#### An overview of Cosmology

CERN Student Summer School
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#### What is Cosmology?

- □ Astrophysics ⇒ detailed description of « small » structures
- □ Cosmology ⇒ Universe as a whole
  - ➤ Is it static? Expanding?
  - ➤ Is it flat, open or closed?
  - > What is it composed of?
  - What about its past and future ?

#### □ Part I: the expanding Universe

- > Hubble law
- Newtonian gravity
- General relativity
- > Friedmann-Lemaître model
- □ Part II: the standard cosmological model
  - > Hot Big bang scenario
  - Cosmological perturbations
  - Cosmological parameters
  - ➤ Inflation & Quintessence

Geometry and abstraction ...

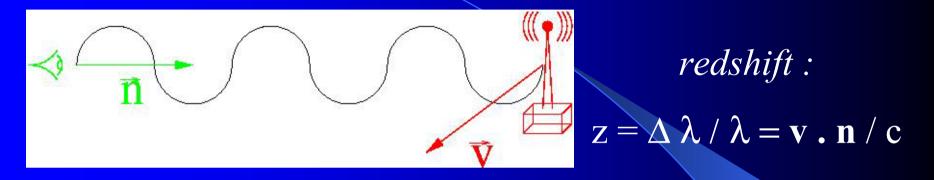
Concrete predictions, results, observations!!

# Part I: The Expanding Universe

#### Part I: (1) - The Hubble Law

- □ First step in understanding the Universe...
  - > first telescopes: observation of *nebulae*
  - > 1750 : T. Wright : Milky way = thin plate of stars?
  - > 1752 : E. Kant : nebulae = other galaxies?
- □ Galactic structure not tested before... 1923!

> 1842 : Doppler effect for sound and light



- ➤ 1868 : Huggins finds redshift of star spectral lines
- ➤ 1868 1920 : observation of many redshift of stars and nebulae
  - random distribution
  - late observations :  $excess \ of \ z > 0 \ for \ nebulae$

#### □ 1920's : Leavitt & Shapley :

➤ cepheids ⇒ period / absolute luminosity relation



absolute luminosity

apparent luminosity

$$1 = dL/ds = L/(4\pi r^2)$$

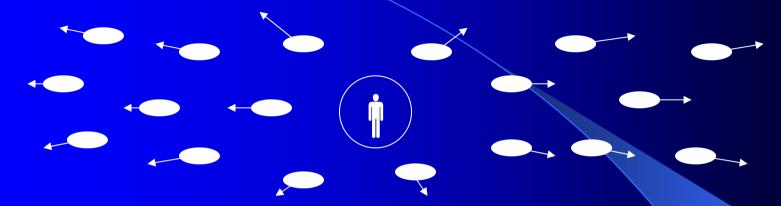
measurement of distances of stars inside the Milky Way (~80.000 lightyear)

- □ 1923 : Edwin Hubble :
  - > 2,50 m telescope at Mount Wilson (CA)
  - > cepheids in Andromeda
  - distance of nearest galaxy = 900.000 lyr
    (in fact 2 Mlyr)
    - *⇒ first probe of galactic structure* !!!!
  - > so : excess of redshifted galaxies

①

Universe expansion???

□ IN GENERAL : expansion ⇒ center



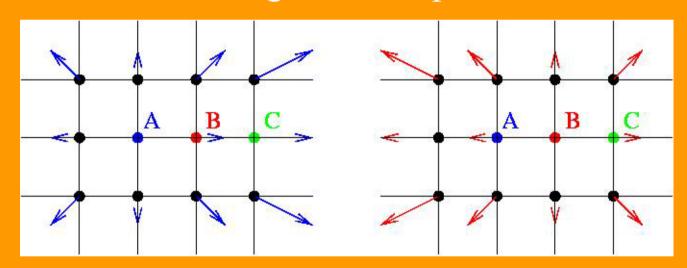
- □ Against « cosmological principle » (Milne):
  - Universe homogeneous ...
  - no privileged point!

QUESTION: is any expansion a proof against homogeneity?

#### ANSWER: not if $\mathbf{v} = \mathbf{H} \mathbf{r} \Leftrightarrow linear expansion$

... like infinite rubber grid stretched in all directions ...

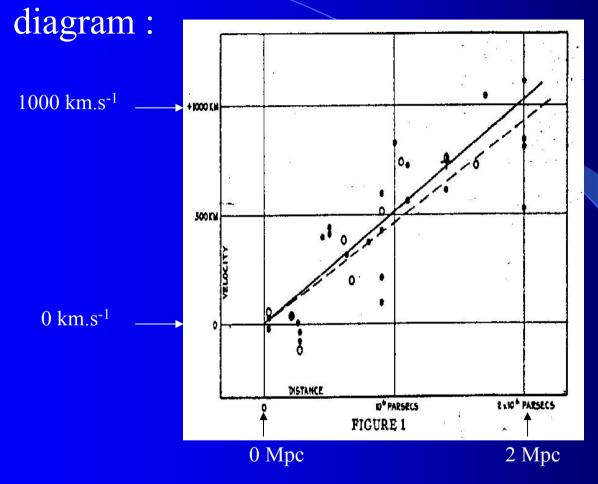
### Proof that *linear expansion* is the only possible homogeneous expansion :



 $\mathbf{v}_{\mathbf{B}/\mathbf{A}} = \mathbf{v}_{\mathbf{C}/\mathbf{B}}$ 

- $\mathbf{v}_{\mathbf{C}/\mathbf{A}} = \mathbf{v}_{\mathbf{C}/\mathbf{B}} + \mathbf{v}_{\mathbf{B}/\mathbf{A}} = 2 \mathbf{v}_{\mathbf{B}/\mathbf{A}} \Rightarrow \text{linearity}$

□ 1929: Hubble gives the first *velocity / distance* 



 $H = v / r = 500 \text{ km.s}^{-1}.\text{Mpc}^{-1}$  for Hubble (  $\cong 70 \text{ km.s}^{-1}.\text{Mpc}^{-1}$  for us ) 1 Mpc =  $3.10^6 \text{ lyr} = 3.10^{22} \text{ m}$ 

## ⇒ THE UNIVERSE IS IN HOMOGENEOUS EXPANSION

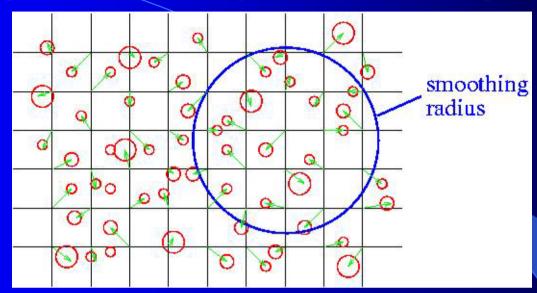
□ 1929 : starting of modern cosmology ...

Remark: what do we mean by

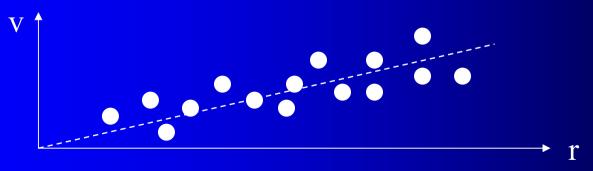
« the Universe is homogeneous »

(cosmological principle)?

> example of structure homogeneous *after smoothing*:



- ➤ local inhomogeneities ⇒ scattering



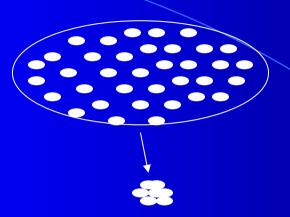
# Part I: (2) – Universe expansion from Newtonian gravity

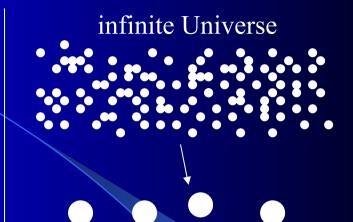
- on cosmic scales, only gravitation
- □ Newton's law = limit of General Relativity (GR)

```
when v \ll c
F = G m_1 m_2 / r^2
OR | speed of object | speed of liberation
```

- > Newton's law should describe expansion at small distances with v = H r << c ...
- but historically, GR proposed the first predictions / explanations !!!

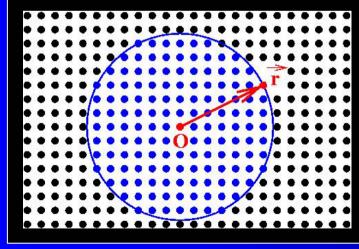
□ Newton: finite Universe





> but how to deal with infinity?

□ Gauss theorem:



$$\dot{\mathbf{r}} = -G \, \mathbf{M}_{\mathbf{r}} / \mathbf{r}^2$$

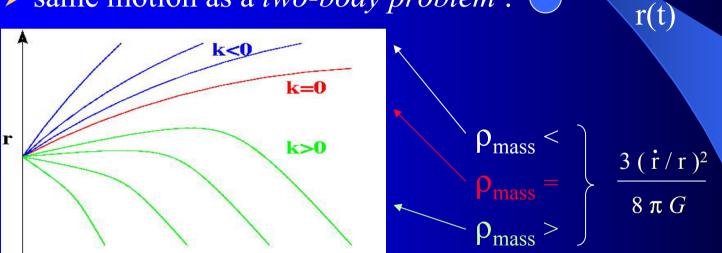
$$M_r = constant = (4/3) \pi r^3 \rho_{mass}$$

$$\Rightarrow$$
  $\dot{\mathbf{r}}^2 = 2 G M_r / r - k$ 

$$= (8/3) \pi G \rho_{\text{mass}} r^2 - k$$

#### Newtonian expansion law : $(\dot{r}/r)^2 = (8\pi G/3) \rho_{mass} - k/r^2$

- $\rho_{\text{mass}}(t) \propto r(t)^{-3}$
- > same motion as a two-body problem :



- $ightharpoonup k \neq 0$   $\Rightarrow$  non-homogeneous expansion ???
- $ightharpoonup v = H r \text{ and } v << c \Rightarrow r < R_H \equiv c / H$

GRAVITY / INERTIA

## Part I: (3) – General Relativity and the Friedmann-Lemaître Universe

Newtonian gravity



invariant speed of light

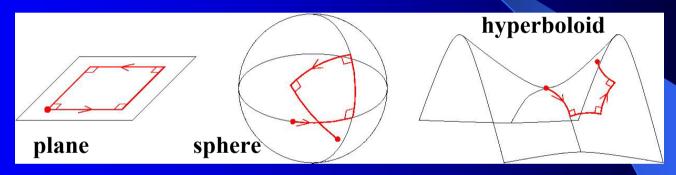
General Relativity

(Einstein 1916)

- $\Rightarrow$  no more  $\Phi_{\text{grav}}$  ( matter distribution  $\Leftrightarrow \Phi_{\text{grav}} \Leftrightarrow \mathbf{E} = \nabla \Phi_{\text{grav}}$ )
- $\Rightarrow$  three basic principles :
  - 1) space-time (t,x,y,z) is curved
  - 2) curvature  $\Leftrightarrow$  matter
  - 3) *free-falling bodies* follow *geodesics*

#### 1) how can we define the *curvature*:

- of a 2-D surface ?
  - embedded in 3D
  - stay in 2D, and use angles :

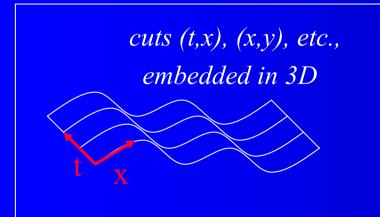


• stay in 2D, and use a scaling law :  $dl(x_1,x_2)$ 

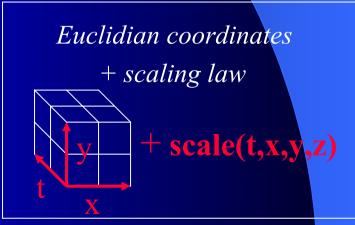


ex: sphere projected on ellipse,  $dl(\theta)$ 

- > of 3-D space?
  - embedded in 4D:  $x_1^2 + x_2^2 + x_3^2 + x_4^2 = R^2$
  - stay in 3D, but provide a scaling law, like on a planisphere: dl(x<sub>1</sub>,x<sub>2</sub>,x<sub>3</sub>)
- > of 4-D space-time?
  - one more dimension
  - time different from space (special relativity: -+++) intuitive representations:



OR



#### 2) curvature $\Leftrightarrow$ matter:



mathematical formulation = Einstein equation

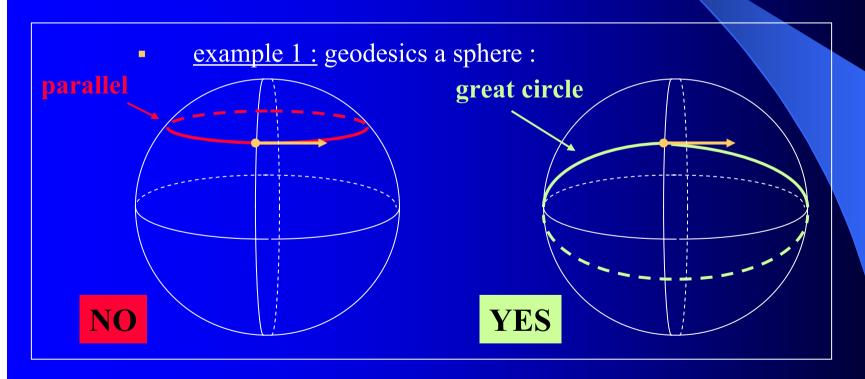
#### 3) free-falling bodies follow geodesics

feel only gravity
(not E.M., etc.)
e.g. galaxies, light...

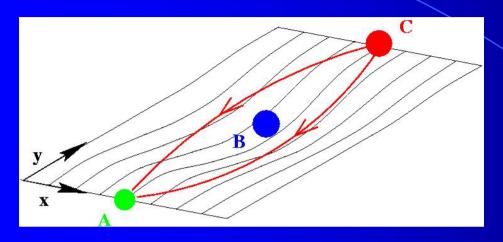
given one point and one direction

 $\Rightarrow$  one single line such that :

 $\forall$  A, B, [AB] = shortest trajectory

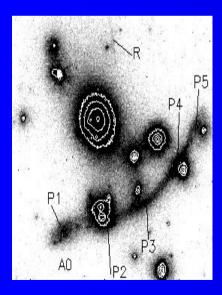


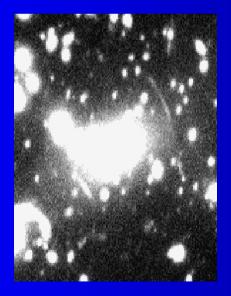
#### • example 2 : gravitational lensing :

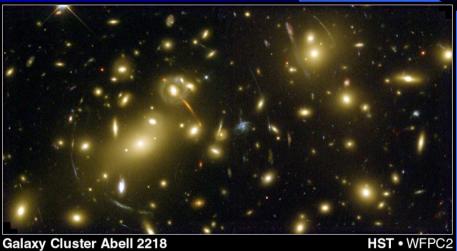


A sees an image of C lensed by B:





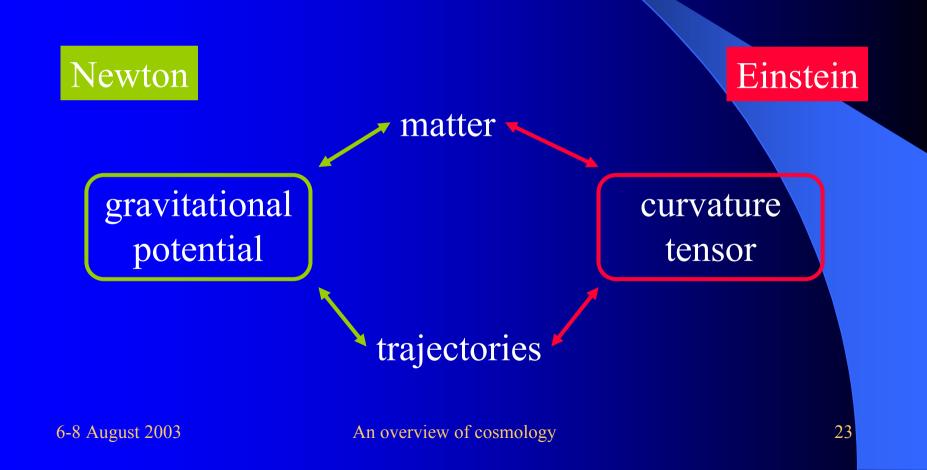




NASA, A. Fruchter and the ERO Team (STScl, ST-ECF) • STScl-PRC00-08

#### □ Newtonian gravity versus G.R.:

two different theories of gravity, i.e. two ways of describing how the presence of matter affects the trajectories of surrounding bodies...



- applying G.R. to the Universe: some history
  - > 1916: Einstein has formulated G.R.
  - ➤ 1917 : Einstein, De Sitter try to build the first cosmological models (PREJUDICE : STATIC / STATIONNARY UNIVERSE)

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> 1922 : A. Friedmann (Ru)
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> 1927 : G. Lemaître (B)

➤ 1933 : { Robertson, Walker (USA)

investigate most general

HOMOGENEOUS, ISOTROPIC,

NON-STATIONNARY

solutions of G.R. equations

- > 1929 : Hubble's law (first confirmation)
- > 1930-65 : accumulation of proofs in favour of FLRW
- > 1965 : CMB discovery : full confirmation

basic principles of G.R.	FLRW solution
space-time is <i>curved</i>	??????
free-falling objects follow geodesics	???????
curvature caused by / related to matter	??????

summary of the situation :

#### **NEWTON**

 $\Rightarrow$  matter distribution  $\Leftrightarrow \Phi_{\text{grav}} \Leftrightarrow$  forces changing the trajectories

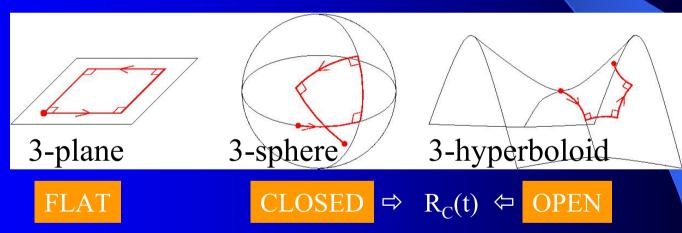
#### General Relativity (GR)

- $\Rightarrow$  three basic principles :
  - 1) space-time (t,x,y,z) is curved
  - 2) free-falling bodies follow geodesics
  - 3) curvature ⇔ matter

Friedmann-Lemaitre model = application of GR to homogeneous Universe

#### 1) the *curvature* of the FLRW Universe:

- Universe space-time (t,x,y,z) curved by its *own* homogeneous matter density  $\rho(t)$
- ► HOMOGENEITY ⇒ decomposition of curvature in :
  - spatial curvatureof (x,y,z) at fixed t3-D space is maximally symmetric :



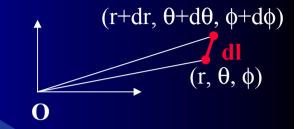
2. 2-D space-time curvature

of  $(t,x) \Leftrightarrow (t,y) \Leftrightarrow (t,z)$  accounts for the expansion

#### □ scale as a function of coordinates?

- $\triangleright$  COMOVING COORDINATES (t, r,  $\theta$ ,  $\phi$ )
- > for Euclidian space:

$$dl^2 = dr^2 + r^2 (d\theta^2 + \sin^2\theta d\phi^2)$$



for 
$$FLRW$$
:  
 $dl^2 = a^2(t) \left[ \frac{dr^2}{(1-k r^2)} + r^2 (d\theta^2 + \sin^2\theta d\phi^2) \right]$ 

- $a(t) \equiv \text{scale factor} \Rightarrow 2\text{-D space-time curvature}$
- k ⇒ spatial curvature

$$-k=0$$
: FLAT

$$-k > 0$$
: CLOSED,  $R_C(t) = a(t) / k^{1/2}$ ,  $0 \le r < 1 / k^{1/2}$ 

$$R_C(t) = a(t) / k^{1/2}$$

$$0 \le r < 1/k^{1/2}$$

$$-k < 0 : OPEN,$$
  $R_C(t) = a(t) / (-k)^{1/2}$ 

$$R_C(t) = a(t) / (-k)^{1/2}$$

#### 2) the geodesics in the FLRW Universe

photons: v = c Ultra-relativistic ordinary matter: v << c Non-relativistic (e.g. galaxies)

1) non-relativistic matter:

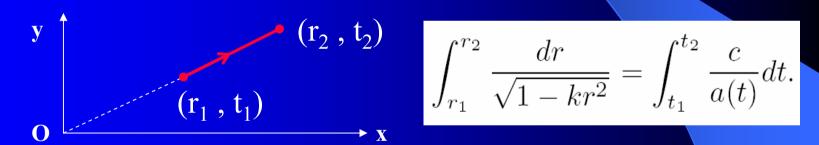
$$dl = 0 \Rightarrow (r, \theta, \phi) = constant$$

- galaxies are still in coordinate space ...
- ... but all distances are proportional to a(t)
  - a(t) gives the expansion between galaxies (although they are still !!!)
- like an inflated rubber balloon with points drawn on its surface ...

2) Relativistic matter: straight line in 3-D space, dl = c dt

$$\Rightarrow c^{2}dt^{2} = a^{2}(t) \left[ \frac{dr^{2}}{(1-k r^{2})} + r^{2} (d\theta^{2} + \sin^{2}\theta d\phi^{2}) \right]$$
EQUATION OF PROPAGATION OF LIGHT

• So  $\Delta 1 = c \Delta t$  is *WRONG*:



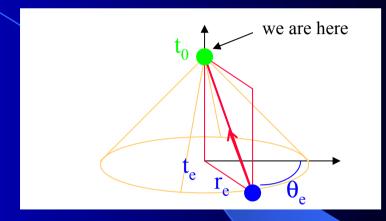
- ⇒ « bending of light in the Universe »
   (one of the two most fundamental equations in cosmology)
- ⇒ various important consequences ...

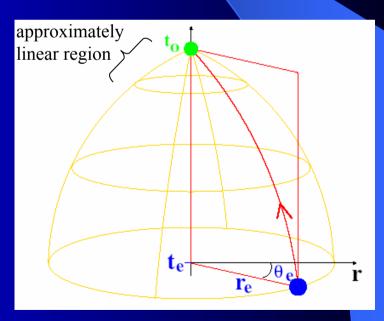
#### definition of the past light-cone :

- > in Euclidian space:
  - $\theta = constant$
  - $r_e = c (t_0 t_e)$
- > in Friedmann universe:
  - $\theta = constant$

$$\int_{r_e}^{0} \frac{-dr}{\sqrt{1 - kr^2}} = \int_{t_e}^{t_0} \frac{c}{a(t)} dt$$

$$\frac{dr}{dt} = -\frac{c}{a(t)} \frac{(1-kr)^{1/2}}{a(t)}$$



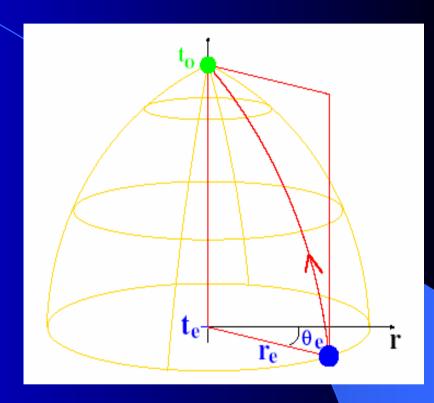


#### observable consequences of propagation of light equation:

#### > the redshift:

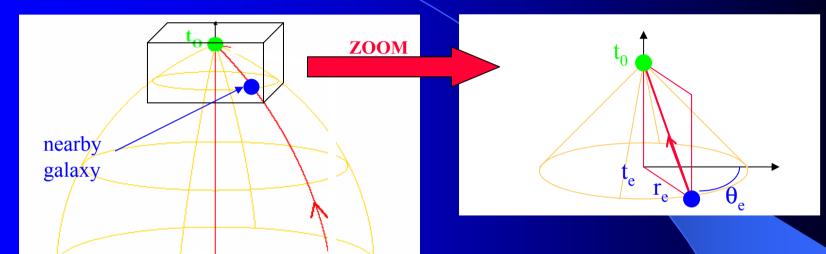
$$z = \Delta \lambda / \lambda = \lambda_0 / \lambda_e - 1$$

$$z = a(t_0) / a(t_e) - 1$$



- remark 1:
  - *Newtonian* :  $z = v / c \le 1$
  - *G.R.* : no limit, as observed ...

remark 2: at short distance, we can recover the *Hubble law* (z = v/c = Hr/c)



$$t_0 - t_e = dt = dl / c$$

• then: 
$$z = \frac{a(t_0)}{a(t_0 - dt)} - 1 = \frac{1}{1 - \frac{\dot{a}(t_0)}{a(t_0)}dt} - 1 = \frac{\dot{a}(t_0)}{a(t_0)}dt = \frac{\dot{a}(t_0)}{a(t_0)}\frac{dl}{c}$$

• so : Hubble parameter = 
$$H(t) = \frac{\dot{a}(t)}{a(t)}$$
 (H<sub>0</sub> = 70 km.s<sup>-1</sup>.Mpc <sup>-1</sup>)

#### 2) the angular diameter-redshift relation



 $d\theta$ 



• Euclidian space:

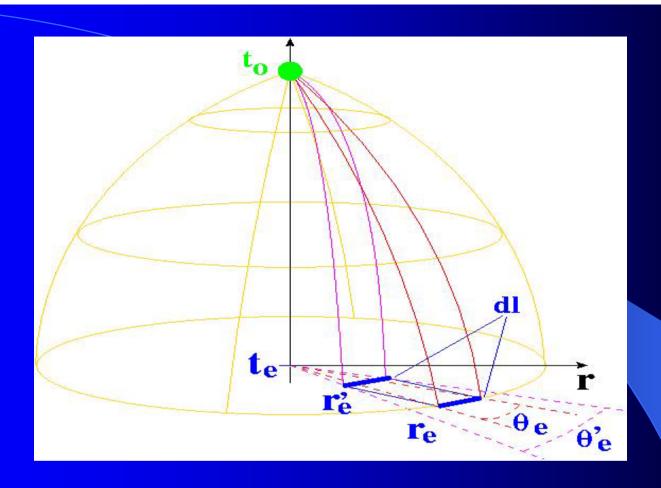
$$dl = r d\theta$$
 with  $r = v/H = z c/H$ 

• *G.R.* :

$$dl = a(t_e) r_e d\theta$$
 with  $r_e$  from

$$dl = a(t_e) r_e d\theta$$
 with  $r_e$  from 
$$\int_{r_e}^{0} \frac{-dr}{\sqrt{1 - kr^2}} = \int_{t_e}^{t_0} \frac{c}{a(t)} dt$$

if dl is known, measurement of  $(d\theta, z) \Rightarrow k$ , a(t)



- CLOSED UNIVERSE: objects seen under larger angle
- > OPEN UNIVERSE : objects seen under smaller angle

#### 3) the luminosity distance-redshift relation



apparent luminosity L absolute luminosity

• Euclidian space:  $l = \frac{L}{4\pi r^2}$ 

$$l = \frac{L}{4\pi r^2}$$

with r = z c / H

• G.R.: 
$$l = \frac{L}{4\pi a^2 (t_0) r_e^2 (1+z)^2}$$

with  $r_{\rho}$  from prop. of light

if L is known, measurement of  $(l, z) \Rightarrow k, a(t)$ 

3) relation between matter and curvature:

## FRIEDMANN LAW

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi\mathcal{G}}{3}\frac{\rho}{c^2} - \frac{kc^2}{a^2}$$

Remark : for non-relativistic matter,  $E = m c^2 \Rightarrow \rho / c^2 = \rho_{mass}$ 

- ⇒ then *Friedmann law* looks similar to *Newtonian expansion law*, but <u>CRUCIAL DIFFERENCES</u>:
  - 1)  $a(t) \neq r(t)$ : very different interpretation
  - 2)  $k \neq 0$  not in contradiction with homogeneity
  - 3) accounts for non-relativistic and relativistic matter

#### NON-RELATIVISTIC

$$v \ll c$$
 $E = m c^2$ 

#### **ULTRA-RELATIVISTIC**

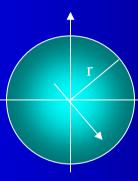
$$v = c$$

$$E = \hbar v = \hbar c / \lambda$$

sphere with

fixed comobile r

fixed particle number



$$V = 4/3 \pi r^3 a(t)^3$$

$$E_V = constant$$
  
 $\rho = E_V / V \propto a(t)^{-3}$ 

$$E_{V} \propto 1 / \lambda \propto 1 / a(t)$$
  
 $\rho = E_{V} / V \propto a(t)^{-4}$ 

FRIEDMANN LAW is the same but DILUTION RATE is different

in fact, in G.R., curvature  $\Leftrightarrow$  matter relation given by

Einstein equation 
$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

in the FLRW solution :

EINSTEIN EQUATION ⇒

Friedmann law

conservation equation:

$$\dot{\rho} = -3 (\dot{a}/a) (\rho + p)$$

1) non-relativistic:

$$\mathbf{v} \ll \mathbf{c} \implies \mathbf{p} \cong \mathbf{0} \implies \dot{\mathbf{p}} / \mathbf{p} = -3 \dot{\mathbf{a}} / \mathbf{a} \implies \mathbf{p} \ll \mathbf{a}(\mathbf{t})^{-3}$$

2) ultra-relativistic:

$$v = c \implies p = \rho / 3 \implies \dot{\rho} / \rho = -4 \dot{a} / a \implies \rho \propto a(t)^{-4}$$

3) in QFT: vacuum with

$$p = -\rho \implies \dot{\rho} = 0 \implies \rho = constant$$

« cosmological constant »

# summary of the situation :

basic principles of G.R.	FLRW solution
space-time is <i>curved</i>	2 types of curvatures : {a(t), k} $dl^2 = a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2\theta \ d\phi^2) \right]$
free-falling objects follow geodesics	<ul> <li>non-relativistic matter : dl = 0</li> <li>ultra-relativistic matter : dl = c dt</li> <li>BENDING OF LIGHT EQUATION</li> </ul>
curvature caused by / related to matter	relation {a(t), k} ⇔ ρ(t)  ♣ FRIEDMANN LAW  + conservation equations

# Part II: The Standard Cosmological Model

#### decomposition of quantities:

**BACKGROUND** 

THEORY OF LINEAR PERTURBATIONS

Initial conditions
INFLATION
part II - 4

HOMOGENEOUS COSMOLOGY part II - 1

LINEAR PERT. part II - 2

NON-LINEAR PERT.

# Part II: (1) – Homogeneous cosmology

- the evolution of the Universe depends :
  - > on SPATIAL CURVATURE
  - > on the density of:
    - - RADIATION : ultra-relativistic particles  $\rho \propto a^{-4}$  $p = \rho / 3$

(photons, massless v's, ...)

MATTER : non-relativistic bodies

$$p = 0$$
  $\rho \propto a^{-3}$  (galaxies, gas clouds, ...)

COSMOLOGICAL CONSTANT A

$$p = -\rho$$
  $\rho = constant = \Lambda c^2 / (8\pi G)$   
(vacuum?...?)

6-8 August 2003

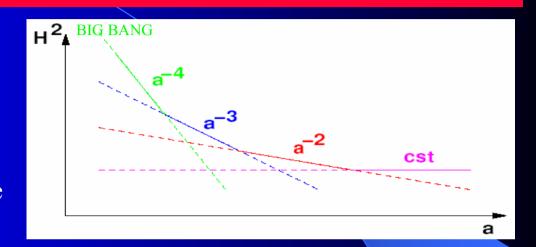
An overview of cosmology

Friedmann

law:

$$H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = \frac{8\pi\mathcal{G}}{3c^{2}}\rho_{R} + \frac{8\pi\mathcal{G}}{3c^{2}}\rho_{M} - \frac{kc^{2}}{a^{2}} + \frac{\Lambda}{3}$$

- most « complete » scenario :
- phases can be skipped,but order cannot change



- > RADIATION DOMINATION :  $a \propto t^{1/2}$  H = 1/2
- MATTER DOMINATION :  $a \propto t^{2/3}$  H = 2 / 3 t
- > CURVATURE DOMINATION :
  - k < 0 (open) :  $a \propto t$  H = 1/t
  - k > 0 (closed):  $\dot{a} \rightarrow 0$ , then  $\dot{a} < 0$  or
- $\triangleright$  VACUUM DOMINATION : a  $\propto$  exp(Λ/3t)<sup>1/2</sup> H → constant

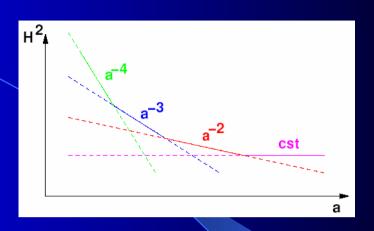
#### Future of the Universe:

- $\triangleright$  if  $\Lambda = 0$ :
  - •if k < 0 or  $k = 0 \rightarrow$  indefinite decelerated expansion
  - •if k > 0  $\rightarrow$  recollapse (BIG CRUNCH)
- $\triangleright$  if  $\Lambda \neq 0$ :
  - $\blacksquare \forall k$

- → indefinite *accelerated* expansion
- > RADIATION DOMINATION :  $a \propto t^{1/2}$  H = 1/2
- > MATTER DOMINATION :  $a \propto t^{2/3}$  H = 2 / 3 t
- > CURVATURE DOMINATION :
  - k < 0 (open) :  $a \propto t$  H = 1/t
  - $k \ge 0$  (closed) :  $\dot{a} \to 0$ , then  $\dot{a} < 0$  or
- ► VACUUM DOMINATION :  $a \propto exp(\Lambda/3t)^{1/2}$  H  $\rightarrow constant$

#### □ the matter budget:

if we can measure  $\{\rho_R, \rho_M, k, \Lambda\}$  today, we can extrapolate back ...



> today:

$$1 = \frac{8\pi \mathcal{G}}{3H_0^2c^2}\left(\rho_{\rm R0} + \rho_{\rm M0}\right) - \frac{kc^2}{a_0^2H_0^2} + \frac{\Lambda}{3H_0^2} \ = \ \frac{\Omega_{\rm R} + \Omega_{\rm M} - \Omega_k + \Omega_{\Lambda}}{\text{MATTER BUDGET EQUATION}}$$

> flatness condition :  $\Omega_0 \equiv \Omega_{\rm R} + \Omega_{\rm M} + \Omega_{\Lambda} = 1$ 

► then: 
$$\rho_R^{\ 0} + \rho_M^{\ 0} + \rho_{\Lambda}^{\ 0} = \frac{3 H_0^2 c^2}{8\pi G} \equiv \rho_c^{\ 0}$$
  $\Rightarrow$   $\Omega_X = \rho_X / \rho_c^{\ 0}$ 

> so far : COSN

COSMOLOGICAL SCENARIOS



4 independent parameters

#### □ COLD or HOT BIG BANG???

> 1929–65: no decisive observation in favour of Friedmann model (apart from accumulation of redshifts)

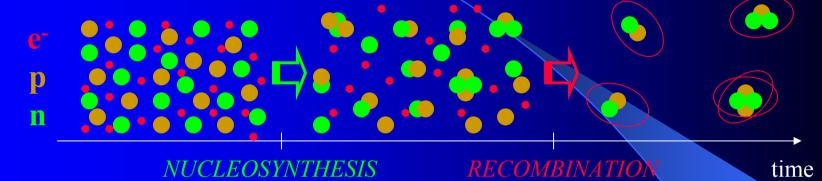
⇒ works in cosmology remain marginal

- but spectacular progress in particle physics...
- > studies based on the *most simple* possible scenario:
  - Universe contains only non-relativistic matter
  - evolution under the laws of nuclear physics between Big Bang and today

♥ COLD BIG BANG SCENARIO

#### □ COLD BIG BANG:

$$\rightarrow$$
 H<sup>2</sup> =  $(8\pi G/3c^2) \rho_M \Rightarrow \rho \propto a^{-3} \propto t^{-2}$ 



#### > NUCLEOSYNTHESIS:

• ensemble of *nuclear reactions* 

• freeze-out due to *expansion* 

## pioneering works on nucleosynthesis:

```
> 1940: Gamow et al. (USSR \rightarrow USA)
```

1964 : Zel'dovitch et al. (USSR)

1965: Hoyle & Taylor (UK)

1965: Peebles et al. (USA)

#### > COLD BIG BANG

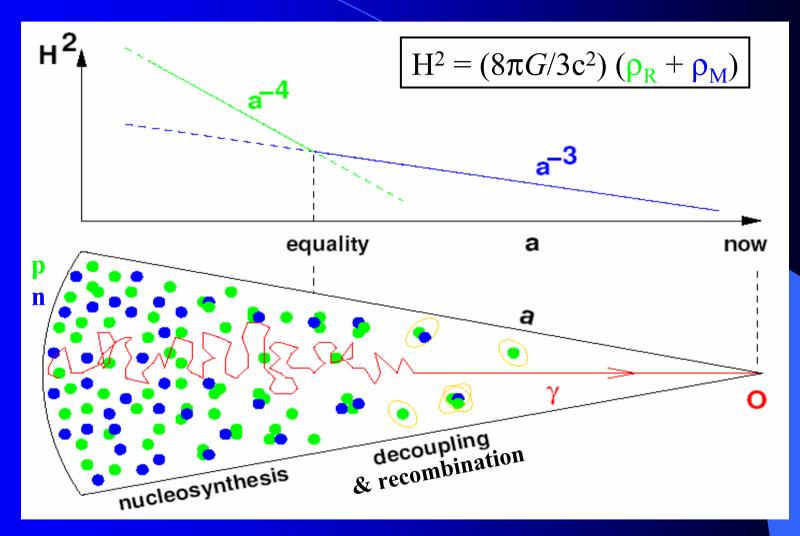
```
⋄ no hydrogen
```

 $\Rightarrow$  need to change  $H(t_{nucleo})$ 

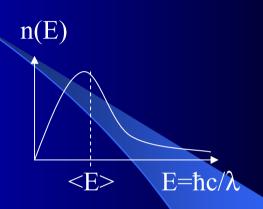
 $\Rightarrow$  add relativistic matter (photons) with  $\rho_R >> \rho_M$ 

♥ HOT BIG BANG !!!

### □ HOT BIG BANG:



- > photon spectrum:
  - before recombination, thermal equilibrium
    - ♦ blackbody spectrum :



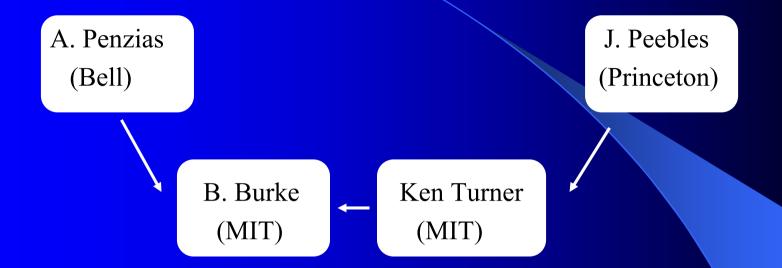
- after recombination, Planck spectrum frozen and redshifted
- so  $T_0 a_0 = T_{\text{nucleo}} a_{\text{nucleo}}$
- Gamow, Peebles et al. :

nucleosynthesis 
$$\Rightarrow$$
 T<sub>0</sub>  $\approx$  1-10 K  $\Rightarrow$   $\lambda_0 \approx$  1-10 mm

## discovery of the Cosmic Microwave Background



#### discovery of the Cosmic Microwave Background



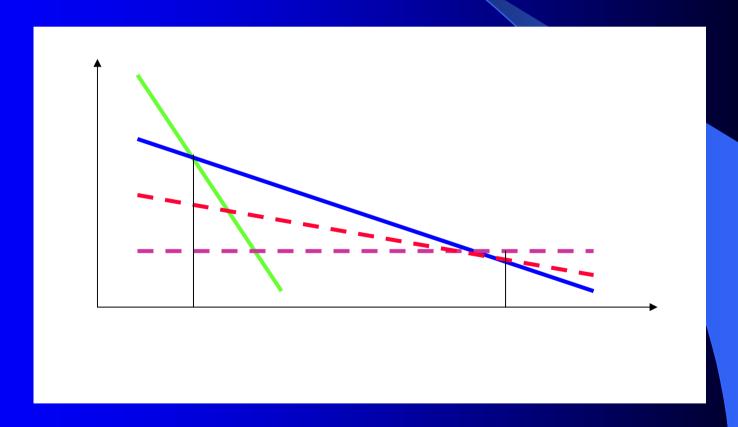
- publications by Penzias & Wilson and Peebles
- Nobel Prize for Penzias & Wilson ...
- confirmation of HOT BIG BANG !!!
- $\bigcirc$  CMB = 25% of TV set noise...
- listen to it at http://:www.bell-labs.com/project/feature/archives/cosmology/

## thermal evolution of the Universe

	t	Т	$\rho^{1/4}$
Planck time inflation?	10 <sup>-36</sup> s		$10^{18}  \mathrm{GeV}$
GUT	10 <sup>-32</sup> s		10 <sup>16</sup> GeV
<b>EW</b> bayogenesis?	10 <sup>-6</sup> s	10 <sup>15</sup> K	100 GeV
quark-hadron	10 <sup>-4</sup> s	$10^{12}\mathrm{K}$	100 MeV
nucleosynthesis	1-1000 s	10 <sup>9-10</sup> K	0.1 - 1 MeV
equality	10 <sup>4</sup> yr	$10^4  \mathrm{K}$	1 eV
decoupling	10 <sup>5</sup> yr	2500 K	0.1 eV
structure formation	$10^5 \text{ yr} \rightarrow t_0$		
today → curvature / A?	$t_0 \approx 13 \text{ Gyr}$	≈ 3 K	$\approx 3.10^{-4} \text{ eV}$

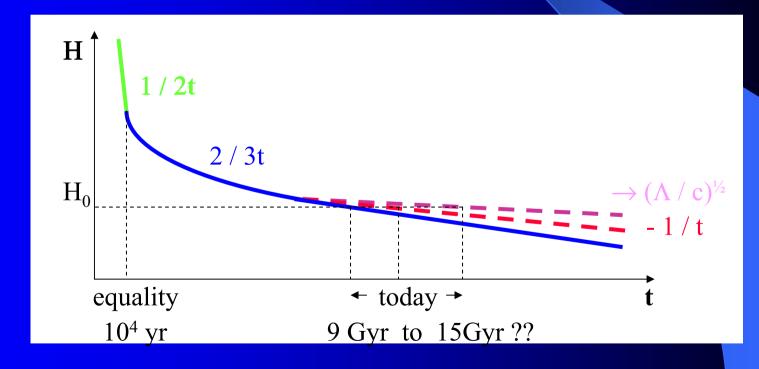
## □ is there a curvature / $\Lambda$ domination today?

- > structure formation  $\Rightarrow$  long-enough M.D.  $\Rightarrow \Omega_{\rm m} \ge 0.2$
- ightharpoonup if  $\Omega_k \sim 1$  or  $\Omega_{\Lambda} \sim 1$ , curvature /  $\Lambda$  domination started recently:



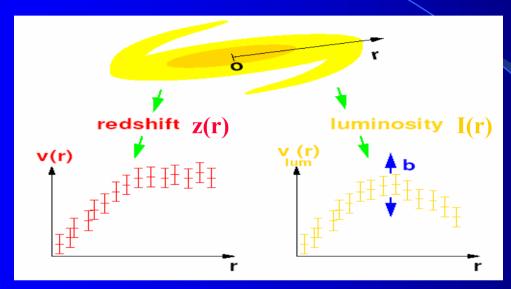
## □ is there a curvature / $\Lambda$ domination today?

- **HOW CAN WE KNOW?** 
  - $\Omega_k$   $\Rightarrow$  change  $\begin{pmatrix} k \\ a(t) \end{pmatrix}$   $\Rightarrow$   $\begin{cases} apparent luminosity / z \text{ relation} \Rightarrow \text{SNIa} \\ angular diameter / z \text{ relation} \Rightarrow \text{CMB} \end{cases}$
  - age of the Universe (measured with redshift of quasars)



#### □ DARK MATTER:

galaxy rotation curves :



$$\rho_{\text{mass}}(\mathbf{r}) = \mathbf{b} \mathbf{I}(\mathbf{r})$$

$$\Delta\Phi_{\rm grav}(r) = 8\pi G \,\rho_{\rm mass}$$

$$v^2(r) = r \left( \partial \Phi_{\text{gray}} / \partial r \right)$$

- $\Rightarrow$  DM halo
- > other compelling evidences

> nature of DM : ( non-luminous baryons ?

Hot dark matter (neutrinos)?

CDM: WIMPS (neutralinos)? axions?

□ so far:

COSMOGICAL SCENARIO



5 INDEPENDENT
PARAMETERS

 $\{\Omega_{\rm R}\,,\,\Omega_{\rm B}\,,\,\Omega_{\rm CDM}\,,\,\Omega_{\Lambda}\,,\,{\rm H}_0\}$ 

## Part II: (2) – cosmological perturbations

- Universe not completely homogeneous even on large scales ...
  - description of matter inhomogeneities ??
    (clusters of galaxies, superclusters ...)
  - description of CMB temperature anisotropies ??
    - ⇒ ∫ information on cosmological scenario measurement of cosmological parameters
- □ this section:
  - overview of mathematical framework
  - > intuitive description of main phenomena

## □ linear perturbation theory:

$$ho_{x}(t, \mathbf{r}) = \overline{\rho}_{x}(t) + \delta \rho_{x}(t, \mathbf{r})$$

#### HOMOGENEOUS BACKGROUND

PERTURBATION

- CMB temperature homogeneity
  - $\Rightarrow$  perturbations linear at least for t < t<sub>dec</sub>

$$\delta_{x}(t, \mathbf{r}) \equiv \delta \rho_{x}(t, \mathbf{r}) / \overline{\rho}_{x}(t) << 1$$

## □ Einstein equations (matter ⇔ curvature)

- background equations :
  - Friedman law
  - conservation equations

$$\{\overline{\rho}_{\gamma}, \overline{\rho}_{\nu}, \overline{\rho}_{CDM}, \overline{\rho}_{B}\} \iff \{a(t), k\}$$

> perturbation equations :

$$\{\delta_{\gamma}, \delta_{v}, \delta_{\text{CDM}}, \delta_{B}\} \Leftrightarrow \{\text{curvature perturbations}\}\$$
 $\approx \Phi(t, \mathbf{r}) \rightarrow \Phi_{\text{grav}} \text{ inside } R_{H}$ 

partial derivative equations

 $\Rightarrow$  Fourier transformation ...

## comoving Fourier space

- $\delta_{x}^{k}(t) \equiv \int dr^{3} e^{-i k.r} \delta_{x}(t, r)$ 
  - comoving Fourier wavenumber k
  - comoving wavelength  $2\pi / k$
  - physical wavelength  $\lambda(t) = 2\pi a(t) / k$
- > independent perturbation equations :

$$\{\delta_{\gamma}^{\ \kappa},\,\delta_{v}^{\ \kappa},\,\delta_{\mathrm{CDM}}^{\ \kappa},\,\delta_{\mathrm{B}}^{\ \kappa}\} \iff \Phi^{\kappa}$$

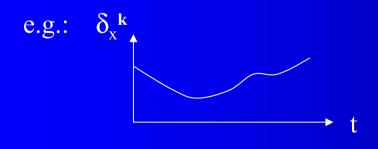
- linear system of ordinary differential equations...
- even at equilibrium, each wavelength redshifted

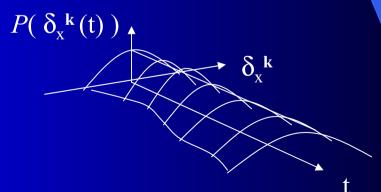
## □ stochastic theory:

- > random initial conditions :  $P(\delta_x^k(t_0))$
- > evolution under differential equation :

$$(d^2/dt^2) \delta_x^{k} + ... (d/dt) \delta_x^{k} + ... \delta_x^{k} = ... \delta_y^{k} + ... \Phi^{\kappa}$$

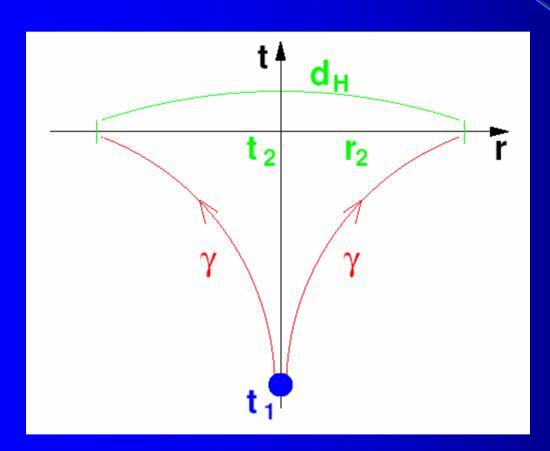
- Early Universe ⇒ gaussian distributions
- linearity : shape  $P(\delta_x^k(t))$  is preserved
- differential equation = evolution of r.m.s.





## □ intuitive description of the evolution :

> definition of the HORIZON:



$$d_H(t_1, t_2) = 2 \int_0^{r_2} dl = 2 \int_0^{r_2} a(t_2) \frac{dr}{\sqrt{1 - kr^2}}$$

#### Propagation of light:

$$\int_{t_1}^{t_2} \frac{c \ dt}{a(t)} = \int_0^{r_2} \frac{dr}{\sqrt{1 - kr^2}}$$



$$d_H(t_1, t_2) = 2a(t_2) \int_{t_1}^{t_2} \frac{c \, dt}{a(t)}$$

> during radiation domination :  $a(t) \propto t^{1/2}$ ,  $R_H = 2 c t$ 

$$d_H(t_1, t_2) = 4 \ c \ t_2^{1/2} [t_2^{1/2} - t_1^{1/2}] \rightarrow 4 \ c \ t_2 = 2 \ R_H(t_2)$$

during matter domination:  $a(t) \propto t^{2/3}$ ,  $R_H = 3/2$  c t

$$d_H(t_1, t_2) = 6 \ c \ t_2^{2/3} [t_2^{1/3} - t_1^{1/3}] \to 6 \ c \ t_2 = 4 \ R_H(t_2)$$

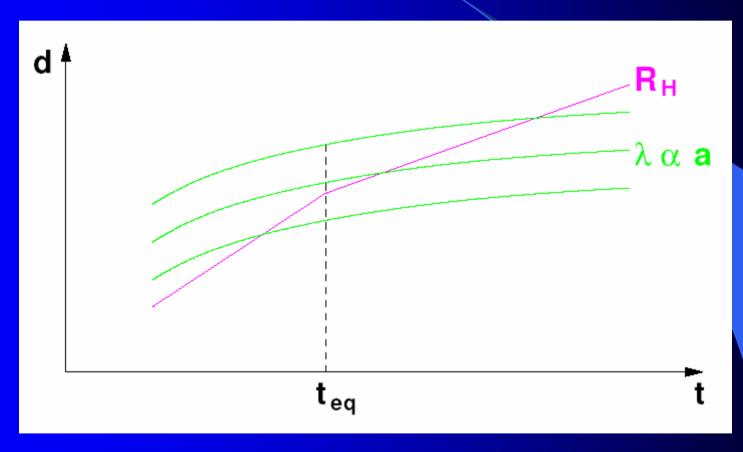
> physical process starting during RD, MD cannot affect  $\lambda(t) \geq R_H(t)$ 

without violating causality...

 $R_H(t) \equiv \text{causal horizon for RD / MD}$ 

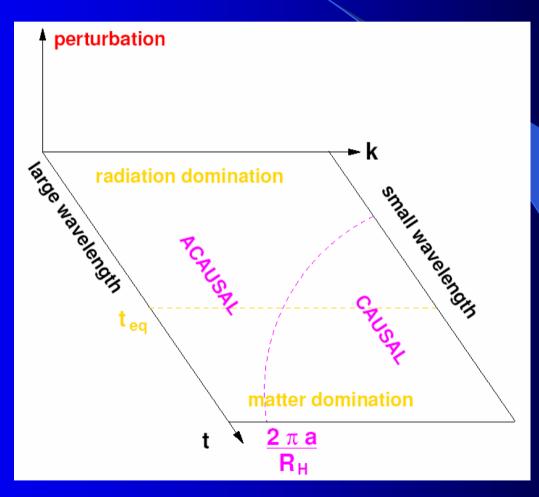
## □ evolution of wavelengths versus R<sub>H</sub>:

$$\lambda(t) = R_H(t) \iff k = 2\pi a(t) / R_H(t)$$



## evolution of wavelengths versus R<sub>H</sub>:

$$\lambda(t) = R_H(t) \iff k = 2\pi a(t) / R_H(t)$$

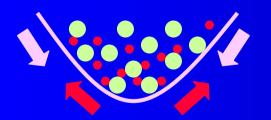


- Planck spectrum  $\Rightarrow \delta_{\gamma}(t, \mathbf{r}) \Leftrightarrow \delta T/T(t, \mathbf{r})$
- 1) before recombination (= decoupling ≈ equality):

```
photons ↔ baryons ↔ grav. pot.

E.M. gravity

tighly coupled fluid
```

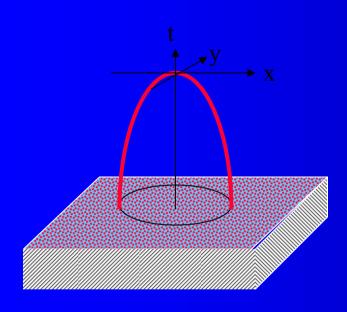


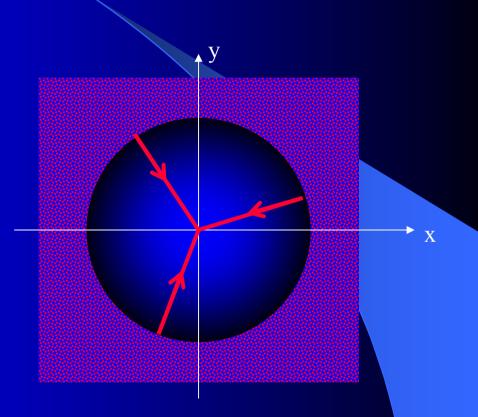
```
gravity « acoustic oscillations » pressure
```

2) after recombination:

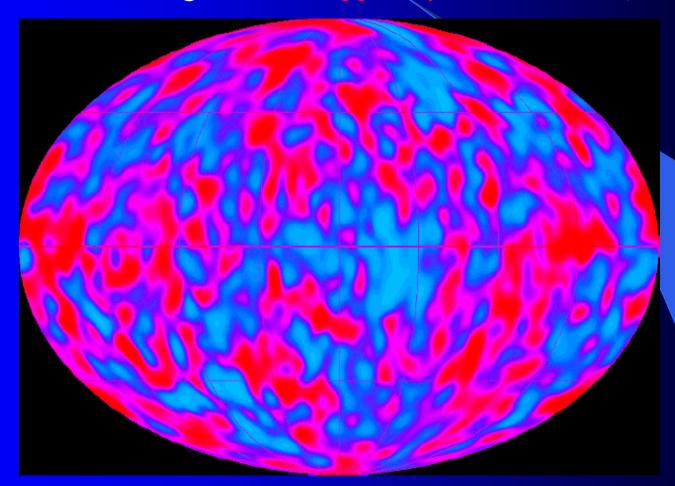
photon decoupled → « free streaming »

last scattering surface :

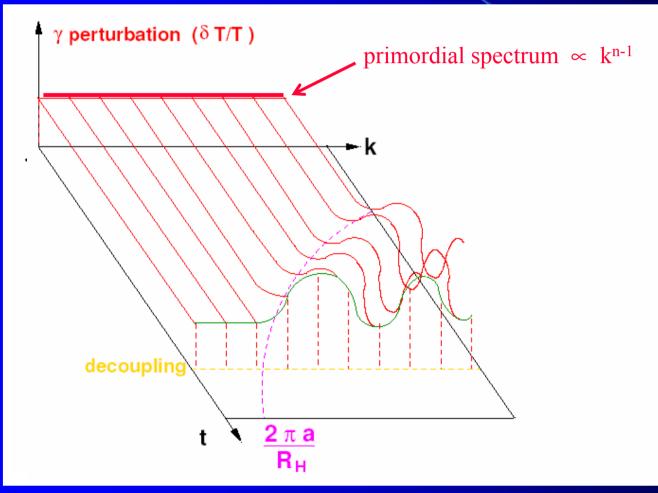




> last scattering surface mapped by COBE DMR (1994):



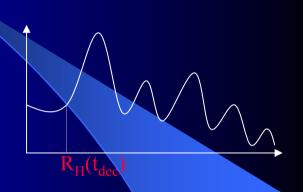
evolution of perturbations :



> spectrum of CMB anisotropies:

observation of CMB

- $^{\ }$   $\delta$ T/T map of last scattering surface
- Fourier spectrum with accoustic peaks



amplitude of the peaks

 $\Rightarrow \Omega_{\rm B}, \Omega_{\rm CDM}, n \dots !!!$ 

 $\Rightarrow$  angle under which  $R_H(t_{dec})$  is seen

angular diameter – redshift relation

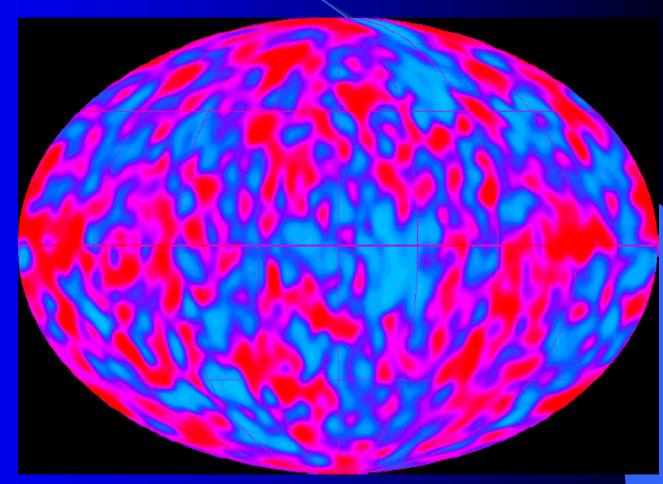
spatial curvature k!!!

### □ PHOTON PERTURBATIONS:

main observations : COBE DMR (1994)

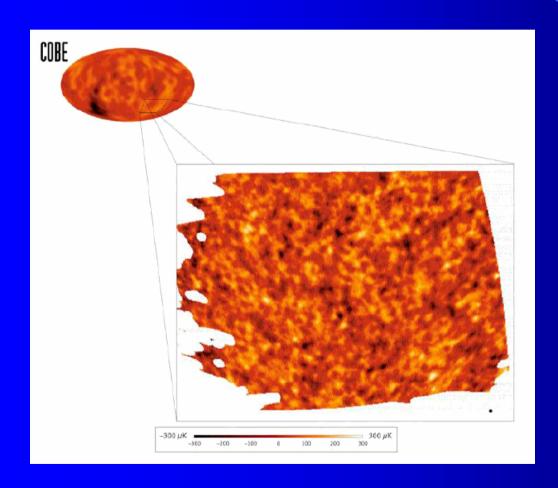
resolution ~ 10°

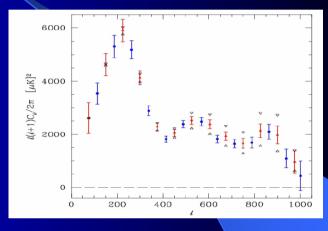
 $\lambda(t) \geq R_{\rm H}(t_{\rm dec})$ 

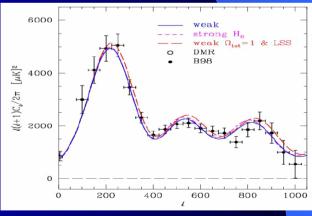


#### □ PHOTON PERTURBATIONS:

main observations: Boomerang (2000)

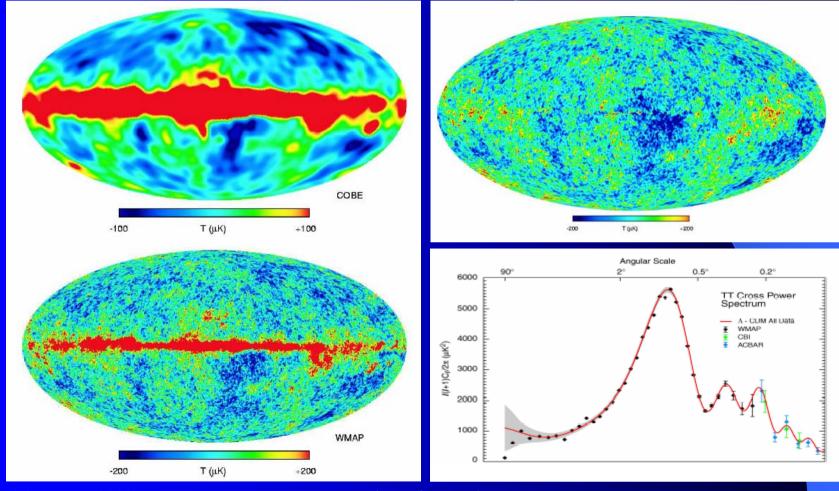






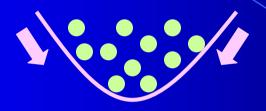
#### □ PHOTON PERTURBATIONS:

main observations: WMAP (February 2003)



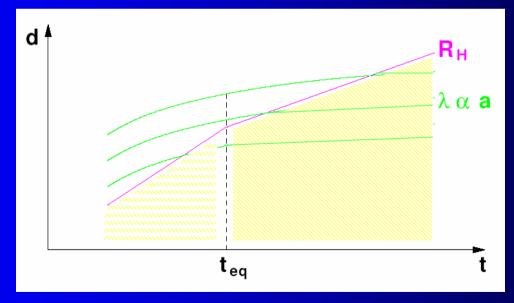
#### □ MATTER PERTURBATIONS :

CDM:



gravitationnal collapse efficient during MD

 $(RD : \Phi_{gray} \text{ follows } \gamma)$ 



#### □ MATTER PERTURBATIONS :

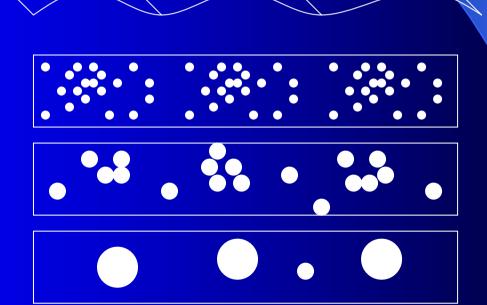
non-linear evolution :

 $\delta_{\text{CDM}}^{k}(t)$ ,  $\delta_{\text{B}}^{k}(t) \sim 1$  first for large k / small  $\lambda$ 

*⇔ hierarchical structure formation :* 

linear theory recovered by smoothing

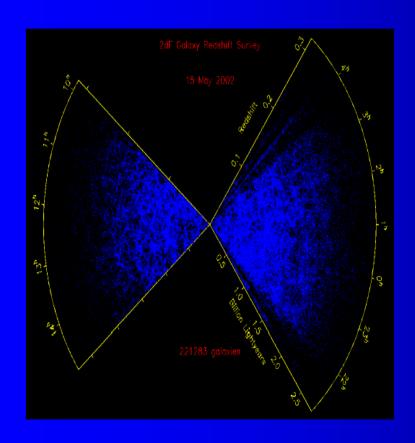
(today: over 30 Mpc)

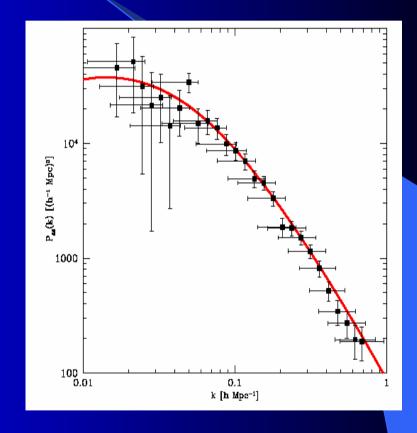


time

### □ MATTER PERTURBATIONS :

bookservations: 2dF redshift survey





## Part II: (3) – cosmological parameters

COSMOGICAL SCENARIO



7 INDEPENDENT

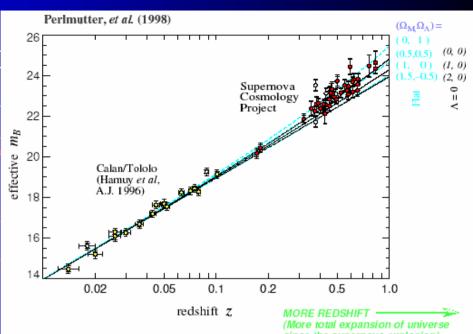
PARAMETERS ... at least !!!

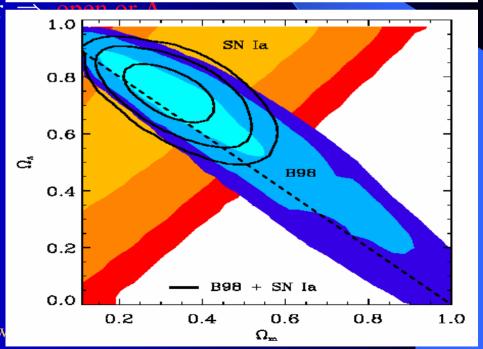
$$\{\Omega_{\rm R}^{\phantom{\dagger}},\Omega_{\rm B}^{\phantom{\dagger}},\Omega_{\rm CDM}^{\phantom{\dagger}},\Omega_{\Lambda}^{\phantom{\dagger}},H_0^{\phantom{\dagger}},A^{\phantom{\dagger}},n\}$$

$$h = \frac{H_0}{100 \text{ km.s}^{-1} \cdot \text{Mpc}^{-1}}$$

#### A short selection of cosmo

- 1) nucleosynthesis
  - $\rho_{B} \Leftrightarrow \rho_{H,\,D,\,He,\,Li}$
- 2) CMB anisotropies
  - position :
  - amplitude :
- 3) age
  - quasars of age ≥ 11 Gyr
- 4) supernovae
  - $\Omega_{\Lambda} \Omega_{M} = 0.5 \pm 0.5$





### A short selection of cosmological tests:

- 1) nucleosynthesis
  - $\rho_{\rm B} \Leftrightarrow \rho_{\rm H,\,D,\,He,\,Li}$

$$\Omega_{\rm B} h^2 = 0.020 \pm 0.002$$
,  $\Omega_{\rm R} \leftarrow \gamma + 3 \text{ v's}$ 

- 2) CMB anisotropies
  - position:

$$\Omega_0 = \Omega_M + \Omega_{\Lambda} = 1.03 \pm 0.05$$

amplitude : 
$$\Omega_{\rm B} h^2 = 0.024 \pm 0.001$$
,  $n = 0.99 \pm 0.04$ 

- age
  - quasars of age  $\geq 11 \text{ Gyr} \Rightarrow \text{open or } \Lambda$
- supernovae

$$\Omega_{\Lambda} - \Omega_{M} = 0.5 \pm 0.5$$

combined: 
$$\Omega_{\rm B} \cong 0.044$$
,  $\Omega_{\rm CDM} \cong 0.23$ ,  $\Omega_{\Lambda} \cong 0.73$ ,  $h \cong 0.71$ 

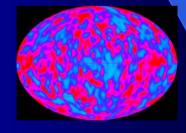
5) large scale structure: perfect agreement

## Part II: (4) – Inflation & Quintessence

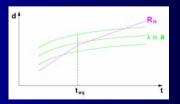
- « early problems » in the Hot Big Bang scenario :
  - > flatness problem:

$$|\Omega_k(t)| = |\rho_0(t)/\rho_c(t) - 1| = \frac{c^2|k|}{a^2H^2} = \frac{c^2|k|}{\dot{a}^2}$$

- $|\Omega_k|$  grows like t (RD) or  $t^{2/3}$  (MD)
- $|\Omega_k| \le 0.1$  at  $t_0 \Rightarrow |\Omega_k| \le 10^{-60}$  at  $t_P$
- horizon problem :
  - causal horizon on CMB maps ~ 1°
     ⇒ 10³ causally disconnected regions

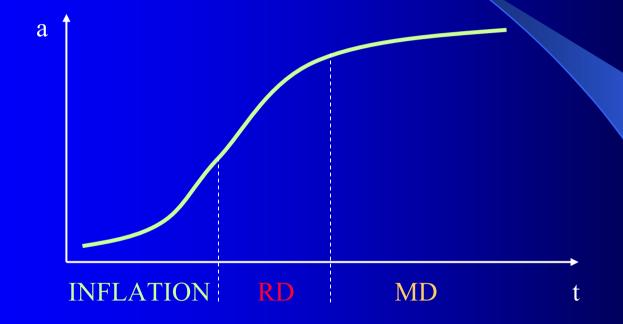


- > origin of fluctuations:
  - initially,  $\lambda \gg R_H \dots$



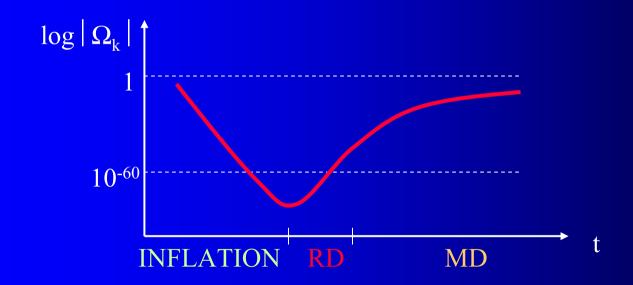
(Guth 79; Starobinsky 79)

defined as an initial accelerated expansion stage:



> solves flatness problem :

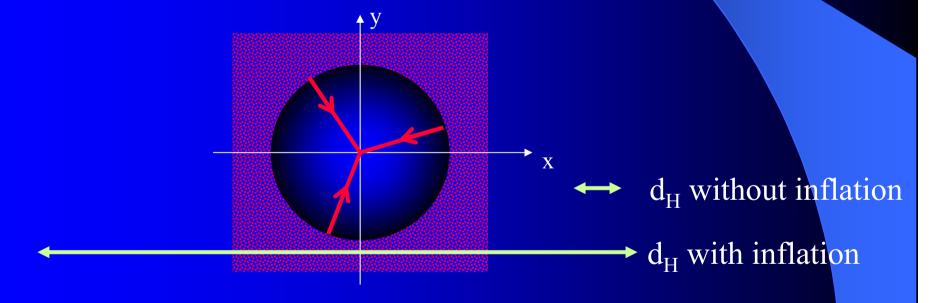
$$|\Omega_k(t)| = |\rho_0(t)/\rho_c(t) - 1| = \frac{c^2|k|}{a^2H^2} = \frac{c^2|k|}{\dot{a}^2}$$



#### > solves horizon problem :

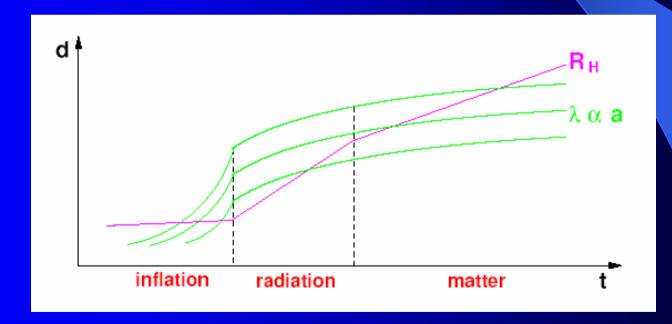
$$d_{H}(t_1,t_2) \cong R_{H}(t_2)$$
 for  $t_1 \ll t_2$  and  $t_1 \in RD, MD$ 

$$d_{H}(t_1,t_2) >>> R_{H}(t_2)$$
 for  $t_1 << t_2$  and  $t_1 \in INFLATION$ 



> solves generation of fluctuations:

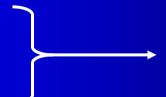
$$\frac{\lambda(t)}{R_H(t)} = \frac{2\pi a(t)}{k} \frac{\dot{a}(t)}{c \; a(t)} = \frac{2\pi \dot{a}(t)}{c \; k}$$



#### > requires:

Friedman law

conservation equation

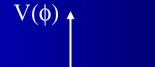


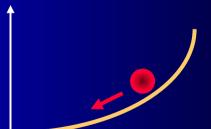
$$a(t) > 0 \iff \rho + 3 p < 0$$

#### > candidates:

- $\Lambda : \rho + 3 p = -2 p < 0$  but inflation forever ...
- slow-rolling scalar field:

$$\begin{cases} \rho = \dot{\phi}^2 / 2 + V(\phi) \\ p = \dot{\phi}^2 / 2 - V(\phi) \end{cases}$$





## □ INFLATION with slow-rolling scalar field:

#### 2 BONUS !!!

> mechanism for generation of cosmological perturbations

quantum fluctuations of scalar field

wavelength
amplification

stochastic background of curvature perturbations



VALIDATED BY OBSERVATIONS

mechanism for generation of first particles : PREHEATING

end of inflation — oscillations of scalar field — particle production

## « late problems » in the Hot Big Bang scenario:

 $\triangleright$  magnitude of  $\Lambda$ :

 $\rho_{\Lambda}^{1/4} \sim 10^{-3} \, \text{eV} \, !!!!$ 

problem for particle physicists :

 $\Leftrightarrow$  killing  $\rho_{\Lambda}^{1/4} \sim \text{MeV / TeV } \dots$ 

problem for cosmologists :

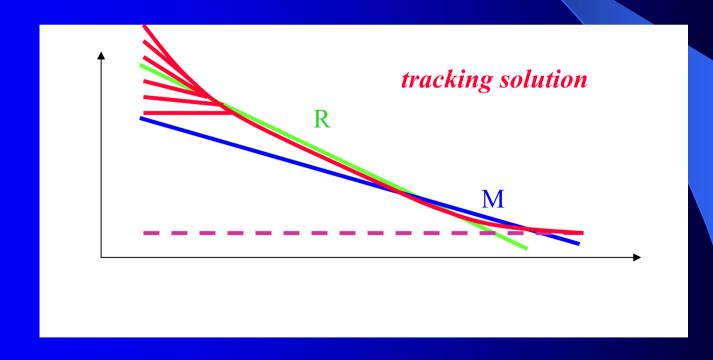
 $\Leftrightarrow$  generating such small number with respect to  $\rho_P^{1/4} \sim 10^{18} \, \text{GeV} \dots$ 

« Cosmic coincidence problem » :

 $\Rightarrow$  why  $\land$  domination today?

- many... unappealing proposals for Dark Energy, e.g.:
  - slow-rolling scalar field : « quintessence »

$$\begin{cases}
\rho = \frac{\dot{\phi}^2}{2} + V(\phi) \\
p = \frac{\dot{\phi}^2}{2} - V(\phi)
\end{cases}$$



## « late problems » in the Hot Big Bang scenario :

 $\triangleright$  magnitude of  $\Lambda$ :

 $\rho_{\Lambda}^{1/4} \sim 10^{-3} \, eV \, !!!!$ 

problem for particle physicists :  $\Leftrightarrow \textit{killing} \ \rho_{\Lambda}^{1/4} \sim \textit{MeV} \ / \ \textit{TeV} \ ...$ 

problem for cosmologists :

solved but  $m \sim 10^{-33} \, eV$  generating such small number with respect to  $\rho_P^{1/4} \sim 10^{18} \, GeV$  ...

« Cosmic coincidence problem » :

 $\Rightarrow$  why  $\land$  domination today?

# CONCLUSION 1

INFLATION is a convincing, predictive theory... but need to relate to particle physics models (GUT, Susy, strings...)

no convincing theory for « DARK ENERGY »

# CONCLUSION 2

- cosmology has remarkable control on:
  - cosmological parameters
  - nucleosynthesis
  - decoupling
  - > structure formation, lensing, etc., etc.
- $\square$  ...but 23 + 73 = 96 % remains MYSTERIOUS !!!

$$dl^{2} = a^{2}(t) \left[ \frac{dr^{2}}{(1-k r^{2})} + r^{2} (d\theta^{2} + \sin^{2}\theta d\phi^{2}) \right]$$

➤ Remark 1: if  $\{k = 0 \text{ AND } a(t) \equiv constant \}$ ,  $\overline{r} = a r \Rightarrow Euclidian space \Rightarrow Newton$ 

$$dI^{2} = a^{2}(t) \left[ \frac{dr^{2}}{(1-k r^{2})} + r^{2} (d\theta^{2} + \sin^{2}\theta d\phi^{2}) \right]$$

#### Remark 2:

- Euclidian :  $dl^2 = dr^2 + r^2 (d\theta^2 + \sin^2\theta d\phi^2) = dx^2 + dy^2 + dz^2$   $\Rightarrow$  does not depend on the choice of origin ...
- FLRW:  $k \neq 0 \Rightarrow dl$  seems to depend on the choice of origin ...

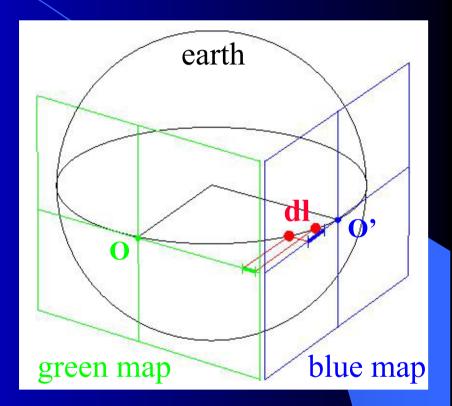


 $\Rightarrow$  do solutions with  $k \neq 0$  violate the assumption of homogeneity ??

NO,  $k \neq 0$  respects homogeneity !!! r changes, but also dr,  $\theta$ , d $\theta$  and d $\phi$ : in fact dl is the same no particular point is priviledged

$$dl^{2} = a^{2}(t) \left[ \frac{dr^{2}}{(1-k r^{2})} + r^{2} (d\theta^{2} + \sin^{2}\theta d\phi^{2}) \right]$$

- Remark 2:analogy with the 2-D mappingof the earth by axial projection:
  - on each map, the scale is a function of r
  - but all points on the sphere are equivalent



... the FLRW model is completely homogeneous !!!

# summary of the situation :

basic principles of G.R.	FLRW solution
space-time is <i>curved</i>	2 types of curvatures : {a(t), k} $dl^2 = a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2\theta \ d\phi^2) \right]$
free-falling objects follow geodesics	????????
curvature caused by / related to matter	????????