

Modified BK and heavy quarks

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1. Motivations

Reliable description of processes towards the region of **very small** values of x and **large** k^2



- **Constrains** on parton distributions implied by HERA data.
- **BFKL** dynamics with subleading corrections.
- Complete **DGLAP** evolution i.e. large k^2 .
- **Non-linear** screening effects built into the evolution equations.
- **Mass** of heavy quark justifies application of presented formalism

2. Extended Balitzkij-Kovchegov (BK) equation

Nonlinear evolution equation for unintegrated gluon distribution.

$$f(x, k^2) = \tilde{f}^{(0)}(x, k^2) + K^1 \otimes f - K^2 \otimes f^2 \quad (1)$$

where

- $\tilde{f}^{(0)}(x, k^2) \rightarrow$ input
- $K^1 \otimes f \rightarrow$ BFKL + running coupling + kinematic constraint on gluon emission $k'^2 < \frac{k^2}{z}$ (dominant subleading BFKL effects) and inclusion of DGLAP evolution
- $K^2 \otimes f^2$ term that takes care for **saturation** effects

$$K^2 \otimes f^2 = \left(1 - k^2 \frac{d}{dk^2}\right)^2 \frac{k^2}{R^2} \int_x^1 \frac{dz}{z} \left[\int_{k^2}^{\infty} \frac{dk'^2}{k'^4} \alpha_s(k'^2) \ln \left(\frac{k'^2}{k^2}\right) f(z, k'^2) \right]^2$$

General structure of modified BK comes from BK under assumption that target is uniform and large in comparison to projectile.

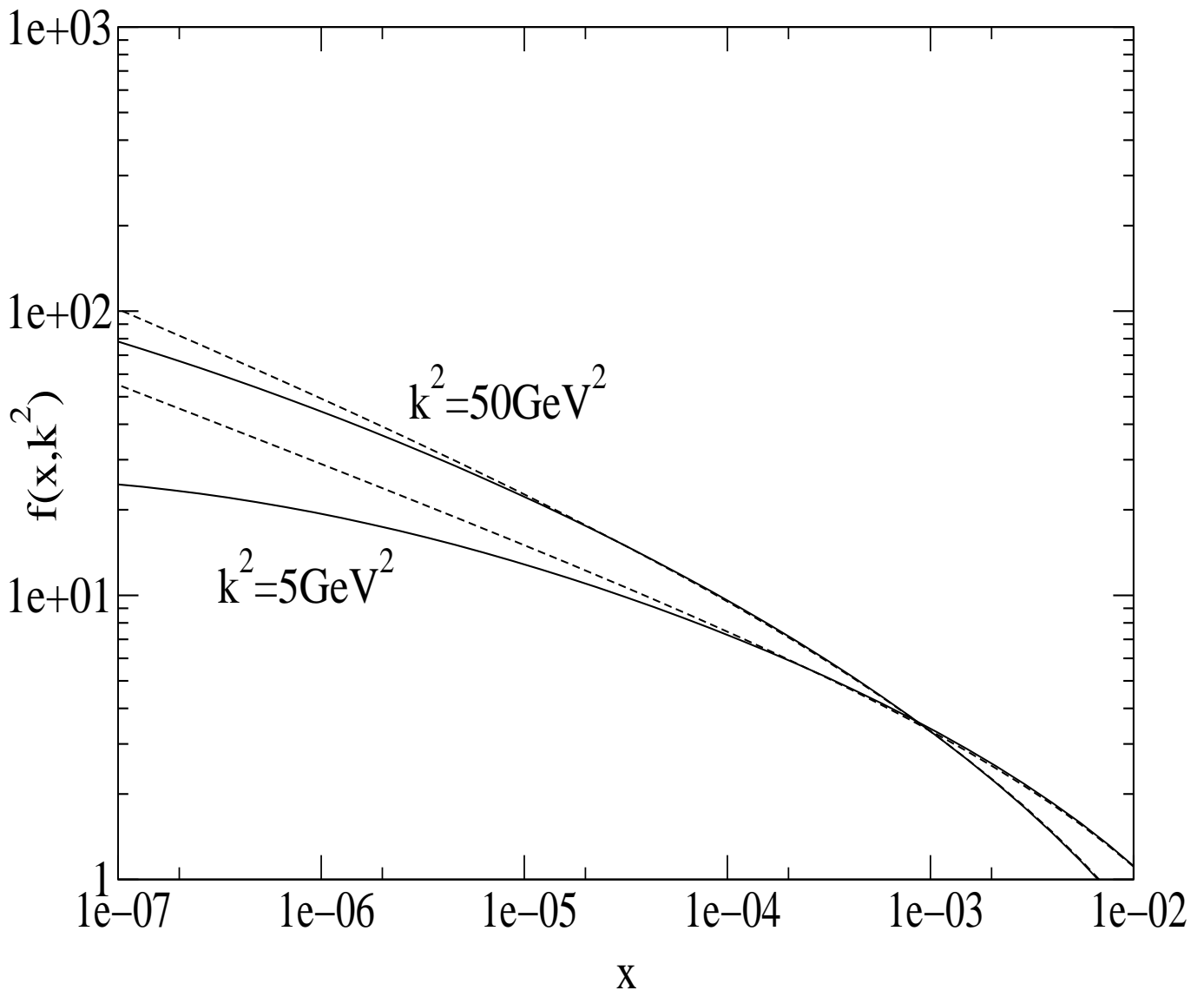


Figure 1: Linear equation vs nonlinear

Origin of suppression of $f(x, k^2)$

- At very small x the **linear** (DGLAP or BFKL) evolution generates strong increase of the gluon distributions for $x \rightarrow 0$ which eventually violates unitarity.
- Non-linear screening corrections tame the growth at small x .
- Emergence of the **saturation scale** $Q_s^2(x) \sim x^{-\lambda}$
- Non-linear effects in the unintegrated gluon distributions are small (eventually negligible) for $k^2 > Q_s^2(x)$.
- Non-linear effects are very strong \rightarrow **saturation** for $k^2 < Q_s^2(x)$
- For $k^2 \ll Q_s^2(x)$, $f \ll f^{linear}$.

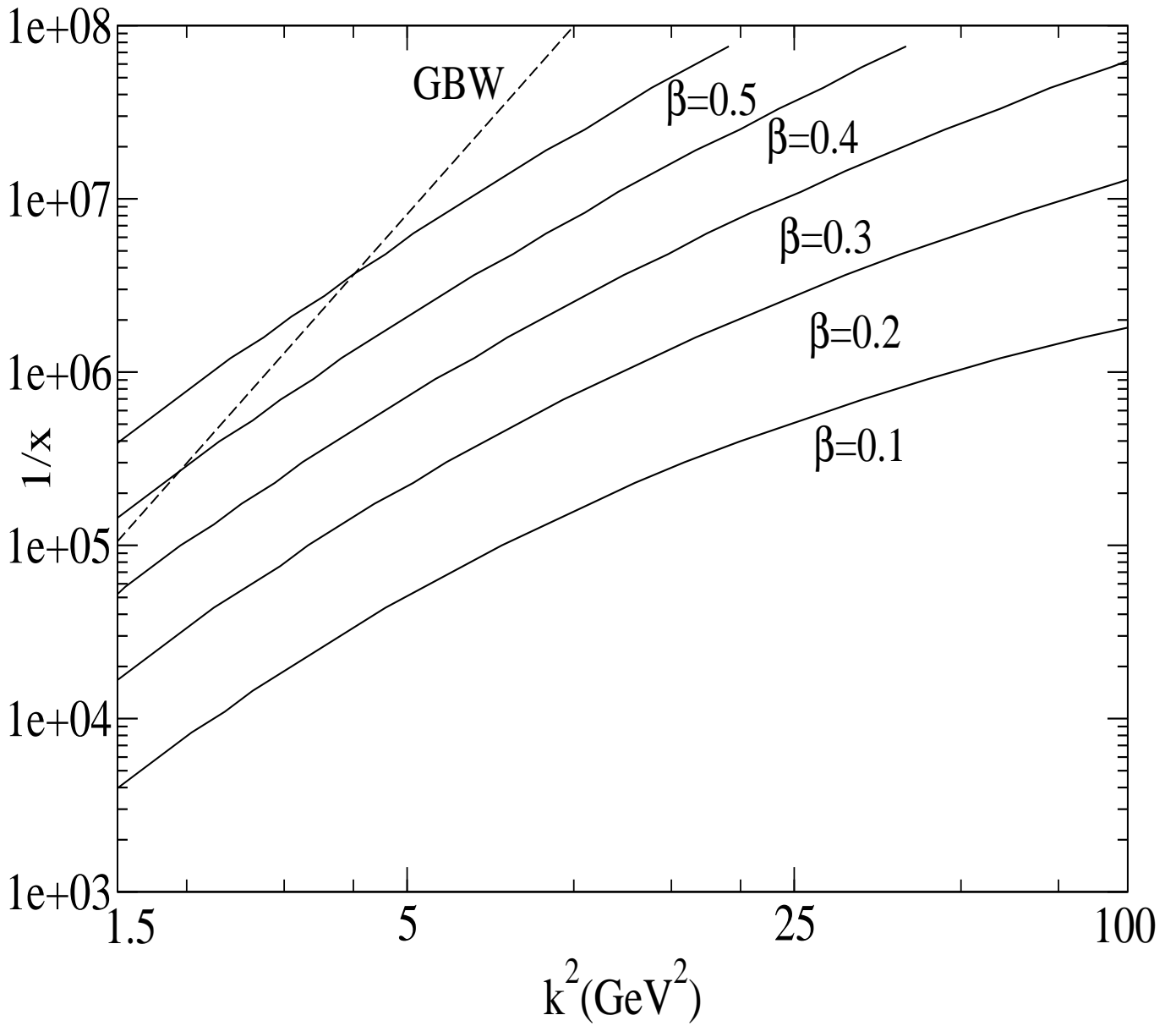
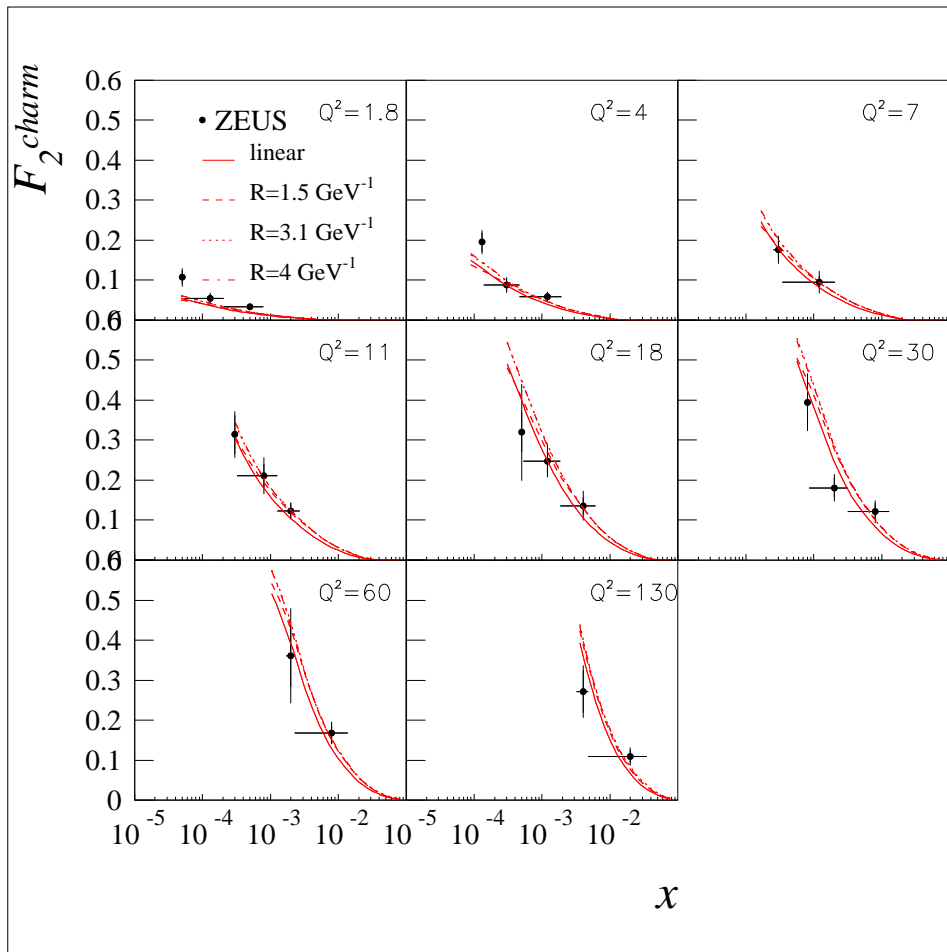


Figure 2: Phenomenologically motivated critical line, $\beta = \frac{f_l(x, k^2) - f_{nl}(x, k^2)}{f_l(x, k^2)}$

4. Heavy quark production

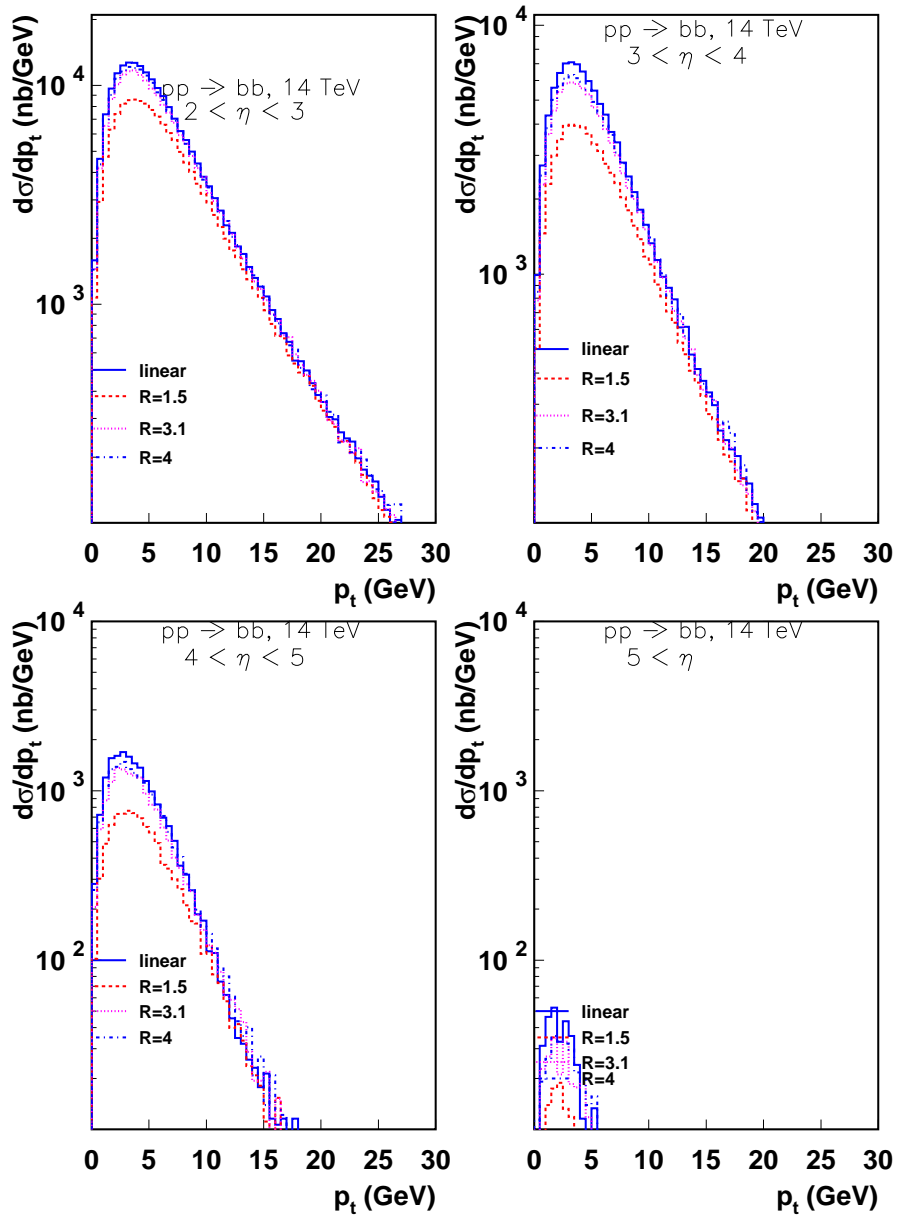
HERA, Modified BK and Cascade give us

$$F_2^{cc}(x, Q^2):$$



For LHC we estimate:

- σ_{diff} for $pp \rightarrow b\bar{b}$



For LHC we estimate:

- $\frac{\sigma_{tot}(R=3.1\text{GeV}^{-1})}{\sigma_{tot}(linear)} = .97$
- $\frac{\sigma_{tot}(R=1.5\text{GeV}^{-1})}{\sigma_{tot}(linear)} = .8$

5. Conclusions

Presented formalism describes data and allows us to predict cross section for production of heavy quarks in situation when partons form very dense system and eventually **saturate**. Saturation is also required to unitarise cross section which is violated by BFKL.