

# SM and SUSY Higgs bosons: Phenomenological aspects for the LHC

Abdelhak DJOUADI (Montpellier/Paris)

## I. A brief introduction

## II. The Higgs boson in the SM

- The production processes in the SM
- Test of the EWSB mechanism at the LHC

## III. The Higgs bosons in the MSSM

- The decoupling regime
- The intermediate regime
- The intense coupling regime
- The SUSY regime
- The extended regime

## IV. Conclusions

# I. Brief Introduction

## The Higgs sector in the SM:

One doublet of complex scalar fields  $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$  with  $Y_\phi = +1$

3 degrees of freedom for  $W_L^\pm, Z_L$  and one is left corresponding to  $H$ .

$M_H$  is the only free SM parameter, the information that we have is

$$114.4 \text{ GeV} \lesssim M_H \lesssim \mathcal{O}(1 \text{ TeV}) \quad \text{and} \quad M_H \lesssim 250 \text{ GeV}$$

But, hierarchy (and also unification) problem in the SM  $\Rightarrow$  SUSY.

## In the minimal SUSY version, the MSSM:

Two doublets of complex scalar fields with opposite  $Y$ ,  $\Phi_1$  and  $\Phi_2$

$\Rightarrow$  5 Higgs particles in the spectrum:  $h, H, A$  and  $H^\pm$ .

Besides the four masses, 2 parameters are needed:  $\tan\beta = \frac{v_2}{v_1}$  and  $\alpha$ .

However, there are strong constraints from SUSY:

- Only two free parameters, taken to be  $M_A$  and  $\tan\beta$  in general.
- Hierarchy in spectrum:  $M_h < M_A < M_H, M_{H^\pm} > M_W, M_h < M_Z$ .
- In the decoupling limit,  $M_A \gg M_Z$  (in practice  $M_A \gtrsim 300 \text{ GeV}$ ):
  - $H, A, H^\pm$  are degenerate  $M_H \sim M_{H^\pm} \sim M_A$  and decouple.
  - $h$  has maximal mass,  $M_h = M_Z |\cos(2\beta)|$  and SM couplings.
  - SM and MSSM Higgs sectors similar but there is a light Higgs.

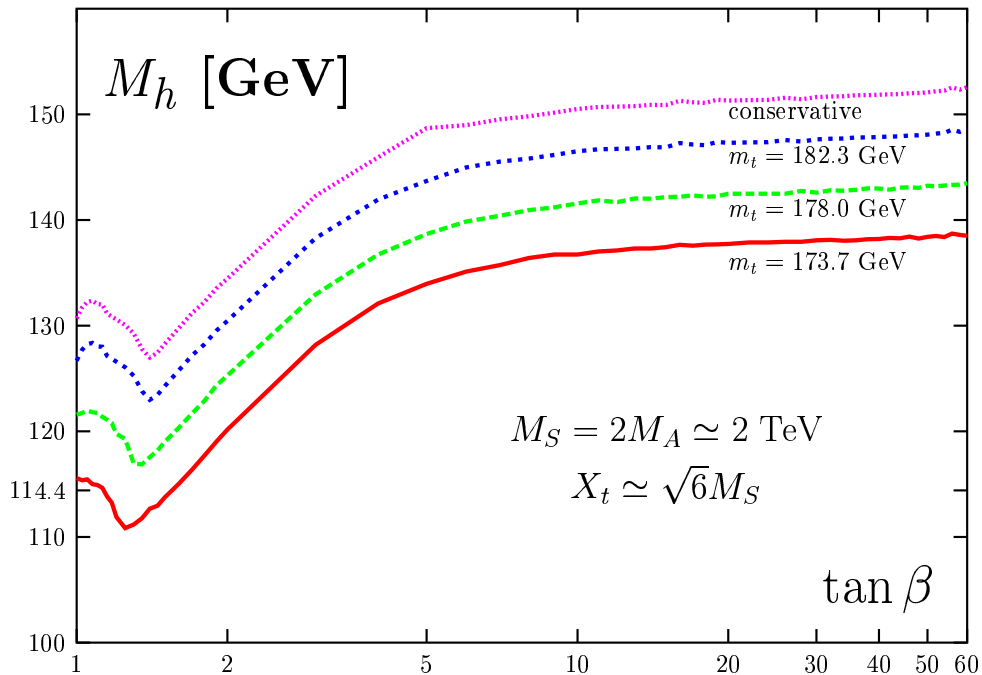
Relations broken by radiative corrections for large  $m_t$ , leading RC:

$$M_h^{\max} \sim M_Z |\cos 2\beta| \left[ 1 - \mathcal{O} \left( \frac{M_Z^2}{M_A^2} \right) \right] + \frac{3 m_t^4}{2 \pi^2 v^2} \left( \ln \frac{M_S^2}{m_t^2} + \frac{X_t^2}{2 M_S^2} - \frac{X_t^4}{12 M_S^4} \right)$$

A large activity for the calculation of the RC in the last 15 years.

Most recent result (June 2004):

- Full one-loop correction with all SUSY loops and  $q^2$  dependence.
- All two-loop RC involving large couplings:  $\alpha_S$  and  $\alpha_t, \alpha_b, \alpha_\tau$ .
- Estimation of all theoretical and experimental uncertainties.
- New Tevatron value of the top quark mass  $M_t = 178.0 \pm 4.3$  GeV.

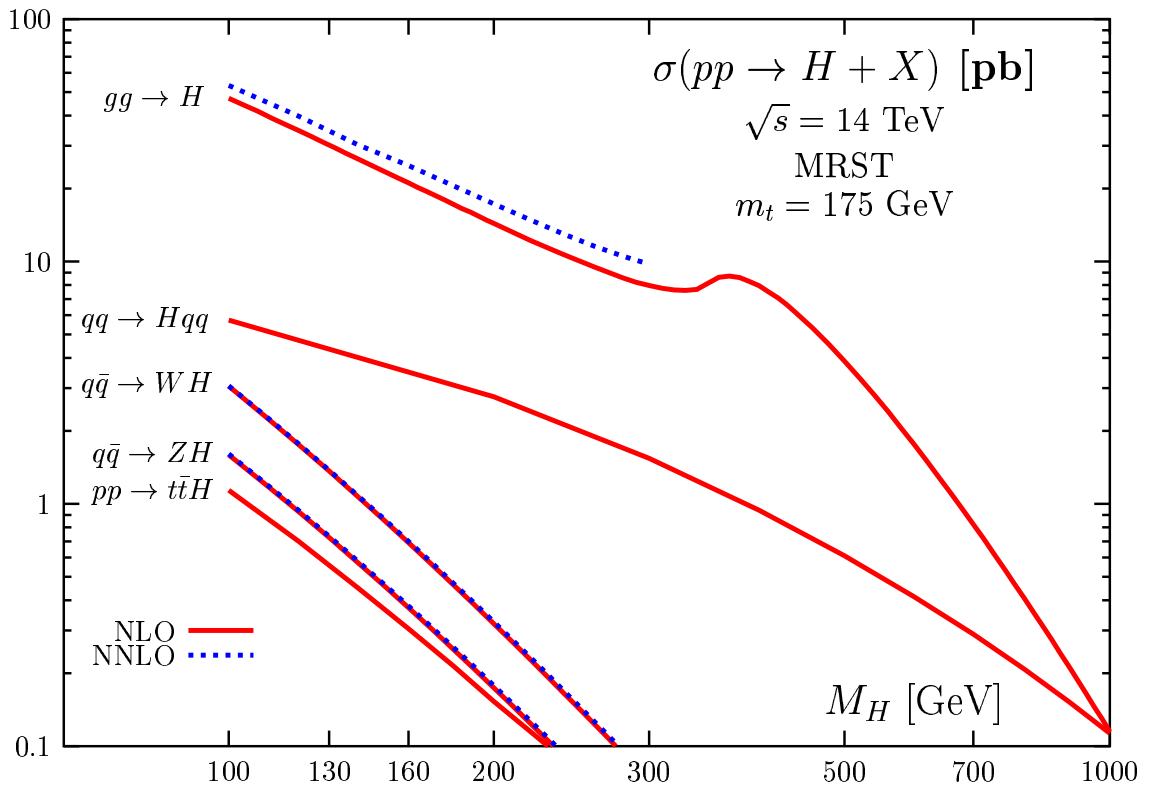
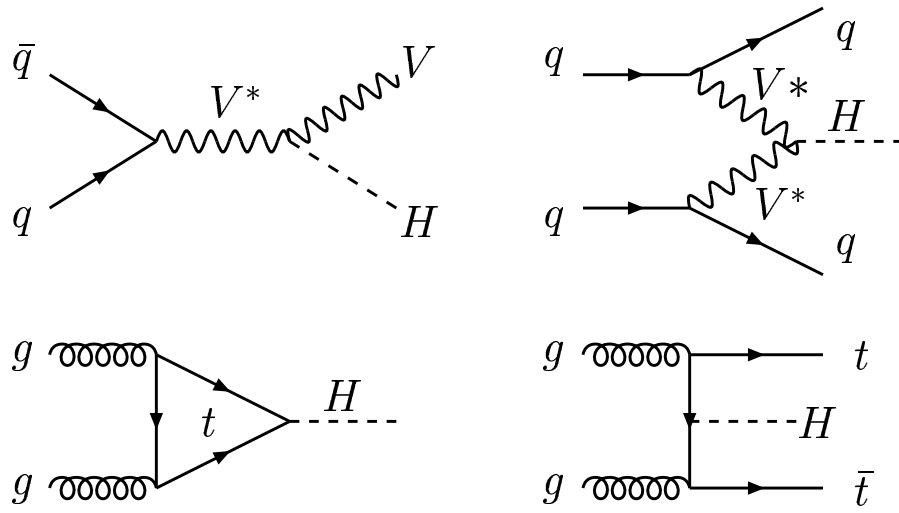


Allanach, AD, Kneur, Porod and Slavich

Consequences :  $M_h^{\max} \sim 150$  GeV ,  $\tan \beta|^{\min} \sim 1$

## II. The Higgs boson in the SM

Production processes and cross sections at the LHC in the SM:



(a)  $gg \rightarrow H$ : dominant production mechanism  $\Rightarrow$  discovery channel:

–  $M_H \lesssim 130$  GeV:  $H \rightarrow \gamma\gamma$

–  $M_H \gtrsim 130$  GeV:  $H \rightarrow WW^{(*)}/ZZ^{(*)} \Rightarrow l^+l^-\nu\nu, l\nu jj/4l^\pm, l^+l^-\nu\nu$

**Theoretical status:**

–  $K_{\text{NLO}} \sim 1.7$  known for arbitrary  $m_t$  and including  $b$ -loop (MSSM).

[Spira, AD, Graudenz, Zerwas \(1995\).](#)

–  $K_{\text{NNLO}} \sim 2$  recently derived in the infinite  $m_t$  limit (also NNLL).

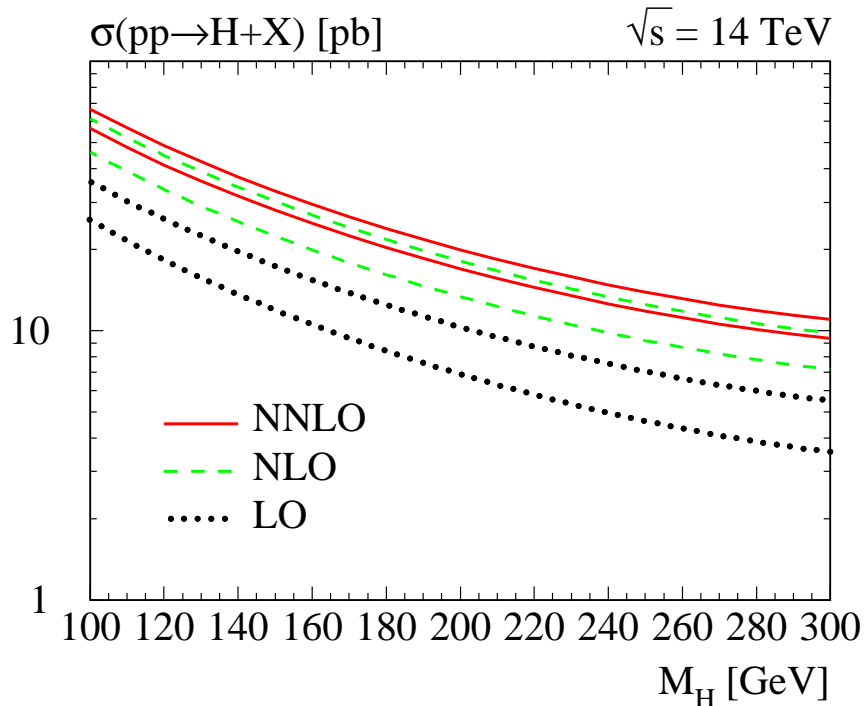
NLO: [AD+Spira+Zerwas; Dawson](#)

NNLO: [Harlander+Kilgore; Anastasiou+Melnikov; Ravindran+Smith+van Neerven](#)

NNLL: [Catani, de Florian, Grazzini, Nason](#)

– Large correction but nice convergence of the perturbative series.

– Scale dependence (measure of higher orders) strongly reduced.



[Harlander, Kilgore](#)

**(b)  $qq \rightarrow qqH$ : most promising channel for couplings measurements**

- Large enough rates,  $\sigma(qq \rightarrow qqH) \sim \mathcal{O}(\text{pb})$  for  $M_H \lesssim 300$  GeV.
- Forward jets (tags) and central Higgs decays (mini-jet veto).
- Rather small backgrounds,  $S/B \sim \mathcal{O}(1)$ , allowing measurements.

Zeppenfeld, Plehn, Rainwater, ...

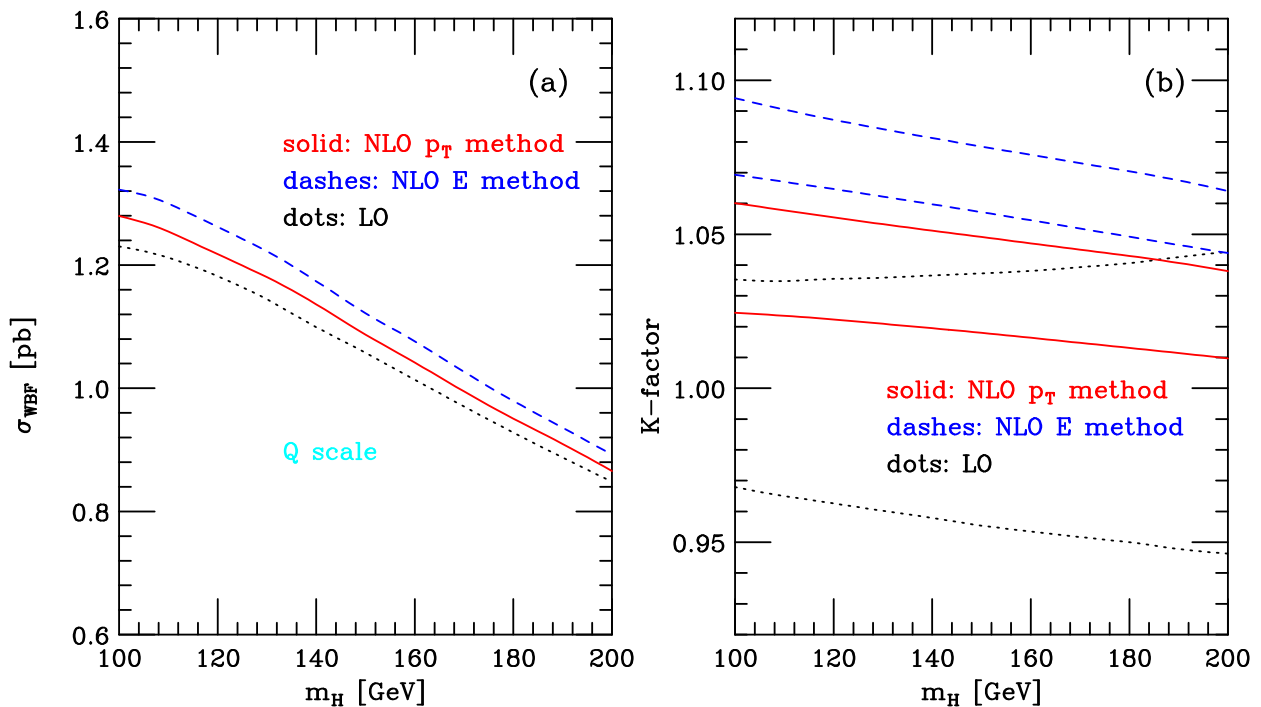
- $H \rightarrow \tau^+\tau^-, WW^{(*)}, \gamma\gamma$  feasible.
- $H \rightarrow ZZ, b\bar{b}, \mu^+\mu^-$  at high  $\mathcal{L}$  and more studies?

**Theoretical status:**

- $K_{\text{NLO}} \sim 1.1$  and very small scale dependence.

Han+Valencia+Willenbrock; Figy+Oleari+Zeppenfeld

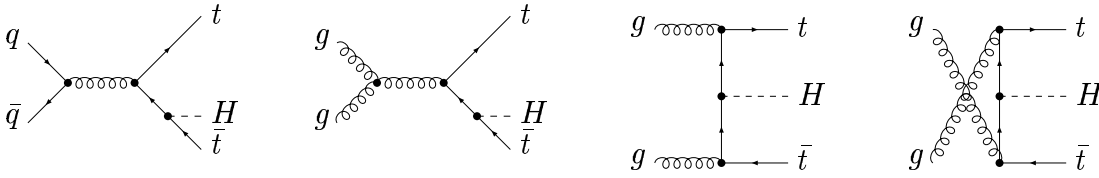
- $K_{\text{NLO}}$  now available for  $p_T$  and  $\eta$  distributions.
- Still questions about backgrounds and central jet veto.



Figy, Oleari and Zeppenfeld

**(c)  $qq/gg \rightarrow ttH$ :**

Most complicated Higgs production channel at hadron colliders:



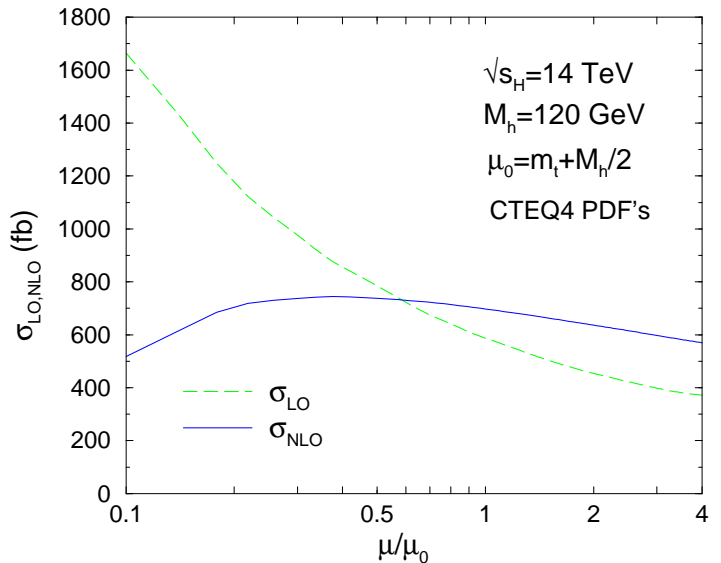
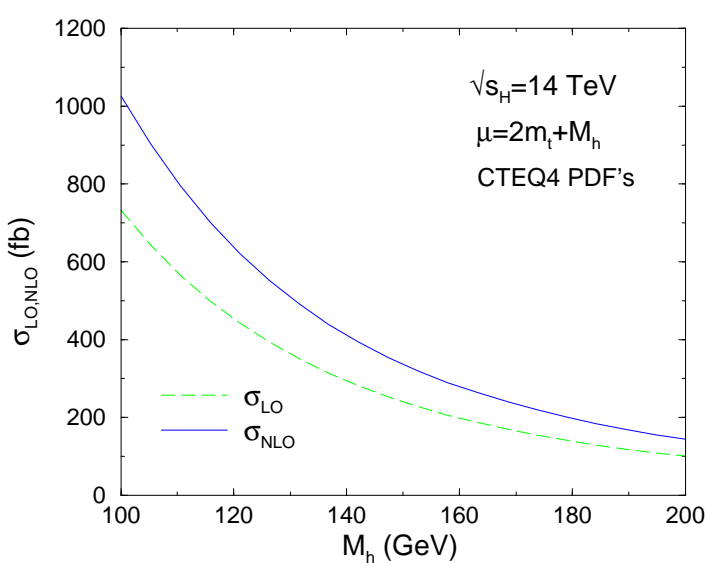
Usefull for measuring the  $ttH$  coupling in  $pp \rightarrow \ell\nu\gamma\gamma$  or  $\ell\nu b\bar{b}$ .

- The **LO** cross section available since 1984 (Kunszt).
- Full **NLO** calculation available only since two years:

Dawson, Orr, Reina, Wackerroth  
 Beenakker, Dittmaier, Kramer, Plumper, Spira, Zerwas

**Main results:**

- **K-factors**  $\sim 1.2$  at the LHC [and 0.8 at the Tevatron!]
- Scale dependence drastically reduced  $\sim 10 - 20\%$ .
- **K-factors** also available in the  $b\bar{b} + \text{Higgs}$  case (MSSM).



(d)  $q\bar{q} \rightarrow HW, HZ$ : the Tevatron channels

This production channel is the most important one at the Tevatron:

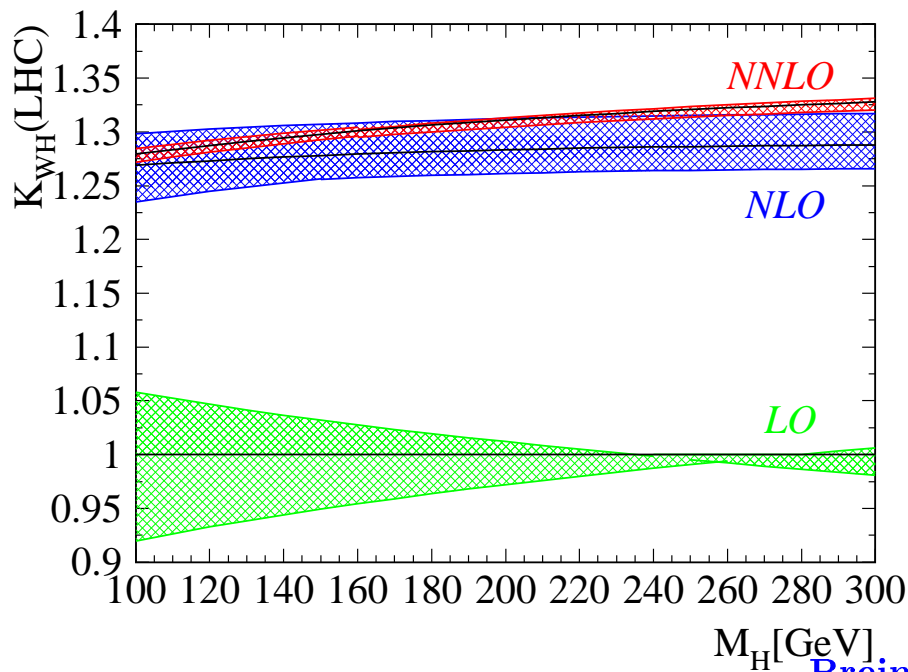
- $M_H \lesssim 130$  GeV:  $H \rightarrow b\bar{b} \Rightarrow \ell\nu b\bar{b}, \nu\bar{\nu} b\bar{b}, \ell^+\ell^- b\bar{b}$ .
- $M_H \gtrsim 130$  GeV:  $H \rightarrow WW^* \Rightarrow \ell^\pm\ell^\pm jj$ .

Up-to-now, plays only a marginal role at LHC (small rates etc...).

However, theoretically very clean channel:

- K-factor  $\sim 1.35$  made recently available at NNLO  
Han+Willenbrock (NLO), Brein+AD+Harlander (NNLO)
- Electroweak corrections are also available ( $-5\%$ ).  
Ciccolini, Dittmaer, Krämer (EW)
- Scale dependence extremely small  $\sim 3 - 5\%$ .
- Small PDF uncertainties (normalisation to DY?).

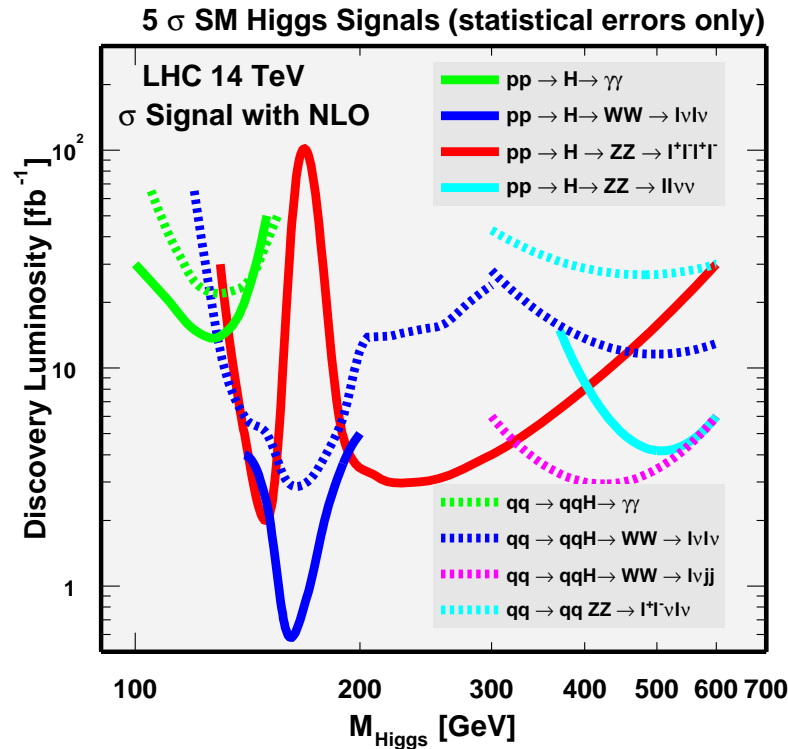
Dittmar, Pauss, Zürcher





In the SM: Higgs detection is in principle guaranteed.

- Many redundant production and detection channels
- $> 5\sigma$  discovery once all channels are combined for low  $\mathcal{L}$ .



The next step is to perform detailed tests of the EWSB mechanism:

- Measure the Higgs mass and total width  
Easy if  $H \rightarrow \gamma\gamma, ZZ \rightarrow 4 \ll$
- Measure the coupling to fermions and gauge bosons (WBF+...).  
Started already; see M. Dührssen for instance
- Measure the Higgs self-coupling and reconstruct  $V_{\text{scal}}$  (!).  
Seems to be hopeless; see Baur, Plehn and Rainwater
- Determine the Higgs spin and parity quantum numbers (!).  
A very challenging issue.

# III. The Higgs bosons of the MSSM

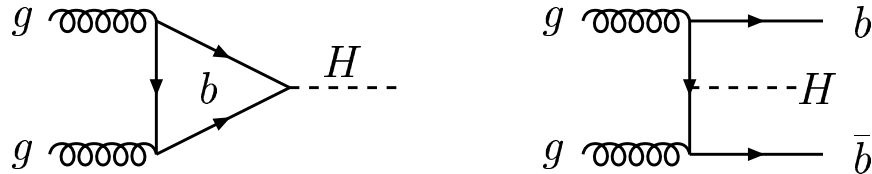
## Production processes and cross sections in the MSSM:

### For the neutral Higgs bosons:

Same channels as in SM apply but rates enhanced/suppressed by different Higgs couplings to fermions and gauge bosons.

**In particular, for  $\tan\beta \gg 1$  and  $M_A \gtrsim 150$  GeV:**

- $b$ -loops dominant in the  $gg \rightarrow H/A$  processes ( $g_{\Phi bb} \propto \tan\beta$ ).
- $b\bar{b}$  final states dominant in  $pp \rightarrow Q\bar{Q} + H/A$  ( $\sigma \propto \tan^2\beta$ ).
- No VV fusion or  $VH$  production for  $A$  (CP) and  $H$  (suppressed).
- $\text{BR}(H/A \rightarrow \gamma\gamma, ZZ)$  very small and  $\text{BR}(H/A \rightarrow b\bar{b}, \tau^+\tau^-) \sim 90\%, 10\%$ .
- For  $h$ : same (or slightly smaller) cross sections as in SM.
- Since  $M_h \lesssim 130\text{--}150$  GeV, rely on  $h \rightarrow \gamma\gamma$  and  $h \rightarrow WW^*$ .



### For the charged Higgs boson:

the production processes, depending on the  $H^\pm$  mass, are:

- $M_{H^\pm} \lesssim m_t$ :  $gg/q\bar{q} \rightarrow t\bar{t} \rightarrow H^\pm + X \rightarrow \tau^\pm\nu + X$ .
- $M_{H^\pm} \gtrsim m_t$ :  $gb \rightarrow H^-t$  and  $gg/q\bar{q} \rightarrow H^-t\bar{b} \rightarrow \tau\nu$  or  $tb$ .

See D.P. Roy

**(a) The decoupling regime:  $M_A \gg M_Z$  and  $\tan \beta \gg 1$**

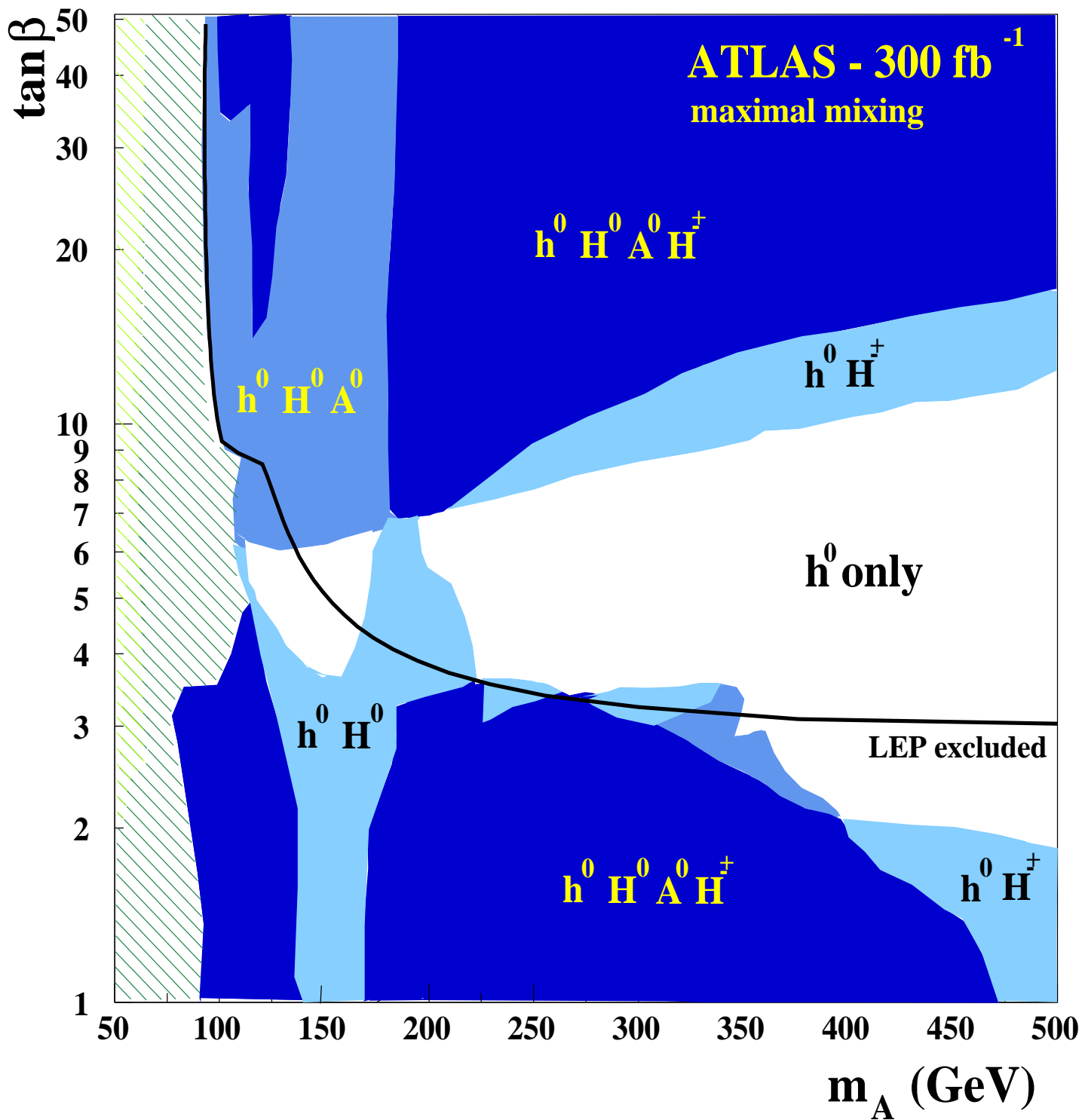
In constrained MSSMs, most of the time in the decoupling regime.

- The  $h$  boson is SM-like with  $115 \text{ GeV} \lesssim M_h \lesssim 150 \text{ GeV}$ 
  - $gg \rightarrow h$  with  $h \rightarrow \gamma\gamma, WW^* \rightarrow \ell\nu\nu, ZZ^* \rightarrow 4\ell$
  - $qq \rightarrow qqh$  with  $h \rightarrow \tau^+\tau^-, \gamma\gamma, WW^* \rightarrow \ell\nu\nu$
  - $q\bar{q} \rightarrow Wh$  and  $q\bar{q}/gg \rightarrow t\bar{t}h$  with  $h \rightarrow \gamma\gamma, b\bar{b}$ .
- $H, A, H^\pm$  are heavy and their search is separated. For  $\tan \beta \gg 1$ :
  - $gb \rightarrow b + H/A$  or  $q\bar{q}/gg \rightarrow b\bar{b} + H/A$  with  $H/A \rightarrow \tau^+\tau^- (\mu^+\mu^-)$ .
  - $gb \rightarrow H^-t$  and  $gg/q\bar{q} \rightarrow H^-t\bar{b}$  with  $H^- \rightarrow \tau\nu$  or  $b\bar{t}$ .

**(b) The intermediate regime:  $M_A \lesssim 500 \text{ GeV}$  and  $\tan \beta \lesssim 5-10$**

- The  $h$  boson is not yet SM-like: more difficult than in SM
  - Couplings to  $WW, ZZ, tt$  suppressed and to  $bb, \tau\tau$  enhanced.
  - All production channels have cross sections smaller than in SM.
  - Decay modes into  $WW, ZZ, \gamma\gamma$  suppressed.
- Situation more tricky for  $H, A, H^\pm$  bosons:
  - Production cross sections not large ( $t$  down but  $b$  not yet up).  
 $\sigma(gg \rightarrow H/A)$  larger than  $\sigma(gg \rightarrow b\bar{b} + H/A)$  or  $\sigma(gg \rightarrow t\bar{t} + H/A)$ .
  - But some special (clean) decays are now revived ( $\tan \beta \sim 1-3$ ):  
 $H \rightarrow hh, A \rightarrow Zh, H^\pm \rightarrow Wh$  and maybe  $H/A \rightarrow t\bar{t}$ .

**Some work to be redone?**



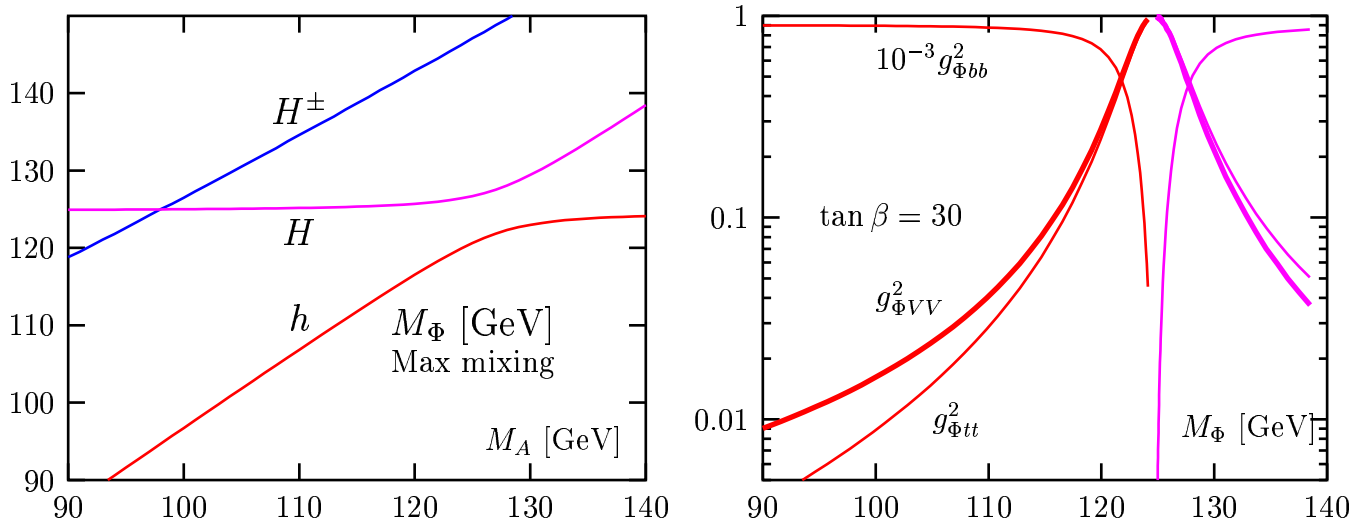
(c) **The intense coupling regime:  $M_A \lesssim 120\text{--}150$  GeV and  $\tan\beta \gg 1$**

**New and/or interesting features:**

- $h, H, A$  light,  $M_h \sim M_A \sim M_H$ , all processes are in principle accessible.
- production cross sections rather large for many channels.
- complete coverage of the MSSM is theoretically possible.

**But many problems and challenges too:**

- The  $h, A, H$  masses are too close to each other.
- The widths are large and must be taken into account.
- The interesting BRs suppressed v.s. SM for both  $h$  and  $H$ .
- The signal for one process is a background for the other.
- The  $bb$  decays are hopeless and the  $\tau^+\tau^-$  resolution too bad.

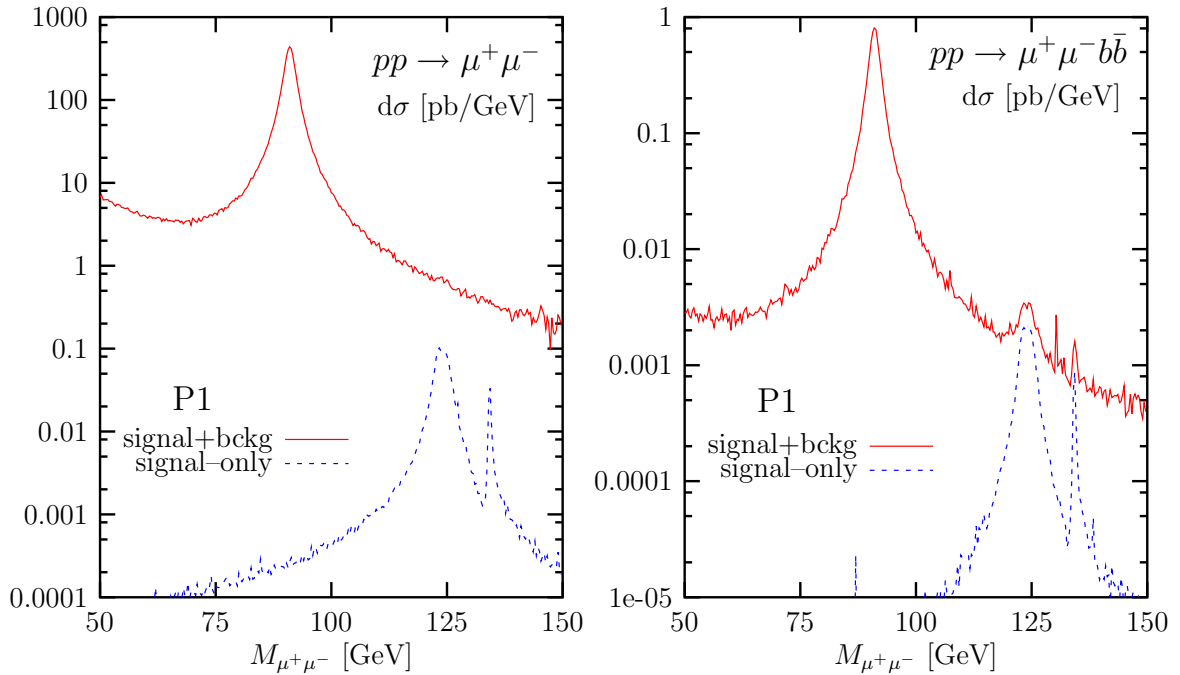


**Go for the  $h, H, A \rightarrow \mu^+\mu^+$  channels:**

**Boos, AD, Nikitenko**

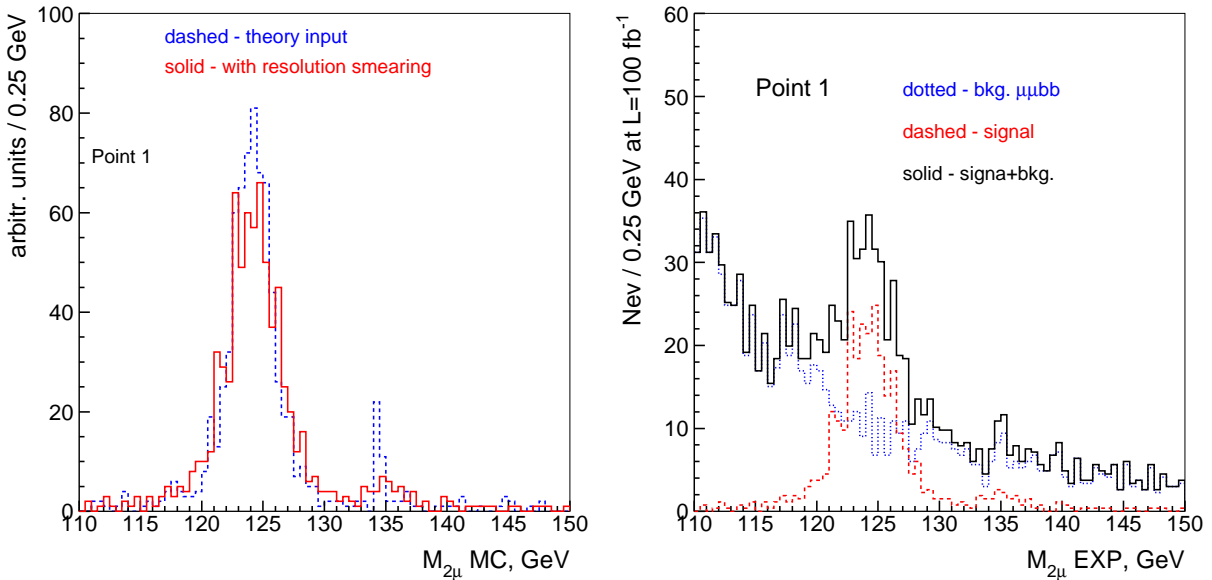
**P1:  $M_A = 125$  GeV,  $\tan\beta = 30 \Rightarrow M_h = 123.3$  GeV,  $M_H = 134.3$  GeV.**

**S and B for:**  $pp(\rightarrow \Phi) \rightarrow \mu^+\mu^-$  and  $pp(\rightarrow \Phi b\bar{b}) \rightarrow \mu^+\mu^-b\bar{b}$ :



–  $gg \rightarrow \Phi \rightarrow \mu^+\mu^-$  seems to be hopeless  $S/B \ll 1$

–  $pp \rightarrow \Phi b\bar{b} \rightarrow \mu^+\mu^-b\bar{b}$  seems **OK: simulation**



– Still difficult to disentangle  $h$  and  $A$ .

– For some other points: situation even more complicated.

**(d) The SUSY regime: not too heavy SUSY particles**

- **SUSY loops** might play an important role in production/decay.

ex: light  $\tilde{t}$  with strong couplings to  $h$  (large stop mixing)

the loops might make  $\sigma(gg \rightarrow h) \times \text{BR}(h \rightarrow \gamma\gamma)$  very small.

- In this case, a complementary  $h$  **SUSY production** channel is:

$pp \rightarrow gg/q\bar{q} \rightarrow h\tilde{t}_1\tilde{t}_1$  with cross section as large as  $\sigma(tth)$  possible.

AD, Kneur, Moutaka

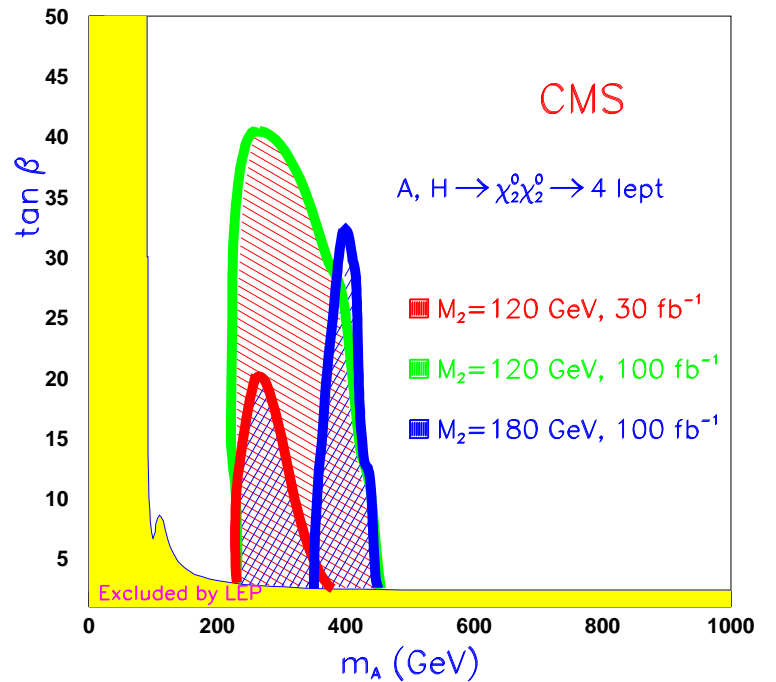
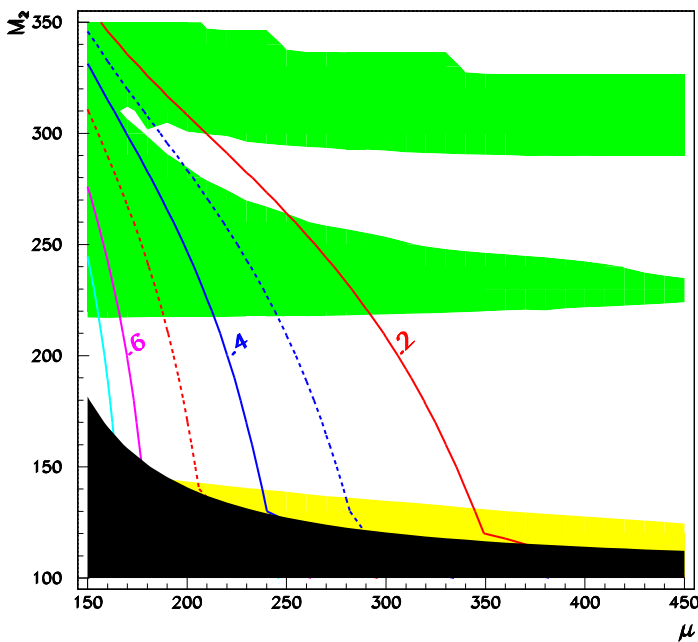
- **SUSY decays**, if allowed, might alter the search strategies:

–  $h \rightarrow \chi_1^0\chi_1^0, \tilde{\nu}\tilde{\nu}$  are still possible in non universal models...

– Decays of  $A, H, H^\pm$  into  $\chi_i^\pm, \chi_i^0$  are possible but can be useful...

Boudjema et al.

Abdullin, Denegri, Moortgat



- **Cascade decays of SUSY particles:**

$\tilde{q}, \tilde{g}$  are copiously produced at the LHC (strongly interacting).

Possible large branching ratios for cascades into MSSM Higgs bosons?

- In the past, the “**little cascade**” has been analyzed for the  $h$  boson:

$$pp \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q}, \tilde{q}\tilde{q}^*, \tilde{q}\tilde{g} \rightarrow \chi_1^\pm, \chi_2^0 + X \rightarrow \chi_1^0 + h(H^\pm, H, A) + X$$

- One can also look at the “**big cascade**” processes for all Higgses:

$$pp \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q}, \tilde{q}\tilde{q}^*, \tilde{q}\tilde{g} \rightarrow \chi_2^\pm, \chi_{3,4}^0 + X \rightarrow \chi_1^\pm, \chi_{1,2}^0 + h, H, A, H^\pm + X$$

- Also possible in some areas of the MSSM parameter space: direct decays of  $\tilde{t}, \tilde{b}$  and for a light  $H^\pm$ , new decays through top quarks.

Processes particularly suited for the intense coupling regime!

Opening the Pandora box: quick analysis in some scenarios:

**Sc 1:**  $m_{\tilde{q}} > m_{\tilde{t}_1} > m_{\tilde{g}}, m_{\chi_{1,2}^\pm, \chi_2^0} - m_{\chi_1^0} < M_\Phi \rightarrow$  **big cascade.**

**Sc 2:**  $m_{\tilde{q}} > m_{\tilde{t}_1} > m_{\tilde{g}}, m_{\chi_{1,2}^\pm, \chi_2^0} - m_{\chi_1^0} > M_\Phi \rightarrow$  **small+big cascade.**

**Sc 3:**  $m_{\tilde{q}} > m_{\tilde{t}_1} > m_{\tilde{g}},$  small  $\mu$  (higgsino):  $\tilde{q} \rightarrow$  **big cascade.**

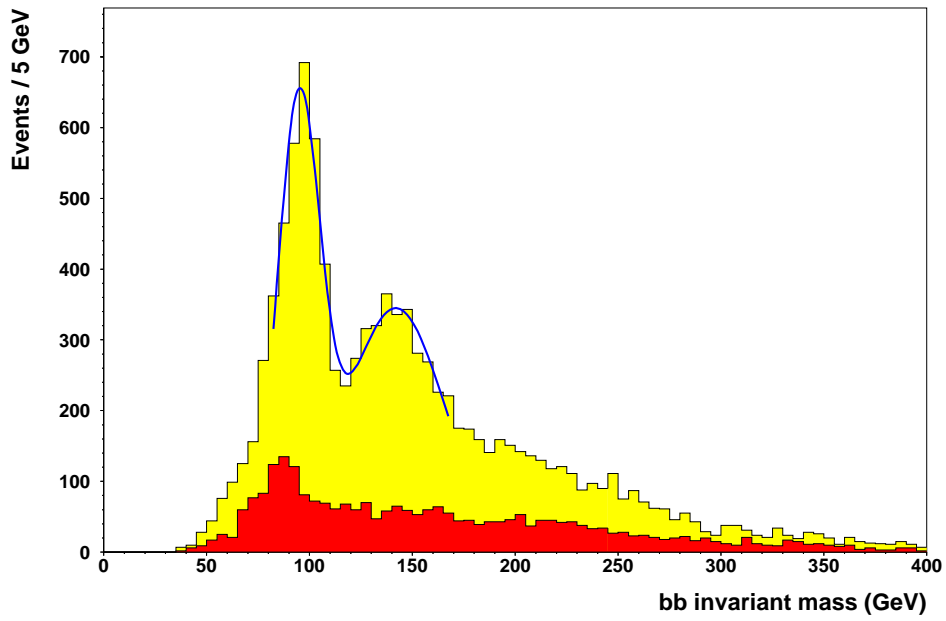
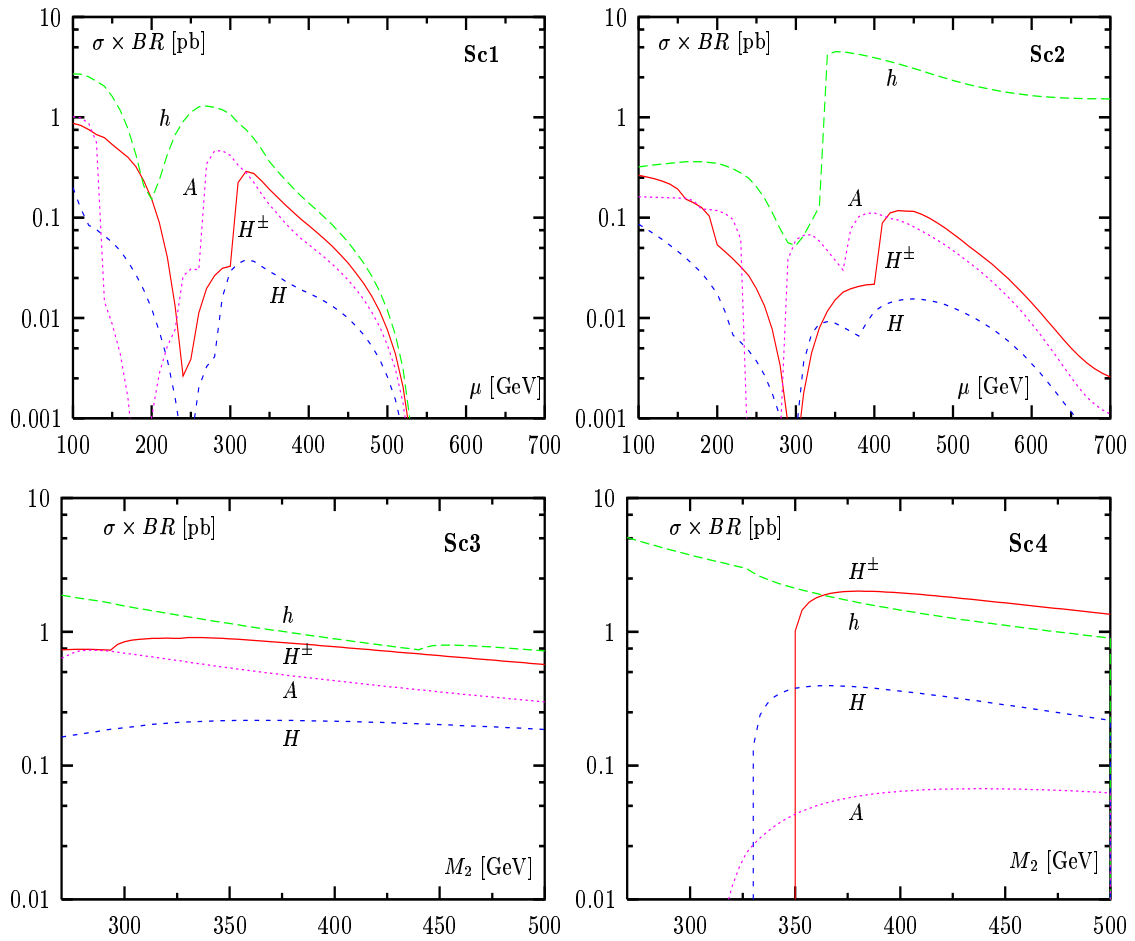
**Sc 4:**  $m_{\tilde{q}} > m_{\tilde{t}_1} > m_{\tilde{g}},$  small  $\mu$  (gaugino):  $\tilde{q} \rightarrow$  **small cascade.**

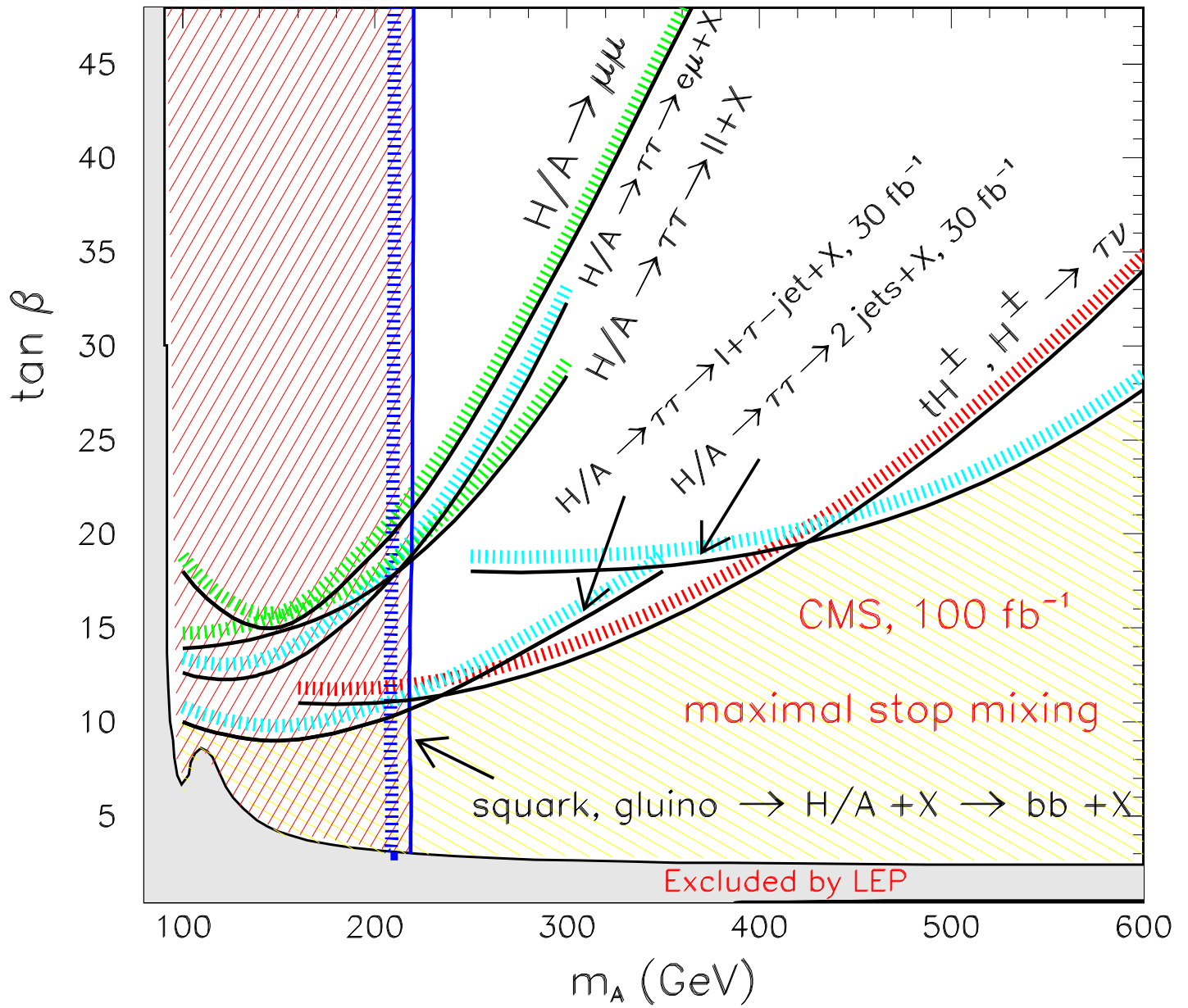
**and it looks very promising....**

More studies are needed in the near future.

See talk by N. Marinelli







**(e) The extended regime:**

The orthodox MSSM (2HDM+CP) might not be the entire truth.

Some SUSY extensions are possible and must be studied:

- CP violation in the Higgs and/or SUSY sectors.
- More complicated Higgs representations (doublet/triplets, etc..).

**Example of possible change: the NMSSM with an additional singlet:**

$$W = \lambda S H_1 H_2 + \kappa/3 \cdot S^3 + \dots$$

- Gives a natural reason for  $\mu \equiv \lambda \langle S \rangle$  to be at the low scale.
- Has less fine tuning than the MSSM....

**Higgs sector extended to include:  $h_1, h_2, h_3, a_1, a_2$  and  $h^\pm$ .**

- Couplings of  $h_1, h_2, h_3$  to  $W/Z$  suppressed compared to  $h$ .
- Constraints from LEP2 do not hold. Ex:  $M_{a_1} \lesssim 50$  GeV possible.
- Not yet a “no-loose theorem” for Higgs discovery at the LHC.

**Ellwanger, Hugonie, Gunion, Moretti**

**A special scenario:  $h_1, h_2$  have SM-like couplings and  $a_1$  is light**

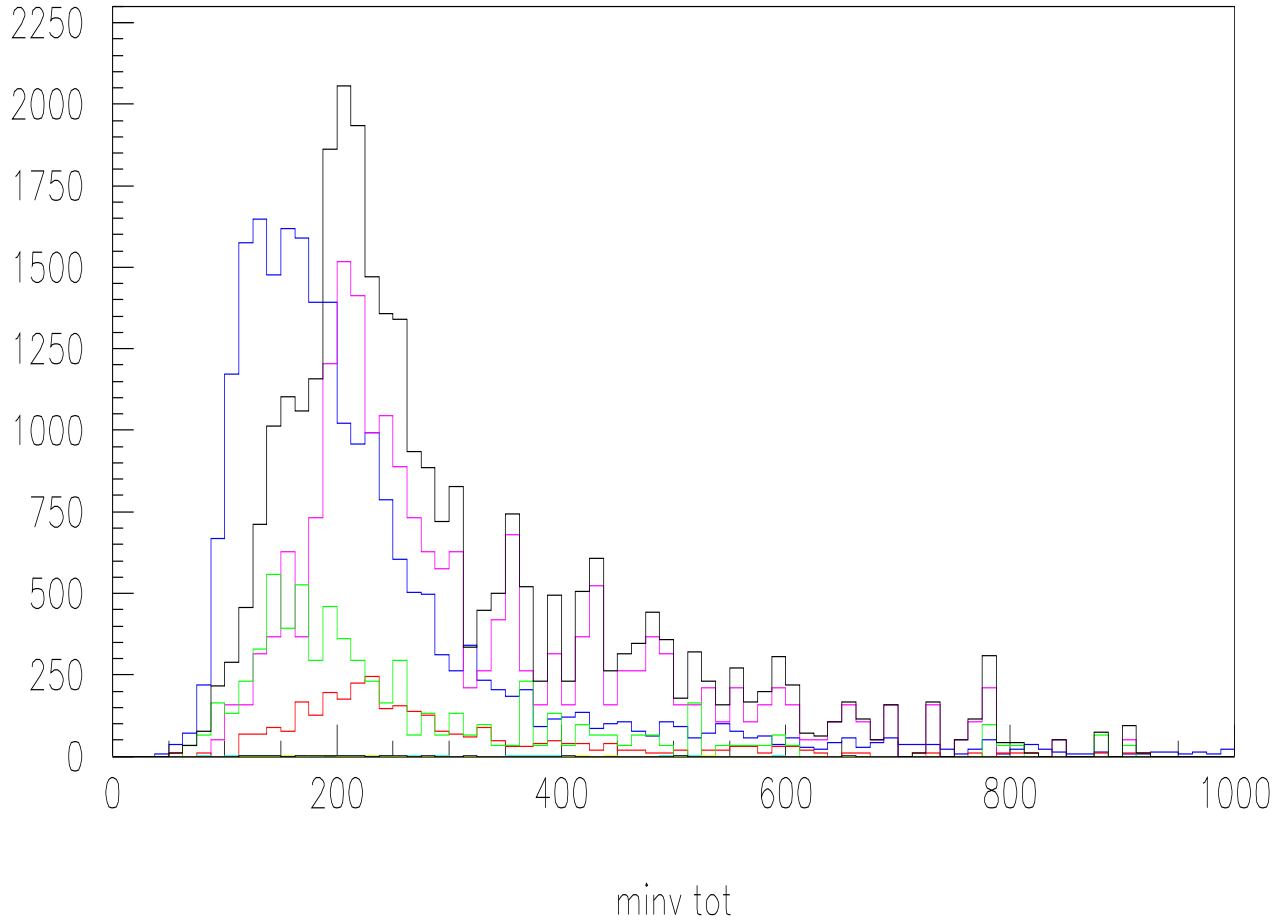
$$pp \rightarrow h_i + X \text{ with } h_i \rightarrow a_1 a_1 \text{ with } a_1 \rightarrow b\bar{b}, \tau^+ \tau^-$$

It has been tried in WBF for  $qq \rightarrow W^* W^* / Z^* Z^* \rightarrow qqh_i \rightarrow b\bar{b} + \tau^+ \tau^-$  in a fast (ATLAS) simulation. Very difficult (if not hopeless)...

$$qq \rightarrow W^*W^*/Z^*Z^*qq \rightarrow qqh_i \rightarrow b\bar{b} + \tau^+\tau^-$$

- $M_{h_1} \sim 150$  GeV,  $M_{a_1} \sim 50$  GeV, usual WBF cuts.
- Assumes  $300 \text{ fb}^{-1}$  luminosity

S. Baffioni



The signal and backgrounds as a function of the invariant mass  $M_{bb\tau\tau}$  in GeV.

- $h_1 \rightarrow a_1 a_1 \rightarrow b\bar{b}\tau^+\tau^- \times 500$ .
- $t\bar{t}$ , —  $\gamma^* \rightarrow e^+e^-, \mu\mu$ , —  $Z \rightarrow \tau^+\tau^-$ ,
- total background.

## IV. Conclusions

### In the SM:

A large (th.+exp.) activity for the search of Higgs bosons at LHC.

Output: Higgs detection is in principle guaranteed for any mass.

The next step is to perform detailed test of the EWSB mechanism:

Measure the mass, total width, couplings to fermions/gauge bosons and self-coupling and determine spin and parity quantum numbers.

### In SUSY extensions:

Also in the MSSM, the light Higgs cannot escape detection and in large areas of parameter space, more than one Higgs can be found.

### But still watch out:

- there are still tricky areas of parameter space.
- complications if SUSY decays+production allowed.
- what about other SUSY non-minimal extensions?

..... Still a lot of work to be done .....