

# Balance Function and $p_t$ Fluctuations at RHIC

- Balance Function
  - $B(\Delta\eta)$ ,  $B(\Delta y)$ ,  $B(q_{inv})$ ,  $B(\Delta\phi)$
- $p_t$  Fluctuations
  - Excitation Function of  $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$ 
    - Au+Au at 20, 130, and 200 GeV
  - $F_{pt}$

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# The Balance Function

$$B(\Delta y) = \frac{1}{2} \left\{ \frac{N_{+-}(\Delta y) - N_{++}(\Delta y)}{N_+} + \frac{N_{-+}(\Delta y) - N_{--}(\Delta y)}{N_-} \right\}$$

Bass, Danielewicz, and Pratt, PRL **85**, 2689 (2000)

## Theoretical expectations for $B(\Delta y)$

PYTHIA representing p+p collisions shows a characteristic width of about 1 unit of  $\Delta y$

Bjorken thermal model representing delayed hadronization shows narrower balance function width

**Nucleon-nucleon** → wide

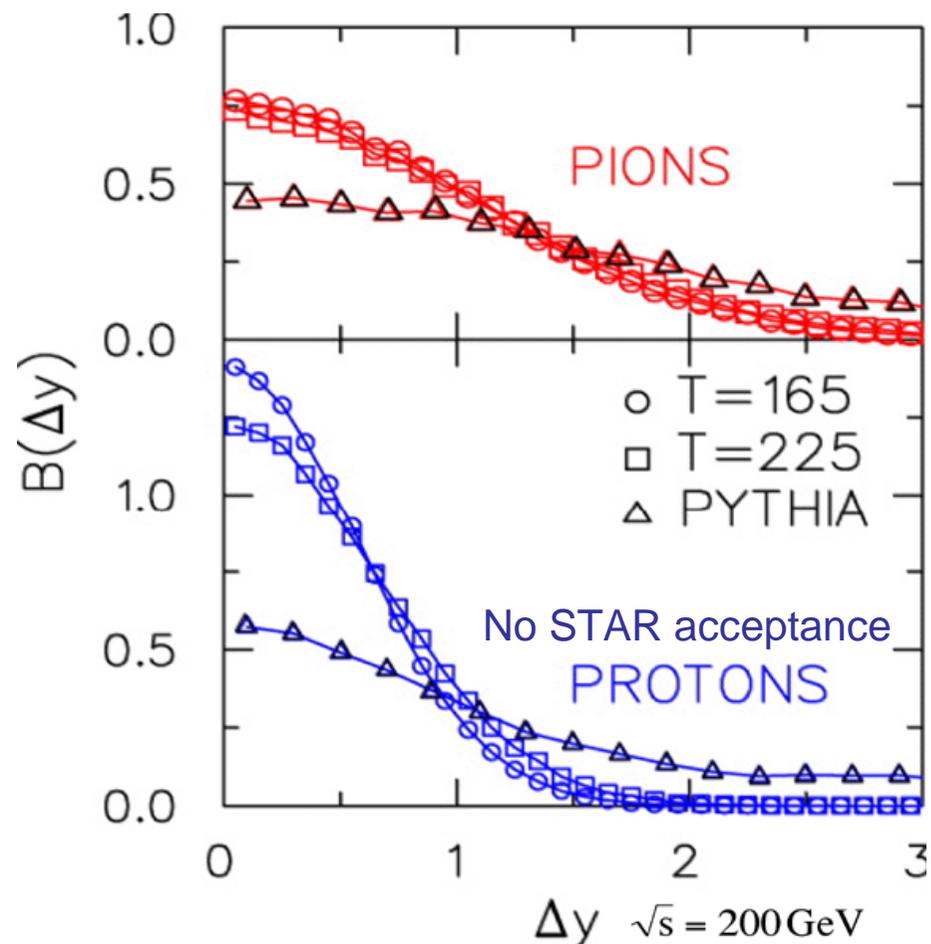
**Delayed hadronization** → narrow

## Experimental considerations

Use  $\Delta\eta$ ,  $\Delta y$ ,  $q_{inv}$ ,  $\Delta\phi$

Centrality dependence for Au+Au and d+Au at 200 GeV

All centralities for p+p at 200 GeV



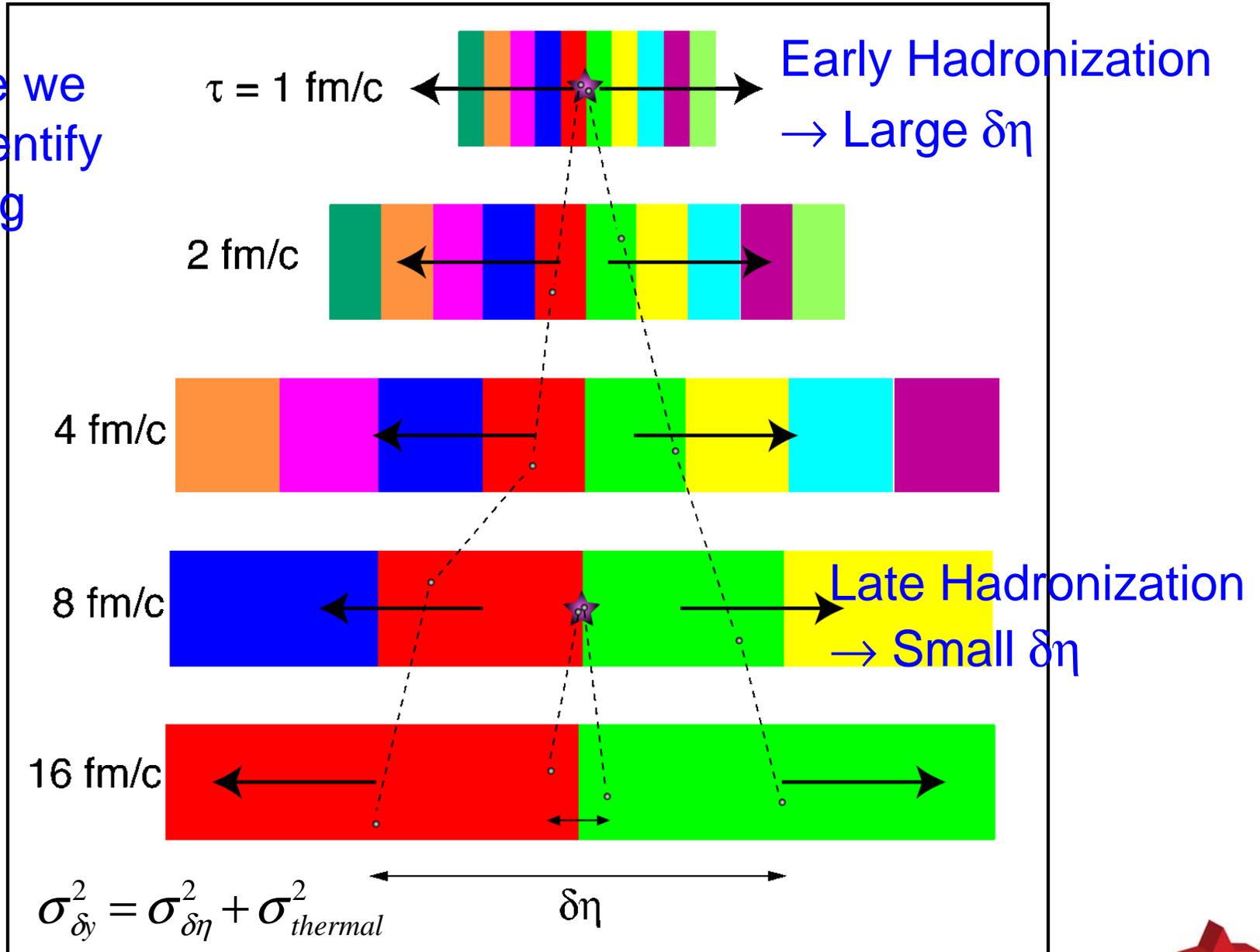
# Delayed Hadronization

- A new observable has been proposed by Bass, Danielewicz, and Pratt [Phys. Rev. Lett. **85**, 2689 (2000)] called the **balance function**
- The basic premise is that charge/anti-charge pairs are created close together in space-time
- If these pairs are created early in the collision, they will be pulled apart in rapidity by longitudinal expansion and will suffer scattering for the duration of the collision, losing their correlation in rapidity
- If instead, the system exists in a deconfined phase for a substantial time, and then the pairs are formed at hadronization, they will experience less expansion and fewer collisions, retaining more of their correlation in rapidity



# Delayed Hadronization

Suppose we could identify balancing charges



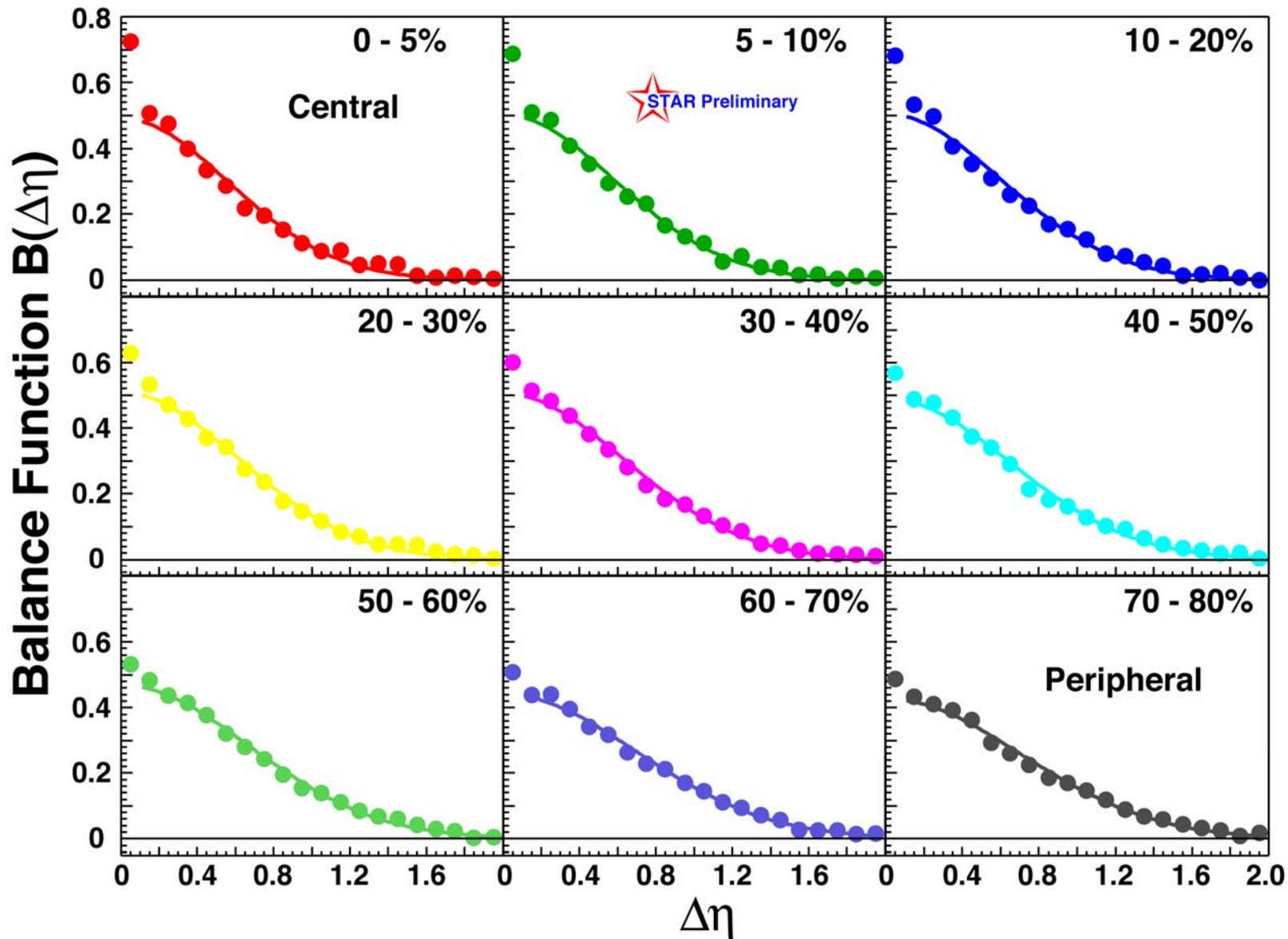
After Scott Pratt



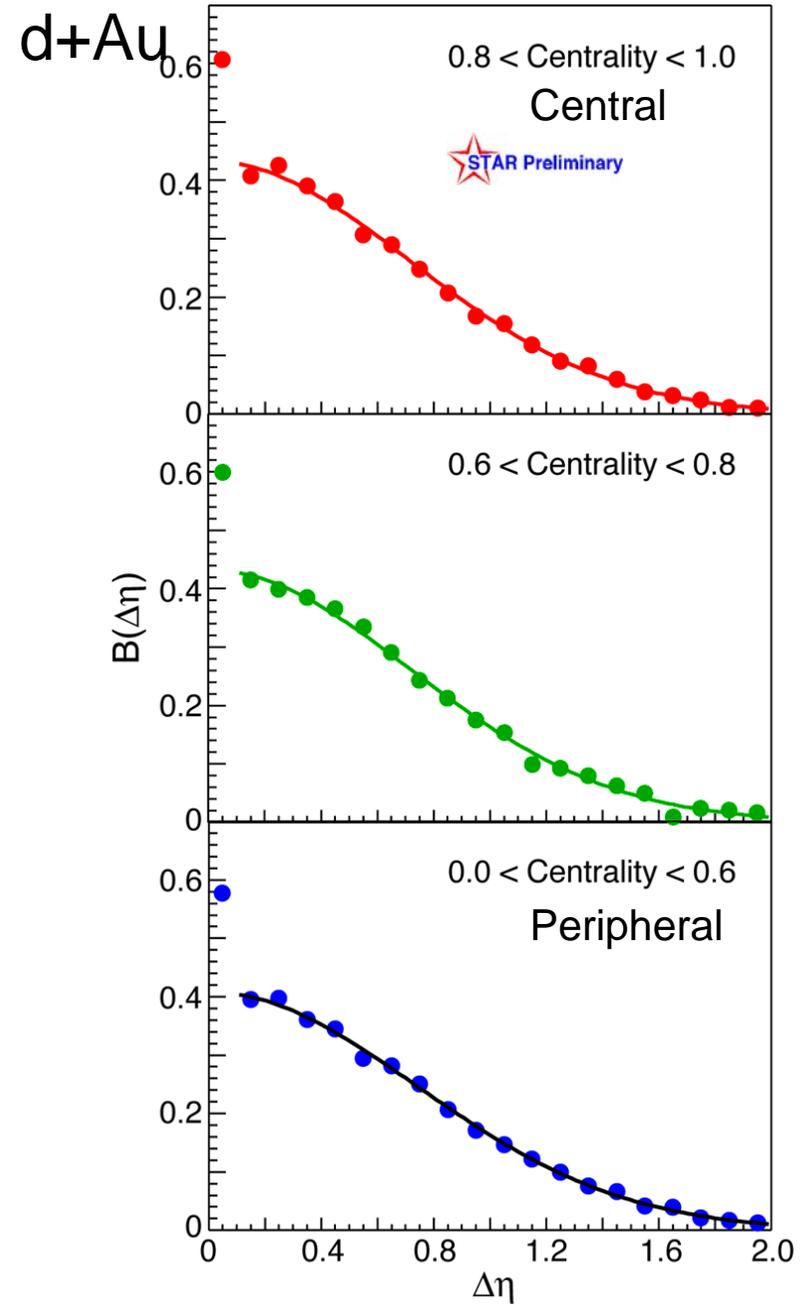
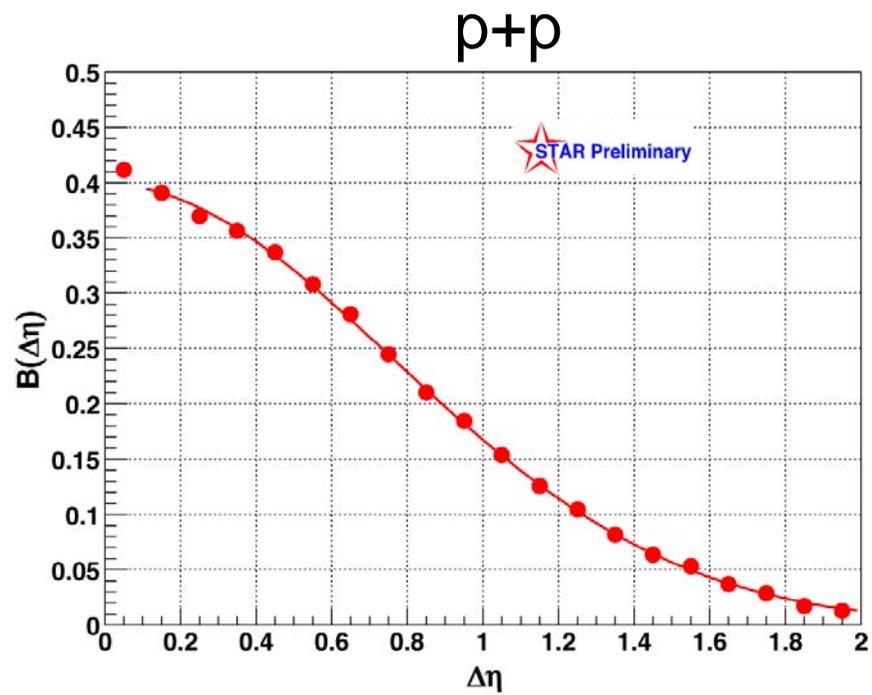
Gary Westfall for STAR



# Balance Function for Au+Au at 200 GeV

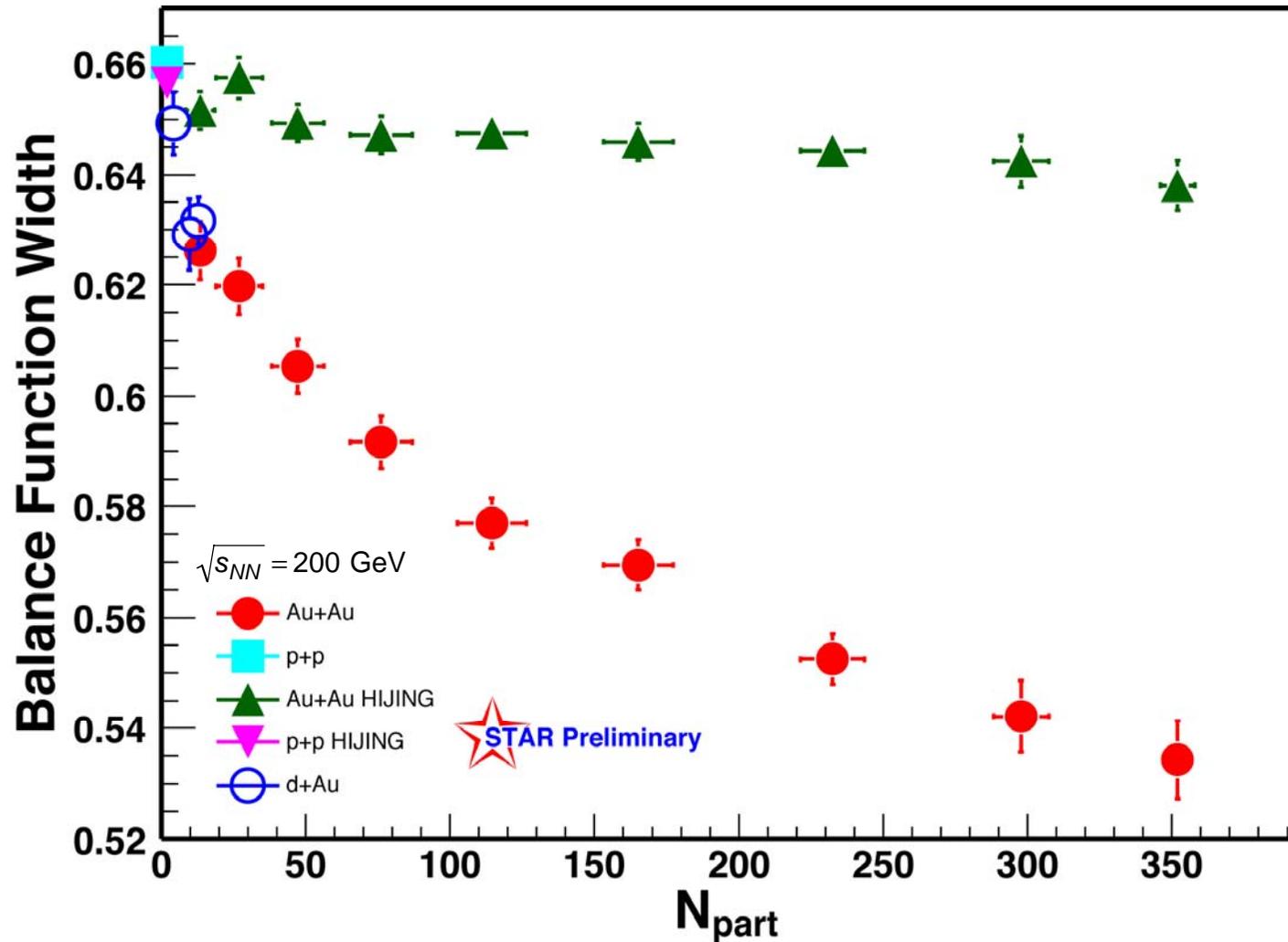


# Balance Function for p+p and d+Au at 200 GeV



# Balance Function Widths

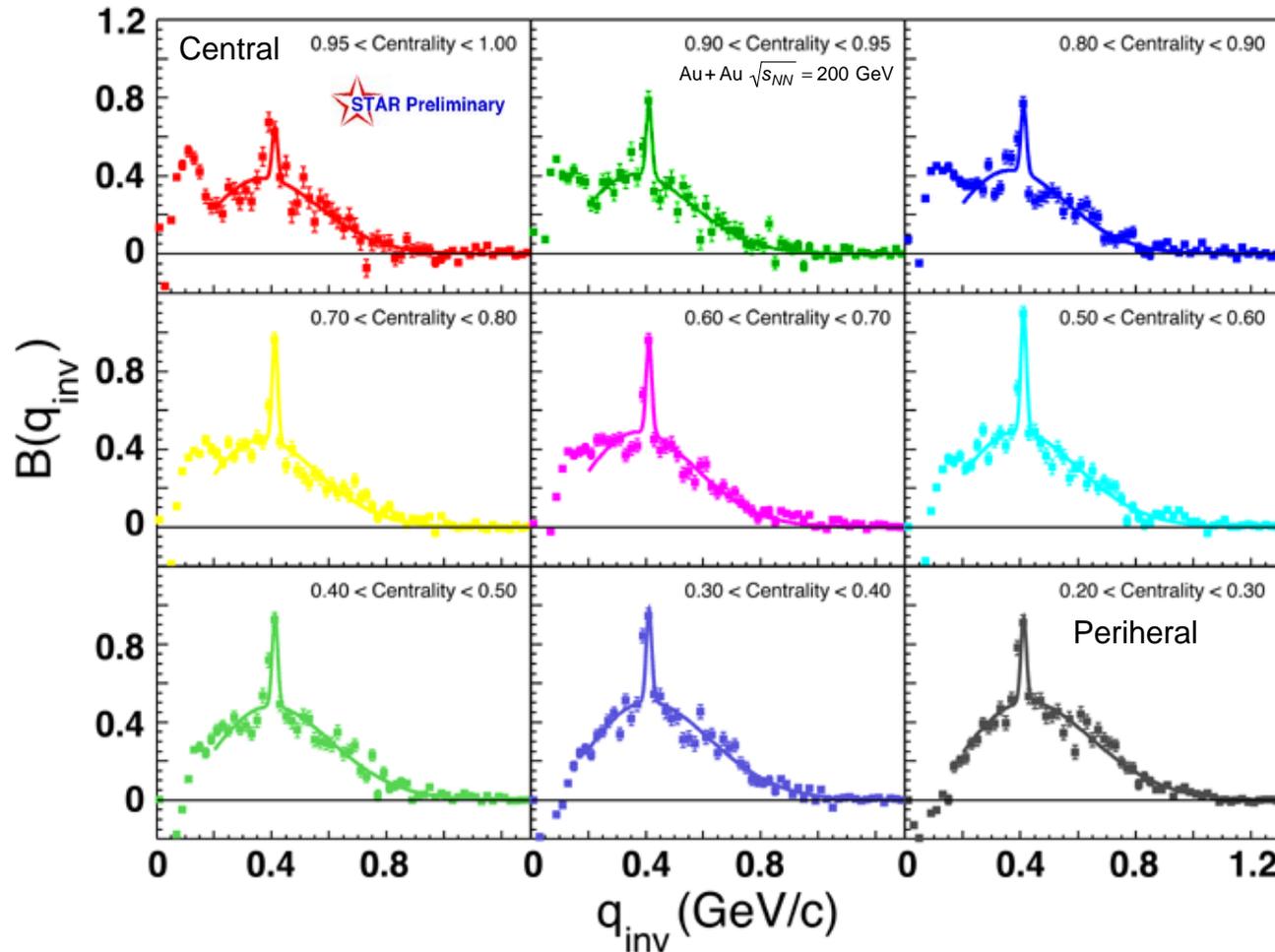
Width is defined as weighted average of  $\Delta\eta$  for  $0.1 \leq \Delta\eta \leq 2.0$



Balance function for Au+Au narrows in central collisions  
HIJING shows little centrality dependence  
Smooth dependence on  $N_{part}$



# Balance Function for Pions using $q_{inv}$



Fits are  
thermal +  
 $K^0$  decay

Thermal  
distribution is

$$\propto q_{inv}^2 e^{-\frac{q_{inv}^2}{2\sigma}}$$

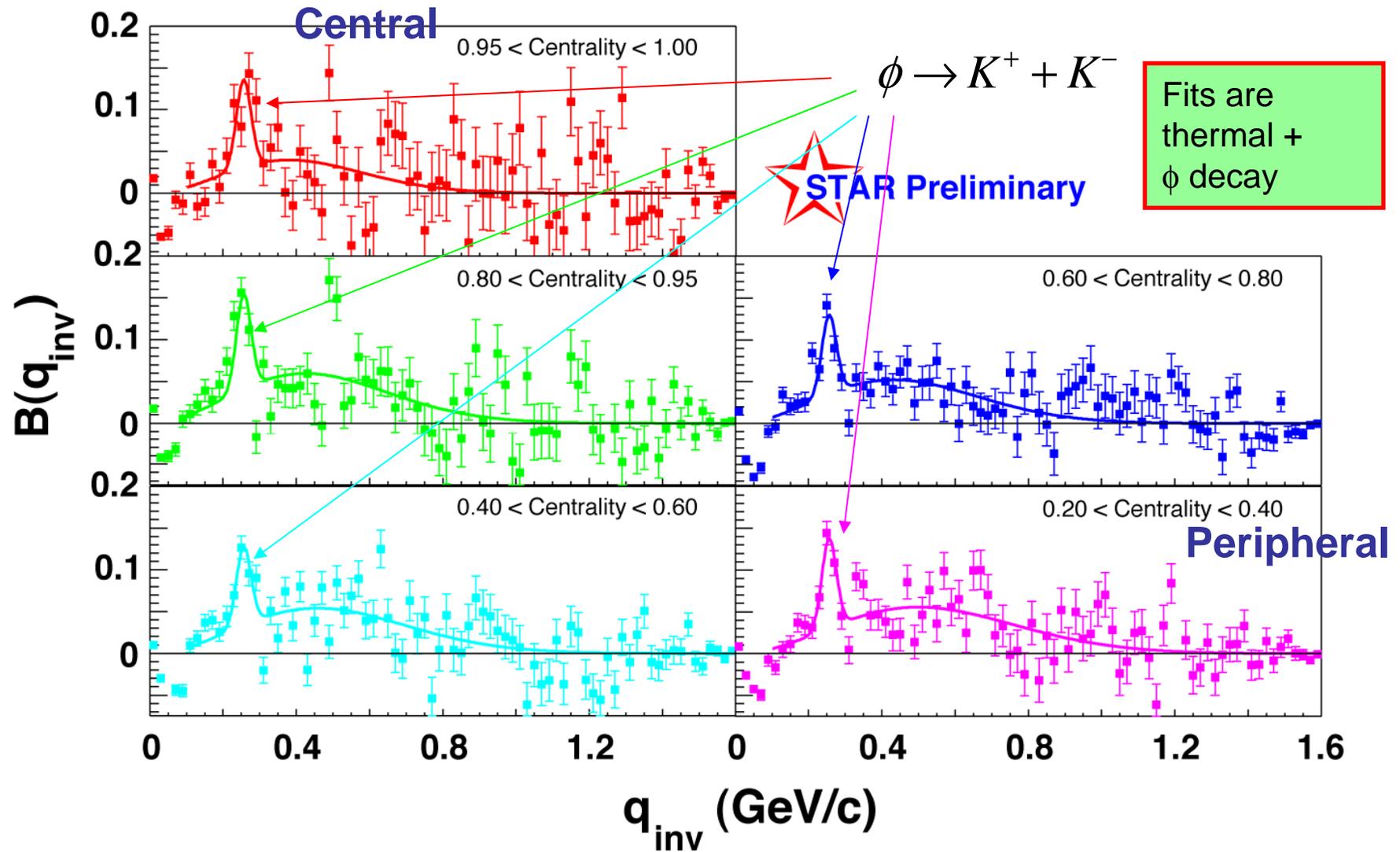
Narrowing of  $B(\Delta\eta)$  may be caused by transverse flow  
Use  $B(q_{inv})$  to remove reference frame dependence  
Allow more direct comparison with thermal models



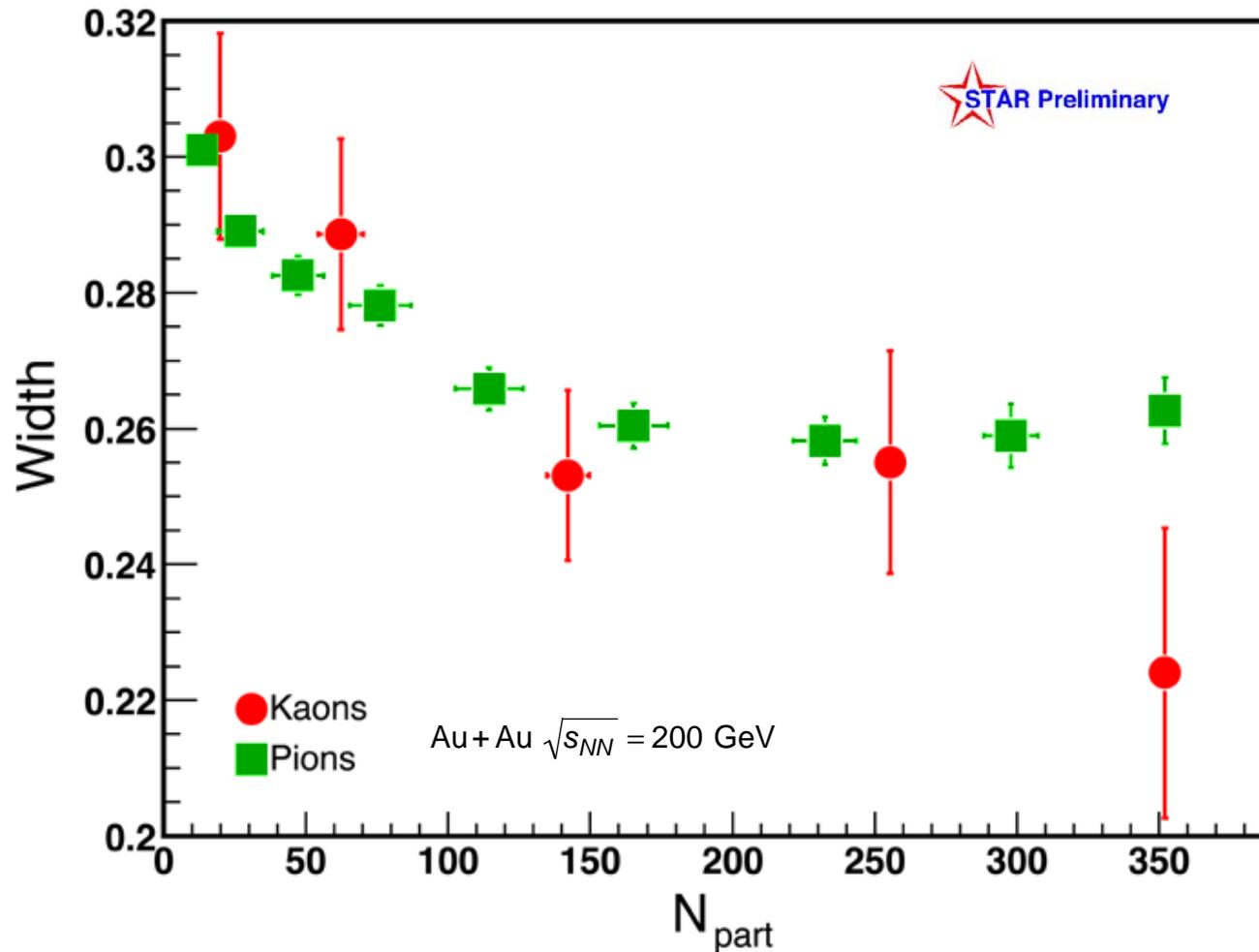
# Balance Function for Kaons using $q_{inv}$

Identified charged kaon pairs ( $K^+, K^-$ )

$p < 0.8$  GeV, no electrons



# Width of Balance Function using $q_{inv}$

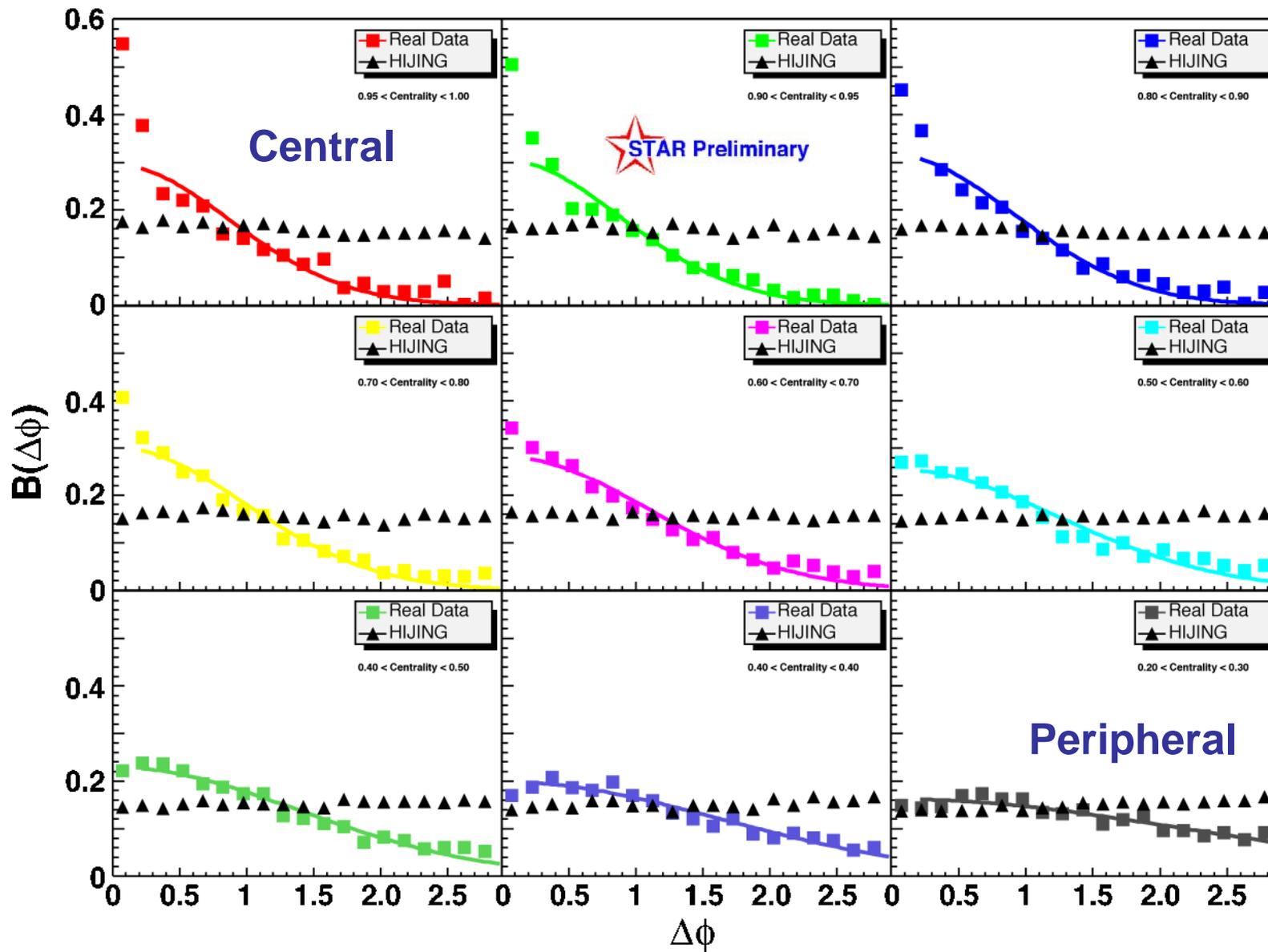


Balance function  $B(q_{inv})$  for pions and kaons narrows in central collisions even when using Lorentz invariant observable



# Balance Function for All Charged Particles, $B(\Delta\phi)$

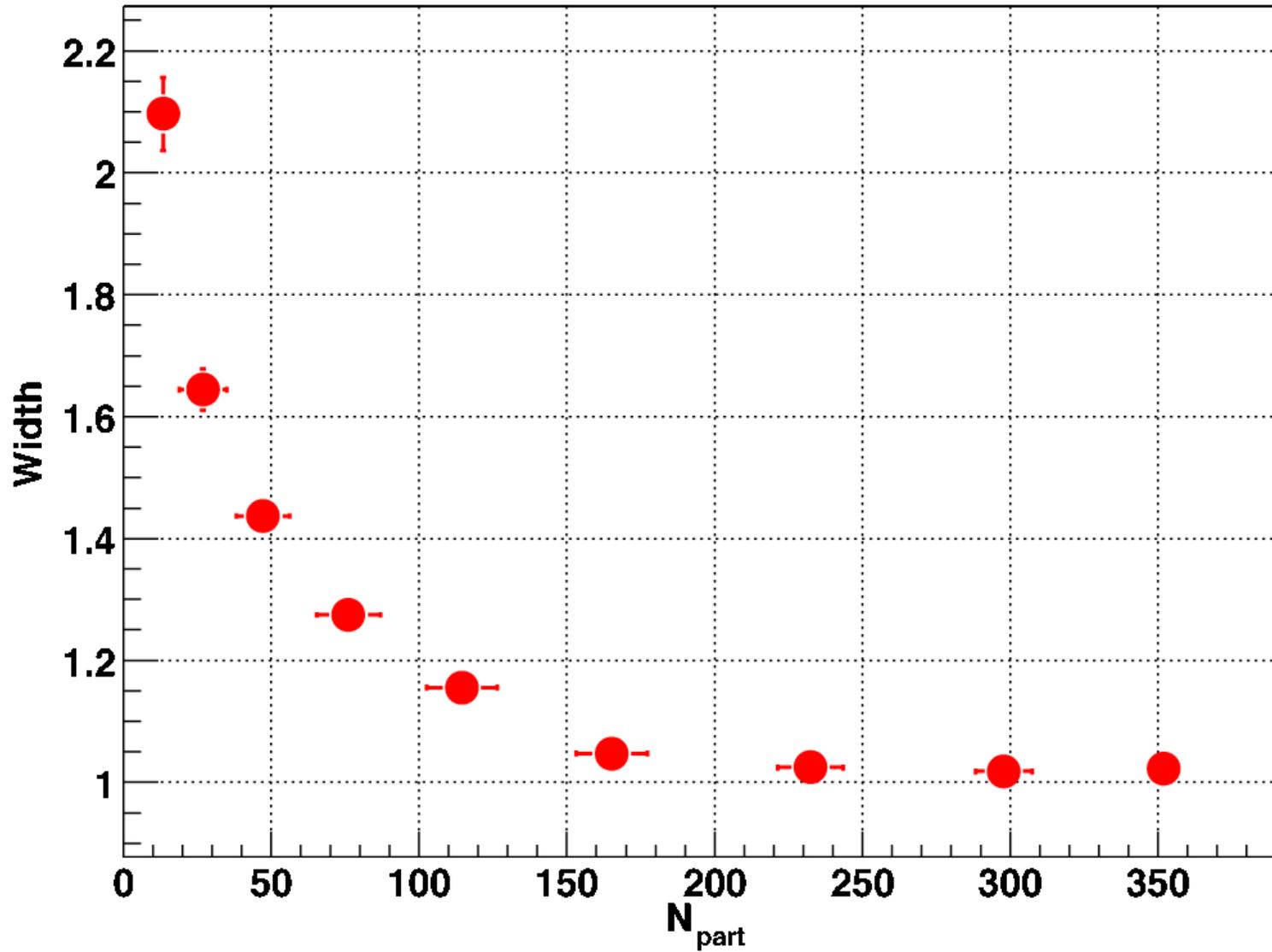
no electrons



# Azimuthal Width

$B(\Delta\phi)$

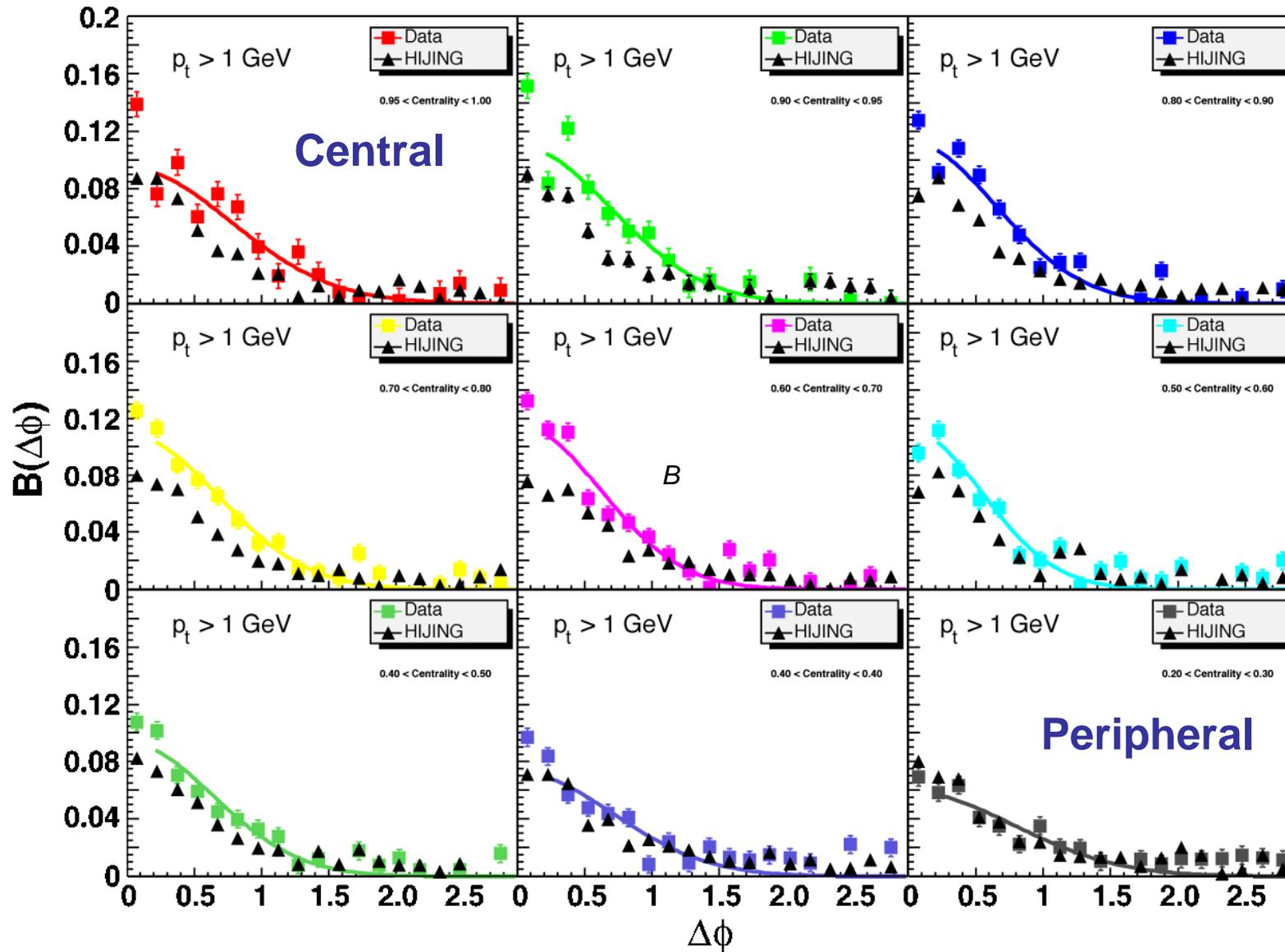
STAR Preliminary



Gary Westfall for STAR

# B( $\Delta\phi$ ) at High $p_t$

STAR Preliminary

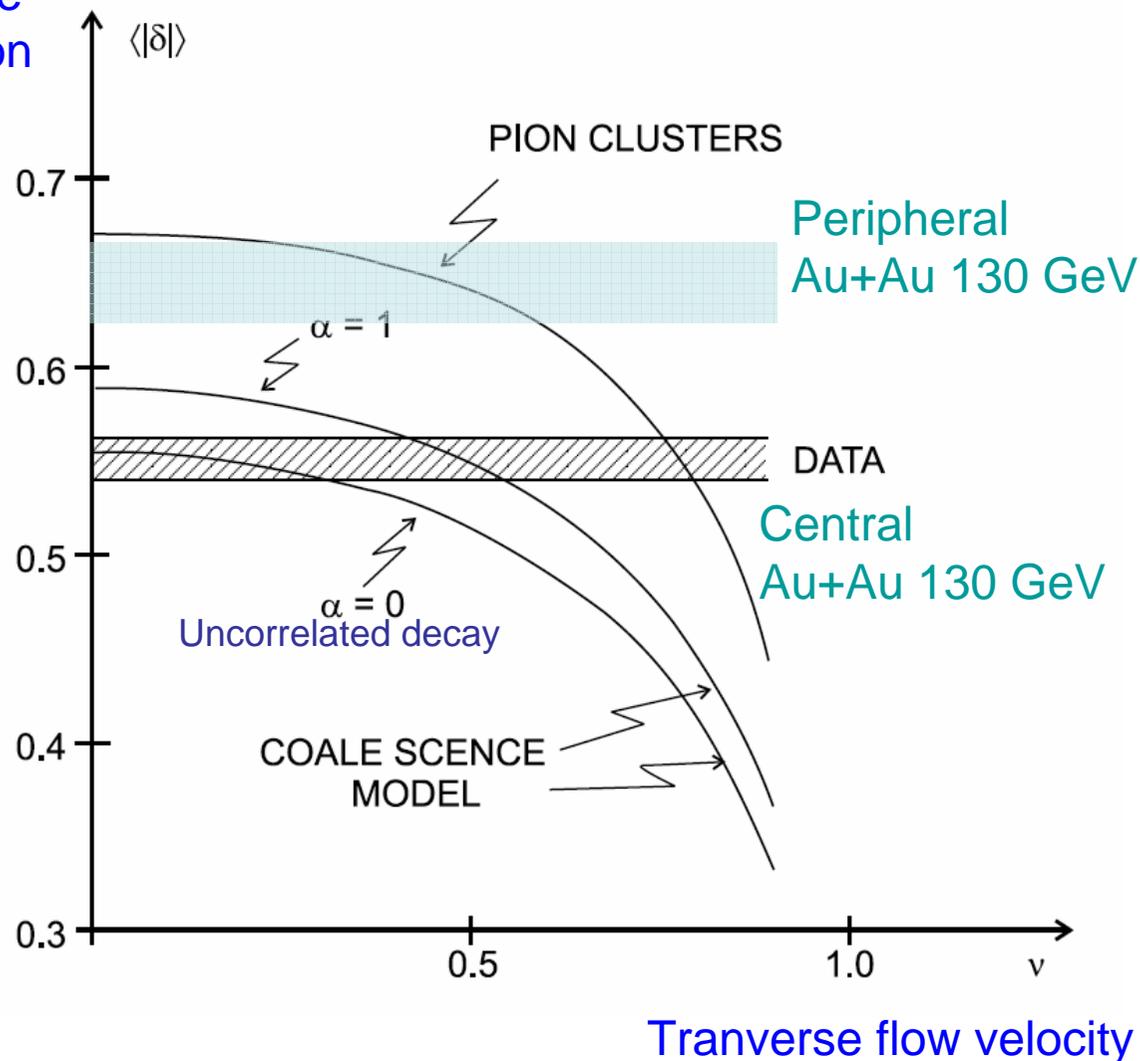


# Quark Coalescence Narrows Balance Function

A. Bialas  
hep-ph/0308245

Shows that measured narrow width in central collisions can be explained using a quark/antiquark coalescence model while hadronic scenario fails

Balance  
Function  
Width



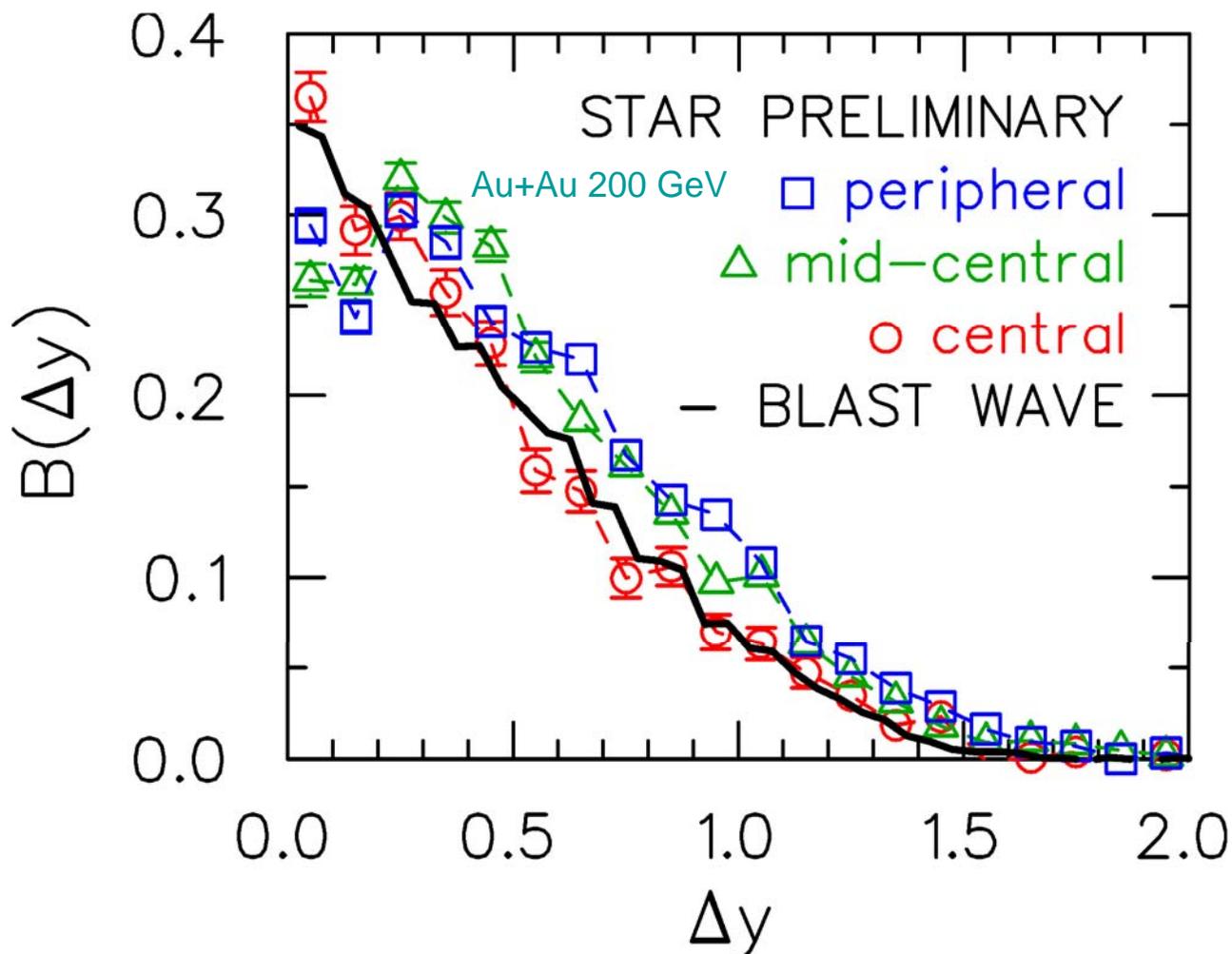
# Theoretical Predictions for the Balance Function

Cheng, Petriconi, Pratt,  
and Skoby, nucl-  
th/0401008

Includes HBT,  
Coulomb, resonances,  
strong interactions,  
radial flow, conservation  
of S,Q,B

Uses STAR acceptance  
filter

The agreement with the  
measured narrow  
balance function in  
central collisions  
suggests that charge  
conservation remains  
highly localized at  
breakup

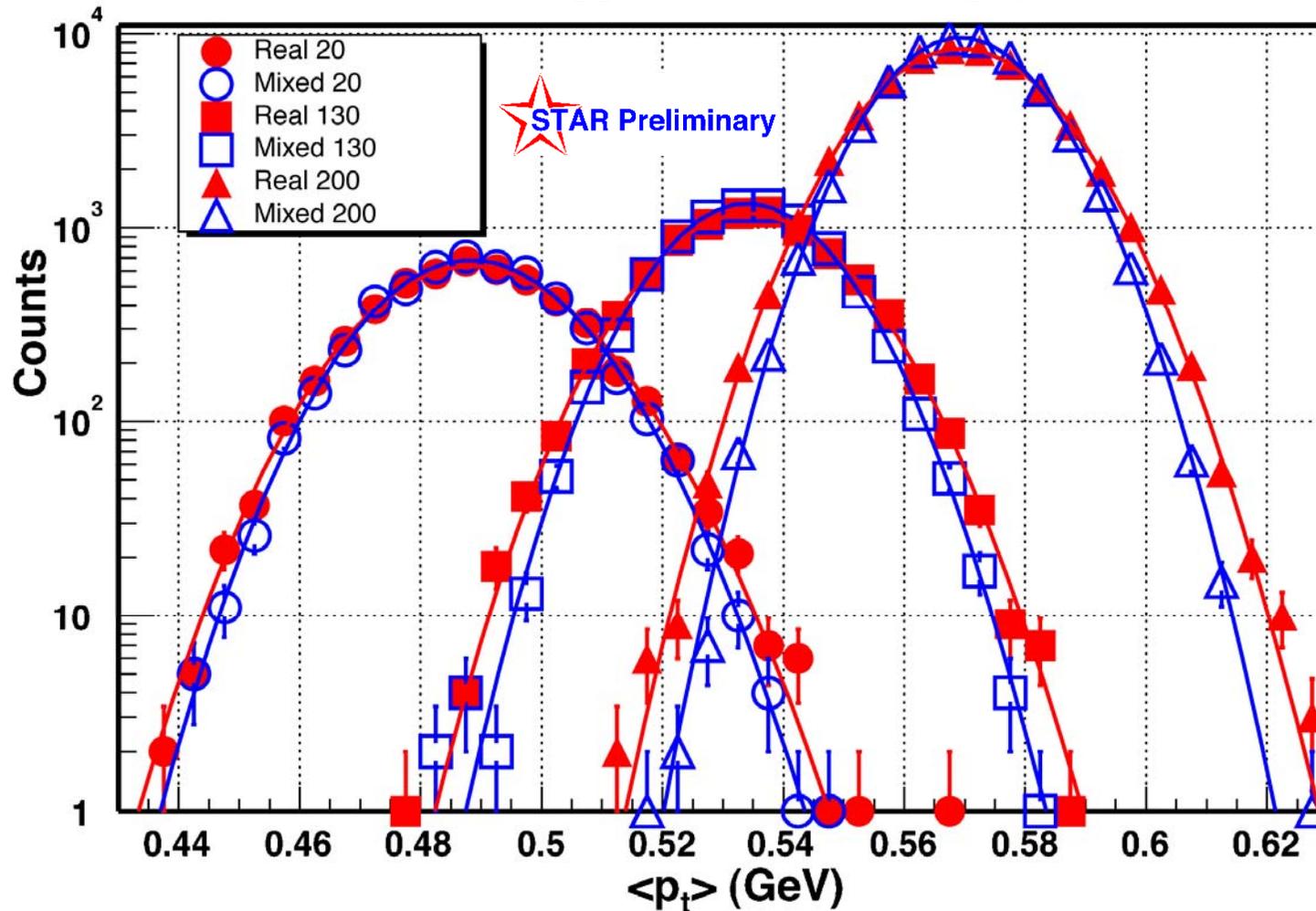


# $p_t$ Fluctuations

- Search for dynamical fluctuations motivated by predictions
  - Fluctuations in energy density due to localized deconfinement
  - Increased fluctuations in energy density due to long range correlations
  - Proximity to tri-critical and critical points would lead to changes in fluctuation patterns
  - Production of DCCs
  - Fluctuations from jet production



# Histograms of $\langle p_t \rangle$



Au+Au at 20, 130, and 200 GeV  
5% most central bin using min bias data,  $|\eta| < 1.0$   
Real is wider than mixed  
→ Dynamical fluctuations



# Definition of $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$

As a function of centrality and acceptance

$$\langle\langle p_t \rangle\rangle = \left( \sum_{k=1}^{N_{event}} \langle p_t \rangle_k \right) / N_{event} \quad \text{where} \quad \langle p_t \rangle_k = \left( \sum_{i=1}^{N_k} p_{t,i} \right) / N_k$$

$$\langle \Delta p_{t,i} \Delta p_{t,j} \rangle = \frac{1}{N_{event}} \sum_{k=1}^{N_{event}} \frac{C_k}{N_k (N_k - 1)}$$

$$C_k = \sum_{i=1}^{N_k} \sum_{j=1, i \neq j}^{N_k} (p_{t,i} - \langle\langle p_t \rangle\rangle) (p_{t,j} - \langle\langle p_t \rangle\rangle)$$

$N_{event}$  = number of events

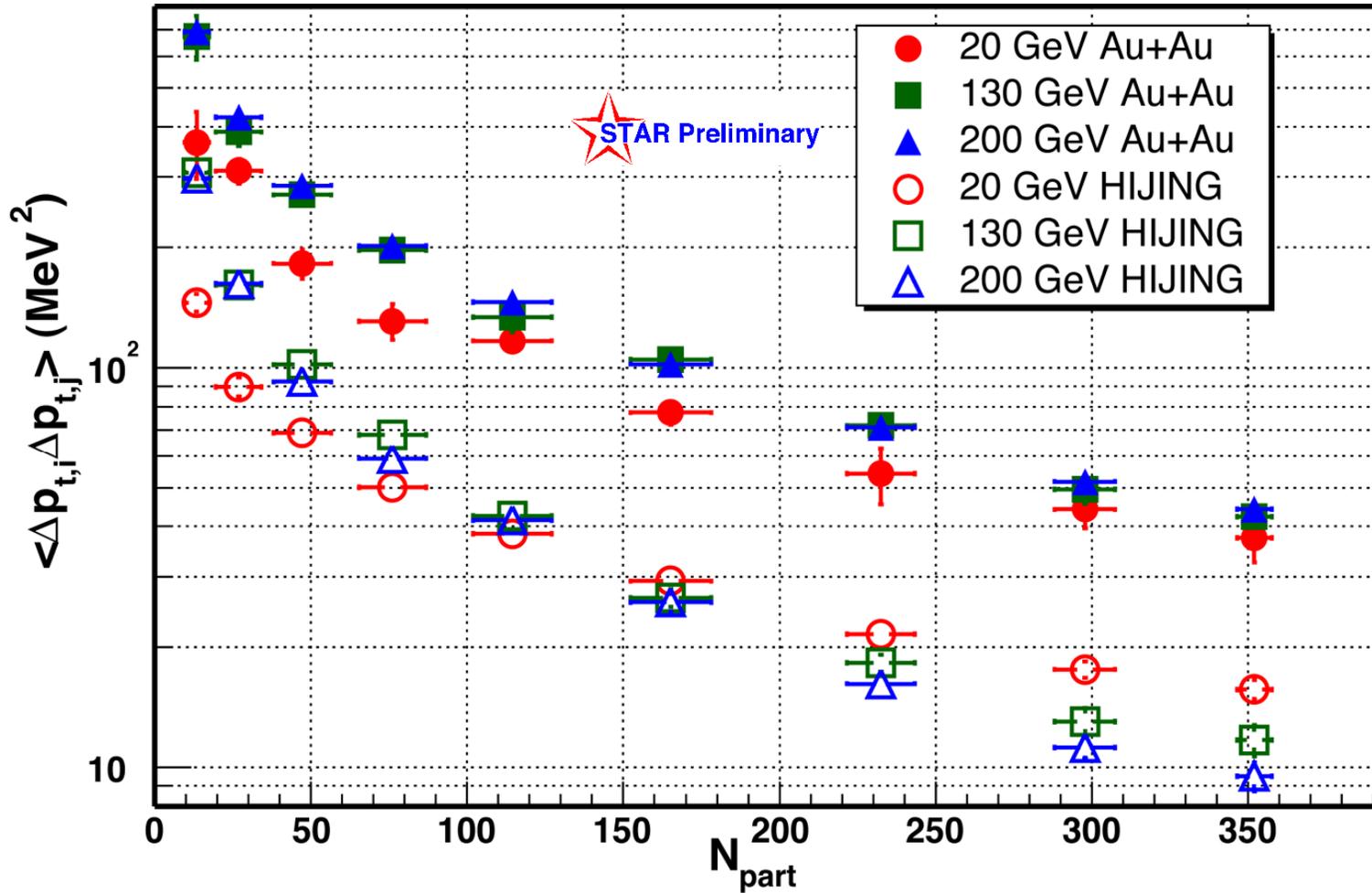
$\langle p_t \rangle_i$  = average  $p_t$  for  $i^{th}$  event

$N_k$  = number of tracks for  $k^{th}$  event

$p_{t,i}$  =  $p_t$  for  $i^{th}$  track in event



# $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$ as a Function of Incident Energy



$$\langle \Delta p_{t,i} \Delta p_{t,j} \rangle = \frac{1}{N_{event}} \sum_{k=1}^{N_{event}} \frac{C_k}{N_k(N_k - 1)}$$

$$C_k = \sum_{i=1}^{N_k} \sum_{j=1, i \neq j}^{N_k} (p_{t,i} - \langle p_t \rangle)(p_{t,j} - \langle p_t \rangle)$$

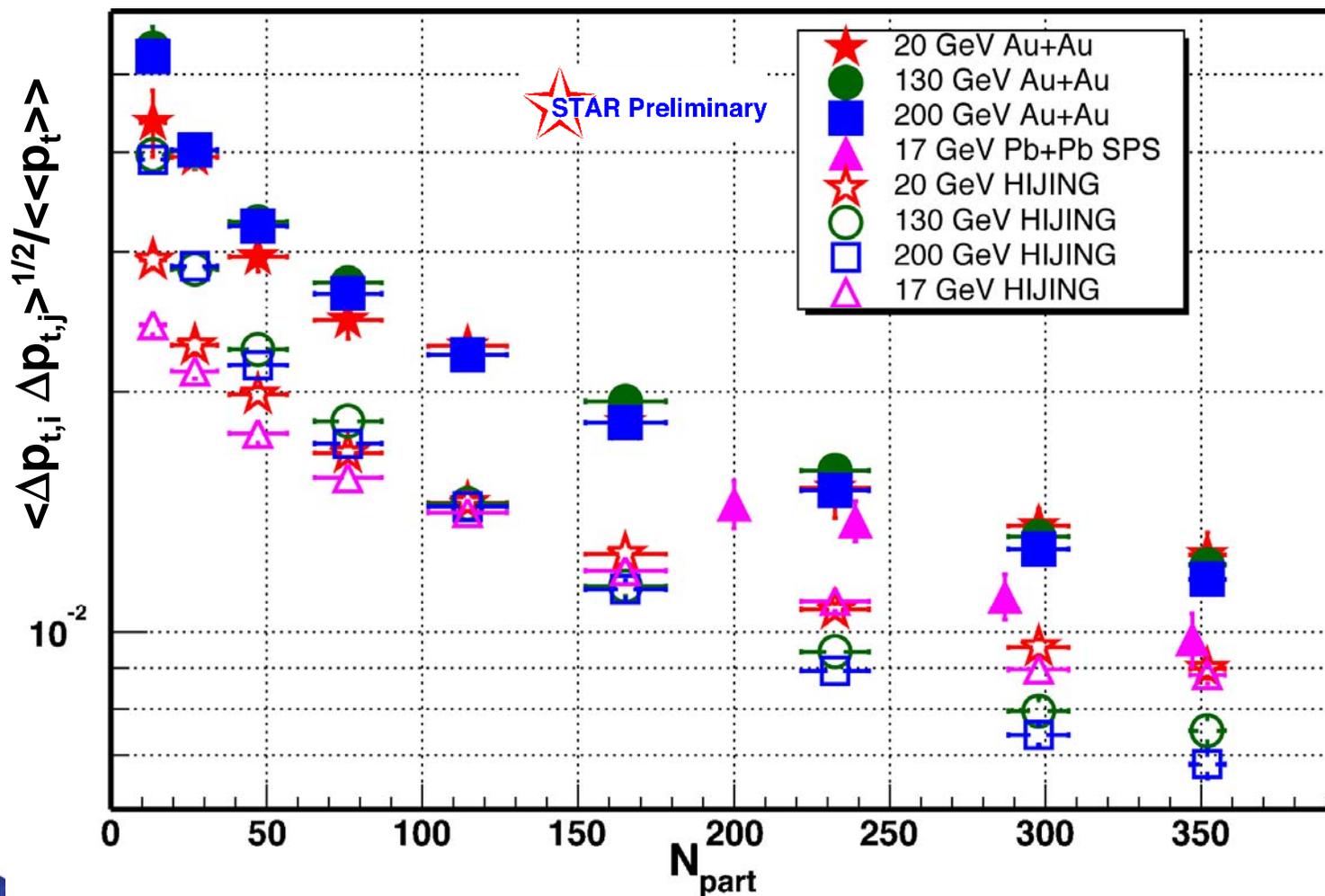
Compare Au+Au at 20, 130, and 200 GeV  
 $|\eta| < 1$   
 Compare with filtered HIJING



# $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle^{1/2} / \langle \langle p_t \rangle \rangle$ as a Function of Incident Energy

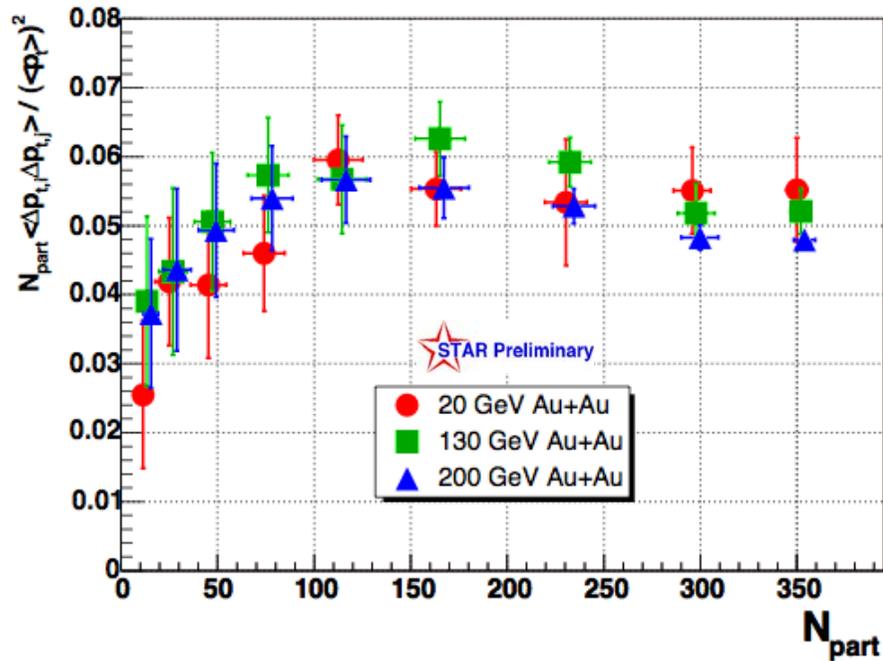
Compare Au+Au at 20, 130, 200 GeV,  $|\eta| < 1$

Compare with CERES result from SPS, 17 GeV Pb+Pb

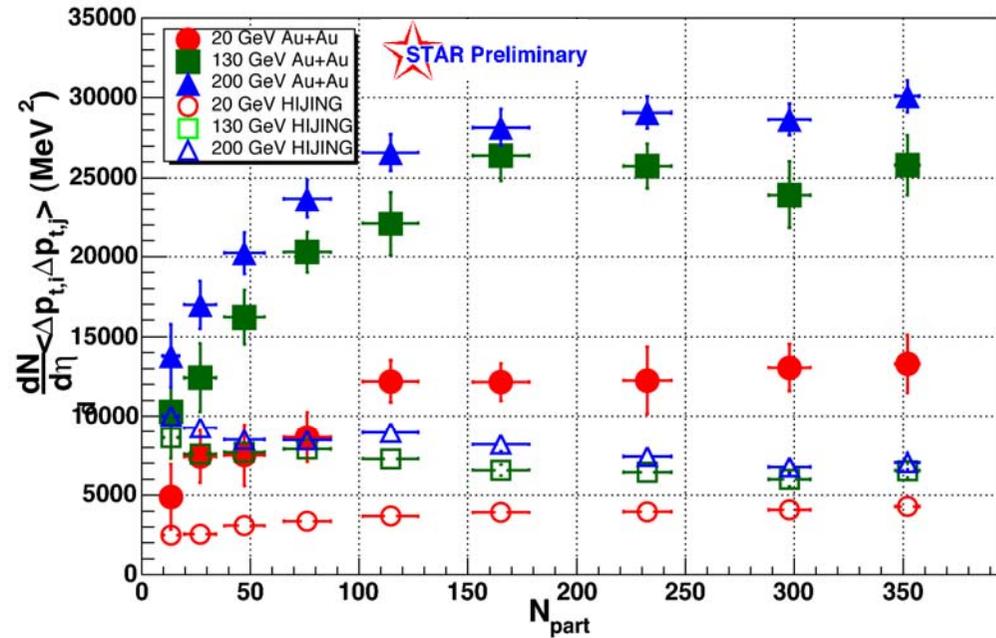


# Different Scaling Methods for $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$

Scale with  $N_{part}$



Scale with  $dN/d\eta$  - compare to  $\Delta\sigma_{pt}$

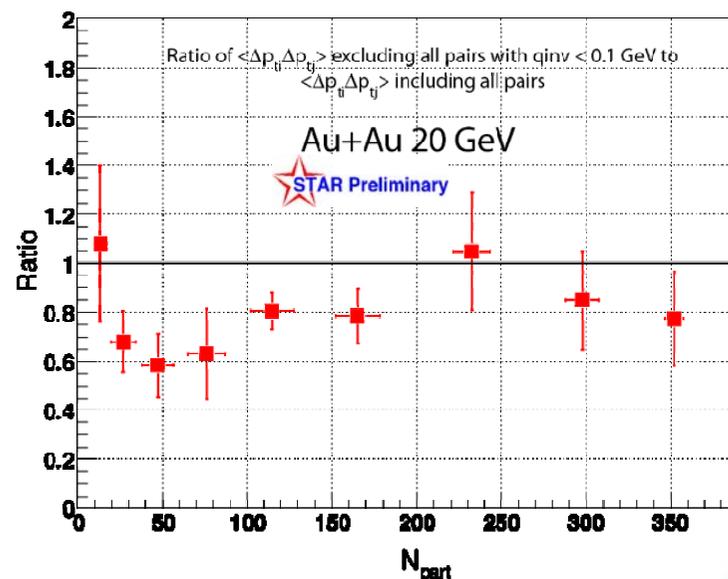
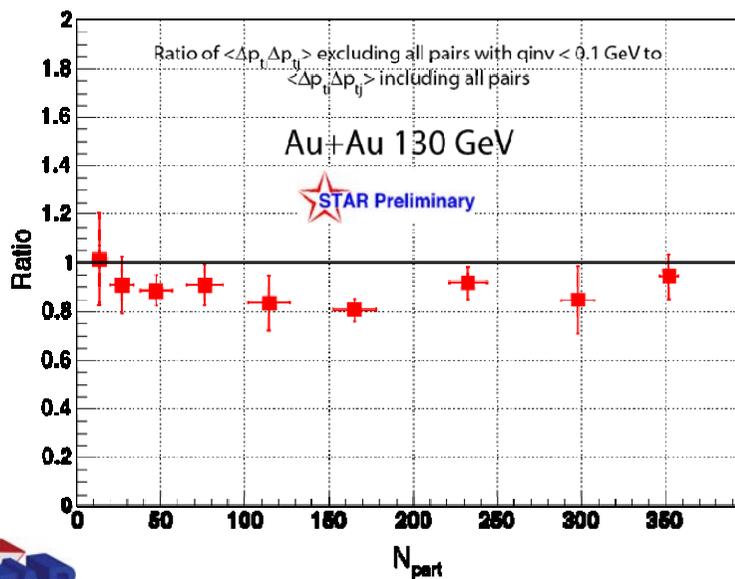
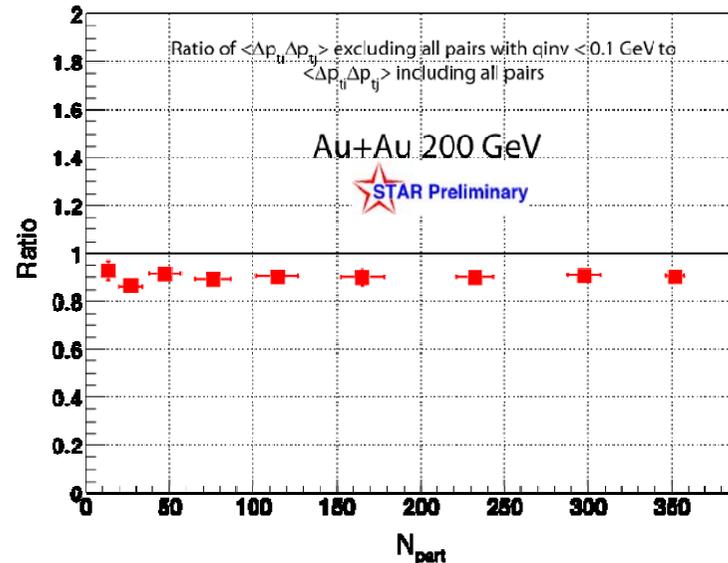
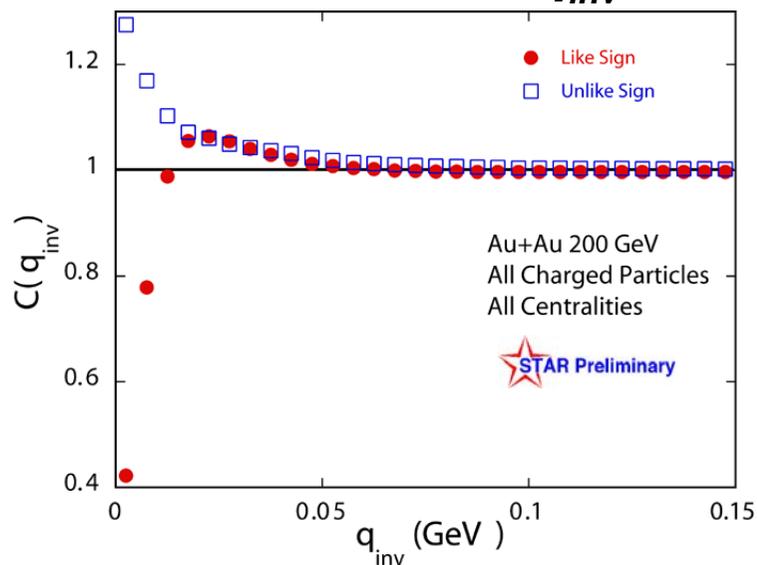


Centrality dependence may be sign of the onset of equilibration in central Au+Au collisions



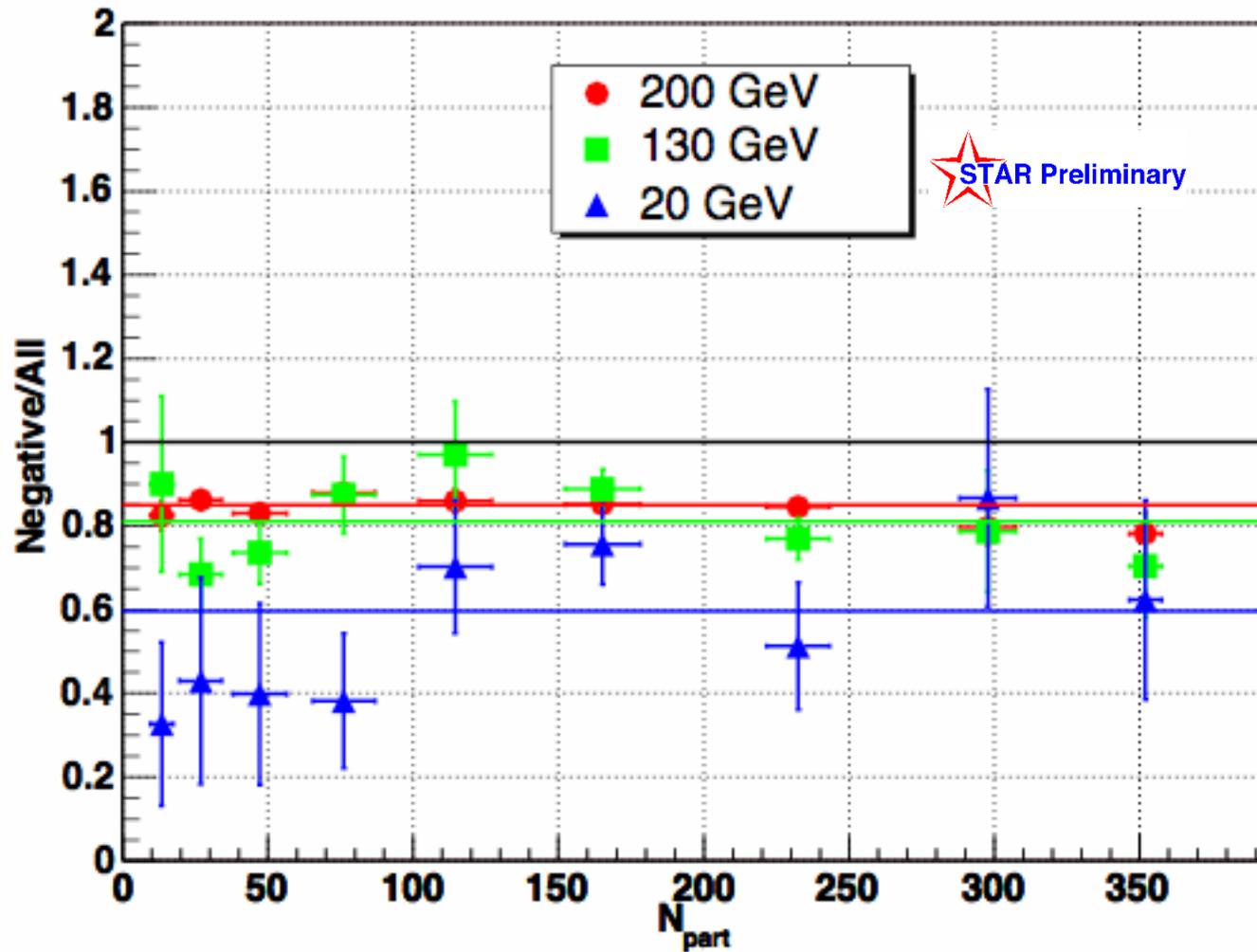
# Effects of Short Range Correlations

Ratio of  $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$  using Pairs with  $q_{inv} < 0.1$  GeV to All Pairs



# Estimate Contribution of Resonances

Ratio of  $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$  using  
Negative Particles Only to All Particles



# Comparison to $F_{pt}$ at 200 GeV

$$\omega_T = \frac{\sigma_T}{\mu_T} \quad F_{pt} = \frac{\omega_{T,real} - \omega_{T,mixed}}{\omega_{T,mixed}} = \frac{\sigma_{T,real} - \sigma_{T,mixed}}{\sigma_{T,mixed}}$$

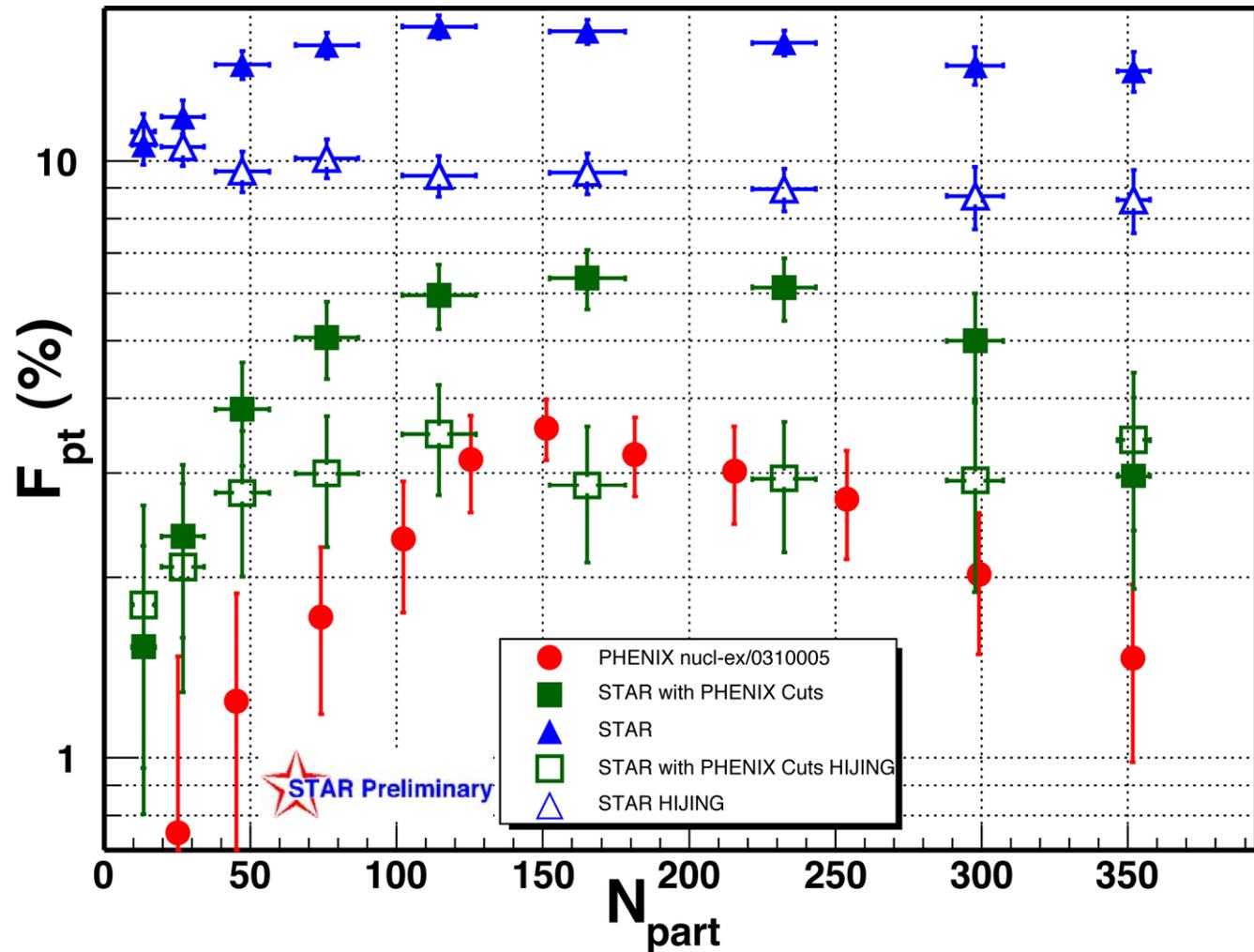
200 GeV Au+Au  
STAR Cuts

$|\eta| < 1.0$   
 $\Delta\phi = 360^\circ$   
 $0.1 < p_t < 2 \text{ GeV}$

200 GeV Au+Au  
STAR with  
PHENIX Cuts

$|\eta| < 0.35$   
 $\Delta\phi = 2 \times 90^\circ$   
 $0.2 < p_t < 2 \text{ GeV}$

Acceptance  
matters



# Conclusions

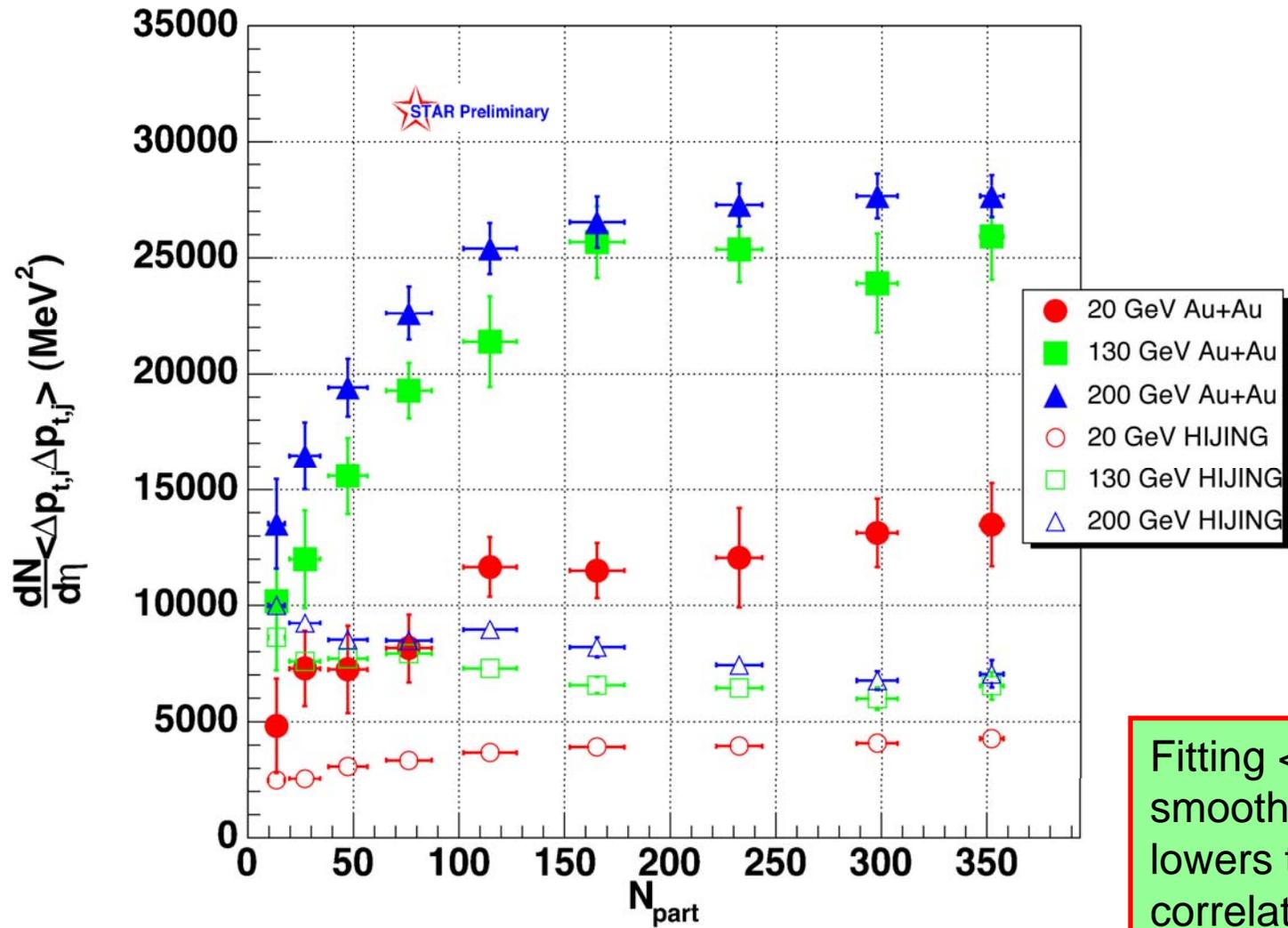
- Balance functions
  - $B(\Delta\eta)$ ,  $B(\Delta y)$ ,  $B(q_{inv})$ , and  $B(\Delta\phi)$  narrow in central Au+Au collisions
  - Narrowing of  $B(\Delta\eta)$  in central collisions consistent with trends of models incorporating late hadronization
- $p_t$  correlations
  - Dynamic correlations observed at Au+Au collisions at 20, 130, and 200 GeV
  - Correlations/particle increase with incident energy
  - Correlations/pair show little incident energy dependence
    - May show onset of equilibration in central Au+Au collisions



# Extra Slides



# Use Fit Method for $\langle \Delta p_{t,i} \Delta p_{t,j} \rangle$



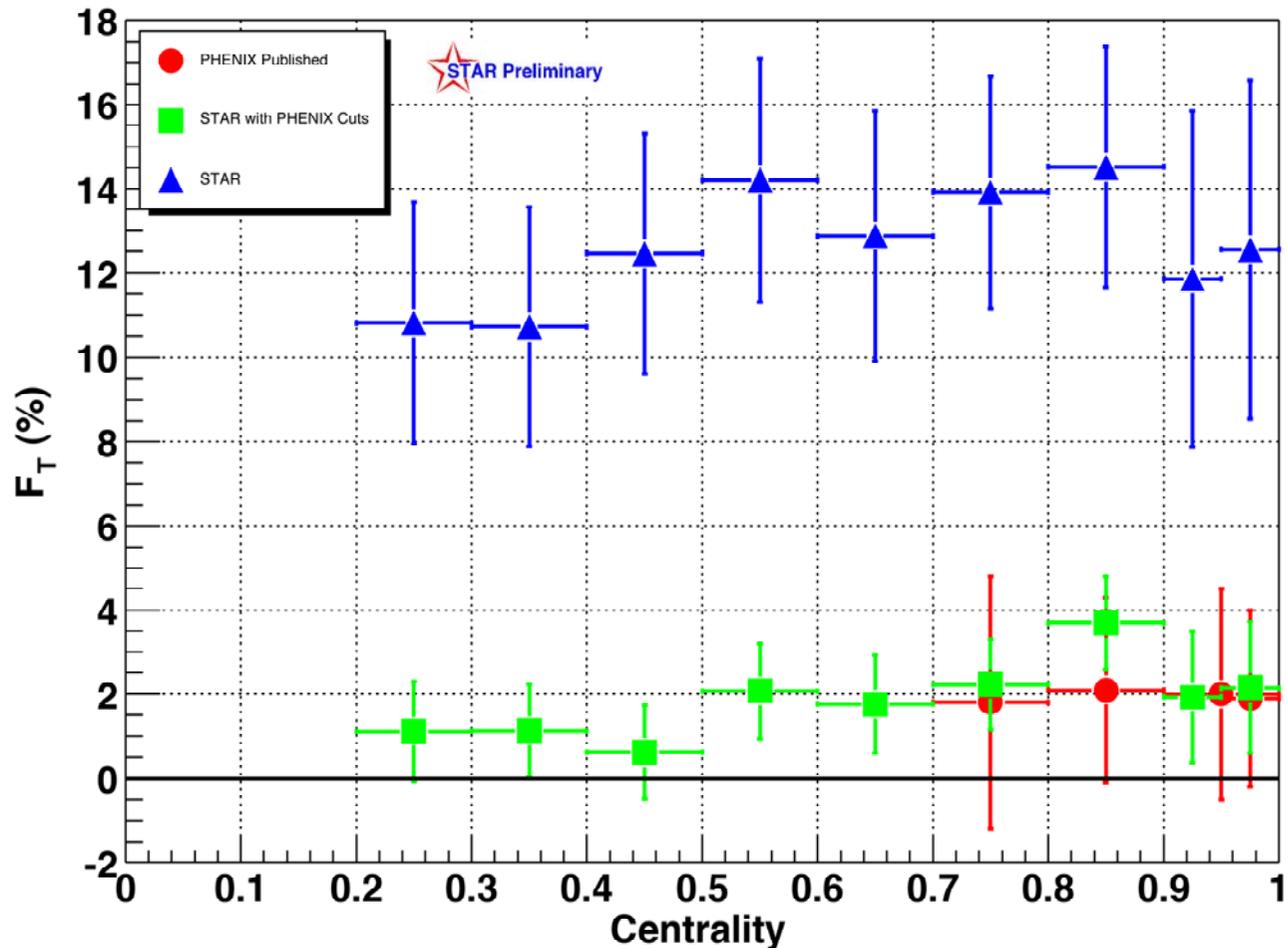
# Comparison to $F_{pt}$ at 130 GeV

$$F_T = \frac{p_T}{Q_T} \quad E_{br} = \frac{\omega_T^{mixed}}{\omega_T^{real} - \omega_T^{mixed}} = \frac{Q_T^{mixed}}{Q_T^{real} - Q_T^{mixed}}$$

130 GeV Au+Au  
STAR Cuts  
 $|\eta| < 1.0$   
 $\Delta\phi = 360^\circ$   
 $0.1 < p_t < 2 \text{ GeV}$

130 GeV Au+Au  
STAR with  
PHENIX Cuts  
 $|\eta| < 0.35$   
 $\Delta\phi = 90^\circ$   
 $0.2 < p_t < 2 \text{ GeV}$

Acceptance  
matters more



# Dynamical Net Charge Fluctuations

Definition:  $V_{+-,dyn} = V_{+-} - V_{+-,stat}$

$$v_{+-} = \left\langle \left( \frac{N_+}{\langle N_+ \rangle} - \frac{N_-}{\langle N_- \rangle} \right)^2 \right\rangle \xrightarrow{\text{Independent Particle (Poisson) Limit}} v_{+-,stat} = \frac{1}{\langle N_+ \rangle} + \frac{1}{\langle N_- \rangle}$$

Measurement: 
$$V_{+-,dyn} = \frac{\langle N_+(N_+ - 1) \rangle}{\langle N_+ \rangle^2} + \frac{\langle N_-(N_- - 1) \rangle}{\langle N_- \rangle^2} - 2 \frac{\langle N_+ N_- \rangle}{\langle N_+ \rangle \langle N_- \rangle}$$

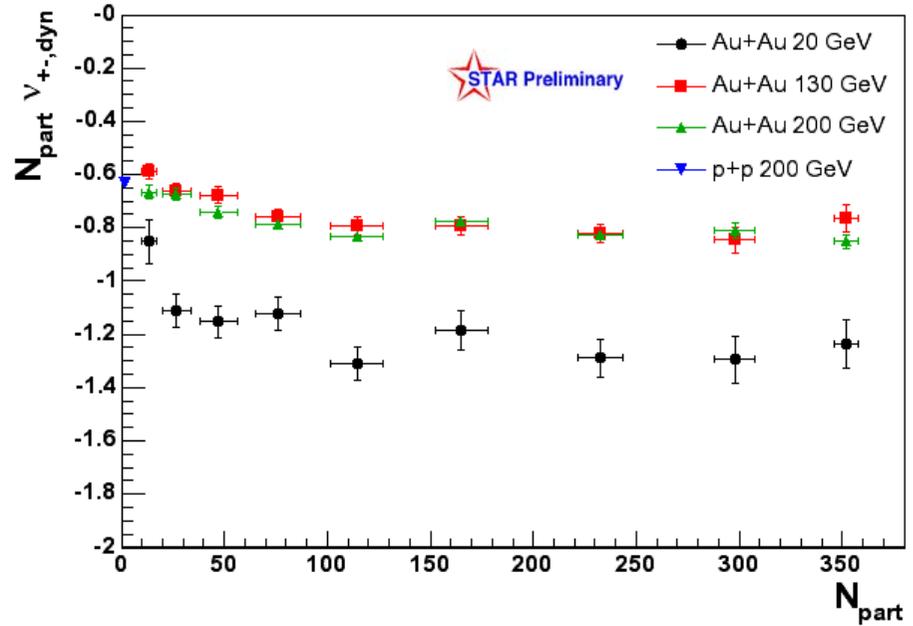
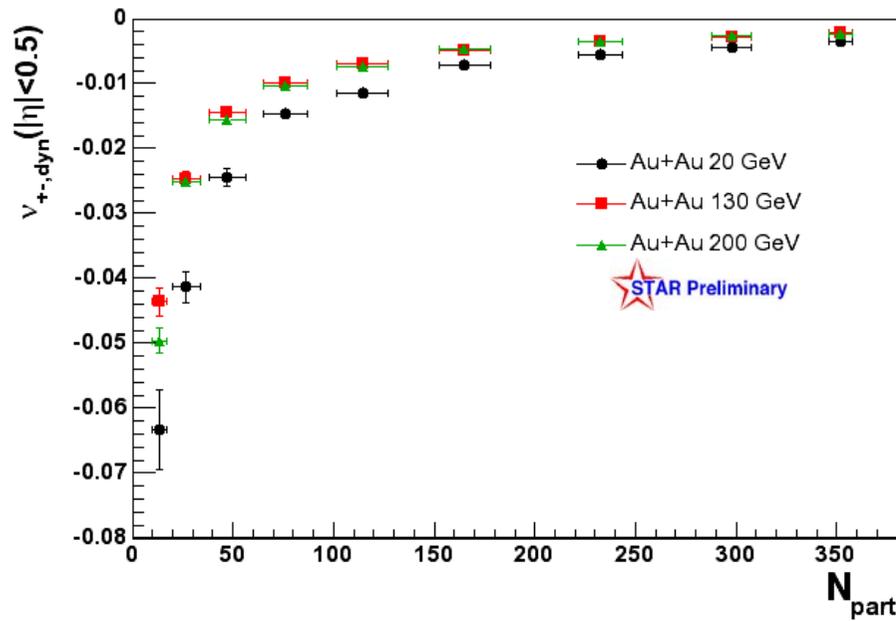
Key Properties:  $V_{+-,dyn} = R_{++} + R_{--} - 2R_{+-}$       Sensitive to two-particle correlations

$V_{+-,dyn} = 0$       Independent Particle Production

$N_{sources} \langle b \rangle v_{+-,dyn}^{A+A} \langle b \rangle = v_{+-,dyn}^{p+p}$       Independent N+N Collisions



# Net Charge Fluctuations - Centrality Dependence



$$V_{+-,dyn} = V_{+-} - V_{+-,stat}$$

$$v_{+-} = \left\langle \left( \frac{N_+}{\langle N_+ \rangle} - \frac{N_-}{\langle N_- \rangle} \right)^2 \right\rangle$$

$$V_{+-,stat} = \frac{1}{\langle N_+ \rangle} + \frac{1}{\langle N_- \rangle}$$

$$V_{+-,dyn} = R_{++} + R_{--} - 2R_{+-}$$

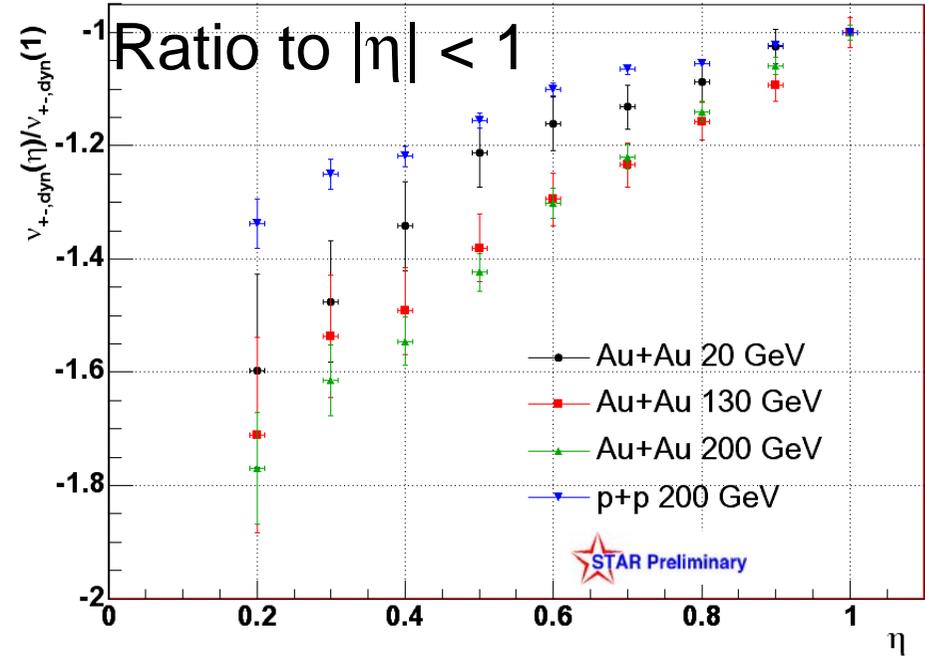
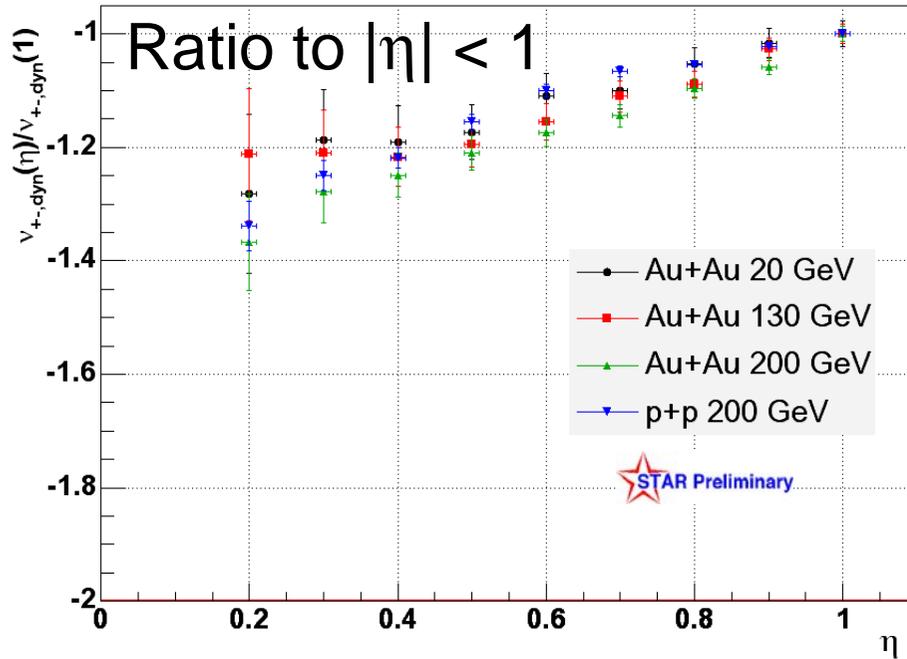
Increased dilution of correlation with increasing  $N_{part}$   
 $|v_{+-,dyn}|$  larger at 20 GeV than 130 and 200 GeV  
 Peripheral Au+Au in agreement with inclusive p+p  
 1/N scaling violation



# Net Charge Fluctuations - Dependence on $\eta$ Acceptance

50-70% Central

0-10% Central



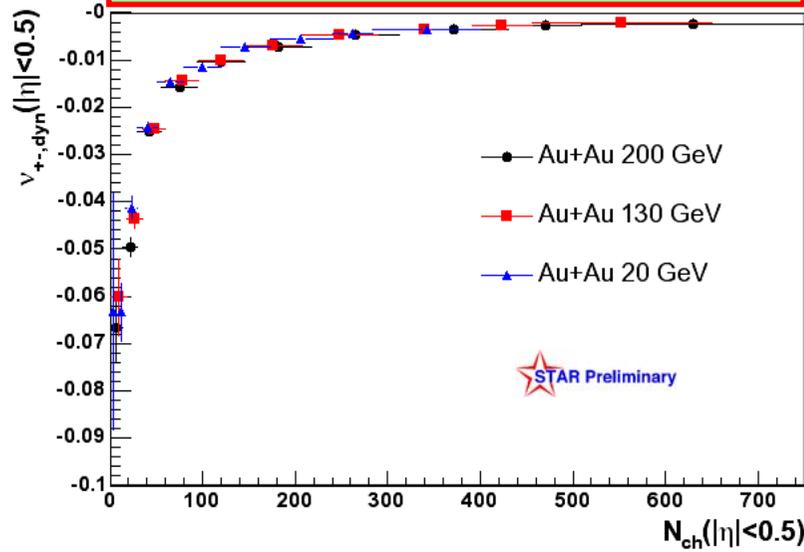
$|\eta|$  acceptance

50-70% Au+Au quite similar to p+p  
 0-10% Monotonic increase of correlation strength with beam energy  
 Flow is important (sensitivity to velocity profile)

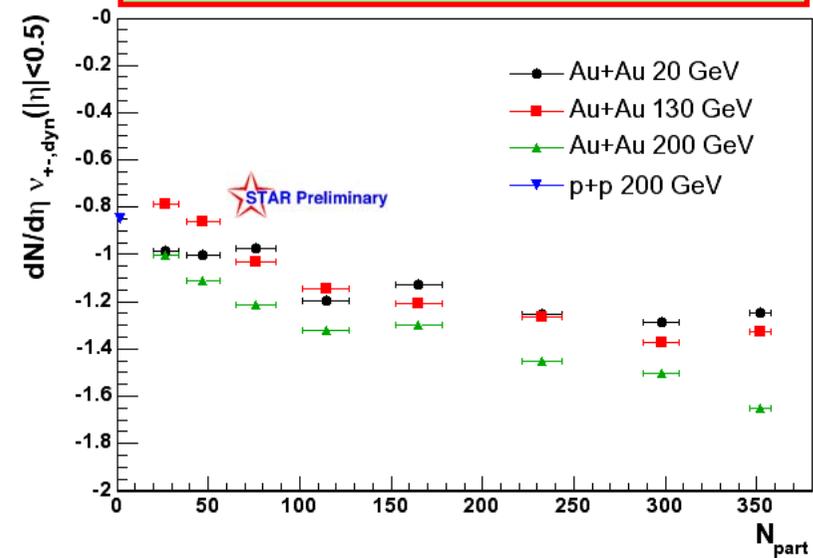


# Net Charge Fluctuations - Scaling

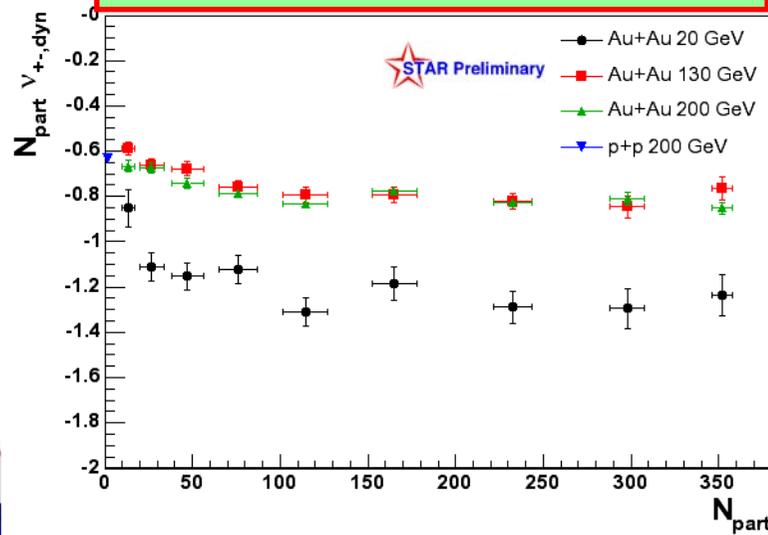
Multiplicity Scaling



$dN/d\eta$  Scaling



Participant Scaling



Binary Scaling

