

# Lecture 4

## The neutrino-mass sector

$$\mathcal{L}^{(\nu\text{-mass})} = L_i \bar{\nu}_{ij} N_j \nu + N_i M_{ij} N_j$$

## 1. $M_{ij} = 0$ DIRAC masses

- Neutrinos are Dirac spinors ( $\nu_L, \nu_R \equiv N^C$ ), like charged fermions
- Lepton number is exactly conserved, like Baryon number
- In the mass-eigenstate basis  $W_\mu \bar{e}_i \gamma_\mu V_{ij} \nu_j$  with  $V$  like  $V$

□ Physical parameters:

3 masses, 3 angles and 1 phase

- A crucial question: Why  $m_\nu$ 's  $\ll m_e, m_q$ 's given the symmetry between neutrinos and charged fermions?

Technically, by  $m_\nu \ll m^u, m^d, m^e$

$$\mathcal{L}^{(\nu\text{-mass})} = L_i \bar{\nu}_{ij} N_j \nu + N_i M_{ij} N_j$$

## 2. $M_{ij} \neq 0$ MAJORANA masses

□ A basic asymmetry between neutrinos and charged fermions

$$\mathcal{L}^{(\nu\text{-mass})} = (\bar{\nu}^T N^T) \begin{pmatrix} 0 & \nu \\ \nu & M \end{pmatrix} \begin{pmatrix} \nu \\ N \end{pmatrix}$$

with  $\bar{\nu}$ ,  $N$  each 3-vectors and  $\nu$ ,  $M$  3x3 matrices

□ a 6x6 mass matrix □ 6 different eigenvalues (in general)

How come, with the same number of d.o.f. as for charged leptons?  $(e, e^c) \quad \square \quad (\nu, N)$

A Dirac spinor:  $e (\uparrow\downarrow) + \bar{e} (\uparrow\downarrow) : 4 \text{ d.o.f.}$

A Majorana spinor:  $\chi_M (\uparrow\downarrow) \quad \square \quad \bar{\chi}_M (\uparrow\downarrow) : 2 \text{ d.o.f.}$

Possible at all because Lepton number is violated

## 2. $M_{ij} \neq 0$ MAJORANA masses - continued

Don't we see 3 neutrinos in  $Z \rightarrow \nu_i \bar{\nu}_i$ ? ( $N_\nu = 2.9841(83)$ )

We see 3 *interacting* neutrinos, in a left helicity state  
(*active*, as opposed to *sterile*)

Why  $m_\nu \ll m_e, m_q$ ? Because  $M \gg \nu$

By diagonalizing the  
mass matrix

$$\begin{aligned} \nu_{light} &= \cos \theta \nu + \sin \theta N & m_l &\approx \frac{\nu^2}{M} & \theta &\approx \frac{\nu}{M} \\ \nu_{heavy} &= -\sin \theta \nu + \cos \theta N & m_h &\approx M \end{aligned}$$

By *integrating out*  
the heavy N's

$$= \nu^T (\nu^T \nu) \frac{1}{M} (\nu \nu)$$

$$\mathcal{L}^{(\nu\text{-mass})} \approx \nu_i m_{ij} \nu_j \quad m = \nu^T M^{-1} \nu^2$$

## 2. $M_{ij} \neq 0$ MAJORANA masses - continued

□ Physical parameters (with 3 light neutrinos hereafter):

$$\mathcal{L}^{lept} = \bar{L} \not{D} L + \bar{e}^c \not{D} e^c + e^T U_L^T m_d^e U_R^T e^c + \square^T V_\square^T m_d^\square V_\square$$

By going to the physical basis:

$$W_\mu \bar{e}_\mu \square_\mu \square \Rightarrow W_\mu \bar{e}_{ph} U_L V_\square^+ \square_\mu \square_{ph} \quad V = U_L V_\square^+$$

Counting intrinsic phases again

$$\begin{aligned} N(\text{phys. phases}) &= n^2 - \frac{n(n-1)}{2} - n^{(*)} = \frac{1}{2}n(n-1) \\ &= 3 \text{ (instead of 1) for } n=3 \end{aligned}$$

(\*unlike the  $(2n - 1)$  of the quark case, because the  $\square$ -phases defined by  $\square^T m_d^\square \square$  )

# The physical parameters precisely defined

**3 masses:**  $m_1, m_2, m_3$ , not by the order (as for charged fermion)  
but by

$$|\Delta m_{32}^2|, |\Delta m_{31}^2| > \Delta m_{21}^2 > 0 \quad \begin{cases} m_3 > m_2 > m_1 & \text{normal} \\ m_2 > m_1 > m_3 & \text{inverted} \end{cases}$$

**3 angles and 1(+2) phases:**

$$V = R_{23}(\theta_{23}) \begin{pmatrix} 1 & & \\ & e^{i\phi} & \\ & & 1 \end{pmatrix} R_{13}(\theta_{13}) R_{12}(\theta_{12}) \begin{pmatrix} 1 & & \\ & e^{i\phi} & \\ & & e^{i\phi} \end{pmatrix}$$

in  $W_\mu \bar{e} V_{\mu} e$

*Majorana phases*

# Means to determine the physical parameters

## A. Absolute spectrum

1. Endpoint of  $\beta$ -decay spectrum  $(Z, A) \rightarrow (Z + 1, A) + e + \bar{\nu}_i$

phase space:  $d\bar{\nu}_i = p_{\bar{\nu}_i} d\bar{\nu}_i (\bar{\nu}_i - m_i - E_e) dE_e$   $\bar{\nu}_i \equiv M_i - M_f$

$$\Rightarrow \frac{d\bar{\nu}_i}{dE_e} = \bar{\nu}_i |V_{ei}|^2 p_{\bar{\nu}_i} (\bar{\nu}_i - m_i - E_e) \quad p_{\bar{\nu}_i} = [(\bar{\nu}_i - E_e)^2 - m_i^2]^{1/2}$$

Currently  $\bar{\nu}_i |V_{ei}|^2 m_i^2 \lesssim (2eV)^2$  (+ info from osc.s:  $m_i|_{Max} \lesssim 2eV$  )

2. Cosmological observations

a - Density power spectrum distorted at small scales

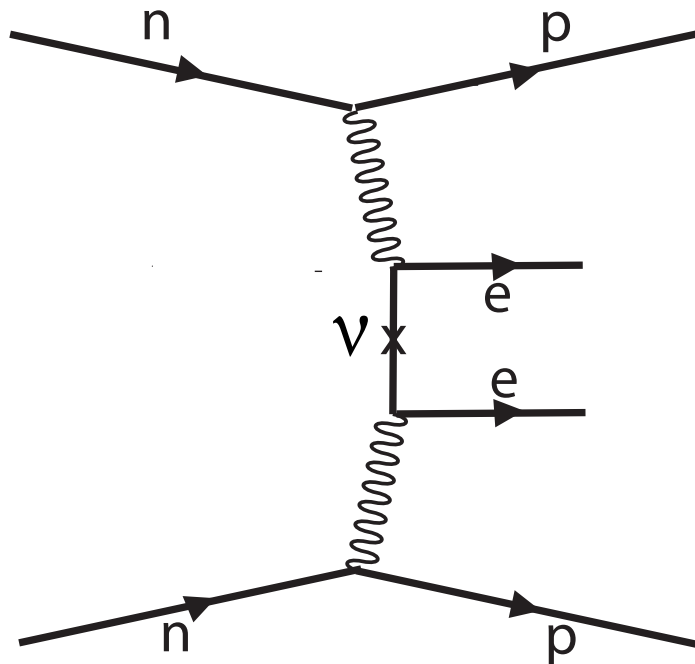
b - CMB distortions

Currently, mostly from a,  $\bar{\nu}_i m_i \lesssim 0.5eV$  (  $m_i|_{Max} \lesssim 0.15eV$  )

# Means to determine the physical parameters (continued)

3.  $\nu\nu_0$ -decay  $(Z, A) \rightarrow (Z + 2, A) + 2e$

only if Majorana, since L violated!



$$A \quad \sum_i V_{ei}^2 m_i M_{nucl} \equiv m_{ee} M_{nucl}$$

$$T_{1/2} \quad \frac{1}{|m_{ee}|^2 |M_{nucl}|^2}$$

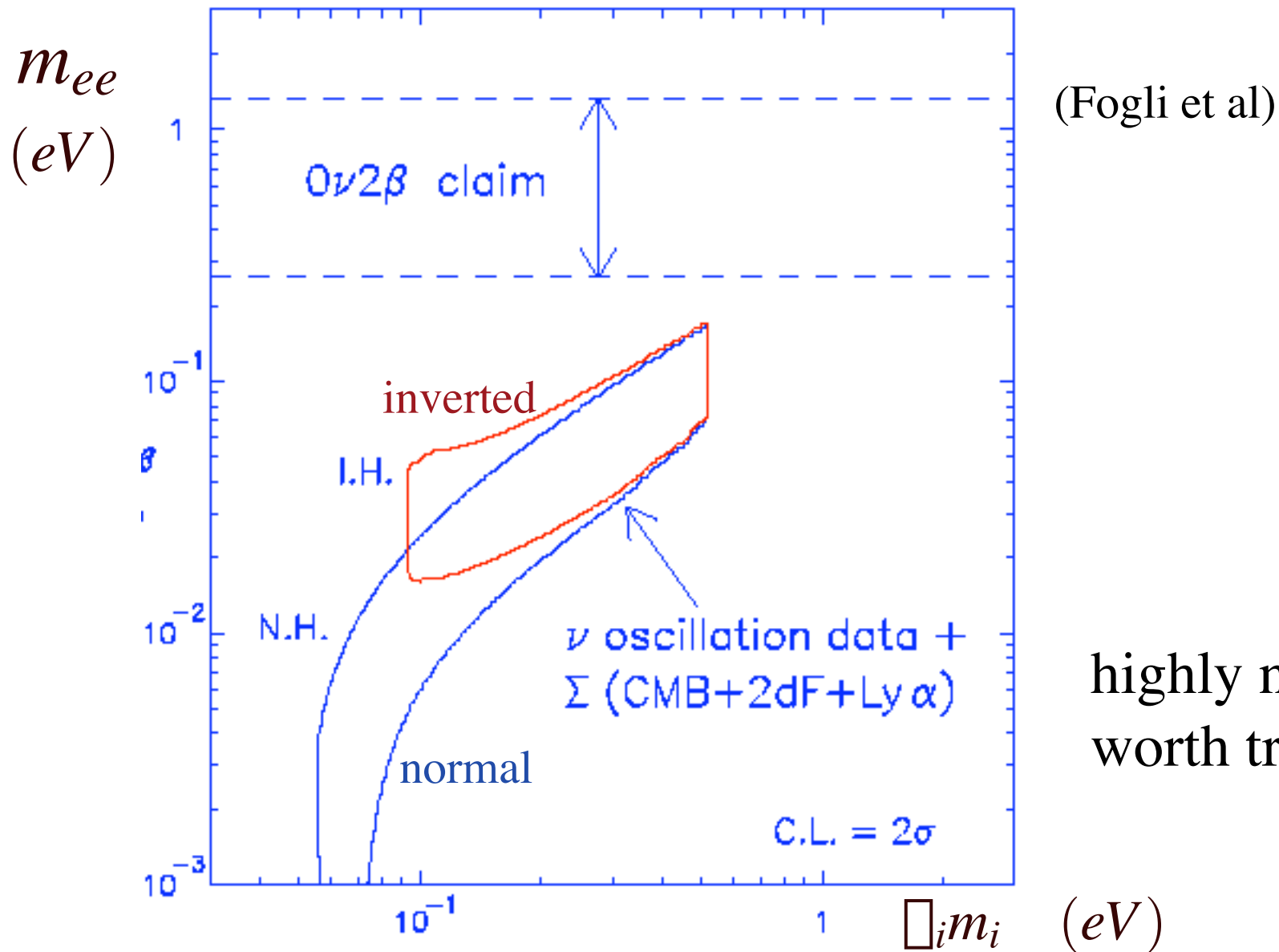
$$m_{ee} = c_{13}^2 (c_{12}^2 m_1 + s_{12}^2 e^{2i\phi} m_2) + s_{13}^2 m_3 e^{2i\phi} \quad (!?)$$

$$\sigma_{th}(M_{nucl}) = O(M_{nucl})$$

Currently: a claimed observation at  $0.17eV < |m_{ee}| < 2.0eV$



### 3. $\beta\beta_{0\nu}$ - decay with oscillation and cosmological data (current 2004)



highly non trivial but worth trying

# Means to determine the physical parameters (continued)

## B. Mixing parameters from neutrino oscillations

The standard picture: Calculate the amplitude for a  $|\nu_{l_\alpha}\rangle$  at  $t=0$  to become a  $|\nu_{l_\beta}\rangle$  at time  $t$

$$|\nu_{l_\alpha}\rangle \xrightarrow[t=0]{R=ct, t} |\nu_{l_\beta}\rangle \quad l_{\alpha,\beta} = e, \mu, \tau$$

$$|\nu_{l_\alpha}, t\rangle = \sum_{i=1,2,3} U_{\alpha i}(t) |\nu_{l_i}\rangle$$

$$i \frac{d}{dt} \nu = V \mathbf{E} V^\dagger \nu \quad \mathbf{E}_{ij} = \left( p + \frac{m_i^2}{2p} \right) \delta_{ij}$$

$$\langle \nu_{l_\beta} | \nu_{l_\alpha}, t \rangle = \sum_{i=1,2,3} V_{i\alpha}^* V_{i\beta} e^{-i \frac{m_i^2 t}{2p}} \quad (!?)$$

up to an irrelevant phase factor

Note: Majorana phases and absolute  $m_i^2$  's

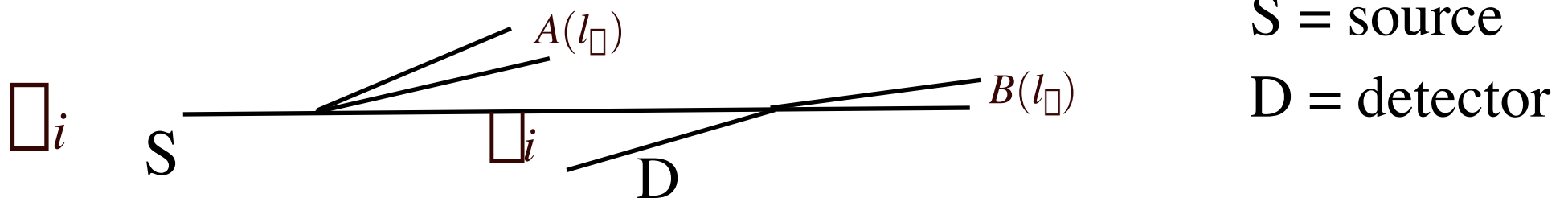
$$\text{irrelevant to } \left| \langle \nu_{l_\beta} | \nu_{l_\alpha}, t \rangle \right|$$

# (Apparent) Paradoxes of the standard picture

1. Why is the  $\square$ -state given a definite  $p$  rather than  $E$ ?
  2. If  $p$  defined, how about coherence?  $v_1 = p/E_1 \neq v_2 = p/E_2$
  3. If  $p$  defined,  $\square x = \dots$ . Why then  $t = R/c$ ?
- etc...

Suggested framework to address them:

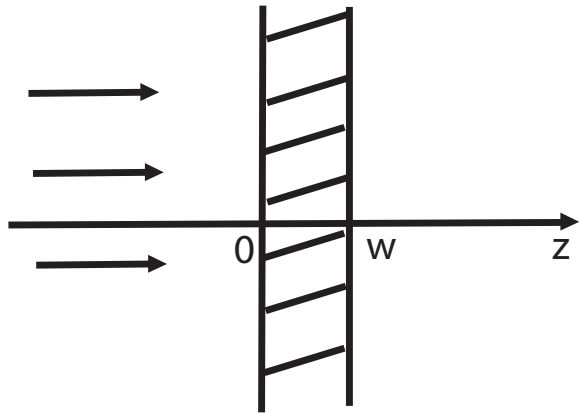
The amplitude of a physical process =



treated in second-order time-dependent perturbation theory (!?)

# Matter effects in $\nu$ -oscillations

$\square$ -propagation in matter: like in optics



$$e^{ipz} \Rightarrow e^{ipz} \left( 1 + iwNf(0) \frac{2\square}{p} \right)$$

$$e^{ipz} \Rightarrow e^{(inpw + p(z-w))} \approx e^{ipz} (1 + iw p(n-1))$$

$$\Rightarrow n - 1 = Nf(0) \frac{2\square}{p^2}$$

$\square$  A modified Schroedinger equation in a medium of electron number density  $N_e$

$$i \frac{d}{dt} \square = (V \mathbf{E} V^\dagger + \mathbf{A}) \square$$

with  $\mathbf{A} = 0$  except for  $\mathbf{A}_{ee} = \pm \sqrt{2} G N_e$  ( $+\square, -\square$ ) (!?)

Nucleons do not contribute to  $\mathbf{A}$  (!?)

# Solving the equation (!?)

(in a simple, but relevant, 2x2 case)

$$E_{matter} = \pm \frac{1}{2} [(\Delta \cos 2\Delta_\nu - \sqrt{2}GN_e)^2 + \Delta^2 \sin^2 2\Delta_\nu]^{-1/2}$$

$$\tan 2\Delta_m = \frac{\Delta \sin 2\Delta_\nu}{\Delta \cos 2\Delta_\nu - \sqrt{2}GN_e}$$

$$\Delta \equiv \frac{\Delta m^2}{2p}$$

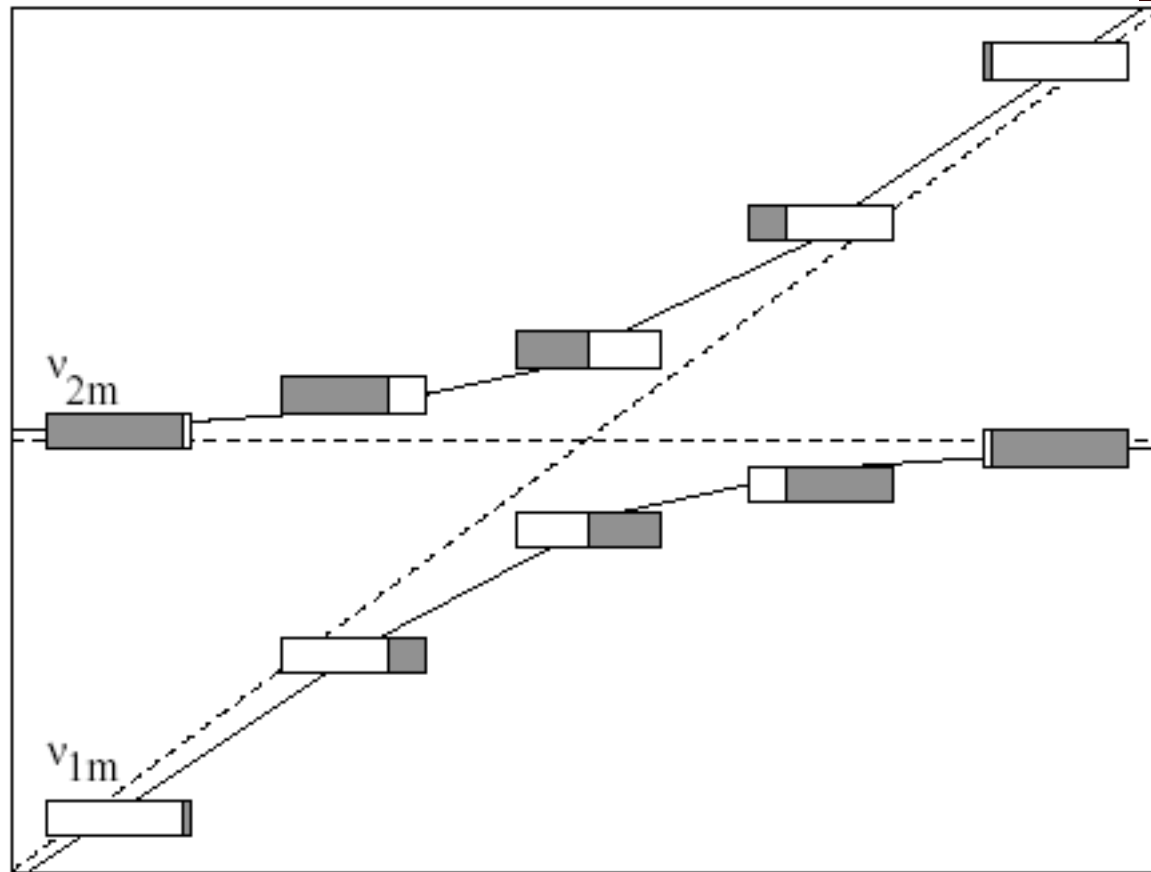
$\Delta_\nu$  = vacuum angle

small  $\Delta_\nu$

*adiabatic*  
conditions

$$\Delta_{res} \ll \Delta R$$

$$\Delta R \frac{dV}{dr} = \Delta E_{min}$$



(Smirnov)

$N_e$

# Current knowledge (2004)

$$|\Delta m_{32}^2| \approx |\Delta m_{31}^2| \approx 2 \cdot 10^{-3} eV^2 \quad (\Delta m_{Atm}^2) \qquad \Delta m_{21}^2 \approx 8 \cdot 10^{-5} eV^2 \quad (\Delta m_{Sun}^2)$$

$$\Delta_{23} = 36^\circ \div 54^\circ \quad (90\% \text{ C.L.})$$

$$\tan^2 \Delta_{12} = 0.40_{-0.07}^{+0.09} \quad (\Delta_{12} = 32^\circ \pm 3^\circ)$$

$$\sin^2 \Delta_{13} < 0.03 \quad (90\% \text{ C.L.})$$

# Current ignorance (2004)

0. Dirac or Majorana?

1. Absolute scale?  $0.05 eV \lesssim m_i|_{Max} \lesssim 0.15 eV$  (if cosmo)

2. Sign of  $\Delta m_{32}^2$  (inverted or normal spectrum)

3.  $\Delta_{13} = ?$

4.  $\Delta = ?$  Need  $\Delta_{13} \neq 0$  to be defined at all