

Distinguishing Non-minimal Higgs sectors via precise measurements of a light neutral Higgs boson

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LCWS2004 Paris, 22 April 2004

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In collaboration with S. Kanemura, Y. Okada and E. Senaha

2 Higgs doublet models

2 Higgs doublets models contain CP even fields and CP odd fields

$$\Phi_1^{\pm} = \begin{pmatrix} \phi_1^{\pm} \\ \phi_1^0 + ia_1 \end{pmatrix} \quad \Phi_2^{\pm} = \begin{pmatrix} \phi_2^{\pm} \\ \phi_2^0 + ia_2 \end{pmatrix}$$

- *CP conserving 2HDM:*

Neutral Higgs mass eigenstates are either pure CP even or pure CP odd:

$$H_i = f(\phi_i^0) \text{ or } H_i = f(a_i)$$

- *CP violating 2HDM:*

Neutral Higgs mass eigenstates are mixtures of CP even and CP odd fields:

$$H_i = f(\phi_i^0, a_i)$$

scalar-pseudoscalar mixing

MSSM Higgs potential

The most general 2HDM potential is given by:

$$\begin{aligned} V = & \mu_1^2(\Phi_1^+\Phi_1) + \mu_2^2(\Phi_2^+\Phi_2) + \lambda_1(\Phi_1^+\Phi_1)^2 + \\ & \lambda_2(\Phi_1^+\Phi_1)^2 + \lambda_3(\Phi_1^+\Phi_1)(\Phi_2^+\Phi_2) + \lambda_4|\Phi_1^+\Phi_2|^2 \\ & + \{m_{12}^2(\Phi_1^+\Phi_2) + h.c\} + [\lambda_5(\Phi_1^+\Phi_2)^2 + h.c] \\ & + [\lambda_6(\Phi_1^+\Phi_1)(\Phi_1^+\Phi_2) + h.c] + [\lambda_7(\Phi_2^+\Phi_2)(\Phi_2^+\Phi_1) + h.c] \end{aligned}$$

In the MSSM tree level potential (V_0) one has:

$$\begin{aligned} \mu_i^2 = & -m_i^2 - |\mu|^2 \quad , \quad \lambda_1 = \lambda_2 = -\frac{1}{8}(g^2 + g'^2) \\ \lambda_3 = & -\frac{1}{4}(g^2 - g'^2) \quad , \quad \lambda_4 = \frac{1}{2}g^2 \quad , \quad \lambda_5 = \lambda_6 = \lambda_7 = 0 \end{aligned}$$

No Explicit CP violation in V_0 (all real parameters)

In 1-loop potential V_1 :

$$\lambda_5, \lambda_6, \lambda_7 \neq 0$$

from radiative corrections from $t, \tilde{t}, b, \tilde{b}$ etc

Effective potential technique employed:

$$\lambda_5, \lambda_6, \lambda_7 = f(\mu, A_t, A_b, M_{SUSY} \dots)$$

$\lambda_5, \lambda_6, \lambda_7$ can be complex if SUSY parameters complex
→ scalar-pseudoscalar mixing

The neutral Higgs boson mass squared matrix in basis (A^0, ϕ_1, ϕ_2) :

$$M_N^2 = \begin{pmatrix} M_P^2 & M_{SP}^2 \\ M_{PS}^2 & M_S^2 \end{pmatrix}$$

where $M_{PS} = (M_{SP})^T$

In terms of λ_i :

$$M_{PS}^2 = v^2 \begin{pmatrix} \text{Im}(\lambda_5) \sin \beta + \text{Im}(\lambda_6) \cos \beta \\ \text{Im}(\lambda_5) \cos \beta + \text{Im}(\lambda_7) \sin \beta \end{pmatrix}$$

In CP conserving MSSM:

- $M_{PS}^2 = (0, 0)$, corresponding to real $\lambda_5, \lambda_6, \lambda_7$
- $M_P^2 = M_A^2$
- Pure CP-even h^0, H^0 and pure CP-odd A^0

In CP violating MSSM (i.e. $\arg(\mu), \arg(A_t) \neq 0$):

- $\text{Im}(\lambda_5), \text{Im}(\lambda_6), \text{Im}(\lambda_7) \neq 0$
- $O^T M_N^2 O = \text{Diag}(M_{H_1}^2, M_{H_2}^2, M_{H_3}^2)$, with $M_{H_1} < M_{H_2} < M_{H_3}$
- H_1, H_2, H_3 are mixed states of CP: $H_i = f(A^0, \phi_1, \phi_2)$

Magnitude of scalar-pseudoscalar mixing

Pilaftsis 98, Pilaftsis,Wagner 99, Choi, Lee,Drees 00

$$M_{PS}^2 = \left(\frac{m_t^4}{v^2} \frac{|\mu||A_t|}{32\pi^2 M_{SUSY}^2} \right) \sin \phi_{CP} \times f(M_{SUSY}, A_t, \mu, \tan \beta)$$

where $\phi_{CP} = \arg(A_t \mu)$

Optimal M_{PS}^2 requires:

- Large $|\mu|/M_{SUSY}$,
- Large $|A_t|/M_{SUSY}$
- moderate to large $\sin \phi_{CP}$

Constraints on $\sin \phi_{CP}$

$\sin \phi_{CP}(=arg(\mu A_t))$ strongly constrained by **EDMs** of the neutron and electron

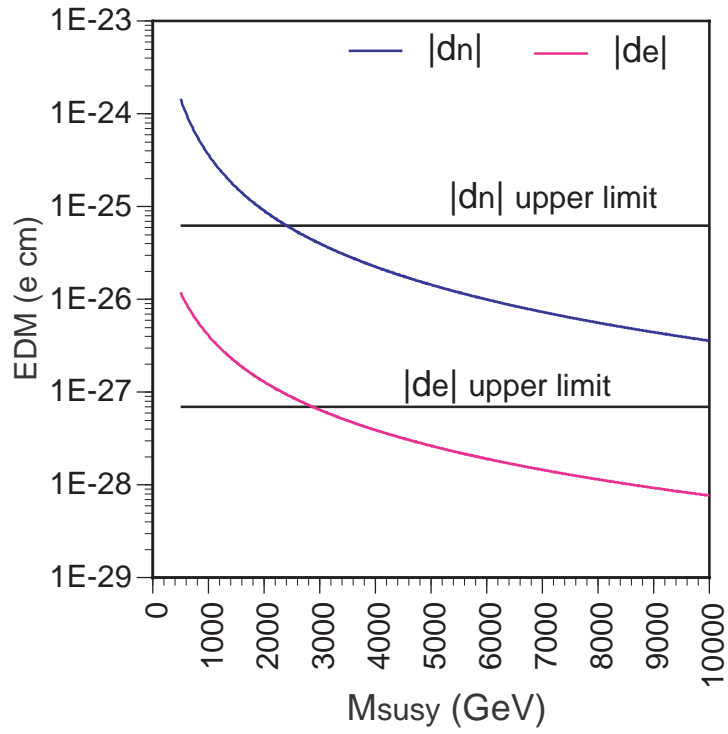
Phenomenology often done for $M_{SUSY} < 1000$ GeV case:

- $arg(\mu)$ strongly constrained by 1-loop EDM diagrams
- $arg(A_t)$ relatively unconstrained by 1-loop EDM \rightarrow 2-loop graphs important (Chang,Keung,Pilaftsis, 99)
- Large $\sin \phi_{CP}$ due to cancellation mechanism (Nath,Ibrahim 98) or fine tuning of phases

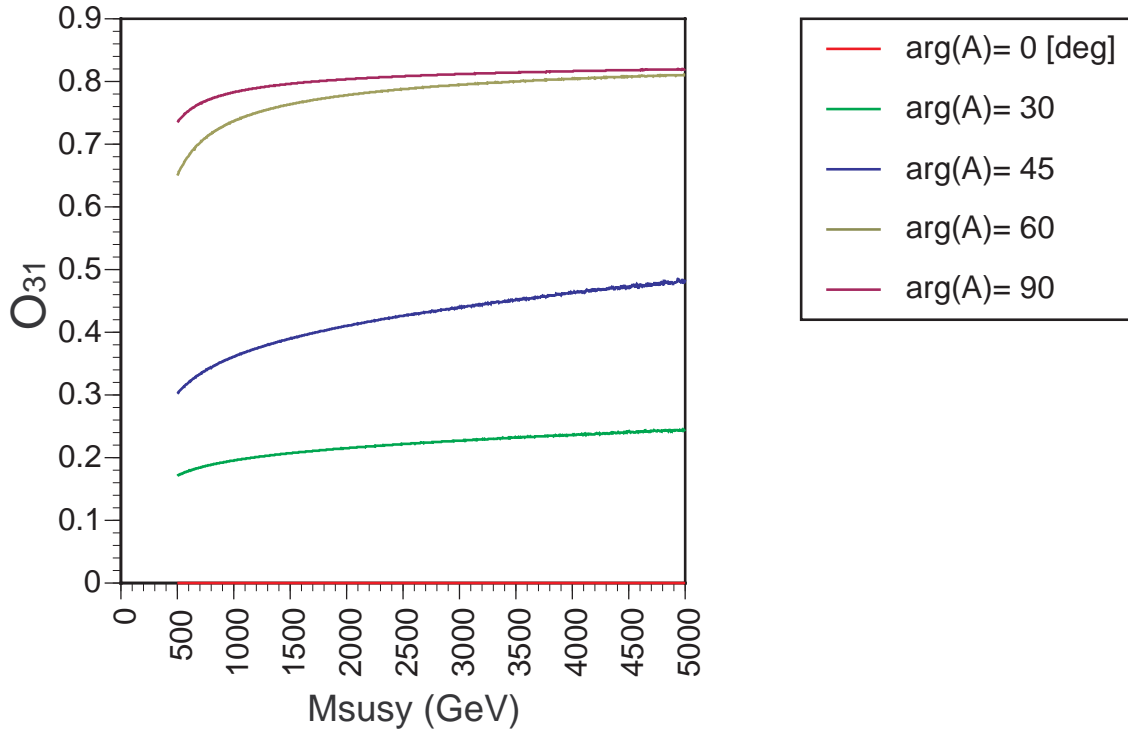
We are interested in $M_{SUSY} > 2000$ GeV case: (Ibrahim 01)

- Constraints on $arg(\mu), arg(A_t)$ are **weakened** $\sim 1/M_{SUSY}^2$
- CP violation in Higgs sector remains as $M_{SUSY} \rightarrow$ large
- Aim to see if pseudoscalar-scalar mixing can be **size-able** even if SUSY particles **out of range** at LHC/Linear Collider

EDMs as a function of M_{susy}



CP-odd component of H_1



Scalar-pseudoscalar mixing in H_1

We focus on H_1

(i) In MSSM with $\phi_{CP} = 0$:

$$H_1 = O_{11}\Phi_1 + O_{21}\Phi_2$$

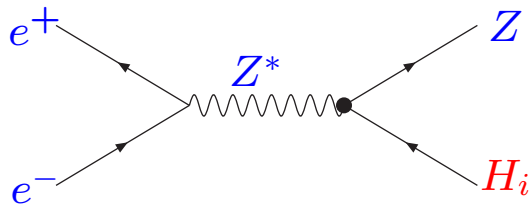
(ii) In MSSM with $\phi_{CP} \neq 0$:

$$H_1 = O_{11}\Phi_1 + O_{21}\Phi_2 + O_{31}A^0$$

	MSSM ($\phi_{CP} = 0$)	MSSM ($\phi_{CP} \neq 0$)
$H_1 b\bar{b}$	$O_{11}/\cos\beta$	$O_{11}/\cos\beta + O_{31}\tan\beta$
$H_1 VV$	$O_{11}\cos\beta + O_{21}\sin\beta$	$O_{11}\cos\beta + O_{21}\sin\beta$
$H_1 c\bar{c}$	$O_{21}/\sin\beta$	$O_{21}/\sin\beta + O_{31}\cot\beta$

Production at e^+e^- colliders

- Offers a clean environment in which to produce neutral Higgs bosons
- Is an ideal place to probe scalar-pseudoscalar mixing



$$\mathcal{L}_{H_i V V} = \frac{gm_W}{c_W^2} \sum_{i=1}^3 C_i H_i Z_\mu Z^\mu$$
$$C_i = O_{1i} \cos \beta + O_{2i} \sin \beta \quad ,$$

- There exists the following **sum rule**:

$$C_1^2 + C_2^2 + C_3^2 = 1$$

Observables sensitive to $\sin \phi_{CP}$

Precision measurements of H_1

(i) $\sigma(e^+e^- \rightarrow H_1 Z)$

(ii) $\text{BR}(H_1 \rightarrow b\bar{b})$

We aim to see if these observables can deviate from the $\sin \phi_{CP} = 0$ case for:

- **Large** $\sin \phi_{CP}$
- $M_{SUSY} > 2000 \text{ GeV}$

Numerical results

We take the following input parameters:

$$\begin{aligned}\widetilde{M}_Q &= \widetilde{M}_t = \widetilde{M}_b = M_{SUSY} = 3\text{TeV}, \\ |A_t| &= 1.5M_{SUSY}, |\mu| = 10M_{SUSY}, \tan\beta = 7 \\ 0 &< \text{Arg}(A_t) < 2\pi\end{aligned}$$

Fig.1: $\sigma(e^+e^- \rightarrow H_1 Z)$ as a function of m_{H_1} at $\sqrt{s} = 500$ GeV, for various ϕ_{CP} :

- $\sigma(e^+e^- \rightarrow H_1 Z)$ can be **very suppressed** for large ϕ_{CP}
- Suppression remains even for $m_{H_1} > 115$ GeV

Fig.2: $\sigma(e^+e^- \rightarrow H_1 Z)$ as a function of m_{H^\pm} at $\sqrt{s} = 500$ GeV, for various ϕ_{CP} :

- Sizeable deviation from SM rate for $\sigma(e^+e^- \rightarrow H_1 Z)$ for $m_{H^\pm} < 230$ GeV

Fig.3: g_{bb} against g_{WW} varying m_{H^\pm} for various ϕ_{CP} :

- g_{bb} for large $\sin\phi_{CP}$ can be much larger than for $\sin\phi_{CP} = 0$ case

Fig.1: $\sigma(e^+e^- \rightarrow H_1 Z)$ as a function of m_{H_1} at $\sqrt{s} = 500$ GeV, for various ϕ_{CP} :

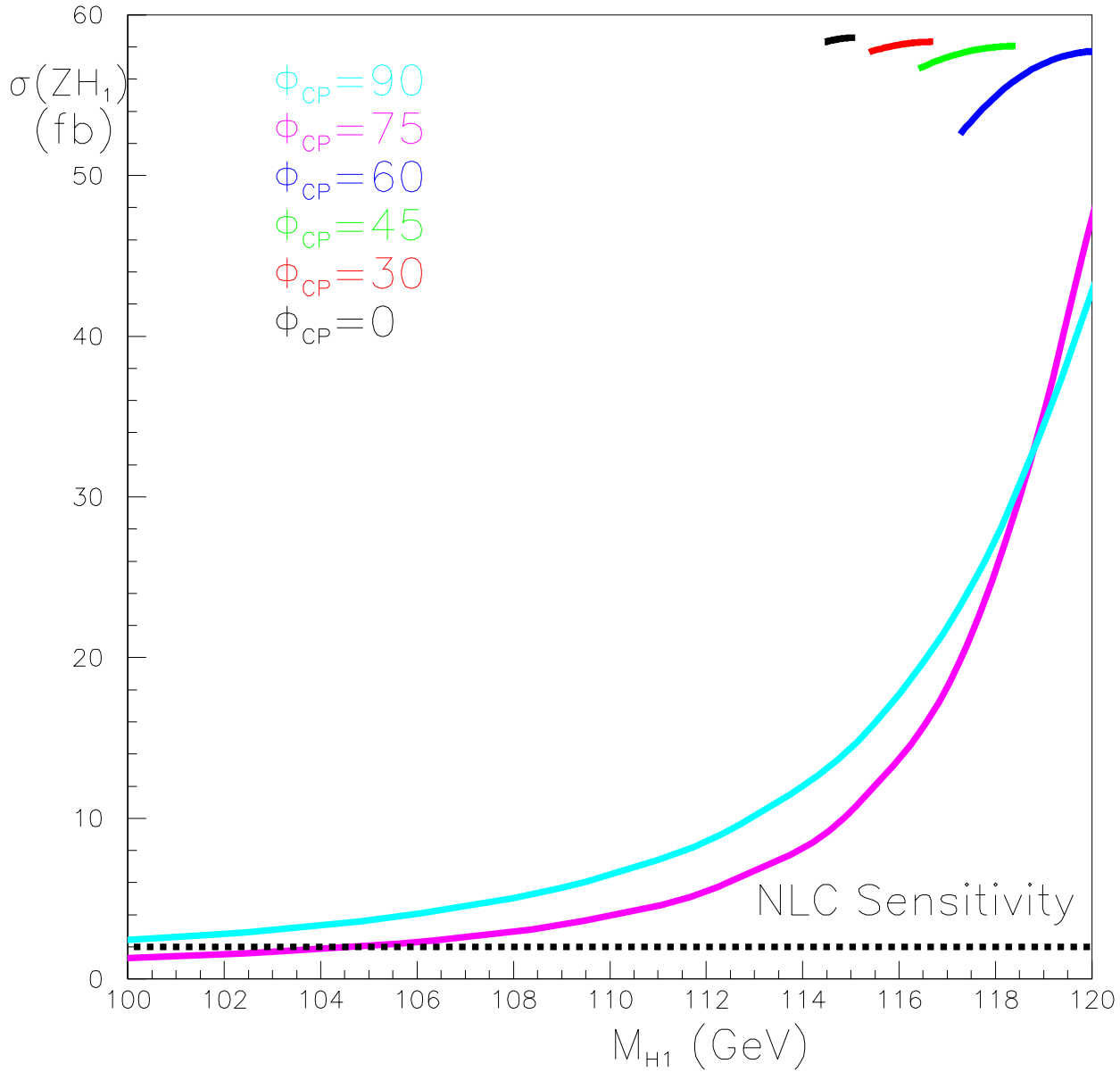


Fig.2: $\sigma(e^+e^- \rightarrow H_1 Z)$ as a function of m_{H^\pm} at $\sqrt{s} = 500$ GeV, for various ϕ_{CP} :

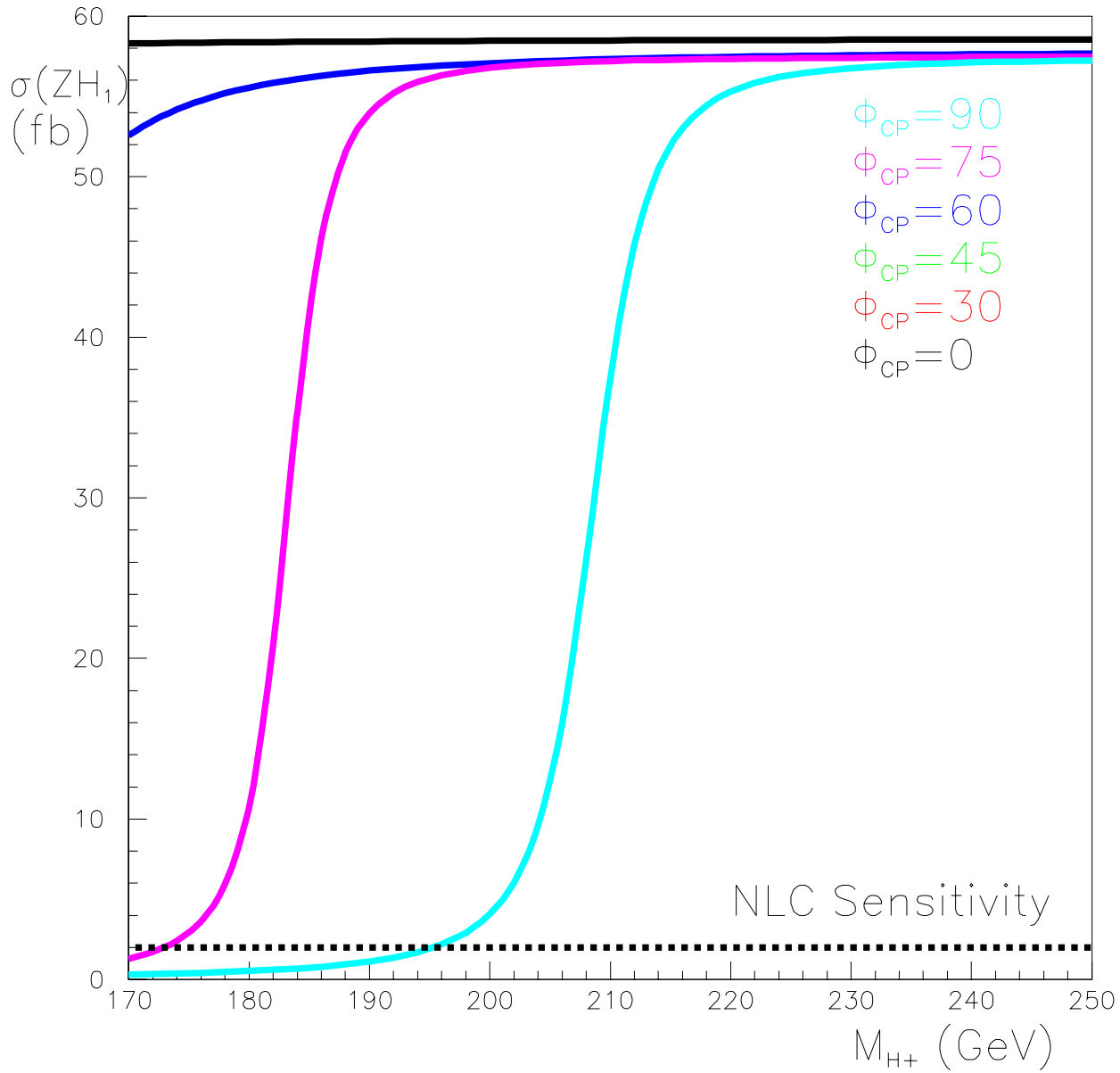
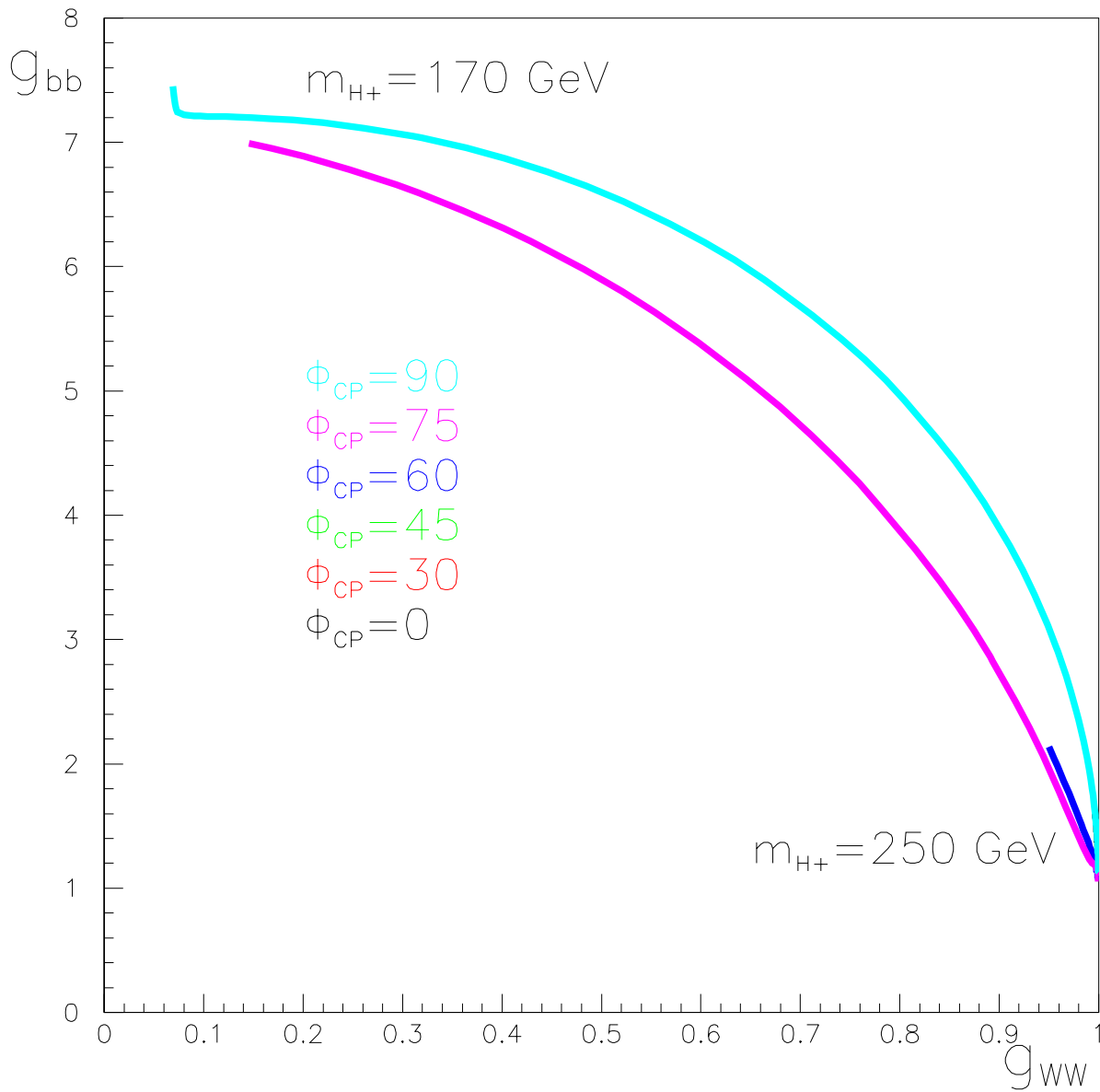


Fig.3: g_{bb} against g_{WW} varying m_{H^\pm} for various ϕ_{CP} :



Conclusions

- SUSY phase $\phi_{CP} \neq 0 \rightarrow$ Scalar-pseudoscalar mixing
- Large A_t and μ required
- $M_{SUSY} > 2000$ GeV relaxes EDM constraints on ϕ_{CP}
- Large scalar-pseudoscalar mixing possible
- $\sigma(e^+e^- \rightarrow ZH_1)$ can be sizeably **suppressed**, even for $m_h \geq 115$ GeV
- Very different from MSSM with $\phi_{CP} = 0$