

# Violation of Particle Anti-particle Symmetry

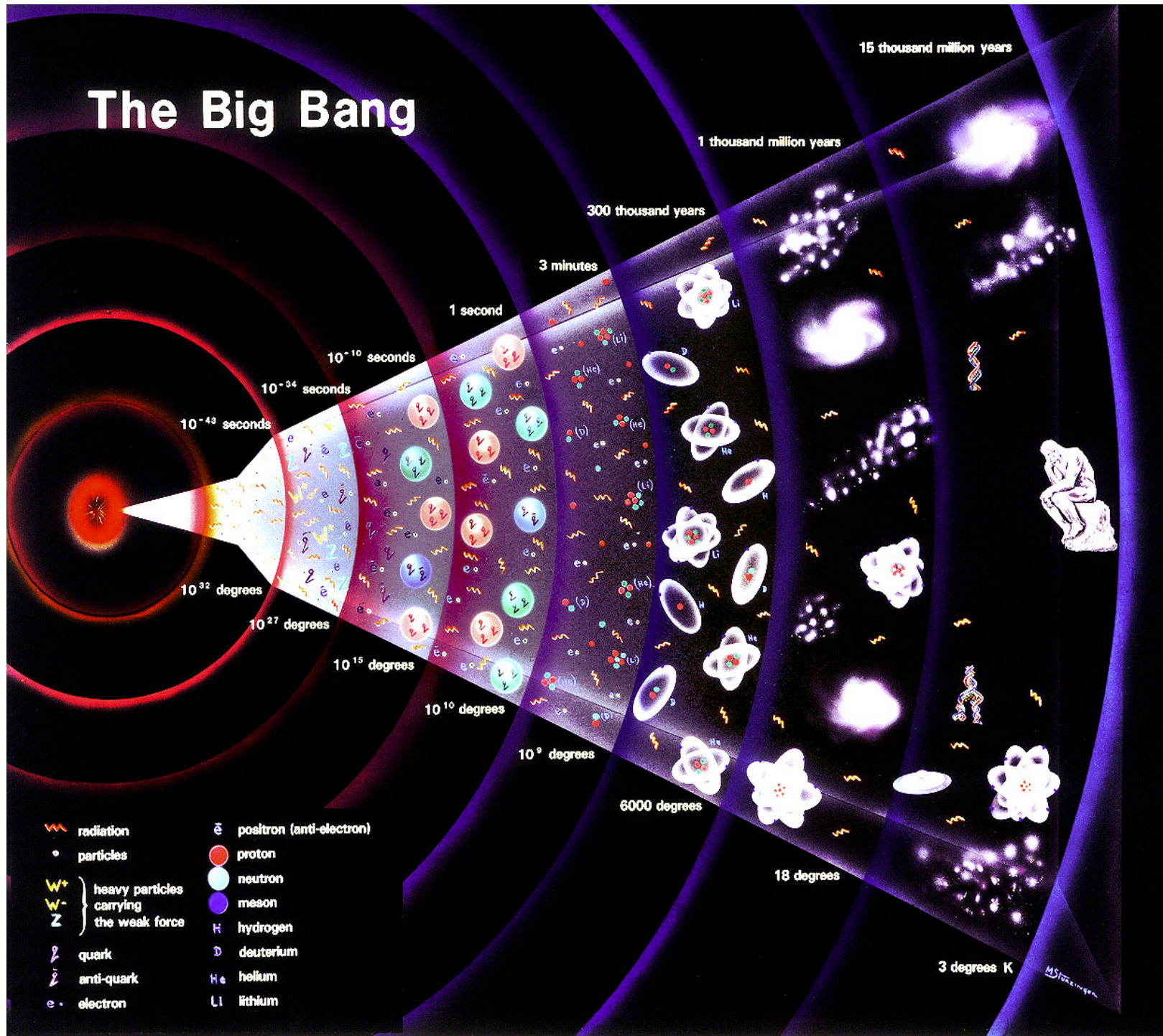
CERN Summer Student Lectures

1, 2 and 5 August 2002

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CERN and University of Lausanne

# The Big Bang



# Contents of the Lecture



I

- 1) Concept of symmetries
- 2) P, T and C transformation
- 3) Conservation of symmetries
- 4) CP violation in the charged kaon system



II

- 5) CP violation in the neutral kaon system
- 6) Kaon interferometer



III

- 7) Standard Model and CP violation (K, B)



IV

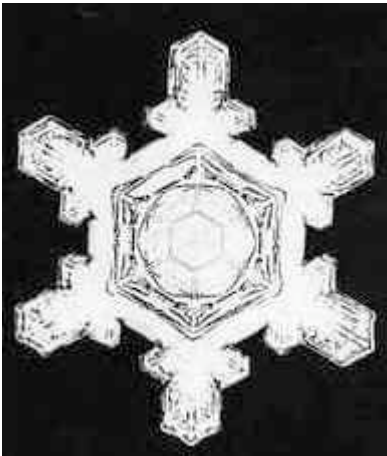
- 8) Baryogenesis and CP violation
- 9) Next experimental steps



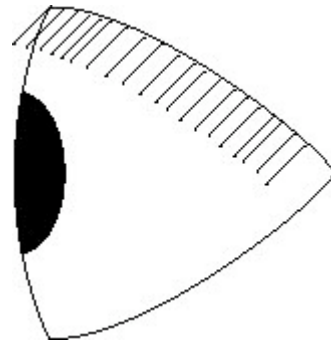
# 1) Concept of symmetries



# Symmetry



Nature

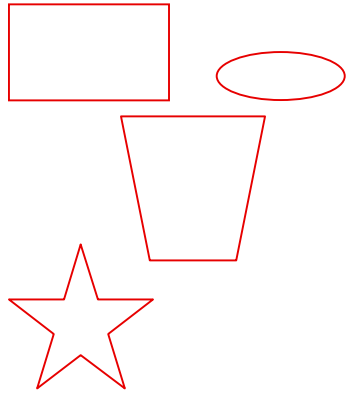



Observation

- regular pattern
- symmetry
- ...

Recognition

# Creating arts



- regular pattern
  - symmetry
  - ...
- 



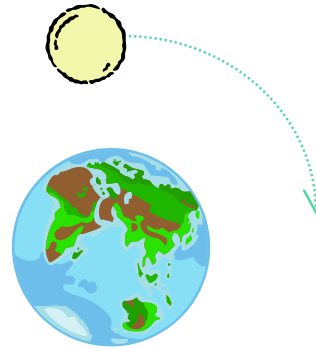
Concept

Realisation

# Natural Science

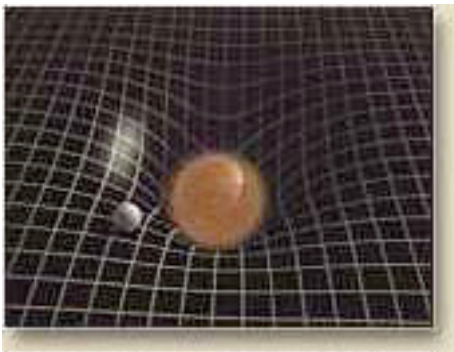


Observation



$$f = G \frac{m_1 m_2}{r^2}$$

Generalising and making physical laws



$$g \square \square g \square \square$$

Extracting more abstract concept



Postulating and predicting phenomena



Observation

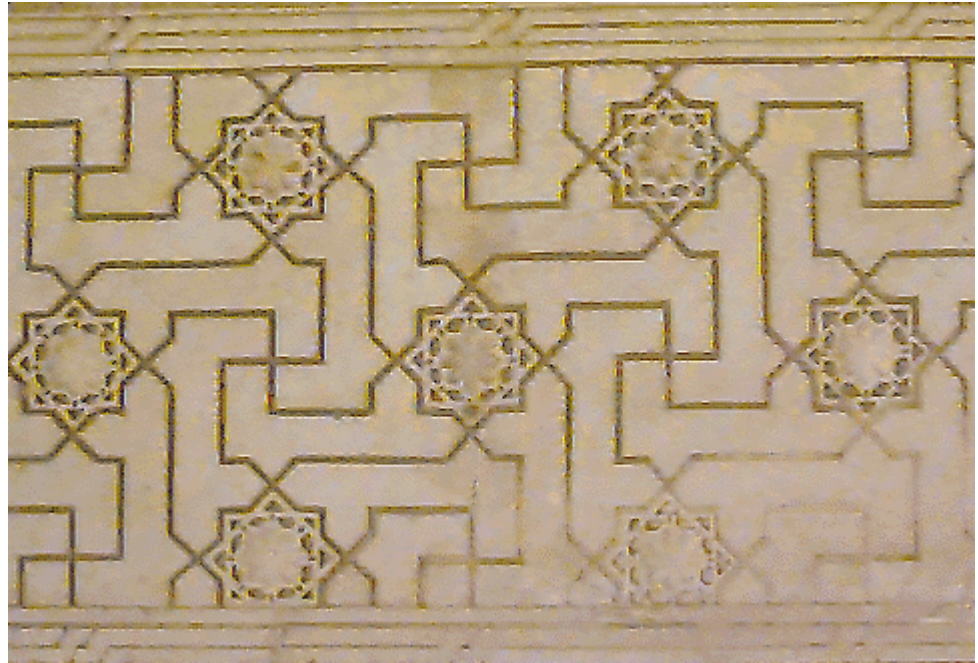




Other examples  
of symmetry







and patterns



Some asymmetry makes... more  
dynamic





more beautiful



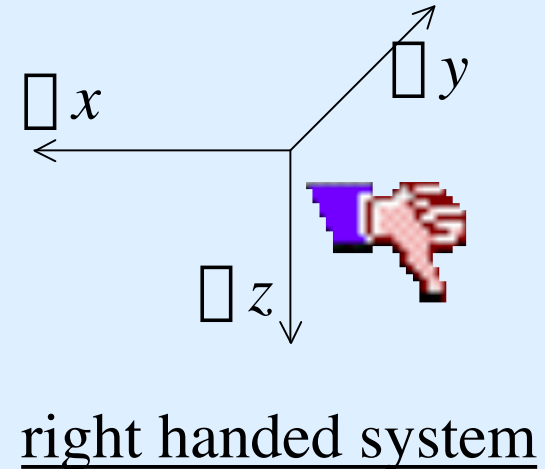
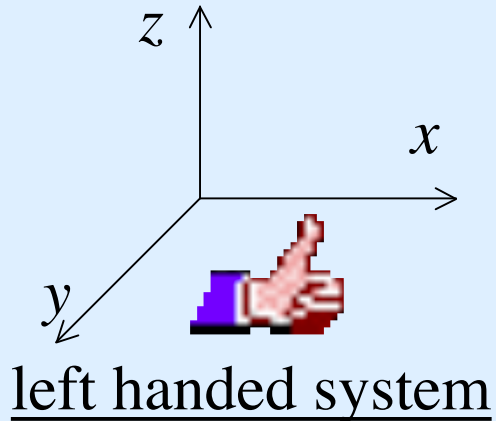
if not too much...?



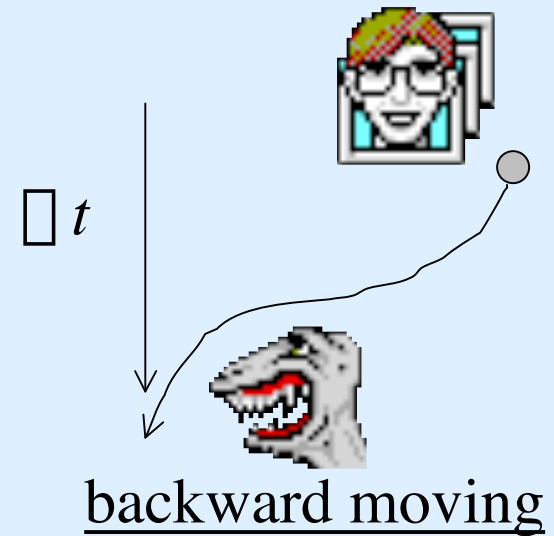
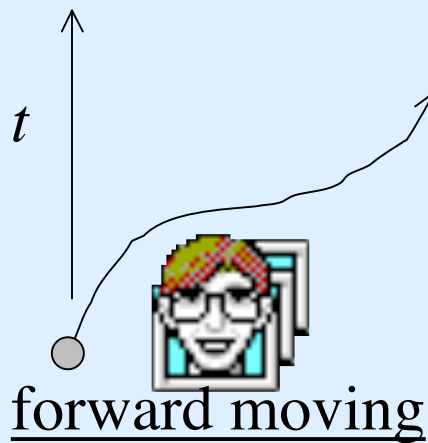


## 2) P, T and C transformation

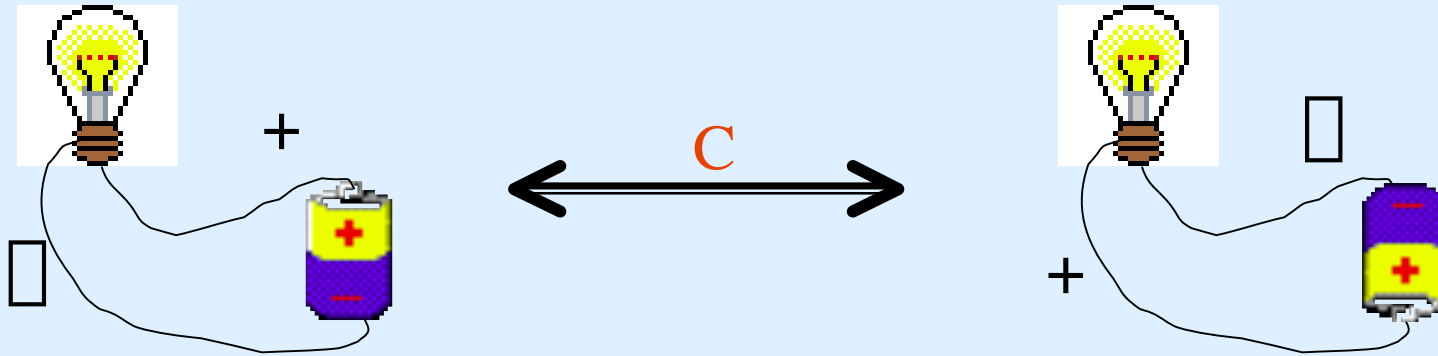
P: parity or space reflection



T: time reversal



## C: charge conjugation



In particle physics reversing internal quantum numbers  
charged states

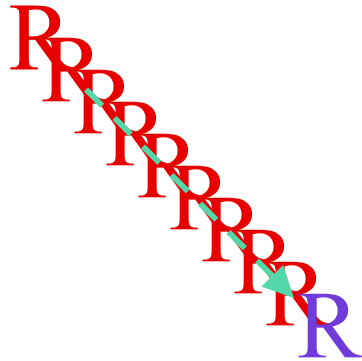
$e^-$ (electron)	$\square$	$e^+$ (positron)
$p$ (proton)	$\square$	$\bar{p}$ (anti proton)
$\pi^+$ (positive pion)	$\square$	$\pi^-$ (negative pion)
$u$ (u quark)	$\square$	$\bar{u}$ (anti u quark)

neutral states

$n$ (neutron)	$\square$	$\bar{n}$ (anti neutron)
$K^0$ (k-zero meson)	$\square$	$\bar{K}^0$ (anti k-zero meson)
$\pi^0$ (neutral pion)	$\square$	$\pi^0$ (neutral pion)

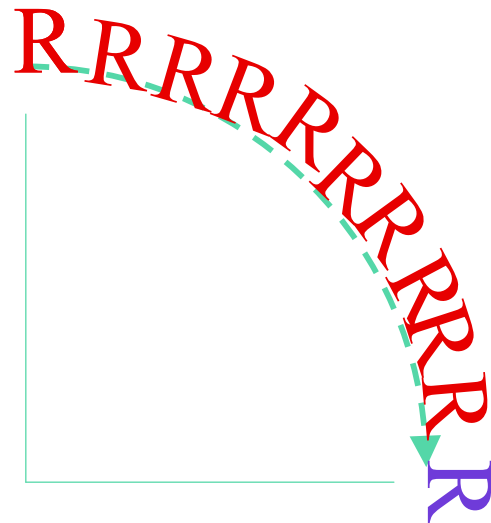
# Discrete and continuous transformations

Translation



continuous

Rotation



continuous

Reflection  
(parity)

R



B

discrete

### 3) Conservation of symmetries

If no difference seen between

“this world” and “space reflected world”



We say:

- parity is conserved,
- P symmetry is conserved,
- world is invariant under P transformation
- etc.

example



More “**professional**” description,

$\hat{H}$  Hamiltonian operator describing the world

$\hat{P}$  Parity transformation operator

$\hat{P}^\dagger \hat{H} \hat{P} = \hat{H}^P$  parity transformation of Hamiltonian

If  $\hat{H}^P \neq \hat{H}$

Parity violation, Parity non-conservation etc. etc.

C and CP



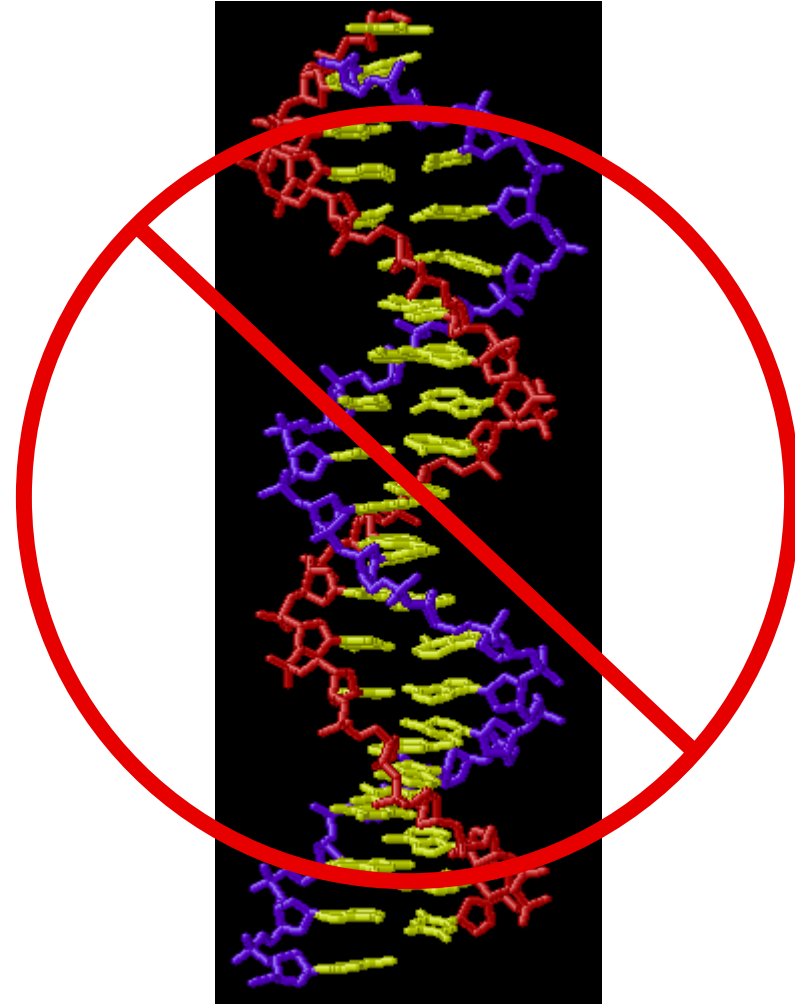
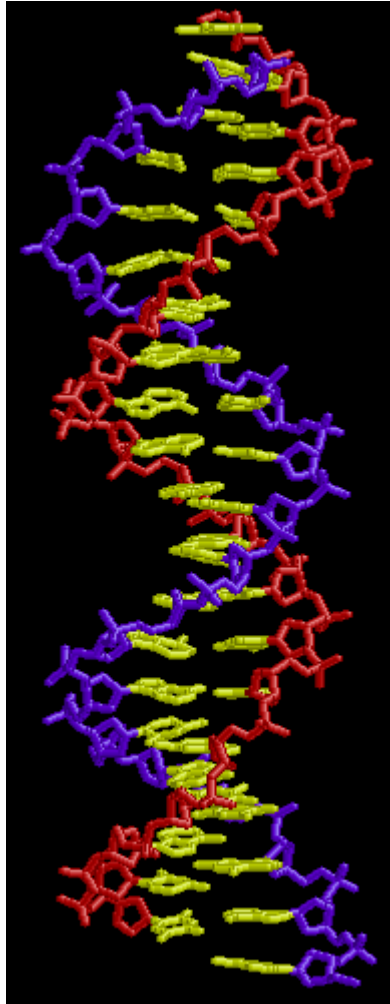


# Violation of Parity



World  $\neq$  Mirror World  
(parity violation)

# Even more with DNA



Parity is fully violated.

to formula



A similar terminology applies to C and T.

Strong and electromagnetic interactions conserve:  
flavour quantum numbers,  
C, P, T, CP, CT, PT and CPT

Particle physics example:

$\pi^0 \rightarrow \pi\pi$  but not  $\pi\pi\pi$

$$\begin{aligned} \pi^0 &= (u\bar{u} + d\bar{d})_{L=0, S=0} \implies C(\pi^0) = +1 \\ \vec{B}, \vec{E} \in \pi\vec{B}, \pi\vec{E} &\implies C(\pi) = -1 \end{aligned}$$

initial state  $C(\pi^0) = +1$ ,

final state  $C(\pi\pi) = +1$ ,  $C(\pi\pi\pi) = -1$

conservation of C in  $\pi^0$  decays

Or... calculating decay amplitudes

$$A_{\mu\mu} = \langle \mu\mu | C^{\dagger} C H C^{\dagger} C | \mu^0 \rangle = \langle \mu\mu | C H C^{\dagger} | \mu^0 \rangle$$
$$= \langle \mu\mu | H | \mu^0 \rangle = A_{\mu\mu}$$

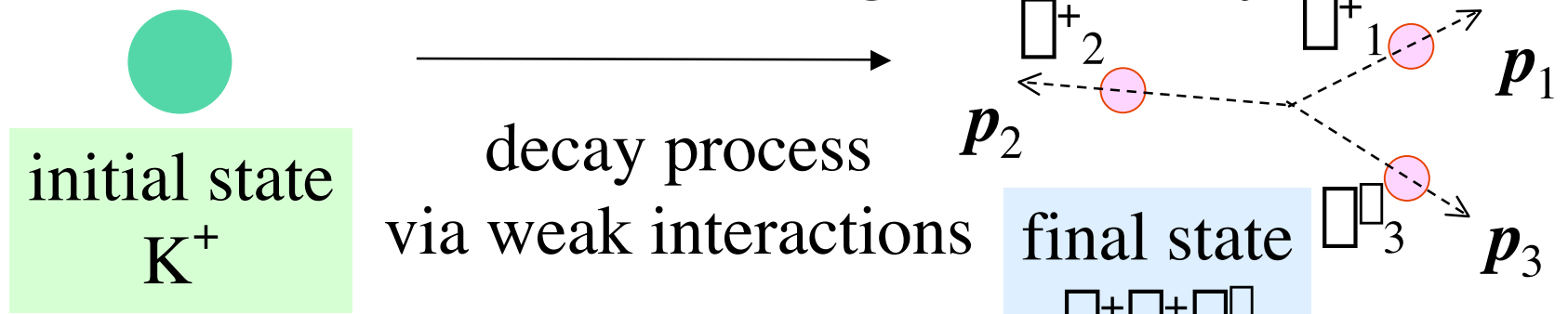
$$A_{\mu\mu} = 0$$

Weak interactions do not conserve:  
flavour quantum numbers,  
C, P, T, CP, CT and PT

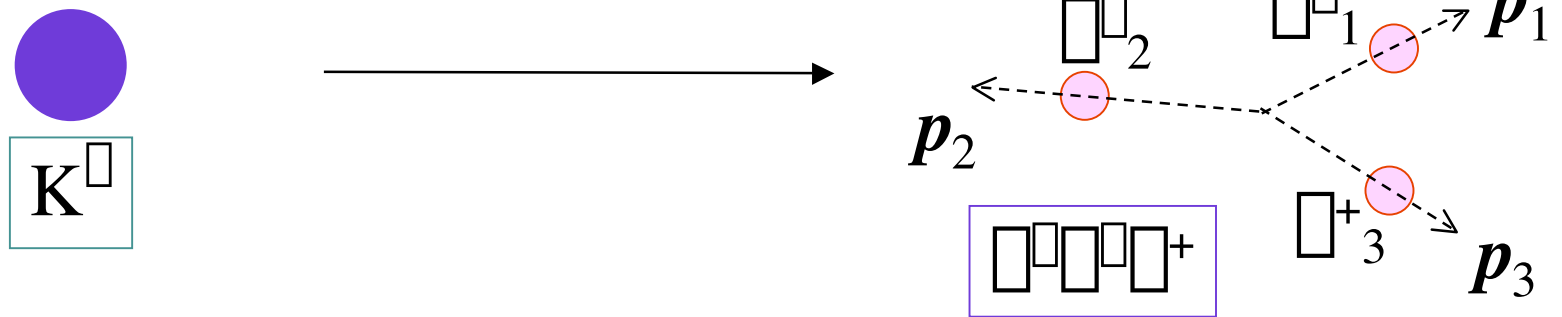
The topic of this lecture series.



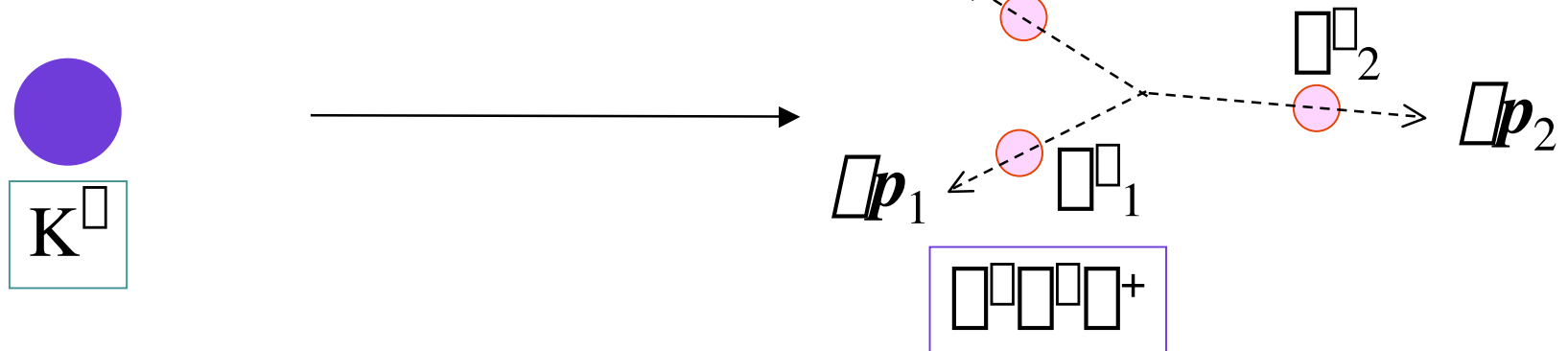
# 4) CP violation in the charged kaon system



## C transformation



## CP transformation



Partial decay width for  $K^+ \rightarrow \pi^+ \pi^+ \pi^0$

$$\Gamma_{K^+ \rightarrow \pi^+ \pi^+ \pi^0} = \int d^3 p_1 \int d^3 p_2 \int d^3 p_3 \Gamma_{\pi_1^+, \pi_2^+, \pi_3^0}(\vec{p}_1, \vec{p}_2, \vec{p}_3)$$

C transformed partial decay width

$$\begin{aligned} \Gamma_{K^+ \rightarrow \pi^+ \pi^+ \pi^0}^C &= \int d^3 p_1 \int d^3 p_2 \int d^3 p_3 \Gamma_{\pi_1^0, \pi_2^0, \pi_3^+}(\vec{p}_1, \vec{p}_2, \vec{p}_3) \\ &\equiv \Gamma_{K^0 \rightarrow \pi^0 \pi^0 \pi^+} \end{aligned}$$

CP transformed partial decay width

$$\begin{aligned} \Gamma_{K^+ \rightarrow \pi^+ \pi^+ \pi^0}^{CP} &= \int \int d^3 p_1 \int d^3 p_2 \int d^3 p_3 \Gamma_{\pi_1^0, \pi_2^0, \pi_3^+}(\pi \vec{p}_1, \pi \vec{p}_2, \pi \vec{p}_3) \\ &= \int d^3 p_1 \int d^3 p_2 \int d^3 p_3 \Gamma_{\pi_1^0, \pi_2^0, \pi_3^+}(\vec{p}_1, \vec{p}_2, \vec{p}_3) \\ &= \Gamma_{K^0 \rightarrow \pi^0 \pi^0 \pi^+} \end{aligned}$$

Partial decay width:  $\Gamma_{K^+ \rightarrow \pi^+ \pi^+ \pi^0}$  and  $\Gamma_{K^0 \rightarrow \pi^0 \pi^0 \pi^+}$   
 are CP and C transformed to each other

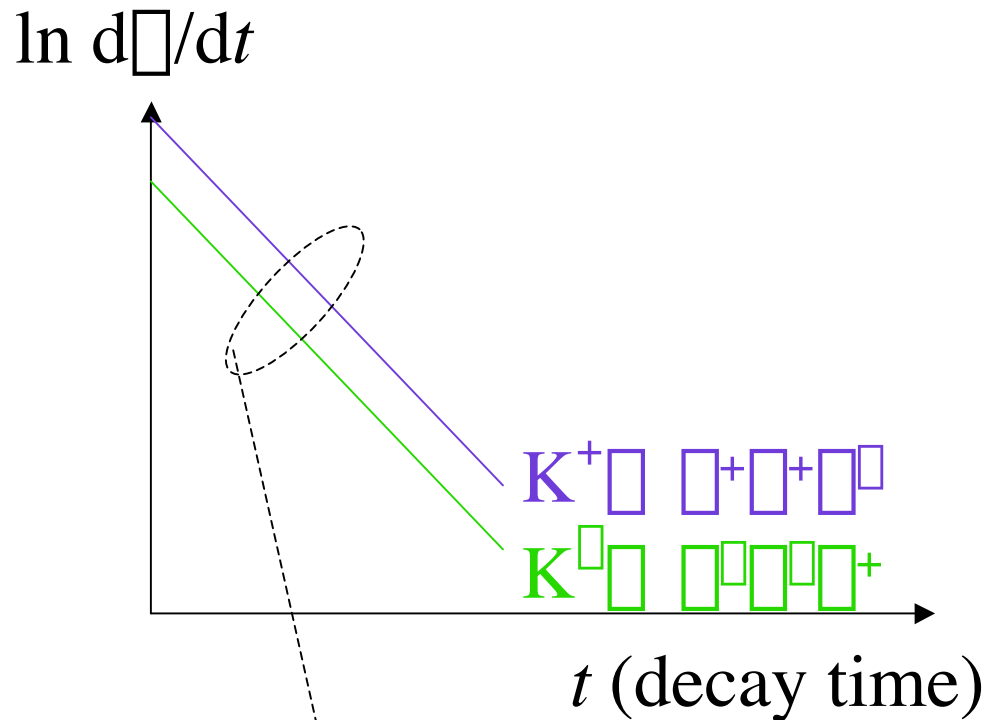
If  $\Gamma_{K^+ \rightarrow \pi^+ \pi^+ \pi^0} \neq \Gamma_{K^0 \rightarrow \pi^0 \pi^0 \pi^+}$  ~~C~~ and ~~CP~~!

**NB: these differences can appear in**  
 $\Gamma$  or  $d\Gamma/dt$

In general,  
~~C~~ and ~~CP~~ are needed in order to generate  
 partial decay widths differences between  
 particles and anti particles.

**Total widths between  $K^+$  and  $K^0$  must be identical**  
**CPT**

If there were  $\cancel{CP}$  and  $\cancel{C}$ , we would observe



**Not seen!**

Exponential decays

( $K^\pm$  have definite masses and decay widths)

Slopes are given by the total decay widths:

CPT theorem guarantees that they are identical.