

## 7) Standard Model and CP violation

Strong interaction

gluons

Electromagnetic interaction

photons

**Weak interaction**

neutral current

charged current:  $W^\pm$

Up type quark

spinor field

$$Q = 2/3$$

$$\mathbf{U} = \begin{pmatrix} u \\ c \\ t \end{pmatrix}$$

Down type quark

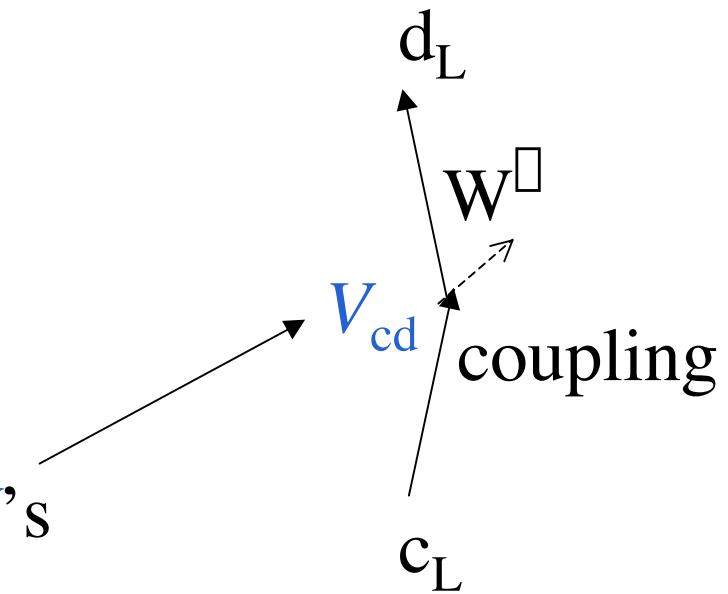
spinor field

$$Q = -1/3$$

$$\mathbf{D} = \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

example

there are  $3 \times 3 = 9$   $V$ 's



## “simplified” way...

- 1) Since CPT is respected,  $\mathcal{CP}$  is like  $\mathcal{T}$
- 2) T transformation is like making complex conjugation:

$$e^{\square iEt} \quad \square \quad T \quad \square \quad e^{iEt}$$

- 3) T transformation to the Hamiltonian operator  $H$

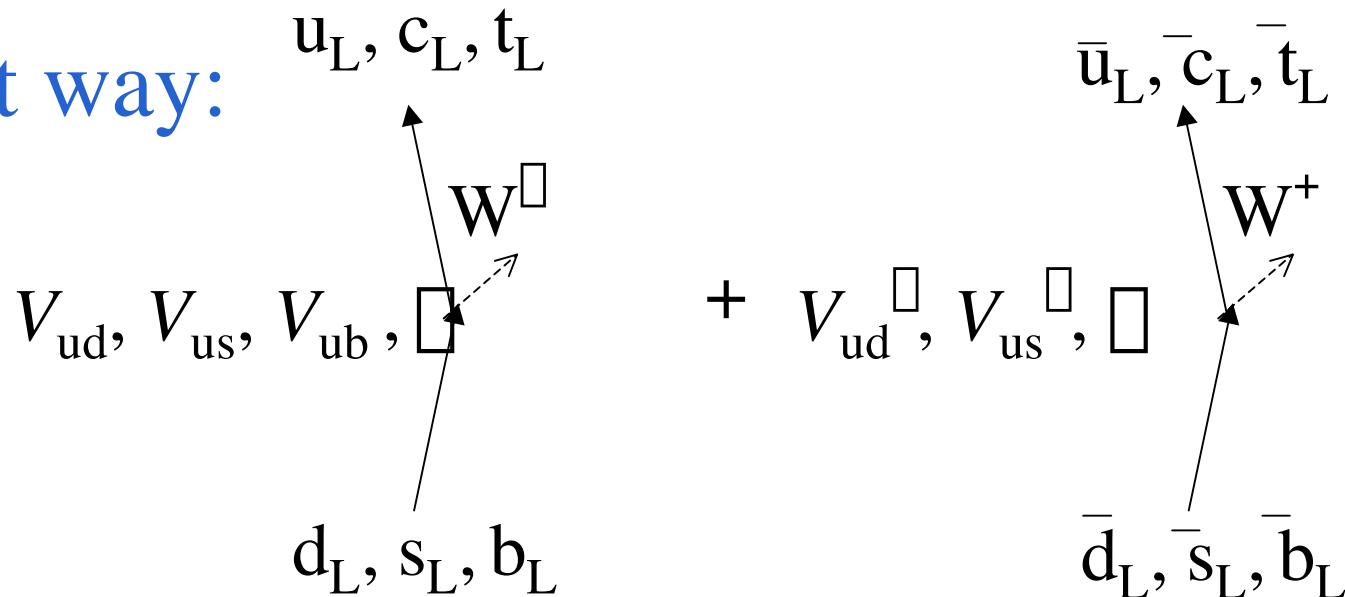
$$H \quad \square \quad T \quad \square \quad H^\square$$

if  $H \neq H^\square$ ,  $\mathcal{T}$  i.e.  $\mathcal{CP}$

Complex coupling  $V^\square \neq V$  generates CP violation,

$V = \{V_{ij}\}$ : Unitary matrix

Correct way:



$$L \propto V_{ij} \bar{U}_i \square(1 \square \square) D_j W_0^\dagger + V_{ij}^\square \bar{D}_i \square(1 \square \square) U_j W_0$$

↑  
CP conjugation

$$L_{\text{CP}} \propto V_{ij} \bar{D}_i \square(1 \square \square) U_j W_0 + V_{ij}^\square \bar{U}_i \square(1 \square \square) D_j W_0^\dagger$$

If  $V_{ij}^\square = V_{ij}$   $L = L_{\text{CP}}$ : i.e. CP conservation

Let us look at now:  $V_{ij} \bar{U}_i \Gamma(1\bar{1}\bar{5}) D_j$

One family

$$V \begin{array}{l} 1 \text{ free phase} \\ 1 \text{ free modula} \end{array} = |V| e^{i\phi}$$

$$|V| e^{i\phi} \bar{u} \Gamma(1\bar{1}\bar{5}) d \longrightarrow |V| \bar{u} \Gamma(1\bar{1}\bar{5}) d$$

Changing u quark phase:  $\bar{u} \rightarrow \bar{u} e^{i\phi}$

Unitarity:  $V^\dagger V = VV^\dagger = E$  (one constraint)

$$|V|^2 = 1$$

0 free phase  
0 free modula

NO CP

Two families  $V = \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix}$

4 free phase  
4 free moduli (or rotation angles)

- 1) The phase of  $V_{ij}$  can be absorbed by adjusting the phase differences between i- and j- quark

$$4 \text{ quarks} = 3 \text{ phase differences}$$

$$4 - 3 = 1 \text{ phase left}$$

- 2) Unitarity  $V^\dagger V = VV^\dagger = E$ : four constraints:

1 off-diagonal constraint for the phase

$$1 - 1 = 0 \text{ phase left}$$

$$\begin{array}{|c|c|} \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array} = \begin{array}{|c|c|} \hline 1 & 0 \\ \hline 0 & 1 \\ \hline \end{array}$$

three constraint for the rest

$$4 - 3 = 1 \text{ rotation angle left}$$

**$V$  is real, i.e. no  $\cancel{CP}$ .**

## Explicit demonstration

$$|V_{ud}|e^{i\beta_{ud}} \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{us}|e^{i\beta_{us}} \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s \\ + |V_{cd}|e^{i\beta_{cd}} \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{cs}|e^{i\beta_{cs}} \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s$$

$\bar{u} \square \quad u e^{i\beta_{ud}}$

$$|V_{ud}| \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{us}|e^{i(\beta_{us}-\beta_{ud})} \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s \\ + |V_{cd}|e^{i\beta_{cd}} \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{cs}|e^{i\beta_{cs}-\beta_{us}+\beta_{ud}} \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s$$

$s \square \quad s e^{i(\beta_{us}-\beta_{ud})}, \quad c \square \quad c e^{i(\beta_{cs}-\beta_{us}+\beta_{ud})}$

$$|V_{ud}| \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{us}| \bar{u} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s \\ + |V_{cd}|e^{i\beta_{cd}} \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) d + |V_{cs}| \bar{c} \begin{array}{c} \square \\ \square \end{array}(1 \square \square) s$$

Out of four quark, three quark phases can be adjusted:  
4 free phase  $\square$  1 free phase

Unitarity:  $V^\dagger V = VV^\dagger = E$  (4 constraints)

$$\begin{pmatrix} V_{ud}^* & V_{cd}^* \\ V_{us}^* & V_{cs}^* \end{pmatrix} \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix}$$

$$V_{ud}^* V_{us} + V_{cd}^* V_{cs} = 0 \quad |V_{ud}| |V_{us}| + |V_{cd}| |V_{cs}| e^{i\alpha} = 0$$

0 free phase:  $\alpha = 0$

$$|V_{ud}| |V_{cd}| |V_{us}| |V_{cs}| = 0$$

$$|V_{ud}|^2 + |V_{cd}|^2 = 1, |V_{us}|^2 + |V_{cs}|^2 = 1$$

1 free modula or rotation angle

$$|V_{11}| = \cos \theta, |V_{22}| = \cos \theta, |V_{12}| = \sin \theta, |V_{21}| = \sin \theta$$

$$V = \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$$

One rotation angle without phase:  $\theta$  NO CP  
(Cabibbo angle)

Three families

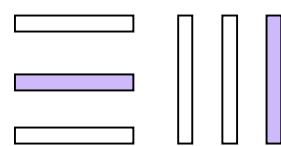
$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

9 free phase

9 free moduli (or rotation angles)

Out of **six** quark, **five** quark phases can be adjusted:

9 free phase  $\square$  4 free phase



1 0 0  
0 1 0  
0 0 1

Out of **nine** unitarity constraints, **three** are for the phases

4 free phase  $\square$  1 free phase

the rest (**six**) are for the rotation angles

9 free rotation angles  $\square$  3 free rotation angles

Three rotation angles with one phase:

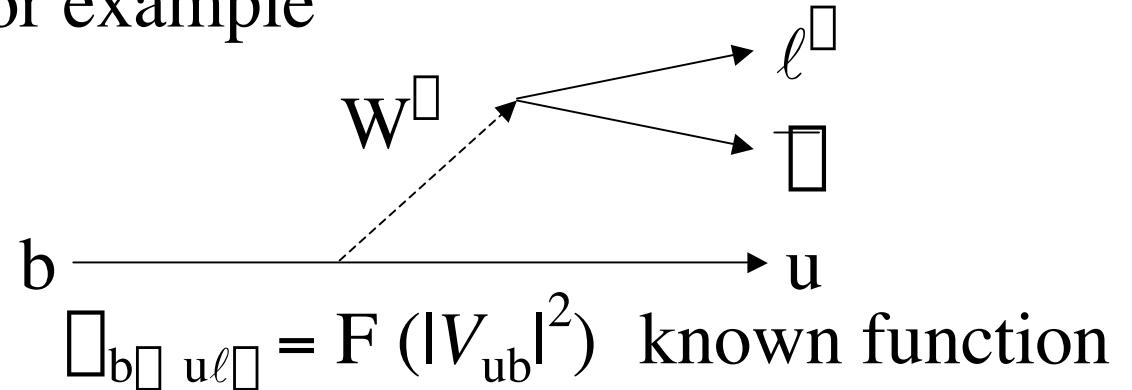
$\square$   **$\mathcal{CP}$  can be generated**

Electroweak theory with 3 families can naturally accommodate CP violation in the charged current induced interactions through the complex Cabibbo-Kobayashi-Maskawa quark mixing matrix  $V$ , with 4 parameters.

To be more qualitative  $\rightarrow$  determination of the CKM matrix

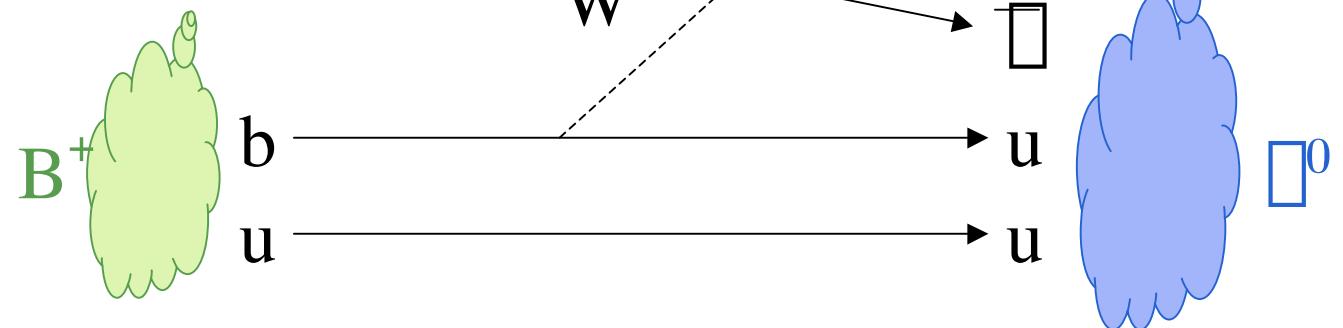
$ V_{ud} $	$ V_{us} $	$ V_{ub} $
$ V_{cd} $	$ V_{cs} $	$ V_{cb} $
$V_{td}$	$V_{ts}$	$ V_{tb} $

From the “tree” level quark decays  
for example



What you observe is “hadrons”. difficult to relate!

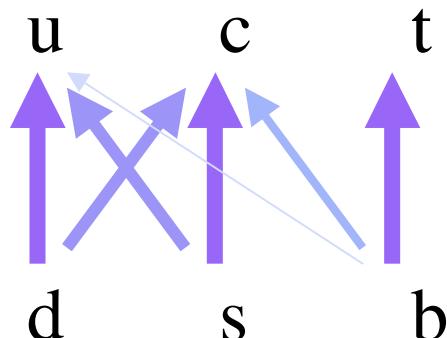
non-perturbative  
strong interactions



The “tree” level measurements give...

$$|V| = \begin{pmatrix} 0.9734 \pm 0.0008 & 0.2196 \pm 0.0023 & 0.0036 \pm 0.0007 \\ 0.224 \pm 0.016 & 0.996 \pm 0.13 & 0.0412 \pm 0.0020 \\ \square & \square & 0.99 \pm 0.29 \end{pmatrix}$$

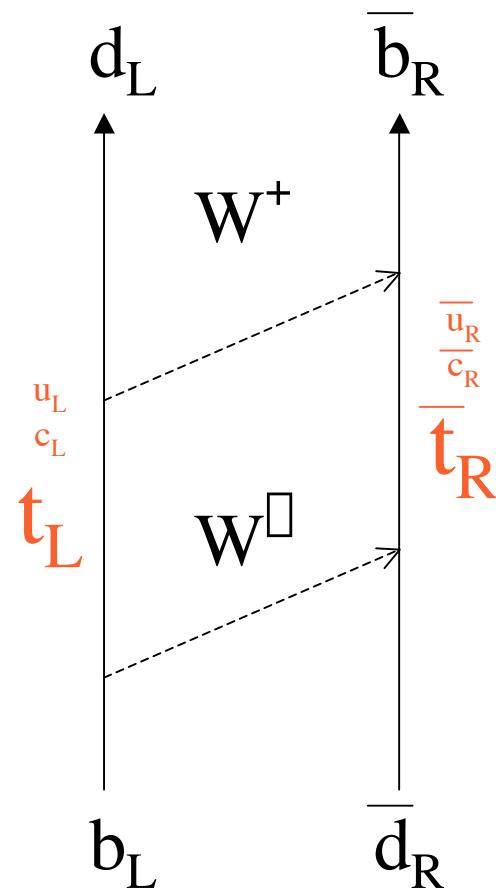
This  $2 \times 2$  sub-matrix (first two generation) is almost unitary.



Errors are all theoretical,  
except  $|V_{tb}|$ .

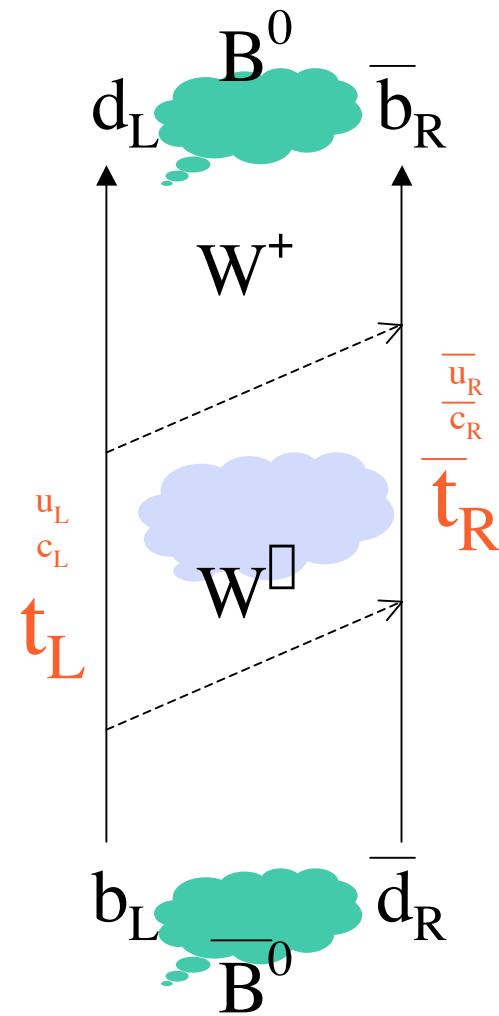
Access to the top quark through virtual process  
 $(b\bar{d}) \square (\bar{b}d)$  oscillation frequency is given by the “box” diagrams

$\square = f(|V_{td}|^2 |V_{tb}|^2)$ : known function



Again non-perturbative strong interactions make them difficult to relate!

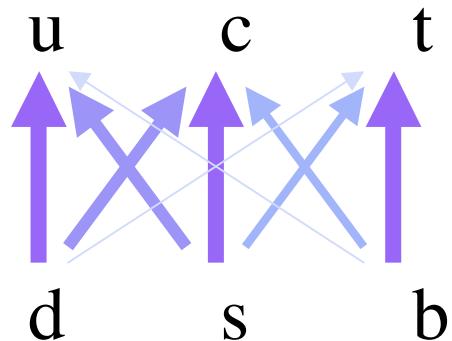
What we measure is  $\bar{B}^0 \square B^0$  oscillation frequency



measurement with “trees”,  
measurements with “box”  
and assuming unitarity

$$|V| = \begin{pmatrix} 0.9734 \pm 0.0008 & 0.2196 \pm 0.0023 & 0.0036 \pm 0.0007 \\ 0.224 \pm 0.016 & 0.996 \pm 0.13 & 0.0412 \pm 0.0020 \\ 0.0083 \pm 0.0016 & & \sim 1 \end{pmatrix}$$

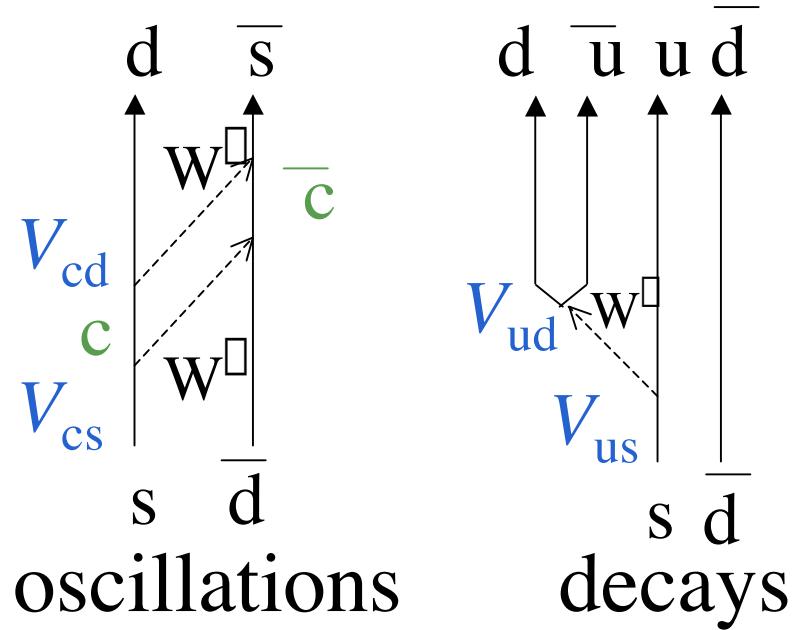
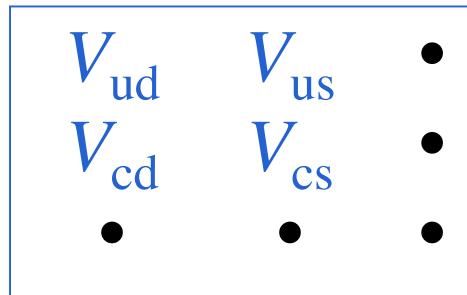
Interesting pattern:



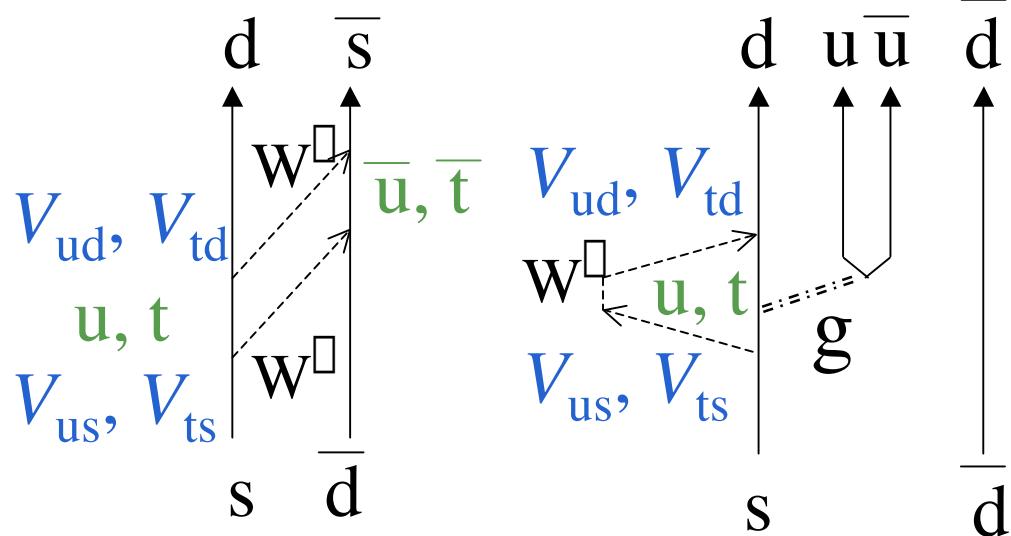
?

Waiting for measurements  
by CDF and D0

Dominant processes  
in  $K^0$  system involve:



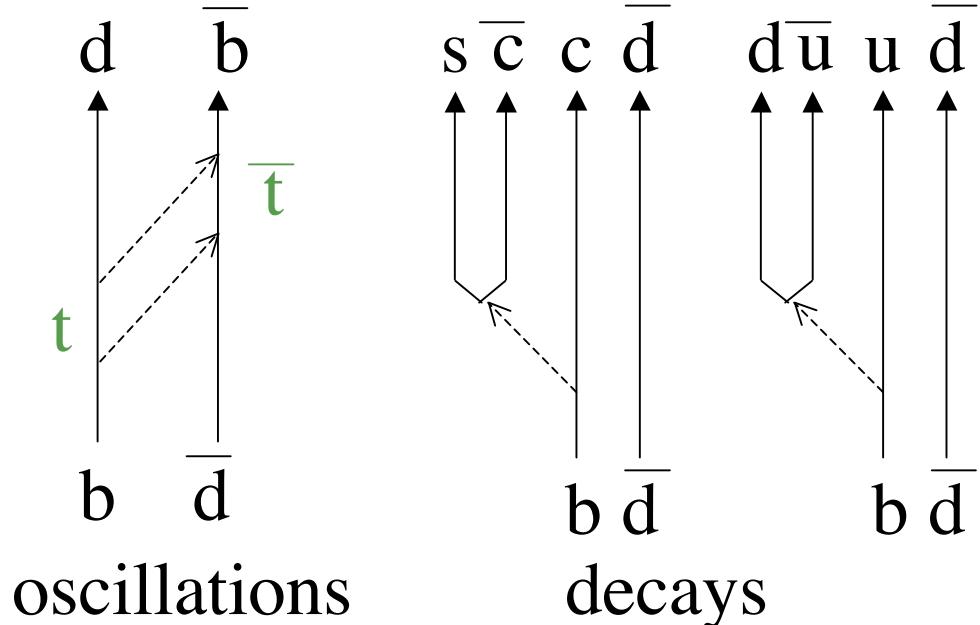
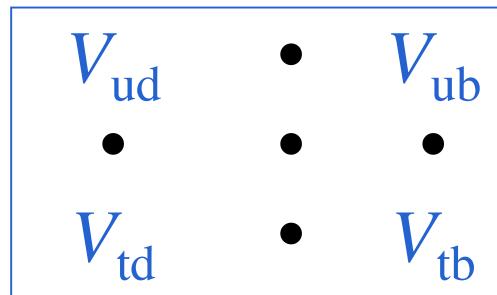
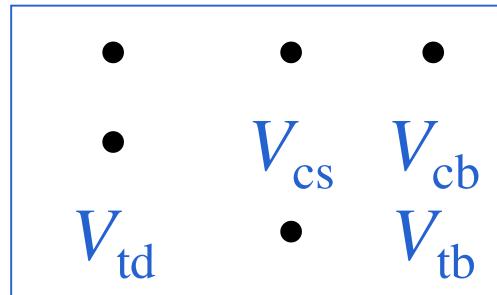
Almost unitary part of  $V \square$  No CP violation.



Much smaller contributions  
to oscillations and decay:  
all **three families** of quarks.

CP violation:  $\text{CP}(K_L) \neq 1$ ,  $|V_{cb}| \neq |V_{cb}^0|$ , but small effect!

Dominant processes in  $B^0$  system involve:



Non unitary part of  $V$   
CP violation is expected to be large.

$$B^0 = (d\bar{b}), \bar{B}^0 = (\bar{d}b)$$

$$\begin{array}{c|c|c} \square & \square & 0 \\ \square m_B & \square 100 & \square m_K \end{array}$$

$$B_{\text{light}} \square J/\square K_S$$

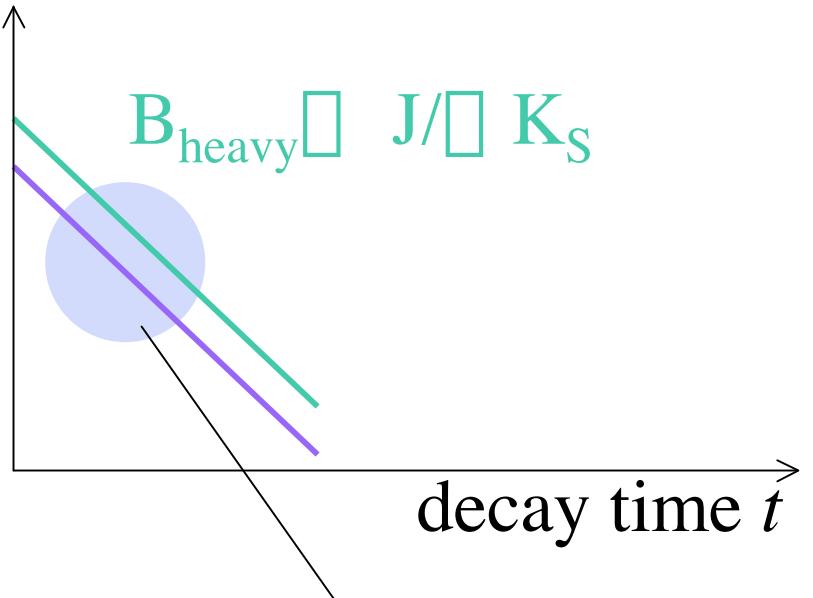
$B \square J/\square K_S$  decays

$$\text{CP}(J/\square K_S) = \square 1$$

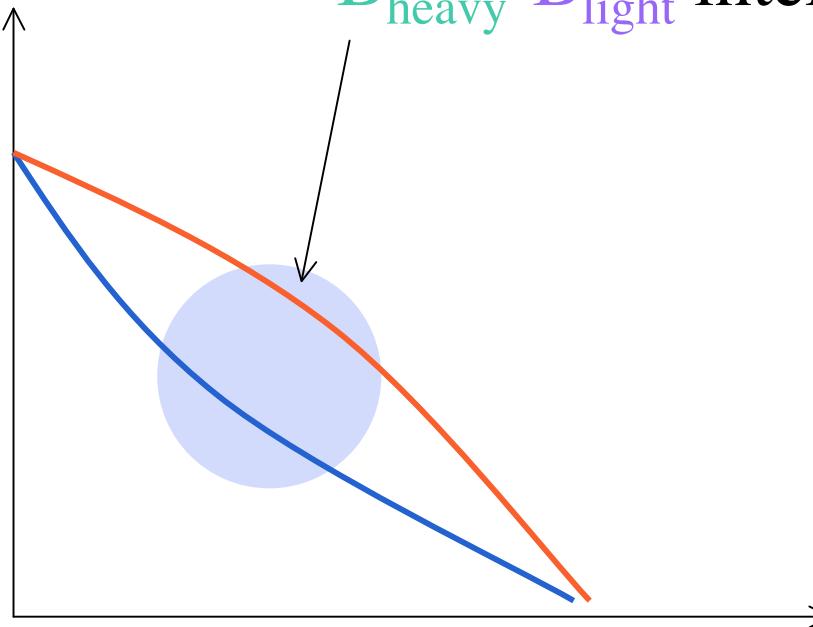
$$B^0 \square J/\square K_S$$

$$\bar{B}^0 \square J/\square K_S$$

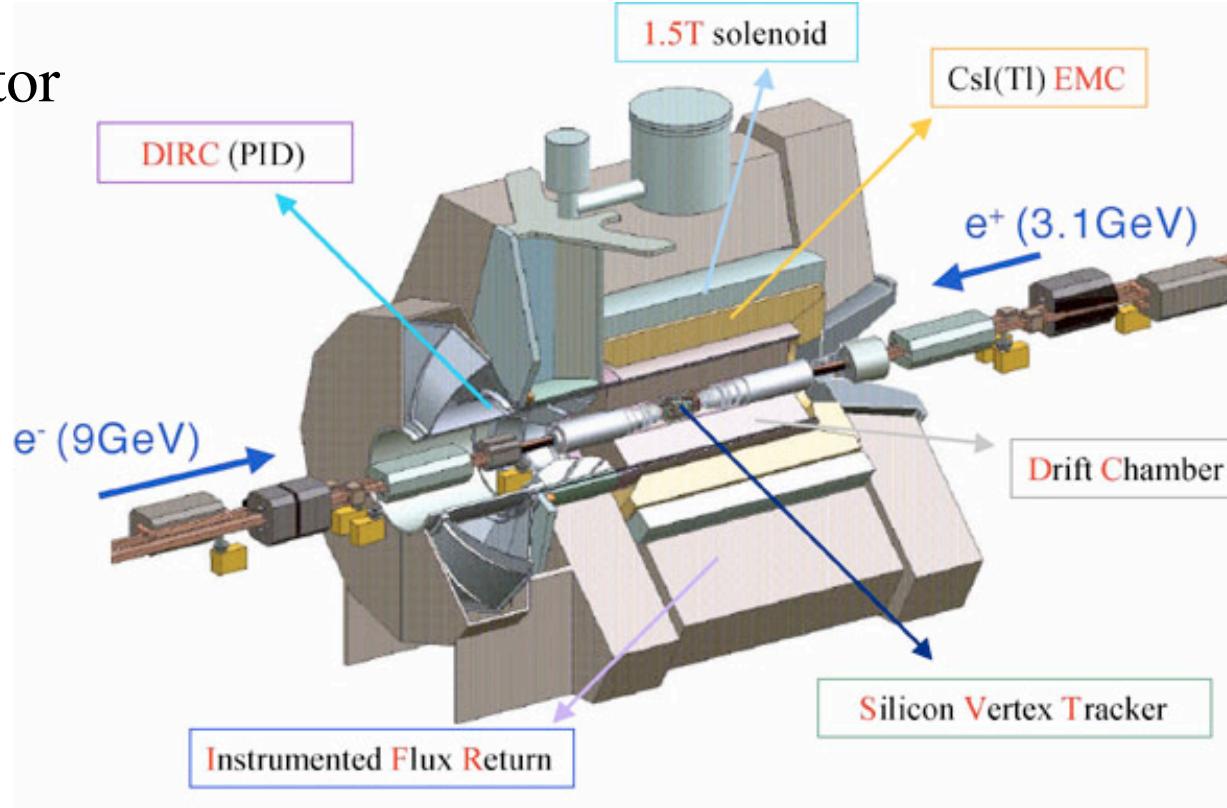
$$dn/dt$$



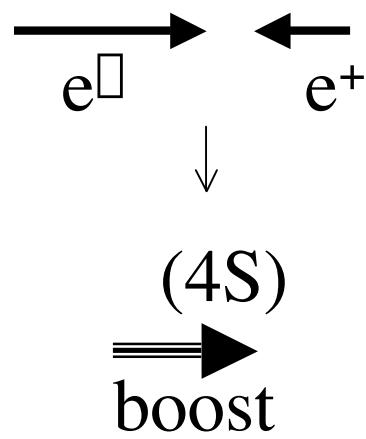
$B_{\text{heavy}} - B_{\text{light}}$  interference



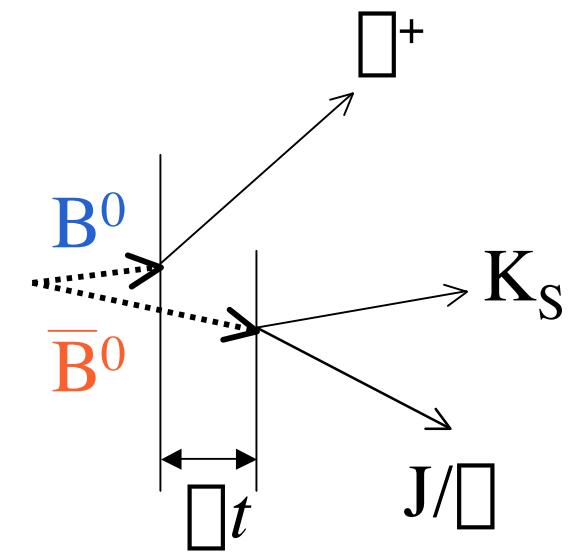
# BABAR detector



in the lab. frame

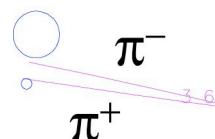
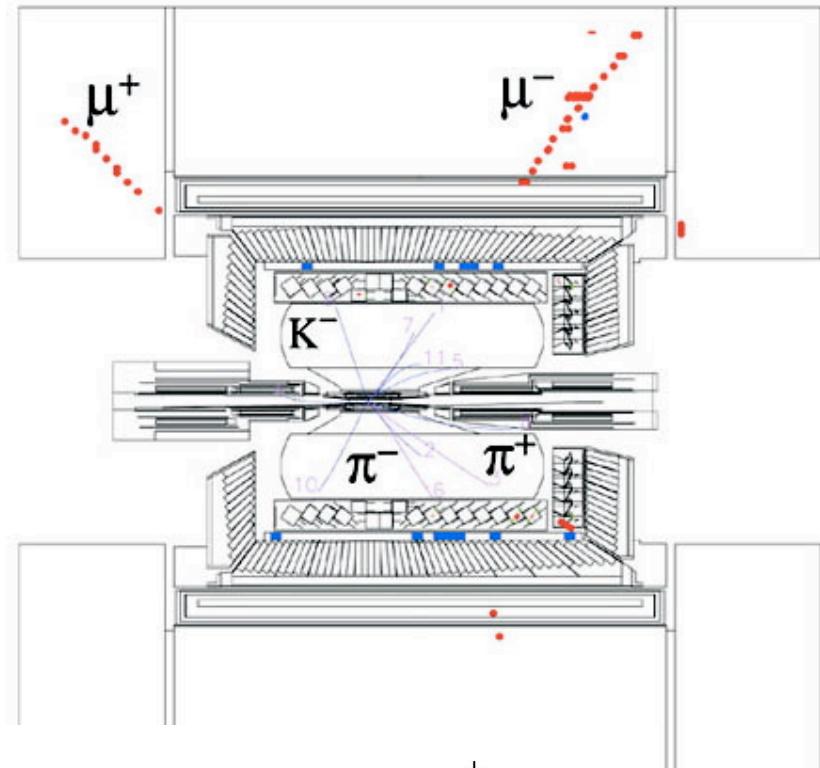
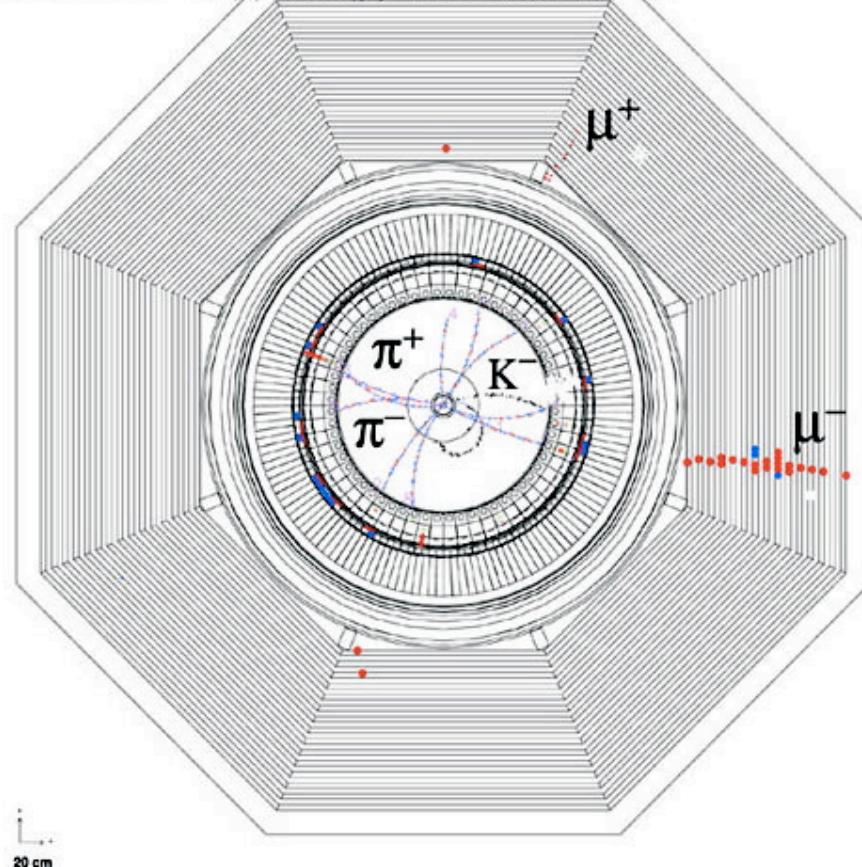


boosted  
 $\bar{B}^0 B^0$  system  
 (remains always  $\bar{B}^0 B^0$ )

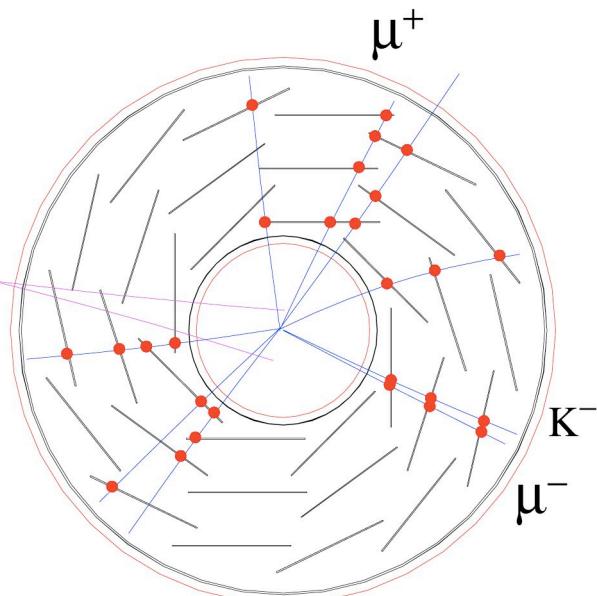


**BELLE**

Exp: 5 Run: 272 Form: 5 Event: 10889  
Ehei: 8.00 Eler: 3.50 Tue Nov 16 23z12z08 1999  
TrgID: 0.0 DetVer: 0. MagID: 0. BField: 1.50. DspVer: 5.04  
Ptot(ch): 10.3 Etot(qm): 0.0 SVD-M: 0 CDC-M: 0 KLM-M: 0

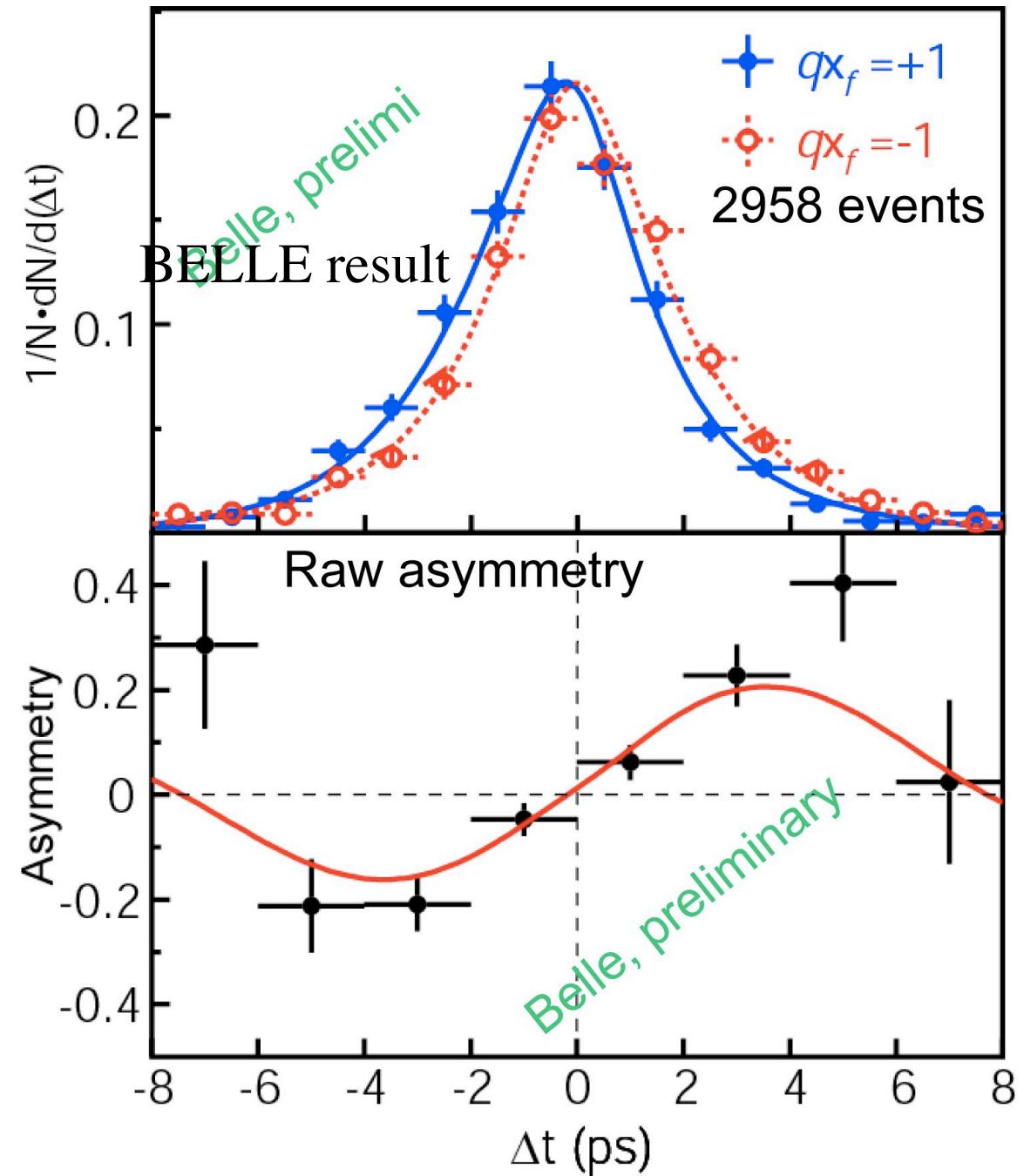


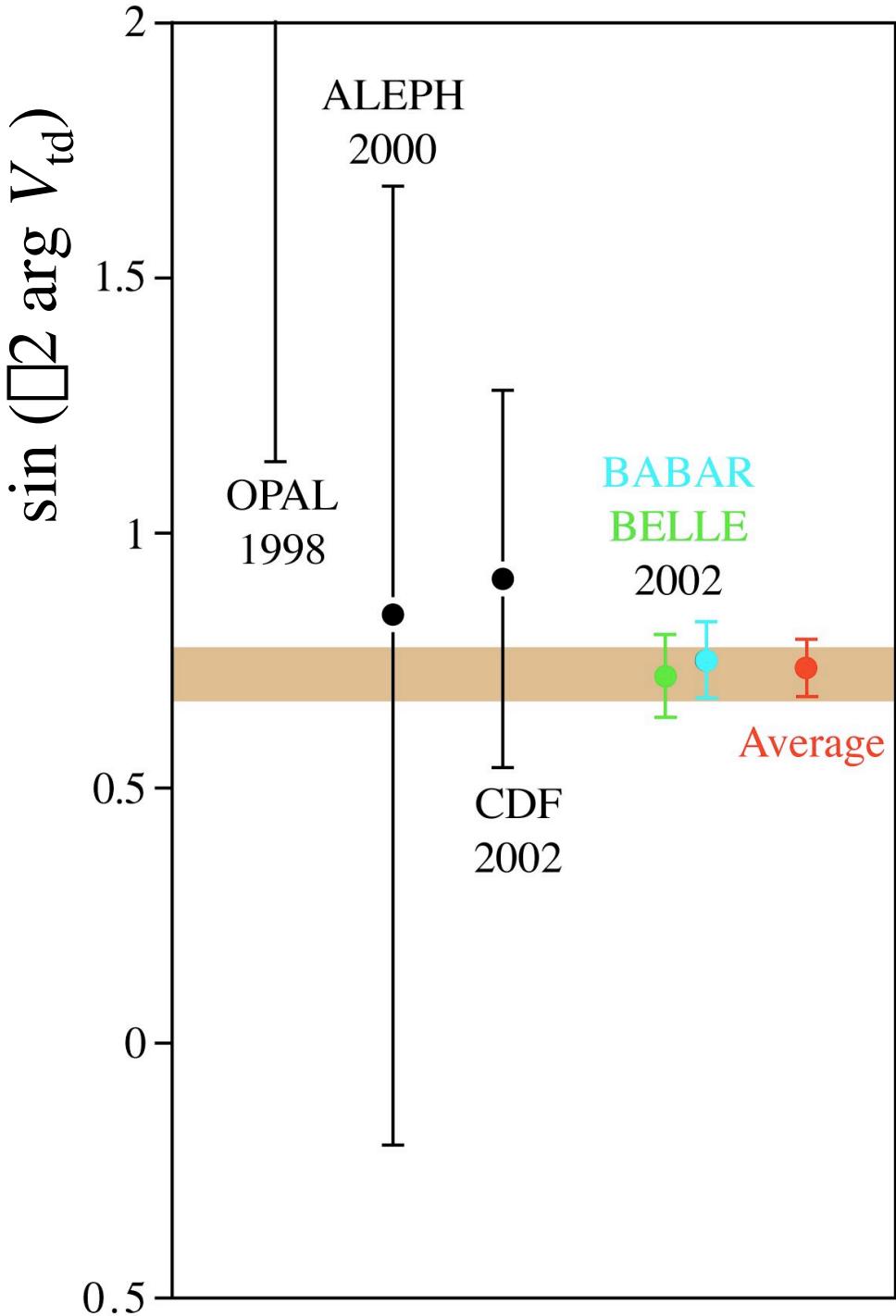
$$\begin{array}{c} \overline{B^0} \rightarrow J/\psi(\ell^+\ell^-) K_S(\ell^+\ell^-) \\ \overline{B^0} \rightarrow K\ell^+ + X \end{array}$$



$B^0$   $J/\psi$   $K_S$   
 $\bar{B}^0$   $J/\psi$   $K_S$

CP violation!





$\arg V_{\text{td}}$  can be extracted  
from CP asymmetry in  
 $B \rightarrow J/\psi K_S$  decays

Standard Model  
prediction

Good agreement  
with the Standard  
Model prediction.

Why do we still worry  
about CP violation?