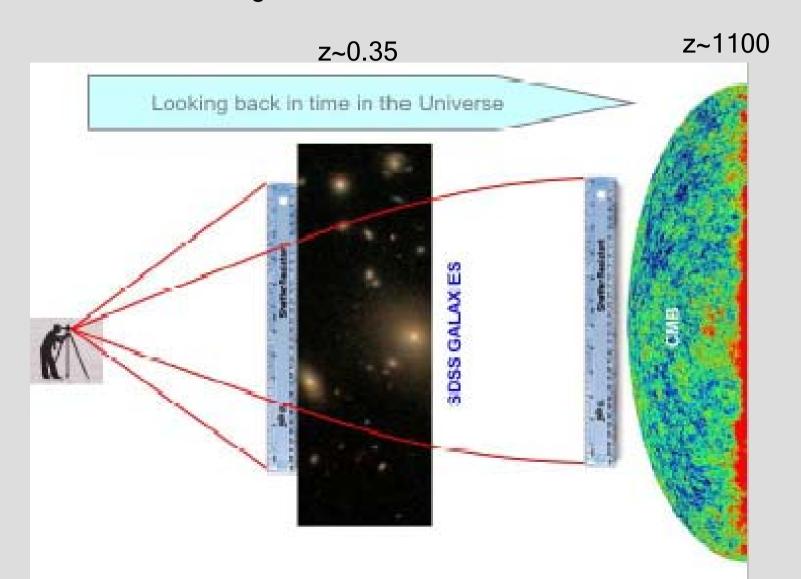




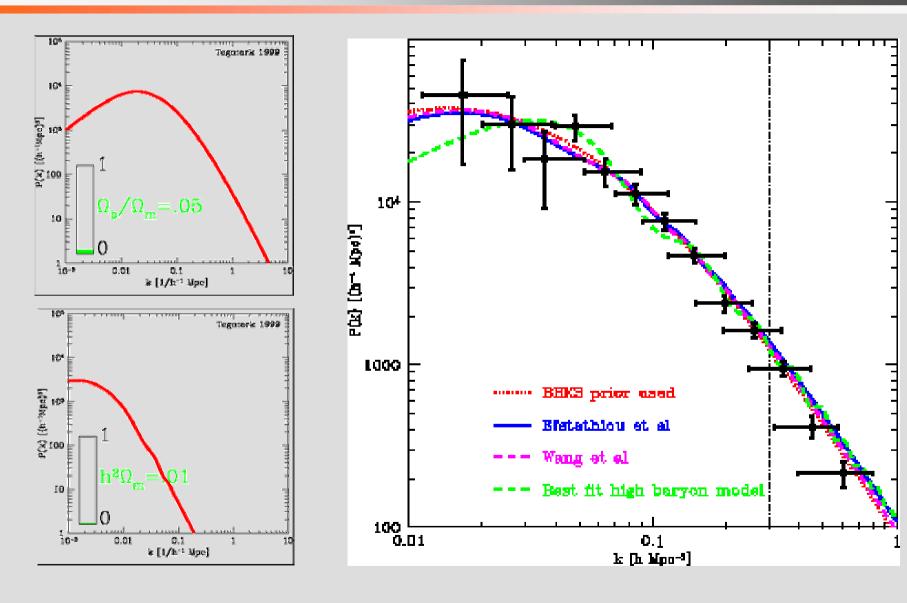


AKTON OSCILLATIONS

Gives a measure of the angular distance



AKK MAIIEK: LJJ EVIDENCE



On cosmological scale all the matter CANNOT be baryonic

AKK MALIEK: SUMMAKT

Evidences

- Galaxy rotation curves
- Galaxy power spectrum
- Dynamics of clusters
- BBN

$\Omega_{ m matter} \sim$	23%,	$\Omega_{\mathrm{baryon}} \sim 4\%$
		.5 5.1

scaling

Scale	Spiral galaxies	clusters	universe
Baryon/DM	8.5±1.5	7.2±2.0	4.83±0.87

Question

Nature of the matter - WIMP

Caveat

Gravity described by GR on these scales

AKK MAIIEK: NAIUKE!

WIMP

Weakly interacting massive particle Abundance depends on m and Γ

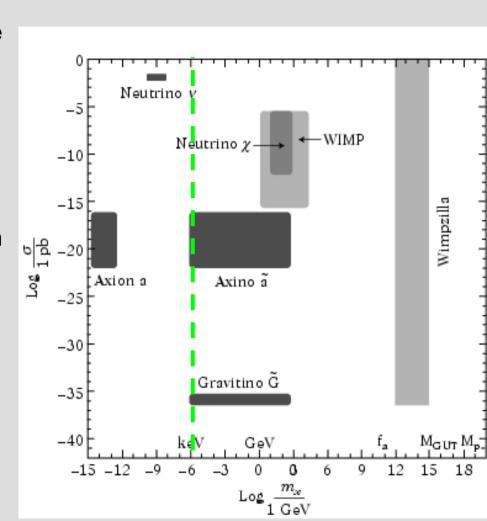
Classification

Hot - cold: mass of the particle at the time of galaxy formation



<u>Thermal-non termal:</u> e.g. neutrinos-neutralinos

Theoretical: existing (neutrino) motivated (LSP, axion)



AKK MAIIEK

Modifying Newton law

$$m\mathbf{a}\mu(a/a_0) = \mathbf{F}$$

$$\frac{\mu(x) = x \quad x \ll 1}{= 1 \quad x \gg 1}$$

Modification in the small acceleration regime

Explain the Tully-Fisher law

New scale

$$\ell_0 = c^2/a_0 \sim 9 \times 10^{26} \text{ m}$$

Difficulties

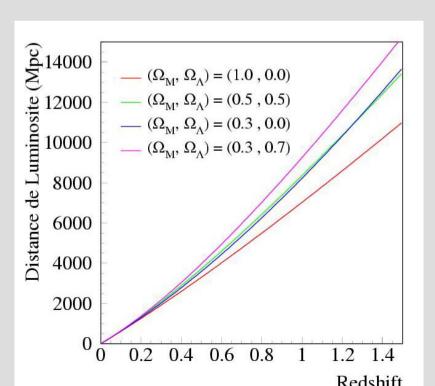
Lacking of a field theory Behavior in other regimes?

UPERNOVAE: PRINCIPLE

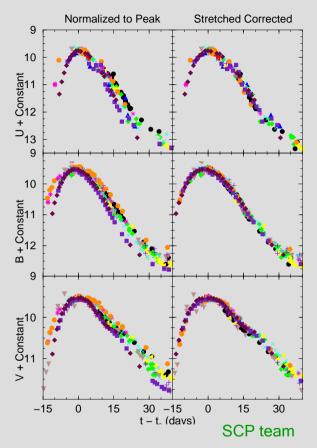
$$\phi_{\rm obs} = \frac{L_{\rm source}(z)}{4\pi D_{\rm lum}^2(z)}$$



$$D_{\text{lum}} \propto (1+z) \int_0^z \frac{dz'}{H(z')}$$



Standard candles?



SNe la exhibit a correlation of Their light (similar to the T-L relation for Cepheids)

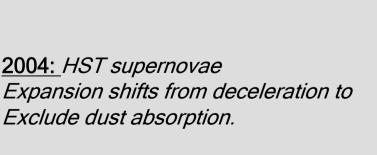
Allow to measure distances with a

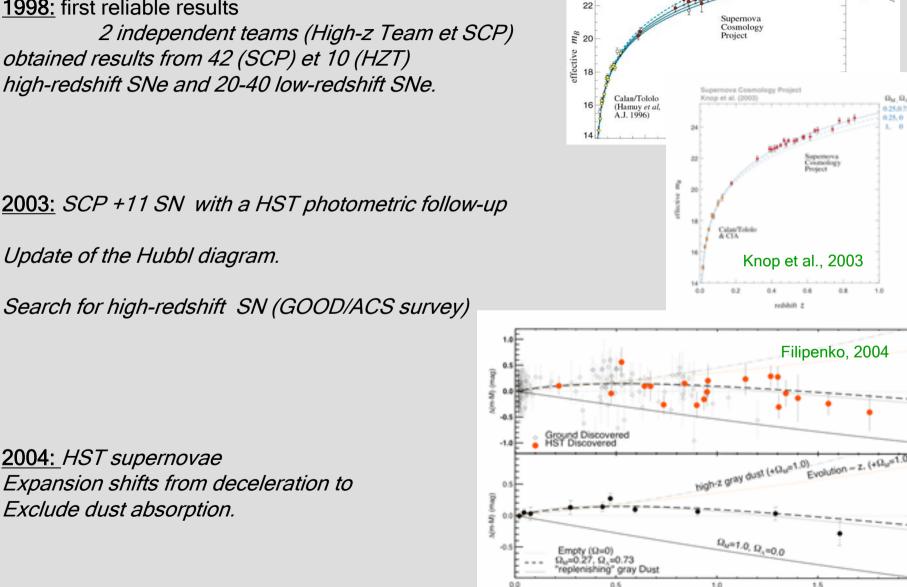
OPEKNOVAE: PKOGKE33E3

1998: first reliable results 2 independent teams (High-z Team et SCP) obtained results from 42 (SCP) et 10 (HZT) high-redshift SNe and 20-40 low-redshift SNe.



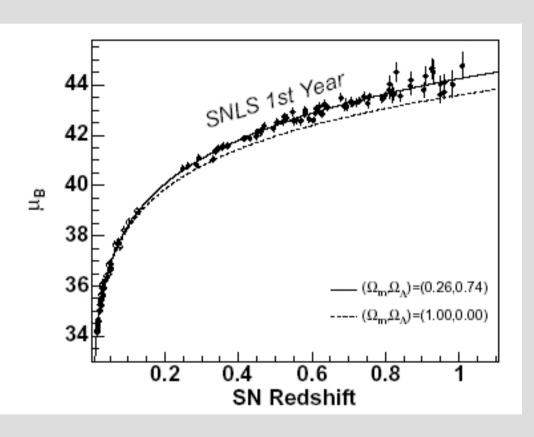
Search for high-redshift SN (GOOD/ACS survey)



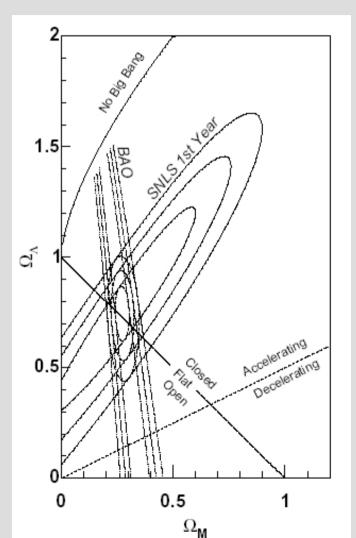


24

Final catalog: 45 close + 71 SNLS SN



For a flat Λ CDM model $\Omega_{\rm M} = 0.263$ +/- 0.042 (stat) +/- 0.032 (syst)



HE ACCELERATION OF THE UNIVERSE

Independently of any theory, we can expand the scale factor as

$$a(t) = a_0 \left[1 + H_0(t - t_0) - \frac{1}{2}q_0H_0^2(t - t_0)^2 + \ldots \right]$$

so that

$$H^2(z)/H_0^2 = 1 + (q_0 + 1)z + \mathcal{O}(z^2)$$
 $q_0 = \Omega_{m0}/2$

The **Hubble diagram** gives

- H₀ at small z
- $-q_0$

Supernovae data (1998+) prove that

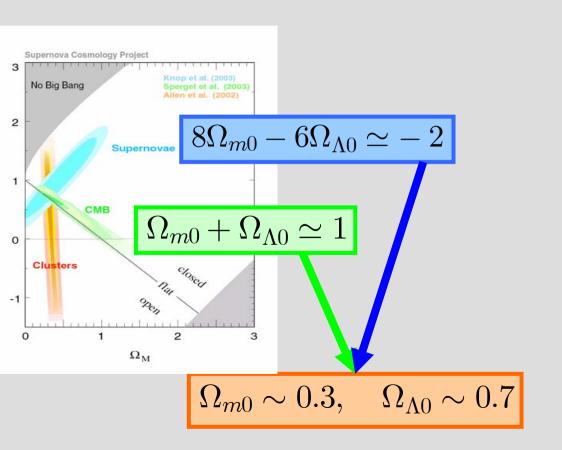


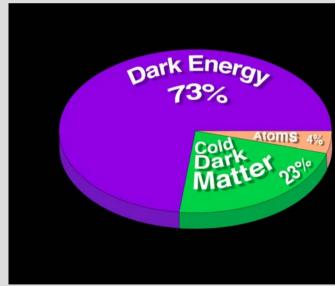
HE ACDM MODEL

The simplest extension is the introduction of a cosmological constant

- Einstein (1917)
- interpretation as vacuum quantum energy
- constant energy density
- well-defined and predictif model.

$$\rho_{\Lambda} = \frac{\Lambda}{8\pi G} = -P_{\Lambda}$$





HE ACDM MODEL

Observationnally, OK with all data

Phénomenologically, very simple (1 parameter)

But

$$\rho_{\Lambda,obs} = \frac{\Lambda}{8\pi G} = H_0^2 M_p^2 = 10^{-47} \text{GeV}^4$$

$$\rho_{\Lambda,th} = M_{\text{fondamental}}^4 > 10^{12} \text{GeV}^4$$

Cosmological constant problem

$$\rho_{\Lambda} > 10^{59} \rho_{\Lambda,obs} !!$$

Today, no solution
Critical problem of fundamental physics

O22IRITITE2

The observed acceleration implies that

$$(
ho + 3P) < 0$$
 singeneral relativity and the Copernician principle hold on cosmological scales

- One must change one of the 3 assumptions of the model
 - 1- The opernician principle is not valid
 - **2-** It exists matter such that ρ +3P<0
 - 3- Gravitation is not described by general relativity on large scales

Nature of the dark energy

FASSES OF WONERS

De nombreux modèles existent,

- liste longue et rébarbative!

On ne peut pas tester les modèles un par un!

- définir des grandes classes de modèles
- peut-on extraire des paramétrisations physiquement motivées pour analyser les données?

But : - donner une idée de la diversité

- avoir une vue des tests permettant de distinguer ces modèles
- définir des stratégies observationnelles
- nécessité d'un modèle au-delà du ΛCDM?

On suppose le problème de la constante cosmologique RESOLU

OPERNICIAN PRINCIPLE

Hypothèses la moins profonde: symétrie de l'espace-temps et non théorie.

Ne touche pas à la relativité / n'invoque pas de nouvelle matière.

Mais la moins étudiée (difficulté technique)

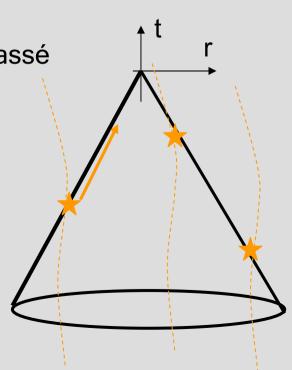
Principe:

on mesure un décalage spectral dégénérescence sur le cône de lumière passé

Principe copernicien:

implique une relation bi-univoque entre redshift et temps d'émission

L'accélération observée pourrait être un effet de structure à grande échelle et d'un mauvais choix de symétries



IEW FORM OF MAILER

Voie la plus étudiée

- Principe : rajouter une nouvelle composante de matière qui
 - 1- domine le contenu de l'univers récemment
 - 2- est telle que ρ +3P<0

Questions à résoudre :

- 1- nature de cette matière
- 2- problème de coïncidence
- 3- compatibilité avec le modèle standard de la physique des particules

Archétype: quintessence

MOINTE 22 FUCE

Champ scalaire évoluant dans un potentiel

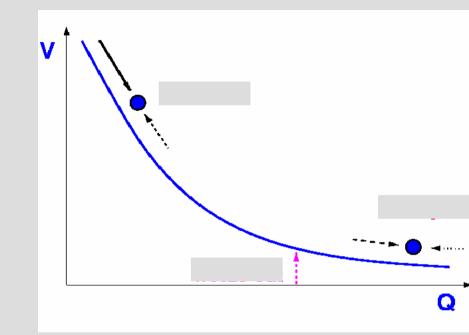
$$\rho = K + V, \quad P = K - V$$
$$-1 \le \frac{P}{\rho} \le 1$$

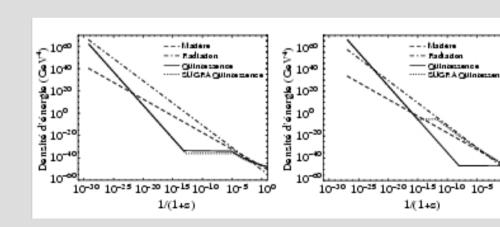
En régime de roulement lent (K<<V)

$$\frac{P}{\rho} = \frac{K-V}{K+V} \simeq -1 + 2\frac{K}{V}$$

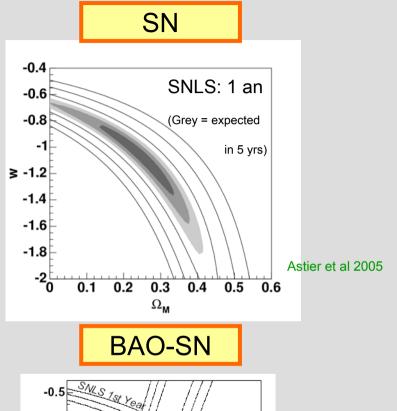
Mécanisme d'attraction

Potentiels justifiables en SUSY





ONSTRAINTS ON A CONSTANT EQUATION OF STATE

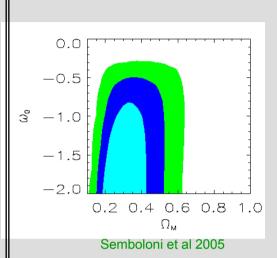


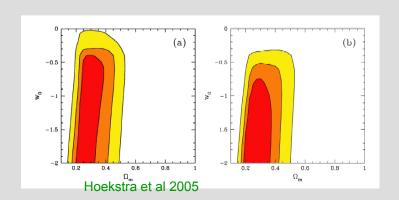
-0.5 SMS 188 Year -1.5 -20 0.2 0.4 0.6

 $\Omega_{\rm M}$

Astier et al 2005 Einsestein et al. 2005

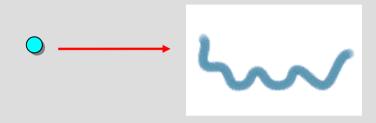
Lensing

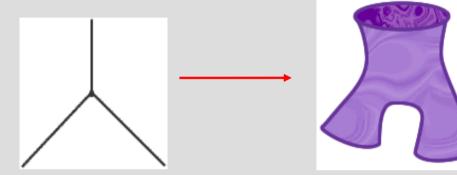




IKING I TEOKT

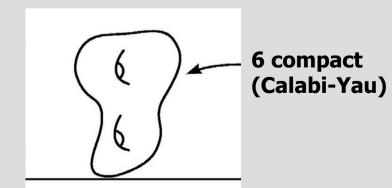
removes the infinities of QFT





- and includes the graviton
- but there is a price!9 space dimensions





IKING INSPIKED COSMOLOGY

Early universe

Extra-dimension dominates at early time / high energy

Models

Pre big-bang models Braneworld model Brane gaz

Phenomenology

Is there a natural mechanism for 3 dimensions to grow very large? Is the observable 3D universe born out of brane-collision? Is there a signature of extra dimensions in observations? Dark matter?

Dark energy?

OW ENERGY LIMIT

Most general theories of gravity include a scalar field beside the metric

Mathematically **consistent** (no ghost, no adynamical field)

Motivated by **superstring**

dilaton in the graviton supermultiplet, **modulii** after dimensional reduction

Only consistent massless field theory to satisfy WEP

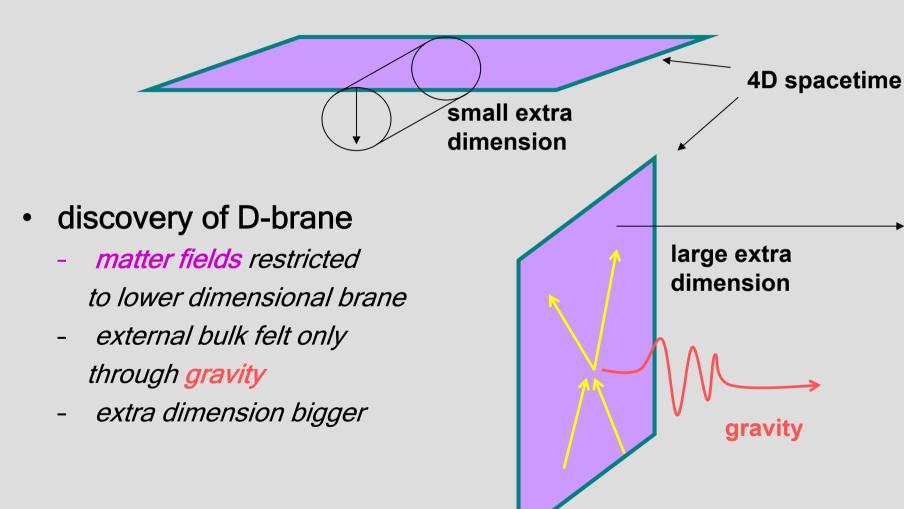
Preserve most symmetries of general relativity

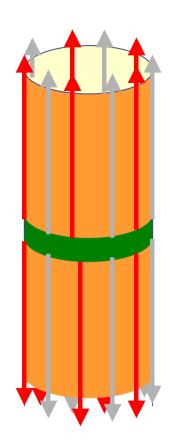
Useful extension of GR (simple but general enough)

$$S = \frac{c^3}{16\pi G} \int \sqrt{-g} \{R - 2(\partial_{\mu}\phi)^2 - V(\phi)\}$$
 *spin 0
 $+ S_m \{\text{matter}, \tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}\}$

Ut DON, I ME 2FF FX I KY-DIWENZIONZ ;

 conventional Kaluza-Klein idea: extra dimensions too small to be seen

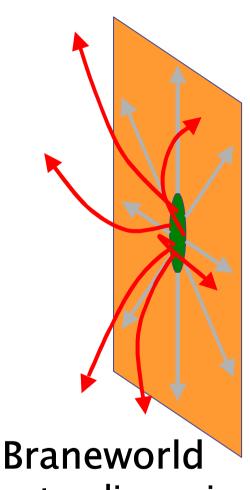




Flux lines of gravity

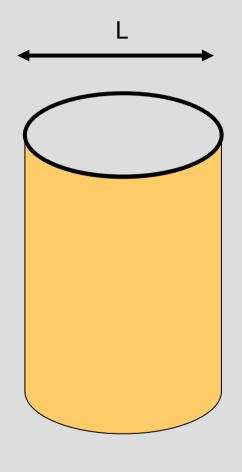
Electric flux line

Kaluza-Klein extra dimensions have SM fields



Braneworld extra dimensions are "dark"

AKGE EXTKA-DIMENSIONS



$$M_4 = (M_{4+n})^{n+2} L^n$$

We can fix n such that $M_{4+n} \sim M_{EW}$.

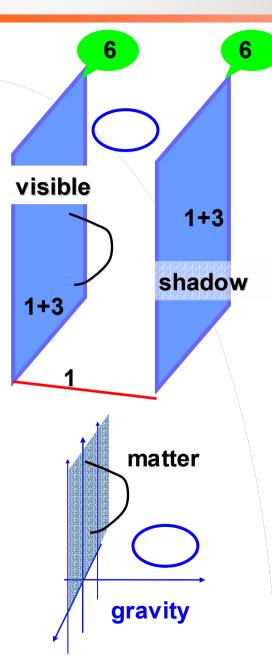
$$L \sim 10^{-17+30/n} \left(\frac{1 \text{ TeV}}{M_{\rm EW}}\right)^{1+2/n}$$

Gravity on small scales.

$$V \propto 1/r^{n+1}$$
 $r < L$

This imposes L<1 mm so that.

N-INEOKY AND NOKAVA-WILLEN MODEL



M theory

1 time + 10 space dimensions

effective 5D braneworld

$$M_5 \sim (M_4^2 / L)^{1/3} << M_4$$

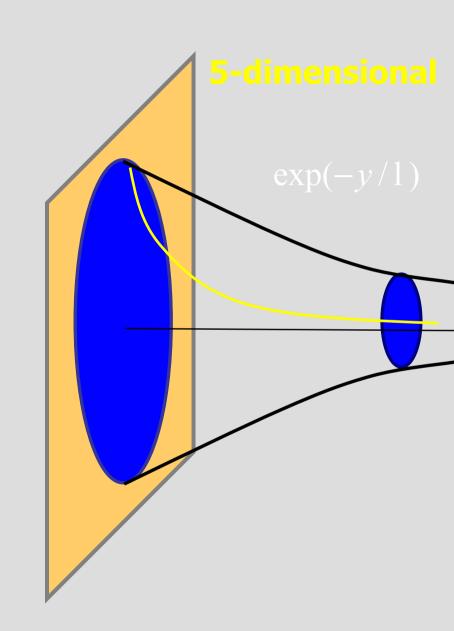
ANVALL-JUNVKUM MUVEL

warped geometry

- extra-dimension shrinks toward infinity
- curvature scale: l

4D GR is recovered on large scales

 $r > 1 \Rightarrow 1 < 0.1$ mm



EVIATION FROM GR IN THE EARLY UNIVERSE

Gravity leakage to the fifth dimension

Gravity becomes effectively five-dimensional

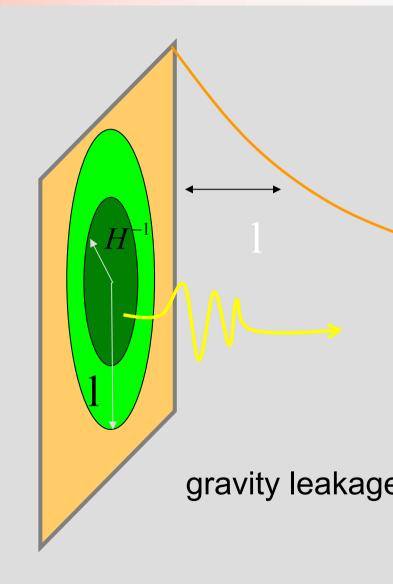
Friedmann equation

$$H^2 = \frac{8\pi G\rho}{3} \left(1 + \frac{\rho}{\lambda} \right) + \frac{\Lambda}{3} - \frac{K}{a^2}$$

In the early universe

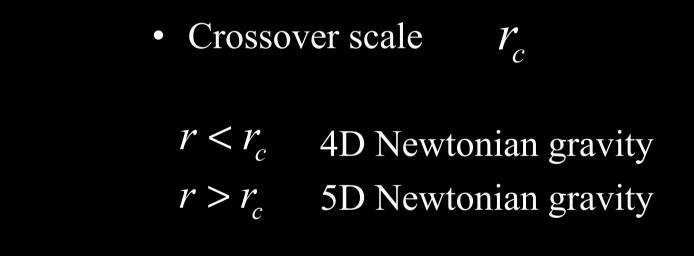
$$H \propto \rho$$

It will modify the phenomenology for inflation



GP MODEL

$$S = \frac{1}{32\pi G r_c} \int d^5 x \sqrt{-\frac{5}{g}} \, g^{(5)} R + \frac{1}{16\pi G} \int d^4 x \sqrt{-g} \left(R + L_m \right)$$



gravity leakage Infinite extra-dimension

EVIATION FROM GR IN THE EARLY UNIVERSE

Friedmann equation

$$H^{2} = \left(\frac{1}{2r_{c}} + \sqrt{\frac{1}{4r_{c}^{2}} + \frac{8\pi G\rho}{3}}\right)^{2} - \frac{K}{a^{2}}$$

Modification in the Infra-red (large distance)

May explain the recent acceleration of the universe

FFECT ON THE FUNDAMENTAL CONSTANTS

In string theory, all dimensionless parameters become VEV of some fields: dynamical.

of some fields: dynamical. e.g.
$$M_4^2={
m e}^{-2\Phi}V_6M_I^8$$
 $g_{YM}^{-2}={
m e}^{-\Phi}V_6M_I^6$ $(+\,c_iM_i)$ $G\propto R^{-D},$ $g_{VM}^{-2}\propto K_i(D)GR^2$

The low energy limit are scalar-tensor theories (dilaton) at tree level.

Dudas (2000)

Loop corrections: need to be understood better couplings are not universal

Couplings are not universal
$$M_4^2=\phi M_H^8, \qquad g_{YM}^{-2}=\phi M_H^6, \qquad \phi=V_6\,\mathrm{e}^{-2\Phi}$$
 $g_{YM}^{-2}=\phi M_H^6-\frac{b_a}{2}(RM_H^2)+\dots$

Phenomenologically: couplings of the quintessence field

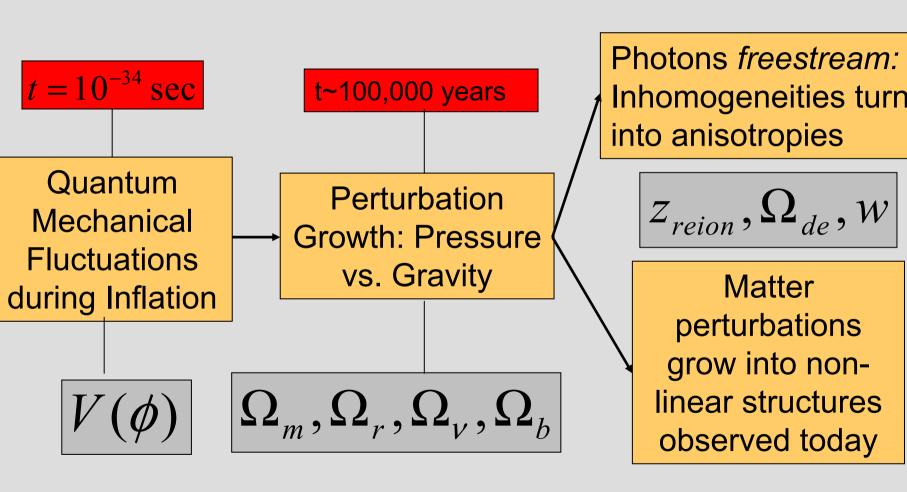
brana modale D. J. J. J. D. C. (2020) 10051

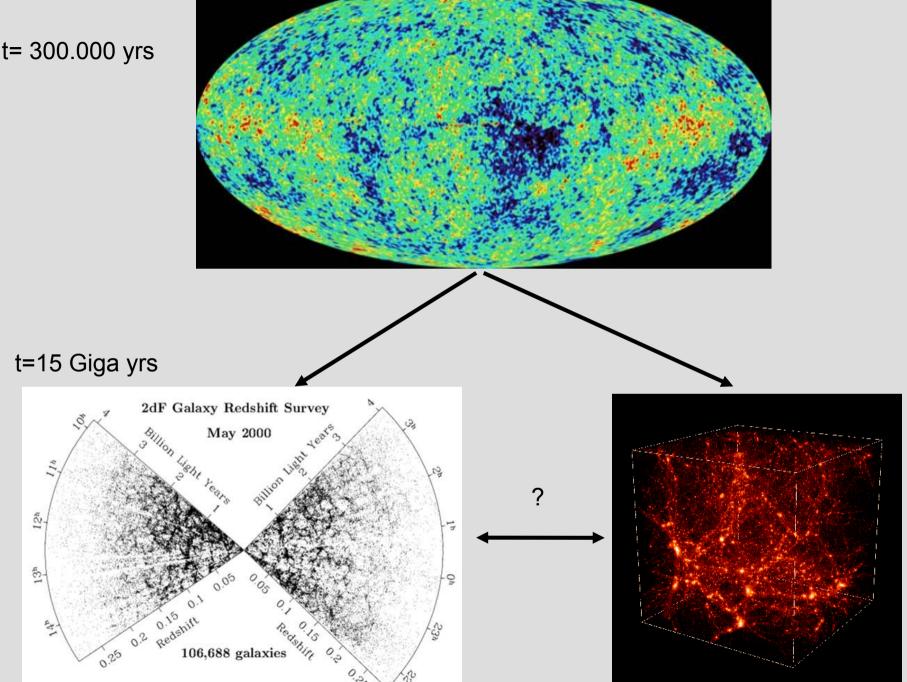
PART VII: CONCLUSIONS

Main topics

- Status of the model
- Open issues and questions

OUEKENI LICIORE OF 21KOCIORE LOKWATION





YCELLENI ZIANDAKD WODEL

The standard model is based on

- relativity
- electromagnetism
- weak interaction
- GUT SUSY
- QFT
- string quantum gravity

Tests

- numerous and successful
- more test of fundamental physics are needed
- what physics behind the parameters

Laboratory

- constraints on extension of the standard models
- Model building

It explains

- the dynamics of the universe
- the origin of LSS

Questions

- dark matter?
- dark energy?
- inflaton?
- constants?

ANDSCAPE PICTURE-THE LIMIT OF EXPLANATION

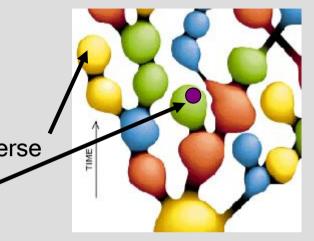
The universe is « machine » producing space

It produces baby-universes

The observable universe is a small patch of a baby-universe

Each baby-universe can have a different physics

(a different vacuum of string theory)



Observable universe < universe < multiverse

Eternal inflation + string theory gives a new picture of the universe

AN WE EXPLAIN THE VALUE OF THE FUNDAMENTAL CIES!

low many numbers:

Standard model of Part. Phys.: 18 + 3 (c, G, h) + 7 (neutrinos) + ... MSSM : 124 parameters

These parameters are **arbitrary** (different values will let the theoretical rchitecture safe).

e.g. a=2/137 will let part. Phys safe but will affect chemistry, biology...

ine tuning among these parameters

e.g. triple alpha...

wo approaches:

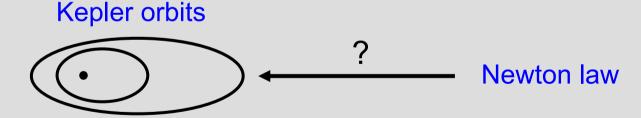
- 1-existence of an attraction mechanism
 - e.g. attraction of scalar-tensor theories toward GR
- 2-anthropic approach

IKING AND CONSTANTS

- 1- fundamental object = string. Characterized by 1 constant (λ)
- 2 4 interactions are unified: 1 coupling constant (g_s)
- 3- Witten (87): all parameters are dynamical variables

But

Theory is defined in D= 9 + 1
One needs to determine the vacuum corresponding to our universe



Equations depend one 1 parameter but solutions on much more.

D=3+1+1: one more parameter (radius of compactification)

D=3+6+1: several hundreds to describe topology and size!

Idea: exploring the landscape of vacua.

HE ANTHKOPIC WAT

Selectionist approach:

goal is not to derived the value of the fundamental parameters but Show that are necessary conditions for some physical phenomena to exist

$$P \Rightarrow C \Longleftrightarrow !C \Rightarrow !P$$

not a finalist approach does not assume that P has to exist!

P = existence of human life (anthropic principle) existence of some life (biotic principle) existence of carbon (carbonic principle)

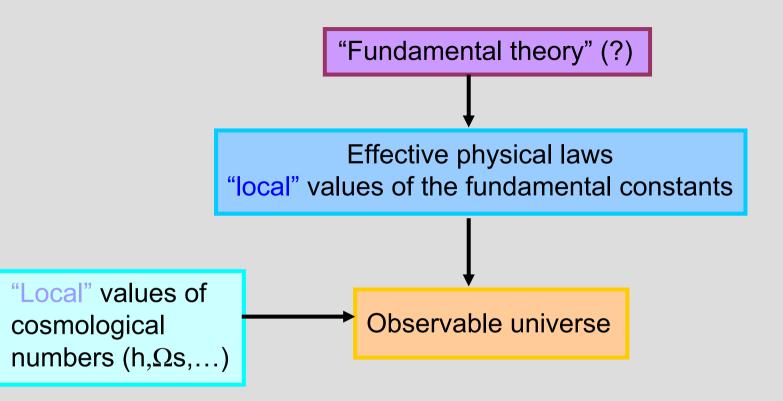
But to apply this principle

- 1- one needs to conceive that other (part of the) Universe may not have the same physics
- 2- find a mechanism to explore the set of possible universes (not trivial: does our universe belong to this set? Fine tuning for discretisation?)

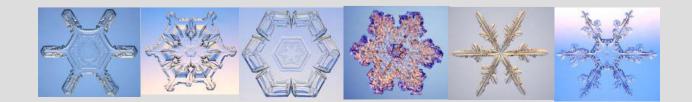
Ideas:

Eternal inflation, landscape...

AN WE EXPLAIN THE VALUE OF THE FUNDAMENTAL CIES!



Which numerical coincidences are consequences of the structure of the fundamental theory and which are accidents?



AND I'LL STOP HEKE

Thank you!