

# Particle ID in high $P_T$ reactions(3)

- oVertex detectors and B-tagging
- oTop -physics
- oParticle ID for  $M_{top}$  measurement
- oParticle ID for SM-Higgs search
- oSUSY, Exotics
- oConclusion

# Vertex detectors (1)

Requirements based on b-jets parameters

- B hadrons lifetime : average of  $\sim 1.6$  ps
- semi-leptonic fraction  $\sim 10\%e$ , and  $10\%m$
- $ct = 470$  microns  $\rightarrow$  impact parameter  $d \sim 100$  microns
- need accuracy :  $< 20$  microns on  $d$

Accuracy limited by

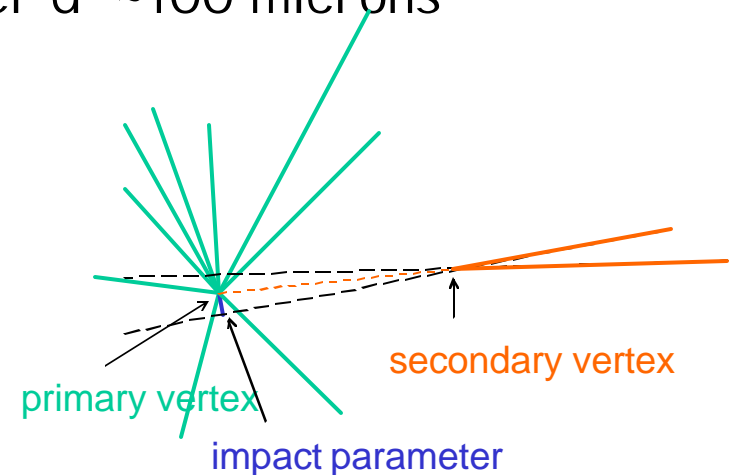
- lever arm,
- granularity,
- number of layers

Solution:

- 3 layer pixel detector
- first layer as close as possible to beam pipe
- single hit accuracy  $< 15$  microns in  $rf$
- equipped with fast electronics

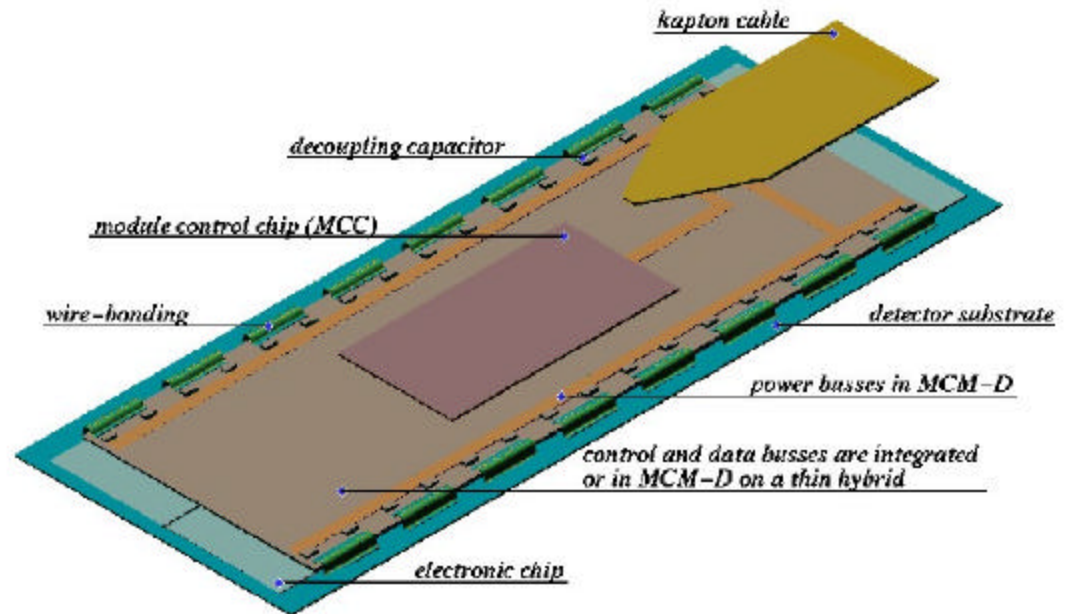
Beware -of radiation damage

- multiple scattering in material
- power dissipation



# Atlas pixels

- 3 cylinders  $R= 5\text{cm} , 10\text{cm}, 13\text{cm}$
- 2 x 3 disks
- Sensor: " n + n oxygenated " ,
- FE Electronics DeepSubMicron
- 82 M pixels 50 x 400 mm
- 2.8%  $X_0/\text{layer}$



Carbon fiber mechanical support  
Length about 80cm

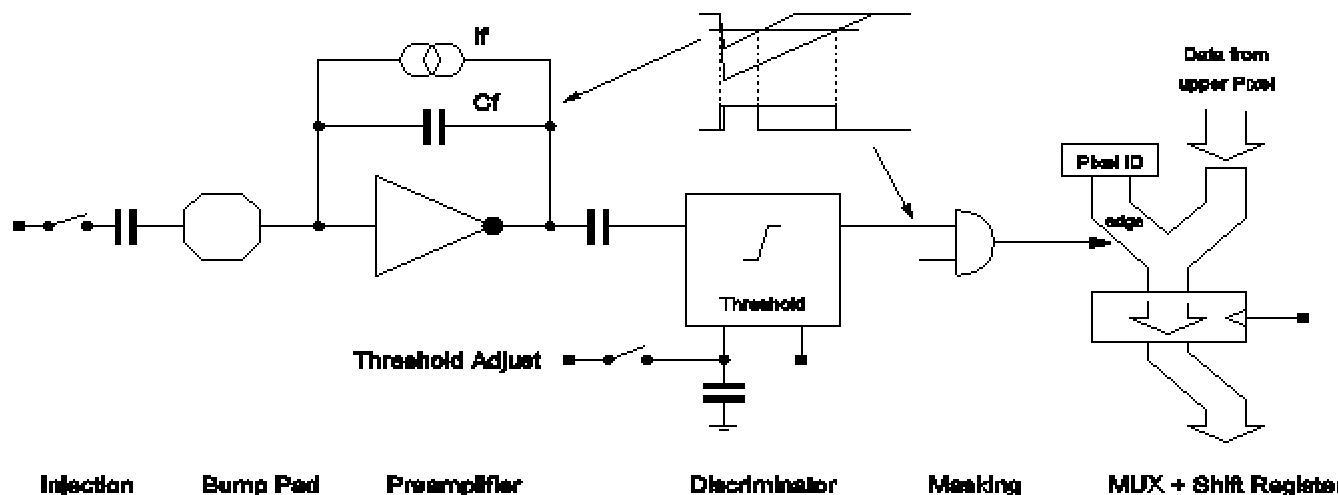
# Pixels electronics

## FE electronics and readout architecture

- 80 million channels to be “looked at” every 25 ns !!
- fast preamp, good  $S(15ke)/N(200e)$  for mips → digital or few bits r/o
- all FE chips proceed in parallel, controlled by local MicroController
- occupancy small ( $\ll 10^{-2}$ ) : logic by columns

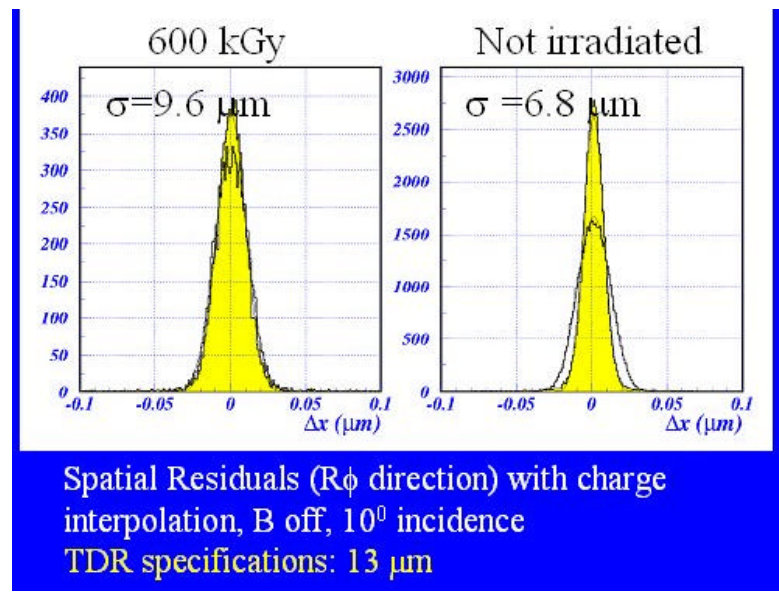
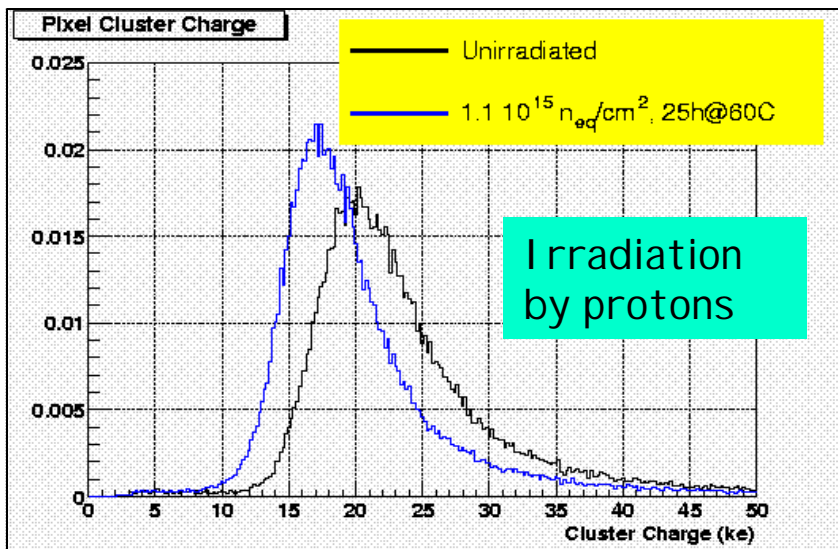
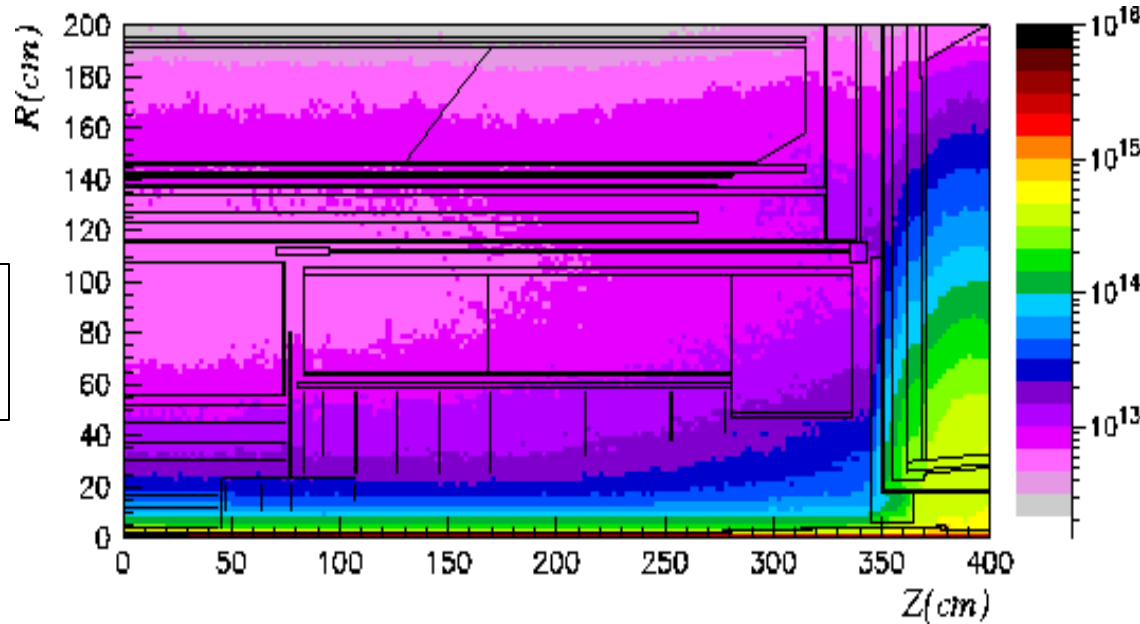
## Schematically:

- each FE chip covers 24 columns of 160 cells  $[400 \times 50 \text{ mm}] = \sim 1 \text{ cm}^2$
- at the bump bond pad is connected preamp + discriminator
- the End of Column logic drains data from columns =  
address of hit pixel(s) in the column + bunch counter
- upon LVL1 signal ( $< 100 \text{ kHz}$ ) the MC scans over its FE chips ( $\sim 12$ )  
gather the hit pixel addresses + bc, keeps those with proper bc, and  
clears the buffers.

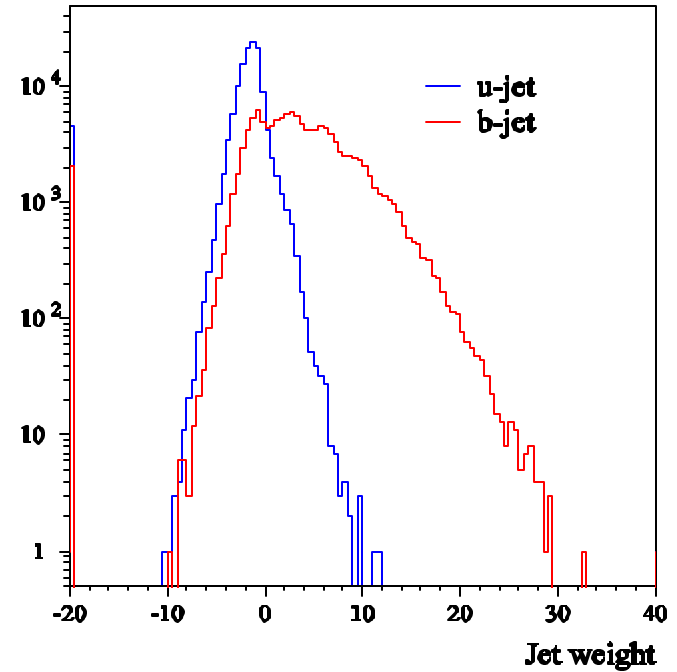
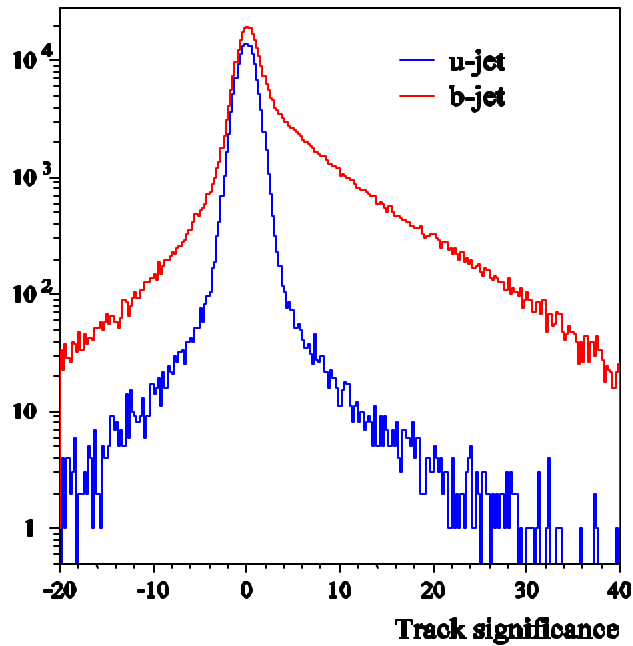


# Pixels: radiation damage

Expected neutron fluence per year at high luminosity <sup>®</sup>



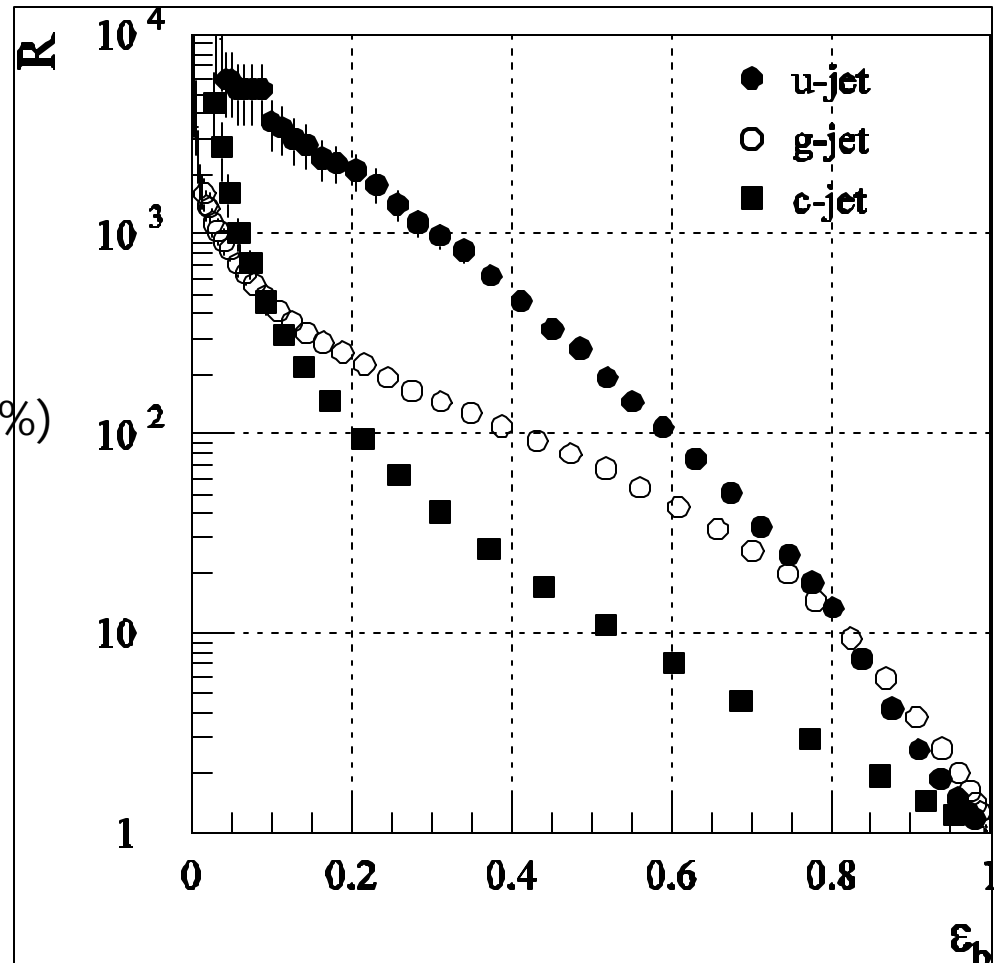
# B-tagging expected performance



- Start from a calorimeter jet
- Reconstruct tracks, and select those with  $p_T > \text{cut}$  ( $\sim 1 \text{ GeV}/c$ ) and  $\Delta R < 0.4$
- Measure  $d$  (with a sign), calculate  $s = d/s$
- Calculate jet weight as  $\sum \log(\text{signif as b} / \text{signif as u})$  from all its tracks
- Adjust cut position
- Tails in u-jets ? secondary interactions,  $V_0$

# B-tagging expected performance

- Full offline reconstruction
- Based on jet weights
- B jet sample from  $H(100) \rightarrow bb$  and  $H(400) \rightarrow bb$
- Background from jet-jet events same  $p_T$  range
- $L = 2 \times 10^{33}$   
(at  $10^{34}$  rejection reduced by  $\sim 20\%$ )
- $R_u \sim 100$  for 60% efficiency
- $R_c$  more modest because  $ct$  ( $D^+$ )=300 mm, ( $D^0$ )=120 mm
- $R_g$  limited by gluon splitting to  $cc$
  
- Soft lepton tag  
- identification of electrons and muons in cone axis adds a few % efficiency each



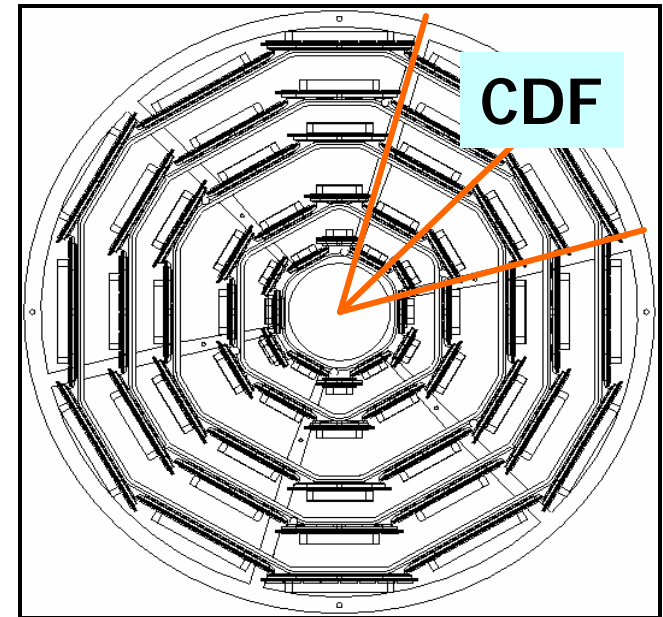
# B-tagging in HLT?

- CdF is using it for selection of “unbiased” B decays
- Starting from LVL1 with 2 tracks  $>2$  GeV/c  $p_T$
- Done in a “hardware oriented” way with a processing time of 25 ms/event

- 12 independent sectors
- 4 layers(out of 5) of micro-strips grouped to 250 mm pitch
- 1 point/tracker ( $f$ ,  $p_T$ ) + hit pattern compared to 32000 masks

- Interesting approach in LHCb/Velo

- In ATLAS-CMS events to be processed, at LVL2 have in general several jets
- Some trial at LVL2-Atlas, not really convincing
- Need “real tracking” to decide if some of the jets are B-jets or not.
- Efficient way to do it is at the Event Filter level, with ~full offline performance





# B-tagging in HLT:CMS example

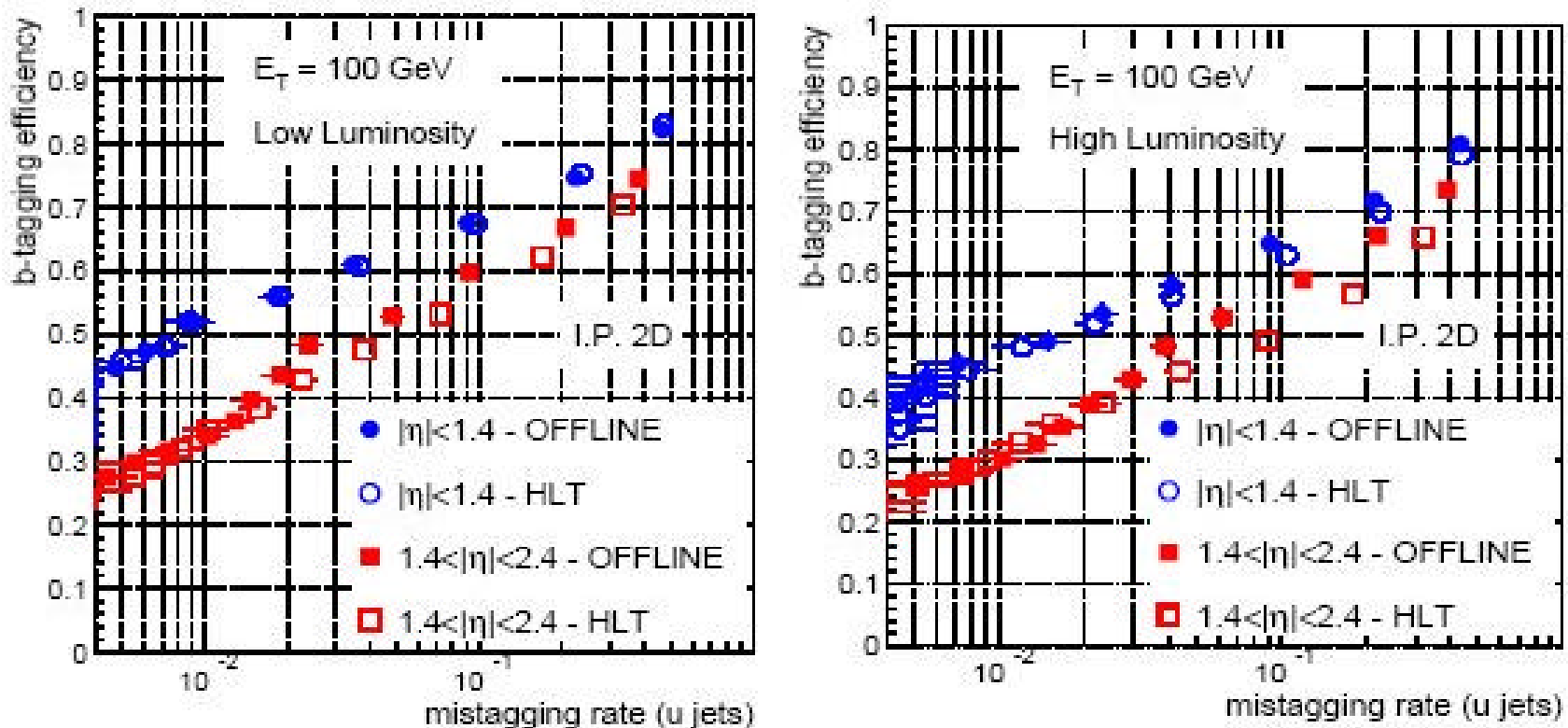
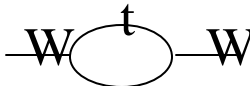


Figure 15-68 Efficiency for the  $b$ -tag versus mistagging rate for jet with  $E_T=100$  GeV in the low (left) and high (right) luminosity scenarios.

- Reconstruct Tracks in a cone defined by Calo-jet
- Ask for at least 2/3 pixel hits.
- Labelled as  $b$ -jet if  $\geq 2$  tracks above thresh significance

# Particle ID for M(top) measurements

M(W) , M(top) and M(Higgs) are linked in the standard model:

$$m(W) = \left( \frac{2 \sqrt{2} G_F}{\sin^2 \theta_W} \right)^{1/2} \left( \frac{1}{\sqrt{1 - \sin^2 \theta_W}} \right)^{1/2} F[M(\text{top}), \log(M(\text{H}))]$$


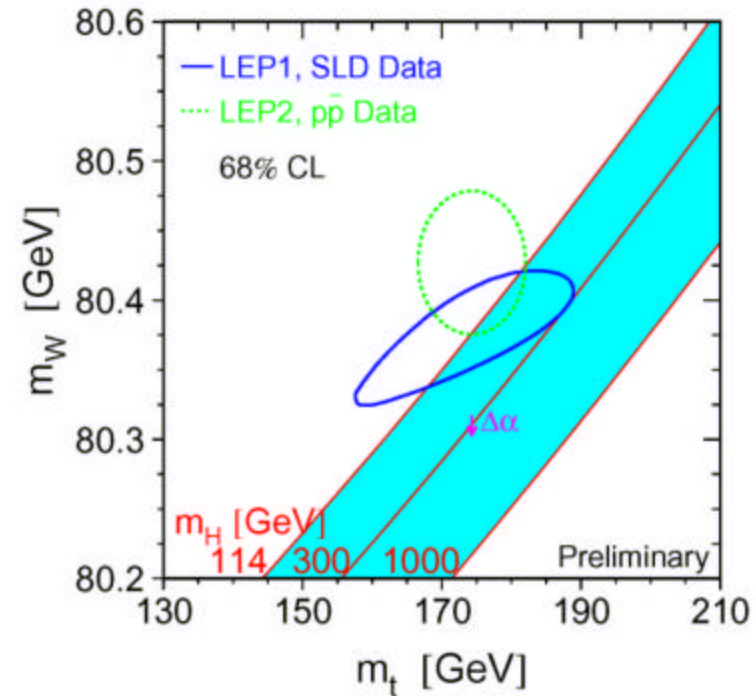
$G_F$ ,  $\alpha_{EM}$ ,  $\sin \theta_W$  known with high precision,  
 precise measurements of M(top) and M(W)  
 constrain M(H) (weakly because of log term)

10 MeV M(W)  $\rightarrow$  ~1 GeV M(top)

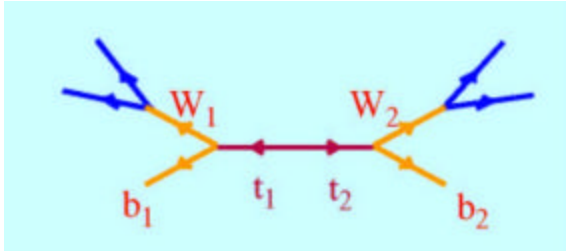
To day (CDF+D0) M(top)=178  $\pm$  5 GeV

Best strategy to measure it at LHC?

[event statistics is not the main problem]



# Particle ID for M(top) measurements



Inclusive tt cross-section ~0.7 nb

t→bW 100%  
 W→ln 11%each , 67% hadronic; no Bs

|                        |   |
|------------------------|---|
| W1→had, W2→had         | : 44% 2 masses fully rec/ huge QCD BG   |
| W1→e/μ, W2→had (+ 2→1) | : 30% 1 mass fully reconstructed        |
| W1→e/μ, ν, W2→e/μ      | : 5% rate; mass not fully reconstructed |
| Others: one W→t        | : not appropriate for precise M meas    |

**Semi-leptonic final state:**  
 expect 2.5 million evts for 10fb<sup>-1</sup>

- one electron or muon → LVL1 Trigger
- two b jets → Rejection of QCD bckg fact 100/jet
- two non-b-jets  $M_{inv}=M(W)$  → Masse scale
- missing  $E_T$  → Rejection of QCD bckg, Reconstruction of 2nd top
- in average 2 or 3 g jets/QCD
- underlying event

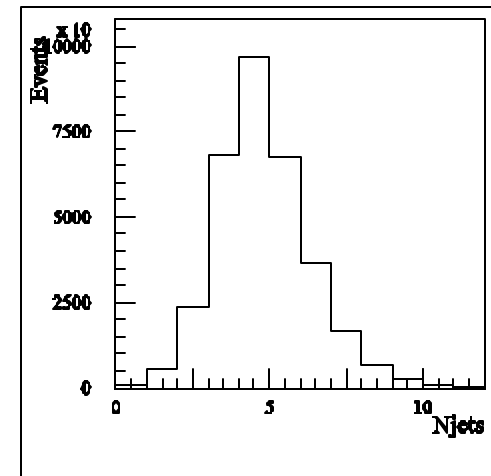
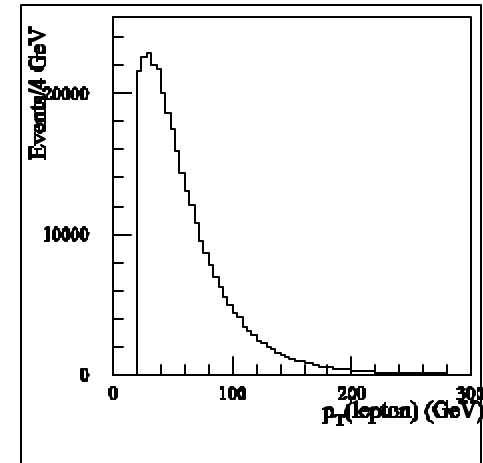
# Particle ID for M(top) measurement

## Selection

- 1 iso lepton,  $p_T > 20$  GeV,  $|\eta| < 2.5$
- $p_{T\text{miss}} > 20$  GeV
- = 4 jets with  $p_T > 40$  GeV,  $|\eta| < 2.5$
- = 2 jets with b-tag

Selection efficiency. = 5%

$\mathcal{P}$  126k events, with S/B ~65

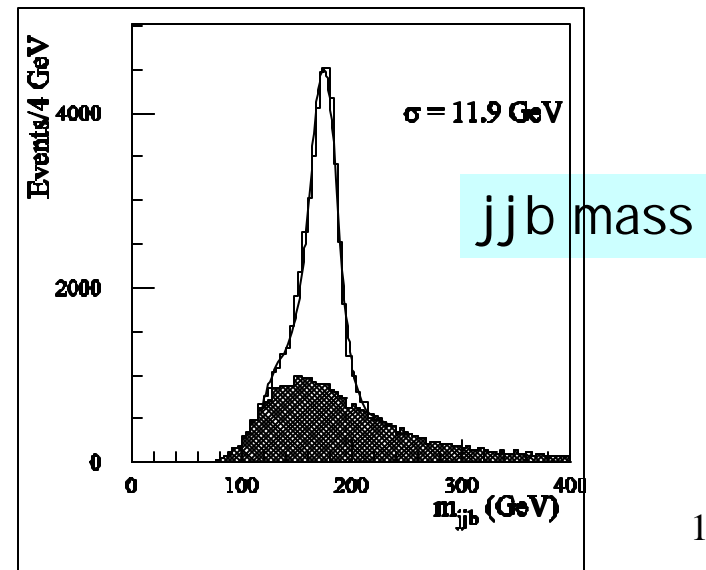
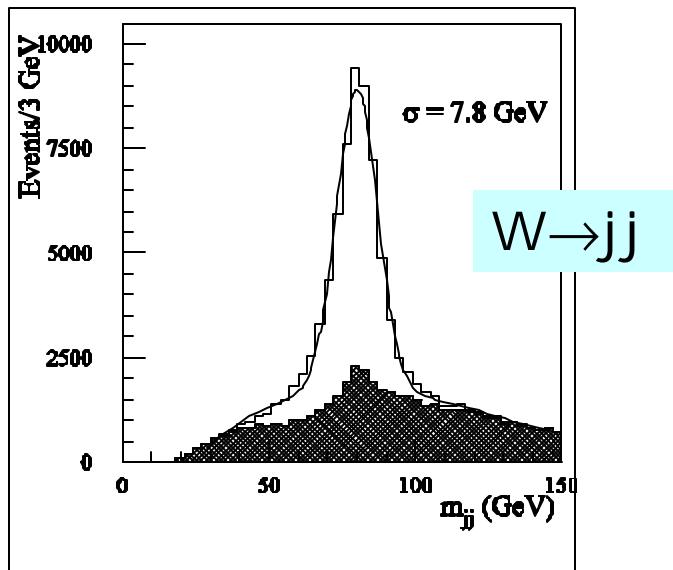


| Process           | $p_T^l > 20\text{GeV}$<br>$E_{T\text{miss}} > 20\text{GeV}$ | As before,<br>plus $N_{\text{jet}} \geq 4$ | As before,<br>plus $N_{\text{b-jet}} \geq 2$ | Events<br>per $10 \text{ fb}^{-1}$ |
|-------------------|---|--|--|------------------------------------|
| $t\bar{t}$ signal | 64.7  | 21.2                                       | 5.0  | 126 000                            |
| W+jets            | 47.9  | 0.1  | 0.002  | 1658                               |
| Z+jets            | 15.0  | 0.05                                       | 0.002  | 232                                |
| WW                | 53.6  | 0.5  | 0.006  | 10                                 |
| WZ                | 53.8  | 0.5  | 0.02   | 8                                  |
| ZZ                | 2.8   | 0.04                                       | 0.008  | 14                                 |
| Total background  |   |  |  | 1922                               |
| S/B               |   |  |  | 65                                 |

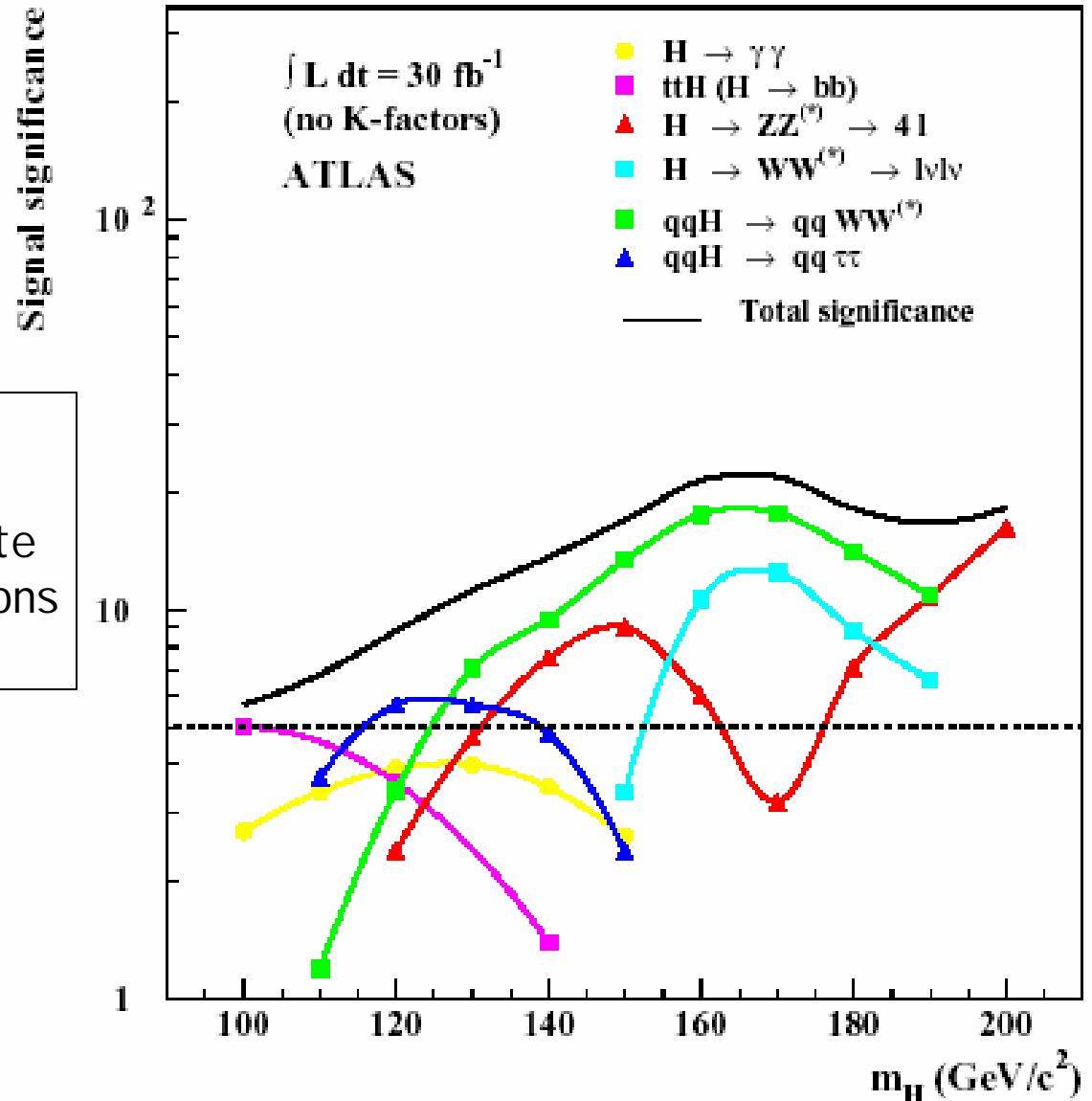
Efficiency of cuts

# W-ID from "top" sample

- ½ leptonic ttbar events contain  $W \rightarrow \text{jet-jet}$  evts with good S/N  
⇒ identified hadronic W decays
- This sample of jets is used to adjust the (non-b) jet -scale, starting from jets normalized using photon-jet evts ( $p_T$  balance)
- B-jets are normalized from photon-B-id jets (only)
- Remaining uncertainties on jet scales dominate the top mass systematic uncertainty
- Hadronic Ws in jj used in other places ( $H \rightarrow WW, \dots$ )

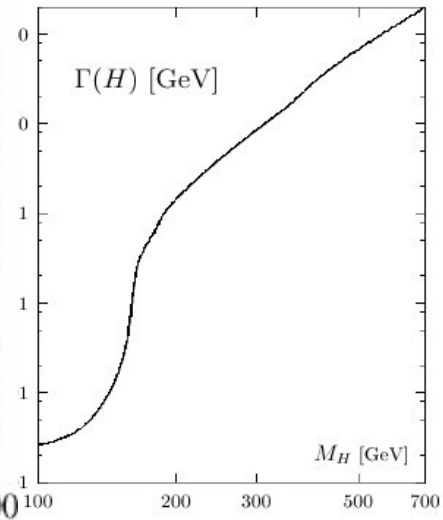
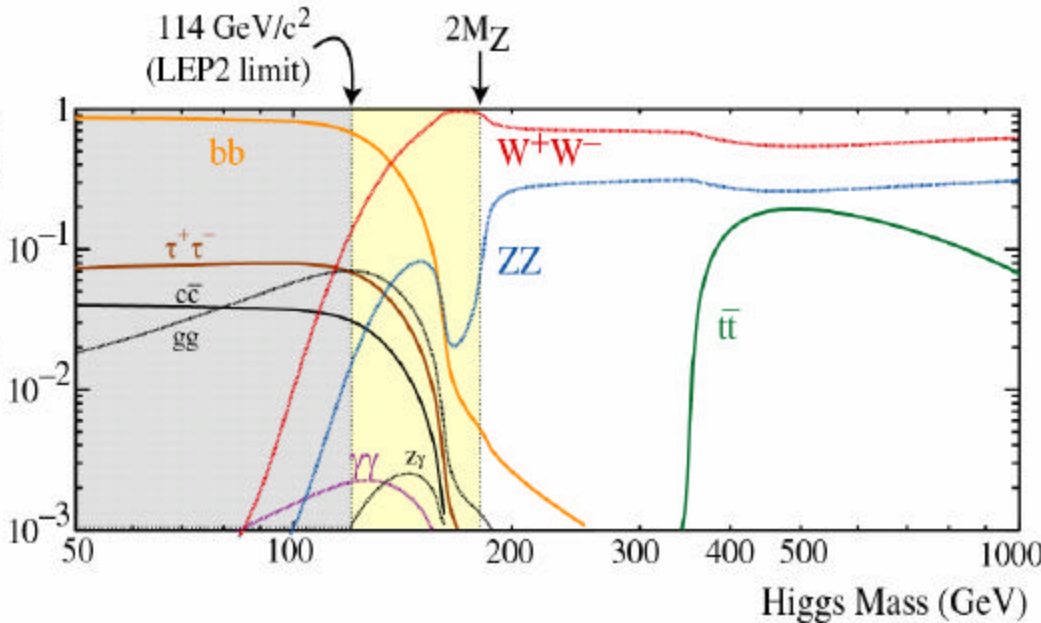
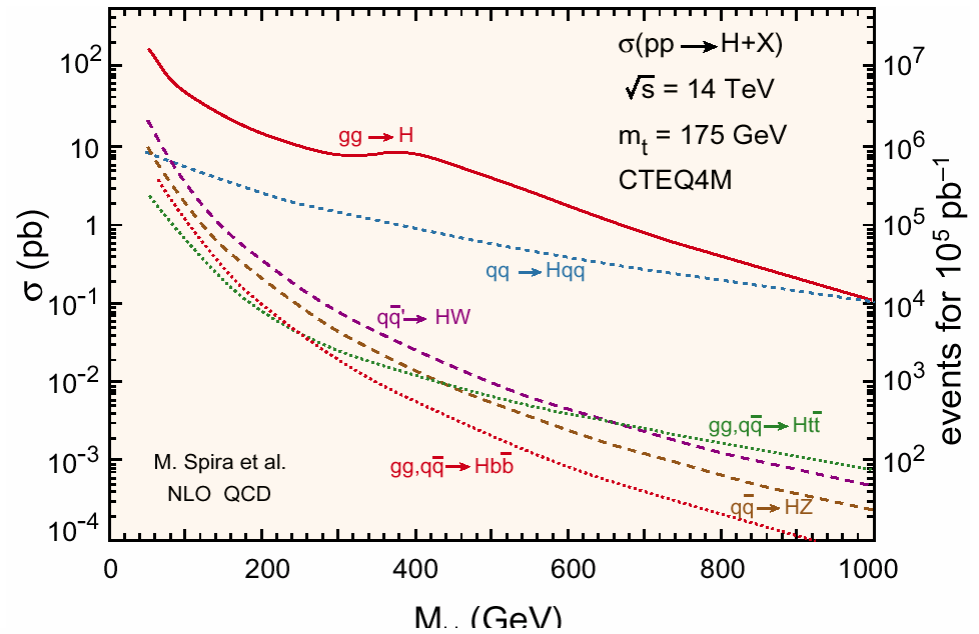
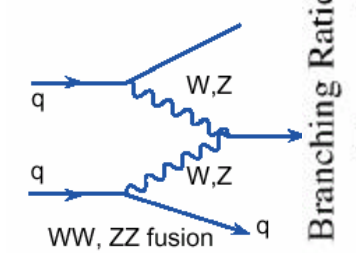
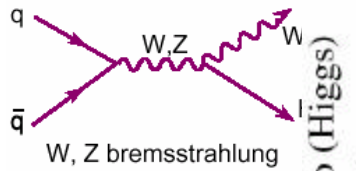
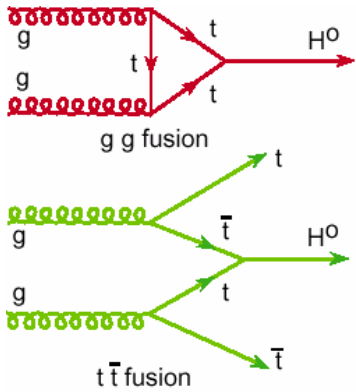


# SM-Higgs search global view



Each channel based  
on well identified  
particles in the final state  
In general containing leptons  
and/or gauge bosons

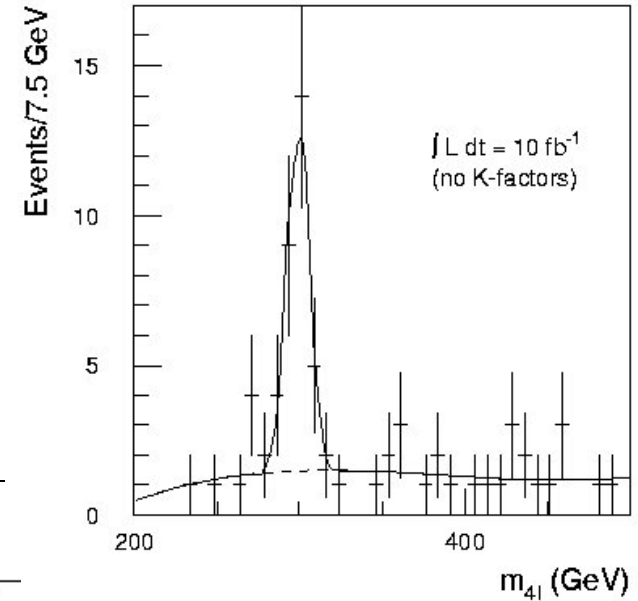
# Production and decay



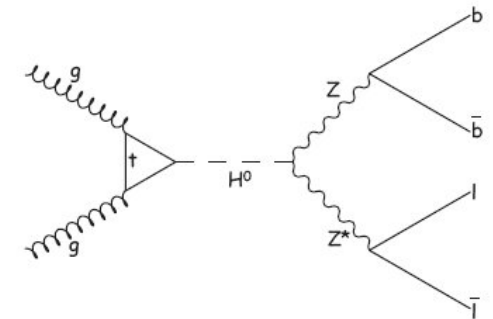
# The ZZ "gold plated" mode

$M(H) > 180$  GeV : ZZ mode open

- "Gold plated channel" with each  $Z \rightarrow ee$  or  $\mu\mu$
- Two  $M(H) = M(Z)$  constraints
- Irreducible background (only) = non resonant ZZ
- Main limitation is rate, when  $M(H)$  increases
- Adding evts with one  $Z \rightarrow tt$  could help (trig/other Z)
- Using  $Z \rightarrow nn$  (\*6  $ee$  or  $\mu\mu$ ) allows to extend search up to 700 GeV



| Higgs mass (GeV)                                     | 200  | 240  | 280  | 320  | 360  | 400  | 500  |
|--|------|------|------|------|------|------|------|
| $BR(H \rightarrow ZZ)$                               | 0.26 | 0.29 | 0.30 | 0.31 | 0.30 | 0.28 | 0.27 |
| $\sigma \times BR$ (fb)                              | 12.4 | 11.2 | 9.6  | 8.9  | 8.7  | 6.8  | 3.2  |
| Signal (no $p_T$ cut)                                | 134  | 127  | 110  | 105  | 105  | 86   | 44   |
| Background (no $p_T$ cut)                            | 74   | 57   | 43   | 33   | 29   | 29   | 17   |
| $S/\sqrt{B}$ (no $p_T$ cut) for $30 \text{ fb}^{-1}$ | 15.6 | 16.8 | 16.8 | 18.2 | 19.3 | 15.9 | 10.7 |

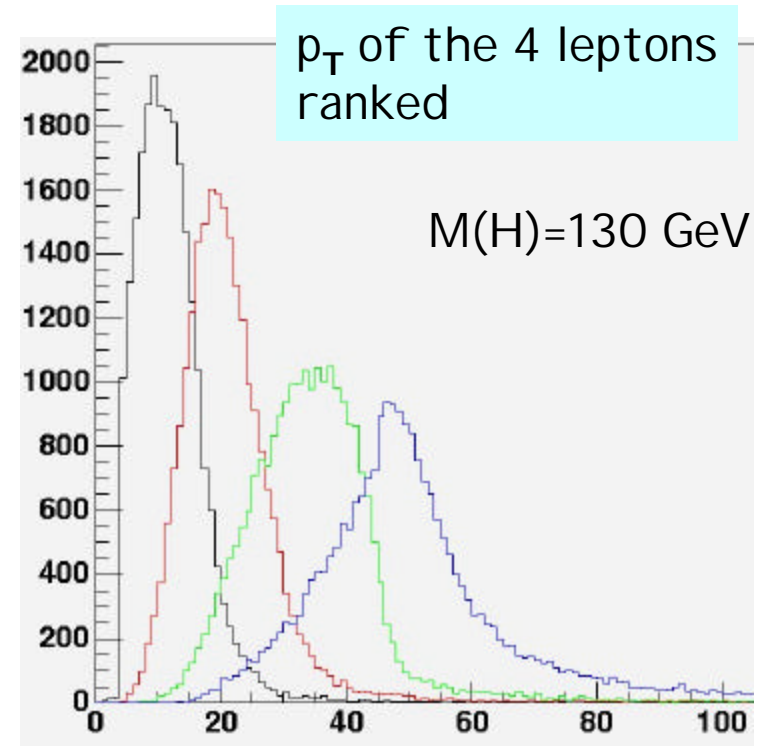


-what about one Z in jet-jet? BG "Z +2 jets" is large....  
 some attempt with one Z in two b jets....



# ZZ\* 130 < M(H) < 180 GeV

- still 4 leptons ee or mm
- only one M(H)=M(Z) constraint remaining
- dip in the BR at 2M(W)
- Zbb main reducible background can be reduced applying a veto on displaced vertices
- important issue : acceptance of the lower PT lepton .Analyses require 2 leptons pT > 20 GeV (trigger) and 2 leptons with pT > 7 GeV (offline)



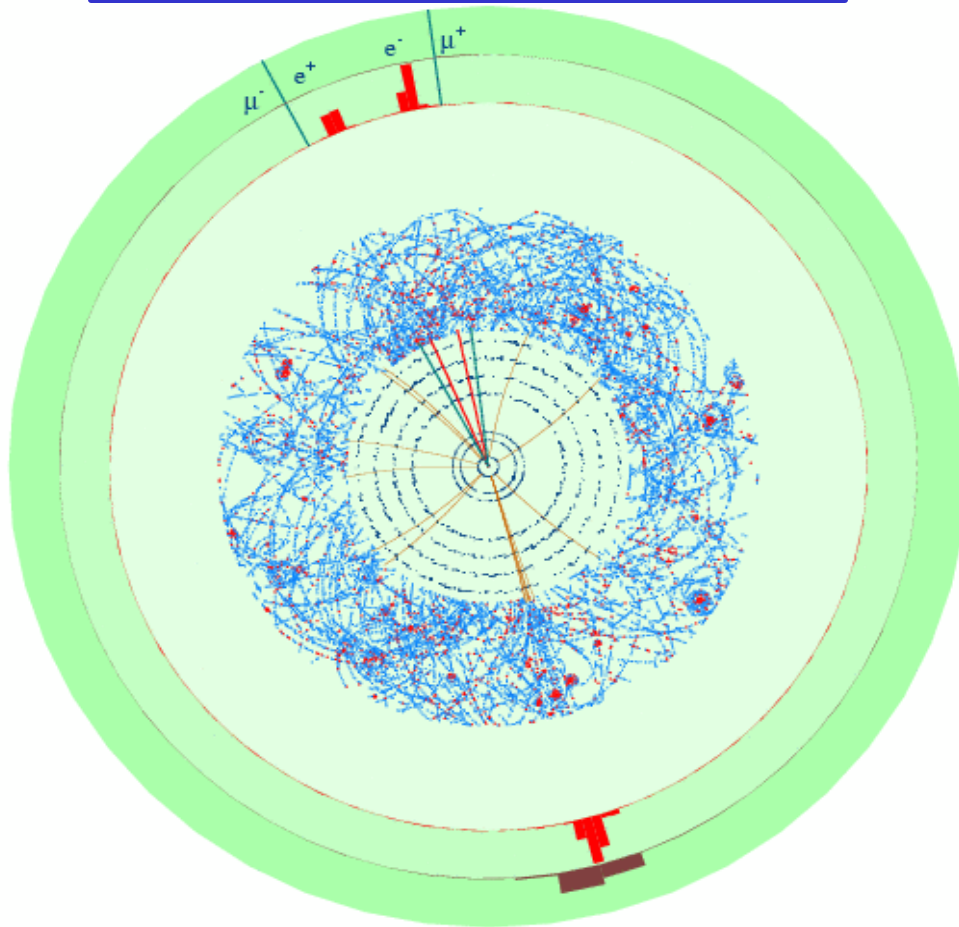
**Table 19-14** Signal and background rates after all cuts and signal significances as a function of  $m_H$ , for  $H \rightarrow ZZ^* \rightarrow 4l$  events and for an integrated luminosity of 30 fb<sup>-1</sup> (low luminosity performance).

| Higgs mass (GeV)              | 120  | 130  | 150  | 170  | 180  |
|-------------------------------|------|------|------|------|------|
| Signal                        | 4.1  | 11.4 | 26.8 | 7.6  | 19.7 |
| $t\bar{t}$                    | 0.01 | 0.02 | 0.03 | 0.02 | 0.02 |
| Zbb                           | 0.08 | 0.12 | 0.19 | 0.17 | 0.19 |
| ZZ*                           | 1.23 | 2.27 | 2.51 | 2.83 | 2.87 |
| ZZ → ττH                      | 0.13 | 0.20 | 0.25 | 0.08 | 0.02 |
| Significance ( $S/\sqrt{B}$ ) | 3.4  | 7.0  | 15.5 | 4.3  | 11.2 |
| Significance (Poisson)        | 2.4  | 4.8  | 15.5 | 3.2  | 11.2 |

30 fb<sup>-1</sup>

ATLAS barrel

$H \rightarrow ZZ^* \rightarrow eemm$  ( $M(H)=130$  GeV)



High luminosity

# Higgs in 2 photons

For  $M(H) < 150$  GeV, most promising channels are:  
 $H \rightarrow gg$  and  $H \rightarrow tt$  in association with 2 forward jets (VBF)

Distinctive features of the  $H \rightarrow gg$  mode are:

- low  $s \times BR$
  - large irreducible background
  - potentially large instrumental background (jet-jet and g-jet)
  - clean signature
  - high invariant mass resolution
- } Defining criteria for EM calorimeters

Ultimate energy resolution , when all calibration/normalization problems are solved:

**CMS** :  $3\% / \sqrt{E} \oplus (200\text{MeV} \oplus \text{pile-up}) / E \oplus 0.55\%$

**ATLAS** :  $10\% / \sqrt{E} \oplus (200\text{MeV} \oplus \text{pile-up}) / E \oplus 0.70\%$

Remember converted photons are somewhat worse

$\rightarrow \Delta M(H) / M(H) \sim 1\%$

At High luminosity the vertex is ambiguous, which worsens the  $M(H)$  resolution in the absence of angular measurements (a place where converted photons<sup>19</sup> help)

# Jet-jet and g-jet rejection

Overall jet rejection obtained in ATLAS MC full simulation, confirmed by test beam (using fine strips):

-1050 for quark jets

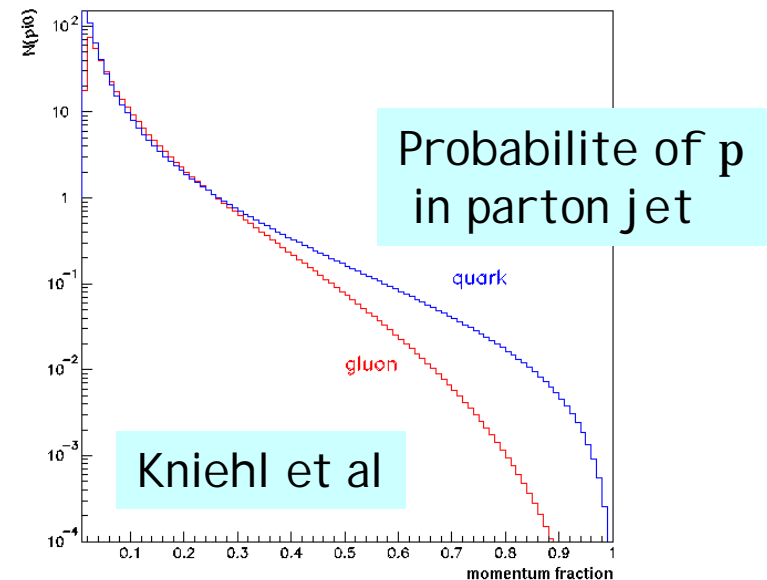
-6000 for gluon jets → Ultimate performance process dependant!

(probability of a high x isolated  $p^0$  is much higher in (MC) quark jet

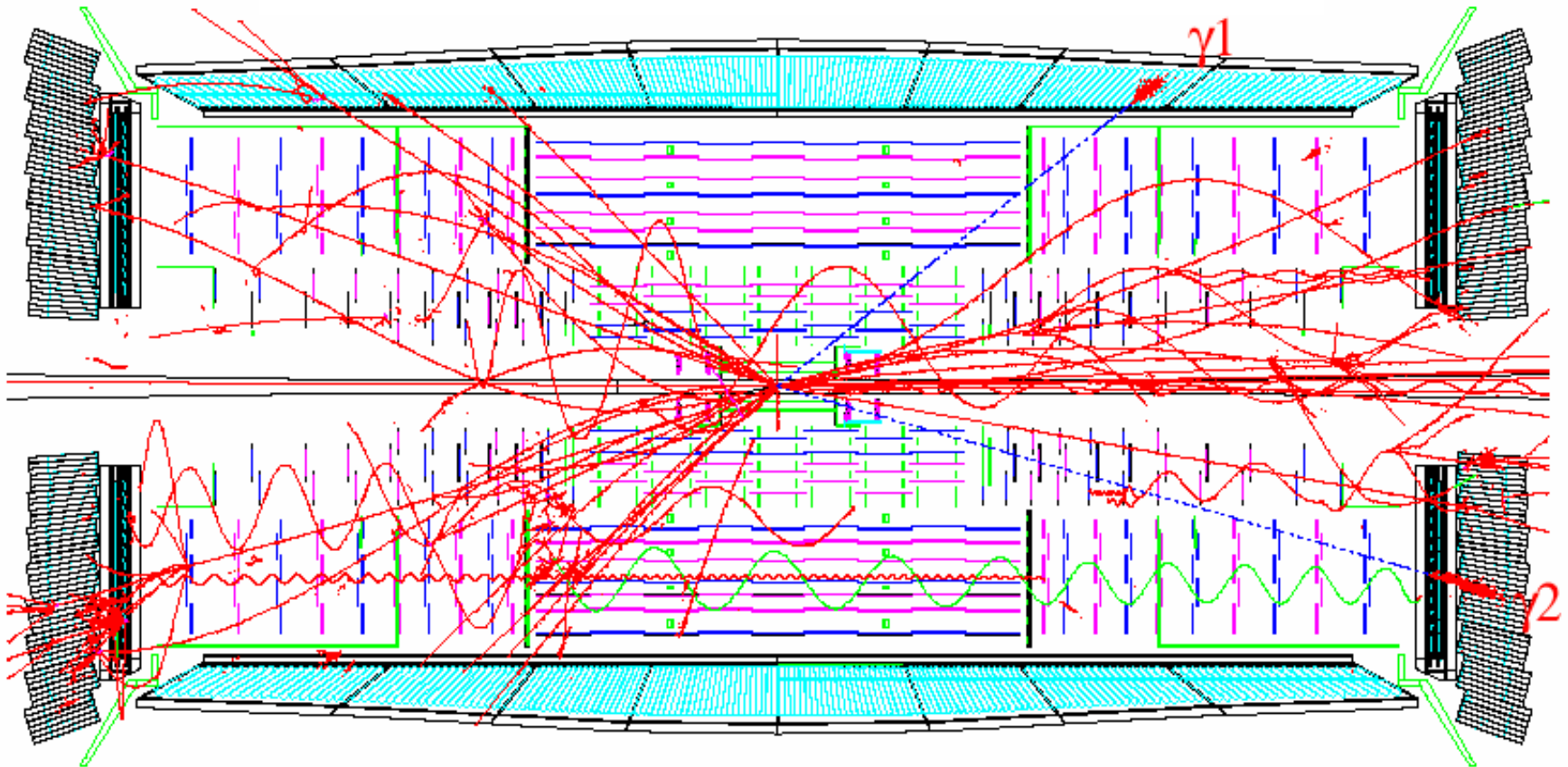
| Higgs mass (GeV)                                | 120    | 130    | 140    | 150    |
|---|--------|--------|--------|--------|
| Signal events (direct production)               | 1190   | 1110   | 915    | 617    |
| Signal events ( $WH, ZH, \bar{t}tH$ production) | 93     | 76     | 58     | 35     |
| $\gamma\gamma$ background                       | 29 000 | 24 700 | 20 600 | 16 900 |
| Jet-jet background                              | 4600   | 4100   | 3550   | 3050   |
| $\gamma$ -jet background                        | 5800   | 4900   | 4100   | 3400   |
| $Z \rightarrow ee$ background                   | -      | -      | -      | -      |
| Stat. significance for $100 \text{ fb}^{-1}$    | 6.5    | 6.5    | 5.8    | 4.3    |
| Stat. significance for $30 \text{ fb}^{-1}$     | 3.9    | 4.0    | 3.5    | 2.6    |

Channel very much dependant on detector  
Ultimate performance, resolution and photon/  
jet identification.

- Main jet-jet Bkg=gluon-gluon rejection of 6000 x 6000 OK
- Main g -jet bkg:g-quark rejection=1000....



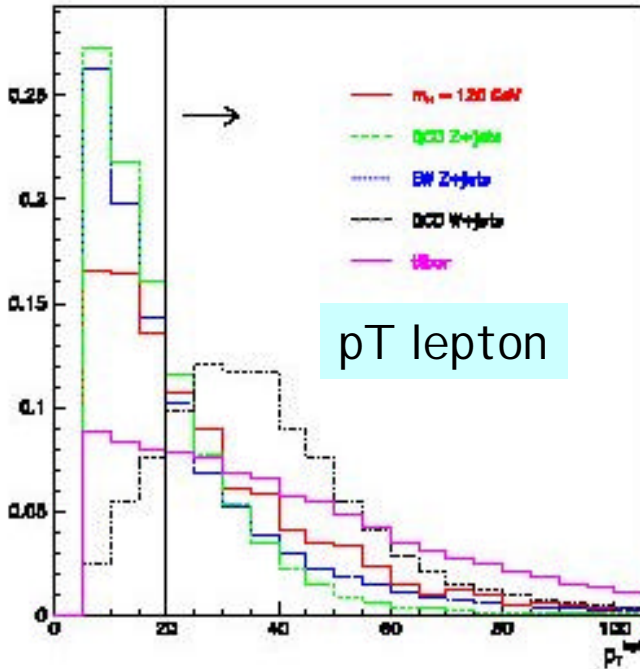
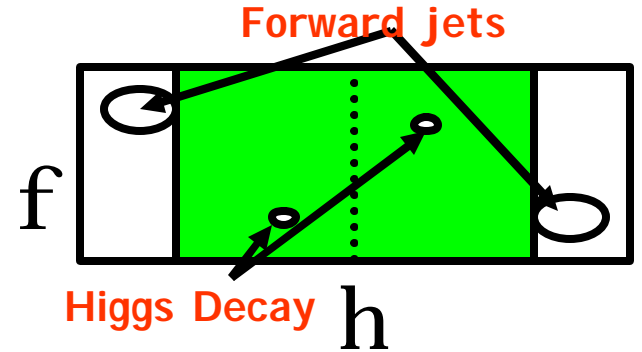
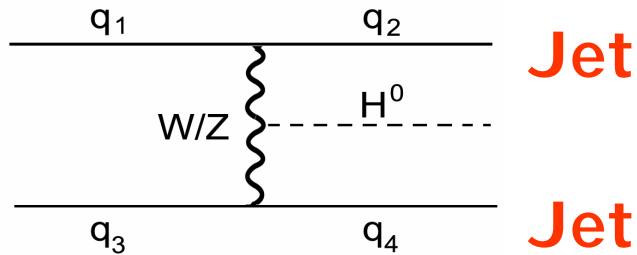
# A $H^{\circ}g$ event display in CMS



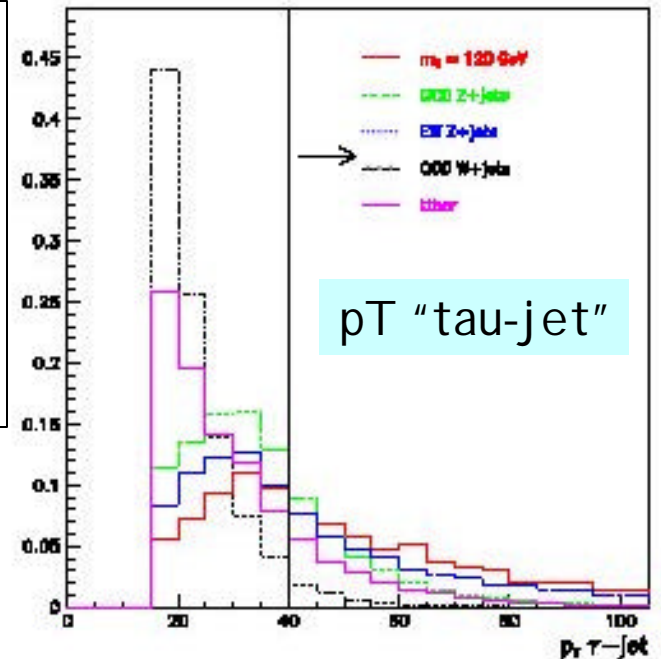
Low luminosity case-1 collision only

# Higgs $\rightarrow \tau\tau$ by VBF

- The  $t\bar{t}$  mode has a larger BR than  $gg$  but a weaker signature and worse resolution
- Asking 2 forward jets (and no other jet) selects **weak process by WW fusion**, which has less background



Trigger  
 Dilepton:  
 $E_{15\mu 10}$   
 Lepton-hadron  
 $p_{T\text{lept}} > 20 \text{ GeV}$   
 or  
 $\text{Tau}35 \text{ xE45}$



# Higgs $\rightarrow \tau\tau$ by VBF

Tagging jets:

$D_h > 4.4$

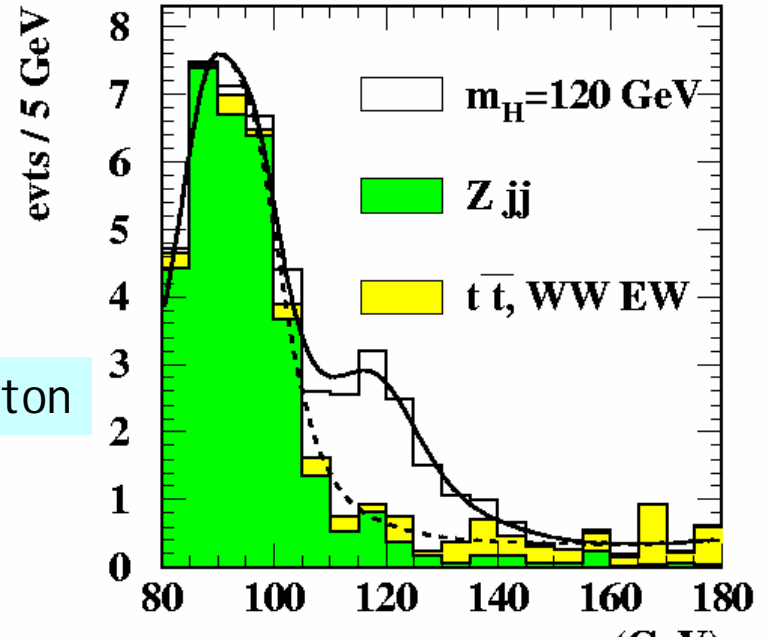
$M(JJ) > 700\text{GeV}$

Backgrounds:

$t\bar{t}$ ,  $Z+2$  jets

Invariant mass:

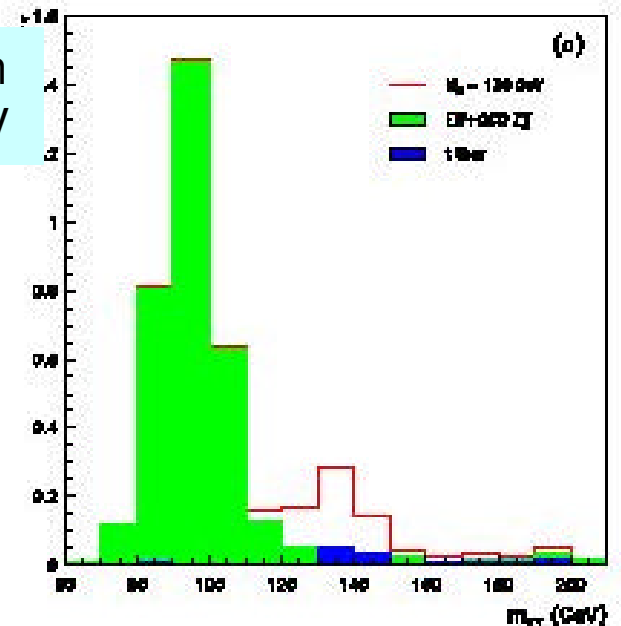
Calculated using the 2 taus and the missing ET



sBR in fb

Hadron-lepton  
 $M(H)=130\text{ GeV}$

| cuts                      | Hqq  | gg fusion | EW Zjj | W+jets            | $t\bar{t}$         |
|---------------------------|------|-----------|--------|-------------------|--------------------|
| preselection cuts         | 28.  | 101.2     | 59.3   | $4.3 \times 10^4$ | $2.07 \times 10^4$ |
| lept. trigger             | 13.7 | 50.31     | 22.0   | $3.4 \times 10^4$ | $1.6 \times 10^4$  |
| $\tau$ cuts.              | 6.18 | 22.65     | 8.03   | $3.2 \times 10^3$ | $4.3 \times 10^3$  |
| Fwd. tag                  | 1.97 | 0.176     | 1.72   | 30.               | 29.7               |
| $\tau$ decay and reconst. | 1.27 | 0.105     | 1.09   | 5.9               | 6.06               |
| $m_T$                     | 1.02 | 0.071     | 0.92   | 0.63              | 1.74               |
| $p_T^{\text{miss}}$       | 0.81 | 0.051     | 0.71   | 0.58              | 1.38               |
| $m_{jj}$                  | 0.71 | 0.028     | 0.69   | 0.37              | 1.01               |
| Jet veto                  | 0.63 | 0.014     | 0.66   | 0.21              | 0.14               |
| Mass Window               | 0.52 | 0.007     | 0.06   | 0.00              | 0.007              |



# SUSY

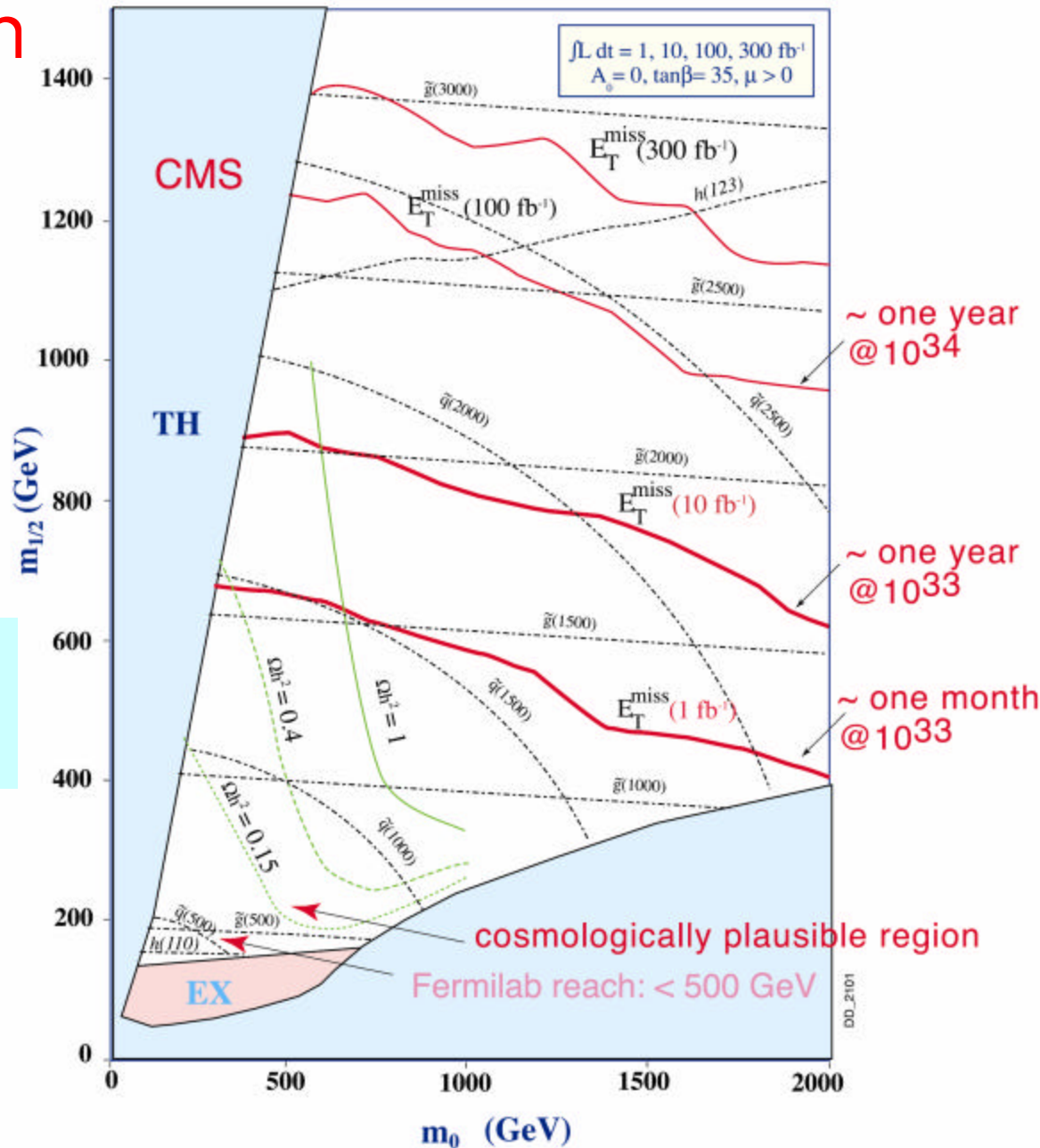
- In the MSSM each fermion  $f_{L,R}$  has a scalar partner  $sf_{L,R}$ , and to each Gauge boson is associated a massless spin  $\frac{1}{2}$  gaugino.
- Among other virtues, these states coming with an opposite sign to normal particles in loop corrections, cancel the quadratic divergence of the Higgs mass, which would otherwise run to "infinity".
- None of these particles has been so far observed  $\rightarrow$  heavy!
- $sq$  and  $sg$  couple to QCD as usual particles, they should be copiously produced at LHC as soon as threshold is passed.
- In the MSSM R-parity =  $(-1)^{3(B-L)+2S}$  is conserved:
  - s-particles are produced in pair
  - the lightest one is stable = neutralino  $\tilde{c}^0$
- Exact signatures depend on the details of the model.  
A common feature is large  $SE_T$  and missing energy (neutralinos)
- The reach of experiments can be evaluated and compared using as parameter  $m_0$  ( $m_{1/2}$ ) the common mass of all scalars (spin  $\frac{1}{2}$ ) at GUT scale



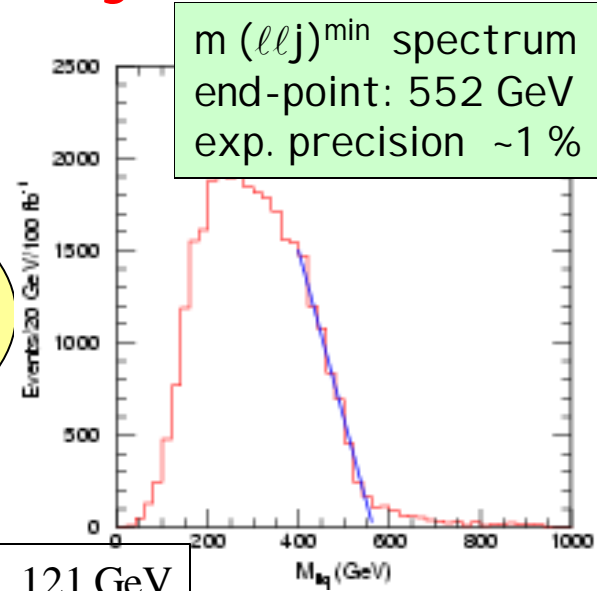
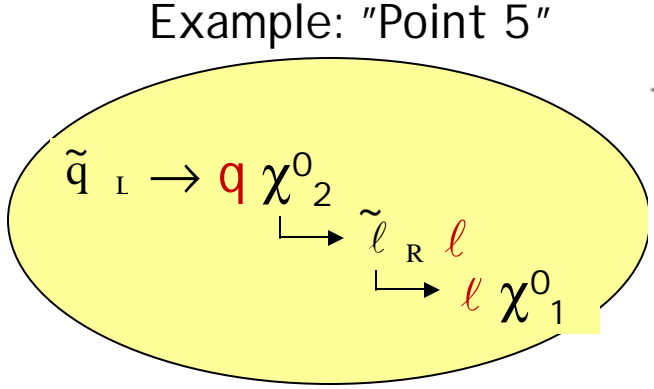
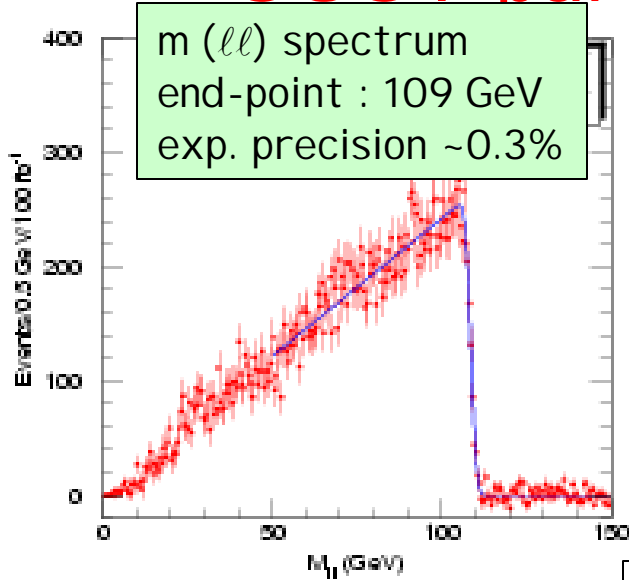
# CMS-SUSY reach

- Rates are high:  
100 evts /day for  
 $m(\text{sq}, \text{sg}) = 1 \text{ TeV}$
- Trigger is easy:  
multijets  $\oplus E_T^{\text{miss}}$

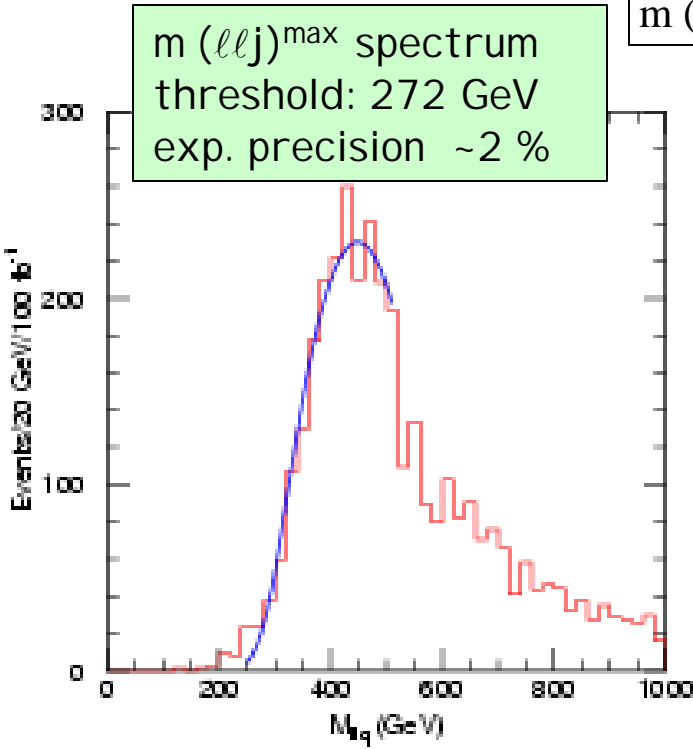
In principle, ...in one month  
at  $10^{33}$ , s-quarks and gluinos  
up to  $\sim 2 \text{ TeV}$  can be discovered



# SUSY parameters from decay chain

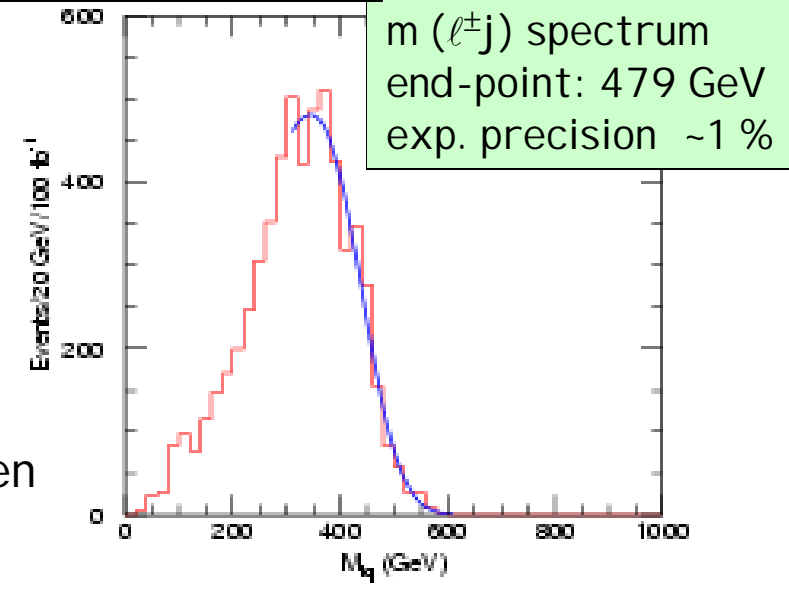


$m(\tilde{q}_L, \chi^0_2, \tilde{l}_R, \chi^0_1) = 690, 232, 157, 121 \text{ GeV}$



100 fb<sup>-1</sup>

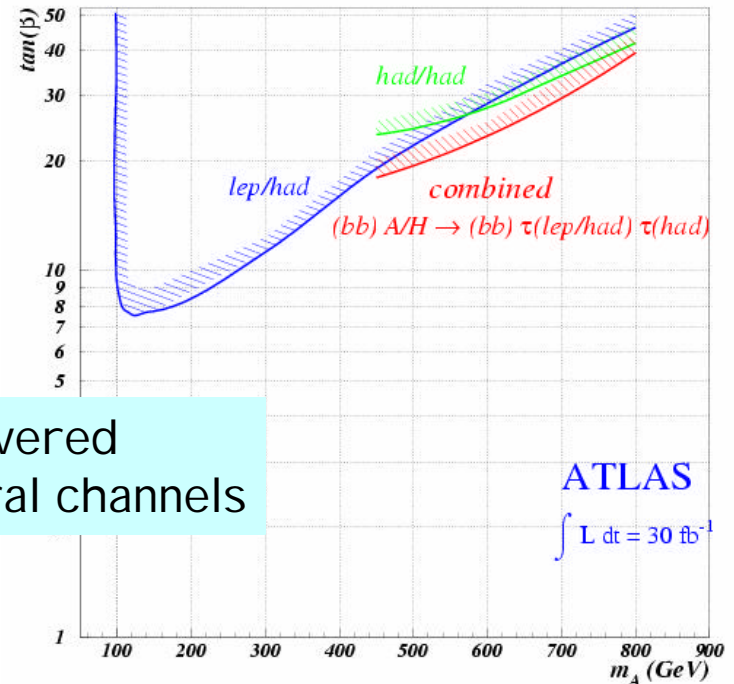
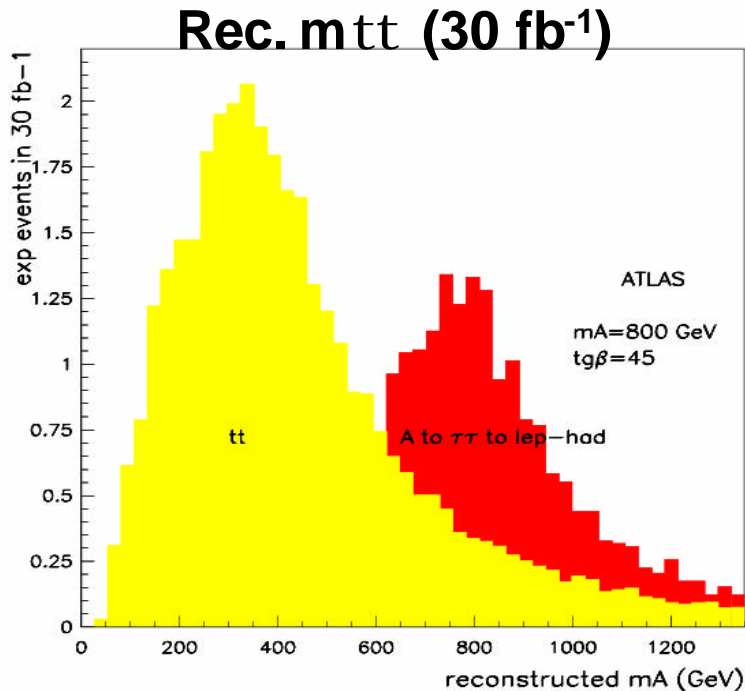
End-points  
 ↓  
 Relation between  
 Masses  
 ↓  
 All masses →



M( $\chi^0_1$ ) to 12% RMS → DM

# A, H decays in tt

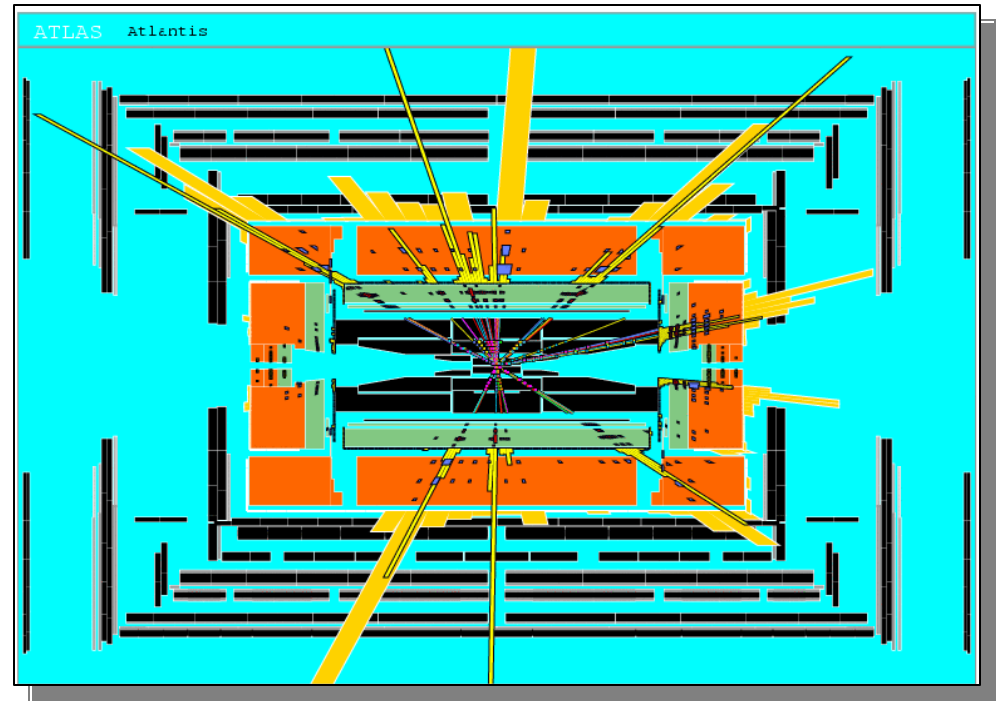
- In the MSSM there are 5 Higgs bosons : h, H, A, H<sup>±</sup>  
The 5 masses depend on 2 parameters: M(A) and tg(b)
- For large M(A) and tg(b) the H and A couplings to bb and t pairs are enhanced  
bb→bbA  
  - ↳ tt is a favored discovery channel (tt BR=11%, s=760 fb for tg(b)=45, M(A)&M(H)=800 GeV)
- LVL1 Triggers = (one tau-jet or lepton)⊕ETmiss



Plane covered  
by several channels

# An example of exotica :Black hole production..

- Another way of stabilizing the Higgs mass is to assume that there are extra space-dimensions, in which case the Planck mass could be as low as a few TeV.
- When  $\sqrt{S}$  reaches  $M_p$  Black Holes are produced, which decay “democratically” to quarks, leptons, Gauge bosons.



A typical event (?) could have:

hadrons/leptons/g,W,Z/Higgs ~ 75%/20%/3%/2%

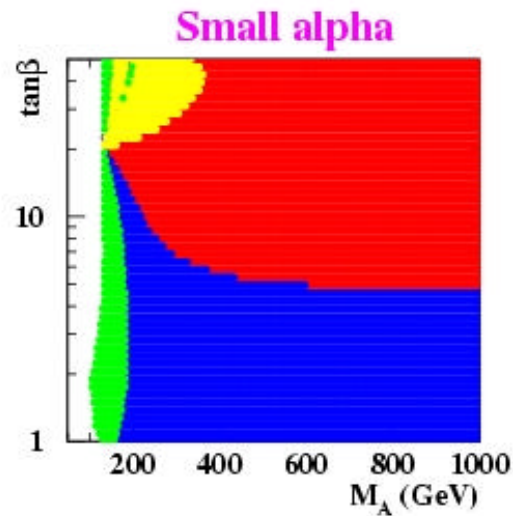
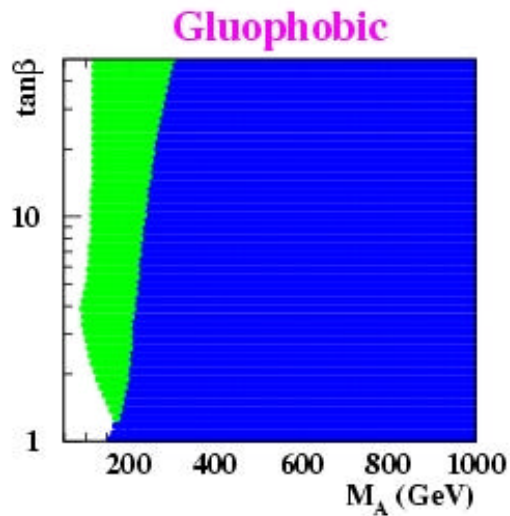
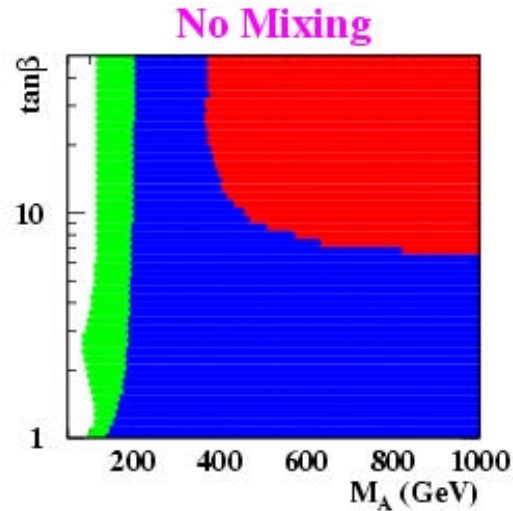
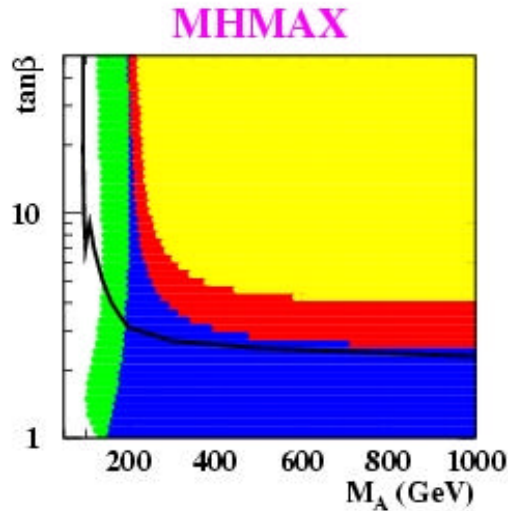
....a dream for particle identification at the LHC

# Conclusion-Summary

- Each of the 4 LHC experiment is a challenge
- Particle identification plays in all a major role
- Detectors have been carefully designed and are being built and tested with a lot of care.
- Tuning the triggers at start-up will be a crucial step
- A major issue-not discussed-will be the capacity of the experiments to “digest” the enormous amount of data they will produce, typically 1 Mbyte\*200Hz to mass storage
- More in 3+e years, when LHC starts...

# Back-up transparency

# Light Higgs boson $h$ : $300 \text{ fb}^{-1}$



**For high lumi.:**  
 $h \rightarrow gg, tth \rightarrow bb$   
 $h \rightarrow ZZ \rightarrow 4l$   
contribute

