



Tracking with Solid State Detectors

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IIb Tracking with Solid State Detectors

2b - Tracking with
Solid State Detectors

- **Lecture 1 - Introduction** C. Joram, L. Ropelewski
- **Lecture 2 - Tracking Detectors** L. Ropelewski, M. Moll
 - **2a) Tracking with Gas detectors**
 - **2b) Tracking with Solid State Detectors** **Michael Moll (CERN - PH/DT2)**
 - Why use Semiconductor Detectors ?
 - How are Silicon Detectors made and how do they work ?
 - Detector types: Microstrip and Pixel Detectors, CCDs
 - Examples: Detectors at LHC
 - Radiation Damage in Silicon Detectors
 - Outlook: Radiation tolerant detectors
 - References
- **Lecture 3 - Scintillation and Photodetection** C. D'Ambrosio, T. Gys
- **Lecture 4 - Calorimetry, Particle ID** C. Joram
- **Lecture 5 - Particle ID, Detector Systems** C. Joram, C. D'Ambrosio

Transparencies: http://cern.ch/ph-dep-dt2/lectures_PD_2005.htm



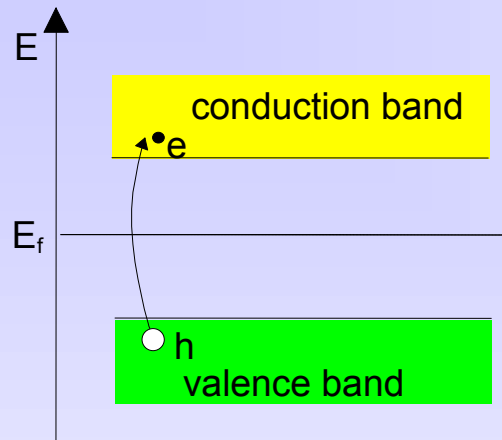
■ Some characteristics of Silicon crystals

- **Small band gap** $E_g = 1.12 \text{ eV} \Rightarrow E(\text{e-h pair}) = 3.6 \text{ eV} (\approx 30 \text{ eV for gas detectors})$
- **High specific density** 2.33 g/cm^3 ; $dE/dx \text{ (M.I.P.)} \approx 3.8 \text{ MeV/cm} \approx 106 \text{ e-h}/\mu\text{m}$ (average)
- **High carrier mobility** $\mu_e = 1450 \text{ cm}^2/\text{Vs}$, $\mu_h = 450 \text{ cm}^2/\text{Vs} \Rightarrow$ fast charge collection ($< 10 \text{ ns}$)
- **Very pure** $< 1 \text{ ppm}$ impurities and $< 0.1 \text{ ppb}$ electrical active impurities
- **Rigidity** of silicon allows thin self supporting structures
- **Detector production by microelectronic techniques**
 \Rightarrow well known industrial technology, relatively low price, small structures easily possible

■ Alternative semiconductors

- Diamond
- GaAs
- Silicon Carbide
- Germanium

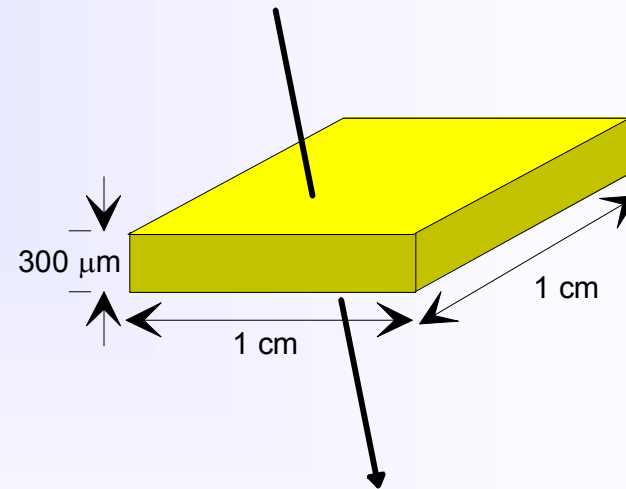
	Diamond	SiC (4H)	GaAs	Si	Ge
Atomic number Z	6	14/6	31/33	14	32
Bandgap E_g [eV]	5.5	3.3	1.42	1.12	0.66
$E(\text{e-h pair})$ [eV]	13	7.6-8.4	4.3	3.6	2.9
density [g/cm^3]	3.515	3.22	5.32	2.33	5.32
e-mobility μ_e [cm^2/Vs]	1800	800	8500	1450	3900
h-mobility μ_h [cm^2/Vs]	1200	115	400	450	1900



In a pure intrinsic (undoped) semiconductor the electron density n and hole density p are equal.

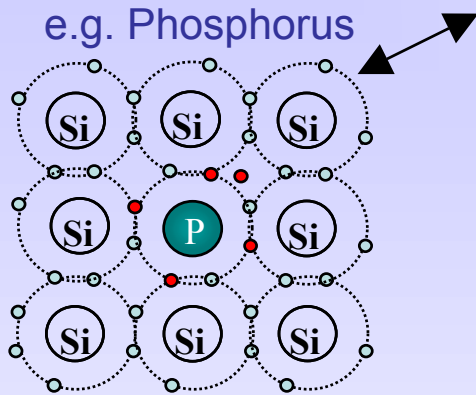
$$n = p = n_i \quad \text{For Silicon: } n_i \approx 1.45 \cdot 10^{10} \text{ cm}^{-3}$$

$4.5 \cdot 10^8$ free charge carriers in this volume,
but only $3.2 \cdot 10^4$ e-h pairs produced by a M.I.P.



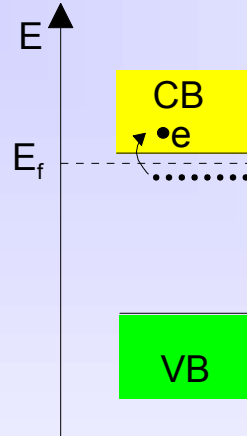
⇒ Reduce number of free charge carriers, i.e. deplete the detector

⇒ **Most detectors make use of reverse biased p-n junctions**



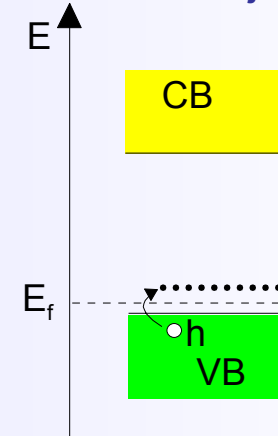
■ Doping: n-type Silicon

- add elements from Vth group
⇒ **donors** (P, As,..)
- electrons are majority carriers



■ Doping: p-type Silicon

- add elements from IIIrd group
⇒ **acceptors** (B,..)
- holes are the majority carriers



■ Resistivity

- carrier concentrations n, p
- carrier mobility μ_n, μ_p

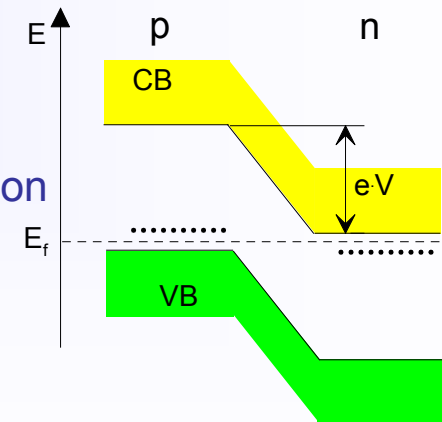
$$\rho = \frac{1}{q_0(\mu_n n + \mu_p p)}$$

	detector grade	electronics grade
doping	$\approx 10^{12} \text{ cm}^{-3}$	$\approx 10^{17} \text{ cm}^{-3}$
resistivity ρ	$\approx 5 \text{ k}\Omega\cdot\text{cm}$	$\approx 1 \text{ }\Omega\cdot\text{cm}$

■ p-n junction

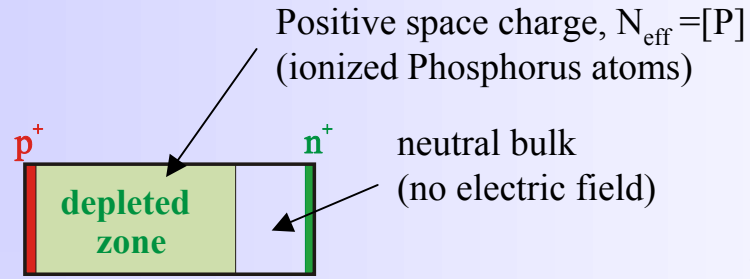
There must be a single Fermi level !

- ⇒ band structure deformation
- ⇒ potential difference
- ⇒ depleted zone

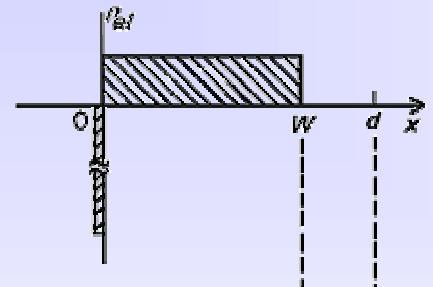


Poisson's equation

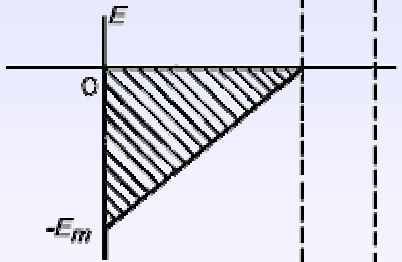
$$-\frac{d^2}{dx^2} \phi(x) = \frac{q_0}{\epsilon\epsilon_0} \cdot N_{eff}$$



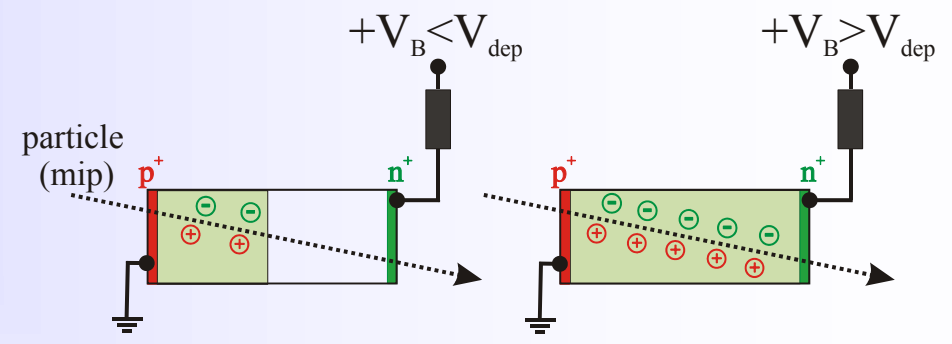
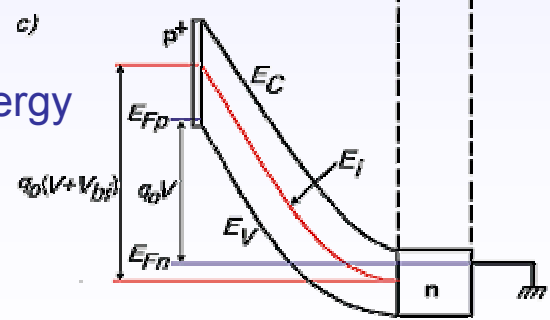
a)
Electrical charge density



b)
Electrical field strength



c)
Electron potential energy



Full charge collection only for $V_B > V_{dep}$!

depletion voltage

$$V_{dep} = \frac{q_0}{\epsilon\epsilon_0} \cdot |N_{eff}| \cdot d^2$$

effective space charge density



Calculation of depletion voltage (diode)

extra slide
not shown

Poisson's equation

$$-\frac{d^2}{dx^2} \phi(x) = \frac{q_0}{\epsilon\epsilon_0} \cdot N_{eff}$$

with $\frac{d}{dx} \phi(x=w) = 0$
 $\phi(x=w) = 0$

$$-\frac{d}{dx} \phi(x) = \frac{q_0}{\epsilon\epsilon_0} \cdot N_{eff} \cdot (x-w)$$

$$\phi(x) = \frac{1}{2} \cdot \frac{q_0}{\epsilon\epsilon_0} \cdot N_{eff} \cdot (x-w)^2$$

w = depletion depth

d = detector thickness

U = voltage

N_{eff} = effective doping concentration

$$C = \frac{dQ}{dU} = \frac{dQ \cdot dw}{dw \cdot dU}$$

$$dQ = q_0 \cdot |N_{eff}| \cdot A \cdot dw$$
$$dw = \sqrt{\frac{\epsilon\epsilon_0}{q_0 |N_{eff}| 2U}} \cdot dU$$

depletion voltage

$$V_{dep} = \frac{q_0}{2\epsilon\epsilon_0} \cdot |N_{eff}| \cdot d^2$$

effective space charge density

$$w(V) = \sqrt{\frac{2\epsilon\epsilon_0}{q_0 |N_{eff}|} \cdot V}$$

$$C(U) = A \cdot \sqrt{\frac{\epsilon\epsilon_0 q_0 |N_{eff}|}{2U}}$$

$$C(w) = \frac{\epsilon\epsilon_0 A}{w}$$



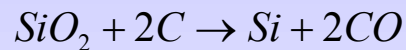
How to make a Float Zone Silicon wafer?

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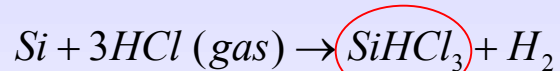
2b - Tracking with
Solid State Detectors

■ Produce a polysilicon rod

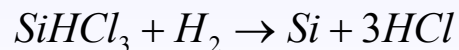
- Melt very **pure sand** (SiO_2) together with coke ($\sim 1800^\circ\text{C}$)



- Grind the “metallurgical grade silicon” (98% Si) and expose it to hydrochloric gas



- **Trichlorsilane** boils at 31.7°C and can thus be distilled and purified
- Deposit silicon in a Chemical Vapour Deposition process



- Cast silicon into a **polycrystalline silicon rod**

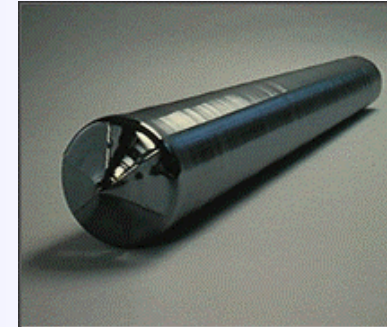
■ Float Zone process

- Using a single Si crystal seed, melt the vertically oriented rod onto the seed using RF power and “pull” the **monocrystalline ingot**



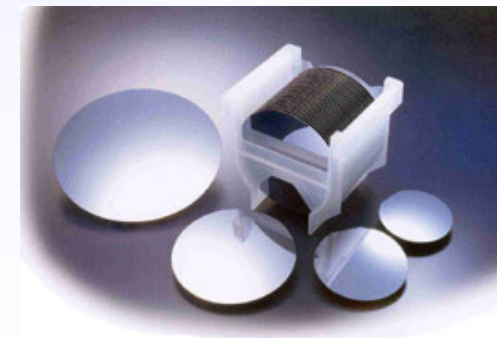
■ Monocrystalline Ingot

- grind into round shape
- make the flat or a notch



■ Wafer production

- Slice the ingot into wafers of $300\text{-}500\ \mu\text{m}$ (diamond saw)
- lapping of wafers
- etching of wafers
- polishing of wafers



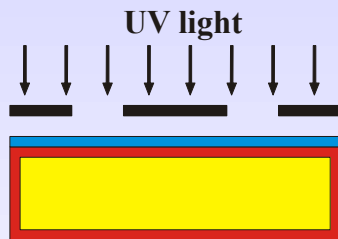
■ A "simple" production sequence (schematic)

n-type silicon

- Polished n-type silicon wafer (typical $\rho \sim 1-10 \text{ K}\Omega\text{cm}$)



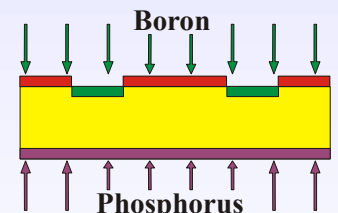
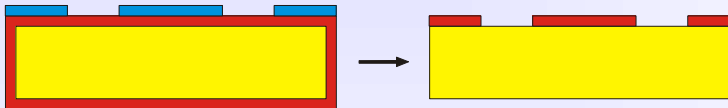
- Oxidation (800-1200°C)



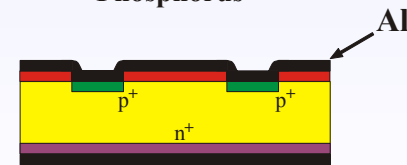
- Photolithography (coat with photo resist; align mask, expose to UV light, develop photoresist);

Etching of oxide

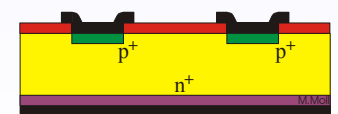
etch



- Doping with boron and phosphorus by implantation (or by diffusion)
Annealing to cure radiation damage and activate dopants
 - p⁺ n junction on front side
 - n n⁺ ohmic contact on back side

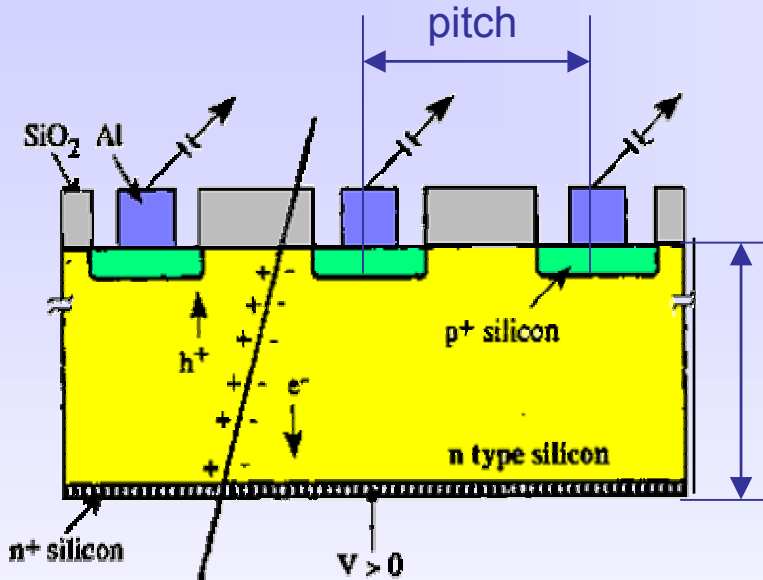


- Aluminize surface (e.g. by evaporation)



- Pattern metal for diode contacts

- Segmentation of the p⁺ layer into strips (Diode Strip Detector) and connection of strips to individual read-out channels gives spatial information



typical thickness: 300 μ m (150 μ m - 500 μ m used)

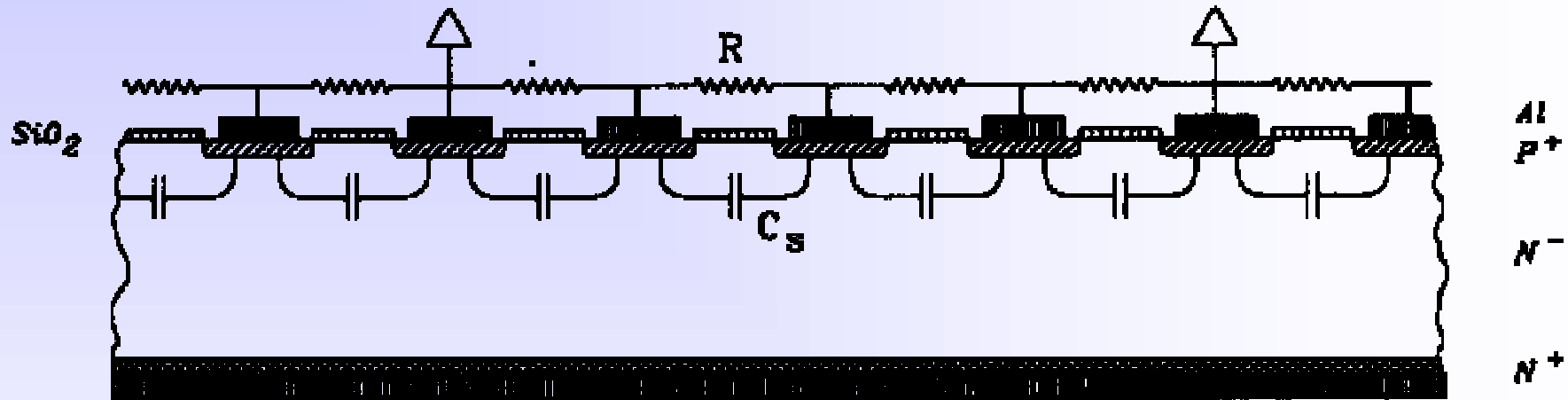
- using n-type silicon with a resistivity of $\rho = 2 \text{ K}\Omega\text{cm}$ ($N_D \sim 2.2 \cdot 10^{12} \text{cm}^{-3}$) results in a depletion voltage $\sim 150 \text{ V}$

- Resolution σ depends on the pitch p (distance from strip to strip)
 - e.g. detection of charge in binary way (threshold discrimination) and using center of strip as measured coordinate results in

$$\sigma = \frac{p}{\sqrt{12}}$$

typical pitch values are 20 μ m– 150 μ m \Rightarrow 50 μ m pitch results in 14.4 μ m resolution

- Analog readout (measurement of signal height) of every strip leads to substantial improvement of position resolution, however not every strip has to be read:



- Charge division readout reduces the number of readout channels as only a fraction of the strips is connected to readout amplifier.
- Charge collected at the interpolation strips is divided between the two neighboring readout channels according to the relative position.



Bias resistor and AC Coupling

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2b - Tracking with Solid State Detectors

■ Bias resistor

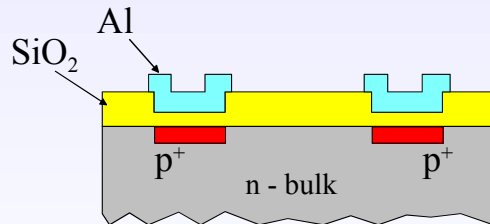
- Need to isolate strips from each other to collect/measure charge on each strip
⇒ high impedance bias connection ($\approx 1\text{M}\Omega$ resistor)

■ Coupling capacitor

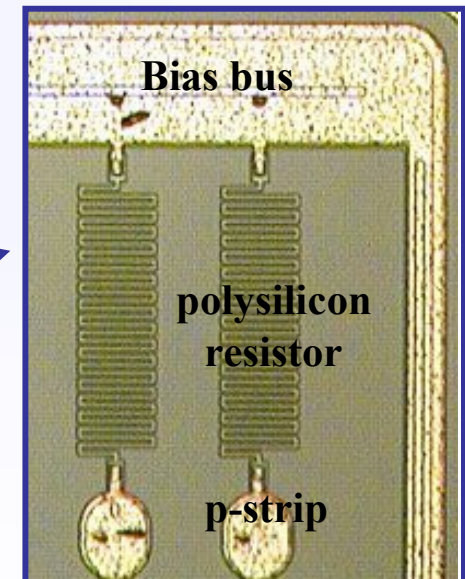
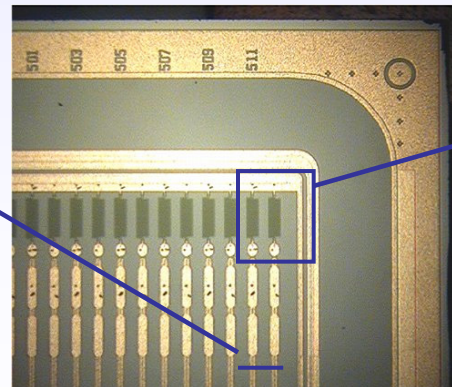
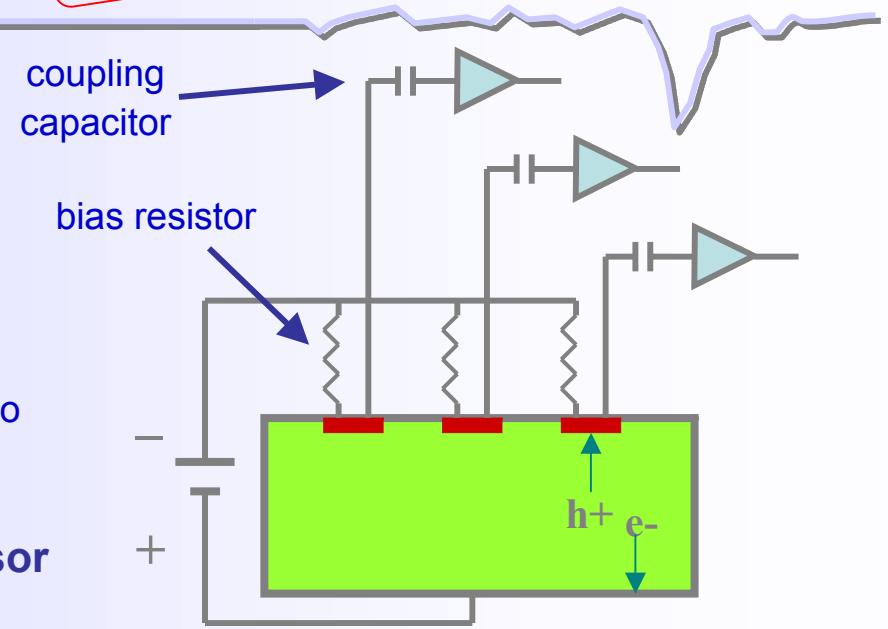
- Couple input amplifier through a capacitor (AC coupling) to avoid large DC input from leakage current

■ Integration of capacitors and resistors on sensor

- Bias resistors via deposition of doped polysilicon
- Capacitors via metal readout lines over the implants but separated by an insulating dielectric layer ($\text{SiO}_2, \text{Si}_3\text{N}_4$).



- ⇒ nice integration
- ⇒ more masks, processing steps
- ⇒ pin holes





The Charge Signal

extra slide
not shown

2b - Tracking with
Solid State Detectors

■ Collected Charge for a Minimum Ionizing Particle (MIP)

• Mean energy loss

dE/dx (Si) = 3.88 MeV/cm
⇒ 116 keV for 300 μ m thickness

• Most probable energy loss

$\approx 0.7 \times$ mean
⇒ 81 keV

• 3.6 eV to create an e-h pair

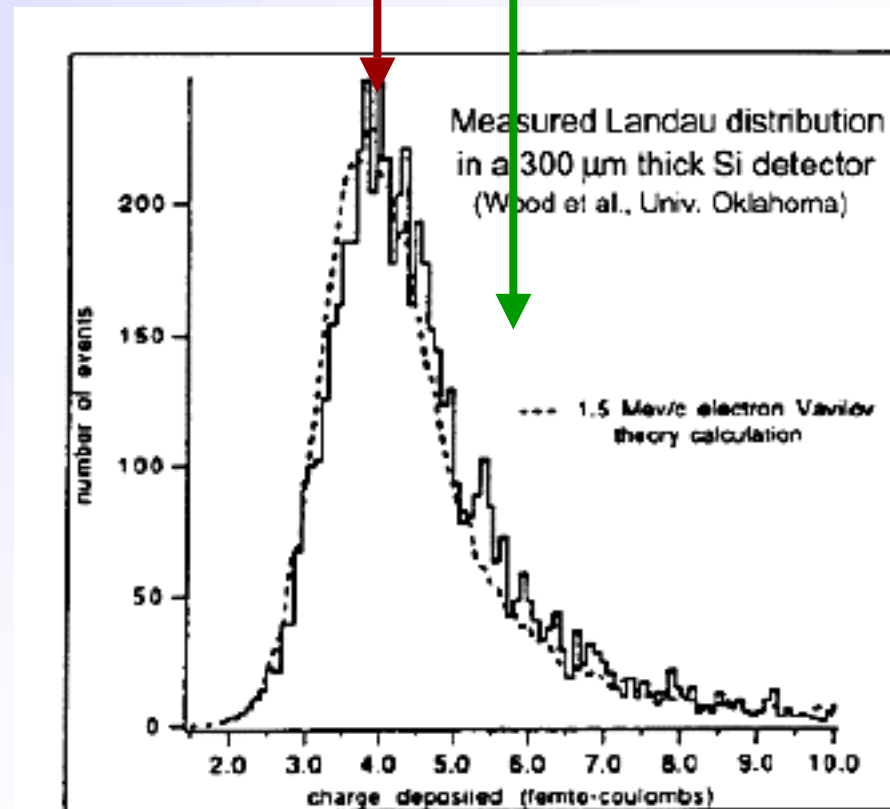
⇒ 72 e-h / μ m (mean)
⇒ 108 e-h / μ m (most probable)

• Most probable charge (300 μ m)

≈ 22500 e ≈ 3.6 fC

Most probable charge $\approx 0.7 \times$ mean

Mean charge



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Signal to noise ratio (S/N)

- **Landau distribution** has a low energy tail
 - becomes even lower by noise broadening

Noise sources: (ENC = Equivalent Noise Charge)

- Capacitance $ENC \propto C_d$

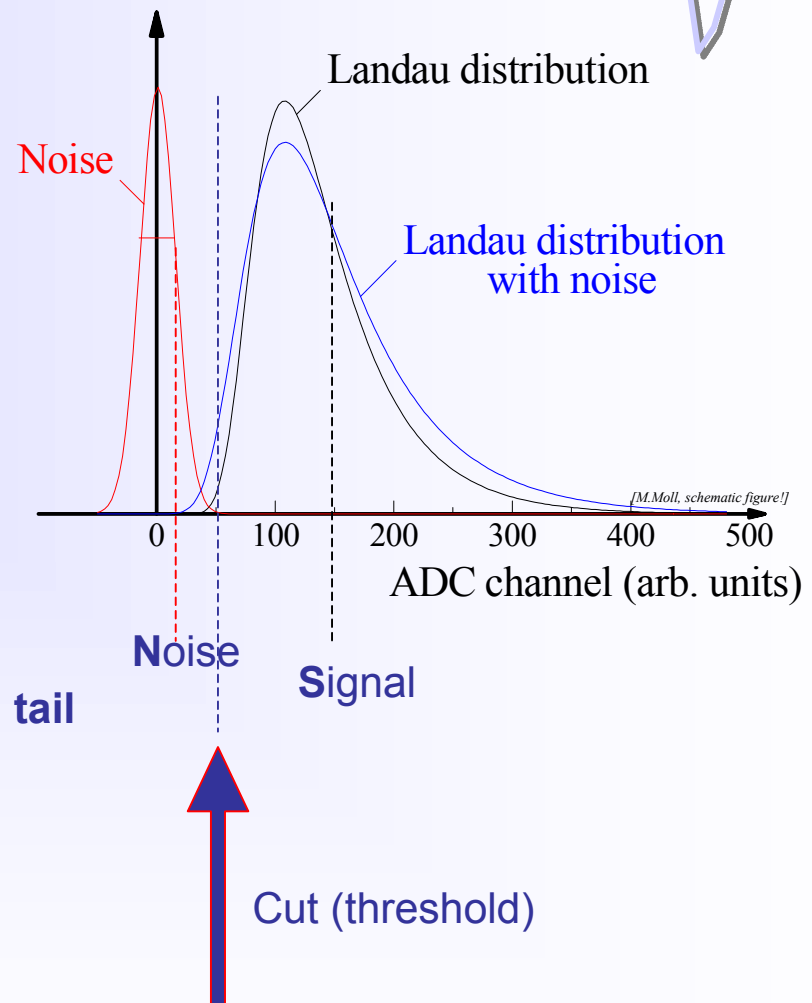
- Leakage Current $ENC \propto \sqrt{I}$

- Thermal Noise (bias resistor) $ENC \propto \sqrt{k_B T / R}$

- **Good hits selected by requiring $N_{ADC} > \text{noise tail}$**
 - If cut too high \Rightarrow efficiency loss
 - If cut too low \Rightarrow noise occupancy

- **Figure of Merit: Signal-to-Noise Ratio S/N**

- **Typical values >10-15, people get nervous below 10.**
Radiation damage severely degrades the S/N.





Charge Collection time and diffusion

extra slide
not shown

2b - Tracking with
Solid State Detectors

Charge Collection time

- Drift velocity of charge carriers $v \approx \mu E$, so drift time, $t_d = d/v = d/\mu E$

Typical values: $d=300 \mu\text{m}$, $E= 2.5 \text{ kV/cm}$,
with $\mu_e= 1350 \text{ cm}^2/\text{V}\cdot\text{s}$ and $\mu_h= 450 \text{ cm}^2/\text{V}\cdot\text{s}$

$$\Rightarrow t_d(e)= 9\text{ns} , t_d(h)= 27\text{ns}$$

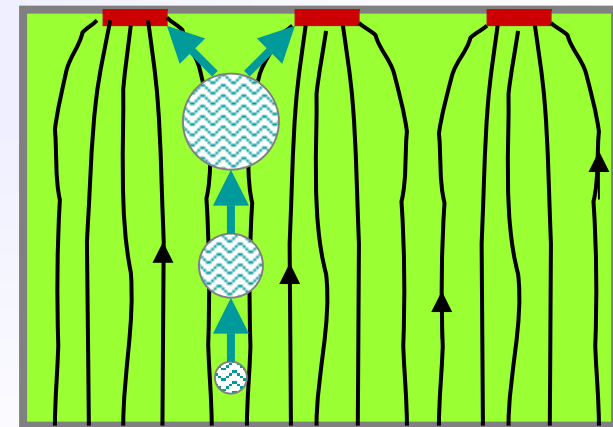
Diffusion

- Diffusion of charge “cloud” caused by scattering of drifting charge carriers, radius of distribution after time td :

$$\sigma = \sqrt{2Dt_d} \quad \text{with diffusion constant } D = \mu kT/q$$

- Same radius for e and h since $t_d \propto 1/\mu$

Typical charge radius: $\sigma \approx 6\mu\text{m}$, could exploit this to get better position resolution due to charge sharing between adjacent strips (using centroid finding), but need to keep drift times long (low field).





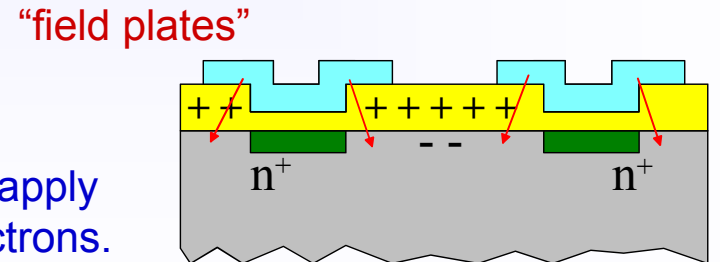
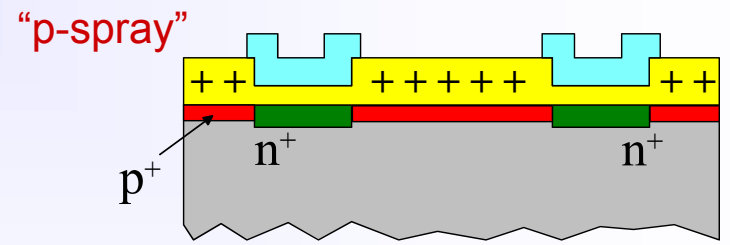
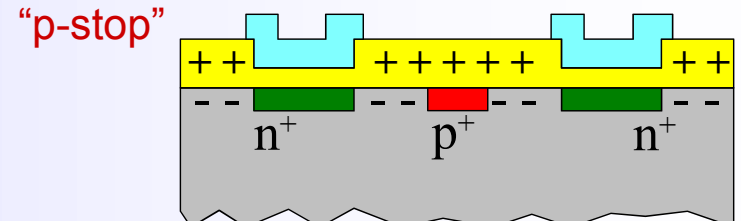
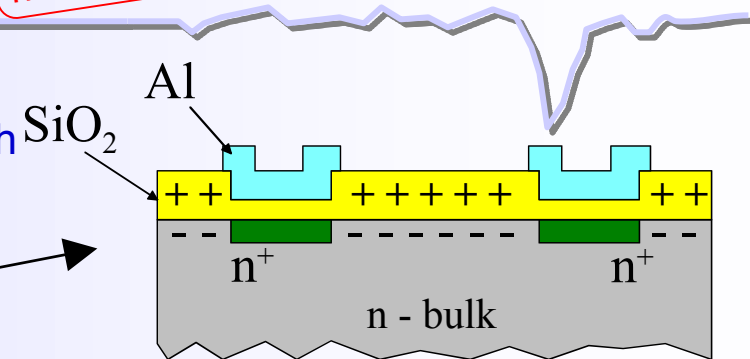
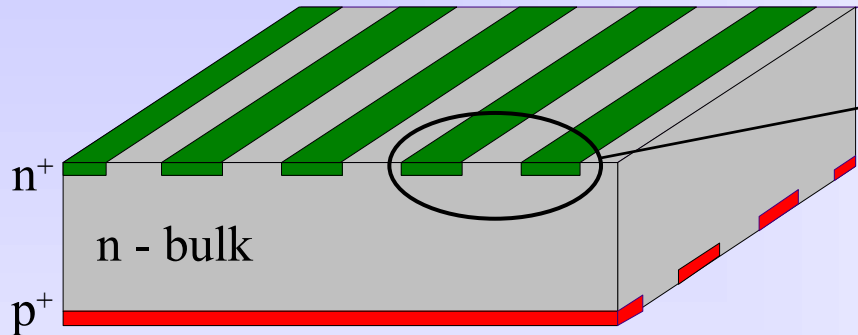
Double sided silicon detectors

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2b - Tracking with Solid State Detectors

■ Get a 2nd coordinate

Put n⁺ and p⁺ strips on opposite sides and read them both



■ Problem: Electron accumulation layer

n⁺-strips are not isolated because of an electron accumulation layer at the Si-SiO₂ interface. This effect is due to the presence of positive charge in SiO₂ layer which attracts electrons.

■ Solution: "Break" accumulation layer

- p-strips in between the n-strips ("**p-stop**")
- moderate p⁺-implantation over all surface ("**p-spray**")
- "**field plates**" (metal over oxide) over the n⁺-strips and apply negative potential with respect to n⁺-strips to repel electrons.

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Detector Module

- **Detector Modules** - “Basic building block of tracking detectors”
 - Silicon Sensors
 - Mechanical support (cooling)
 - Front end electronics and signal routing (connectivity)

■ Example: ATLAS SCT Barrel Module

SCT = SemiConductor Tracker
 ASICS = Application Specific Integrated CircuitS
 TPG = Thermal Pyrolytic Graphite

• Silicon sensors (x4)

- 64 x 64 mm²
- p-in-n, single sided
- AC-coupled
- 768 strips
- 80µm pitch/12µm width

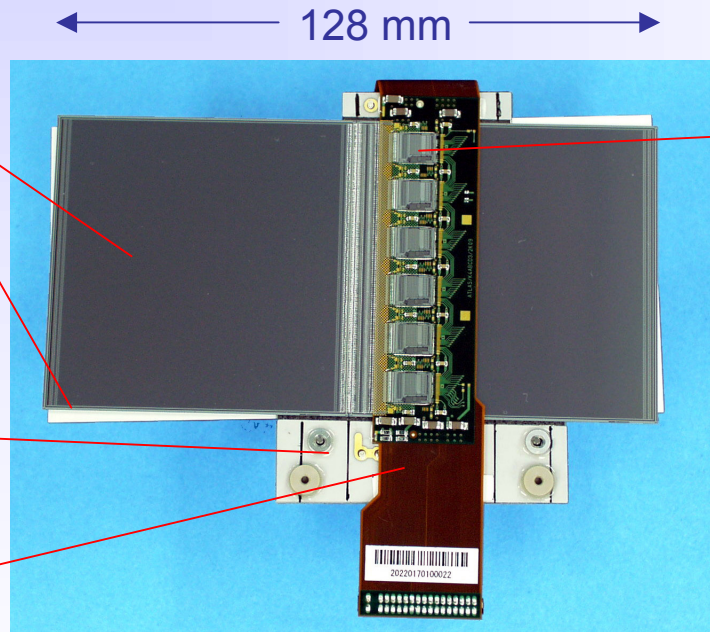
• Mechanical support

- TPG baseboard
- BeO facings

• Hybrid (x1)

- flexible 4 layer copper/kapton hybrid
- mounted directly over two of the four silicon sensors
- carrying front end electronics, pitch adapter, signal routing, connector

$$\sigma(r\phi) \sim 16 \mu\text{m}, \sigma(z) \sim 850\mu\text{m} \text{ [NIMA538 (2005) 384]}$$



• ASICS (x12)

- ABCD chip (binary readout)
- DMILL technology
- 128 channels

• Wire bonds (~3500)

- 25 µm Al wires

■ ATLAS – SCT

- 15.552 microstrip sensors
- 2.112 barrel modules
- 1.976 forward modules
- 61 m² silicon, 6.3·10⁶strips

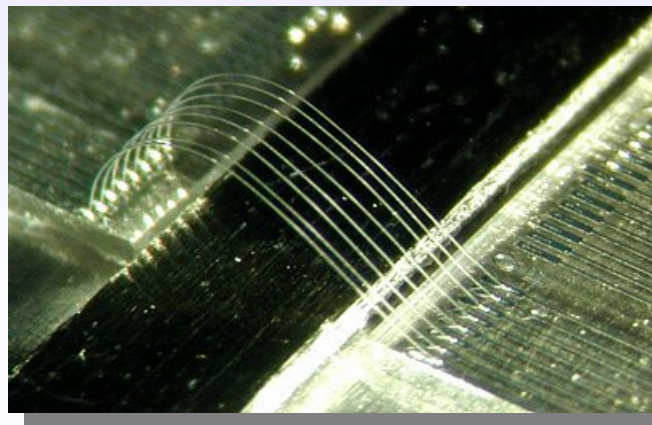
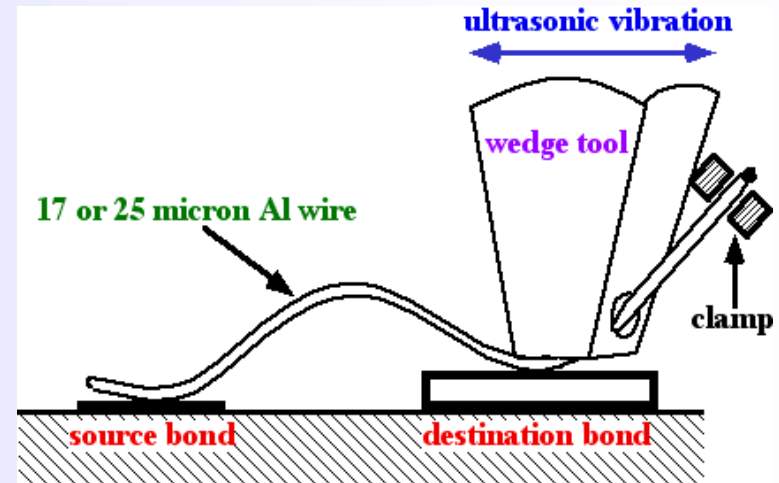


Wire bonding

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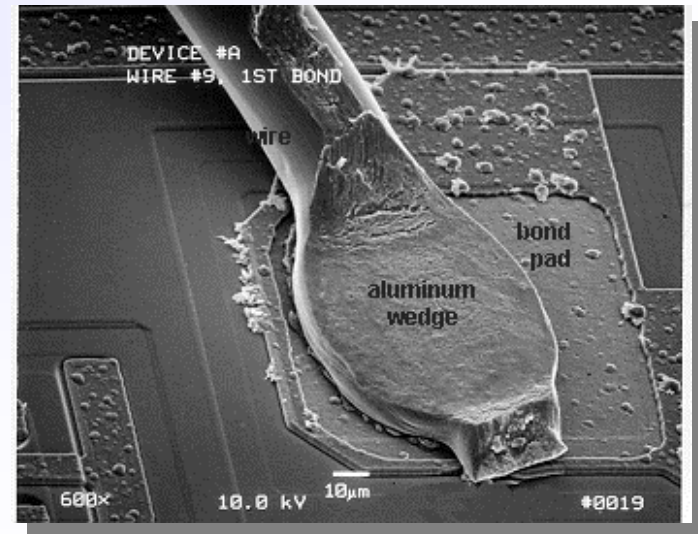
2b - Tracking with Solid State Detectors

- Uses ultrasonic power to vibrate needle-like tool on top of wire. Friction welds wire to metallized substrate underneath.
- Can easily handle 80 μ m pitch in a single row and 40 μ m in two staggered rows (typical FE chip input pitch is 44 μ m).
- Generally use 25 μ m diameter aluminum wire and bond to aluminum pads (chips) or gold pads (hybrid substrates).
- Heavily used in industry (PC processors) but not with such thin wire or small pitch.



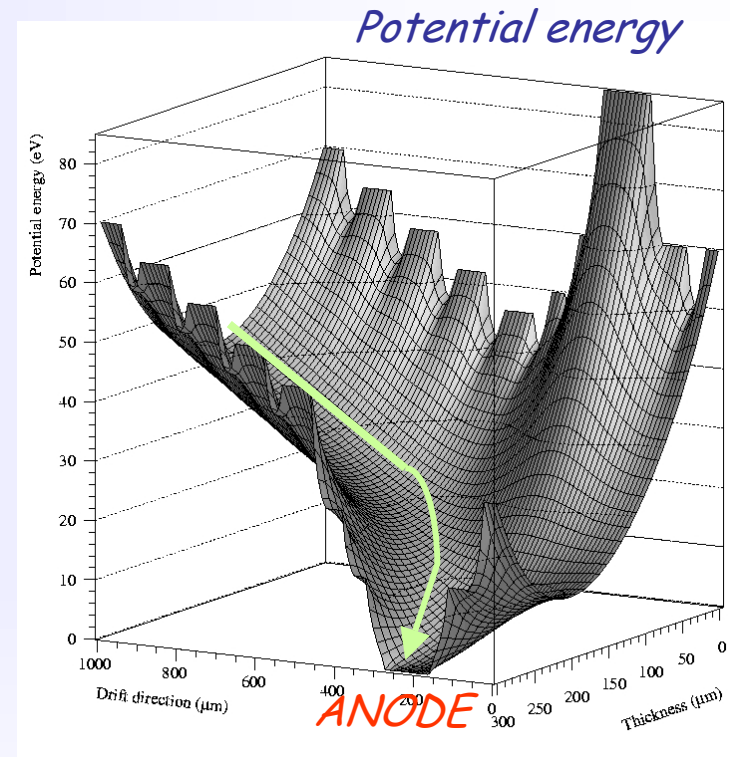
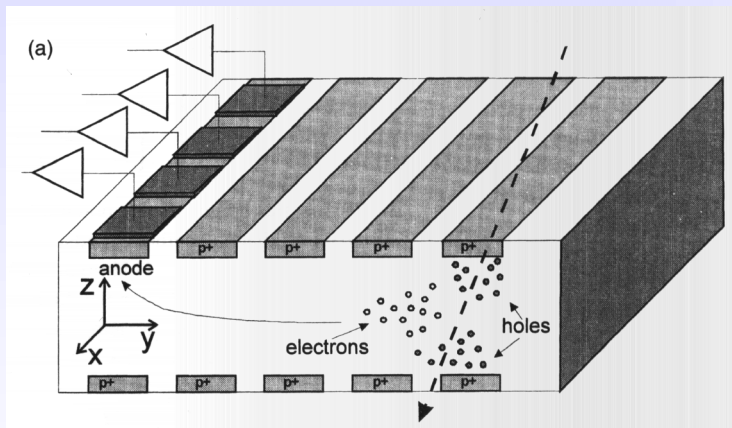
Microscope:
connect sensor to fan-out circuit

Electron microscope:
bond "foot"



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- Principle of sideways depletion (as for DEPFET sensors)
- p^+ segmentation on both sides of sensor
- complete depletion of wafer from segmented n^+ anodes located at one side of sensor
- electrons drift parallel to substrate surface to n^+ anodes
- voltage divider network (resistors) for p-strips to provide uniform drift field



- Need to ensure good material uniformity, low defect rates, good drift field homogeneity, precise voltage dividing on p-strips and good temperature control.
- HEP: Implemented for STAR at RHIC and for ALICE at LHC



Hybrid Pixel Detectors



CERN Globe is a 3 million times bigger bump!

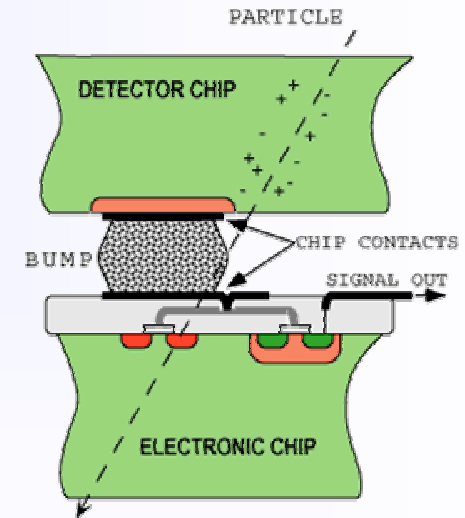
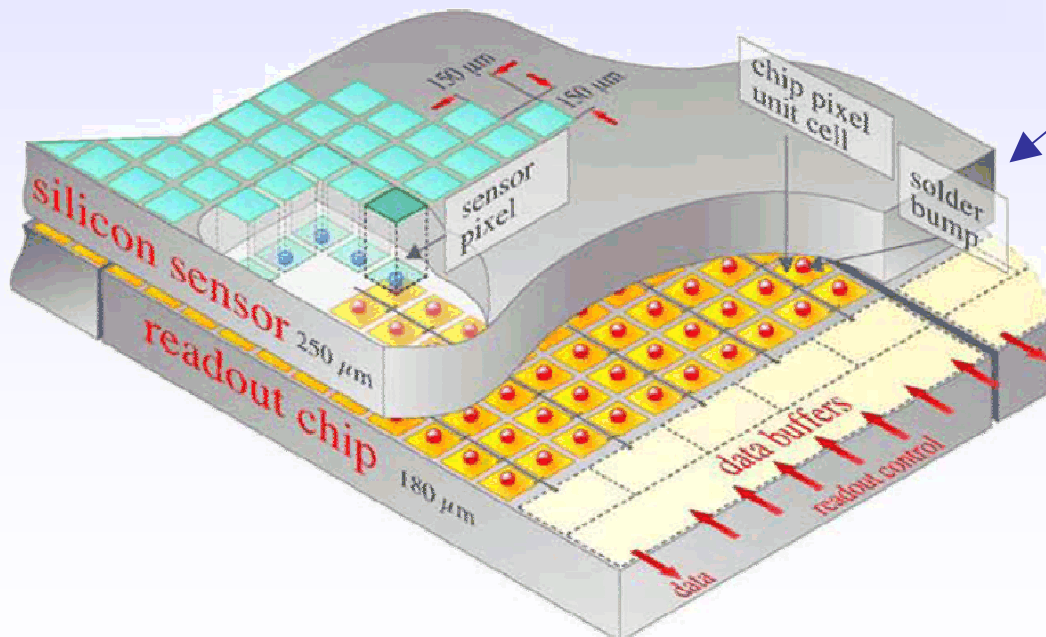
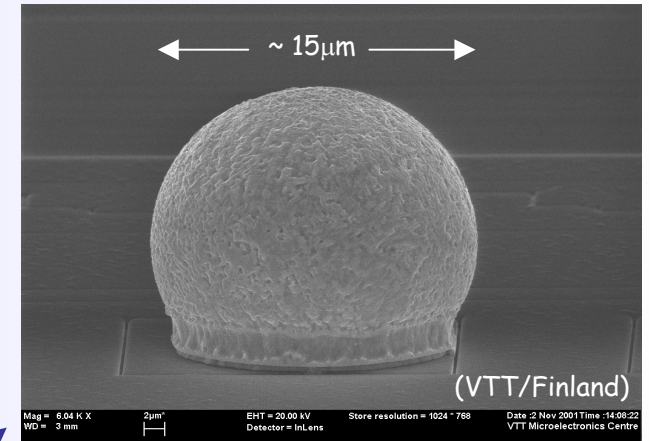


2b - Tracking with Solid State Detectors

■ HAPS – Hybrid Active Pixel Sensors

- segment silicon to diode matrix with high granularity (⇒ true 2D, no reconstruction ambiguity)
- readout electronic with same geometry (every cell connected to its own processing electronics)
- connection by “bump bonding”
- requires sophisticated readout architecture
- Hybrid pixel detectors will be used in LHC experiments: ATLAS, ALICE, CMS and LHCb

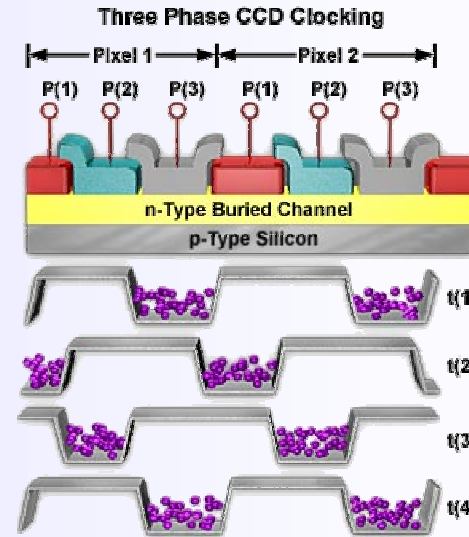
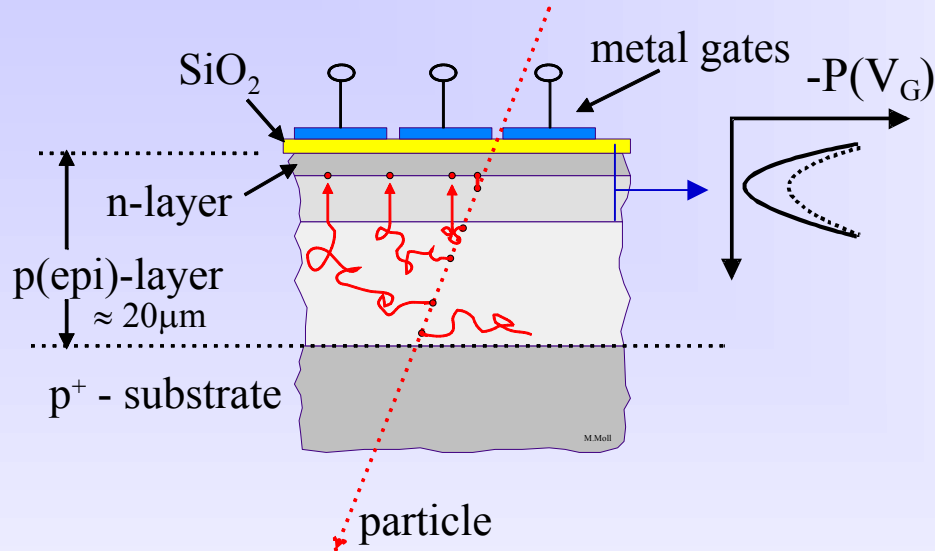
Solder Bump: Pb-Sn



Flip-chip technique

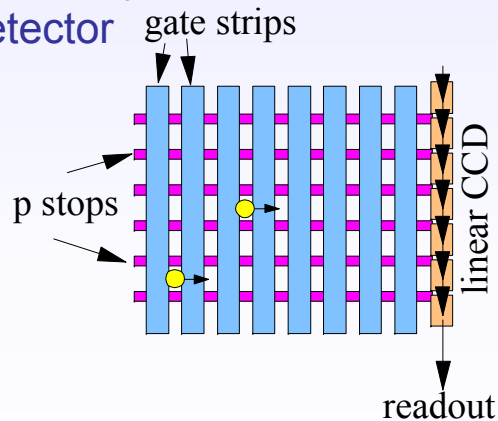
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(1) MOS structure with segmented metal layer;
Charge is captured in a potential well.



(2) Readout: Shift electrons towards anode by periodic variation of 3 potentials

(3) Create an array of pixel for a 2D detector



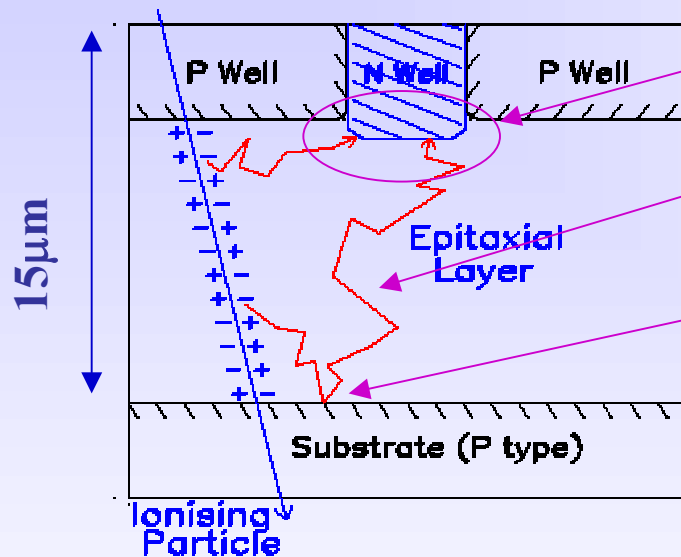
Pixel CCD

- needs only few readout channels
- small charge ($\approx 2000 e$) \Rightarrow needs cooling
- long readout time, active during readout
- sensitive to radiation damage

\Rightarrow applicable for low rate experiment without high intensity radiation field

Monolithic detectors

- readout electronics directly within sensor material (same epi layer)



- charge collected at n-well / p-epi diode
- thermal diffusion of free charge
- reflection at potential barriers between areas with different doping concentration
- no depletion voltage applied
⇒ potential formed by different doping concentrations only

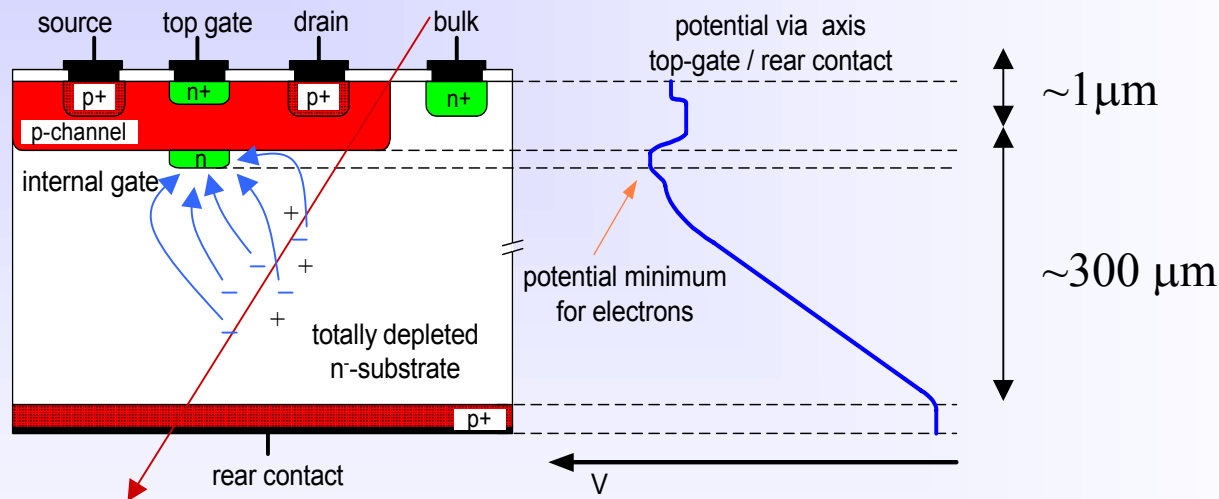
- no connections needed to electronics (e.g. no bumps)
- very small sizes achievable



DEPFET - DEP(leted)F(ield)E(ffect)T(ransistor)

2b - Tracking with Solid State Detectors
extra slide not shown

- FET integrated on high resistivity bulk, bulk sideward depleted
- electrons collected in potential minimum at internal gate
 - transistor current modulated by collected charge
 - charge removed by reset mechanism (clear)
- switch on/off by (external) top gate to read out



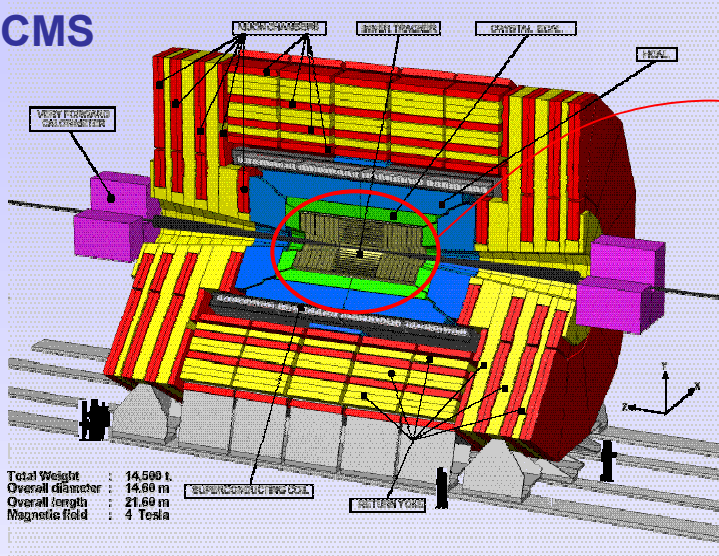
- amplification of charge at the position of collection \Rightarrow no transfer loss
- full bulk sensitivity, bulk can be thinned down to 50 μm if needed
- non structured entrance window (backside)
- very low input capacitance \Rightarrow very low noise



Example from LHC: The CMS tracker

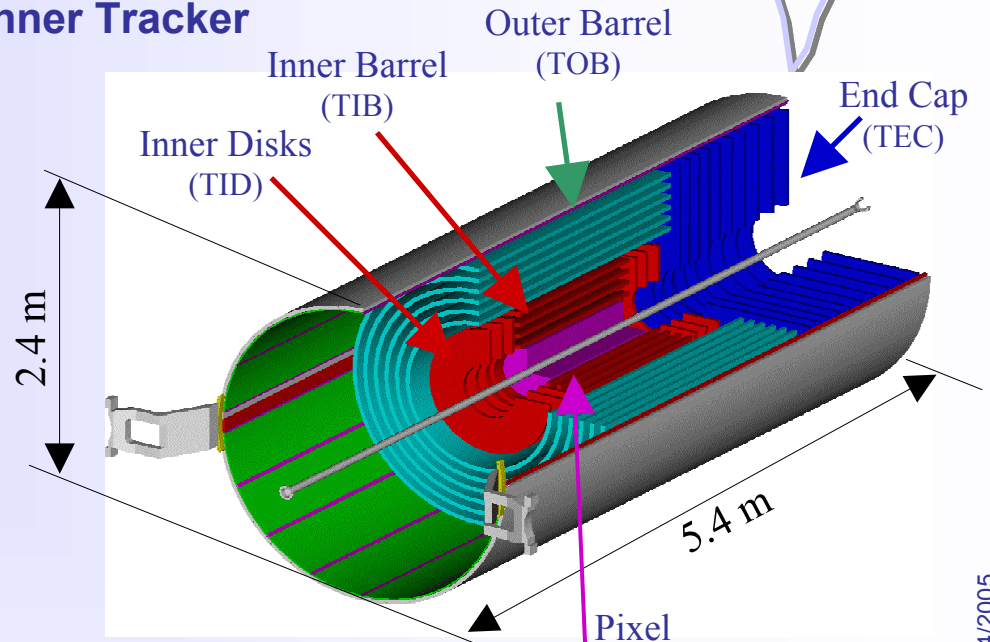
2b - Tracking with Solid State Detectors

CMS



Total Weight : 14,500 t
Overall diameter : 14.00 m
Overall length : 21.00 m
Magnetic field : 4 Tesla

Inner Tracker



CMS - Currently the Most Silicon

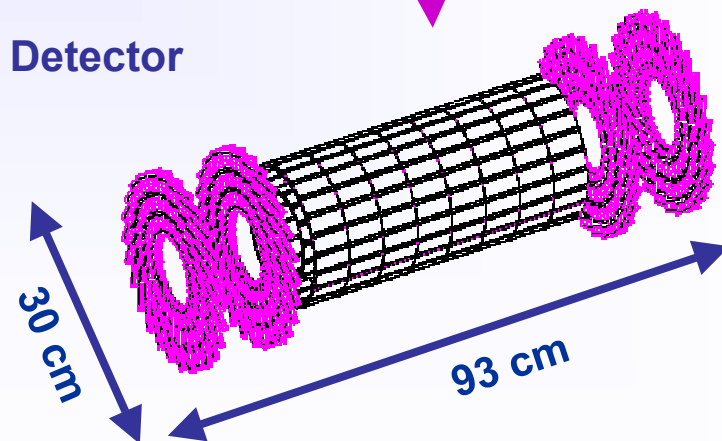
Micro Strip:

- ~ 214 m² of silicon strip sensors
- 11.4 million strips

Pixel:

- Inner 3 layers: silicon pixels (~ 1m²)
- 66 million pixels (100x150μm)
- Precision: $\sigma(r\phi) \sim \sigma(z) \sim 15\mu\text{m}$
- Most challenging operating environments (LHC)

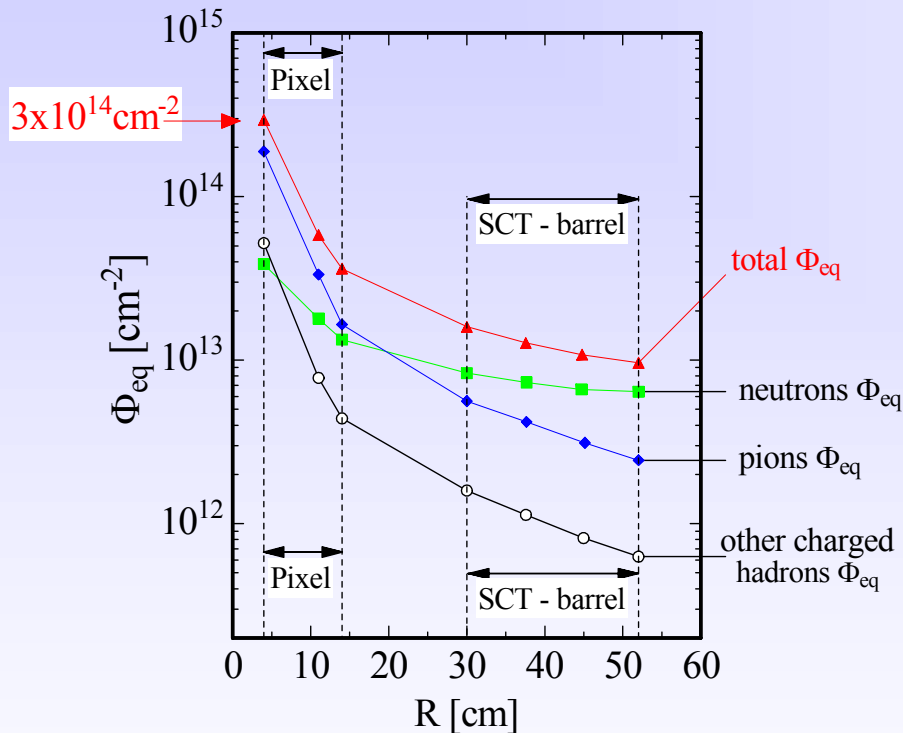
Pixel Detector





Example: ATLAS

- Fluences per year at full Luminosity



- Pixel detector: up to $\Phi_{eq} \approx 3.5 \cdot 10^{14} \text{ cm}^{-2} / \text{year}$
- Dominating type of particle is different for pixel (pions) and strip detectors (neutrons)

LHC silicon detectors:

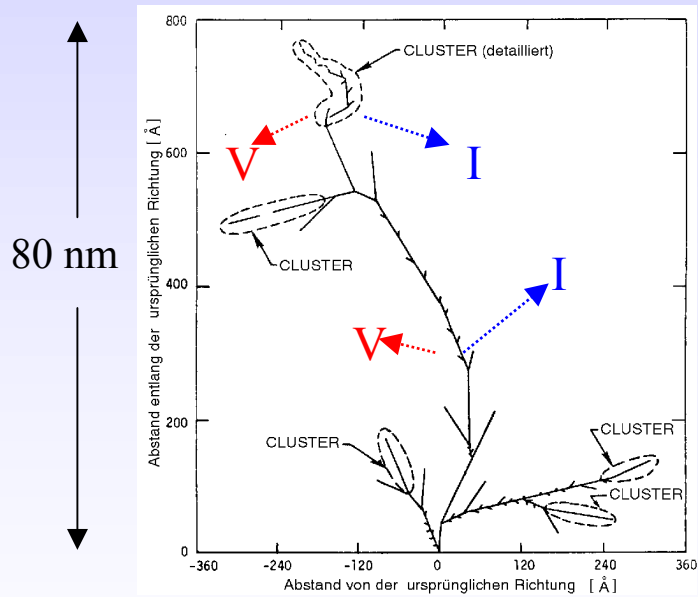
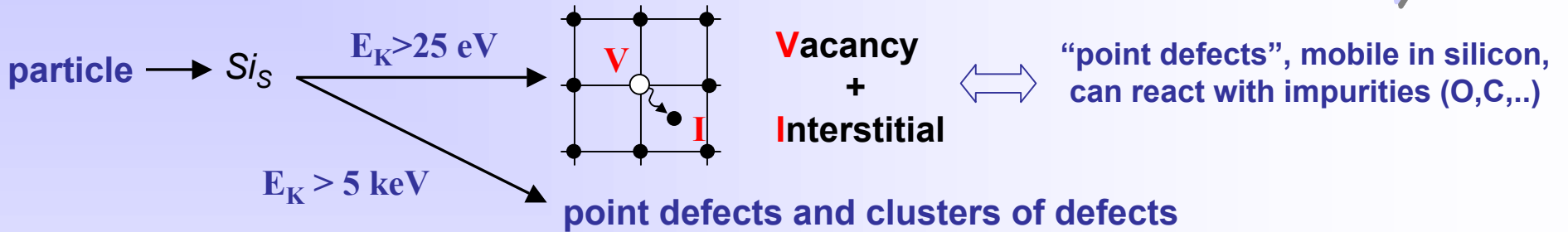
- All detectors have been extensively tested and developed for radiation tolerance and are expected to survive the LHC radiation environment.
- Some experiments have already foreseen upgrades (e.g. LHCb Velo after 3 years).

Super LHC

- upgrade of LHC to 10 x higher Luminosity
 - ⇒ 10 x higher radiation levels
 - ⇒ Radiation damage will become a critical issue!
- ⇒ New, radiation tolerant detectors needed!

- What is radiation damage ?
- How to cope with it ?

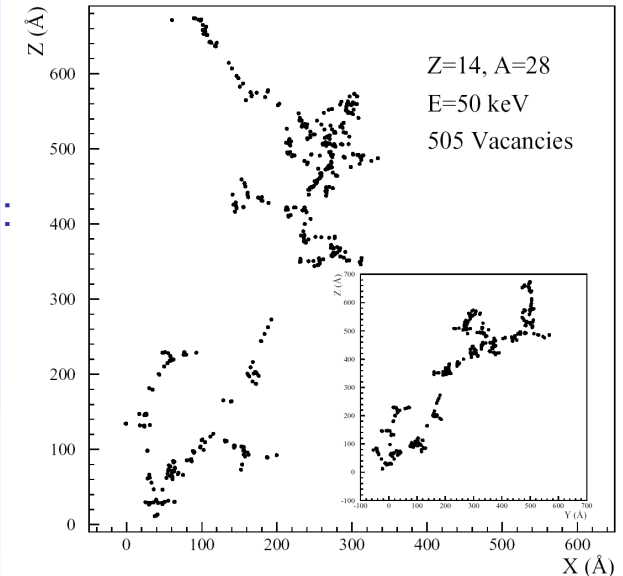
Damage to the silicon crystal: Displacement of lattice atoms



Distribution of vacancies created by a 50 keV Si-ion in silicon (typical recoil energy for 1 MeV neutrons):

← **Schematic** [Van Lint 1980]

Simulation [M.Huhtinen 2001] →



Defects can be electrically active (levels in the band gap)

- capture and release electrons and holes from conduction and valence band

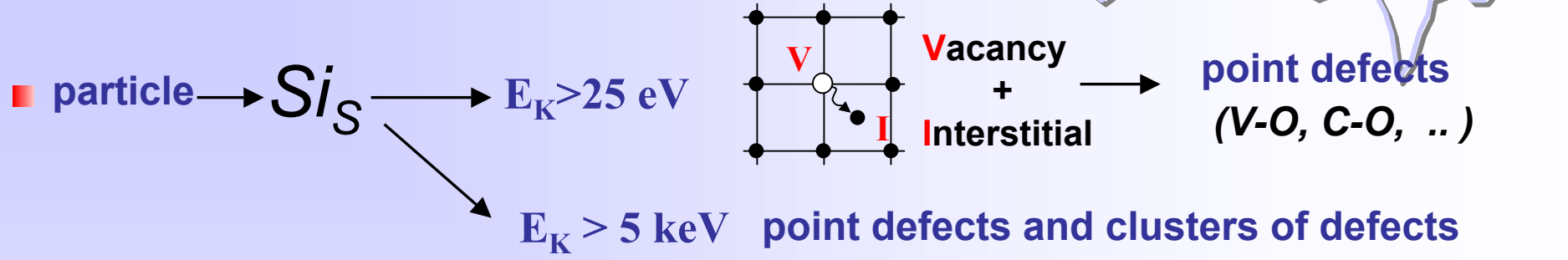
⇒ can be charged - can be generation/recombination centers - can be trapping centers



Radiation Damage: Particle dependence

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2b - Tracking with Solid State Detectors



^{60}Co -gammas
Compton Electrons
with max. $E_\gamma \approx \text{MeV}$
(no cluster production)

Electrons
 $E_e > 255 \text{ keV}$ for displacement
 $E_e > 8 \text{ MeV}$ for cluster

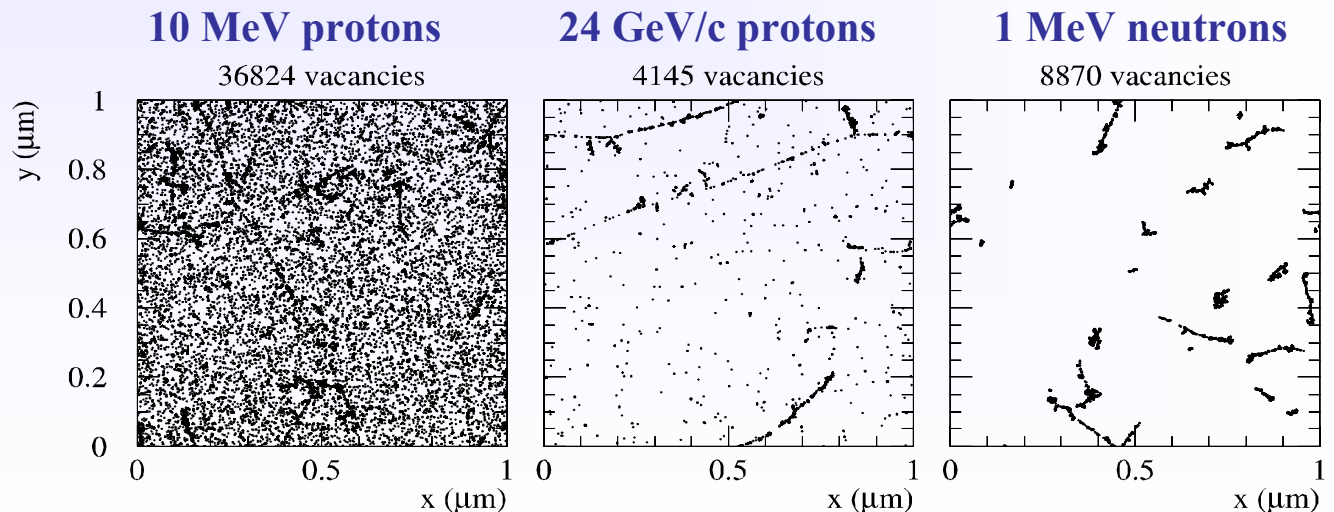
Neutrons (elastic scattering)
 $E_n > 185 \text{ eV}$ for displacement
 $E_n > 35 \text{ keV}$ for cluster

only point defects ↔ **point defects & clusters** ↔ **mainly clusters**

■ Simulation:

Initial distribution of vacancies in $(1\mu\text{m})^3$ after 10^{14} particles/cm 2

[Mika Huhtinen NIMA 491(2002) 194]



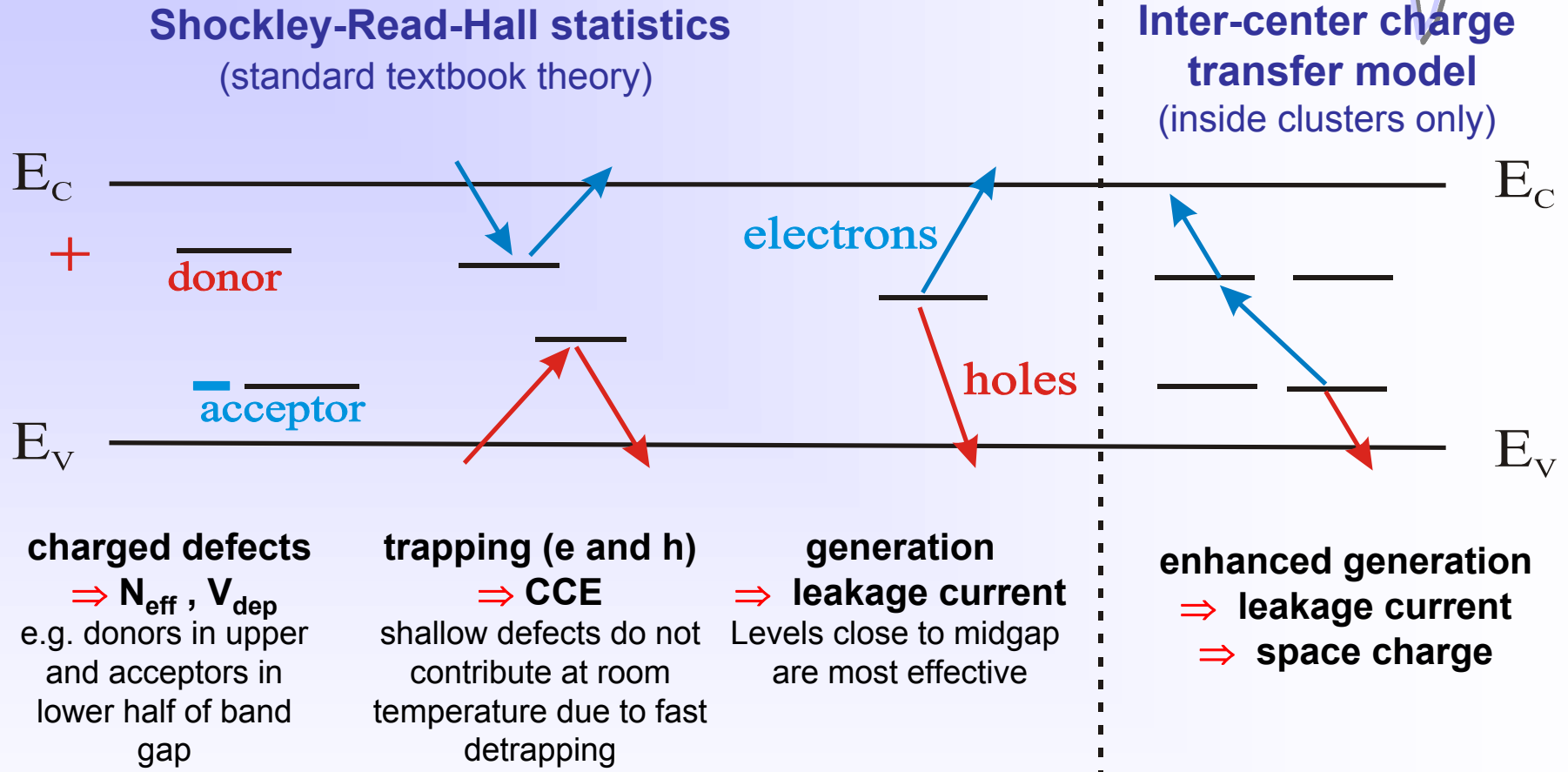
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Impact of defects on detector properties

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2b - Tracking with Solid State Detectors



■ Impact on detector properties can be calculated if all defect parameters are known:

$\sigma_{n,p}$: cross sections

ΔE : ionization energy

N_t : concentration



Radiation Damage in Silicon Sensors

2b - Tracking with
Solid State Detectors

■ Two general types of radiation damage:

- **Bulk (Crystal) damage** due to **Non Ionizing Energy Loss (NIEL)**
- Displacement Damage –

- I. Change of **depletion voltage** (higher operation voltage, underdepletion)
⇒ constant cooling needed to avoid reverse annealing
- II. Increase of **leakage current** (increase of shot noise, thermal runaway)
⇒ needs cooling of sensors during operation
- III. Decrease of **charge collection efficiency**
due to underdepletion and increased trapping

- **Surface damage** due to **Ionizing Energy Loss (IEL)**
- accumulation of positive in the oxide (SiO_2) and the Si/SiO_2 interface –

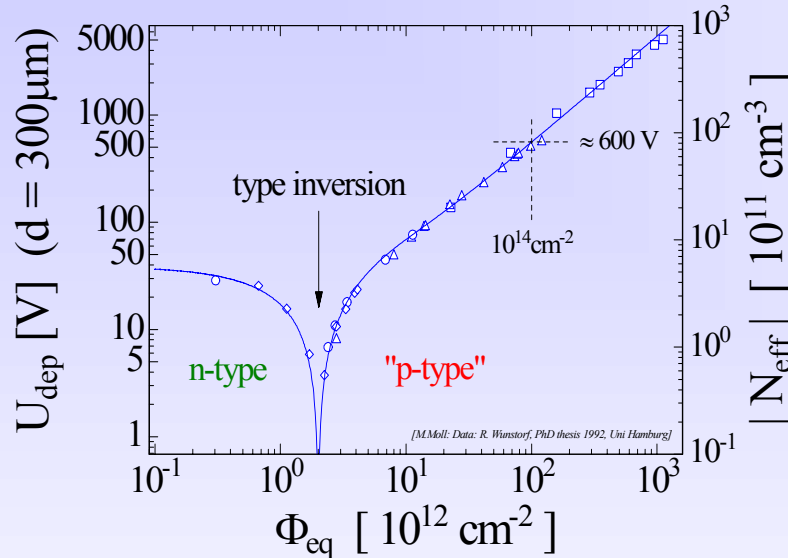
affects: interstrip capacitance (noise factor), breakdown behavior
and other structures depending on near-surface effects

■ Signal/noise ratio is the quantity to watch

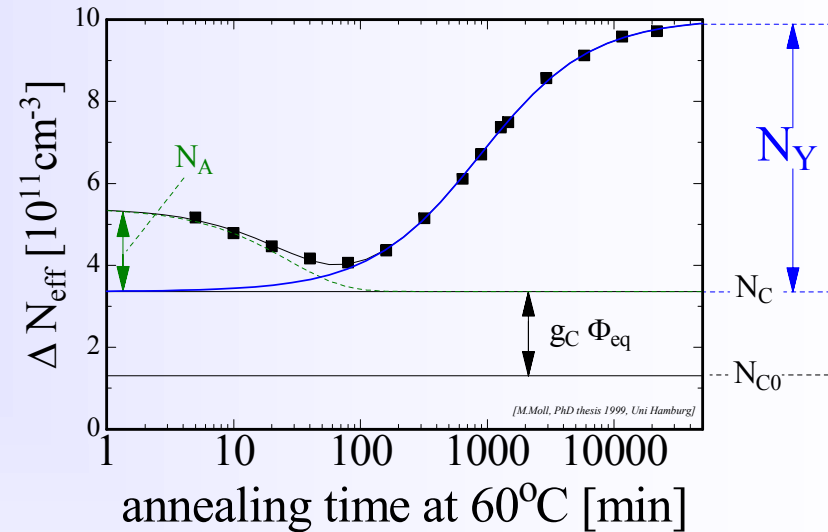
⇒ **Sensors can fail from radiation damage !**

Change of Depletion Voltage $V_{\text{dep}} (N_{\text{eff}})$

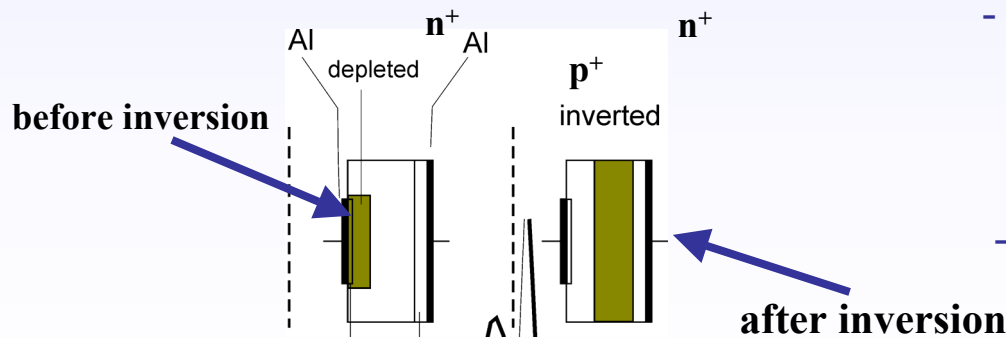
.... with particle fluence:



.... with time (annealing):



- “Type inversion”: N_{eff} changes from positive to negative (Space Charge Sign Inversion)



- Short term: “Beneficial annealing”
- Long term: “Reverse annealing”
- time constant depends on temperature:
 - ~ 500 years (-10°C)
 - ~ 500 days (20°C)
 - ~ 21 hours (60°C)
- Consequence: Detectors must be cooled even when the experiment is not running!

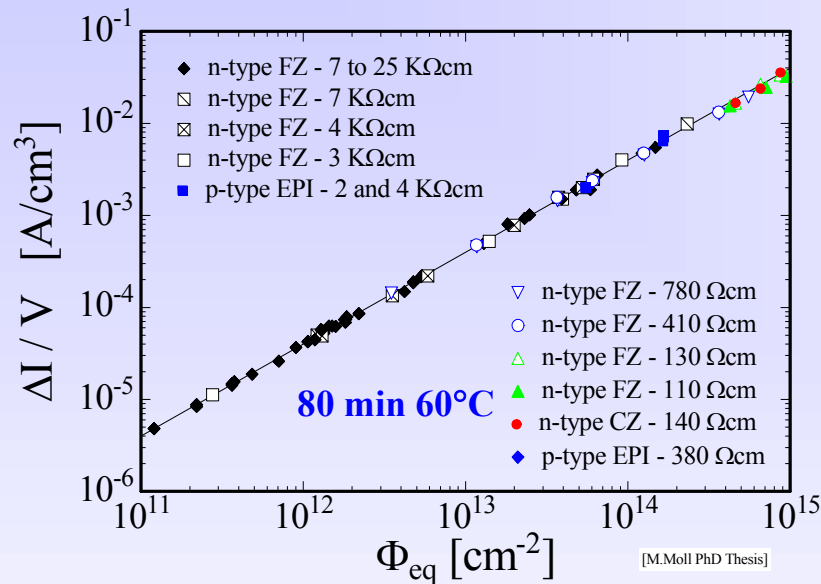


Radiation Damage – II. Leakage Current

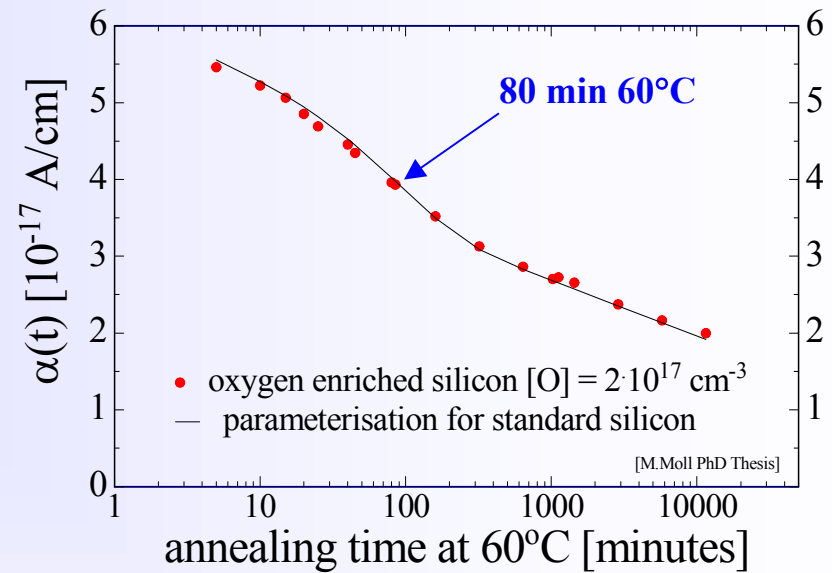
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2b - Tracking with Solid State Detectors

Change of Leakage Current (after hadron irradiation) with particle fluence:



.... with time (annealing):



• Damage parameter α (slope in figure)

$$\alpha = \frac{\Delta I}{V \cdot \Phi_{eq}}$$

Leakage current per unit volume and particle fluence

- α is constant over several orders of fluence and independent of impurity concentration in Si
 \Rightarrow can be used for fluence measurement

- Leakage current decreasing in time (depending on temperature)
- Strong temperature dependence

$$I \propto \exp\left(-\frac{E_g}{2k_B T}\right)$$

Consequence:

Cool detectors during operation!
 Example: $I(-10^\circ\text{C}) \sim 1/16 I(20^\circ\text{C})$



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Deterioration of Charge Collection Efficiency (CCE)

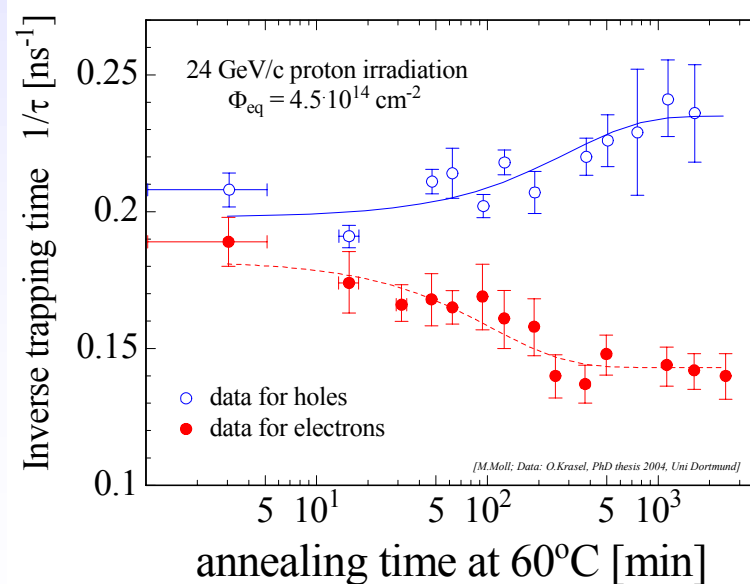
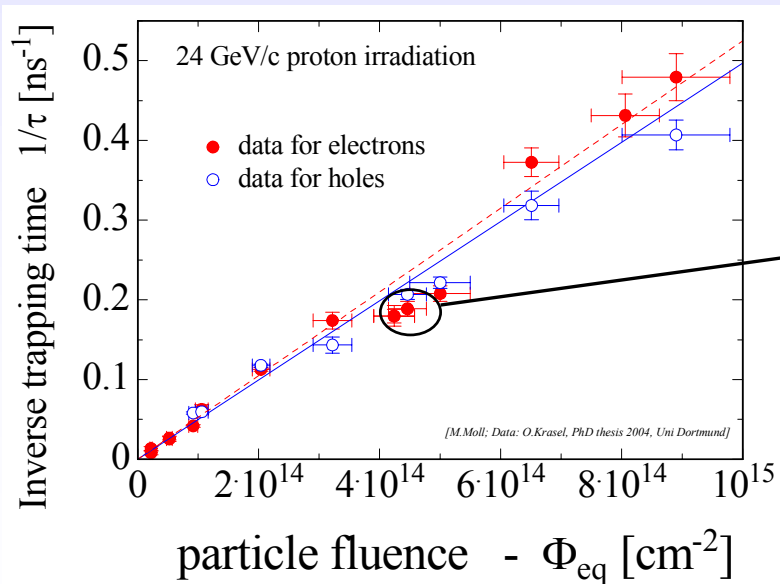
2 mechanisms: - **Trapping** of electrons and holes

- **Underdepletion** (loss of active detector volume due to increase of V_{dep})

Trapping is characterized by an effective trapping time τ_{eff} for electrons and holes:

$$Q_{e,h}(t) = Q_{0e,h} \exp\left(-\frac{1}{\tau_{eff\ e,h}} \cdot t\right) \quad \text{where} \quad \frac{1}{\tau_{eff\ e,h}} \propto N_{defects}$$

Increase of inverse trapping time ($1/\tau$) with fluence and change with time (annealing):





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Scientific Strategies:

I. Material engineering

II. Device engineering

III. Variation of detector operational conditions

■ Defect Engineering of Silicon

Deliberate incorporation of impurities or defects into the silicon bulk to improve radiation tolerance of detectors

- **Needs:** profound understanding of radiation damage (*microscopic defects, macroscopic parameters, dependence on particles type and energy, defect formation kinetics and annealing processes*)
- **Examples:**
 - Oxygen enriched silicon
 - Hydrogen enriched silicon
 - Pre-irradiated silicon

■ New Materials (other semiconductors than Si)

- Diamond, Silicon Carbide (SiC), ...

■ New detector designs

- **Examples:**
 - p-type silicon detectors (n-in-p)
 - thin detectors, epitaxial detectors
 - 3D and Semi 3D detectors

■ Cryogenic operation of detectors

Operate detectors at 100-200K to reduce the charge loss ("Lazarus effect")

Active CERN R&D collaborations:

- *RD50 "Radiation hard semiconductor devices for very high luminosity colliders"*
- *RD42 "CVD Diamond Radiation Detectors"*
- *RD39 "Cryogenic Tracking Detectors"*



New Material: Oxygen enriched silicon – DOFZ

2b - Tracking with Solid State Detectors

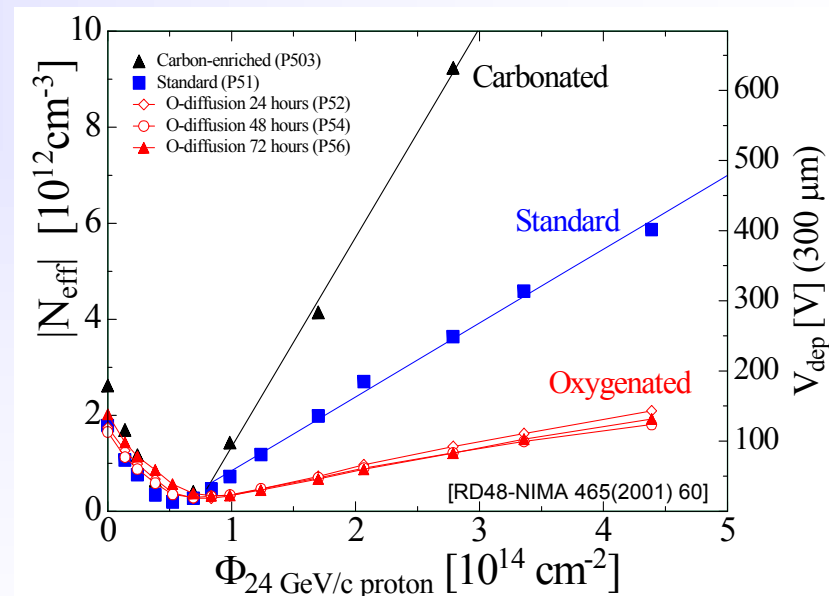
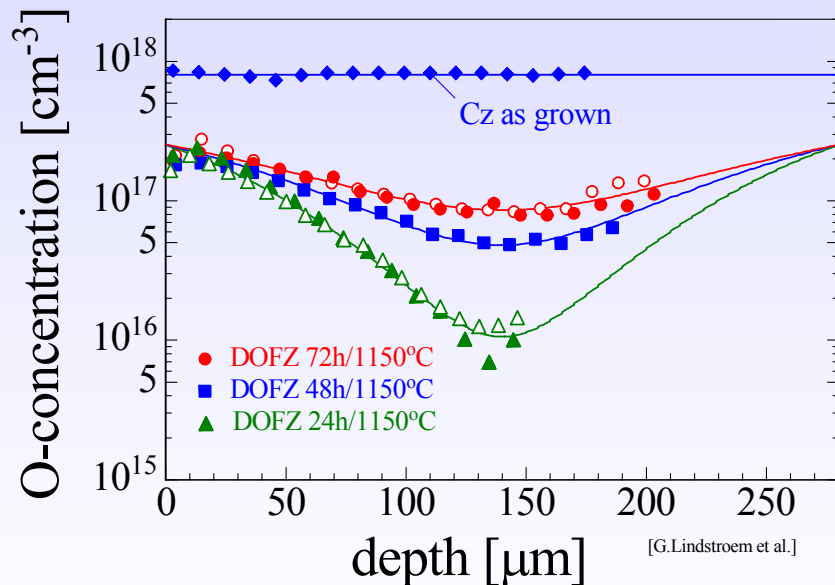
DOFZ (Diffusion Oxygenated Float Zone Silicon)

- 1982 First oxygen diffusion tests on FZ [Brotherton et al. J.Appl.Phys.,Vol.53, No.8.,5720]
- 1995 First tests on detector grade silicon [Z.Li et al. IEEE TNS Vol.42,No.4,219]
- 1999 Introduced to the HEP community by CERN - RD48 (ROSE-Collaboration)



Very long oxidation (e.g. 48h at 1150°C) increases the oxygen content in silicon

Strong improvement after charged hadron irradiation observed

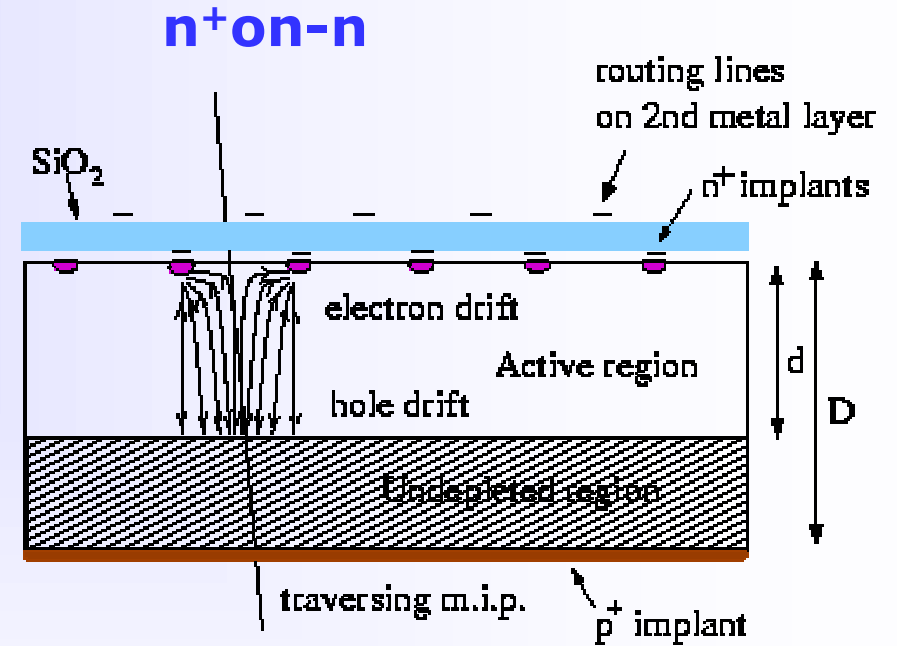
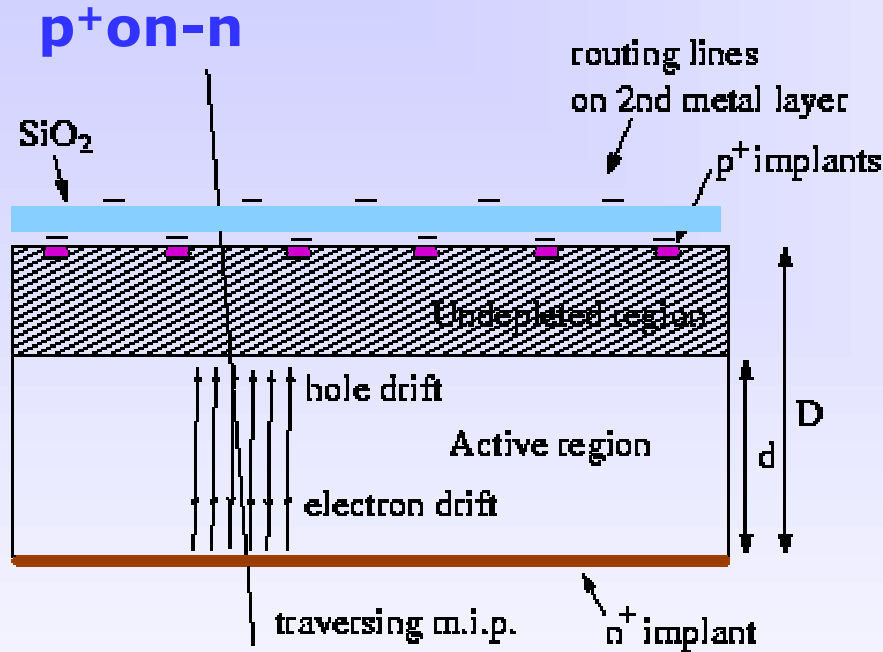


- 2005: DOFZ silicon used for the ATLAS and CMS Pixel detectors
- 2005: Other types of oxygen rich silicon under investigation: Czochralski Si, epitaxial Si

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n-type silicon after type inversion:



p-on-n silicon, under-depleted:

- Charge spread – degraded resolution
- Charge loss – reduced CCE

n-on-n silicon, under-depleted:

- Limited loss in CCE
- Less degradation with under-depletion
- Collect electrons (fast)

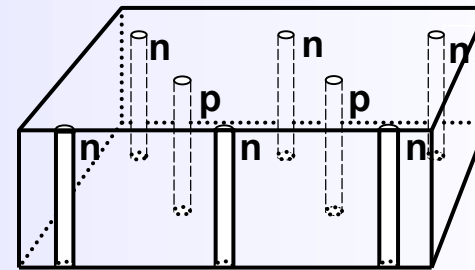


New detector concepts: 3D detectors

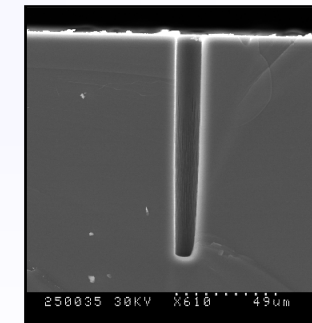
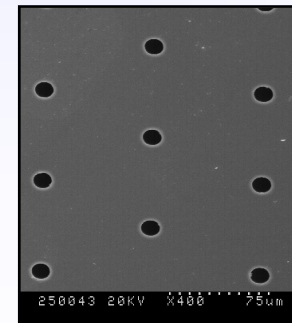
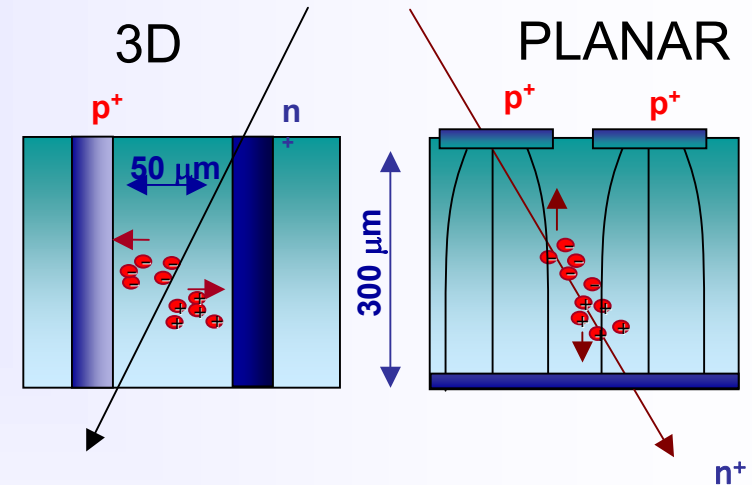
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2b - Tracking with Solid State Detectors

- **Electrodes:**
 - narrow columns along detector thickness-“3D”
 - diameter: $\approx 10\mu\text{m}$; distance: 50 - 100 μm
- **Lateral depletion:**
 - lower depletion voltage needed ($V_{\text{dep}} \sim d^2$)
 - radiation tolerant or thick detectors possible
 - fast signal (≈ 3.5 ns measured)
- **Processing of detectors:**
 - complex fabrication: Holes have to be made and filled with electrodes (DRIE etching, Laser drilling, Photo Electro Chemical etching); present aspect ratio (depth to diameter) $\approx 30:1$
 - possibility to implement narrow dead regions at edges “edgeless detectors”
- **Application:**
 - detectors still under development!
 - option for LHC experiments upgrade ?



[Proposed: S.I. Parker et al., NIMA 395 (1997) 328]



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Besides references given on the transparencies, the following sources have been used:

■ Books

- **Gerhard Lutz**, “**Semiconductor Radiation Detectors**”, Springer, ISBN 3-540-64859-3
- S.M.Sze, “Physics of Semiconductor Devices”, John Wiley & Sons, ISBN 0-471-05661-8

■ Articles

- Anna Peisert, “Silicon microstrip detectors”, Instrumentation in High Energy Physics, World Scientific, 1992
- Michael Moll, “Radiation Damage in Silicon Particle Detectors”, PhD thesis, DESY, December 1999
- Geoffrey Hall, “Semiconductor particle tracking detectors”, Rep.Prog.Phys. 57 (1994) 481-531

■ Lectures and Presentations

- **Alan Honma**, “**Silicon Detectors**”, Nato Advanced Study Institute, Virgin Islands, 06/2002, <http://cern.ch/honma/>
- **Christian Joram**, “**Particle Detectors**”, CERN, Summer Student Lectures June 2003
- Paula Collins, “Recent Detector R&D and operational experience”, IWORID07, Riga, September 2003
- Gerhard Lutz, “Semiconductor Radiation Detectors”, Louvain, Seminar, June 2002
- Marcello Mannelli, “Tracking at the LHC: The CMS example”, CERN Academic Training, March 2005
- Pierre Jarron, “Microelectronics, Nanoelectronics, Monolithic Pixel Detectors” CERN Academic Training, Jan. 2004
- Hans Dijkstra, “Overview of Silicon Detectors”, Vienna conference, VCI 2001
- Volker Adler, “The TESLA Vertex Detector” ZEUS Student Seminar, Jan.2004
- Daniela Bortoletto, “An introduction to semiconductor detectors”, Vienna conference, VCI 2004
- A list of conferences about Solid State Detectors and Radiation Damage: <http://cern.ch/mmoll/links/conferences.htm>
- Vertex 2004 conference: <http://sucimaweb.dipscfm.uninsubria.it/vertex04/>