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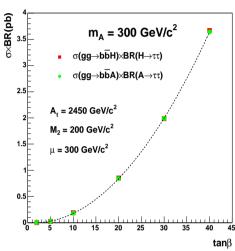
# Estimating the precision of a tan $\beta$ determination with $H_{SUSY} \rightarrow \tau \tau$ in CMS

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## Motivation

- •Tan $\beta$  is one of the most important parameters in MSSM, it enters in all sectors of the theory HDECAY×PPHTT
- Value of  $tan\beta$  can be measured at LHC with
  - Sfermion or neutralino sector at low  $\tan\beta$
  - The Higgs sector using event rates at large  $\tan\beta$
  - ► It may be possible to measure the value of  $tan\beta$  from the width  $\Gamma(H \rightarrow \mu\mu)$

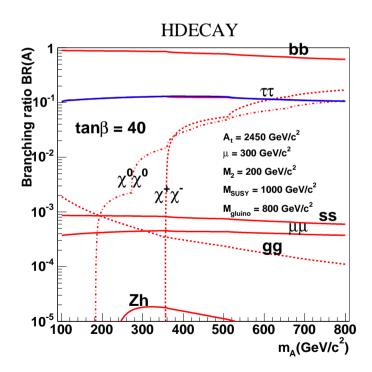


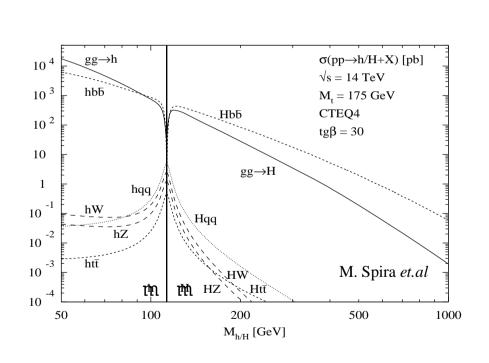
- In MSSM Higgs couplings to down type fermions tanp enhanced
  - > Exploited in  $gg \rightarrow H_{SUSY}/gg \rightarrow bbH_{SUSY}, H_{SUSY} \rightarrow \tau\tau$
  - > Dominant parts of the cross section are proportional to  $\tan^2\beta_{-}eff^{*)}$
  - Sumpling ratio BR( $H_{SUSY} \rightarrow \tau \tau$ ) ~ constant as a function of tan $\beta$
- $\Rightarrow$  Uncertainty of the tan $\beta$  measurement only half of the uncertainty of the rate measurement

\*)  $\tan\beta_{eff}$  will later be called as  $\tan\beta$ 

• At large  $\tan\beta$  ( $\tan\beta > 10$ ) the Higgs production is predominantly in association with two b quarks (90%), loop mediated gluon fusion has less importance

 $_{g}$  D





 $_{g}$   $\mathcal{U}$ 

h, H, A

g g

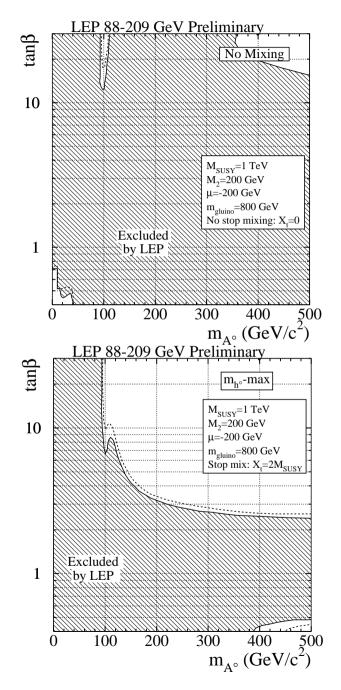
Selected SUSY scenario: LEP benchmark scenario max  $m_h$ 

SUSY parameters chosen

$$M_{SUSY} = 1 \text{ TeV}$$
$$A_t = \sqrt{6} M_{SUSY}$$
$$\mu = 300 \text{ GeV}$$
$$M_2 = 200 \text{ GeV}$$
$$m_{gluino} = 800 \text{ GeV}$$

Stop mixing has only small effect on the observability of the  $H_{SUSY} \rightarrow \tau\tau$  channels

At large  $m_{A,} H_{SUSY} \rightarrow \tau \tau$  sensitive to gaugino sector, especially if tan $\beta$  not very high. Affects mostly the fully hadronic final state due lower tan $\beta$  reach.



### Simulation methods

Event generation: pythia6, toprex+tauola for tt and tW backgr Signal cross section and branching ratios from PPHTT/HDECAY (M. Spira *et. al.*)

Studies with parameterized fast simulation, partly with full simulation

Full simulations for detector sensitive issues: tau trigger, tau identification, mass reconstruction, b tagging, impact parameter resolution

Full simulation results used as efficiencies or they are parameterized

### Final states and main backgrounds

H $\rightarrow \tau \tau \rightarrow e\mu + X$ Branching ratio BR( $\tau \tau \rightarrow e\mu$ )~6.3%

 $H \rightarrow \tau \tau \rightarrow l + jet + X$ Branching ratio BR( $\tau \tau \rightarrow lj$ )~45.6%

Main backgrounds:					
Ζ,γ*→ττ	(all channels)				
Z,γ*→ll	(11)				
tt	(all channels)				
tW	(all channels)				
bb	(eµ,ll,lj)				
W+jet	(lj,jj)				
QCD	(jj)				

 $H \rightarrow \tau \tau \rightarrow ll + X$ Branching ratio BR( $\tau \tau \rightarrow ll$ )~12.5%

 $H \rightarrow \tau \tau \rightarrow jet+jet+X$ Branching ratio BR( $\tau \tau \rightarrow jj$ )~41.5%

Background suppression after HLT: lepton isolation
τ-jet identification
τ impact parameter
b-tagging
jet veto
positive neutrino energy

# Trigger

### $H{\rightarrow}\tau\tau{\rightarrow}e\mu{+}X$

#### Single muon trigger, eff = 0.85

trigger threshold effect  $0.9 \times$ muon reco efficiency $0.97 \times$ calorimetric isolation 0.97

#### $H {\rightarrow} \tau \tau {\rightarrow} ll {+} X$

# Single muon trigger+di-electron trigger, eff = 0.82

di-electron trigger: (eff/electron)

trigger threshold effect 0.95× Level-1 electron efficiency 0.872× Level 2.5 electron efficiency 0.946

#### $H \rightarrow \tau \tau \rightarrow l + jet + X$

#### Single muon trigger + $e+\tau$ jet trigger, eff = 0.73

e+tjet trigger: electron trigger threshold effect  $0.95 \times$ Level 1 electron efficiency  $0.872 \times$ HLT electron efficiency  $0.77 \times$ tau trigger threshold effect 0.95

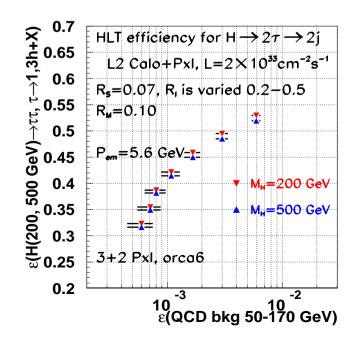
Kinematic cuts selected above the trigger thresholds, 20 GeV for leptons and 45 GeV for  $\tau$  jets

### $H \rightarrow \tau \tau \rightarrow jet+jet+X$

Di- $\tau$  trigger, eff = 0.38

L1 efficiency  $0.9 \times HLT$  efficiency 0.42

QCD rejection vs HLT efficiency



# $\tau$ jet identification

- Level-1 and HLT  $\tau$  trigger thresholds (95% points) 1 $\tau$  jet:86 GeV, 2 $\tau$  jet:59 GeV, e- $\tau$  jet:45 GeV
- Offline selections:
  - leading track  $p_T > 40 \text{ GeV}$
  - 1 or 3 tracks in cone 0.04 around the leading track
  - no other tracks with  $p_T > 1$  GeV in isolation cone 0.4
- Very efficient against QCD, bb, W+jet backgrounds:

suppression of QCD di-jet events ~1000 per jet 9% efficiency per H $\rightarrow \tau\tau \rightarrow 2$ jet event

including trigger+offline  $\tau$  selections (m<sub>H</sub> = 500 GeV)

# b tagging in gg→bbH/bbA

- B jets in bbH/bbA soft and distributed over wide rapidity range
- $\Rightarrow$  B tagging efficiency low even with perfect ip measurement
- Tagging method:

counting significant tracks with  $p_T > 1$  GeV,  $\sigma_{ip} > 2$  (signed).

• B tagging efficiency ~35% for genuine b's with ~1% mistagging rate

Single or double b tagging?

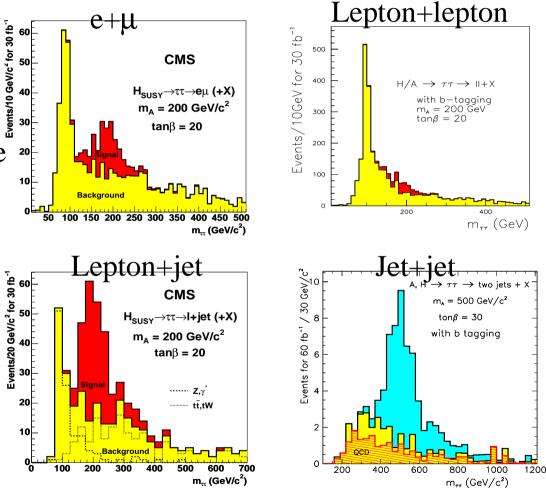
 $gg \rightarrow bbH/bbA, H/A \rightarrow \tau\tau \rightarrow lepton+jet+X,$ number of selected events for 30fb<sup>-1</sup>, all selection cuts  $m_A = 200 \text{ GeV}, \tan\beta = 20 \text{ N}_S \text{ N}_B \text{ signif } \text{sqrt}(\text{N}_S + \text{N}_B)/\text{N}_S$ 1b-tagging+jet veto 157 70 18.8 9.6% 2b-tagging 9 44 1.3 80.9%

• Best results with one tagged b per event

- allows central jet veto to be used to suppress efficiently tt backgr.

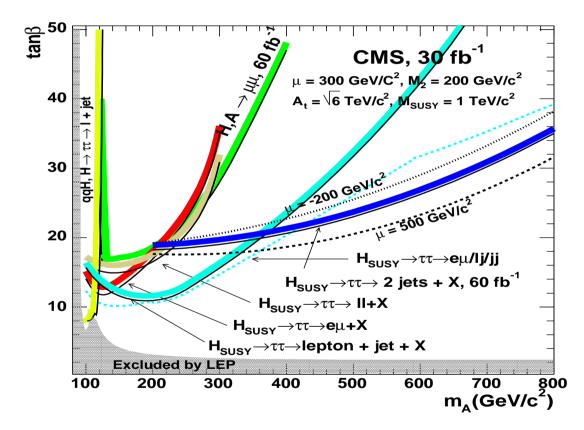
## $H_{SUSY} \rightarrow \tau \tau$ , reconstructed Higgs boson mass

- Neutrinos in the final state
- collinear approximaton used  $\frac{1}{2}$  ...
- Mass reconstruction possible if  $\tau$ 's not back-to-back:  $\Delta \phi(\tau_1, \tau_2) < 175^\circ$  required
- Positive neutrino energy required
- Suppresses tt background: all neutrinos not emitted along the tau



## $5\sigma$ discovery contours

30 fb<sup>-1</sup> at low luminosity ( $2 \times 10^{33}$  cm<sup>-2</sup>s<sup>-1</sup>) max m<sub>h</sub> SUSY scenario



 $H_{SUSY} \rightarrow \tau \tau \rightarrow X$  most promising channel for discovering heavy neutral MSSM Higgs boson at large tan $\beta$ 

Uncertainty of tanβ measurement At large tanβ  $\sigma$ -tan<sup>2</sup>β×X, subleading tanβ dependence small, can be absorbed into tanβ<sub>eff</sub> N<sub>s</sub> = tan<sup>2</sup>β×X×L×ε<sub>sel</sub>

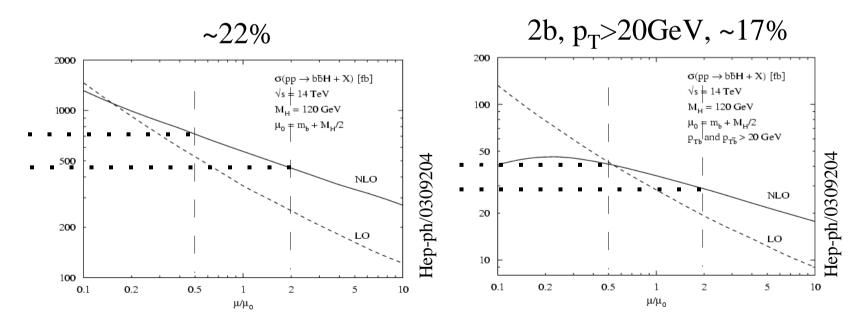
> $tan\beta = tan\beta_0 \pm \Delta stat \pm \Delta syst$ Max error:  $\Delta tan\beta/tan\beta = \frac{1}{2}(\Delta N_S/N_S + \Delta L/L + \Delta X/X)$   $= \frac{1}{2}(sqrt(N_S+N_B)/N_S + \Delta L/L + \Delta X/X)$

Luminosity error assumed  $\Delta L/L \sim 5\%$ Theoretical error  $\Delta\sigma/\sigma \sim 20\%$ ,  $\Delta BR/BR \sim 3\%$ 

Uncertainties of the background and signal selection efficiency, and the accuracy of the SUSY parameter measurement not yet taken into account. (Uncertainty of the selection (s+b) efficiency expected ~5%)

### Theoretical uncertainty

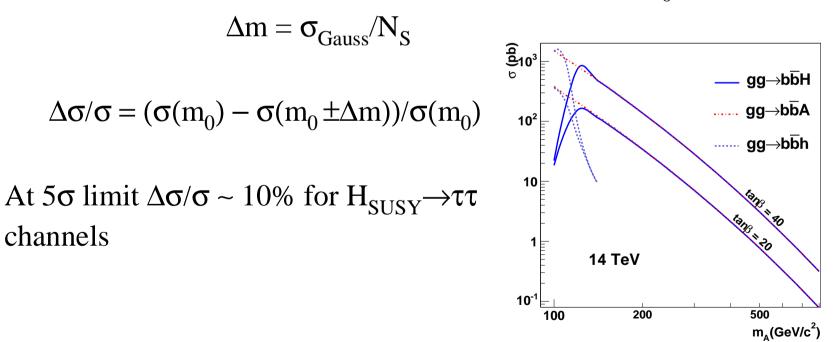
- Uncertainty of the NLO cross section for gg→bbH/A 20-30% Dittmaier, Kramer, Spira, hep-ph/0309204
- The error depends on the transverse momentum range of the spectator quarks:
  - if both associated b's have  $p_T > 20$  GeV in bbH/A, the error reduces to 10-15%
- We use 1b-tagging and 20% theoretical error



### Uncertainty of the mass measurement

Production cross section depends on Higgs mass, which is measured with some accuracy. This induces errors on the cross section.

Mass resolution from Gaussian fit to the signal,  $m = m_0 \pm \Delta m$ 



### Uncertainty of the SUSY parameters measurement

- In addition to mA and tan $\beta$ , Higgs sector sensitive to  $A_t, \mu, M_2, M_{SUSY}$
- These parameters need to be measured using e.g. global fit to SUSY parameters. The error of this measurement gives error to the Higgs production rate ( $\sigma \times BR$ )
- The error is still unknown \*)
- To get a feeling how large the errors might be, let's take e.g. 20% uncertainty of the SUSY parameters

	$\sigma(gg \to b\bar{b}H) \times BR(H \to \tau\tau)$						
	$\tan\beta = 20$	0	$\tan\beta = 40$				
	$m_{\rm A} = 200 \; {\rm GeV}/c^2$	$m_{\rm A} = 500  {\rm GeV}/c^2$	$m_{\rm A} = 200 \; {\rm GeV}/c^2$	$m_{\rm A} = 500 \text{ GeV}/c^2$			
$\mu\pm 20\%$	(+1.4 -1.1)%	(+13.7 -25.9)%	(+2.0 -2.3)%	(+6.9 -11.3)%			
$M_2\pm 20\%$	(+0.055 -0.26)%	(+9.5 -18.2)%	(+0.013 -0.066)%	(+2.9 -6.3)%			
$A_t\pm 20\%$	(+1.3 -0.022)%	(+0.82 +0.19)%	(+0.65 -0.26)%	(+0.27 +0.53)%			
$M_{SUSY}\pm20\%$	(+1.6 -1.0)%	(+1.4 -0.78)%	(-1.8 +2.7)%	(-1.5 +2.7)%			

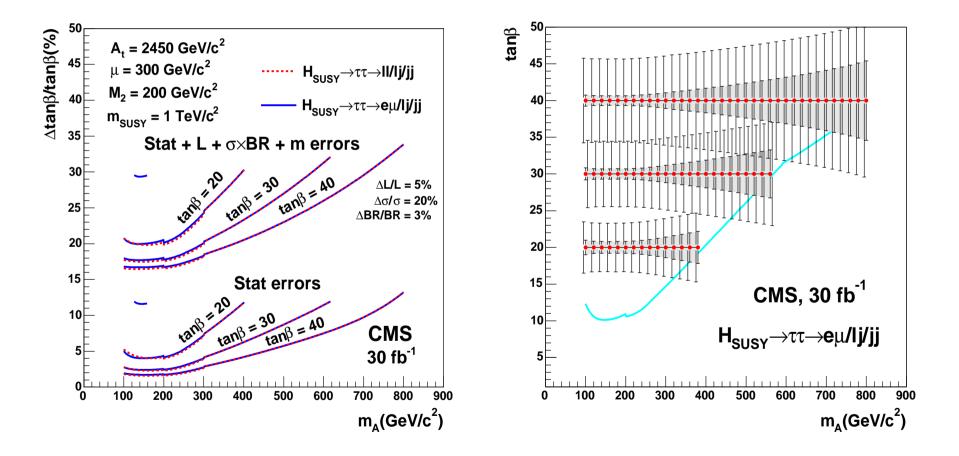
<sup>\*)</sup> some initial studies will be published in LHC/LC document on global fit

## Uncertainty of tan $\beta$ measurement, 30fb<sup>-1</sup>

	$m_{\rm A} = 200  {\rm GeV}/c^2$		$m_{\rm A} = 200  {\rm GeV}/c^2$		$m_{\rm A} = 500  {\rm GeV}/c^2$		$m_{\rm A} = 500 {\rm GeV}/c^2$	
$30  {\rm fb}^{-1}$	$\tan\beta = 20$		$\tan\beta = 30$		$\tan\beta = 30$		$\tan\beta = 40$	
	$\Delta$ stat	$\Delta \sigma(\Delta m)$	$\Delta$ stat	$\Delta \sigma(\Delta m)$	$\Delta$ stat	$\Delta \sigma (\Delta m)$	$\Delta$ stat	$\Delta \sigma(\Delta m)$
$H/A \rightarrow \tau \tau \rightarrow e\mu$	8.95%	4.82%	4.85%	3.27%	-	-	-	-
$H/A \to \tau \tau \to \ell \ell$	7.96%	3.50%	4.08%	2.37%	-	-	-	-
$H/A \rightarrow \tau \tau \rightarrow \ell j$	4.81%	2.46%	2.84%	1.65%	-	-	8.40%	4.82%
$H/A \rightarrow \tau \tau \rightarrow jj$	13.7%	4.73%	8.25%	3.21%	12.4%	5.82%	8.45%	4.44%
Combined	4.05%	1.99%	2.35%	1.34%	9.09%	4.28%	5.96%	3.26%
eµ+ℓj+jj	$\Delta tan \beta / tan \beta$		$\Delta  aneta/ aneta$		$\Delta tan \beta / tan \beta$		$\Delta tan \beta / tan \beta$	
	20.1%		17.7%		27.4%		23.3%	
Combined	3.94%	1.85%	2.24%	1.25%	9.09%	4.28%	5.96%	3.26%
l ℓℓ+ℓj+jj	$\Delta tan \beta / tan \beta$		$\Delta  aneta/ aneta$		$\Delta  aneta/ aneta$		$\Delta tan \beta / tan \beta$	
	19.9%		17.5%		27.4%		23.3%	

## Uncertainty of $tan\beta$ measurement

Small error bars (gray): stat error only Large error bars: total error



# Conclusions

- The precision of a tan $\beta$  determination is estimated using  $H_{SUSY} \rightarrow \tau \tau$  in CMS with 30fb<sup>-1</sup>
- The uncertainty includes
  - Statistical error
  - Error of the mass measurement
  - Theoretical error (cross section and branching ratio)
  - Luminosity error
- What is <u>not</u> included:
  - Uncertainty of the SUSY parameter measurement  $(\mu, M_2, ...)$
  - Uncertainty of the signal selection and background determination
- With signal significance >5 $\sigma$ , the study gives better than 35% accuracy for the tan $\beta$  determination at the benchmark point considered
- Errors are dominated by theoretical uncertainty