

ASACUSA

Atomic Spectroscopy And Collisions Using Slow Antiprotons

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ASACUSA

Atomic Spectroscopy And Collisions Using Slow Antiprotons

Goal

Test CPT (matter-antimatter symmetry) to the highest-possible precision -

How ?

'Weigh' the antiproton (proton mass = antiproton mass?)

Use the **antiprotonic helium atom** as our 'scale' - so far -

Collaborating institutes and funding



| | | |
|---|------------------|--|
|  | Tokyo RIKEN | MEXT, Japan RIKEN |
|  | Aarhus | Danish natural science foundation, ISA |
|  | RMKI Debrecen | OMFB TeT OTKA |
|  | CERN | |
| MoU to be signed | | |



Part 1

ASACUSA 2000 - 2004

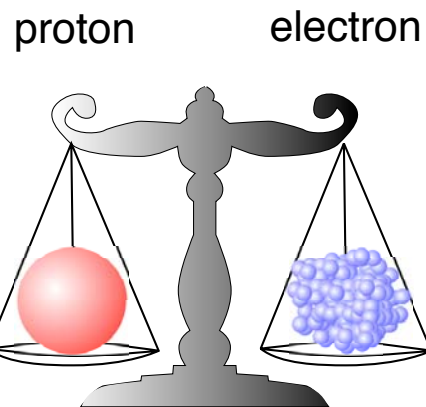
CPT 'theorem'



- ≡ CPT 'theorem' -
Physics laws unchanged by the simultaneous exchange of
C(particle \leftrightarrow antiparticle), P(left \leftrightarrow right), T(future \leftrightarrow past)
- ≡ If CPT is OK, **particle mass = antiparticle mass**
- ≡ Extensions of the standard model accommodate CPT violation.
- ≡ CPT violation, if discovered, has a profound impact on the basic understanding of nature (but the magnitude must be very small)
- ≡ CPT must be experimentally tested to the highest-possible precision
- ≡ ASACUSA compares **proton mass vs antiproton mass**

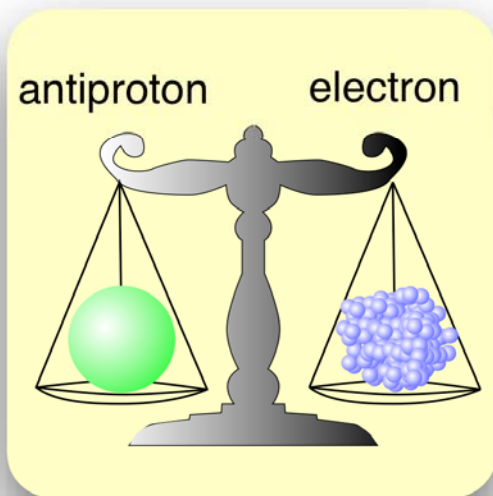
The best baryonic CPT test

Weighing the antiproton



≡ proton-electron mass ratio is known to high precision

$$m_p/m_e = 1836.15267261 \pm 0.00000085 (4.6 \times 10^{-10})$$



≡ ASACUSA measures antiproton-electron mass ratio, (and then compares it with the proton-electron ratio) but how?

Hydrogen spectrum and electron mass



Famous Balmer lines of hydrogen

The photon frequency is

$$\nu(n, n') = Rc \left(\frac{1}{n^2} - \frac{1}{n'^2} \right)$$

$$R = \frac{m_e e^4}{8c\epsilon_0^2 h^3}$$

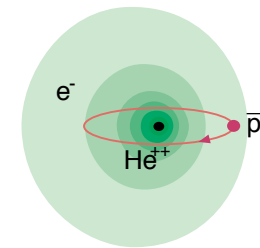
where R is the Rydberg constant.

Antiprotonic helium

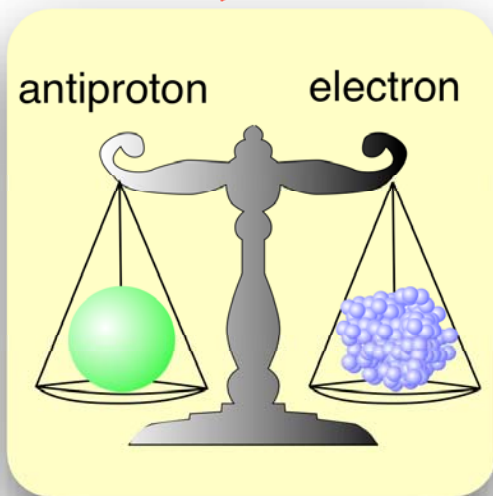


For antiprotonic helium ($\bar{p} + e^- + \alpha$),

$$\nu(n, n') = R\bar{c} \frac{M_{\bar{p}}}{m_e} Z_{\text{eff}}^2 \left(\frac{1}{n^2} - \frac{1}{n'^2} \right)$$



Z_{eff}^2 : helium charge, shielded by electron (calculated by theory)



≡ **Rc** is known to an astounding precision of 6.6×10^{-12} .

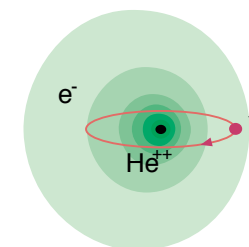
≡ By measuring ν , antiproton mass can be determined.

(to be more precise)

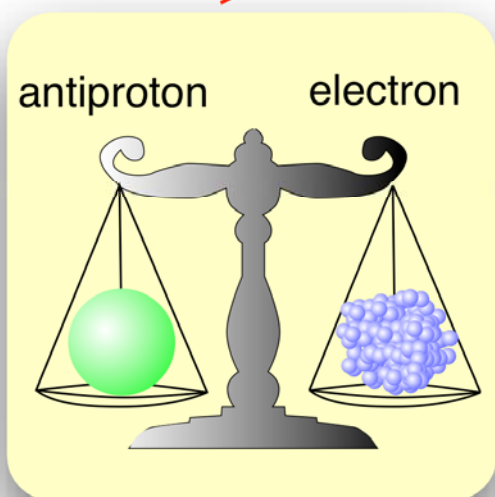


For antiprotonic helium ($\bar{p} + e^- + \alpha$),

$$\nu(n, \ell; n', \ell') = Rc \frac{M_{\bar{p}}^*}{m_e} Z_{\text{eff}}^2 \frac{Q_{\bar{p}}^2}{e^2} \left(\frac{1}{n^2} - \frac{1}{n'^2} \right)$$



$Z_{\text{eff}}^2(n, \ell; n', \ell')$: state-dependent helium charge, shielded by electron (calculated by 3-body QED)



≡ By measuring ν , and by combining it with $Q_{\bar{p}}/M_{\bar{p}}$ measured by the TRAP group (at LEAR), antiproton mass and charge can be determined (PDG)

Q/M by TRAP @ LEAR



TRAP III (Phys.Rev.Lett. **82** (1999) 3198)

$$\frac{Q_{\bar{p}}/M_{\bar{p}}}{Q_p/M_p} + 1 = 0.9(9) \times 10^{-10}$$

Note: A polarization force shifts the H^- ion's cyclotron frequency (Nature **430**, 58 (2004)); a corrected value is $1.6(9) \times 10^{-10}$

Best baryonic CPT limit



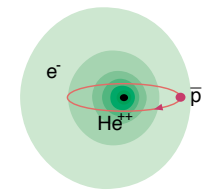
≡ Antiproton 'weighed' to 10^{-8} precision using the antiprotonic helium laser spectroscopy - the **best** (baryonic) CPT limit.

From The Review of Particle Physics (2004)

$$\left| \frac{m_p - m_{\bar{p}}}{m_p} \right|$$

A test of *CPT* invariance. Note that the comparison of the \bar{p} and p charge-to-mass ratio, given in the next data block, is much better determined.

| VALUE | CL% | DOCUMENT ID | TECN | COMMENT |
|---|-----|--------------------|------|--|
| $<1.0 \times 10^{-8}$ | 90 | ¹ HORI | 03 | SPEC $\bar{p}e^{-} {}^4\text{He}$ and $\bar{p}e^{-} {}^3\text{He}$ |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | |
| $<6 \times 10^{-8}$ | 90 | ¹ HORI | 01 | SPEC $\bar{p}e^{-} \text{He}$ atom |
| $<5 \times 10^{-7}$ | | ² TORII | 99 | SPEC $\bar{p}e^{-} \text{He}$ atom |



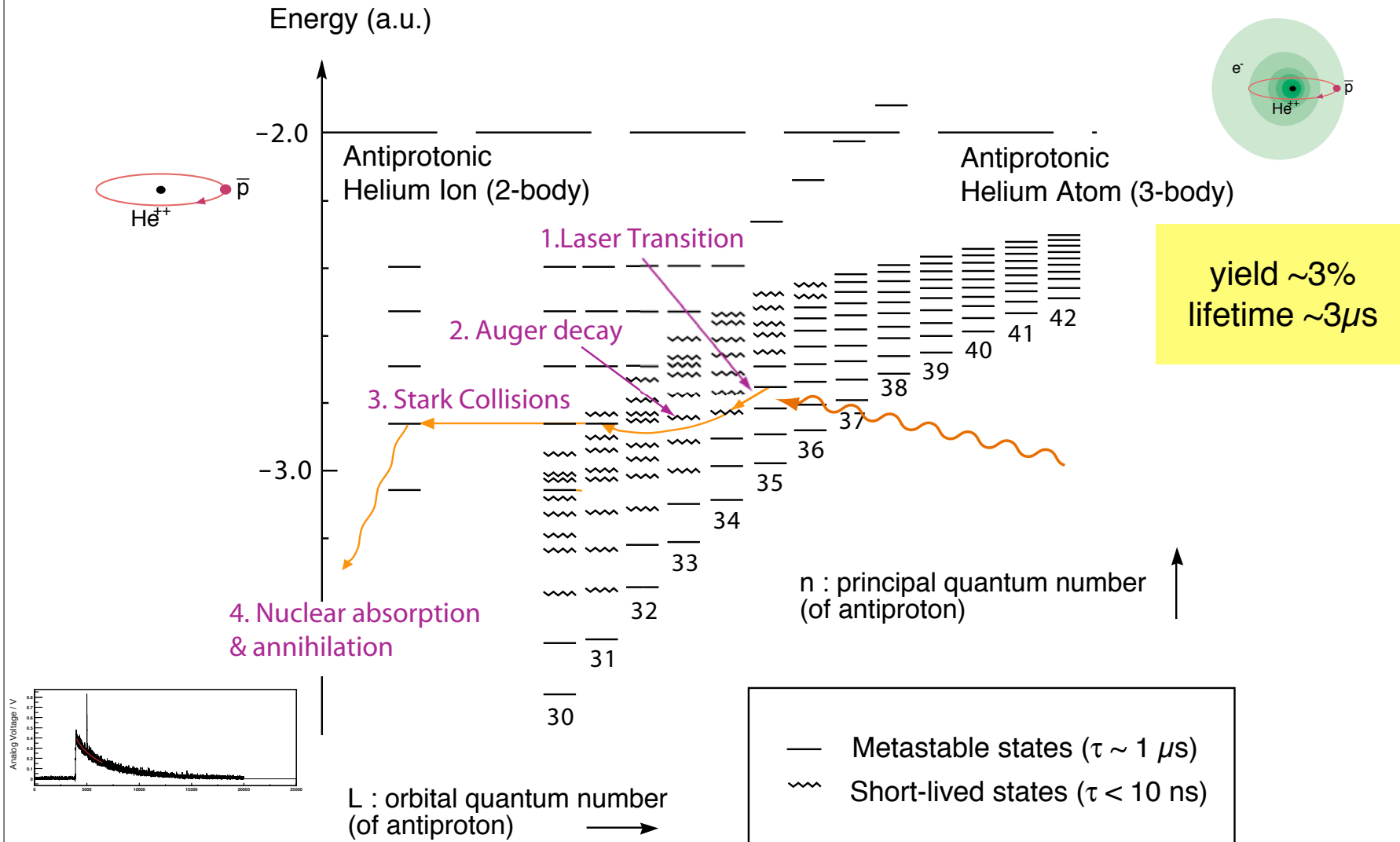
Antiprotonic helium

antiprotonic helium vs antihydrogen

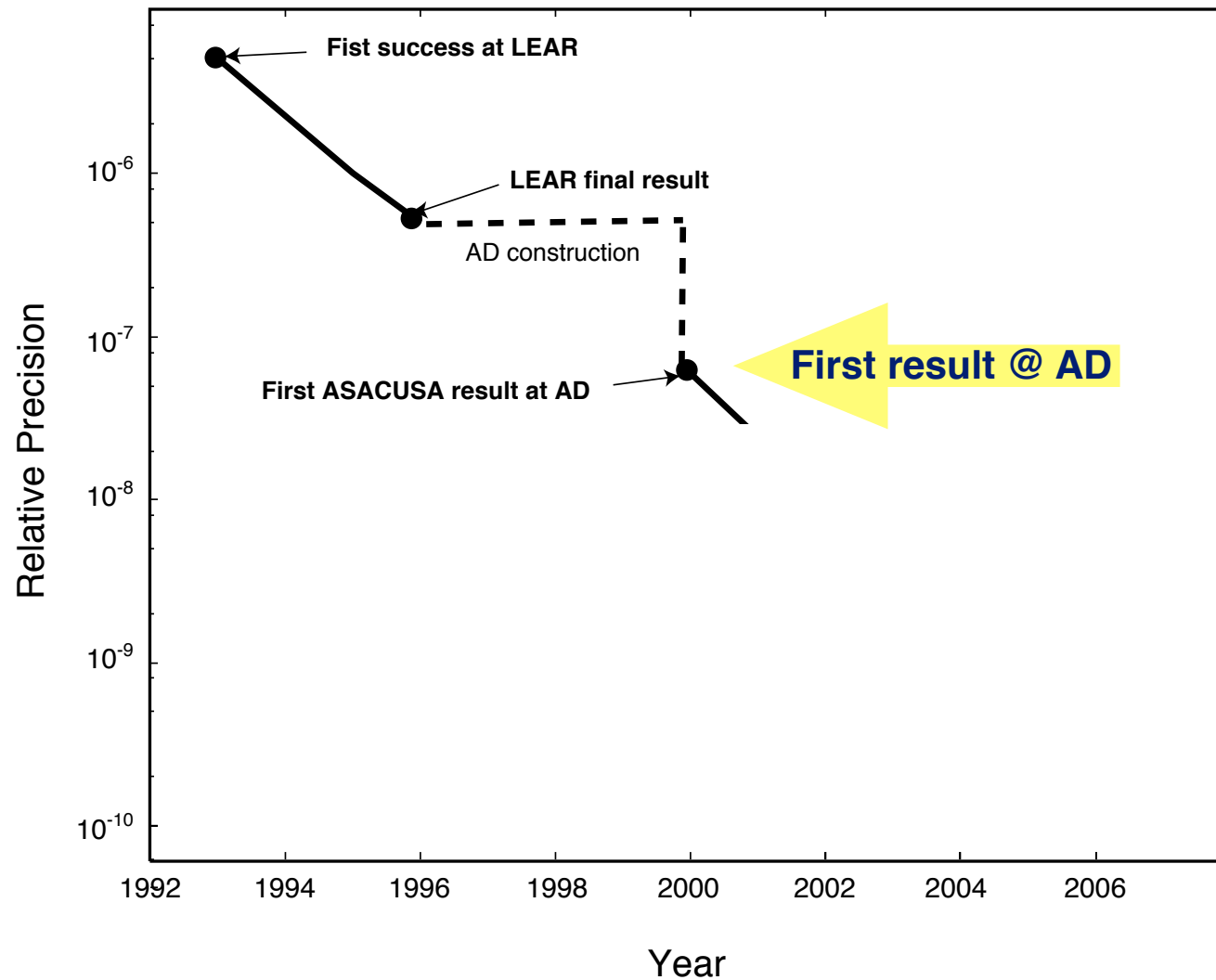


| | antiprotonic helium atom | Antihydrogen |
|---------------------------|--|--|
| System | 3-body | 2-body |
| Theory | Difficult, accuracy ~ppb? | not needed if H and HBar are compared |
| Production | Easy, can be abundantly produced | demonstrated, but not so many in the ground state yet? |
| Cold? | Yes, < 10 K | ??? |
| Ultimate precision | $10^{-10} \sim 12$ | $10^{-14} \sim 18$ |
| Outlook | Will continue to provide the best baryonic CPT test for some more time | Future hopeful |

Antiprotonic helium - a closer look



Weighing the antiproton, progress



60 ppb

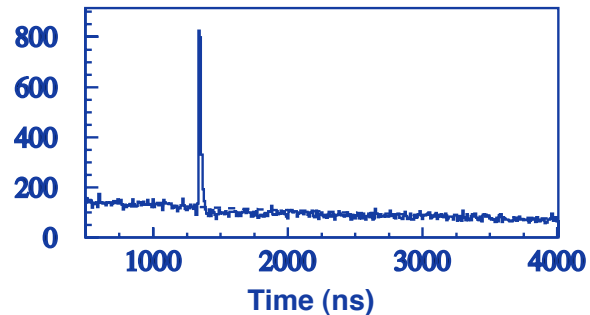
using 5 MeV beam
(ASACUSA phase 1)

NOTE

Every x10 improvement
requires new developments

from LEAR (500 ppb)
to ASACUSA phase 1 (60 ppb)

LEAR offered slow extraction

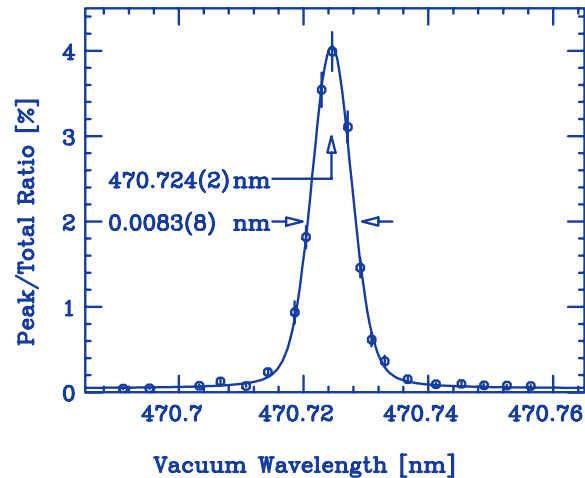


≡ LEAR - **slow** extraction, event by event counting

≡ good event identification

≡ 1 atom, 1 laser shot, ~ 300 Hz

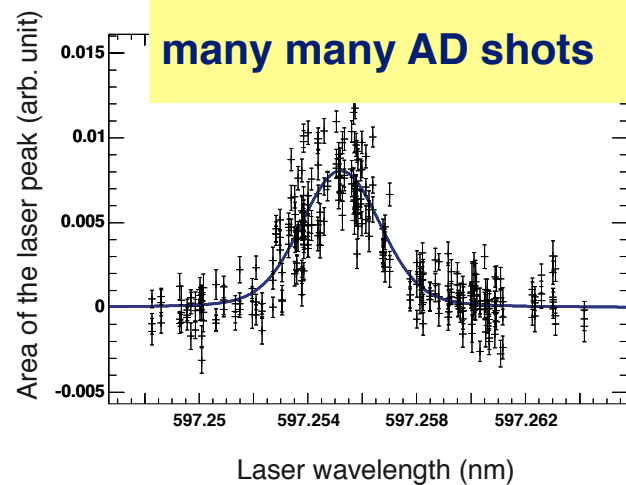
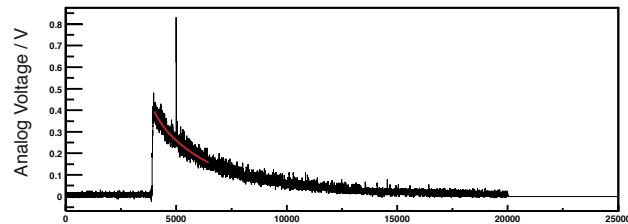
≡ rapid frequency scan so as to minimize (average out) systematic errors



AD - fast extraction



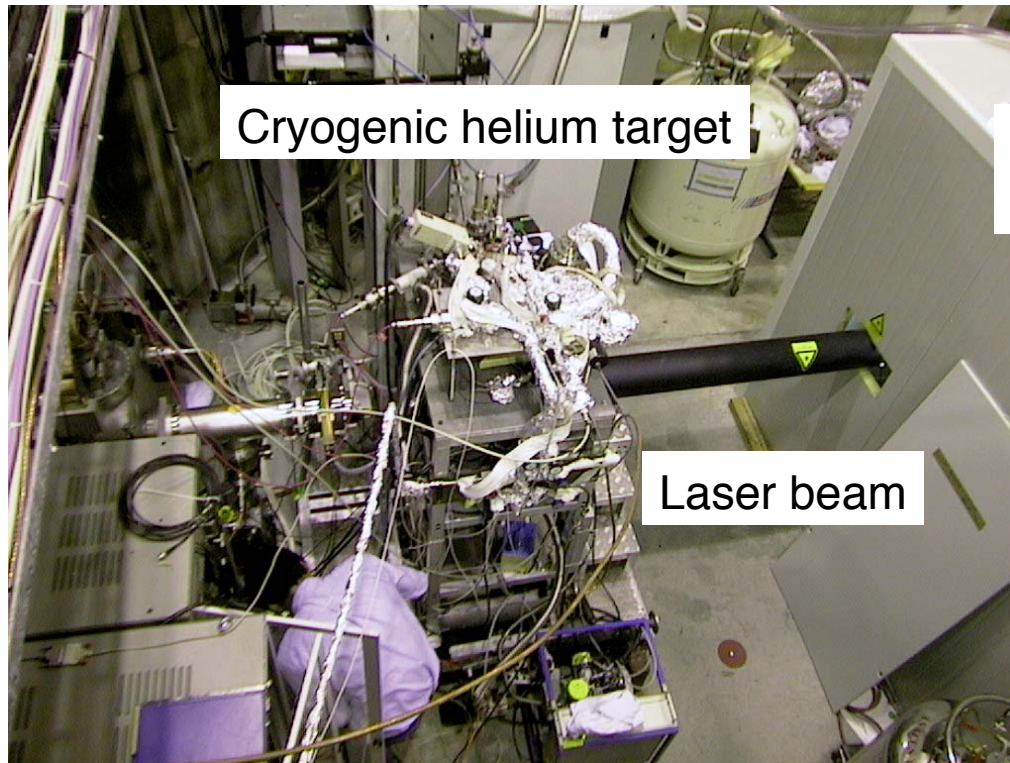
Measured in a
single AD shot



- ≡ AD - **fast** extraction, 3×10^7 pbar, 10^6 atoms, 1 laser pulse
- ≡ Each laser-frequency point \leftrightarrow 1 AD shot
- ≡ Good control of systematics over ~ 8 hours is crucial
 - ≡ target condition
 - ≡ laser shot-to-shot power stability, frequency stability
 - ≡ antiproton beam intensity, position

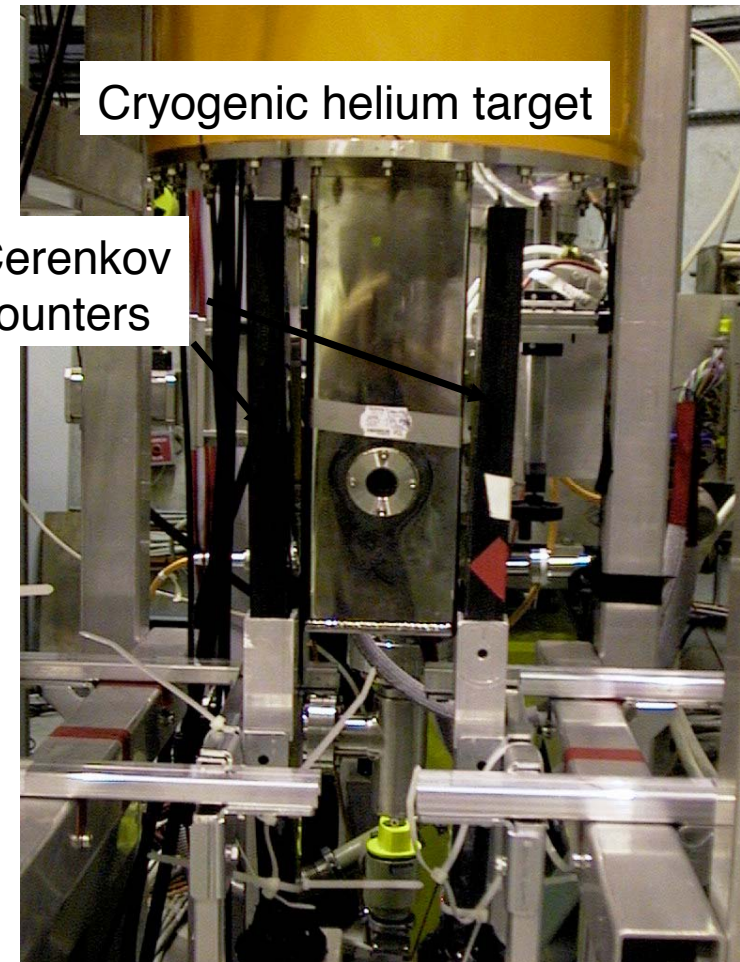


Phase 1 (5 MeV beam)



Cryogenic helium target

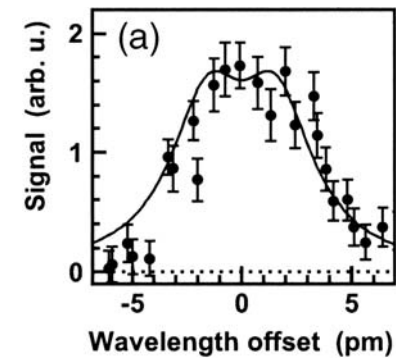
Laser beam



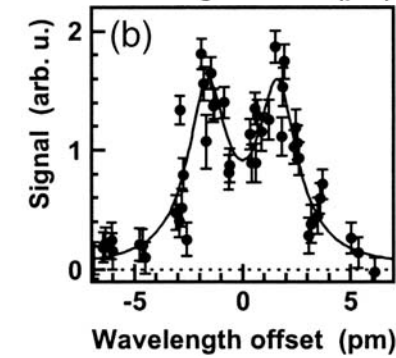
Cryogenic helium target

Cerenkov
counters

Pulsed laser system



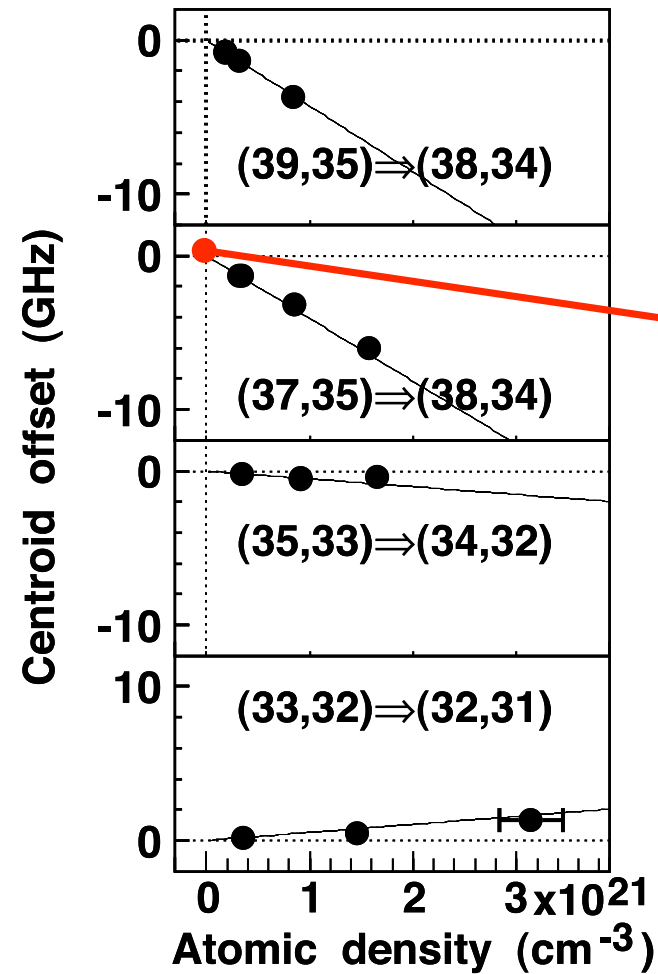
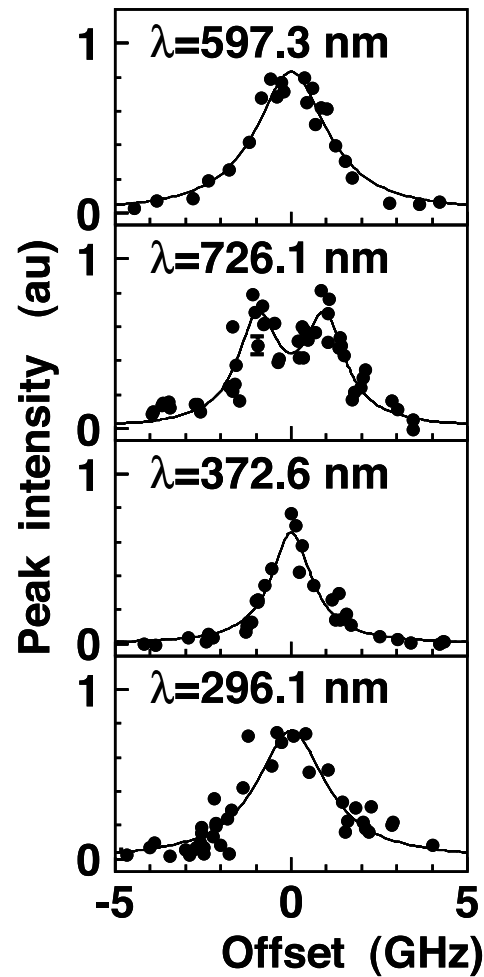
LEAR



AD

hyperfine structure can be resolved (more later)

Collisional frequency shift correction



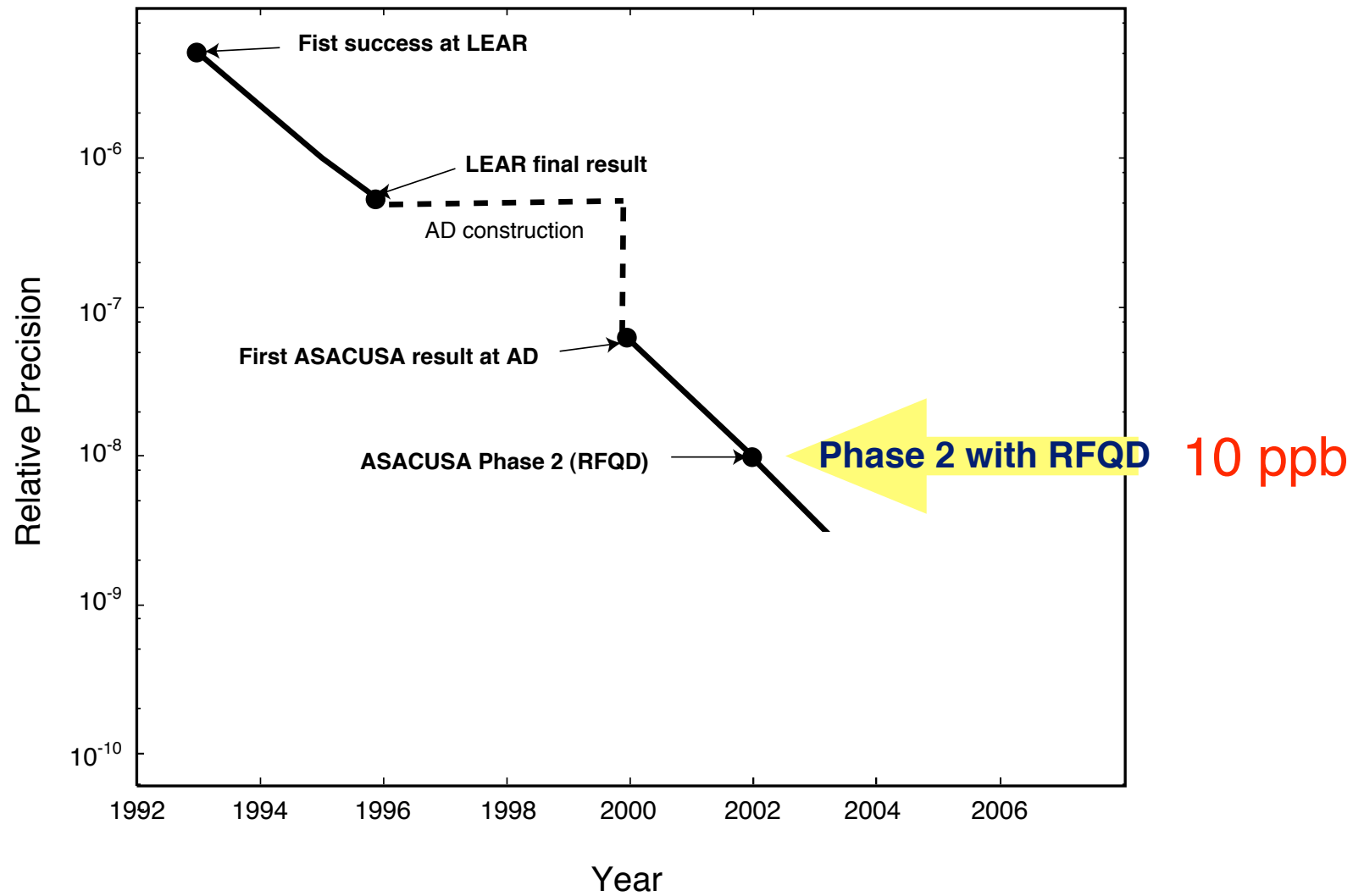
Antiprotons stopped in dense (~ 1 bar, 0.5 K) target, zero-density extrapolation needed

60 ppb

from ASACUSA phase 1 (60 ppb)

to ASACUSA phase 2 (10 ppb)

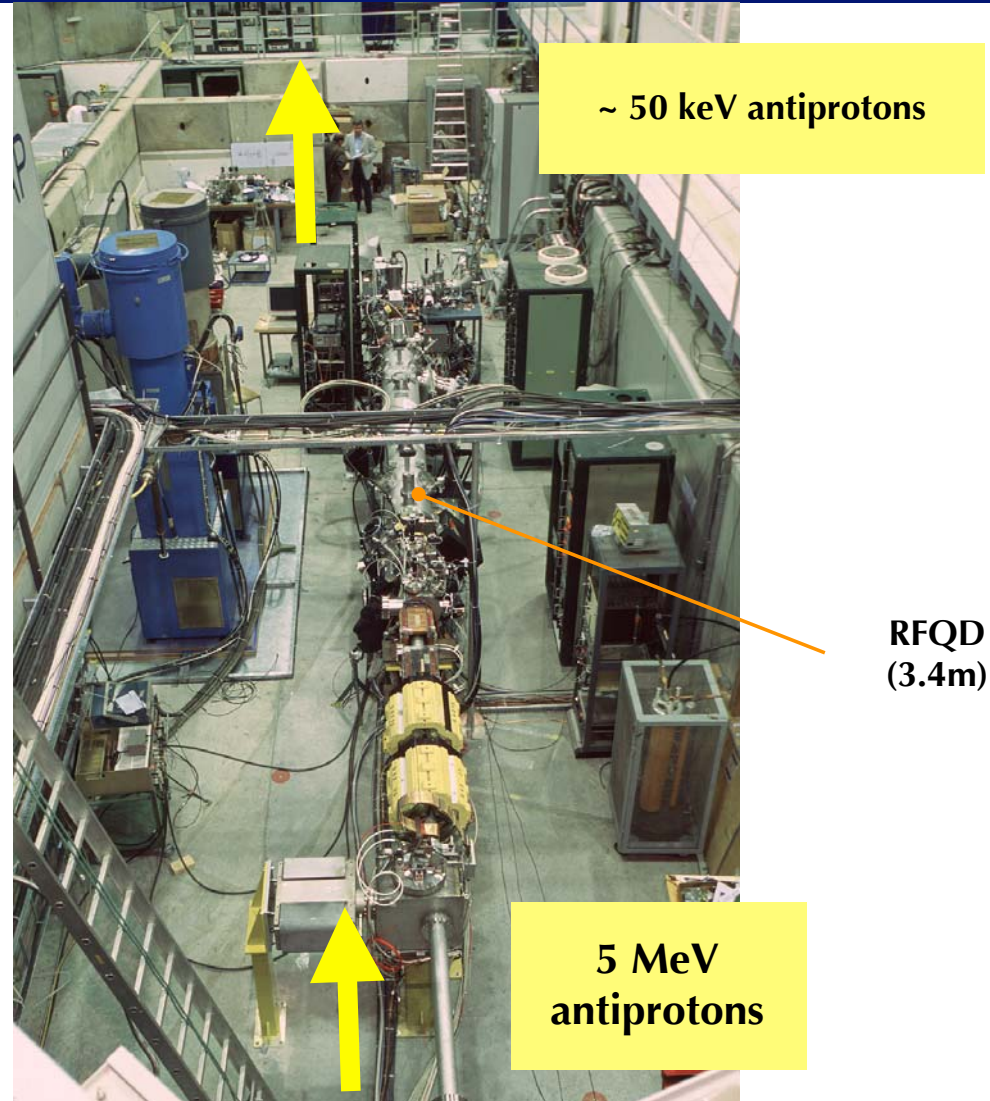
RFQD is essential



RFQD - an inverse linear accelerator



- ≡ RFQD (radio-frequency quadrupole decelerator) is an 'inverse accelerator'
- ≡ Antiprotons can be decelerated from 5 MeV to 50 keV
- ≡ With the 50 keV beam, antiprotonic helium can be produced in a 'near-vacuum' condition



Beam energy & target density



| | Method | Beam Energy | Typical target density | Physics output |
|---------|-----------|-------------|------------------------------------|---|
| Phase 1 | AD direct | 5 MeV | 10^{21}cm^{-3} | $\bar{\rho}\text{He}$ - 60 ppb $\bar{\rho}\text{He}$ HFS interaction with H_2/D_2 |
| Phase 2 | RFQD | <100 keV | $10^{16} - 10^{18} \text{cm}^{-3}$ | $\bar{\rho}\text{He}$ - 10 ppb dE/dx to <10 keV |

5 orders of magnitude

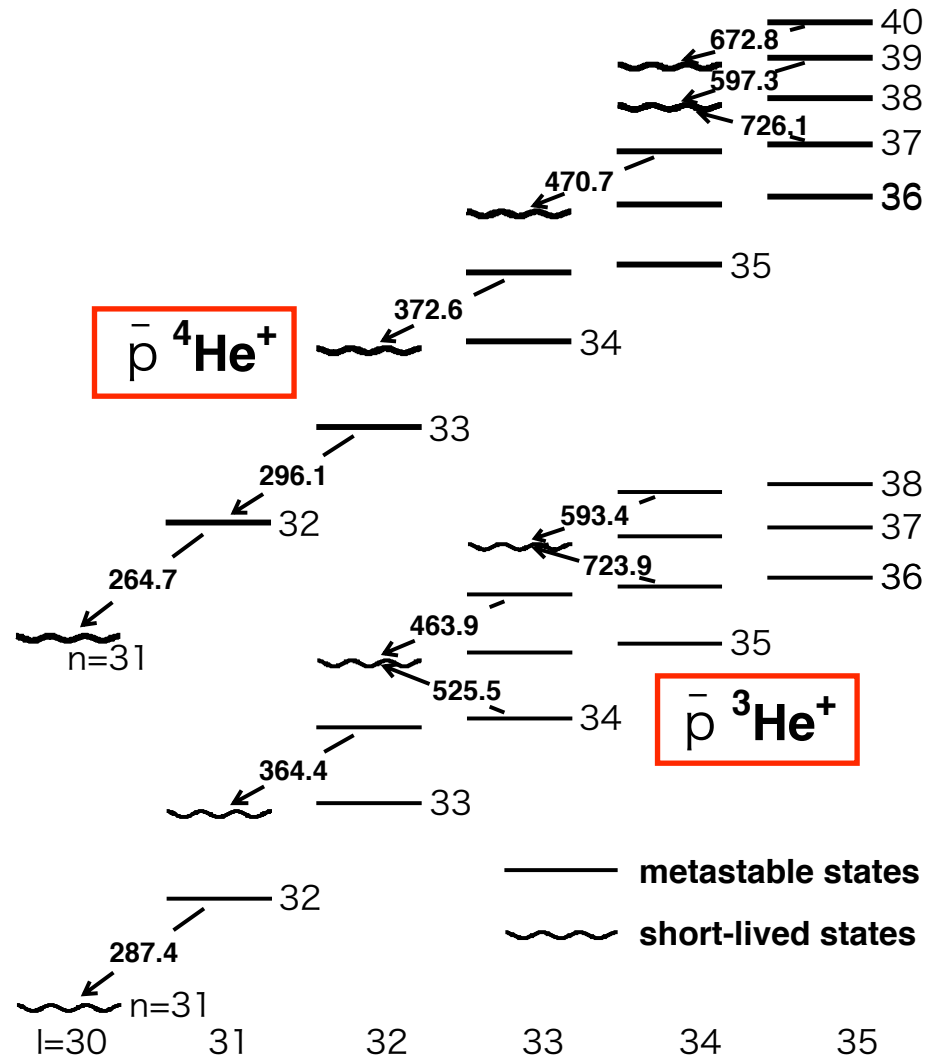
RFQD eliminates collisional effects



| | Phase 1 | Phase 2 | Notes |
|-------------------------------------|---------------------|---------|--|
| Natural width | 0.1 - 100000 MHz | ← | |
| Collisional Shift | ~500 MHz | <1 MHz | Shift is state dependent, difficult systematics |
| Collision width | ~500 MHz | ~1 MHz | |
| Doppler width | ~500 MHz | ← | Peak center can be determined to ~1/100 of the width |
| Laser band width (beaware of chirp) | 800 ~2000 MHz | ← | |
| Calibration | 10 - 60 MHz | ← | |
| Achieved precision | 60 ppb | 10 ppb | |

for a typical transition, 5 MHz ↔ 10 ppb

13 transitions were measured



and the results were compared with
3-body QED theoretical calculations

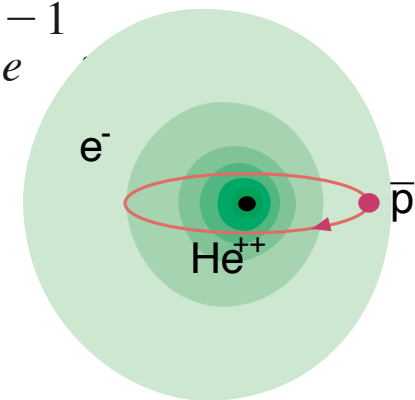
Theory - non-relativistic H



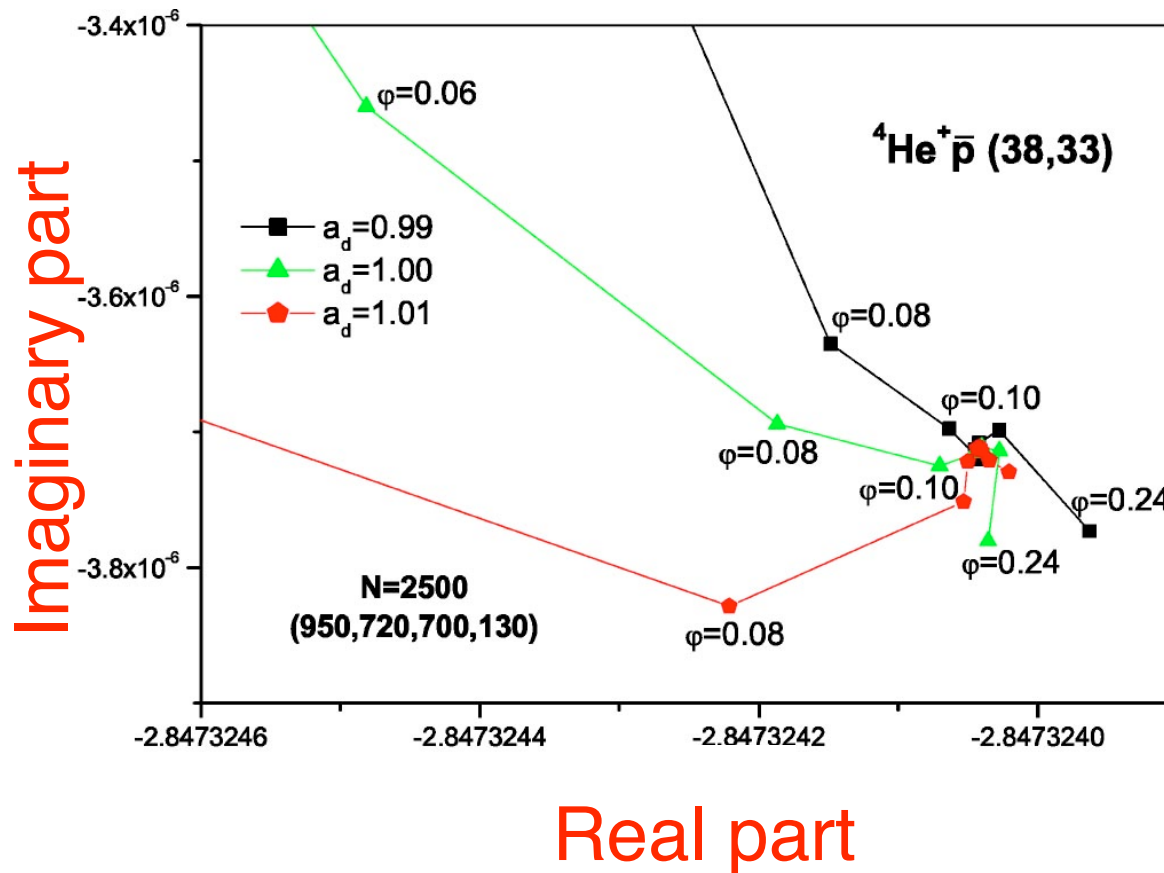
$$H = T + V$$
$$= -\frac{1}{2\mu_1} \nabla_{\mathbf{R}}^2 - \frac{1}{2\mu_2} \nabla_{\mathbf{r}}^2 - \frac{1}{M_{\text{He}}} \nabla_{\mathbf{R}} \cdot \nabla_{\mathbf{r}} - \frac{2}{R} - \frac{2}{r} + \frac{1}{|\mathbf{R} - \mathbf{r}|},$$

antiproton electron

$$\mu_1^{-1} = M_{\text{He}}^{-1} + M_X^{-1}, \quad \mu_2^{-1} = M_{\text{He}}^{-1} + m_e^{-1}$$



Complex coordinate rotation method



≡ not true bound states

≡ careful treatment of Auger decay is needed

≡ complex eigenvalues calculated by using the “complex coordinate rotation” method

add relativistic correction (~ 100 ppm)



$$H = T + V$$

$$= -\frac{1}{2\mu_1} \nabla_{\mathbf{R}}^2 - \frac{1}{2\mu_2} \nabla_{\mathbf{r}}^2 - \frac{1}{M_{\text{He}}} \nabla_{\mathbf{R}} \cdot \nabla_{\mathbf{r}} - \frac{2}{R} - \frac{2}{r} + \frac{1}{|\mathbf{R} - \mathbf{r}|},$$

$$\mu_1^{-1} = M_{\text{He}}^{-1} + M_X^{-1}, \quad \mu_2^{-1} = M_{\text{He}}^{-1} + m_e^{-1},$$

$$E_{rc} = \alpha^2 \left\langle -\frac{\mathbf{p}_e^4}{8m_e^3} + \frac{4\pi}{8m_e^2} [Z_{\text{He}} \delta(\mathbf{r}_{\text{He}}) + Z_{\bar{p}} \delta(\mathbf{r}_{\bar{p}})] \right\rangle.$$

add self energy (~15 ppm)



$$H=T+V$$

$$= -\frac{1}{2\mu_1}\nabla_{\mathbf{R}}^2 - \frac{1}{2\mu_2}\nabla_{\mathbf{r}}^2 - \frac{1}{M_{\text{He}}}\nabla_{\mathbf{R}}\cdot\nabla_{\mathbf{r}} - \frac{2}{R} - \frac{2}{r} + \frac{1}{|\mathbf{R}-\mathbf{r}|},$$

$$\mu_1^{-1} = M_{\text{He}}^{-1} + M_X^{-1}, \quad \mu_2^{-1} = M_{\text{He}}^{-1} + m_e^{-1},$$

$$E_{rc} = \alpha^2 \left\langle -\frac{\mathbf{p}_e^4}{8m_e^3} + \frac{4\pi}{8m_e^2} [Z_{\text{He}}\delta(\mathbf{r}_{\text{He}}) + Z_p^-\delta(\mathbf{r}_p^-)] \right\rangle.$$

$$\begin{aligned} E_{se} &= \frac{4\alpha^3}{3m_e^2} \left[\ln \frac{1}{\alpha^2} - \ln \frac{k_0}{R_\infty} + \frac{5}{6} - \frac{3}{8} \right] \langle Z_{\text{He}}\delta(\mathbf{r}_{\text{He}}) + Z_p^-\delta(\mathbf{r}_p^-) \rangle \\ &+ \frac{4\alpha^4}{3m_e^2} \left[3\pi \left(\frac{139}{128} - \frac{1}{2} \ln 2 \right) \right] \langle Z_{\text{He}}^2\delta(\mathbf{r}_{\text{He}}) + Z_p^2\delta(\mathbf{r}_p^-) \rangle \\ &- \frac{4\alpha^5}{3m_e^2} \left[\frac{3}{4} \right] \langle Z_{\text{He}}^3 \ln^2(Z_{\text{He}}\alpha)^{-2} \delta(\mathbf{r}_{\text{He}}) \\ &+ Z_p^3 \ln^2(Z_p^-\alpha)^{-2} \delta(\mathbf{r}_p^-) \rangle, \end{aligned}$$

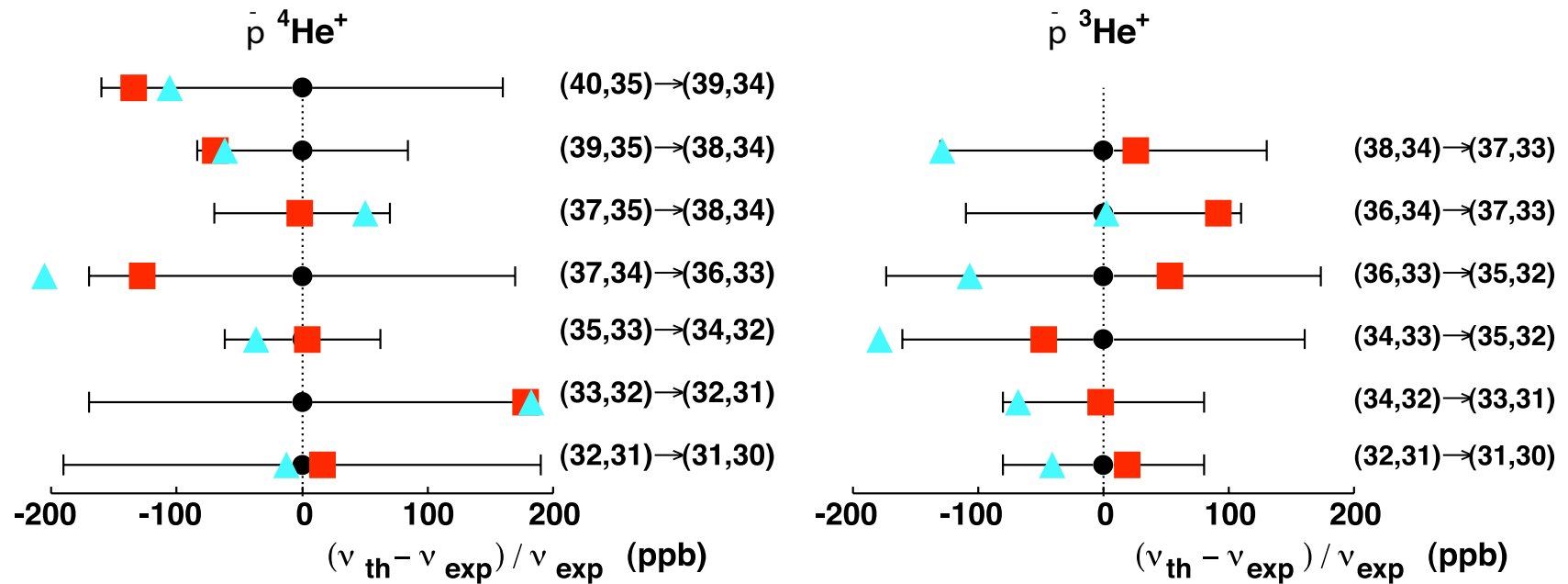
Breakdown of various contributions



(37,34) → (38,33) example (Korobov)

| | | | |
|--------------|---|-----------------|--|
| E_{nr} | = | 420 158 166(20) | |
| E_{rc} | = | -43 753(30) | |
| E_{rc-QED} | = | 360 | |
| E_{se} | = | 5 929(5) | |
| E_{vp} | = | -189 | ~ 2 ppb - 100 ppb, depending on the Auger lifetime |
| E_{kin} | = | -4 | |
| E_{ret} | = | -65 | |
| E_{fsc} | = | 4 | |
| E_{total} | = | 420 120 448(40) | MHz |

Theory vs Experiment



Two theory calculations (▲ and ■) compared with experiment ●
▲ and ■ differ up to about 100 ppb

In order to improve, we must



- ≡ Reduce the experimental error bar by an order of magnitude
 - ≡ One of the two theoretical calculations turn may turn out to be wrong
- ≡ Urge theorists to work harder

Part 2

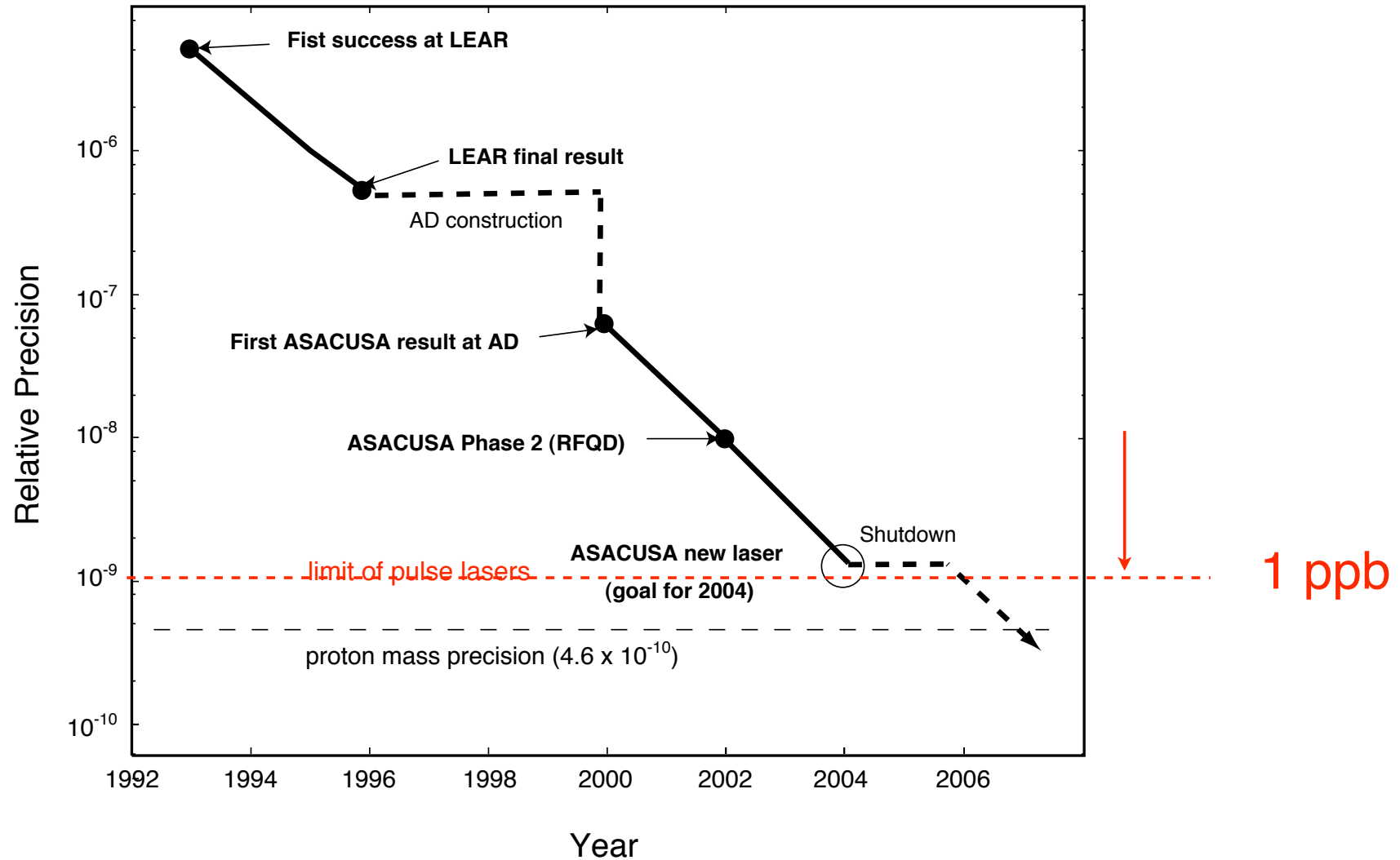
ASACUSA FUTURE

| Physics Goals | Measurements | | Method | Status & Outlook | |
|--|-------------------------------|--|---|--|--|
| CPT tests: high-precision determination of antiproton mass, charge, magnetic moment and magnetic structure using various spectroscopic methods | 3-body system | Antiprotonic helium atom laser & microwave spectroscopy | RFQD +low-density target | New high-precision laser system. Two-photon spectroscopy will enable ultimate accuracy. | Proton mass is known to 0.46 ppb. The goal is to measure antiproton mass to a similar precision. |
| | 2-body systems | Antiprotonic helium ion laser spectroscopy | RFQD +low-density target | Free from theoretical ambiguities. | The magnetic properties of antiproton poorly known. The goal is to compare the magnetic structure of proton and antiproton. |
| | | Antihydrogen ground-state HFS microwave spectroscopy | RFQD + Paul trap + Two-tone Paul trap (point source of cold antihydrogen) | Superconducting Paul traps being developed first Hbar production test in 2006. | |
| | | | RFQD + Penning trap + cusp trap (possible source of polarized antihydrogen beam) | Cusp trap being developed. Proton + electron test to be done in Japan in 2004-2005. | |
| | Auxiliary measurements | Collisional behavior of very low energy antiprotons | RFQD + gas/solid targets 100 eV beam extracted from Penning trap | Ready to start measurements in 2006. Potential "users" of the extracted beam | |

from ASACUSA 2002 (10 ppb)

to 1 ppb

going to sub-ppb

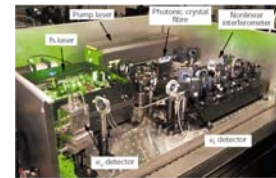


We need a new laser system



| | Phase 1 | Phase 2 | New Laser |
|---|-------------------------|---------|----------------------------------|
| Natural width | 0.1 - 100000 MHz | ← | ← |
| Collisional Shift | ~500 MHz | <1 MHz | ← |
| Collision width | ~500 MHz | ~1 MHz | |
| Doppler width | ~500 MHz | ← | Split by ~1/100 |
| Laser band width beaware of chirp | 800~2000 MHz | ← | < 20 MHz (pulse amplified CW) |
| Calibration | 10 - 60 MHz | ← | ~0 (frequency comb) |
| Achieved precision | 60 ppb | 10 ppb | work in progress |

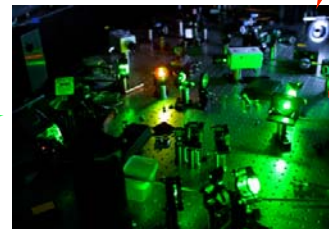
New Laser System



Frequency Comb: calibration with atomic clock precision

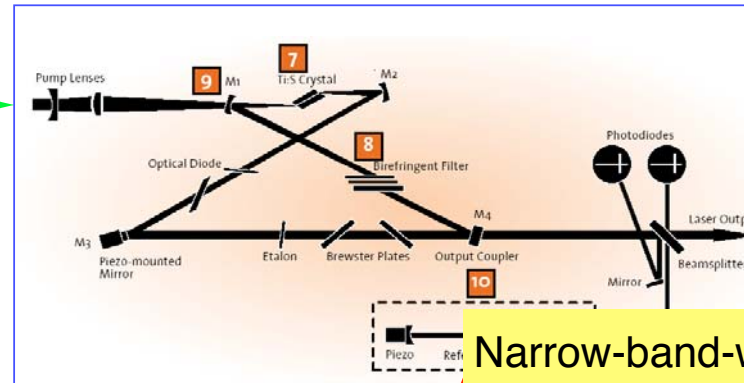
Frequency lock/scan unit

Infinity (Pulsed Nd:YAG)



Pulse amplification: narrow-band high-power

MBR-110 (CW Ti:S)

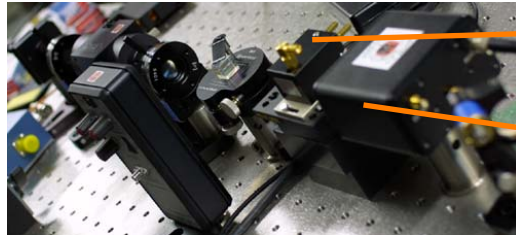


Narrow-band-width CW laser locked to the freq. comb

Chirp compensation / measurement essential

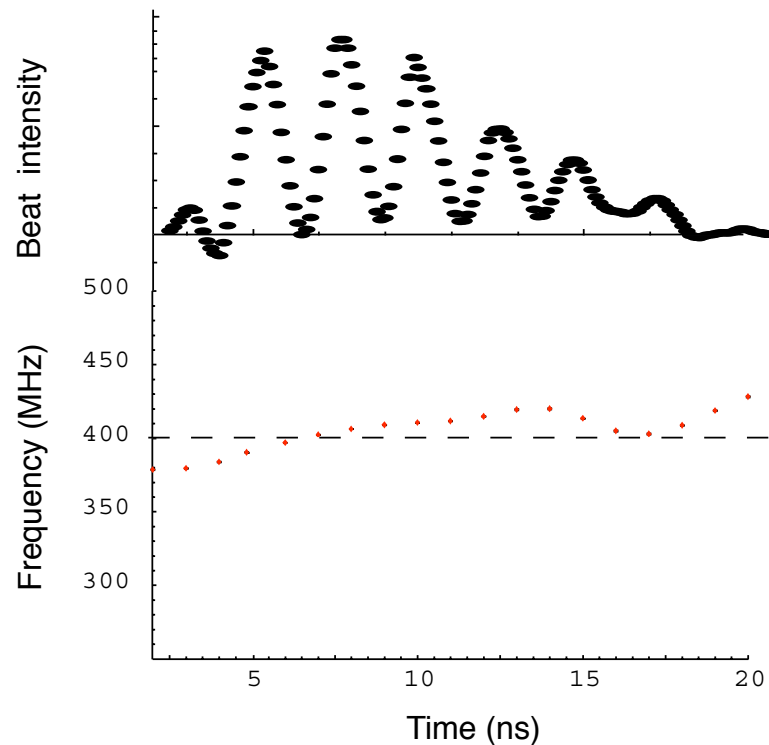
(Photodiode + log amp + digital oscilloscope)

Chirp measurement



Acousto-optic modulator
Shifts CW laser by +400 MHz

Measure “beat note” of
the shifted CW laser and
pulsed laser

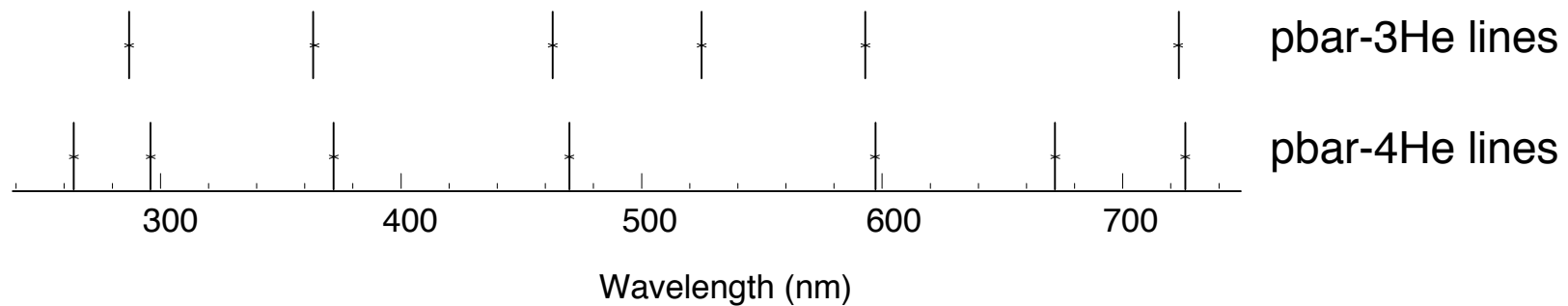


Most of these methods are not new;
e.g., chirp measurement/
compensation was done by
Jungmann et al. for muonium 1s-2s,
but doing this at AD for many
different colors is still a challenge

wavelengths of the resonance lines



UV  IR

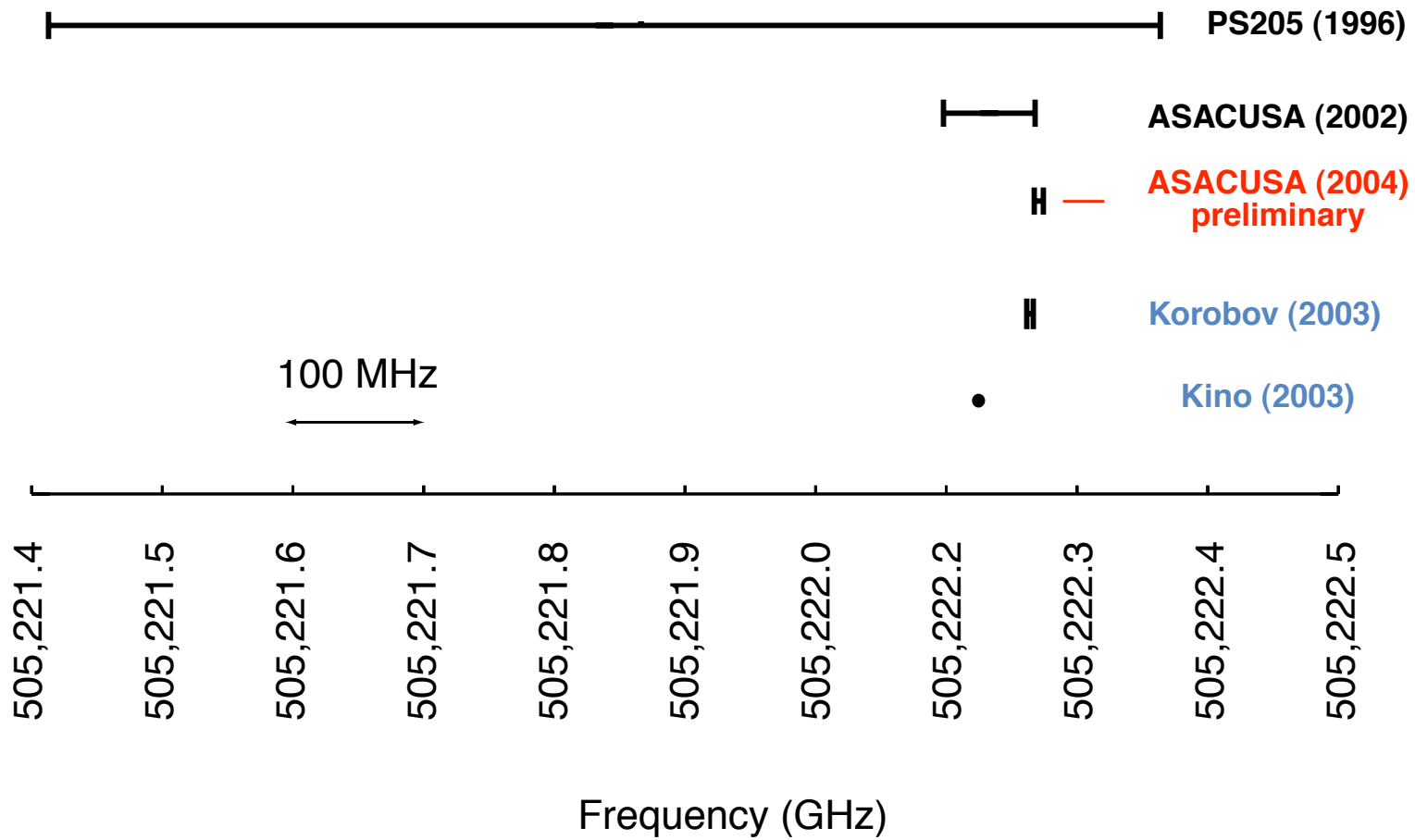


it is nontrivial to achieve highest accuracy for these many transitions

Preview of the 2004 result



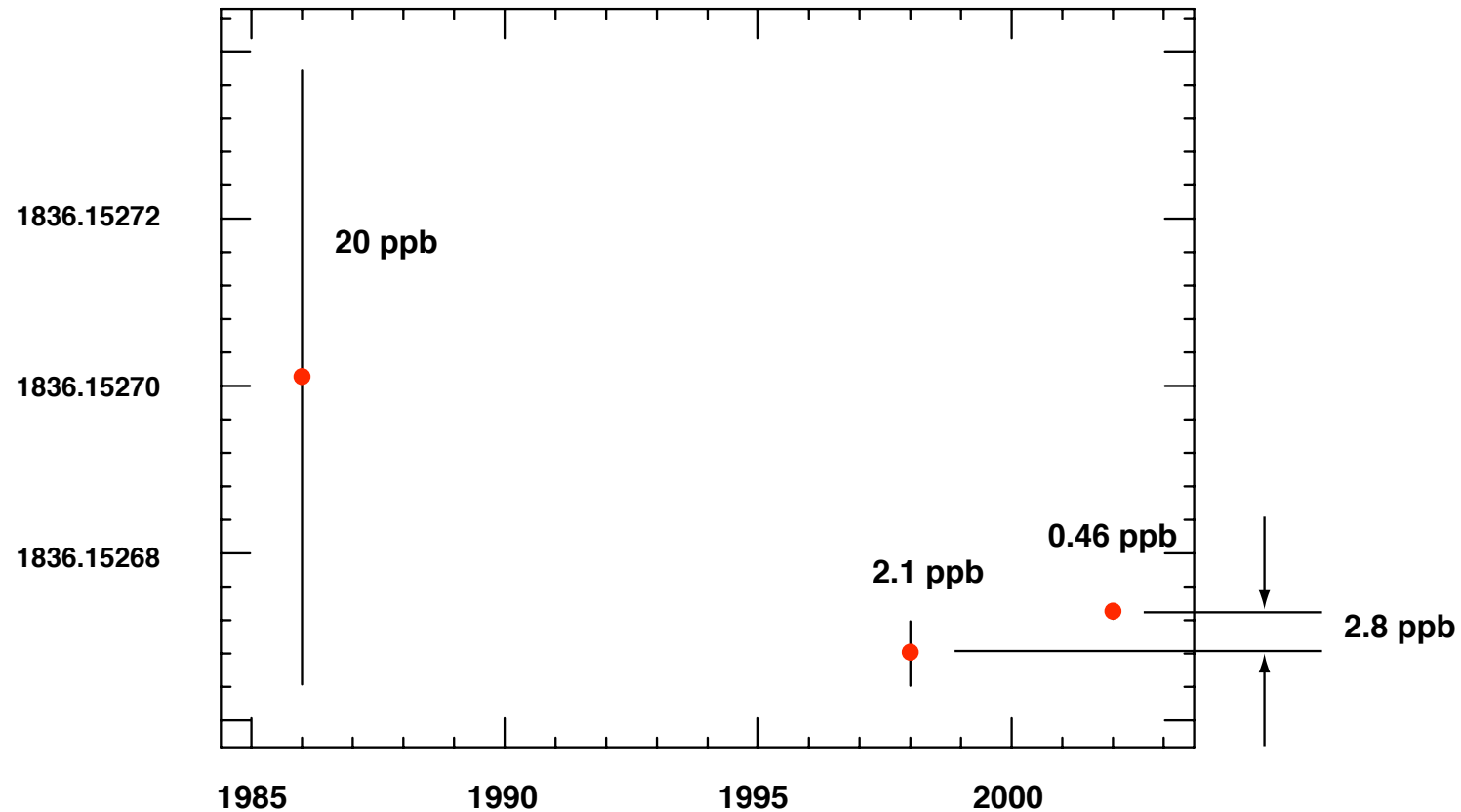
593-nm (505,222 GHz) resonance in helium3



Meanwhile the proton mass has moved

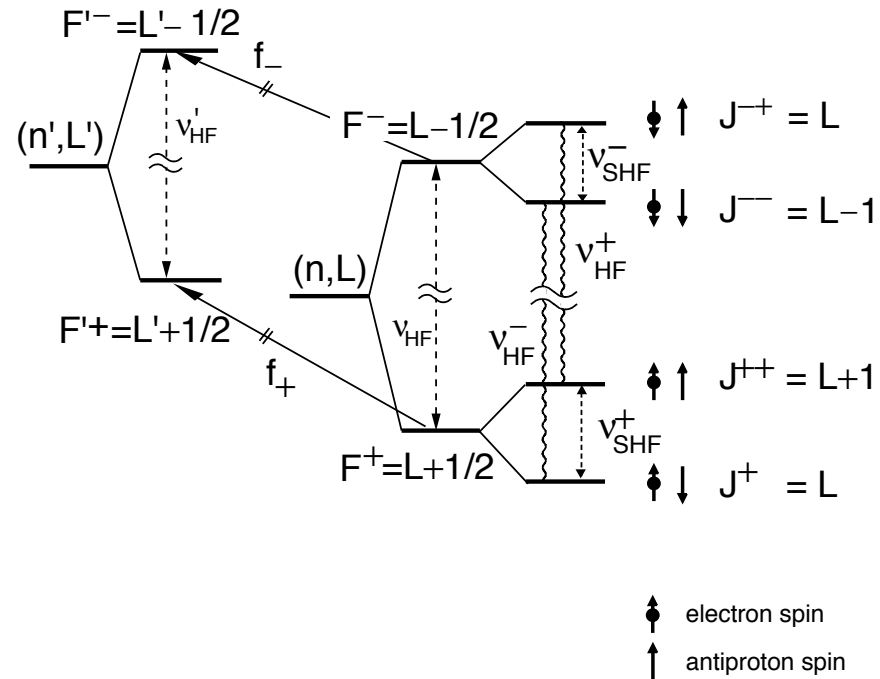
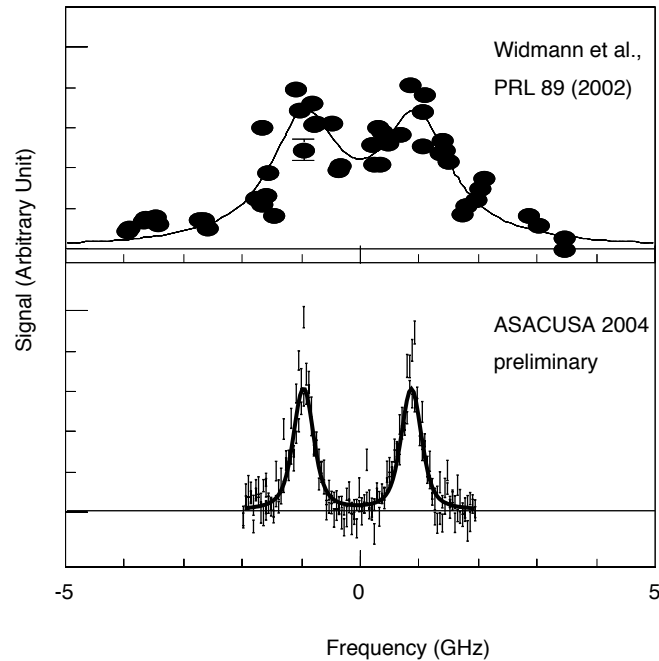


m_p/m_e vs CODATA years



note: alpha mass/proton mass known to 0.13 ppb

HFS → magnetic moment



- ≡ HFS measurement, 726-nm laser + 13GHz microwave, so far limited by laser
- ≡ with the new laser, accuracy improvement possible
- ≡ antiproton μ known only to 0.3%, ASACUSA 2001 was 1.6%
- ≡ In 2006 we will measure antiproton μ to $\ll 0.1\%$

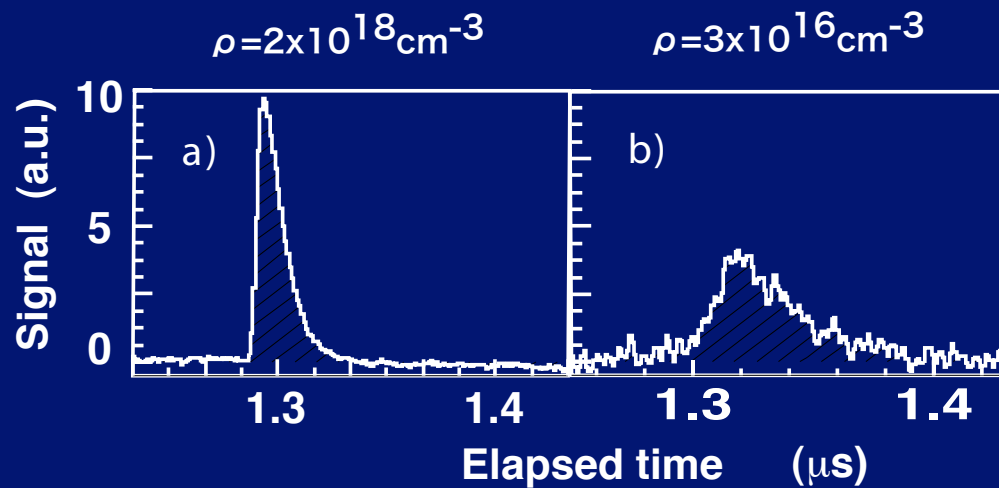
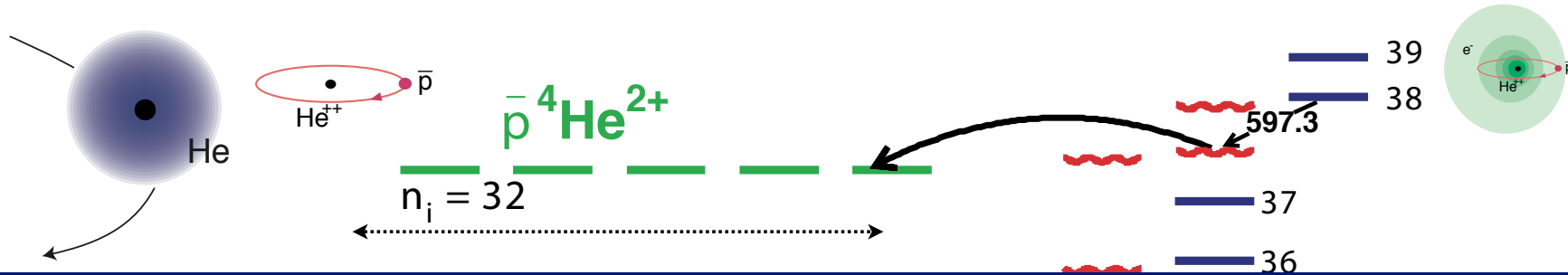
| Physics Goals | Measurements | | Method | Status & Outlook | |
|---|--------------------------------------|---|--|---|---|
| <p>CPT tests:</p> <p>high-precision determination of antiproton mass, charge, magnetic moment and magnetic structure using various spectroscopic methods</p> | <p>3-body system</p> | <p>Antiprotonic helium atom laser & microwave spectroscopy</p> | <p>RFQD+low-density target</p> | <p>New high-precision laser system.</p> <p>Two-photon spectroscopy will enable ultimate accuracy.</p> | <p>Proton mass is known to 0.46 ppb.</p> <p>The goal is to measure antiproton mass to a similar precision.</p> |
| | | <p>Antiprotonic helium ion laser spectroscopy</p> | | <p>RFQD+low-density target</p> | |
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difficulties with the 3-body calculations



- ≡ Reduce the experimental error bar by an order of magnitude
 - ≡ One of the two theoretical calculations turn may turn out to be wrong
- ≡ Urge theorists to work harder
- ≡ Try 2-body system(s)

at very low density, Stark rate is reduced



The first observation
of cold, long-lived
antiprotonic helium
ions

- Change of the decay slope of the laser “spike” is due to the prolongation of antiprotonic helium ION lifetime (Stark rate is reduced)

Laser spectroscopy of antiprotonic helium ions



- ≡ Antiprotonic helium ion:
 - ≡ Almost pure classical Bohr atom. No relativistic correction, no QED, no strong interaction, no hyperfine; practically no theoretical error
 - ≡ Already cold (guaranteed to be thermalized to <10 K - parent is a thermalized 3-body atom)
 - ≡ In a well-defined circular orbit $(n, l) = (N, N-1)$
 - ≡ Lifetime long enough for laser spectroscopy

plan was to try this already in 2004
but had to be deferred due to PS-AD failures

| Physics Goals | Measurements | | Method | Status & Outlook | |
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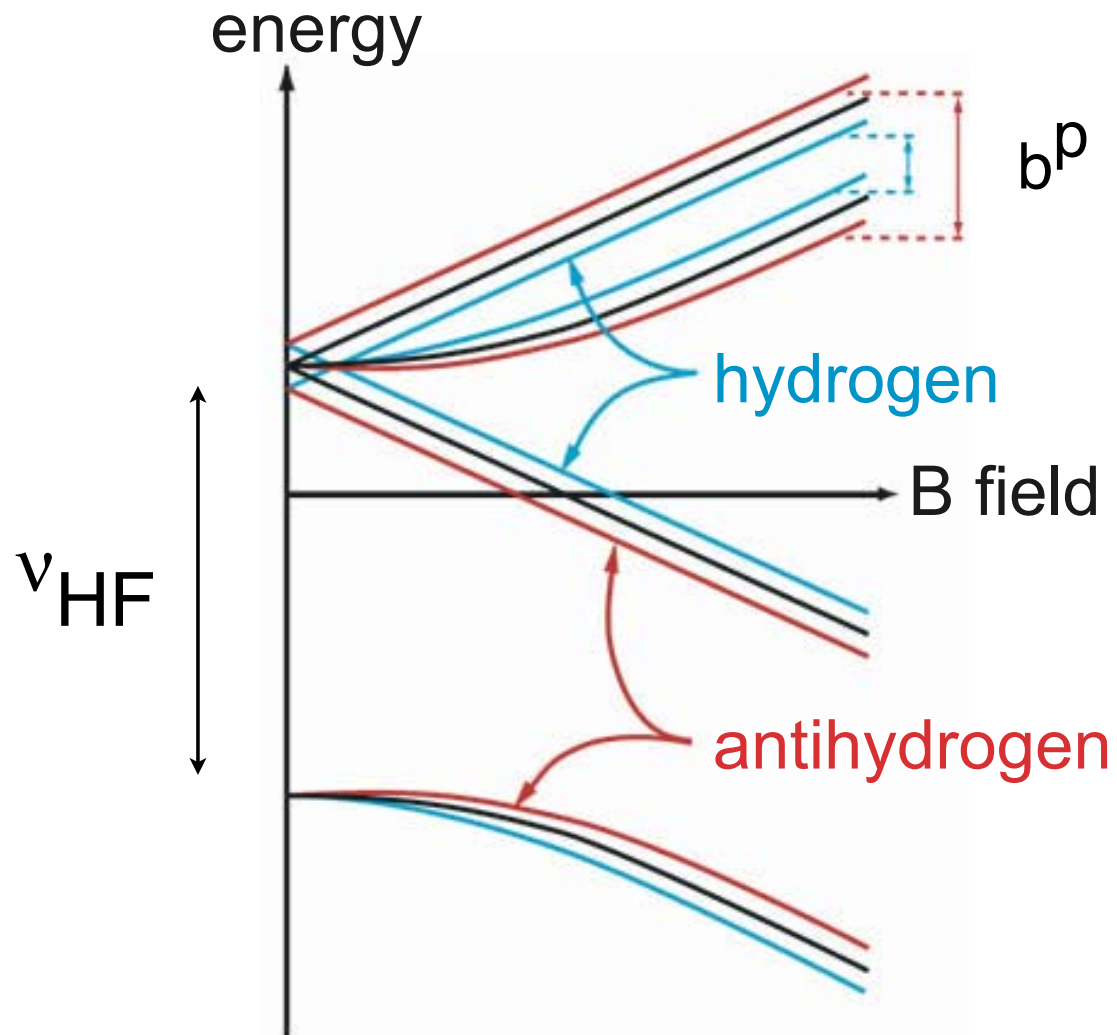
Why yet another antihydrogen experiment?



| | GSHFS | 1s-2s |
|----------------------------|---------------------------|-------------------------------------|
| method | beam + microwave | atom trapping + 2-photon transition |
| sensitivity to CPTV | can directly probe "b" | in free H, no CPTV sensitivity |

$$\begin{aligned}
 & (i\gamma^\mu D_\mu - m_e - a_\mu^e \gamma^\mu - b_\mu^e \gamma_5 \gamma^\mu \\
 & - \frac{1}{2} H_{\mu\nu}^e \sigma^{\mu\nu} + ic_{\mu\nu}^e \gamma^\mu D^\nu + id_{\mu\nu}^e \gamma_5 \gamma^\mu D^\nu) \psi = 0.
 \end{aligned}$$

Breit-Rabi diagram



μ_p and R_M



GS-HFS

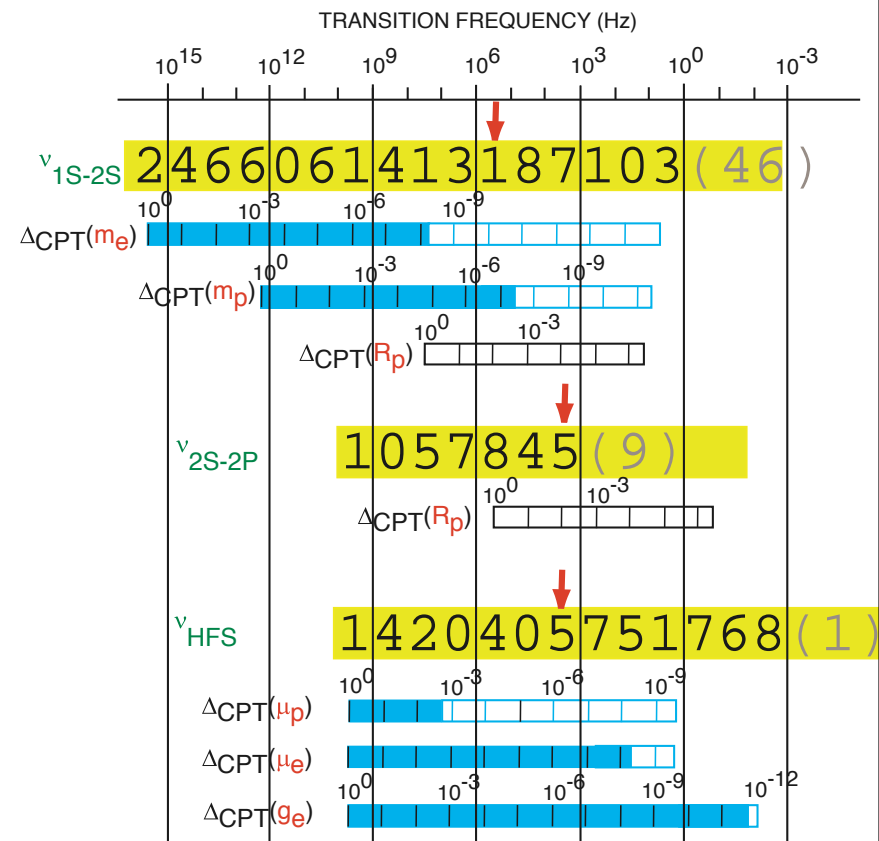
- Proton magnetic moment μ_p
- μ_e
- Proton magnetic radius R_M

Theory

- R_p and R_M

$$\nu_{HF} = \frac{16}{3} \left(\frac{M_p}{M_p + m_e} \right)^3 \frac{m_e}{M_p} \frac{\mu_p}{\mu_N} \alpha^2 Ry$$

$$\Delta\nu(\text{Zemach}) = \nu_F \frac{2Z\alpha m_e}{\pi^2} \int \frac{d^3p}{p^4} \left[\frac{G_E(p^2)G_M(p^2)}{1+\kappa} - 1 \right]$$

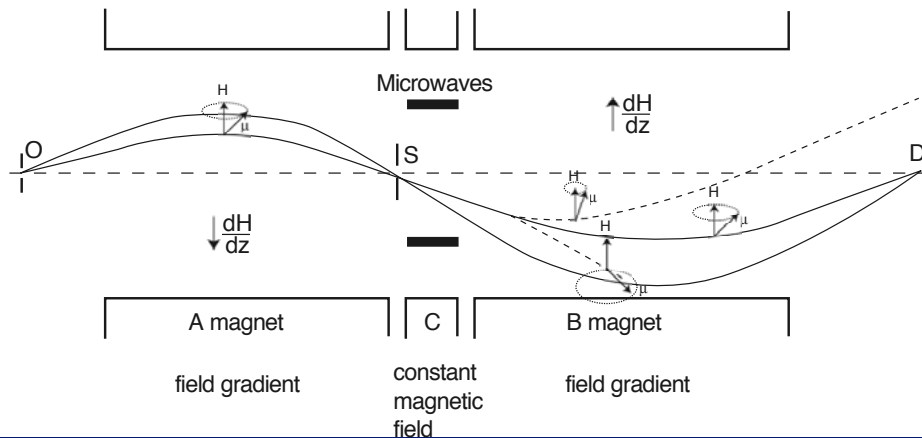


History of Hydrogen HFS



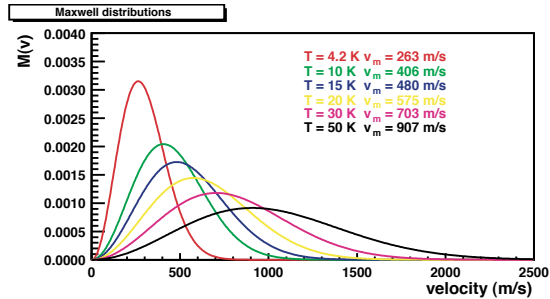
| | | | |
|---------|---------------------------------------|---------------------|--|
| 1936 | Simple atomic beams | $\sim 5\%$ | |
| 1947 | Atomic beams plus microwave resonance | 4×10^{-6} | discovery of anomalous magnetic moment of e^- |
| 1950 | | 4×10^{-8} | |
| 1960-70 | Hydrogen maser | 6×10^{-13} | not possible for antimatter |

Molecular Beam Resonance Setup, I.I. Rabi et al., Phys. Rev. 55, 526 (1923)



N.B. HFS spectroscopy of trapped antihydrogen does **not** necessarily lead to high precision due to the inhomogeneous magnetic field inside the trap

focus - resonate - analyze



Assumed: 15 K source, 1.2T on the pole

Acceptance is 10^{-4}

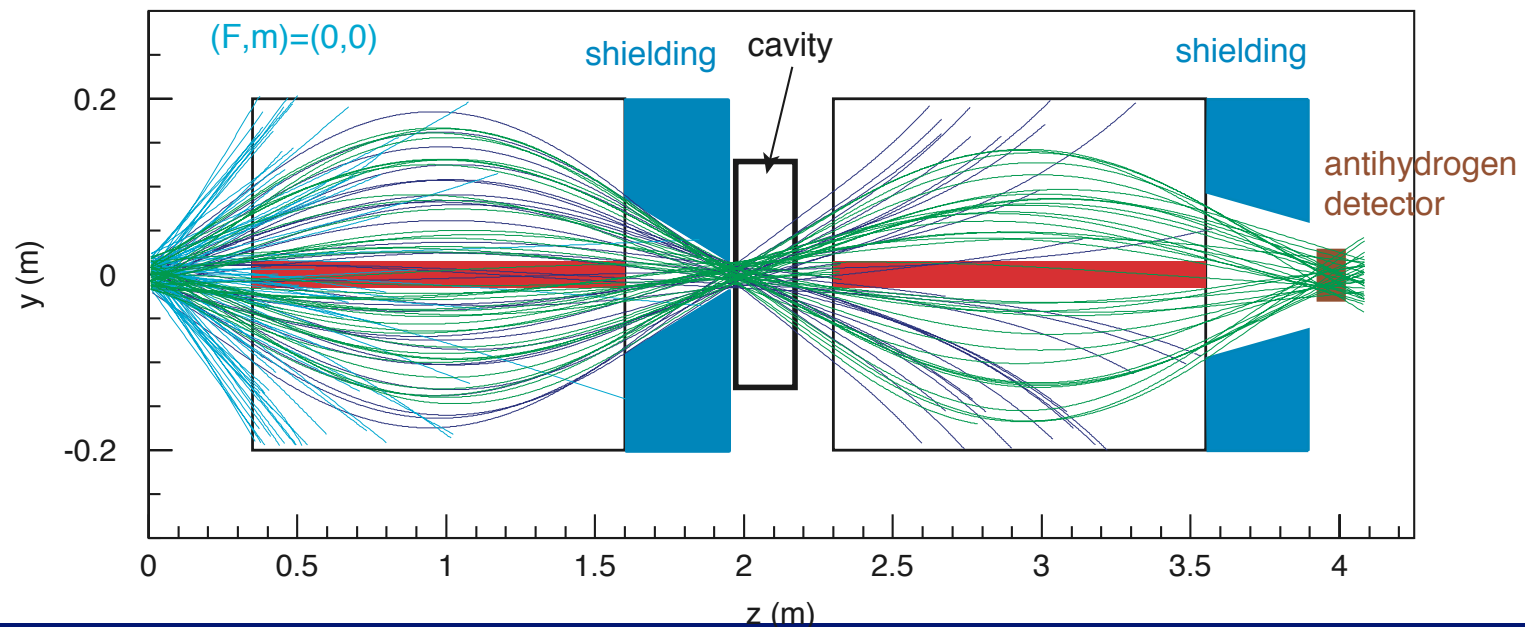
Resonance curve width ~ 2 kHz, 2 ppm

split the line by 1/100 $\rightarrow 10^{-8}$

(F,m)=(1,1) without spin flip

(F,m)=(1,1) with spin flip

(F,m)=(0,0)



How to make a point source

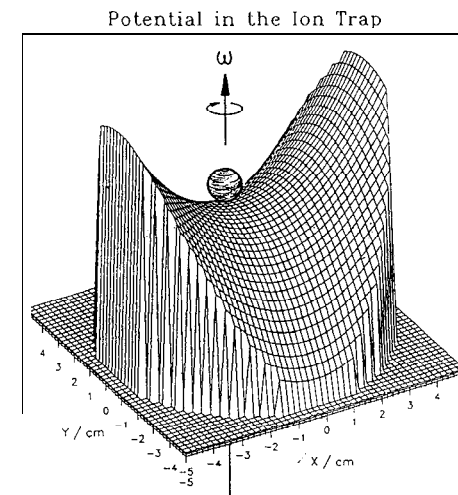


- ≡ atomic-beam geometry works best if the source is point like
- ≡ antihydrogen source size produced in nested Penning trap $\sim 1\text{cm}^3$ - too large
- ≡ limited access (optical & extraction)
- ≡ why not some other methods?

RF (Paul) Trap



- ≡ High-frequency RF (Paul) trap
 - ≡ Paul traps used for high-precision atom studies, but not usually to store a large number of particles
 - ≡ Try to catch, cool, and store a large number of protons, electrons, antiprotons, positrons in Paul traps
 - ≡ Needs high frequency, high field, use superconducting technology



Linear Paul trap, Q-mass, RFQD

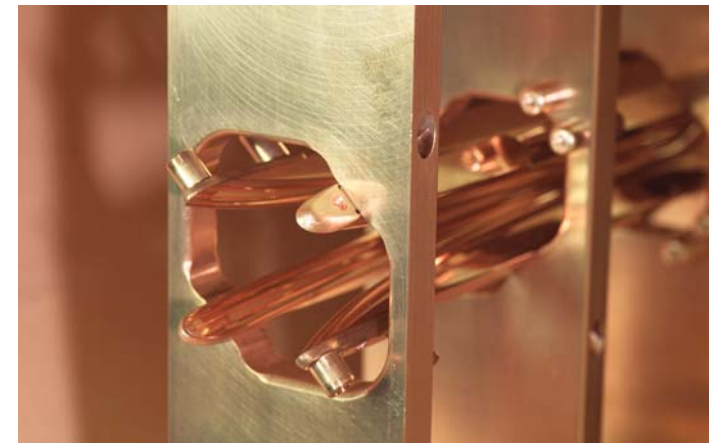
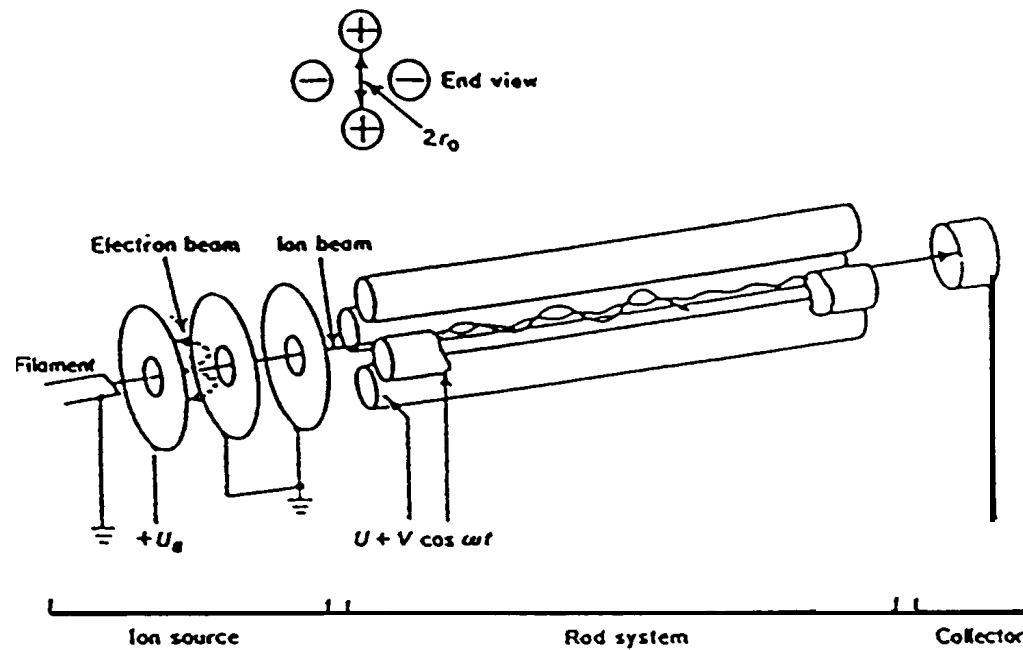


Figure 4. Schematic view of the quadrupole mass spectrometer or mass filter.

why not apply the RFQD technology to antimatter trapping?

Particle cooling in the trap is the key



- ≡ RFQD has no beam cooling
- ≡ No magnetic field → synchrotron cooling is not effective
- ≡ We use resistive cooling (dump energy in an external register)
- ≡ This requires careful optimization of the cavity L & C

Cooling considerations



- ≡ Resistive cooling damps the center-of-motion of the particle cloud.
 - * Effectively damps the secular and micro-motions simultaneously.
 - * Diminishes the particle excursion = size of the cloud.

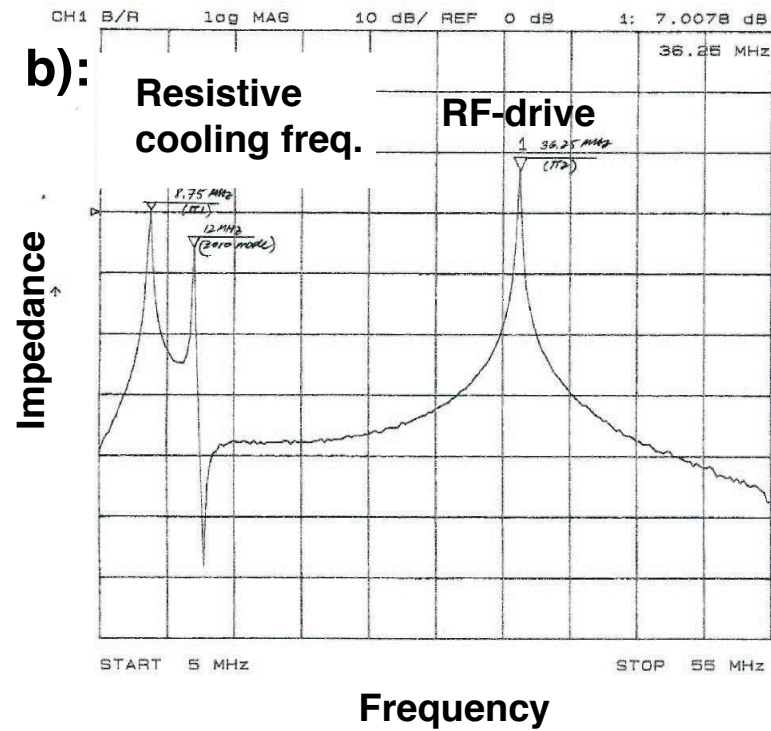
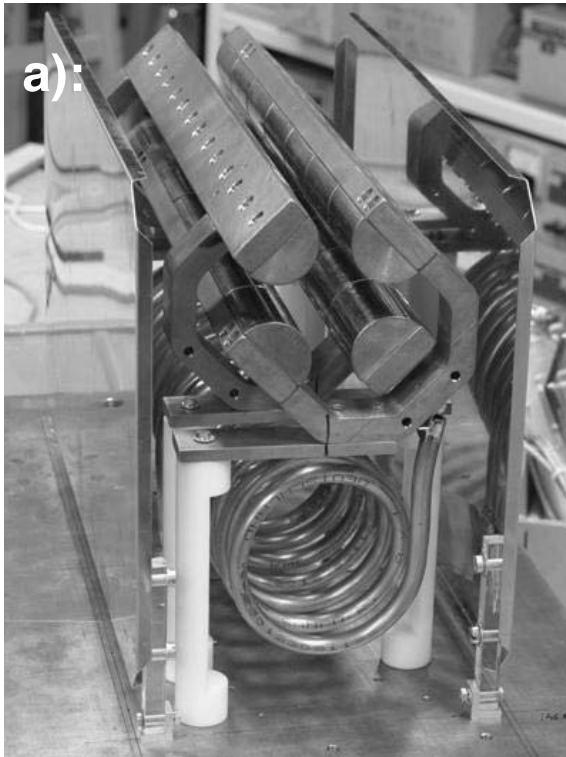
- ≡ Speed of cooling is independent of particle number.
 - Inversely proportional to C and d.
 - Inversely proportional to the square root of the Pseudopotential.

$$\tau_{\text{cool}} = g \frac{8\sqrt{2}mr_0^2\omega(1+s)^2C_{\text{elec}}}{q(1-s)^2Q} = \frac{16\sqrt{2me}\Theta r_0(1+s)^2C_{\text{elec}}}{q^2(1-s)^2Q}.$$

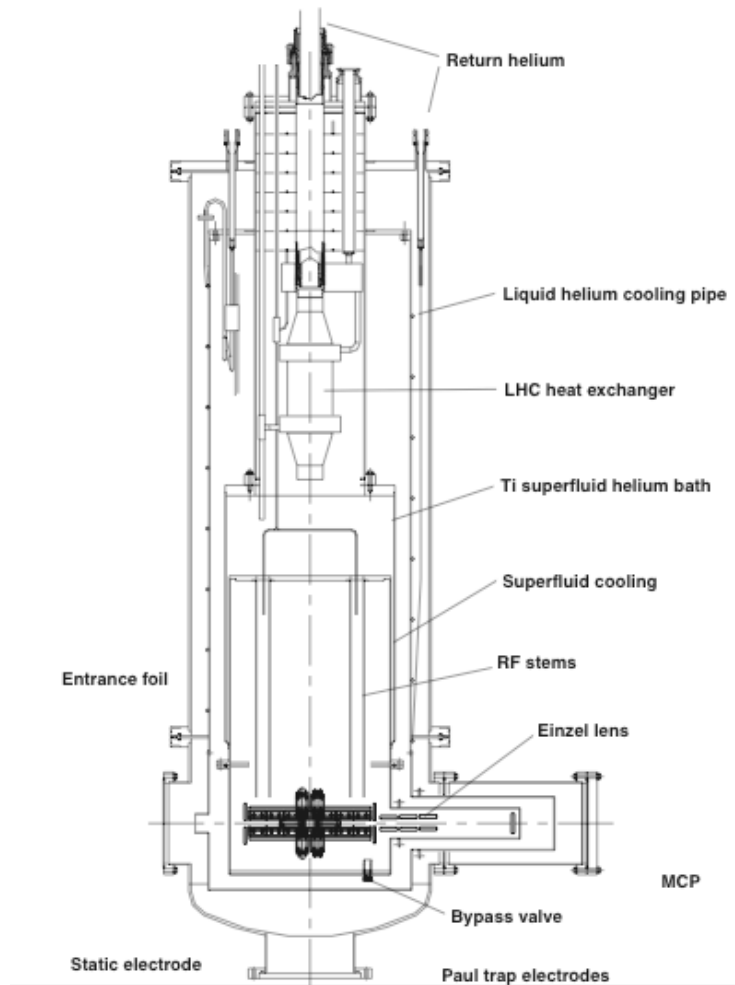
- ≡ Resistive cooling stops when space-charge effects become dominant (the bottom of the harmonic potential becomes flat) and the particles do not move with a definite “secular frequency” any more.
 - * They begin to move chaotically, like the particles in a gas.
 - * The edges of this gas cloud is heated by the radiofrequency.
 - * Particles near the cloud center are Debye-screened from RF.
 - * The heating strength depends on the distance from the trap center.
 - * The size of the cloud = surface temperature depends on N.

$$T = \frac{1}{k_B} \left(\frac{9meN^2(2\omega_r^2 + \omega_z^2)}{128\pi^2} \right)^{1/3} \quad \lambda = \sqrt{\frac{k_B T}{m(2\omega_r^2 + \omega_z^2)}}.$$

Paul trap R&D in progress

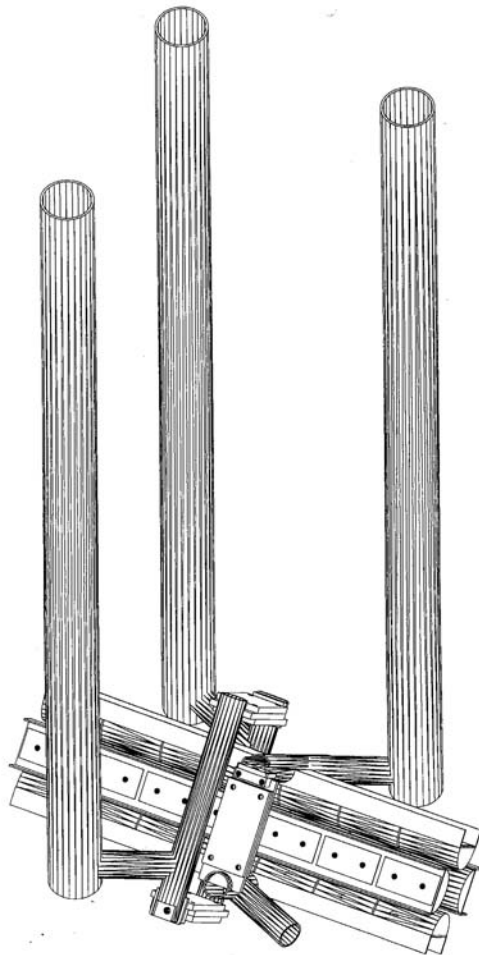


Superconducting Linear Paul Trap



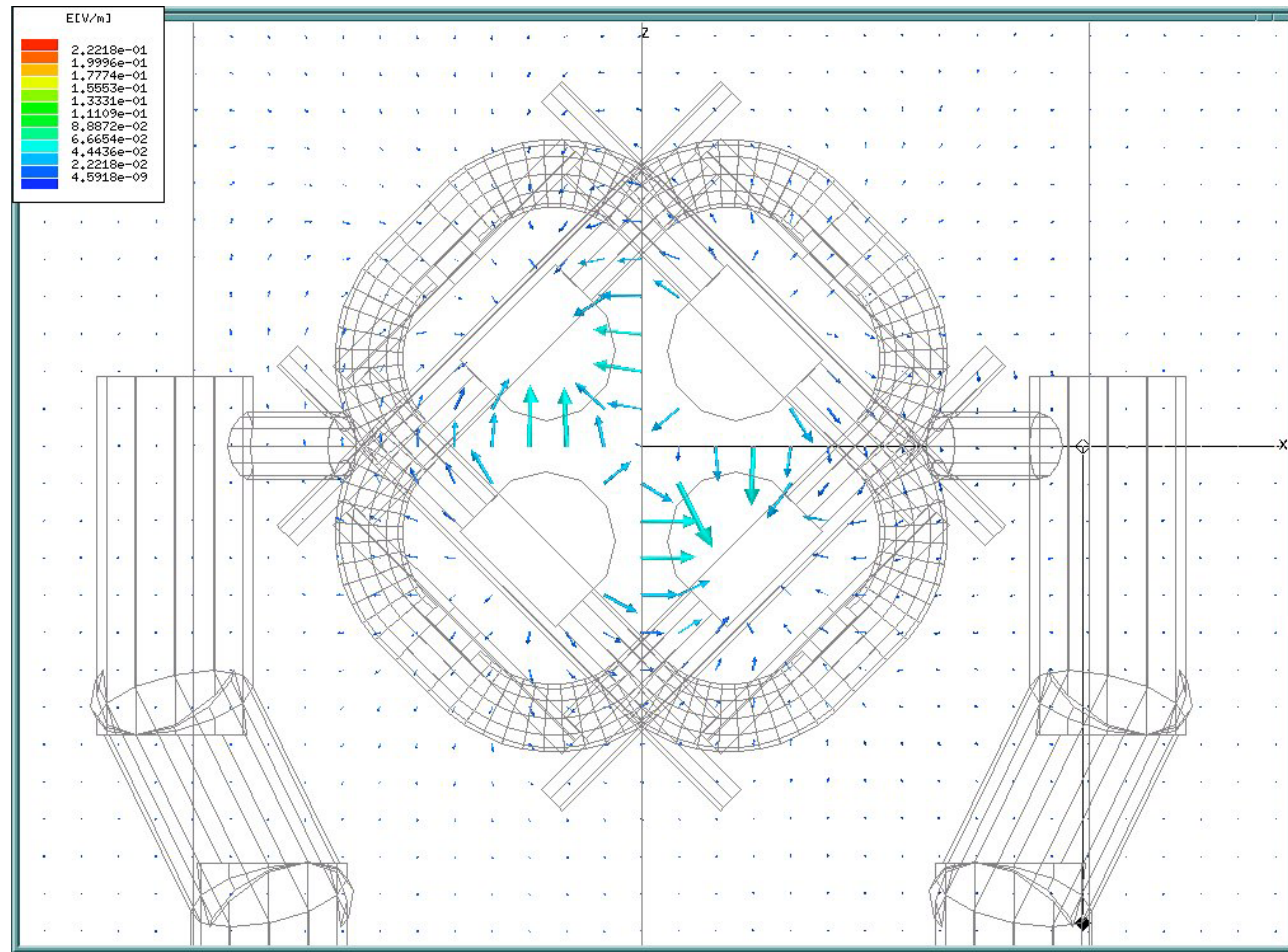
- ≡ Engineering Drawing being completed
- ≡ This model will be tested using protons
- ≡ Cooling to superfluid 1.6 K to avoid microphonics (no bubbles).
- ≡ Inductive RF feedthrough (3cm diam, center cooled).
- ≡ Two pickups to observe quadrupole and dipole resonances.
- ≡ 0.7 micron-thick biaxially-oriented monomar entrance window.

Some specifications

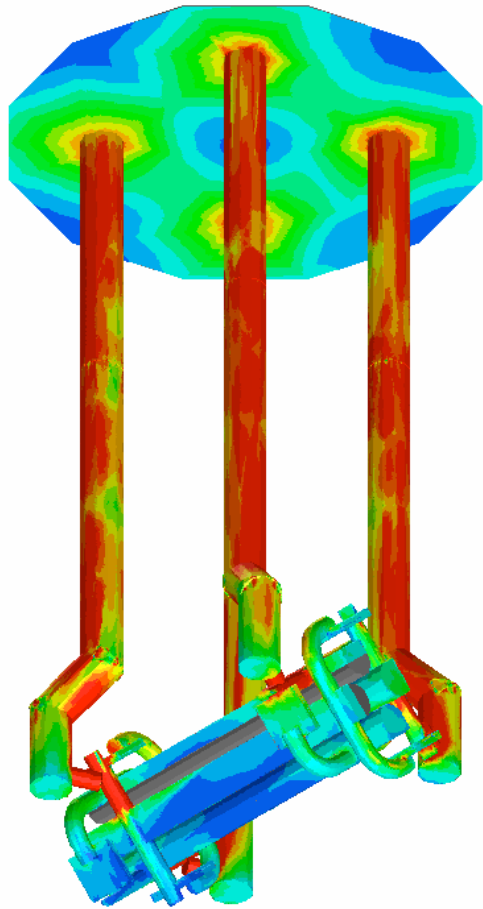


| | |
|------------------------------|-----------------------------|
| Radius: | 3 cm |
| Length: | 15 cm |
| RF drive: | 30 MHz |
| Fluctuations: | +/- 10 kHz |
| Maximum voltage: | 60 kV |
| Electrode capacitance: | 15 pF |
| Dipole cooling capacitance: | 300 pF |
| End ring electrodes voltage: | 6 kV |
| Pseudopotential: | 4.5 kV |
| q-value: | 0.85 |
| Secular frequency: | 11.4 MHz |
| (Anti)protons trapped: | 1 million per shot? |
| (Anti)proton density: | 4 million / mm ³ |
| Cloud diameter: | 40 micron |
| RF heating at surface: | 200 K |
| Q-value of cooling: | 10 ⁵ |
| Cooling time constant: | 10 seconds |

Detailed RF simulations done



Cooling (cryogenic) calculations

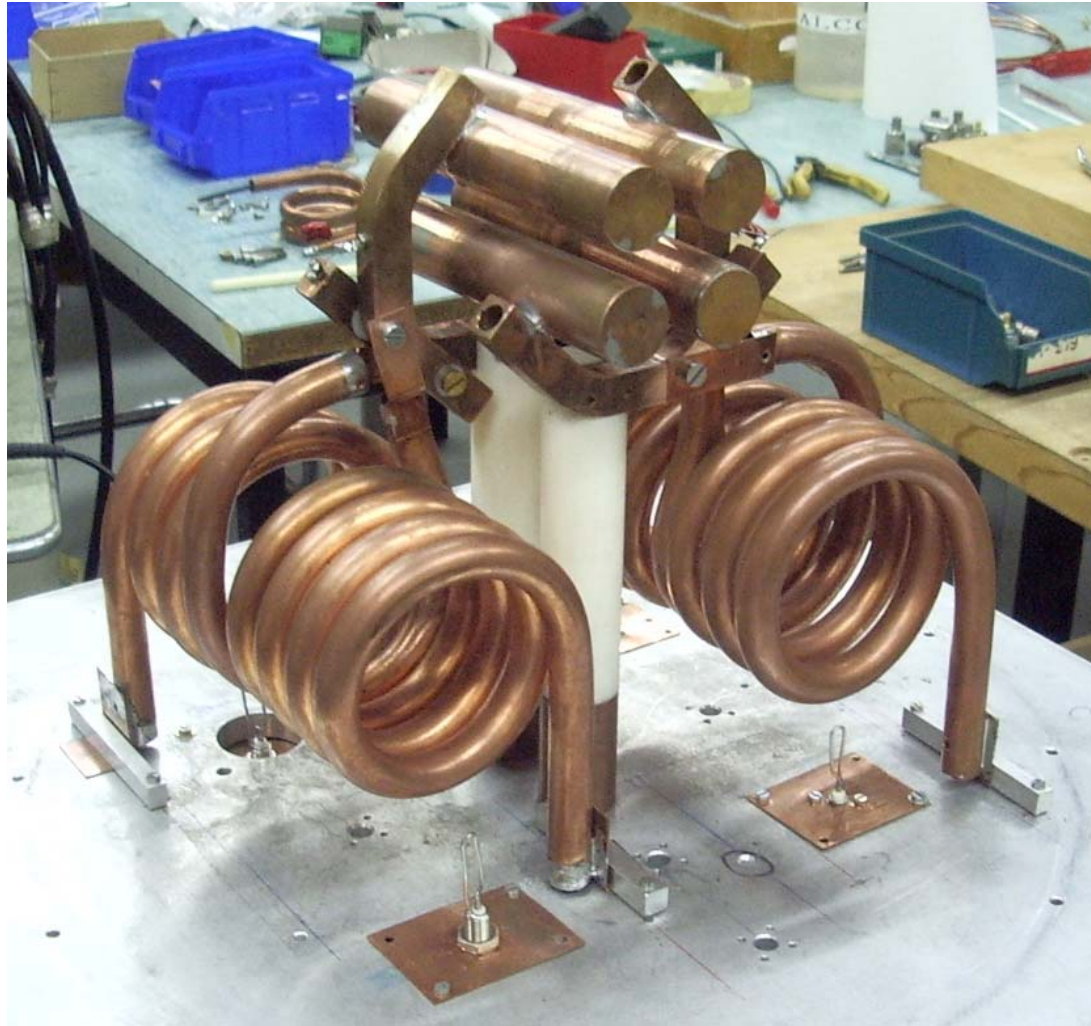


- ≡ Cool by flowing superfluid helium through stems.
- ≡ Simulation shows 0.5-1 W heat dissipation per stem.

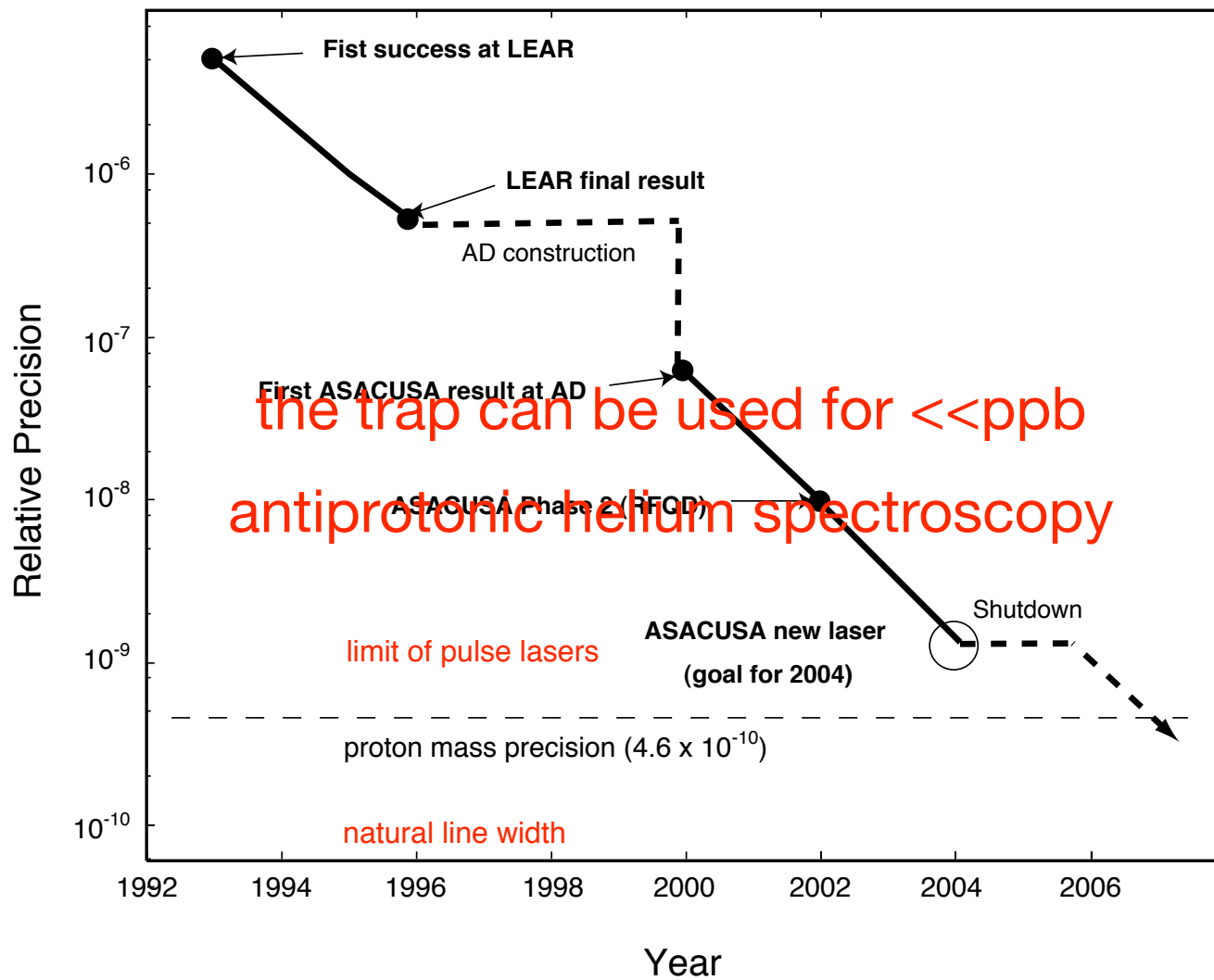
Copper model (version 3)



Copper model version 5



Larger L (inductance)
required for efficient
resistive cooling - now
use coils (but must
reduce microphonics)



to reach \ll ppb

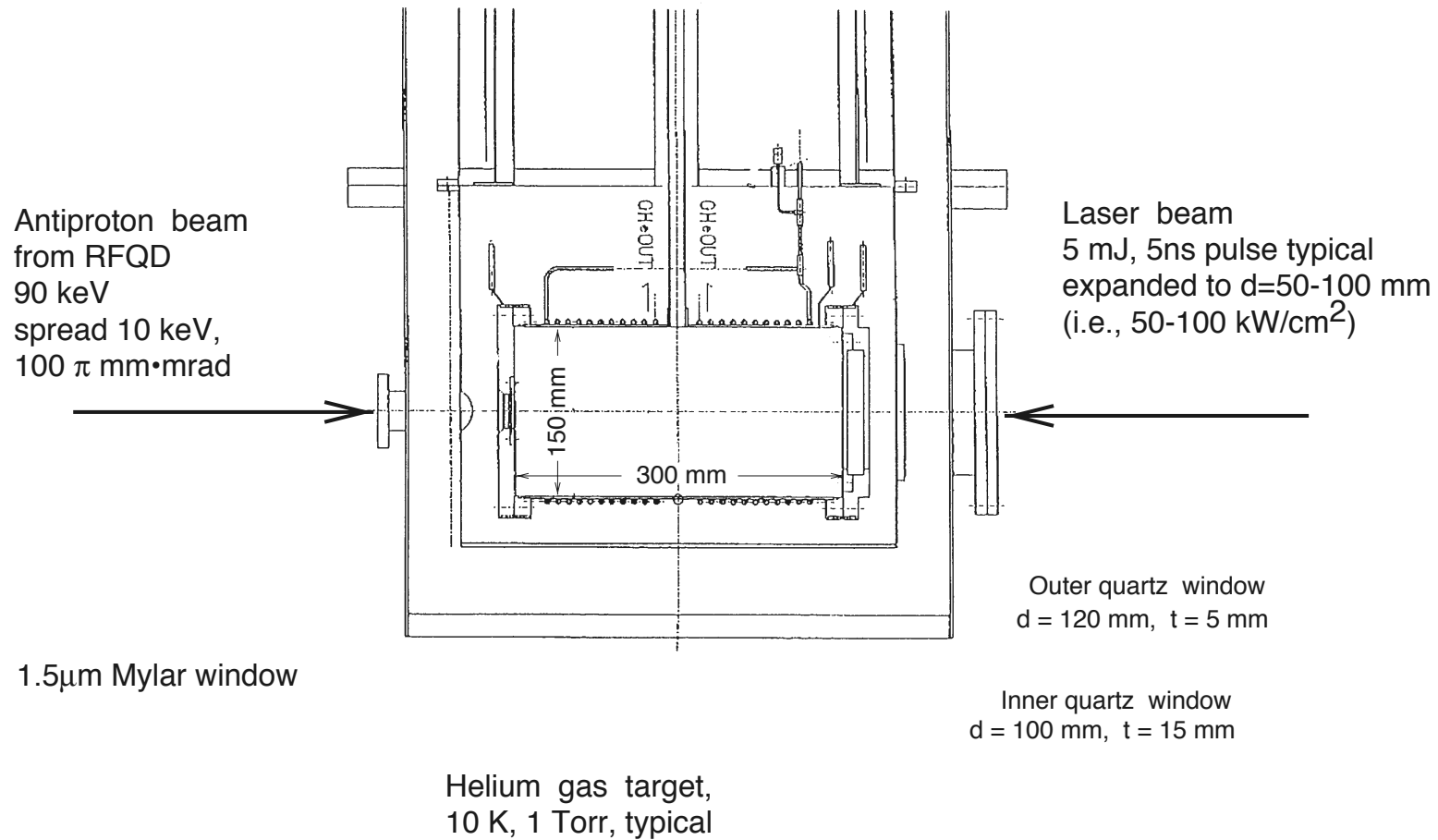


- ≡ Eliminate Doppler width (~ 500 MHz) using two-photon transitions
- ≡ Eventually, the use of pulsed laser has to be abandoned (chirp < 1 MHz is difficult)

the present target - HUGE



pulsed laser is unavoidable now



antiprotonic helium in a Paul trap



- ≡ with antiprotons compressed in the Paul trap (filled with dilute helium gas), antiprotonic helium atoms can be produced in a tiny volume
- ≡ then, CW-laser 2-photon spectroscopy will become possible
- ≡ natural line width ~ 160 kHz ~ 0.16 ppb, split the line by $1/100 \rightarrow 10^{-12}$ will be possible

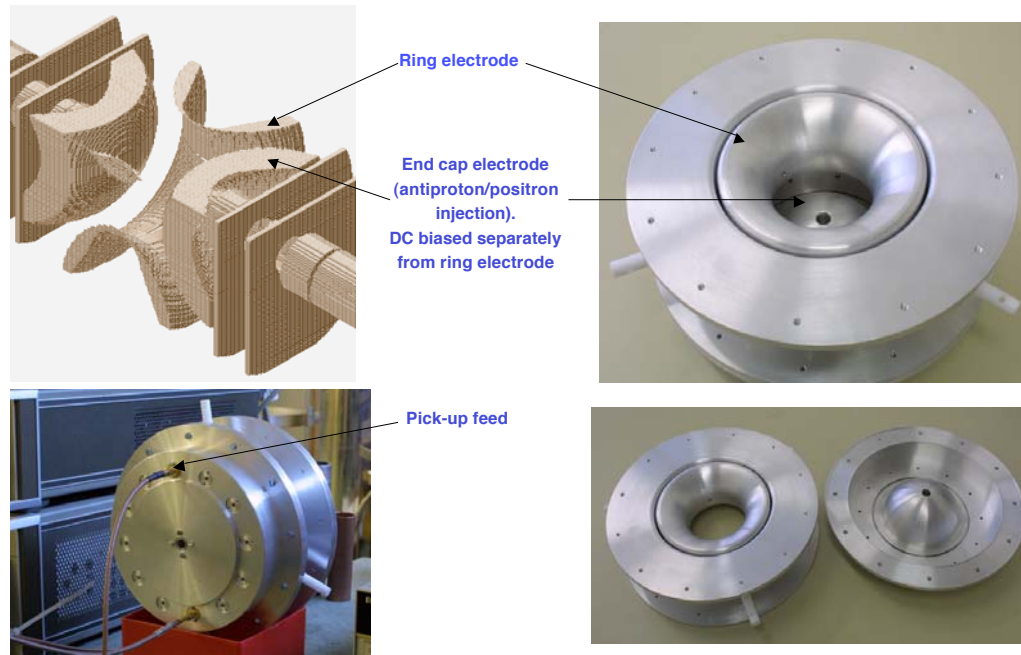
back to antihydrogen

how to make antihydrogen in a Paul trap?

Innovative two-tone trap

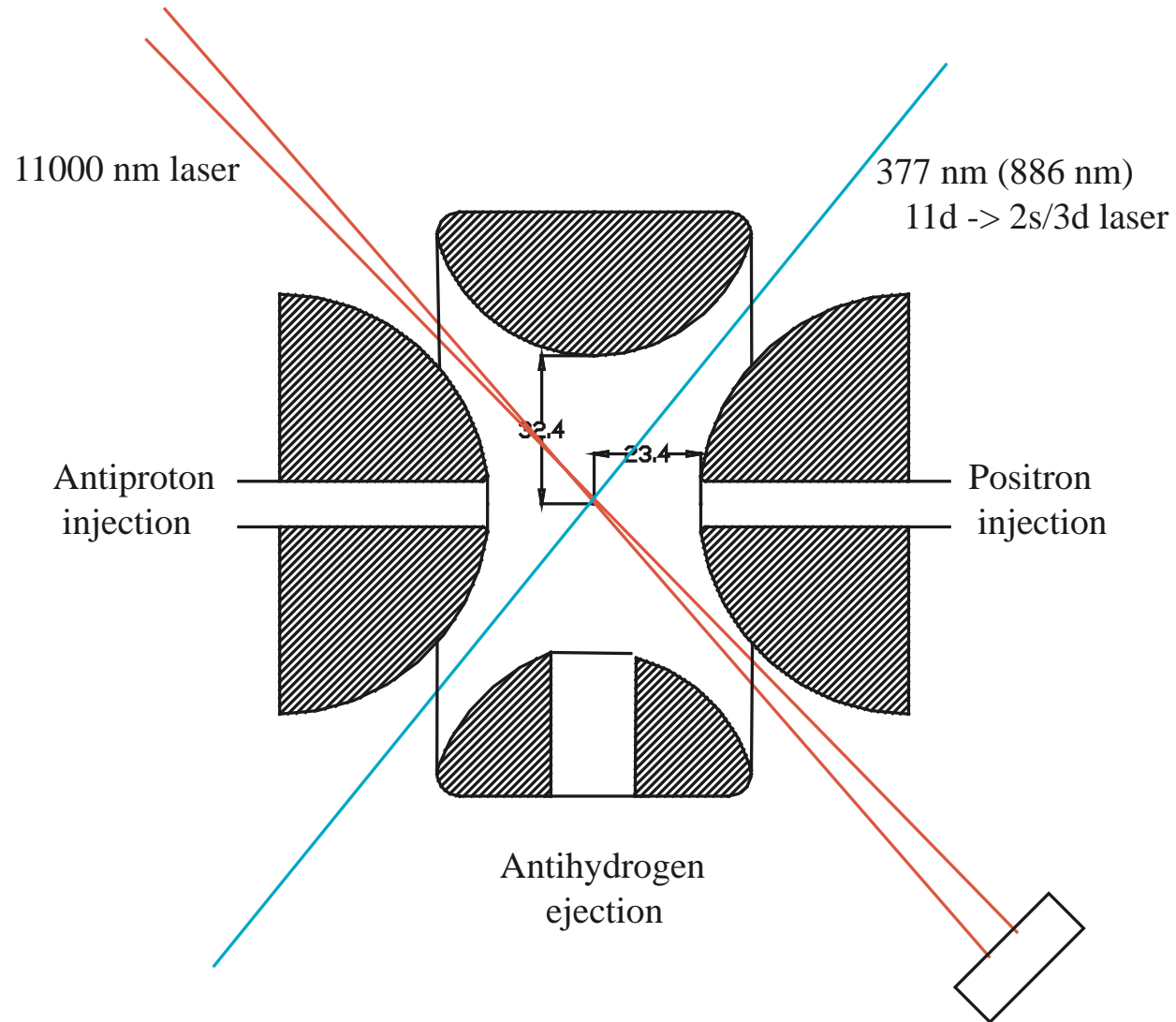


Antihydrogen production trap



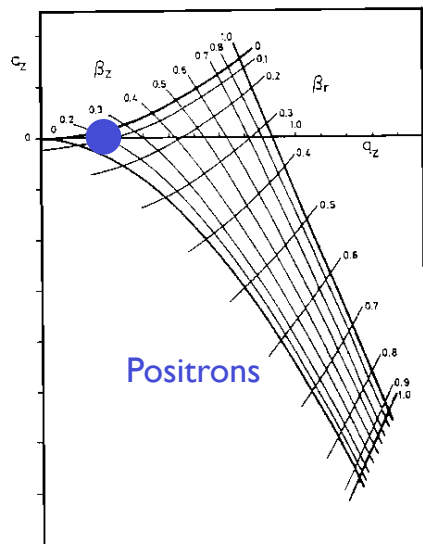
- ≡ two-frequency trap (3GHz to confine positrons, 2MHz for antiprotons)
- ≡ RF characterization done, mechanical & cryogenic implications studied

Two tone trap + laser



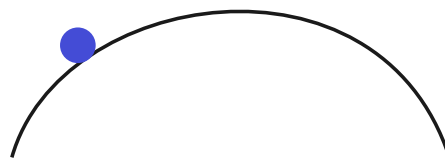


e^+ in the two-tone trap

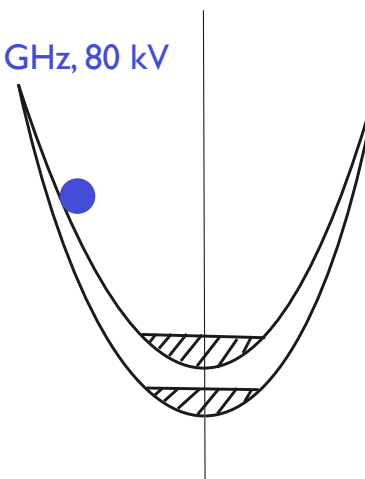


| | |
|-------------------------------|---------------------------------|
| Positron q-value: | 0.15 |
| Positron oscillatory a-value: | $\pm 1.7 \times 10^{-4}$ |
| Positron pseudopotential: | 2100-2200 V |
| Positron secular motions: | 78 \pm 0.4, 156 \pm 0.5 MHz |
| Number of positrons trapped: | 10000 |
| Positron density: | $5 \times 10^8 / \text{cm}^3$ |
| Positron cloud diameter: | 300 micron |
| Temperature at cloud surface: | 50 K |

RF=1 MHz, 40 V

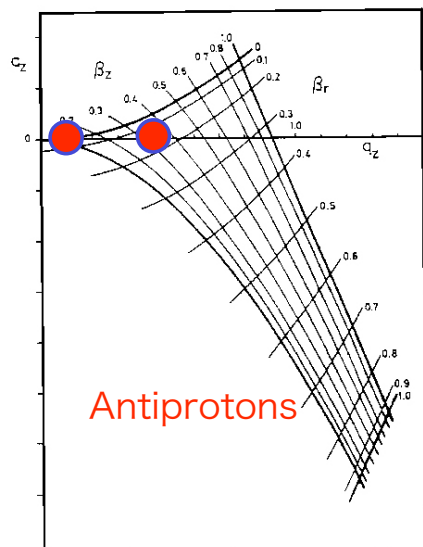


RF=3 GHz, 80 kV



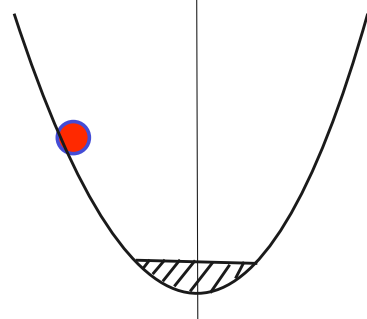


Pbar in the 2-tone trap

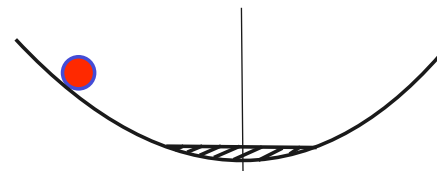


| | |
|--------------------------------|-------------------------------|
| Antiproton q-values: | $0.37, 8 \times 10^{-5}$ |
| Antiproton pseudopotential: | 3 V, 1.2 V |
| Number of antiprotons trapped: | 10000 |
| Antiproton density: | $7 \times 10^5 / \text{cm}^3$ |
| Antiproton cloud diameter: | 3 mm |
| Temperature at cloud surface: | 50 K |

RF=1 MHz, 40 V



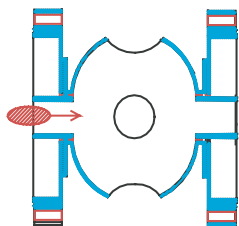
RF=3 GHz, 80 kV



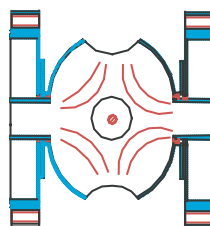
Antihydrogen production in the two-tone trap



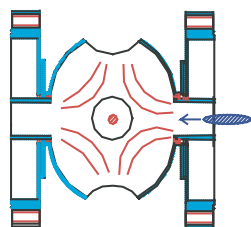
1): Antiproton injection



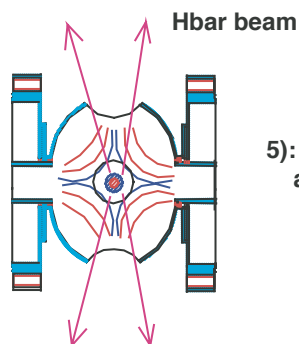
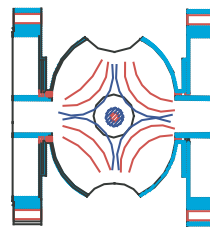
2): Antiproton trapping by 7 MHz RF field



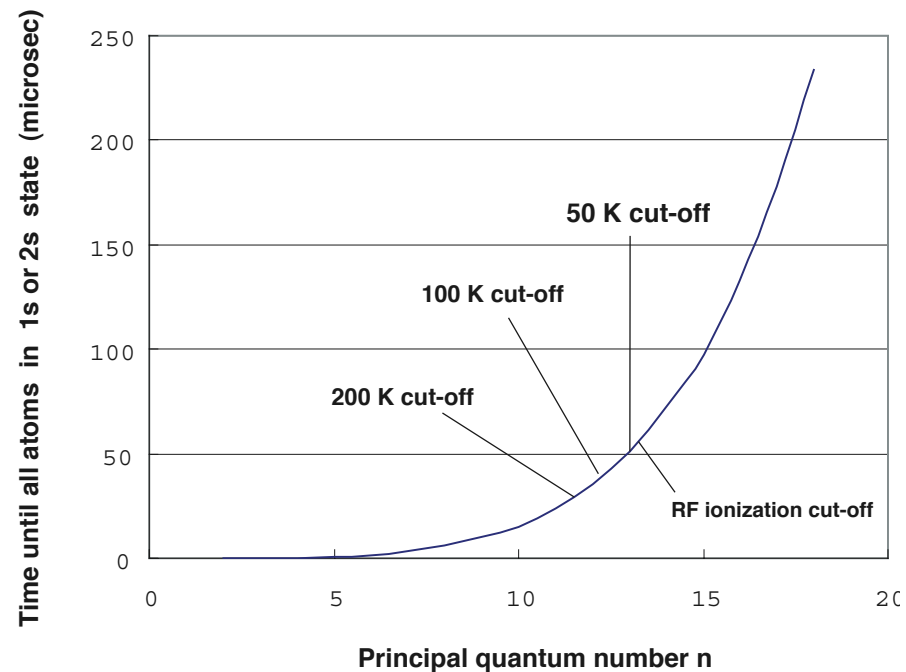
3): Positron injection



4): Positron trapping by 3 GHz RF field



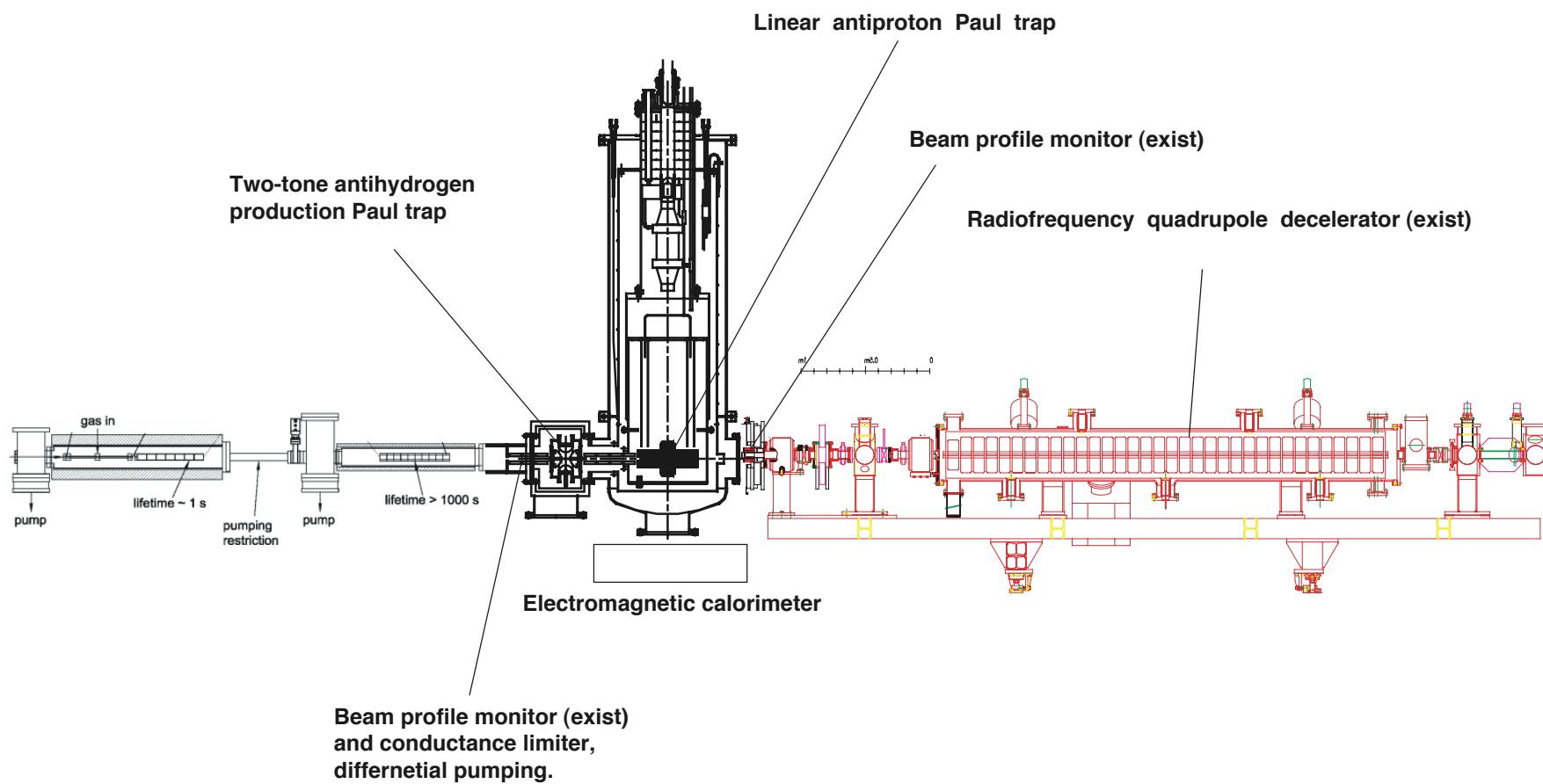
5): Antihydrogen production and extraction

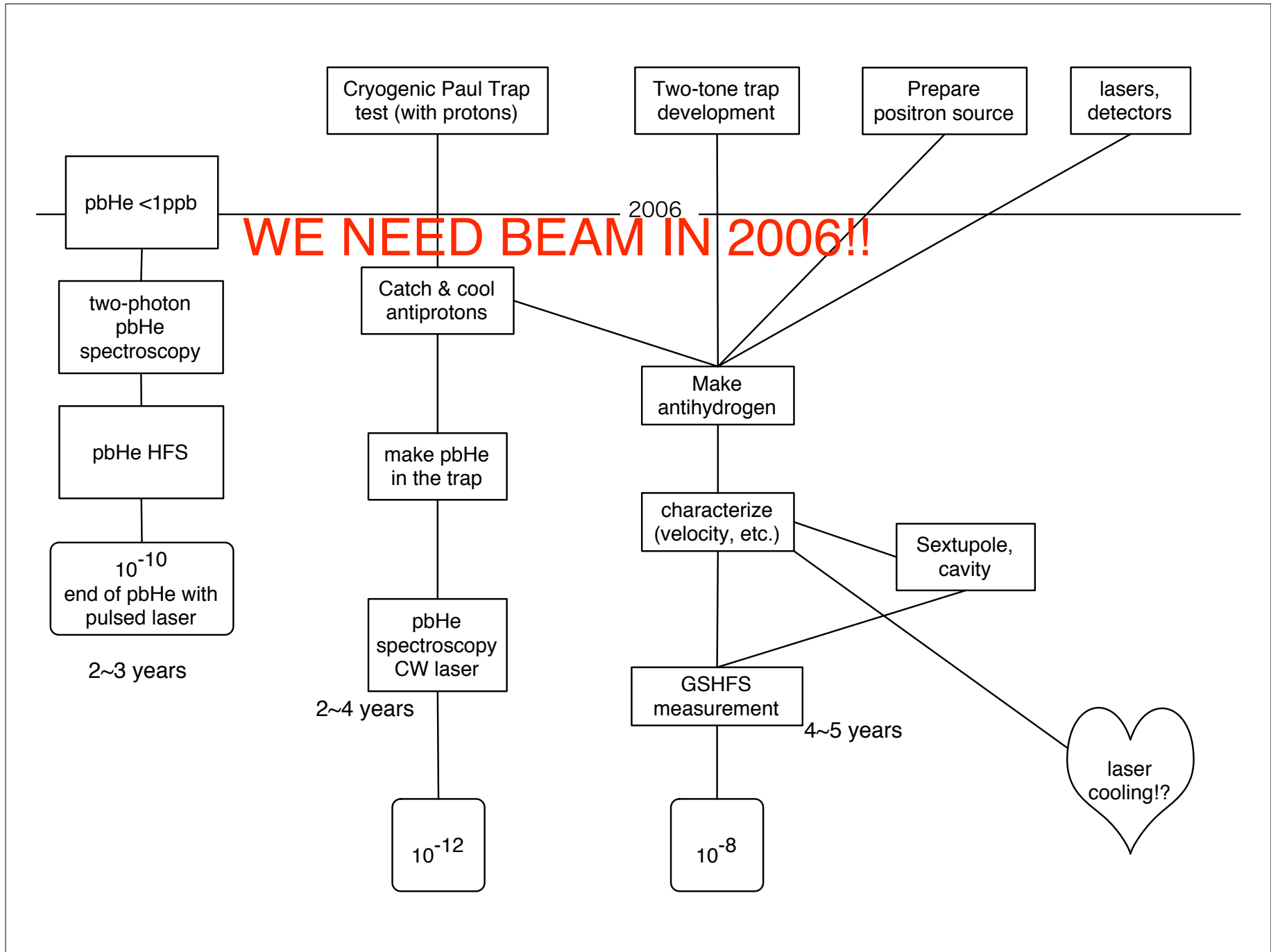


Only ground-state (or 2s) antihydrogen are emitted



A possible setup



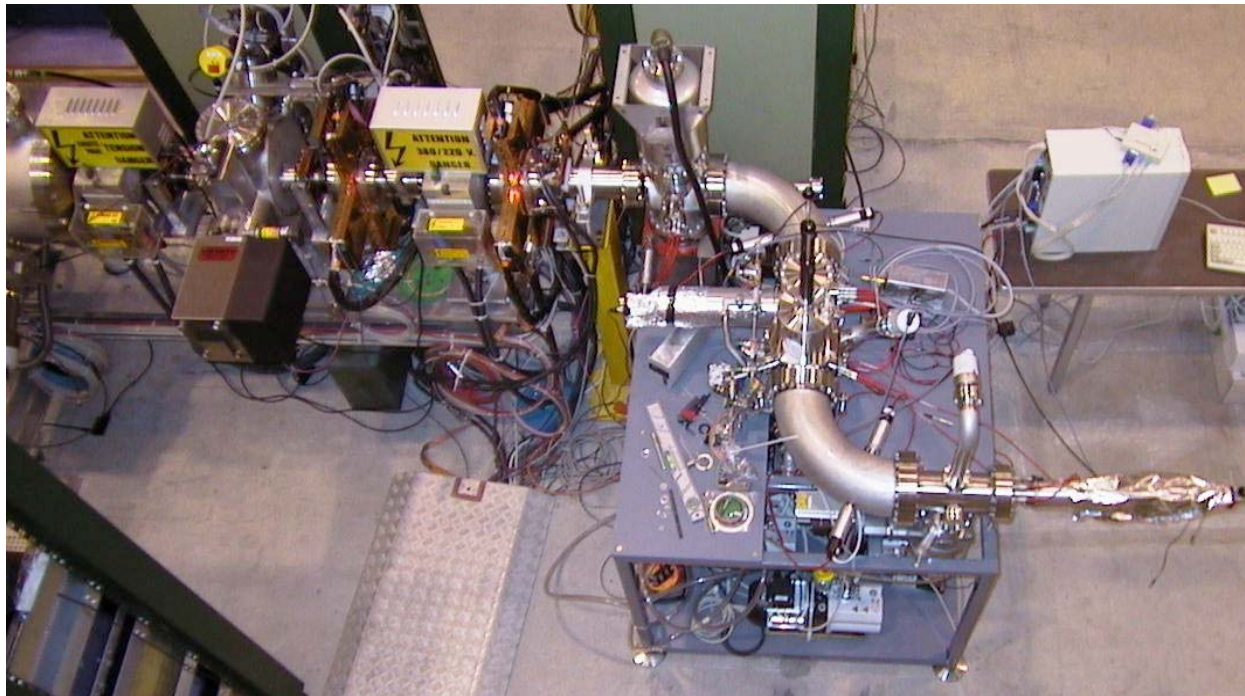
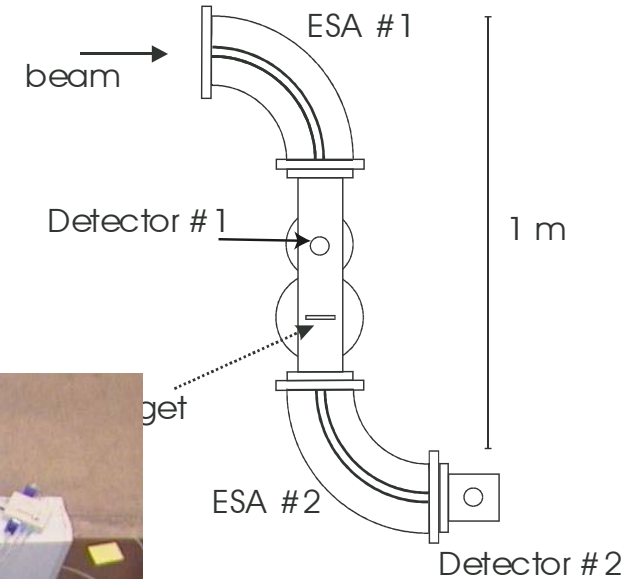


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Antiproton dE/dx measured to \sim keV



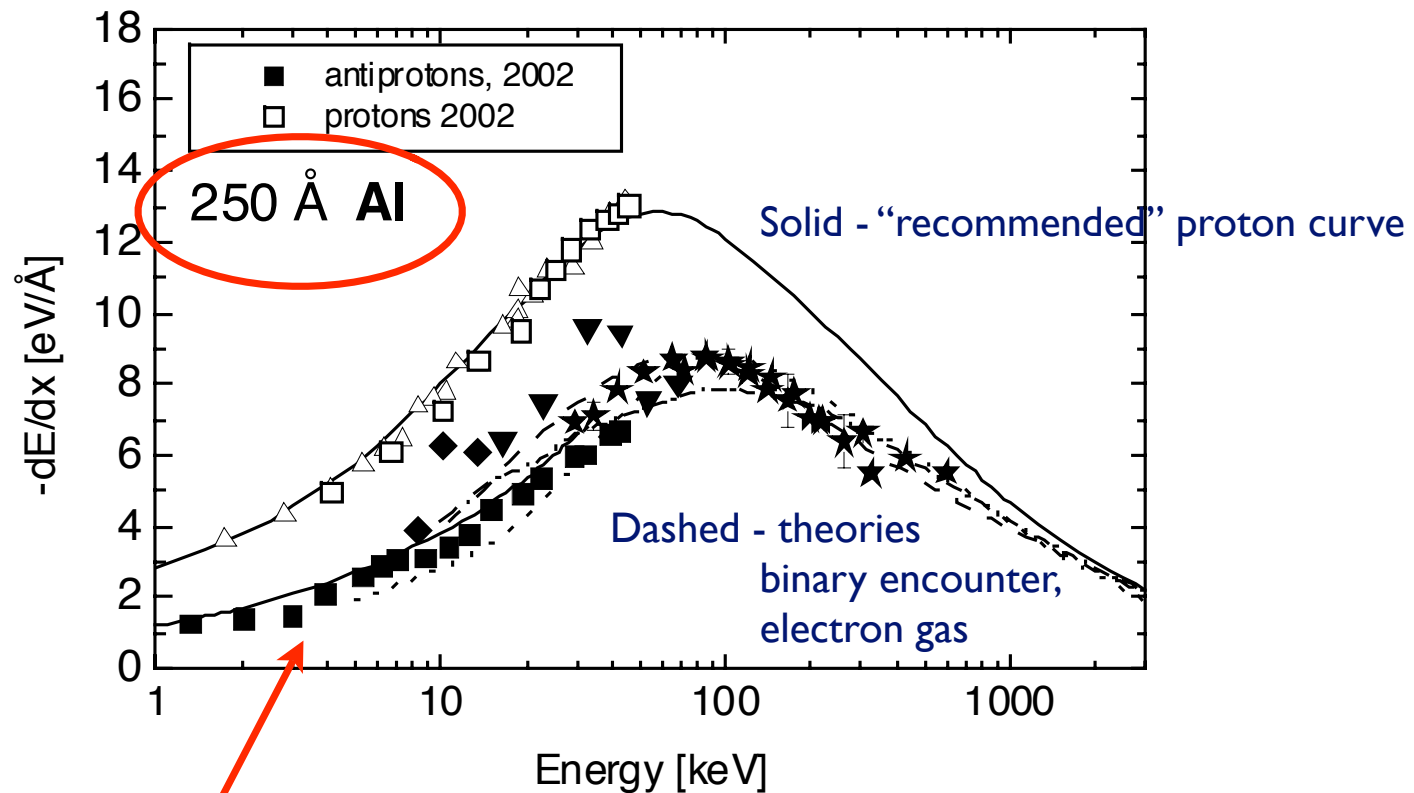
From RFQD



Electrostatic
Analyzer (ESA)



Antiproton dE/dx measured to \sim keV

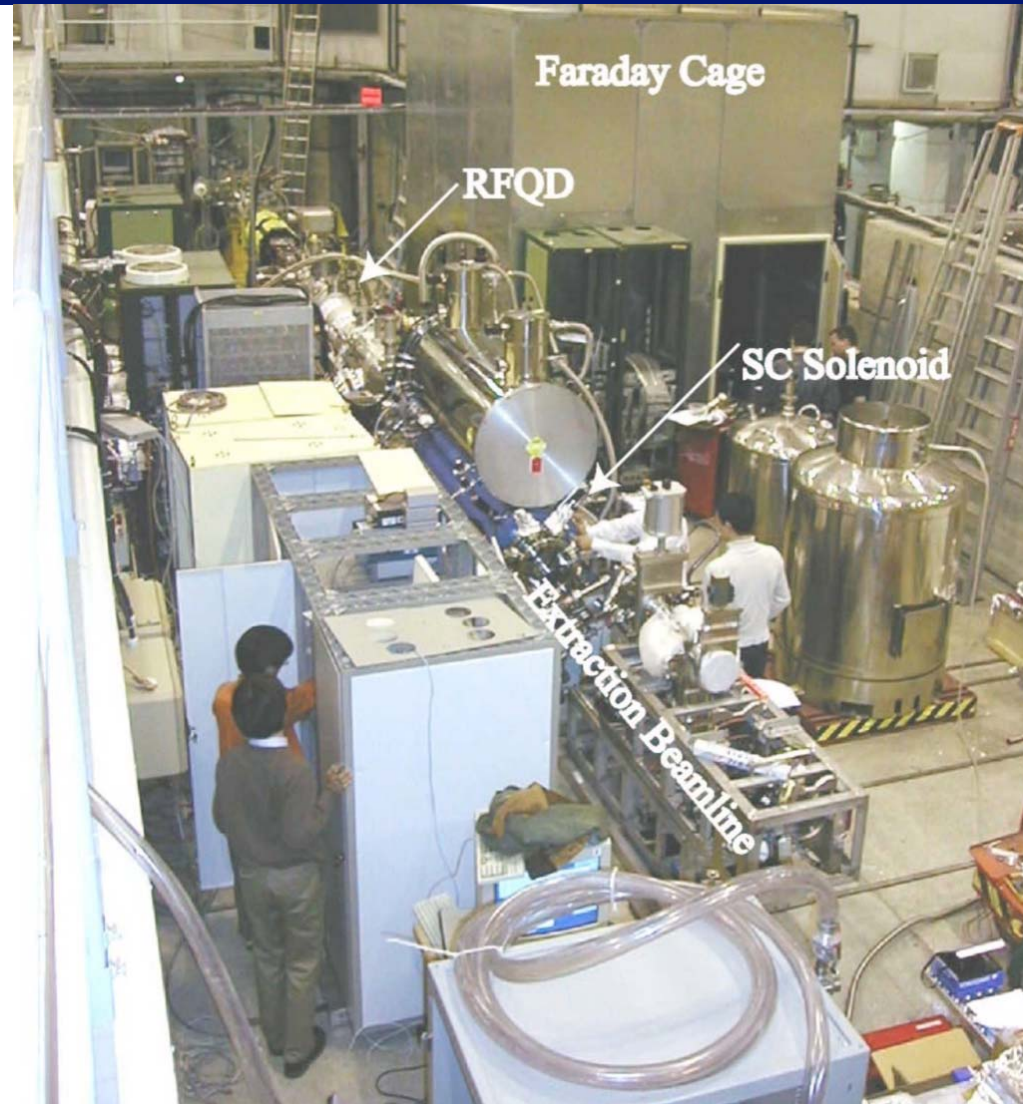


≡ Velocity proportionality at low energy (expected from electron gas model) evident in metal foils

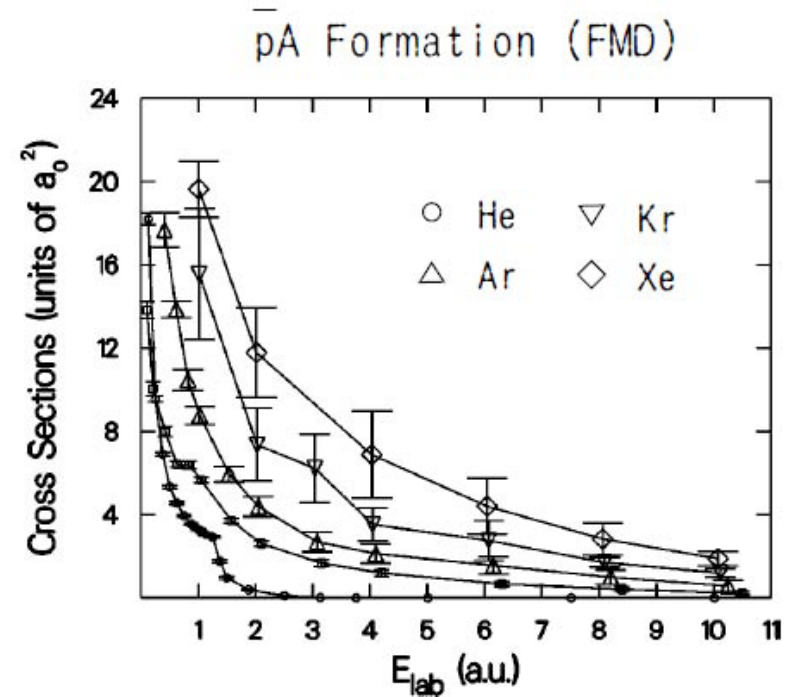
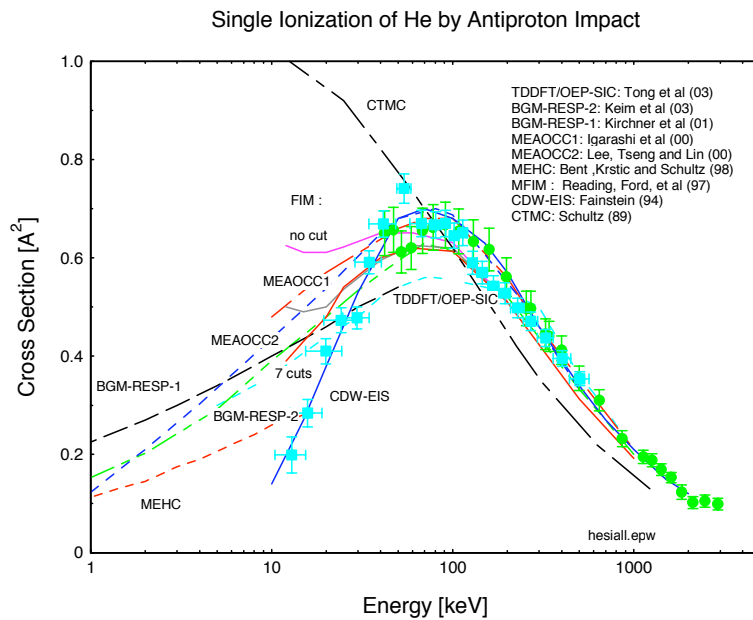
Phase 3 (use TRAP, <100 eV)



- ≡ Recently succeeded to extract a large number of antiprotons from the trap
- ≡ details by Y. Yamazaki



Ionization, antiprotonic atom formation



In the ASACUSA original proposal, but were deferred,
waiting for the phase-3 beam development

“phase 3 beam”



- ≡ some experiments planned for FLAIR (slow extraction) may already be explored at AD

Summary

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| <p>CPT tests:</p> <p>high-precision determination of antiproton mass, charge, magnetic moment and magnetic structure using various spectroscopic methods</p> | 3-body system | Antiprotonic helium atom laser & microwave spectroscopy | RFQD +low-density target | <p>New high-precision laser system.</p> <p>Two-photon spectroscopy will enable ultimate accuracy.</p> | <p>Proton mass is known to 0.46 ppb.</p> <p>The goal is to measure antiproton mass to a similar precision.</p> |
| | 2-body systems | Antiprotonic helium ion laser spectroscopy | RFQD +low-density target | Free from theoretical ambiguities. | <p>The magnetic properties of antiproton poorly known.</p> <p>The goal is to compare the magnetic structure of proton and antiproton.</p> |
| | | Antihydrogen ground-state HFS microwave spectroscopy | RFQD + Paul trap + Two-tone Paul trap (point source of cold antihydrogen) | Superconducting Paul traps being developed first Hbar production test in 2006. | |
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promising future**