

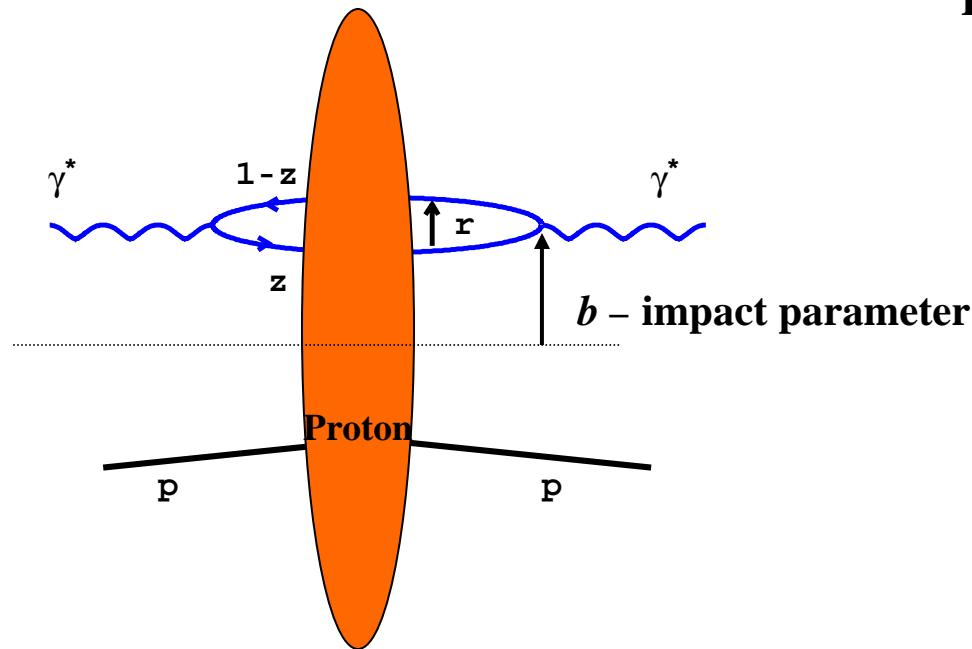
Multiple Interaction in DIS from AGK rules

Evaluation within a Dipole Model

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DESY
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HERA-LHC Workshop

Impact Parameter Dipole Saturation Model

H. Kowalski
 D. Teaney
[hep-ph/0304189](https://arxiv.org/abs/hep-ph/0304189)



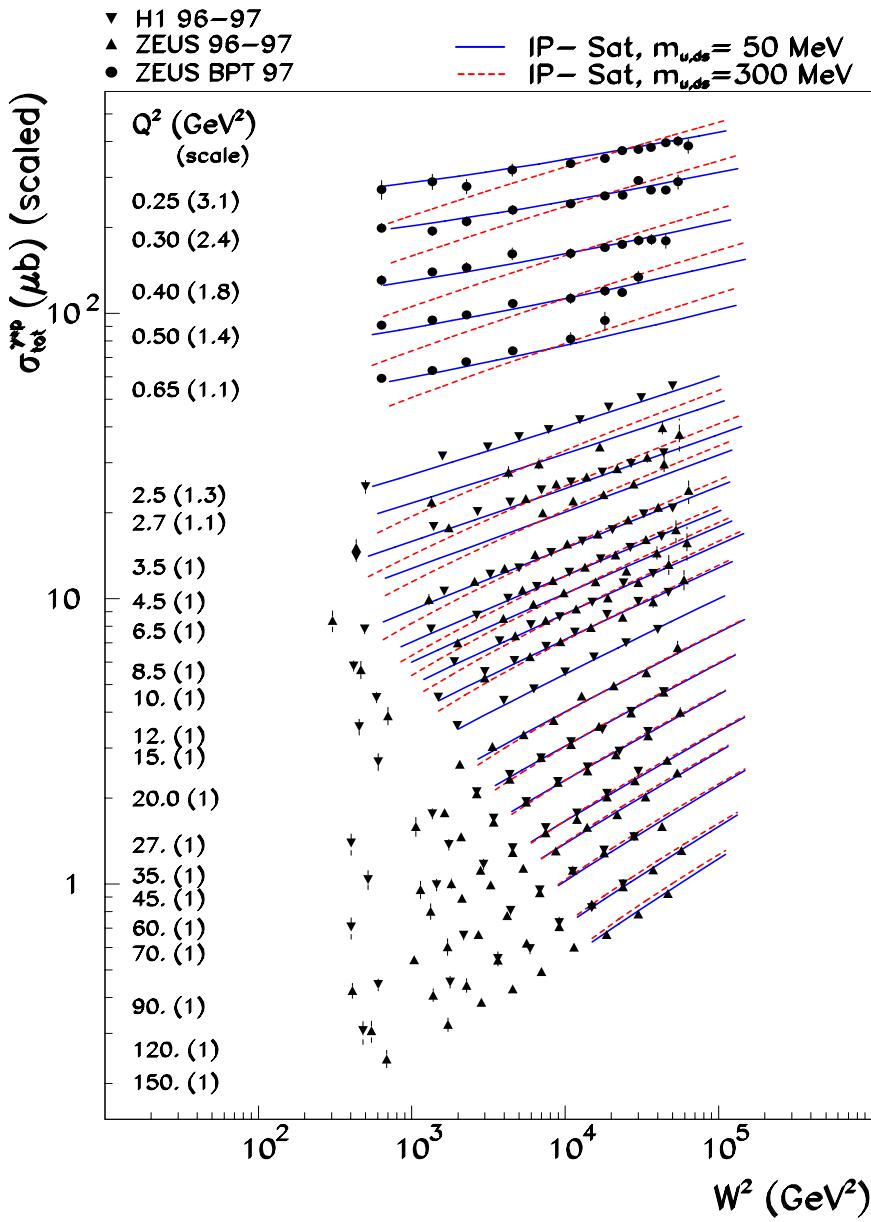
$$\sigma_{tot}^{\gamma^* p} = \int d^2 \vec{b} d^2 \hat{\vec{r}} \int_0^1 dz \Psi(Q^2, z, \vec{r})^* \sigma_{q\bar{q}}(x, r^2, b) \Psi(Q^2, z, \vec{r})$$

$$\frac{d\sigma_{qq}(x, r)}{d^2 b} = 2 \cdot \left\{ 1 - \exp\left(-\frac{\pi^2}{2 \cdot 3} r^2 \alpha_s(\mu^2) x g(x, \mu^2) T(b)\right) \right\}$$

$$T(b) = \frac{1}{2\pi B_G} \exp(-b^2/2 B_G)$$

$T(b)$ - proton shape

$$\sigma^{\gamma^* p} = \sum_f \int d^2\vec{r} \int d^2b \int_0^1 dz \Psi_f^*(Q^2, z, \vec{r}) \left[1 - \exp\left(-\frac{\pi^2}{2 \cdot 3} r^2 \alpha_s x g(x, \mu^2) T(b)\right) \right] \Psi_f(Q^2, z, \vec{r})$$



$$\mu^2 = \frac{C}{r^2} + \mu_0^2$$

$$xg(x, \mu_0^2) = A_g \left(\frac{1}{x}\right)^{\lambda_g} (1-x)^{5.6}$$

$Q^2 > 0.25 \text{ GeV}^2$
 $m_u = 0.05 \text{ GeV}$
 $m_c = 1.30 \text{ GeV}$

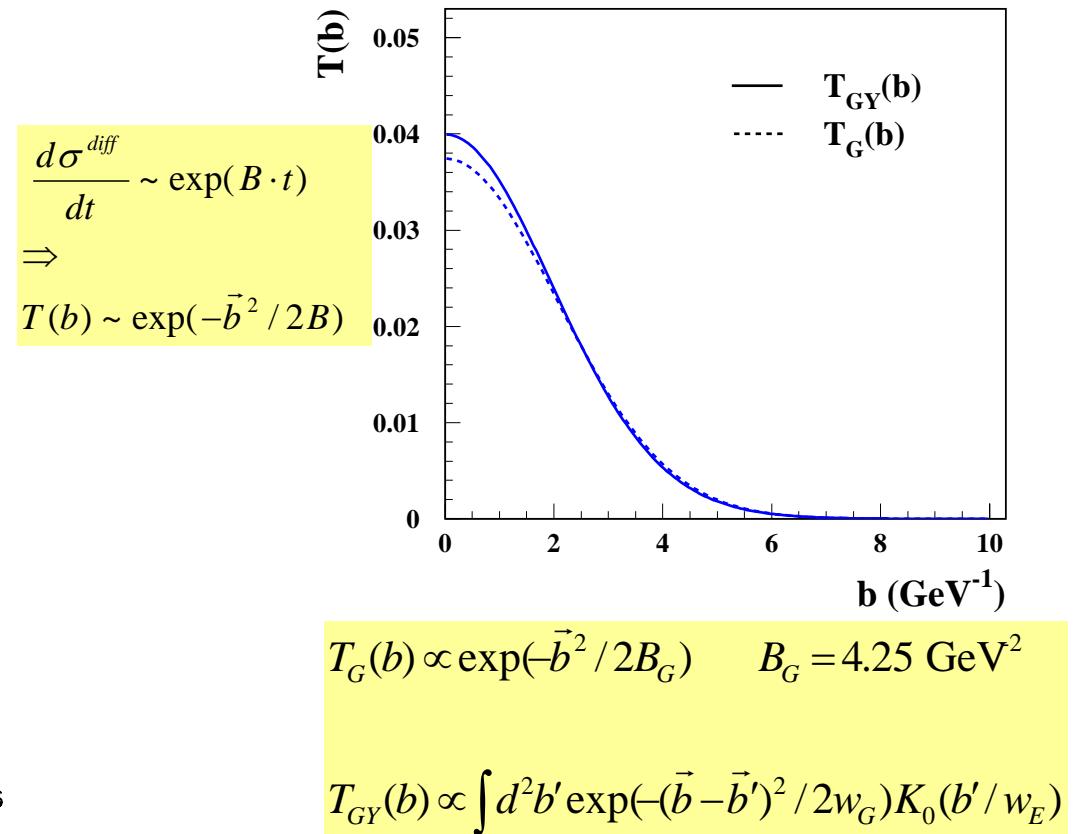
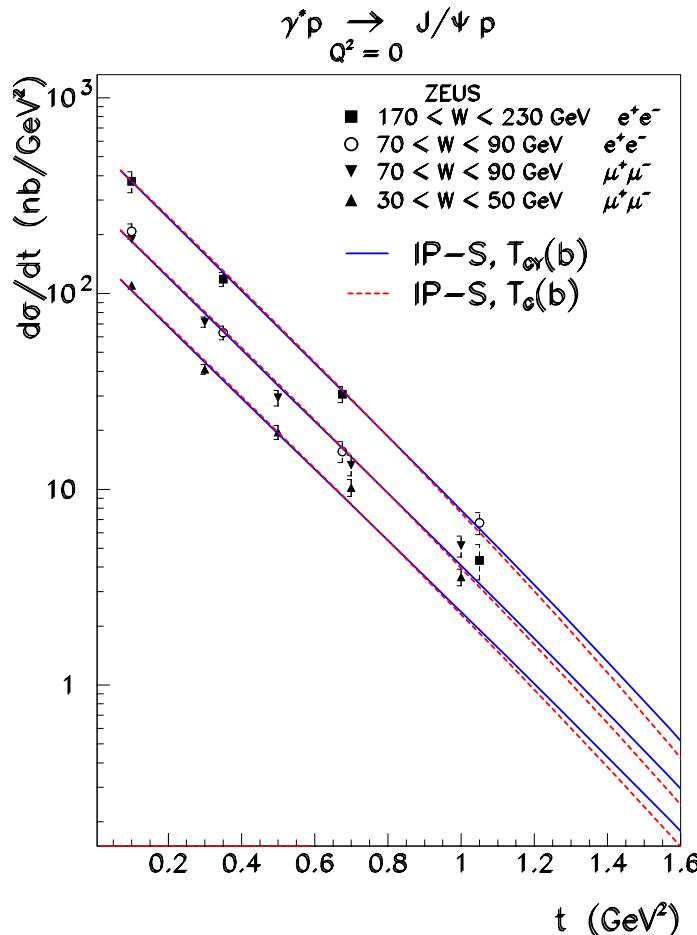
Fit parameters
 $\lambda_g = -0.12$ $C = 4.0$
 $Q_0^2 = 0.8 \text{ GeV}^2$
 $\chi^2/N = 0.8$

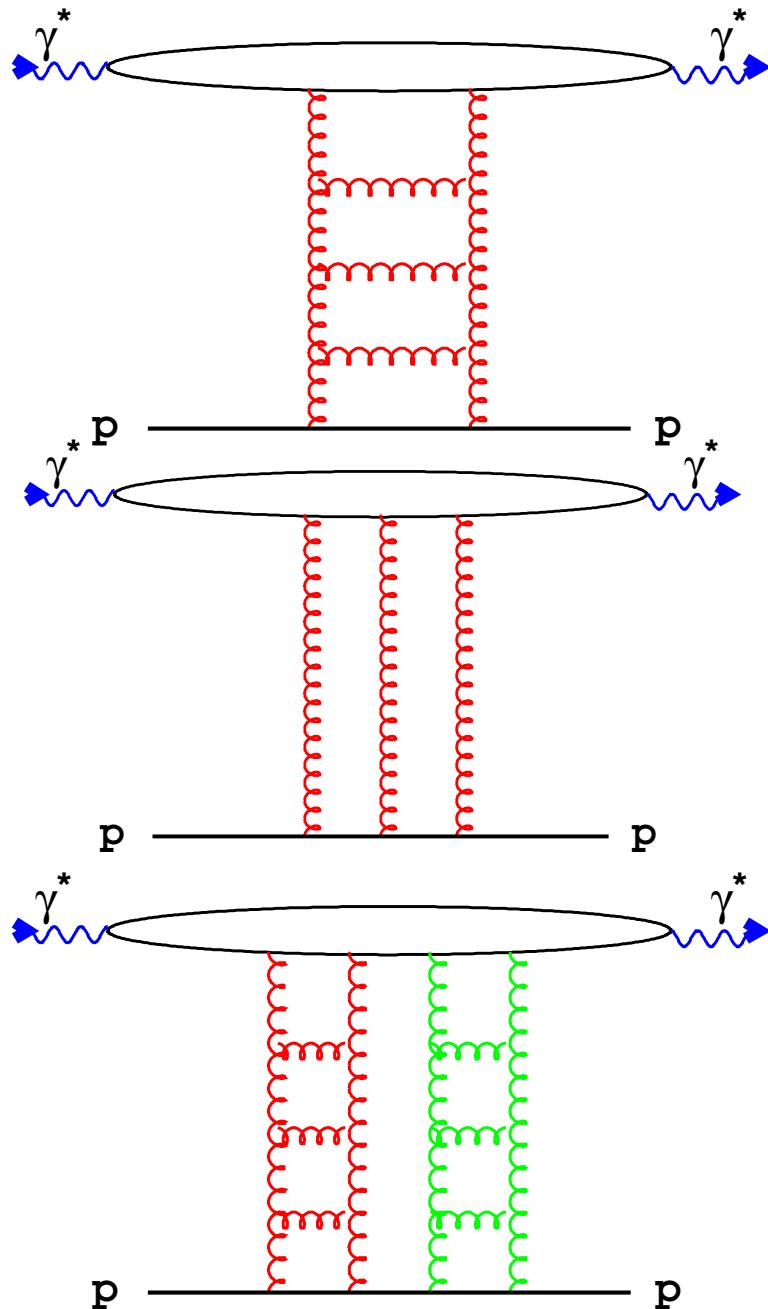
$x < 10^{-2}$

t-dependence of the diffractive cross sections determines the \vec{b} distribution

$$t = -\vec{\Delta}^2 \quad \vec{\Delta} \text{ - transv. momentum (2-d)} \quad \vec{b} \text{ - impact parameter (2-d)}$$

$$\frac{d\sigma_{VM}^{r^* p}}{dt} = \frac{1}{16\pi} \left| \int d^2\vec{r} \int d^2b e^{-i\vec{b}\vec{\Delta}} \int_0^1 dz \Psi_{VM}^*(Q^2, z, \vec{r}) 2 \left\{ 1 - \exp\left(-\frac{\pi^2}{2 \cdot 3} r^2 \alpha_s x g(x, \mu^2) T(b)\right) \right\} \Psi(Q^2, z, \vec{r}) \right|^2$$





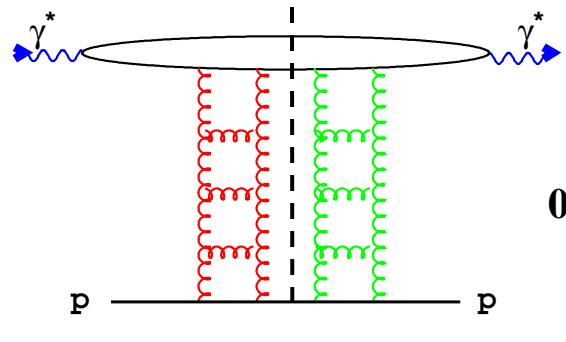
t-channel picture

**Color singlet dominates over octet
in the 2-gluon exchange amplitude
at high energies**

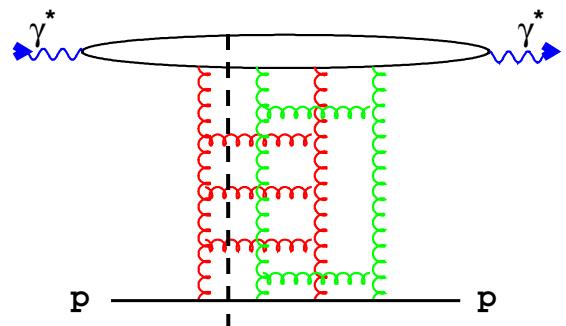
**3-gluon exchange amplitude is suppressed
at high energies**

**2-gluon pairs in color singlet (Pomerons)
dominate the multi-gluon QCD amplitudes
at high energies**

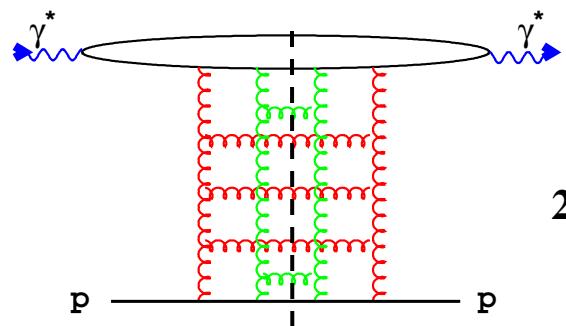
2-Pomeron exchange in QCD



0-cut

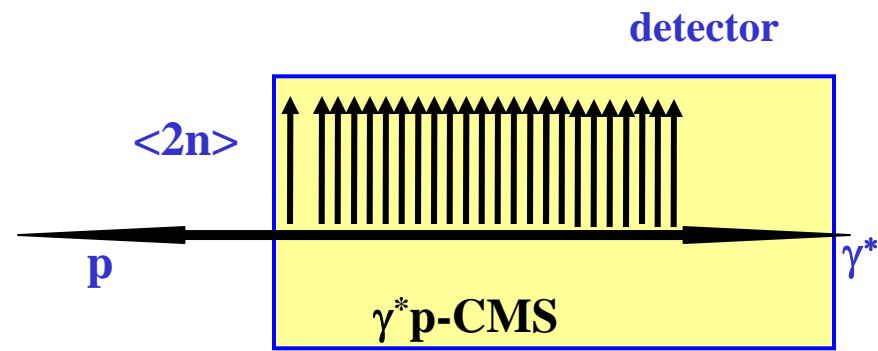
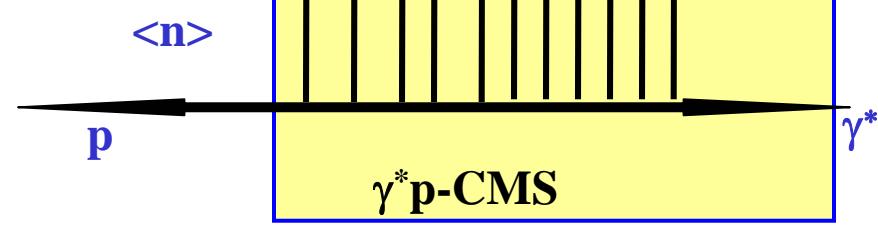
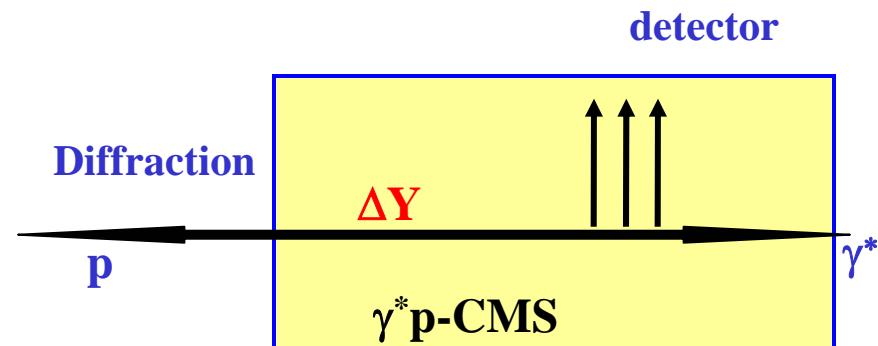


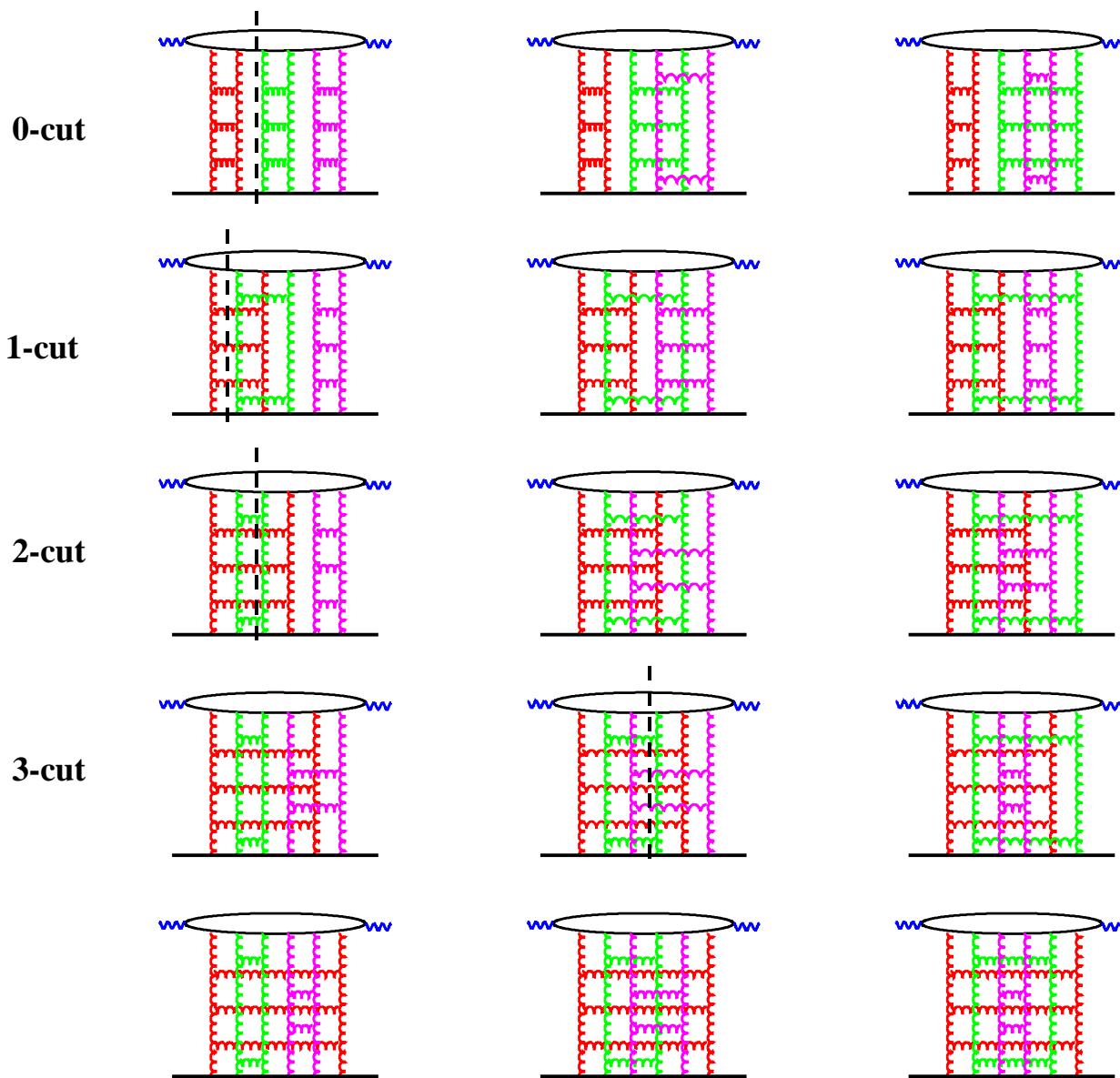
1-cut



2-cut

Final States (naïve picture)





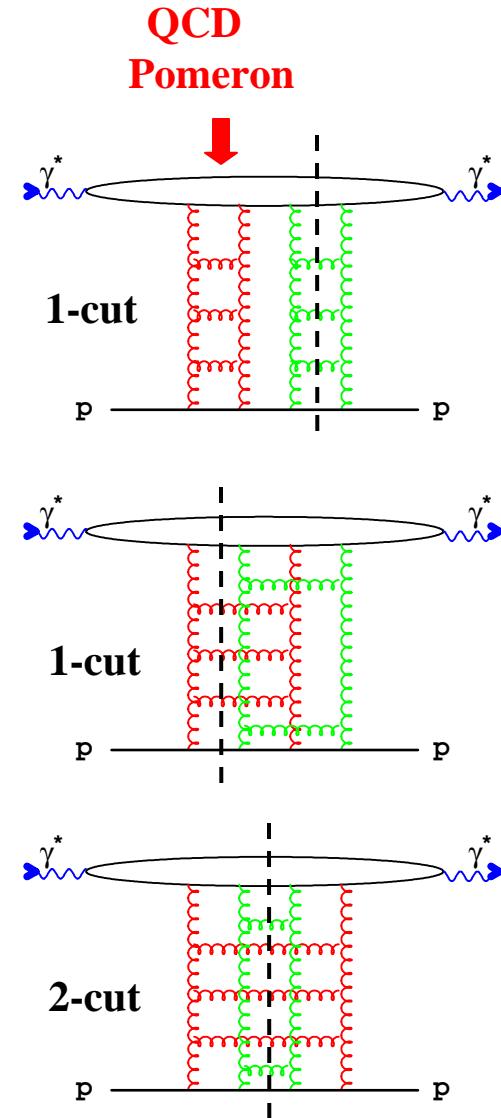
AGK Rules

The cross-section for k-cut pomerons:

Abramovski, Gribov, Kancheli
Sov. J., Nucl. Phys. 18, p308 (1974)

$$\sigma_k = \sum_{m=k}^{\infty} (-1)^{m-k} 2^m \frac{m!}{k!(m-k)!} F^{(m)}$$

$F^{(m)}$ – amplitude for the exchange of
 m Pomerons



AGK Rules in the Dipole Model

Total cross section

Mueller-Salam (NP B475, 293)

$$\sigma_{tot} = 2 \sum_{m=1}^{\infty} (-1)^{m-1} F^{(m)}$$

Dipole cross section

$$\frac{d\sigma}{d^2b} = 2(1 - \exp(-\Omega/2)) = 2 \sum_{m=1}^{\infty} (-1)^{m-1} \left(\frac{\Omega}{2}\right)^m \frac{1}{m!}$$

Amplitude for the exchange of m pomerons in the dipole model

→ $F^{(m)} = \left(\frac{\Omega}{2}\right)^m \frac{1}{m!}$ $\Omega = \frac{\pi^2}{N_C} r^2 \alpha_s(\mu^2) x g(x, \mu^2) T(b)$ KT model

AGK rules

$$\frac{d\sigma_k}{d^2b} = \sum_{m=k}^{\infty} (-1)^{m-k} 2^m \frac{m!}{k!(m-k)!} F^{(m)}$$

Dipole model

$$\frac{d\sigma_k}{d^2b} = \sum_{m=k}^{\infty} (-1)^{m-k} 2^m \frac{m!}{k!(m-k)!} \left(\frac{\Omega}{2}\right)^m \left(\frac{1}{m!}\right) = \frac{\Omega^k}{k!} \sum_{m=k}^{\infty} (-1)^{m-k} \frac{\Omega^{m-k}}{(m-k)!}$$



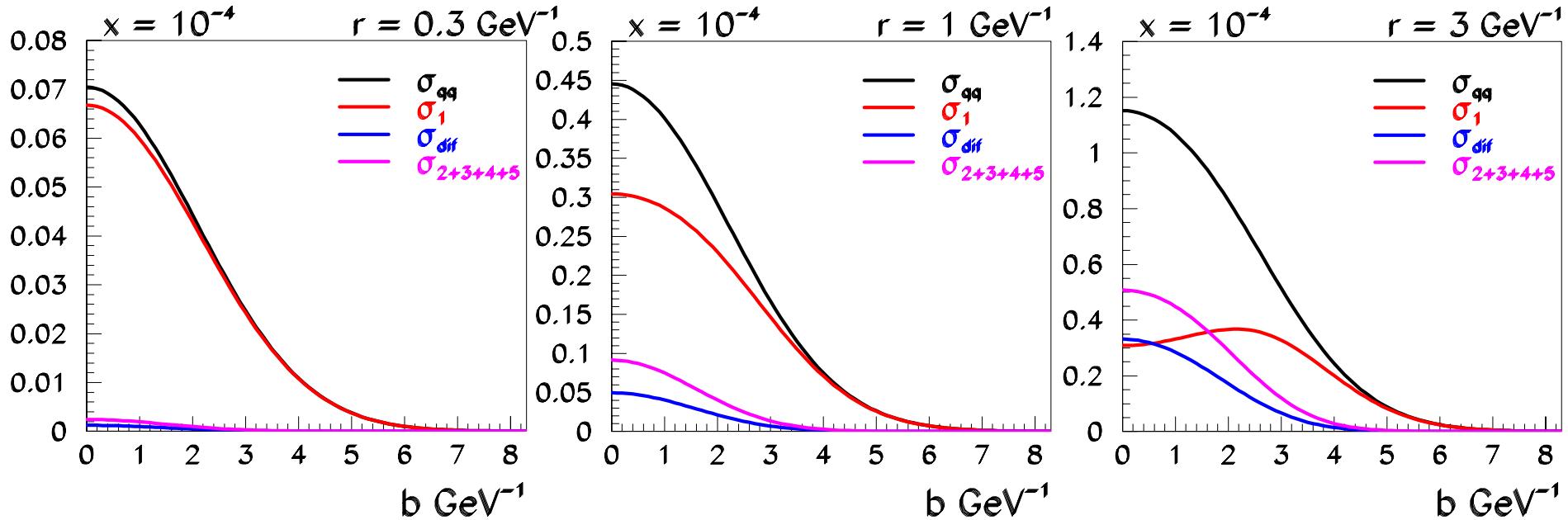
$$\frac{d\sigma_k}{d^2b} = \frac{\Omega^k}{k!} \exp(-\Omega)$$

Diffraction from AGK rules

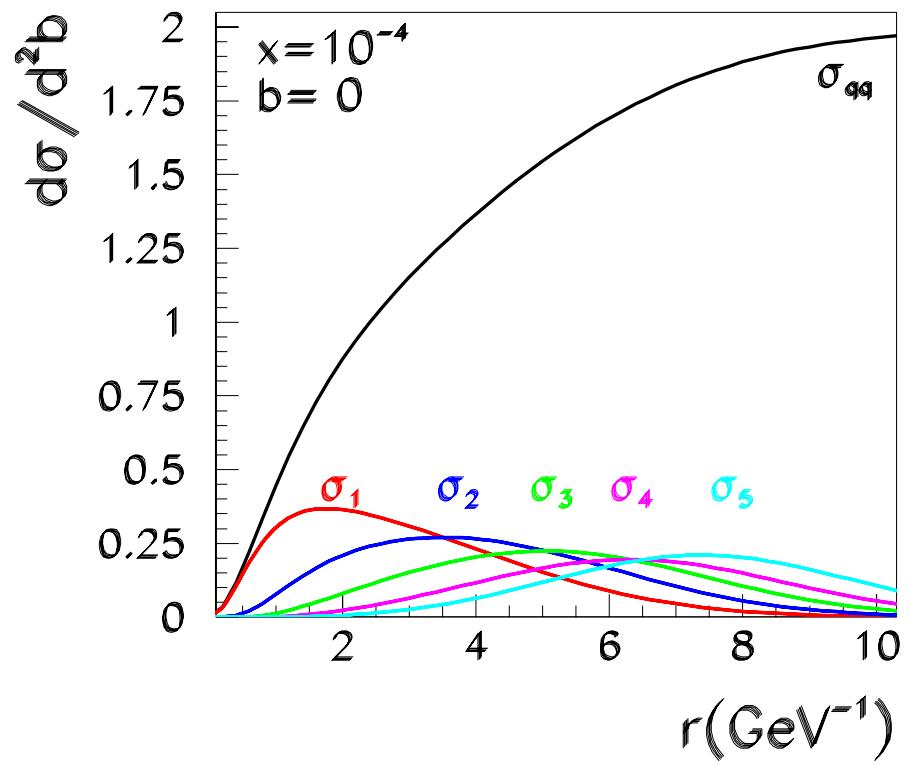
$$\begin{aligned} \frac{d\sigma_{diff}}{d^2b} &= \frac{d\sigma_{qq}}{d^2b} - \sum_{k=1}^{\infty} \frac{d\sigma_k}{d^2b} = 2(1 - \exp(-\Omega/2)) - (1 - \exp(-\Omega)) \\ &= 1 - 2 \exp(-\Omega/2) + \exp(-\Omega) = (1 - \exp(-\Omega/2))^2 \end{aligned}$$

$$\frac{d\sigma_{qq}}{d^2b} = 2 \cdot \left\{ 1 - \exp\left(-\frac{\Omega}{2}\right) \right\}$$

$$\frac{d\sigma_k}{d^2b} = \frac{\Omega^k}{k!} \exp(-\Omega)$$



$$\Omega = \frac{\pi^2}{N_C} r^2 \alpha_s(\mu^2) x g(x, \mu^2) T(b)$$

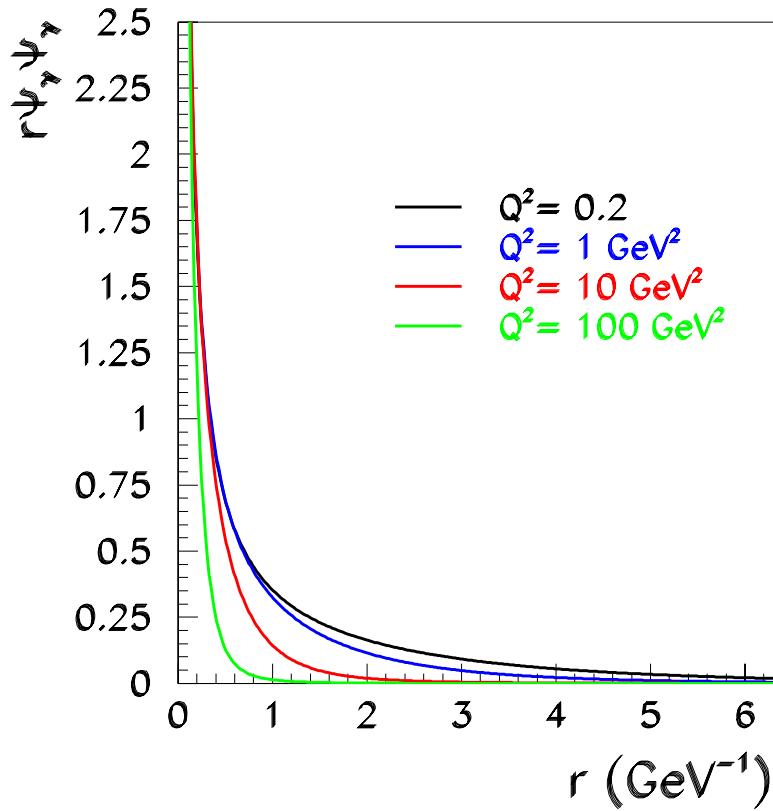


$$\sigma_{T,L}^{\gamma^* p}(x, Q^2) = \int_0^\infty d^2 \vec{r} \int_0^1 dz \sum_f \left| \Psi_{T,L}^f(\vec{r}, z, Q^2) \right|^2 \sigma_{qq}(x, r)$$

$$\left| \Psi_T^f(\vec{r}, z, Q^2) \right|^2 = \frac{3\alpha_{em}}{2\pi^2} e_q^2 \{ [z^2 + (1-z)^2] \varepsilon^2 K_1^2(\varepsilon r) + m_q^2 K_0^2(\varepsilon r) \}$$

$$\left| \Psi_L^f(\vec{r}, z, Q^2) \right|^2 = \frac{3\alpha_{em}}{2\pi^2} e_q^2 \{ 4Q^2 z^2 (1-z)^2 K_0^2(\varepsilon r) \} \quad \varepsilon^2 = z(1-z)Q^2 + m_q^2$$

$\gamma^* \gamma^*$ Overlap



$$K_1(\varepsilon r) = 1/\varepsilon r$$

for $\varepsilon r \ll 1$



$$Q^2 \sim 1/r^2$$

$$K_1(\varepsilon r) = \sqrt{\pi/2x} \exp(-\varepsilon r)$$

for $\varepsilon r \gg 1$

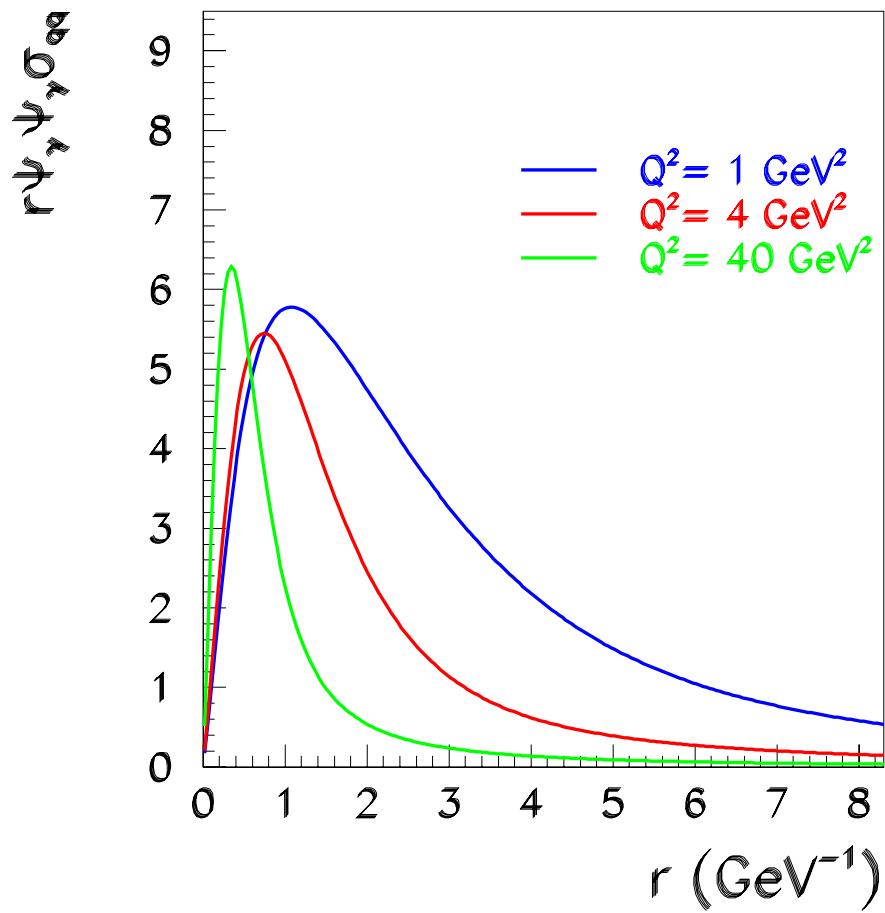


$$\exp(-m_q r)$$

All quarks

$$2\pi r \int_0^1 dz \sum_f \left| \Psi_{T,L}^f(\vec{r}, z, Q^2) \right|^2 \sigma_{qq}(x, r)$$

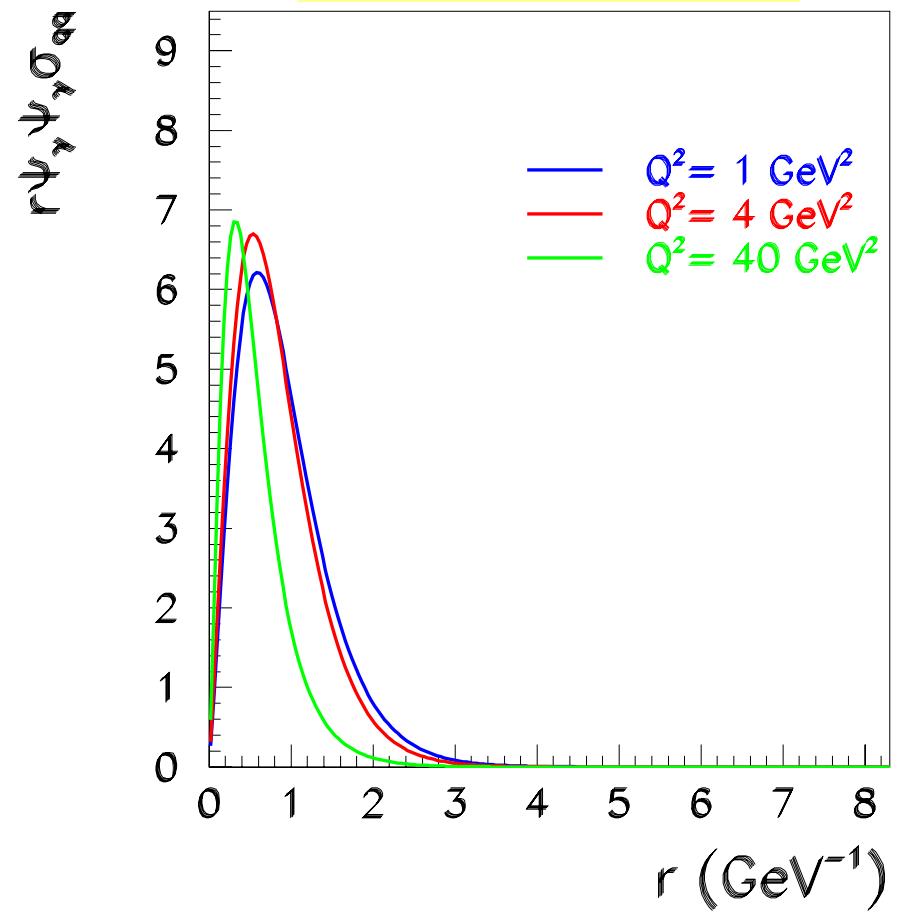
$$m_{u,d,s} = 100 \text{ MeV} \quad m_c = 1.3 \text{ GeV}$$

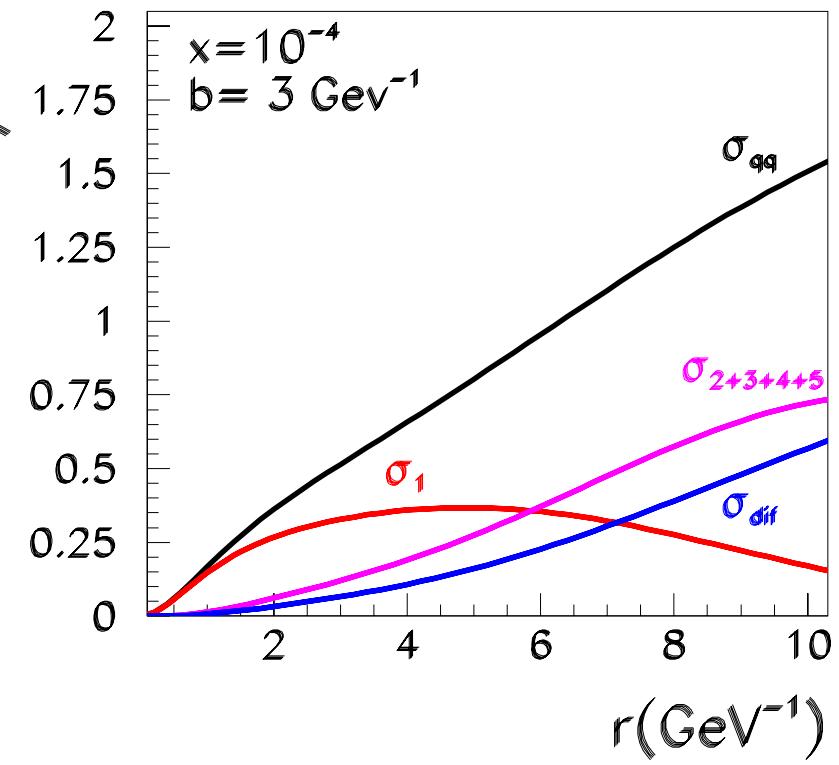
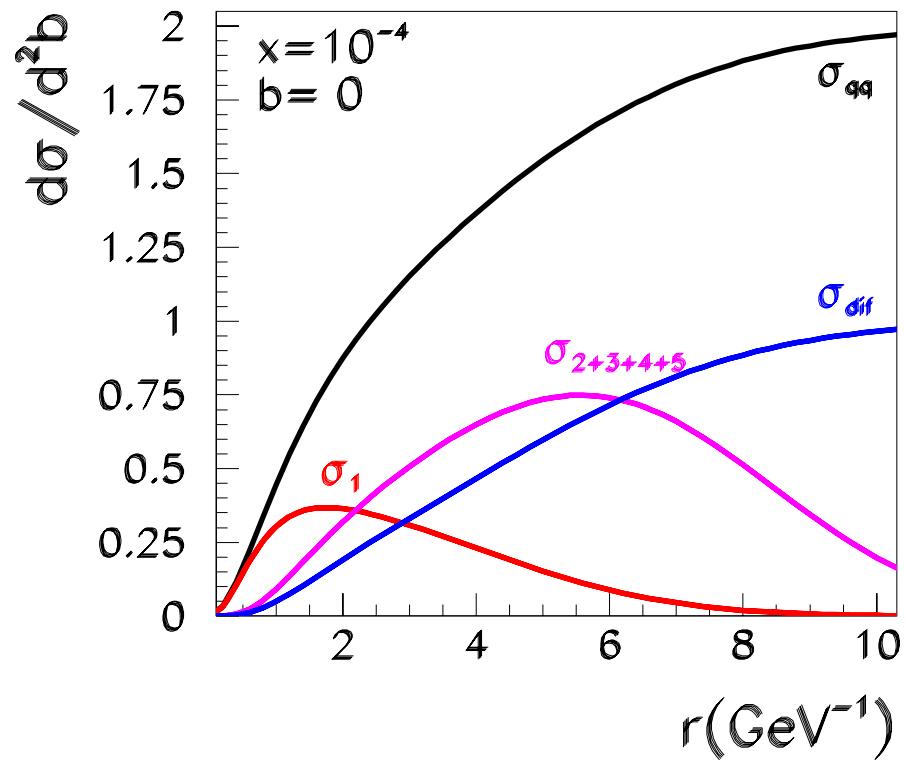


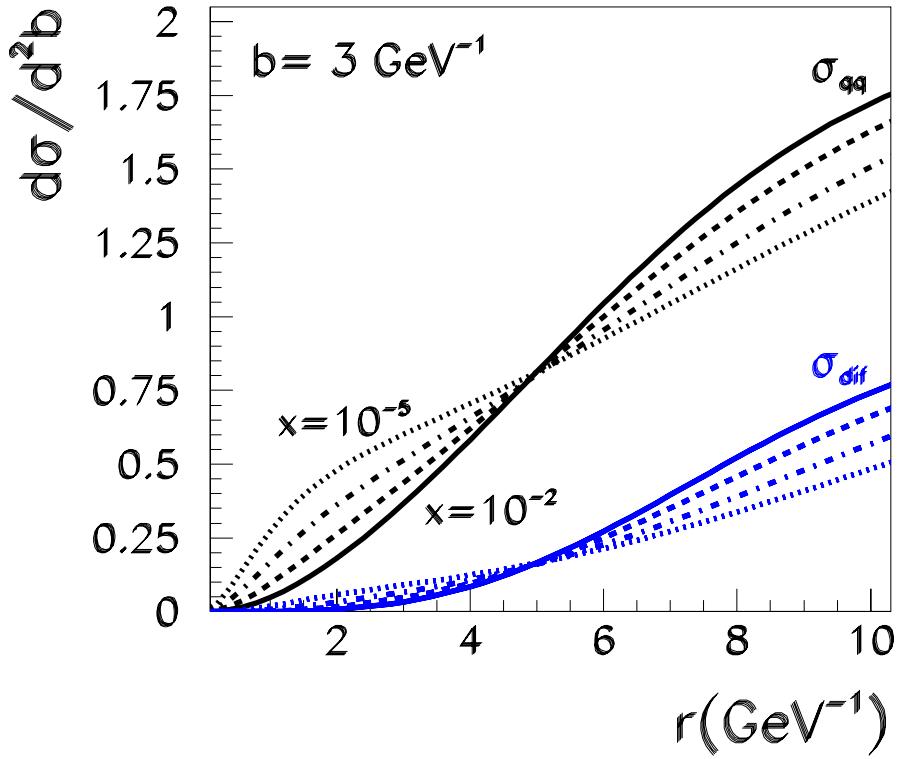
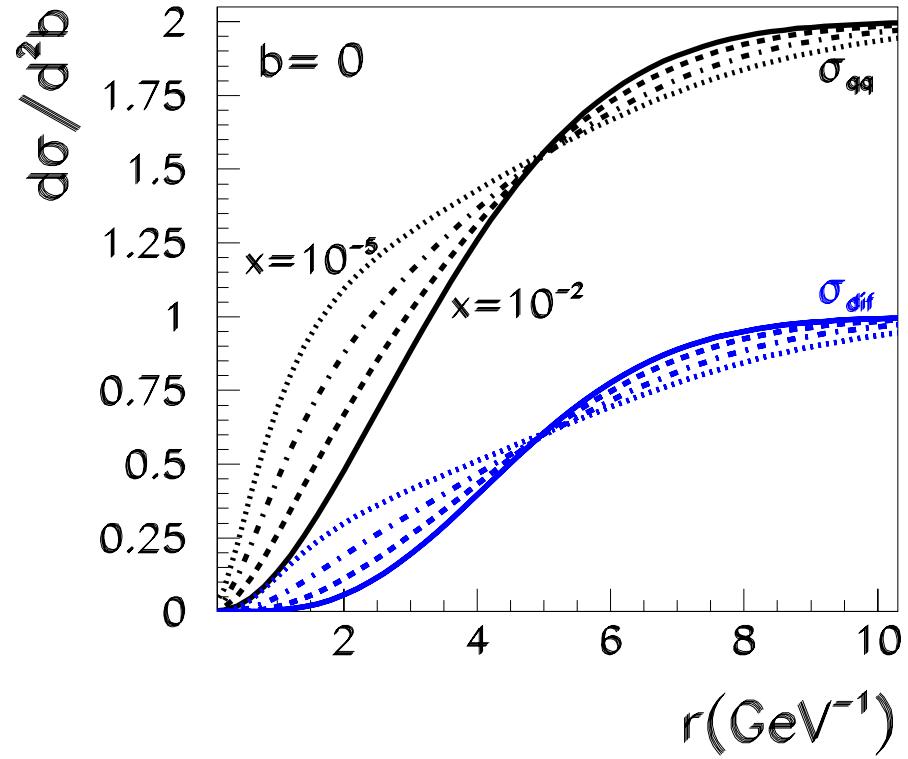
Charmed quark

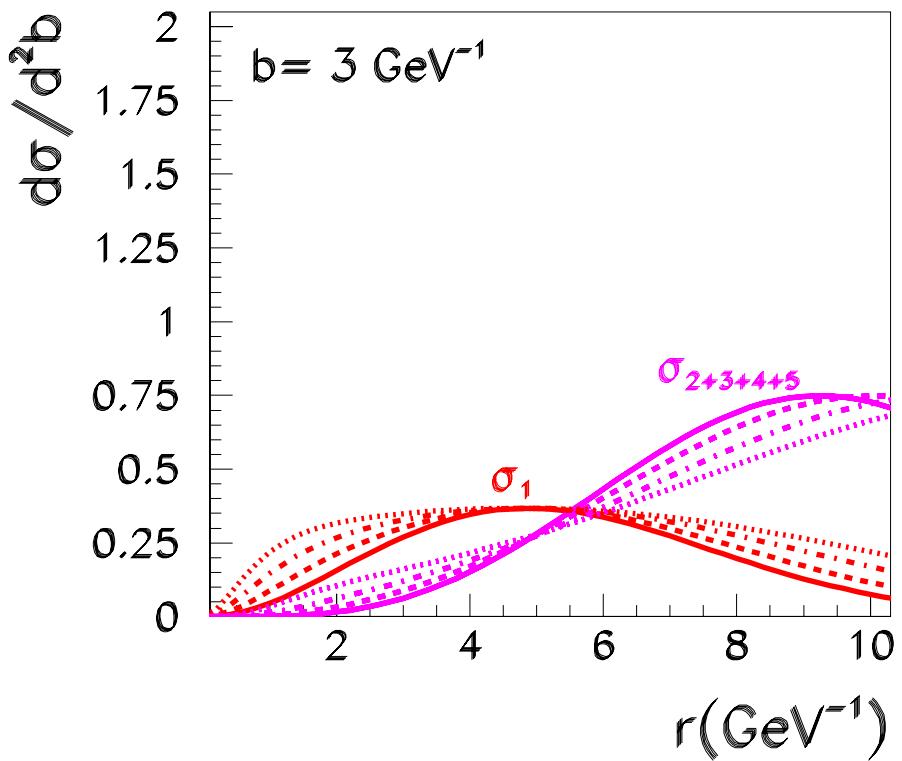
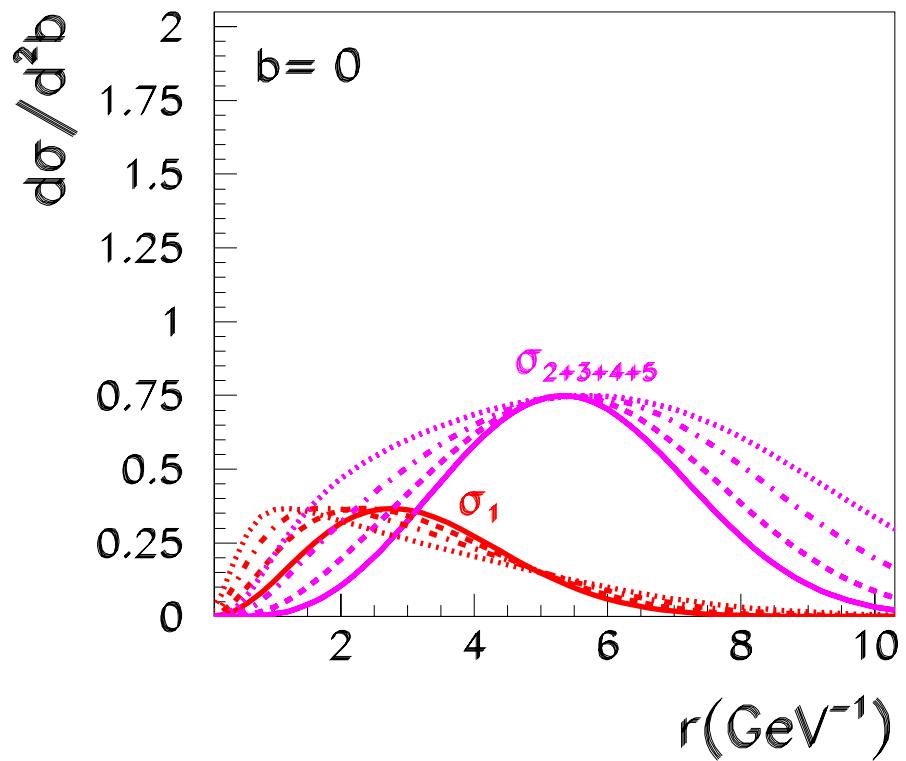
$$2\pi r \int_0^1 dz \sigma_{qq}(x, r) \Psi_{T,L}^c(\vec{r}, z, Q^2)^2$$

$$m_c = 1.3 \text{ GeV}$$



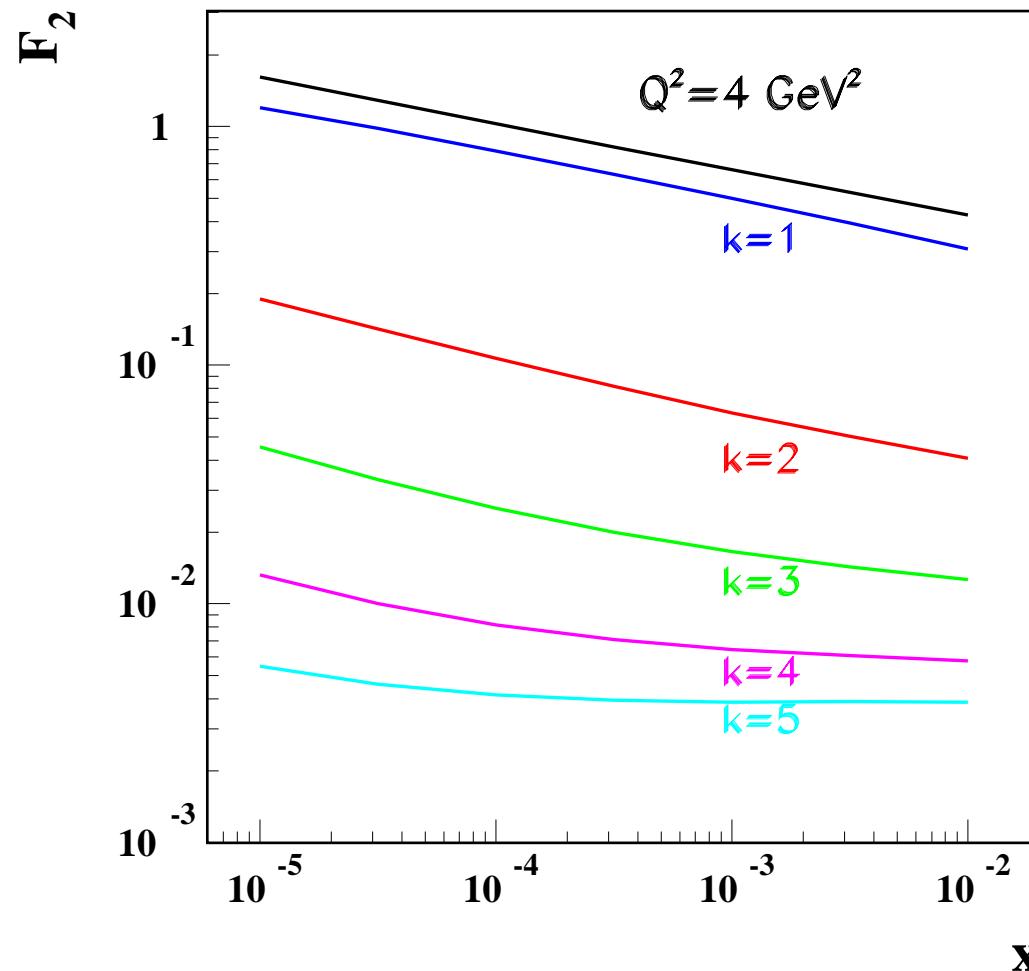


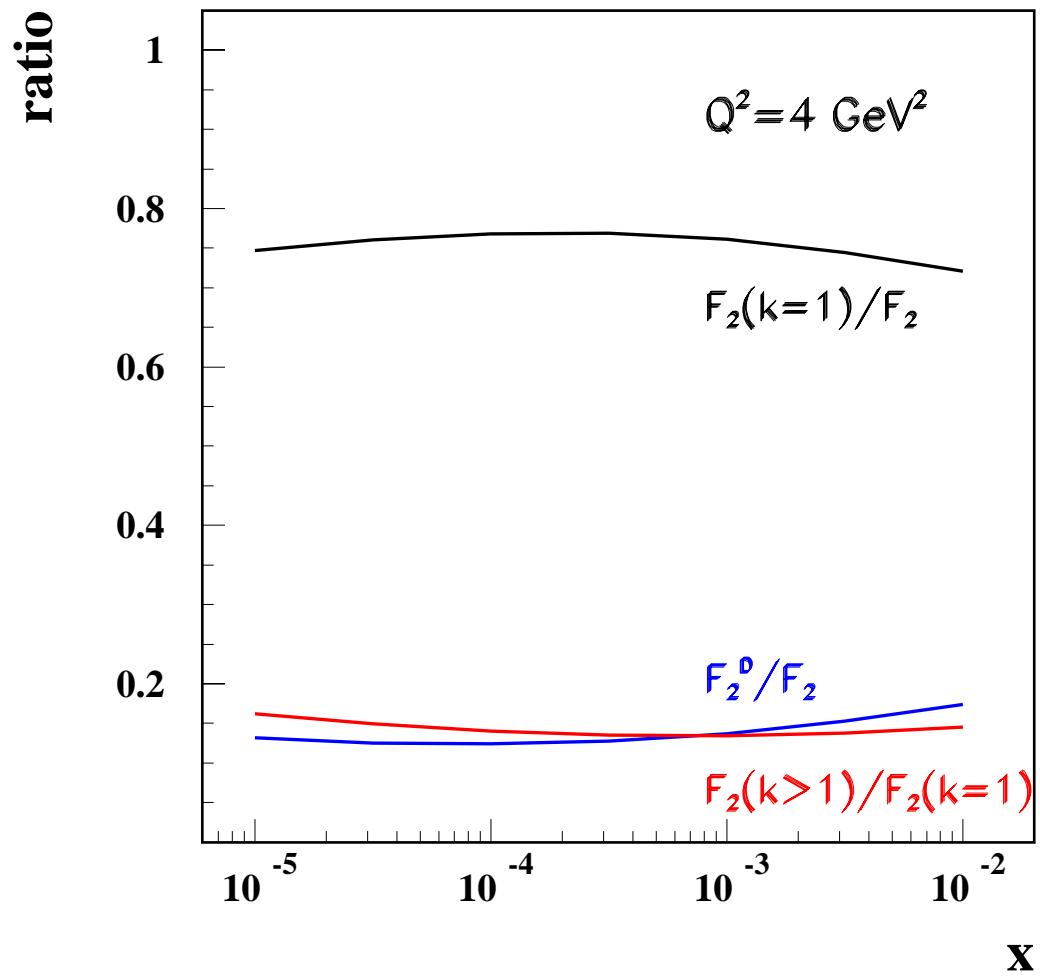


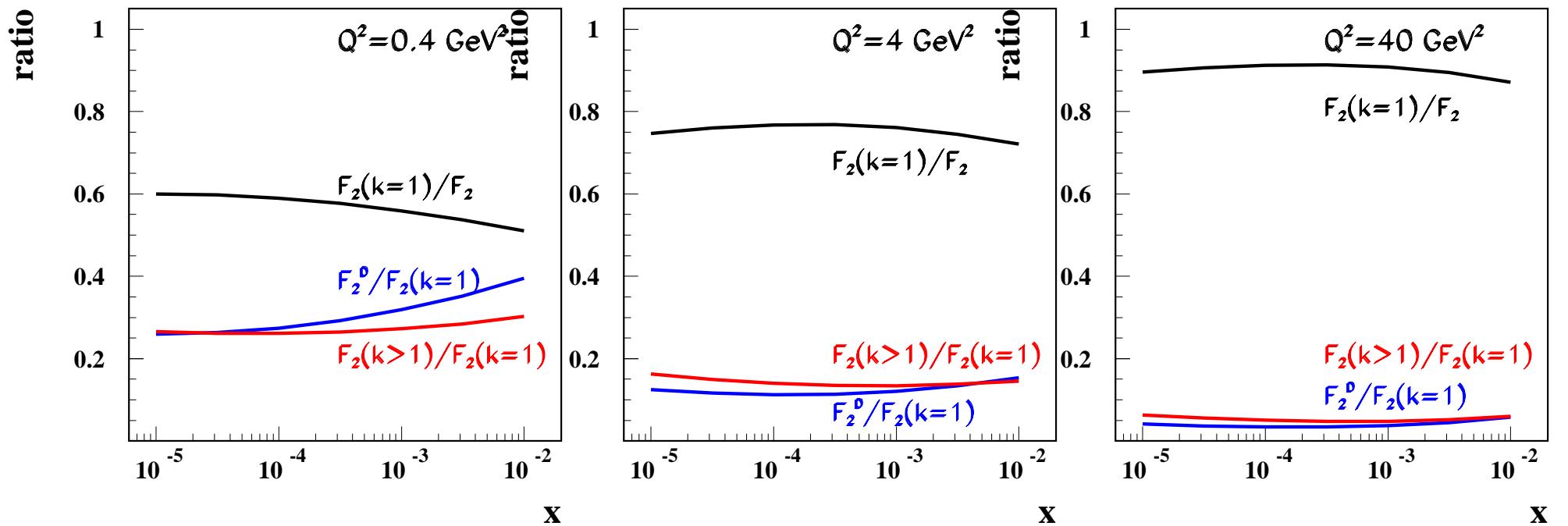


$$\sigma^{\gamma^* p} = \sum_f \int d^2\vec{r} \int d^2b \int_0^1 dz \Psi_f^*(Q^2, z, \vec{r}) 2 \left[1 - \exp(-\frac{\Omega}{2}) \right] \Psi_f(Q^2, z, \vec{r})$$

$$\sigma_k^{\gamma^* p} = \sum_f \int d^2\vec{r} \int d^2b \int_0^1 dz \Psi_f^*(Q^2, z, \vec{r}) \frac{\Omega^k}{k!} \exp(-\Omega) \Psi_f(Q^2, z, \vec{r})$$

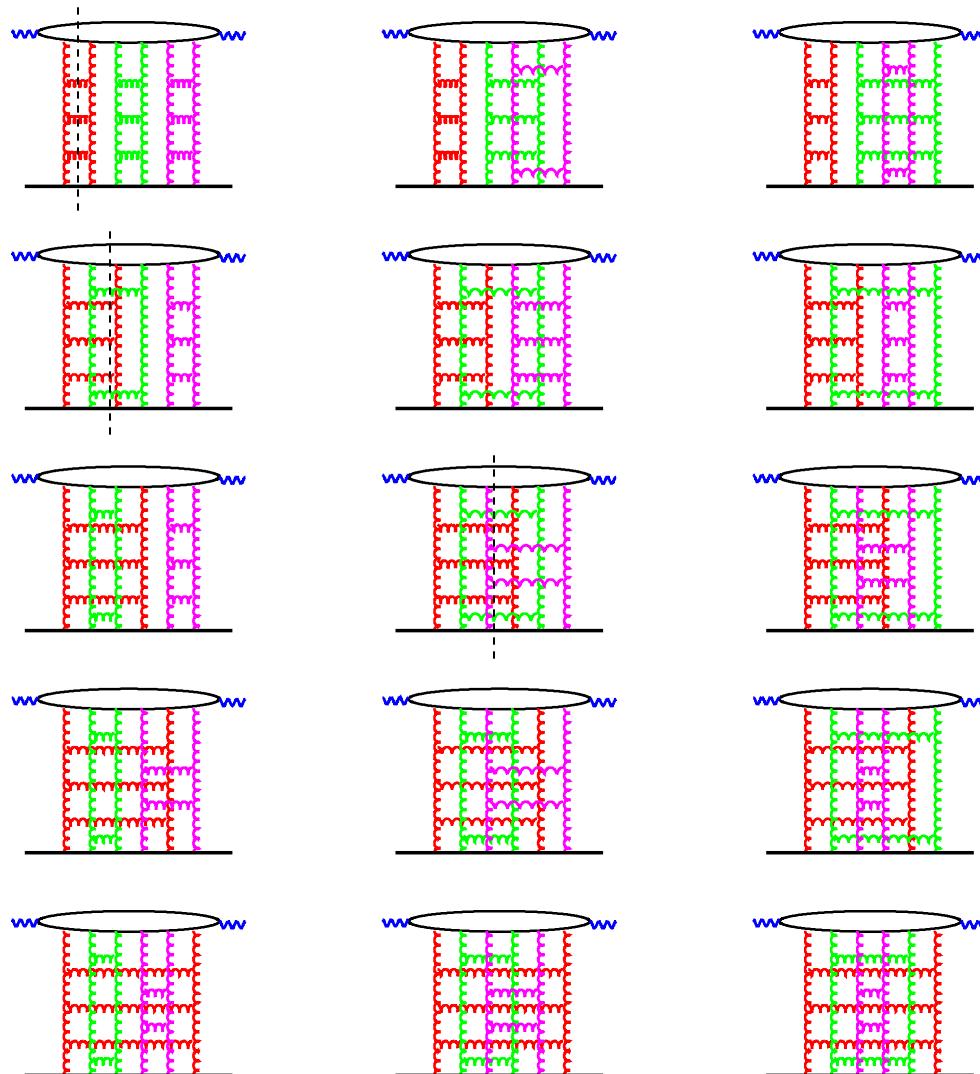




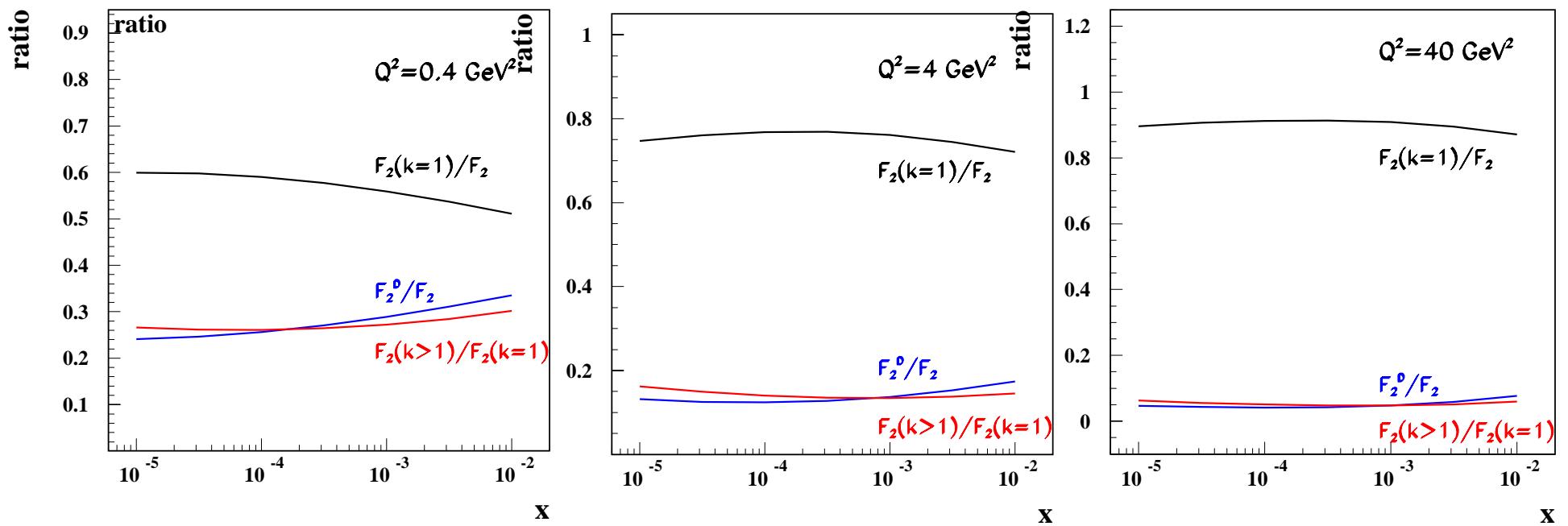


Note: AGK rules underestimate the amount of diffraction in DIS

**Outlook: Monte Carlo Feynman diagrams for 3, 4, 5 cut-pomerons
Sum up contribution of uncut-pomerons to infinity**



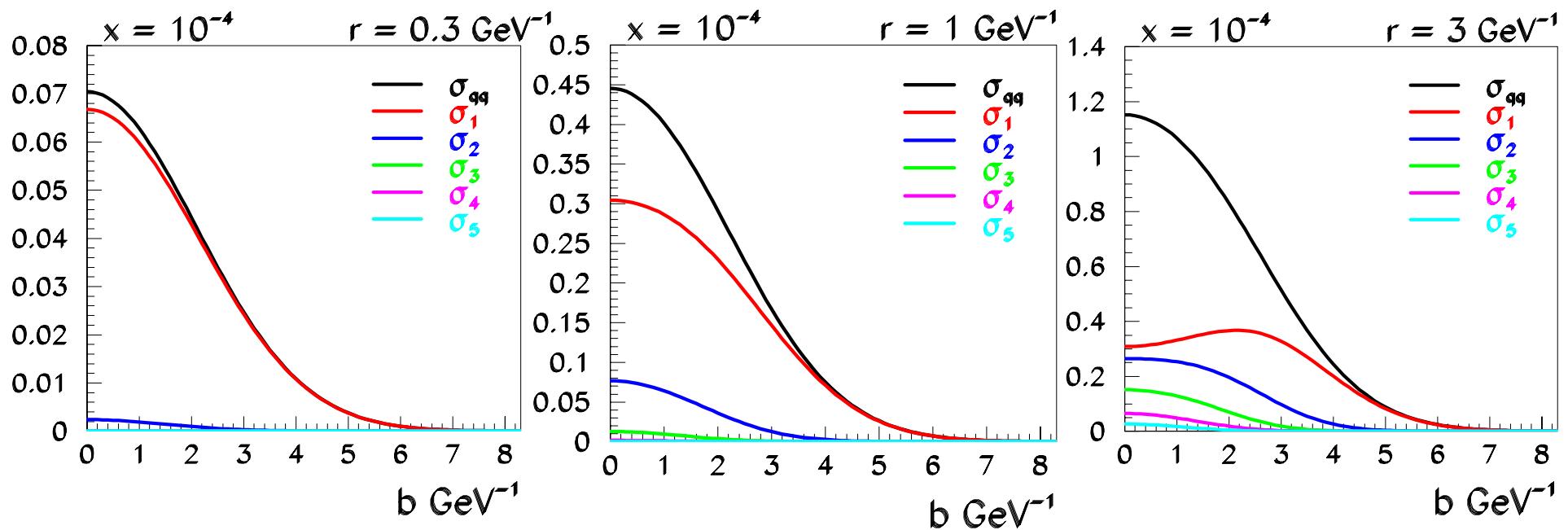
END



Note: AGK rules underestimate the amount of diffraction in DIS

$$\frac{d\sigma_{qq}}{d^2b} = 2 \cdot \left\{ 1 - \exp\left(-\frac{\Omega}{2}\right) \right\}$$

$$\frac{d\sigma_k}{d^2b} = \frac{\Omega^k}{k!} \exp(-\Omega)$$



$$\Omega = \frac{\pi^2}{N_C} r^2 \alpha_s(\mu^2) x g(x, \mu^2) T(b)$$