

Lecture 2

The gauge sector

The tree level predictions

$v, g, g' \Rightarrow$ the boson masses, M_W, M_Z , their self-interactions and their interactions with matter, all read off (!?) from \mathcal{L}

$$gW_\mu^+ J_\mu^- + h.c., \quad J_\mu^- = \bar{u} \gamma_\mu d + \bar{\nu} \gamma_\mu e$$

$$eA_\mu J_\mu^{em}, \quad J_\mu^{em} = 2/3 \bar{U} \gamma_\mu U - 1/3 \bar{D} \gamma_\mu D - \bar{E} \gamma_\mu E, \quad e = g \sin \theta$$

$$\frac{g}{\cos \theta} Z_\mu J_\mu^Z, \quad J_\mu^Z = \bar{\psi} \gamma_\mu (T_3 - Q \sin^2 \theta) \psi$$

$$U = u + (u_c)^c \equiv U_{Dirac}$$

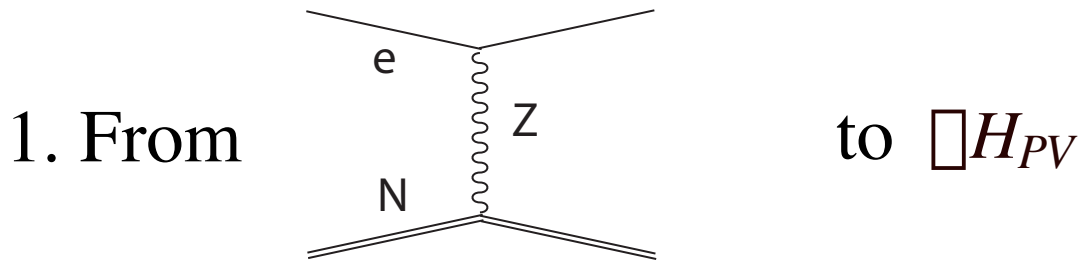
Highly predictive: a great variety of phenomena from

$l_{max} \approx 10^{-8} cm$ (Atomic Parity Violation) to

$l_{min} \approx 10^{-16} \div 10^{-17} cm$ (HERA, LEP, TEVATRON)

[Note, however, that $l_{min} \approx G^{1/2}$]

An example: Atomic Parity Violation



$\square H_{PV}$ must be local and proportional to G

$$\square \quad \square H_{PV} = \frac{G}{\sqrt{2}m_e} Q_W \square \cdot \square \square^3(\mathbf{r}) \quad \text{or} \quad A_{PV} = \frac{G}{\sqrt{2}m_e} Q_W \square \cdot \mathbf{q}$$

2. From the Z-exchange diagram and \mathcal{L}

$$A_{PV} = \frac{G}{\sqrt{2}} (\bar{E} \square_{\mu} \square E) (c_u \bar{U} \square_{\mu} U + c_d \bar{D} \square_{\mu} D) \quad \begin{aligned} c_u &= -1/2 + 4/3 \sin^2 \square \\ c_d &= 1/2 - 2/3 \sin^2 \square \end{aligned}$$

so that, by comparison, in the NRL (!?)

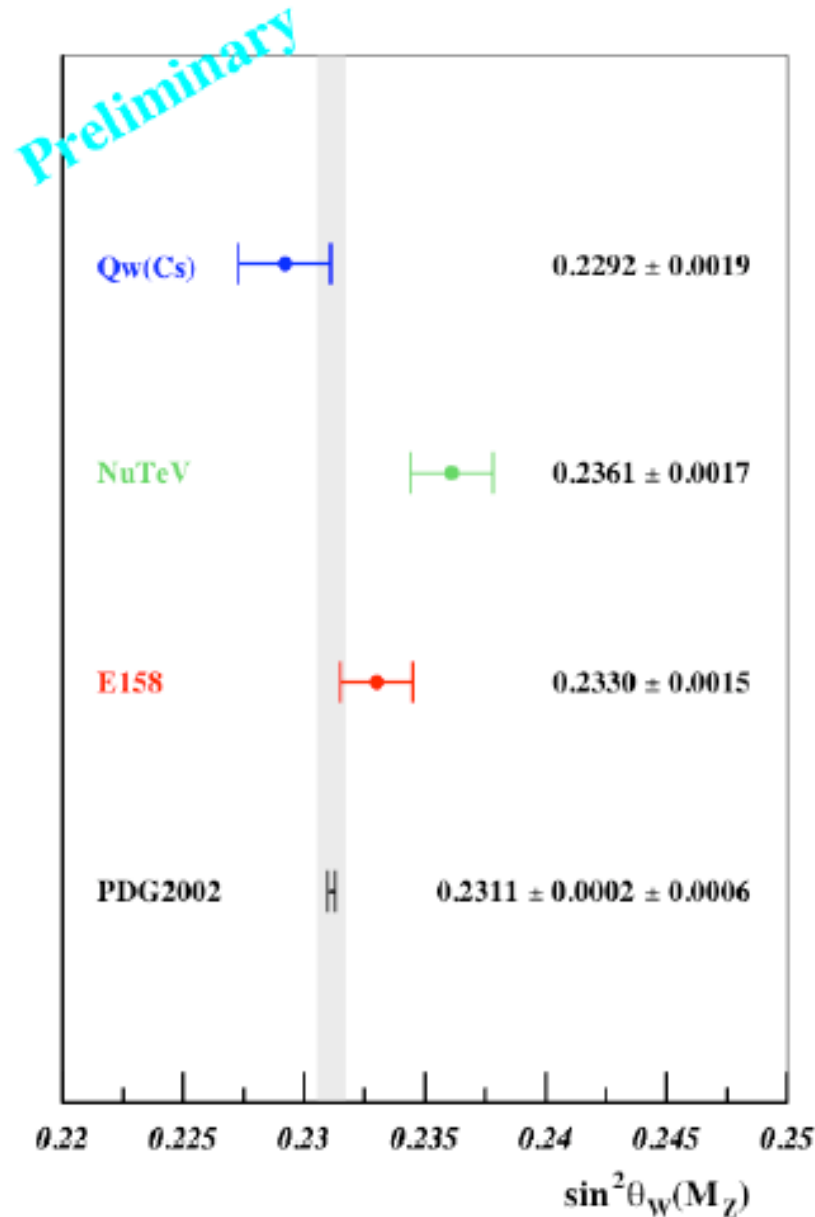
$$Q_W = -2(c_u n_u + c_d n_d) = (2 - 4 \sin^2 \square) Z - A$$

3. From the measured S-P mixing in Ce_{55}^{133} induced by $\square H_{PV}$

$$Q_W|_{exp} = 72.69(48)$$

$$Q_W|_{th} = 73.19(3)$$

3 low-energy neutral-current measurements



APV

□ N-scattering

Moller scattering

(Langacker)

Radiative effects

The prototype examples:

$$| \quad \frac{(g-2)_e}{2} \equiv a_e = a_e(QED) + a_e(\cancel{had}) + a_e(\cancel{EW})$$

$$a_e(QED) = \sum_n C_n^e \left(\frac{\hbar}{m_e c \lambda}\right)^n$$

From $\alpha^{-1}(at.int.) = 137.03600030(100) \quad \frac{\hbar c}{m_e c^2 \lambda} = 7.4 \cdot 10^{-9}$

$$a_e(th) = 1159652182.3(8.5) \cdot 10^{-12} \quad \text{versus}$$

$$a_e(exp) = 1159652185.9(3.8) \cdot 10^{-12}$$

Remarkable, but a pure QED effect (contained in \mathcal{L})

Radiative effects

The prototype examples (continued):

2 $\frac{(g-2)_\mu}{2} \equiv a_\mu$ probes distances $\frac{m_e}{m_\mu}$ smaller than a_e so that $\frac{\Delta a_e}{\Delta a_\mu}|_{univ} = \left(\frac{m_e}{m_\mu}\right)^2$

$$a_\mu = a_\mu(QED) + a_\mu(had) + a_\mu(EW)$$

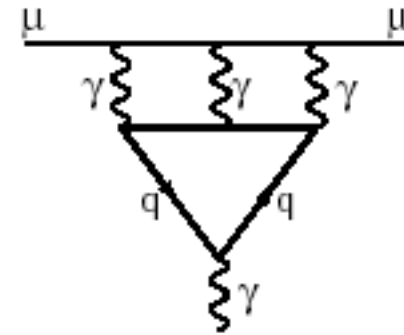
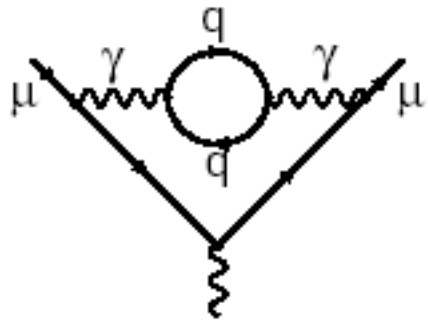
$$a_\mu(exp) = 116592080(60) \cdot 10^{-11}$$

$$a_\mu(EW) = 154(3) \cdot 10^{-11}$$

$$\frac{\Delta a_\mu}{a_\mu} = 5 \cdot 10^{-7}$$

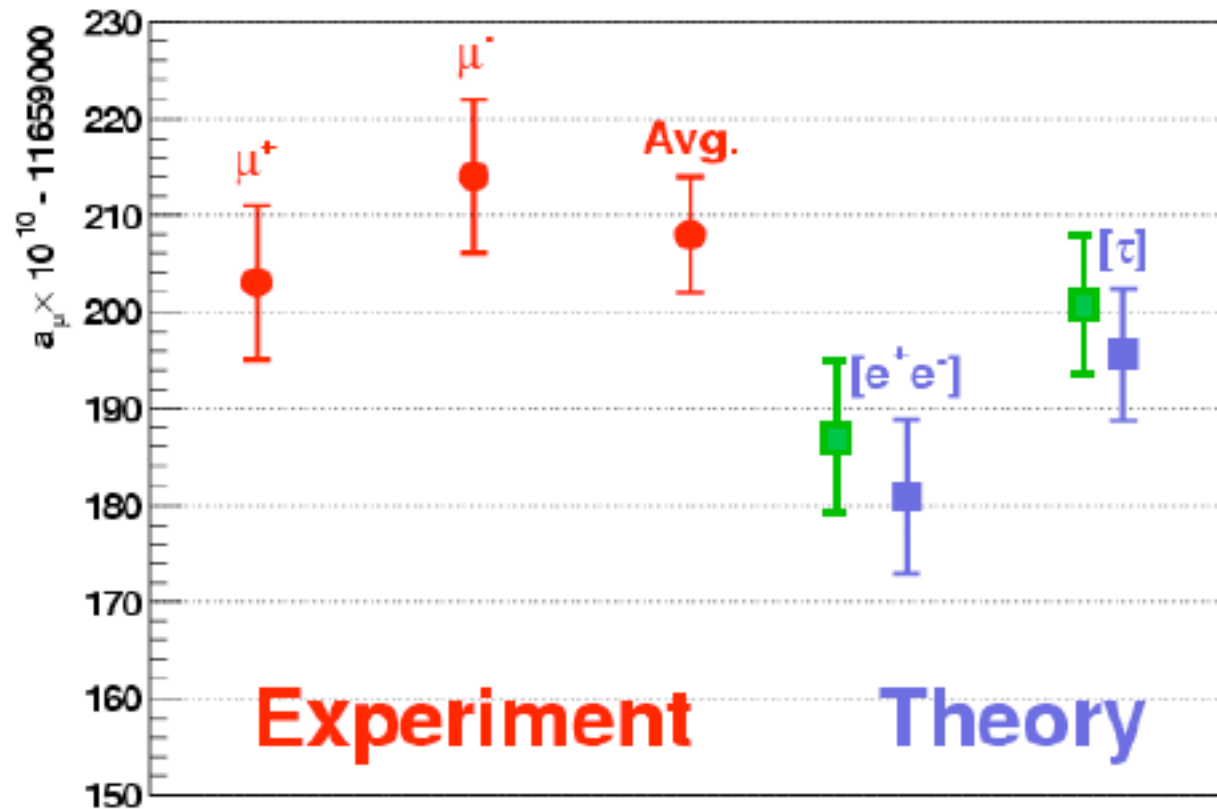
$$a_\mu(exp) - a_\mu(th) = \begin{bmatrix} 208 \pm 100 \\ 61 \pm 100 \end{bmatrix} \cdot 10^{-11} \quad \text{depending on which data one infers } a_\mu(had) \text{ from: upper (e+e-), lower (}\square\text{-decays)}$$

More sensitive, yes, but also more uncertain even apart from the \square / e+e- problem



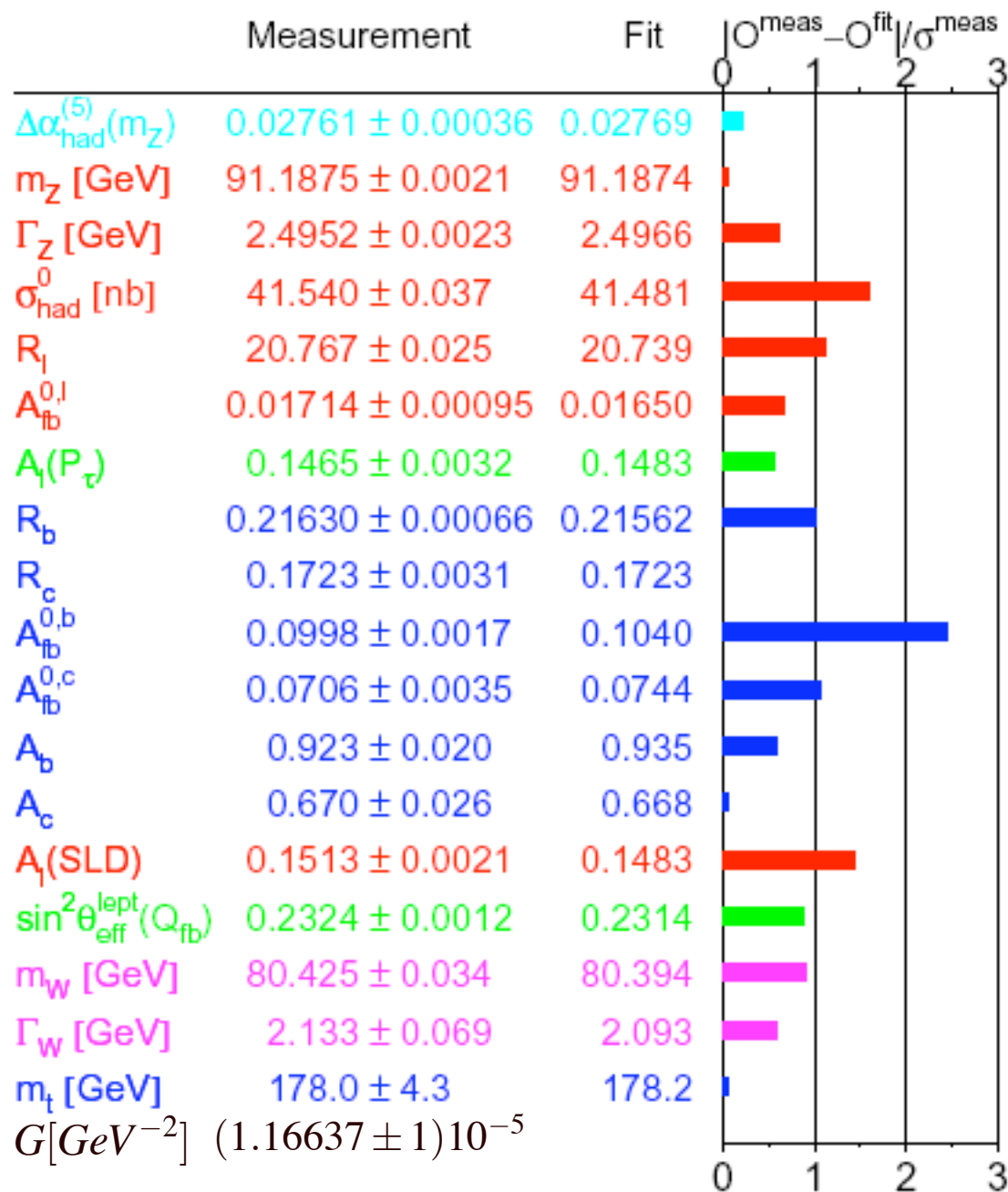
$$a_{\mu}^{\text{had,LO}} = \begin{cases} 6963(62)(36) \times 10^{-11} & e^+e^- \text{ based} \\ 7110(50)(8)(28) \times 10^{-11} & \tau \text{ based} \end{cases}$$

$$a_{\mu}^{\text{LbL}} = 136(25) \times 10^{-11}$$



(Vainshtein)

The ElectroWeak Precision Tests (Z and W prop.s)



parameters:

v, g, g' (at tree level)

m_t, m_h (from 1 loop on)

precision often better
than 10^{-3}

LEPEWWG 2004

In fact the EWPT bring together:

A. The gauge sector $(\frac{g^2}{4\pi}, \frac{g'^2}{4\pi})$


B. The flavor sector, through $\langle Q_{3t} \rangle$ $(\frac{\langle Q_t^2 \rangle}{4\pi} = \frac{Gm_t^2}{8\pi^2\sqrt{2}})$

C. The EWSB sector, mostly through $\frac{g^2}{4\pi} \log m_h$

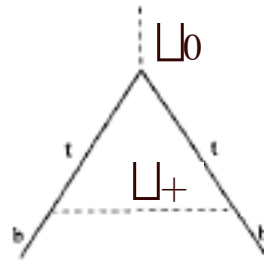
About A - In principle not different from the standard $\frac{e^2}{4\pi}$ expansion of QED, but with exchanged W's and Z's

About B - In fact, these effects (the dominant part) can be most easily computed with $g = g' = 0$, hence in a theory of top/bottom quarks, the Higgs h and the (unphysical) Goldstones (π 's)

The leading corrections of type B(!?)



$$\Rightarrow \frac{M_W^2}{M_Z^2 \cos^2 \theta} \equiv \rho = \frac{Z_2^+}{Z_2^0} = 1 + 3x$$

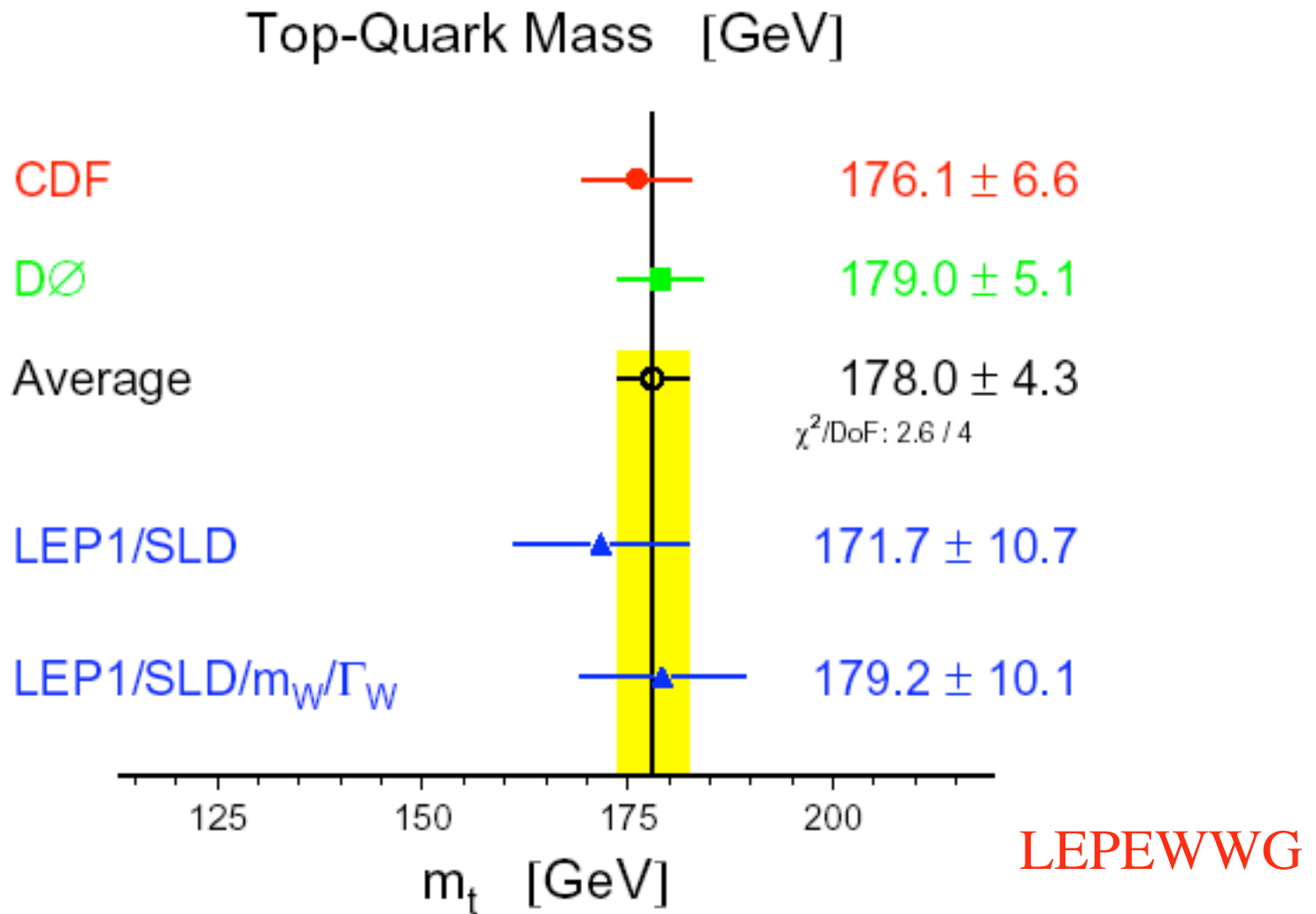


$$\Rightarrow \rho_{V_\mu}(Z \rightarrow b\bar{b}) \equiv \frac{g}{2 \cos \theta} \rho_{V_\mu} \frac{1 + \rho_b}{2} \quad \rho = \frac{Z_1}{Z_2^b} = -2x$$

where $x = \frac{Gm_t^2}{8s^2\sqrt{2}} \approx 0.5\%$

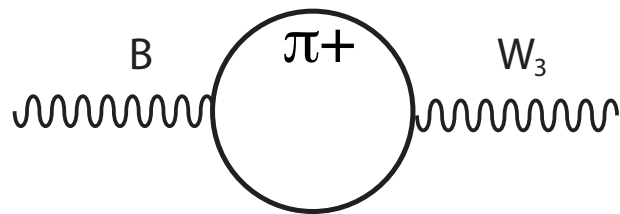
\Rightarrow clearly visible effects, used to get a range of top masses before the actual discovery (in 1993 $m_t = 120 \div 160 GeV$), now almost a background

Current comparison (2004)



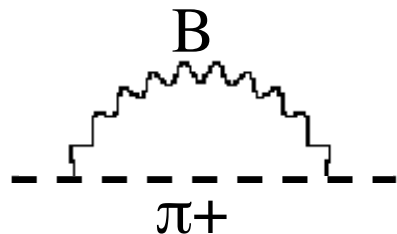
About the Higgs mass dependence

In the limit of infinite Higgs mass, ($m_h^2 = 4\mu v^2 : \mu \rightarrow \infty$) divergences appear: log's at 1 loop, quadratic at 2 loops. In the perturbative regime ($\mu \ll 4\mu$ or $m_h \ll 1-2$ TeV) the log's dominate, with 2 effects only (!?):



$$= \mu W_{3B}(p^2) g_{\mu\nu} + p_\mu p_\nu - \text{terms}$$

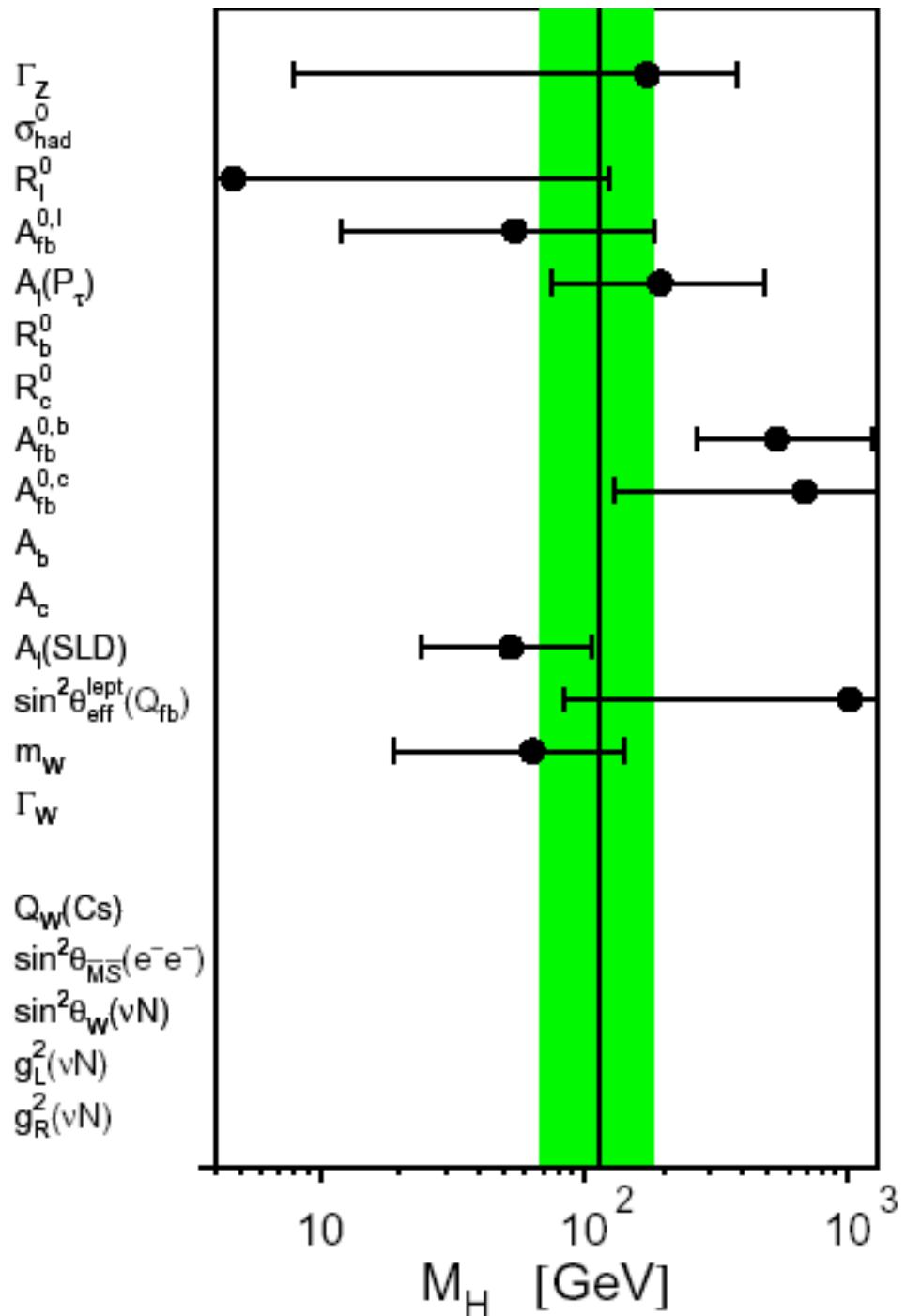
$$\Rightarrow \hat{S} \equiv \mu'_{W_3B}(0) = \frac{\mu}{24\mu \sin^2 \mu} \log \mu$$



$$= \mu_+(p^2)$$

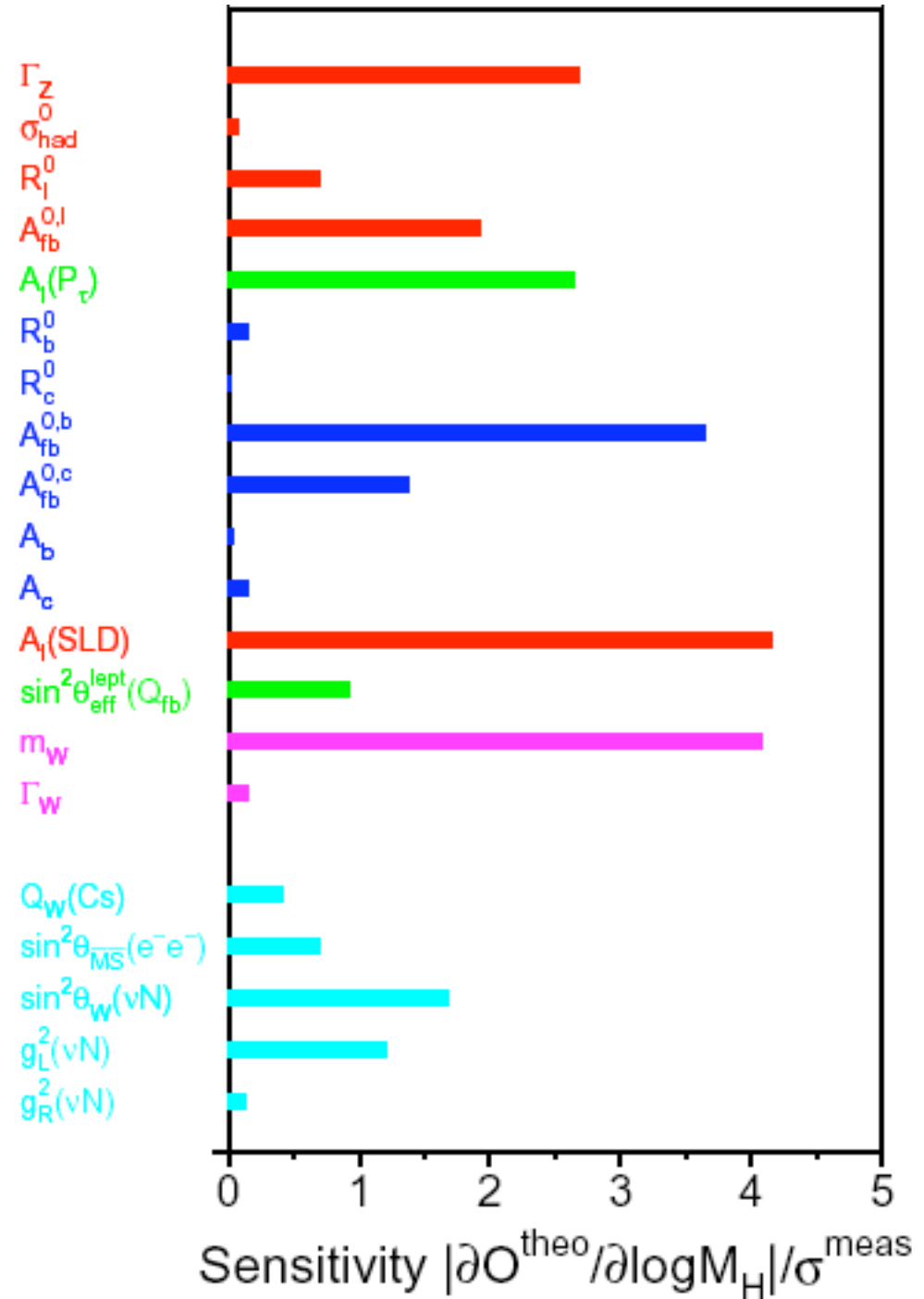
$$\Rightarrow \hat{T} \equiv \mu'_+(0) = -\frac{3\mu}{8\mu \sin^2 \mu} \log \mu$$

which spread in the various observables with $\log \mu \rightarrow \log m_h$



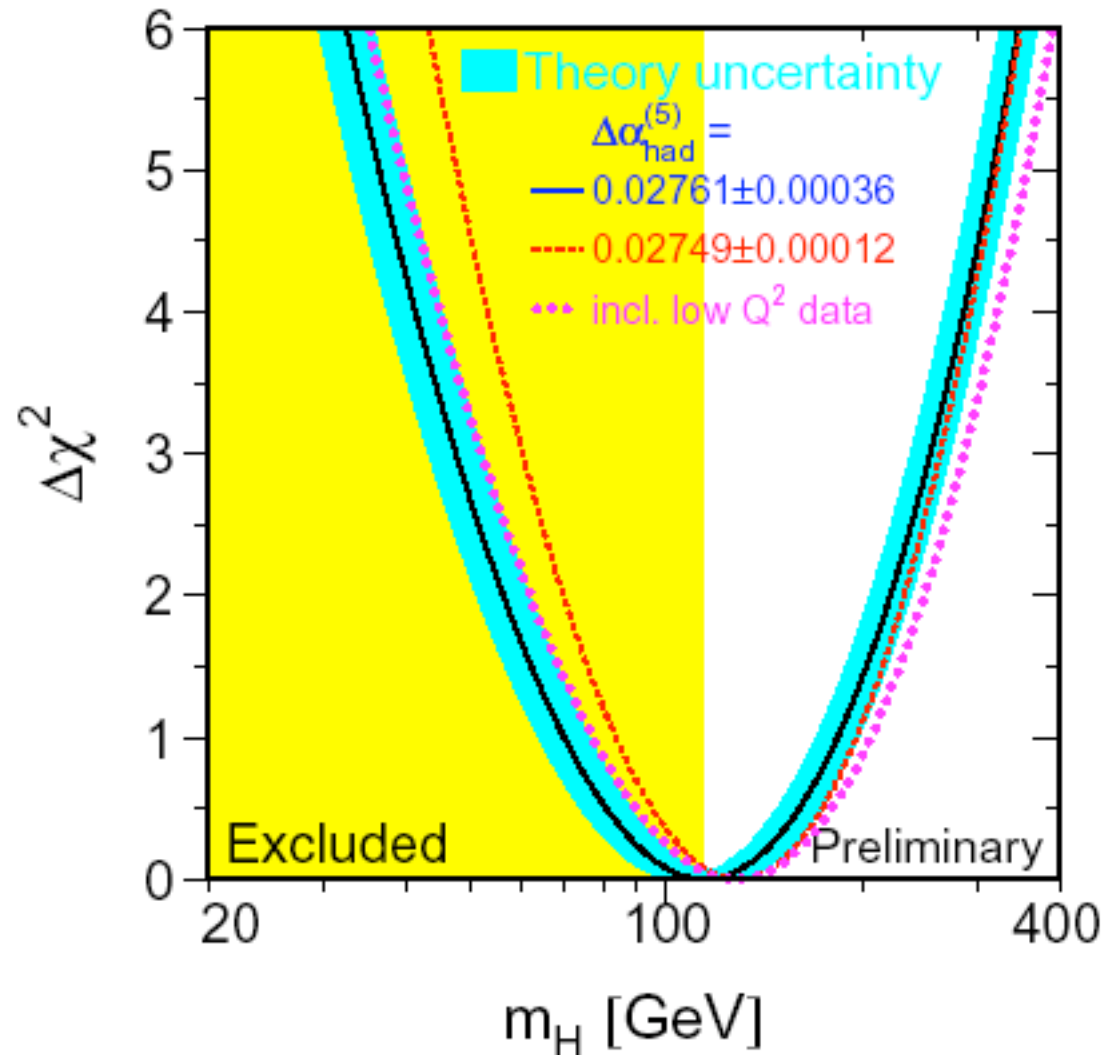
Riccardo Barbieri

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ElectroWeak Interactions: Theory 2004

The Higgs mass indirect determination in the SM



LEPEWWG

(more on this later in Lecture 5)

The foreseeable (?) future in EWPT

At present 3 main (relatively) uncertain parameters: $m_t, m_h, \Delta(M_Z)$
 \square 4 more precise measurements needed for a better test

	now	LHC	LC	Giga-Z
$\square \sin^2 \square_{eff} (10^{-5})$	16 (?)	15	?	1.3
$\square M_W [MeV]$	34	15	10	7
$\square M_t [GeV]$	4.3	1.0	0.2	0.1
$\Rightarrow \frac{\square m_h}{m_h}$	60%	15-20%	10-15%	5-10%

(A future improved $\square(M_Z)$?)