

Parton Recombination at all p_T

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Hard Probes 2004

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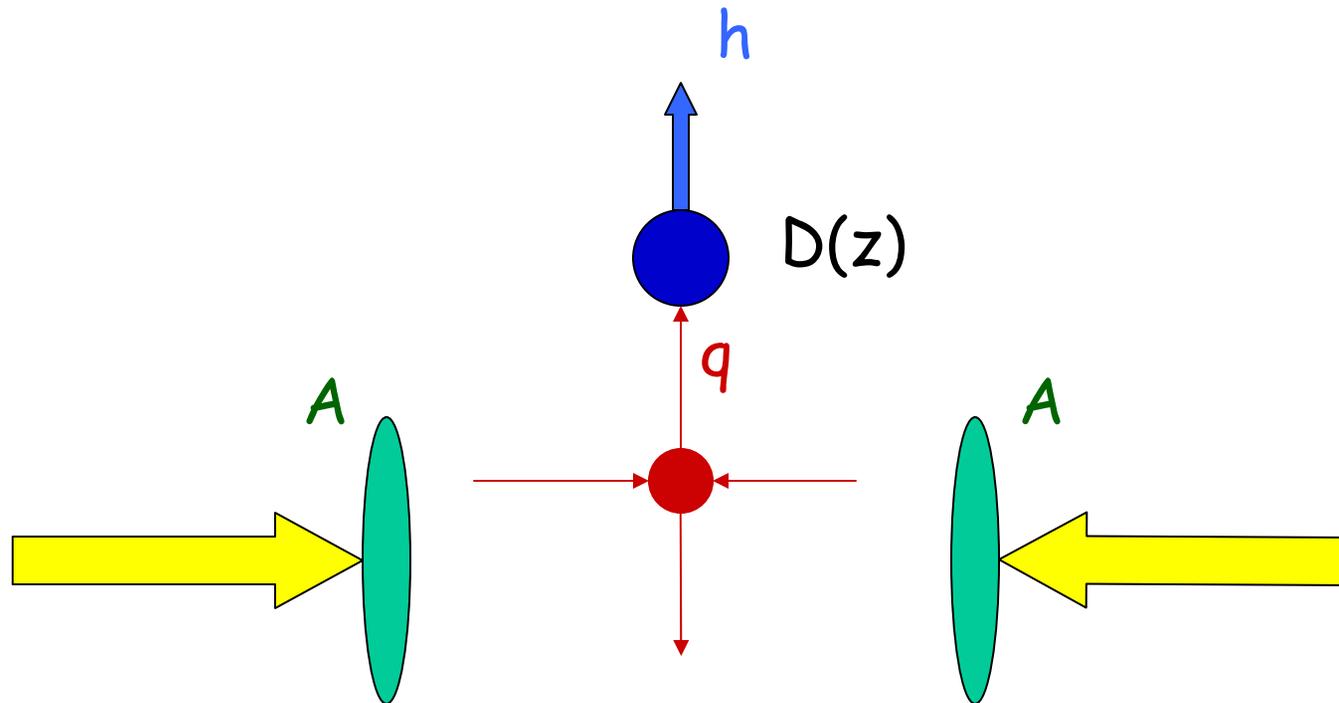
Outline

Why fragmentation is inadequate at intermediate p_T in heavy-ion collisions.

How parton recombination resolves a number of puzzles.

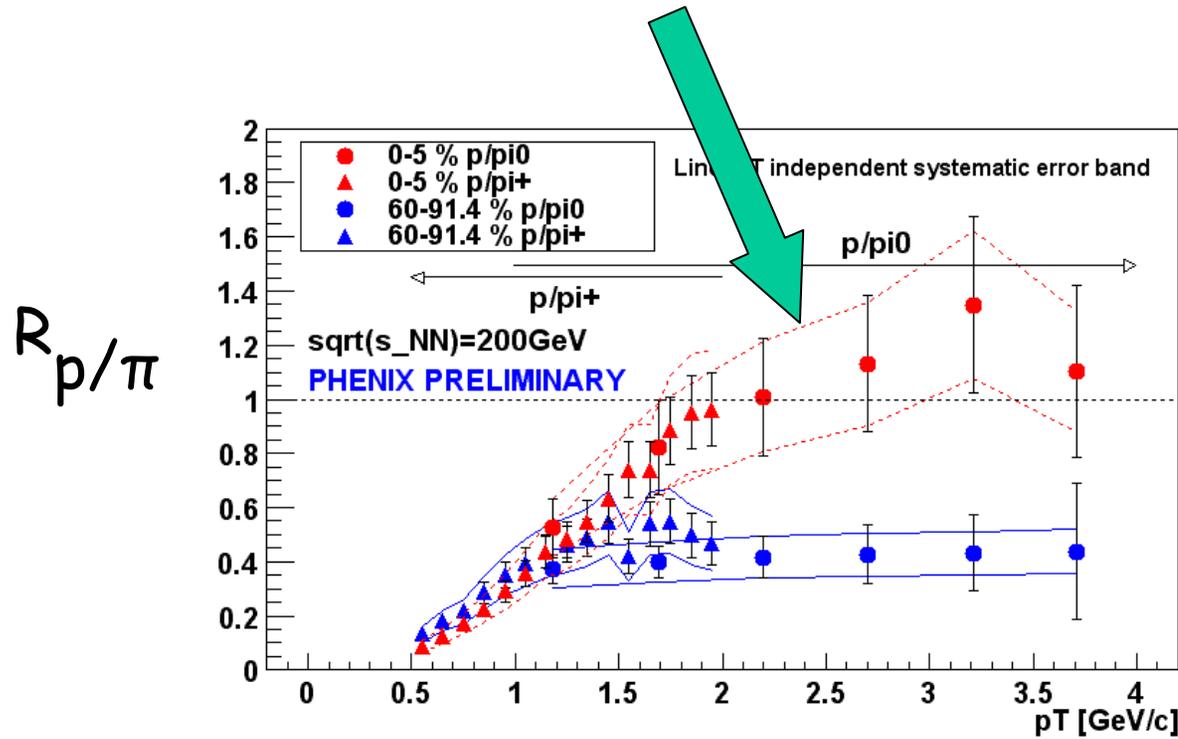
- Shower partons
- Inclusive distributions at all p_T in Au+Au and d+Au collisions (Cronin effect)
- Dihadron correlations
- Forward production in d+Au collisions

Conventional approach to hadron production at high p_T



There are numerous evidences against this framework: hard scattering \times fragmentation in heavy-ion collisions.

Exhibit #1 $R_{p/\pi} > 1$



Not possible in fragmentation model:

$$D_{p/q} \ll D_{\pi/q}$$

$$\frac{D_{p/q}}{D_{\pi/q}}$$

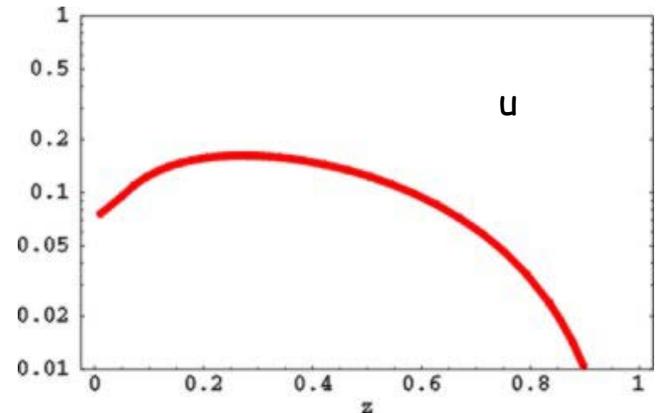
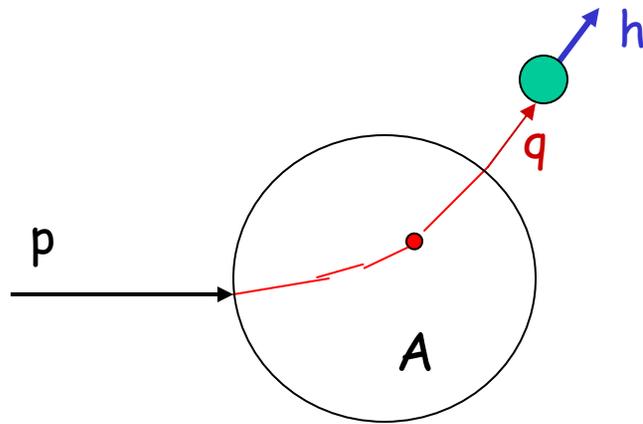


Exhibit #2 in pA or dA collisions

Cronin Effect

Cronin et al, Phys.Rev.D (1975)



$$\frac{dN}{dp_T} (pA \rightarrow \pi X) \propto A^\alpha, \quad \alpha > 1$$

k_T broadening by multiple scattering in the initial state.

Unchallenged for ~30 years.

If the medium effect is before fragmentation, then α should be independent of $h = \pi$ or p

$$\boxed{\alpha_p > \alpha_\pi} \longleftrightarrow \boxed{R_{CP}^p(p_T) > R_{CP}^\pi(p_T)}$$

$$R_{CP}^p(p_T) > R_{CP}^\pi(p_T)$$

PHENIX d+Au PRELIMINARY

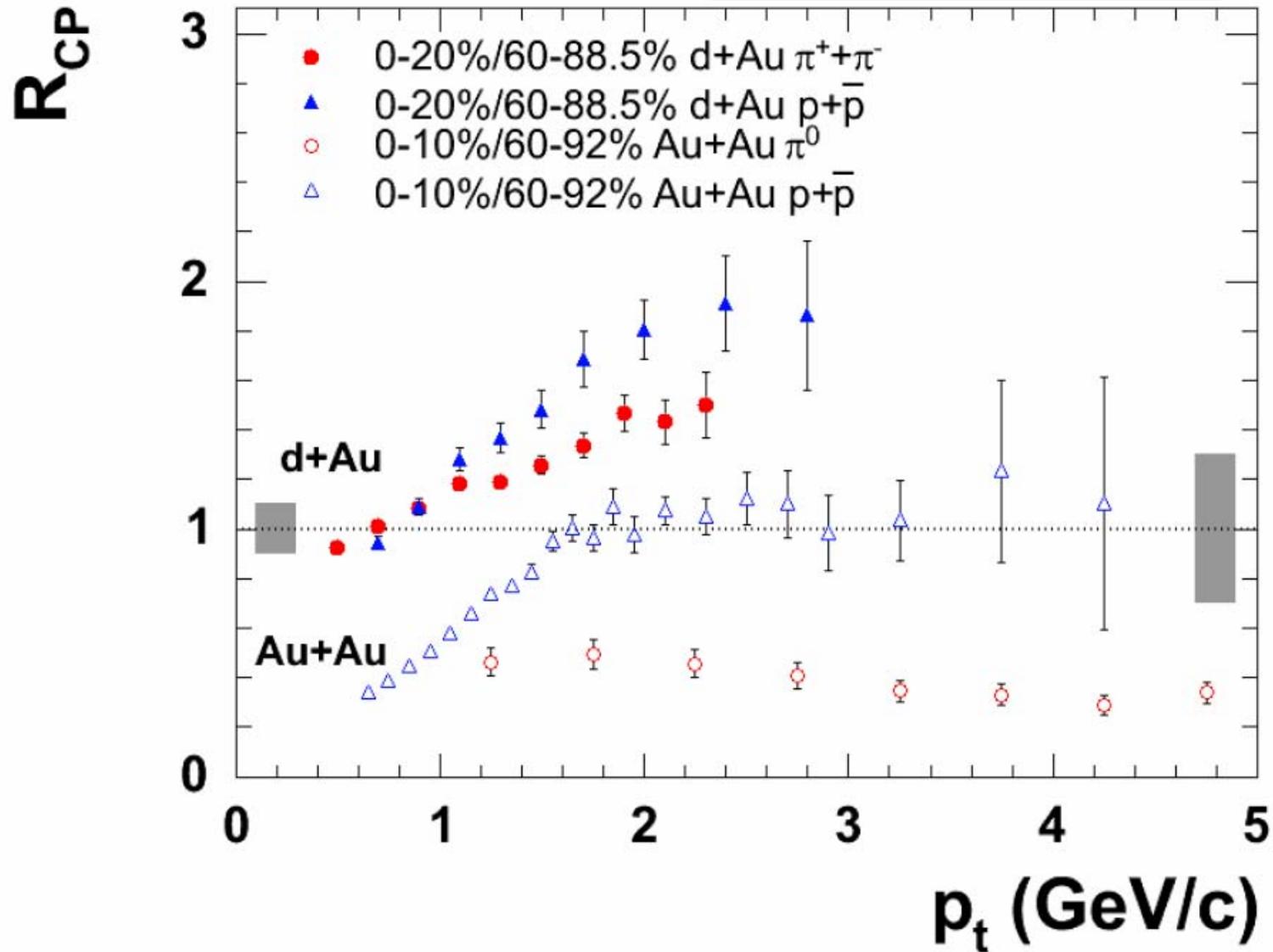
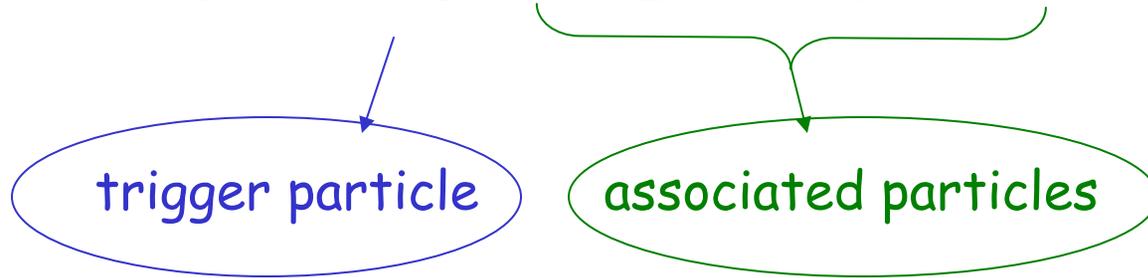
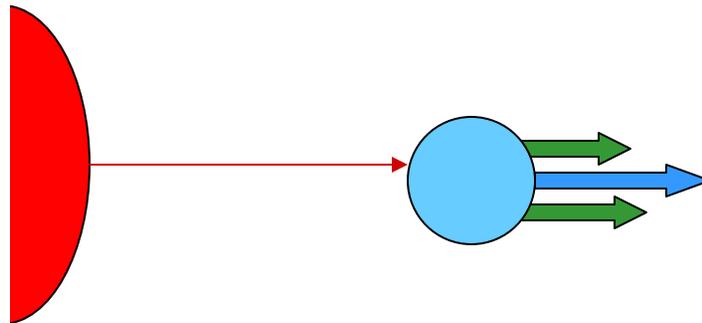


Exhibit #3 Jet structure

Hard parton \Rightarrow jet $\{ \pi(p_1) + \pi(p_2) + \pi(p_3) + \dots \}$



The distribution of the associated particles should be *independent* of the medium if fragmentation takes place in vacuum.



Fuqiang Wang (STAR) Quark Matter 2004

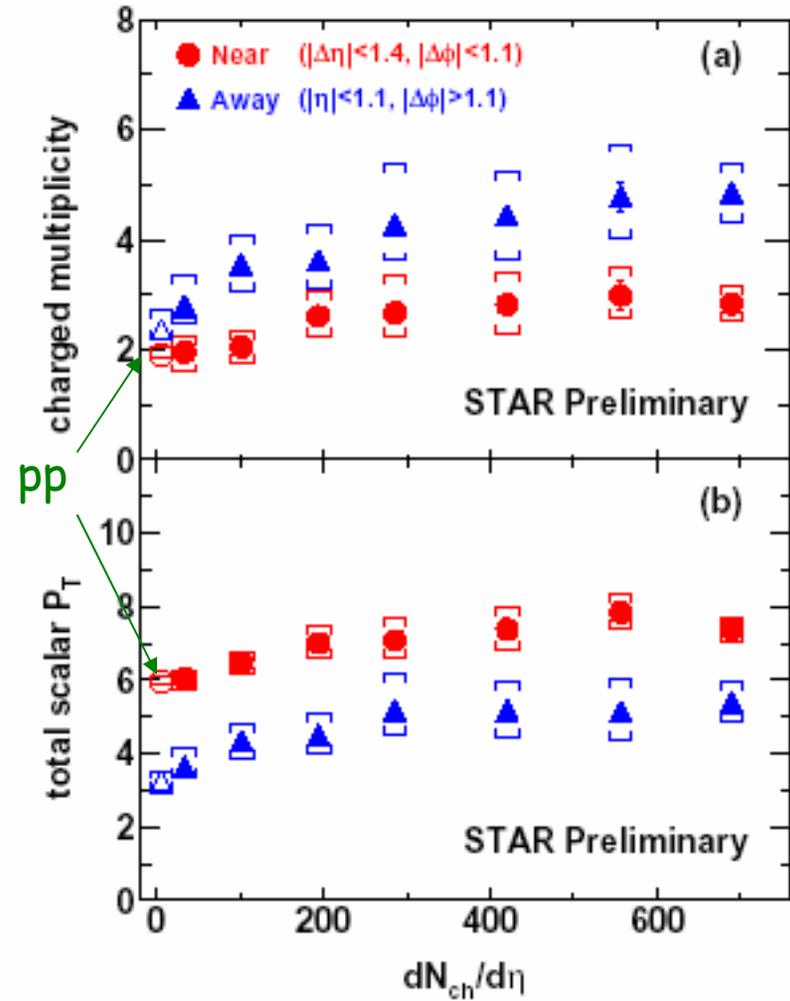
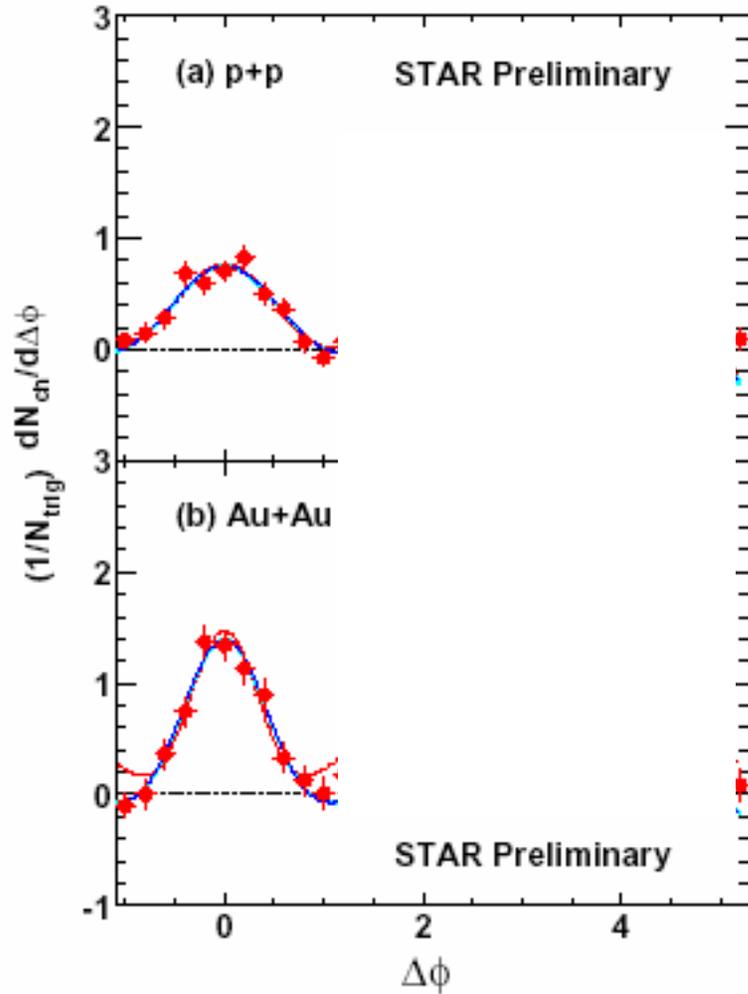
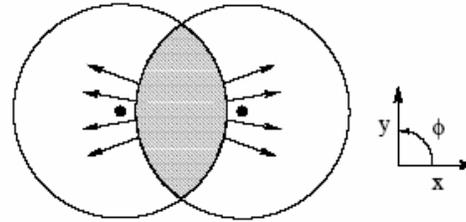


Exhibit #3

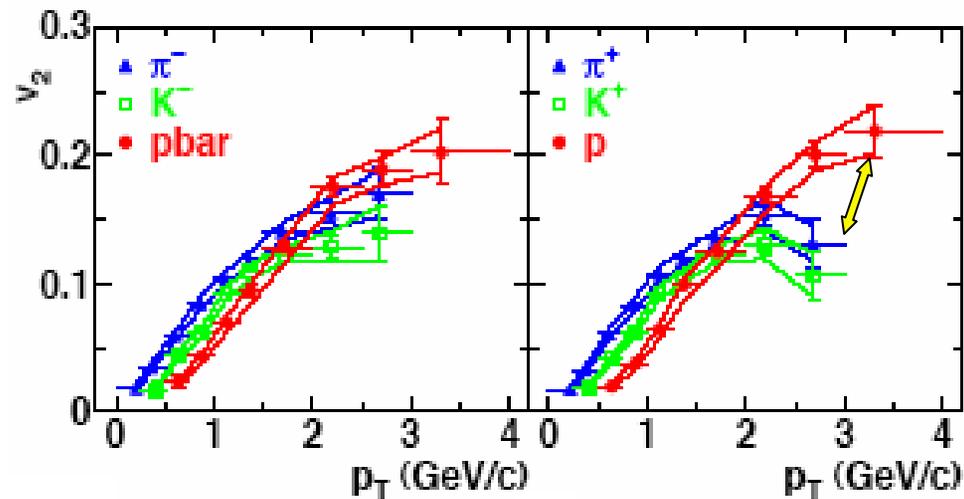
Jet structure for Au+Au collisions is different from that for p+p collisions

Exhibit #4

Azimuthal anisotropy



v_2 : coeff. of 2nd harmonic in ϕ distribution



$$v_2(p) > v_2(\pi) \quad \text{at } p_T > 2.5 \text{ GeV}/c$$

Exhibit #5 Backward-forward asymmetry at intermed. p_T in d+Au collisions

backward
forward

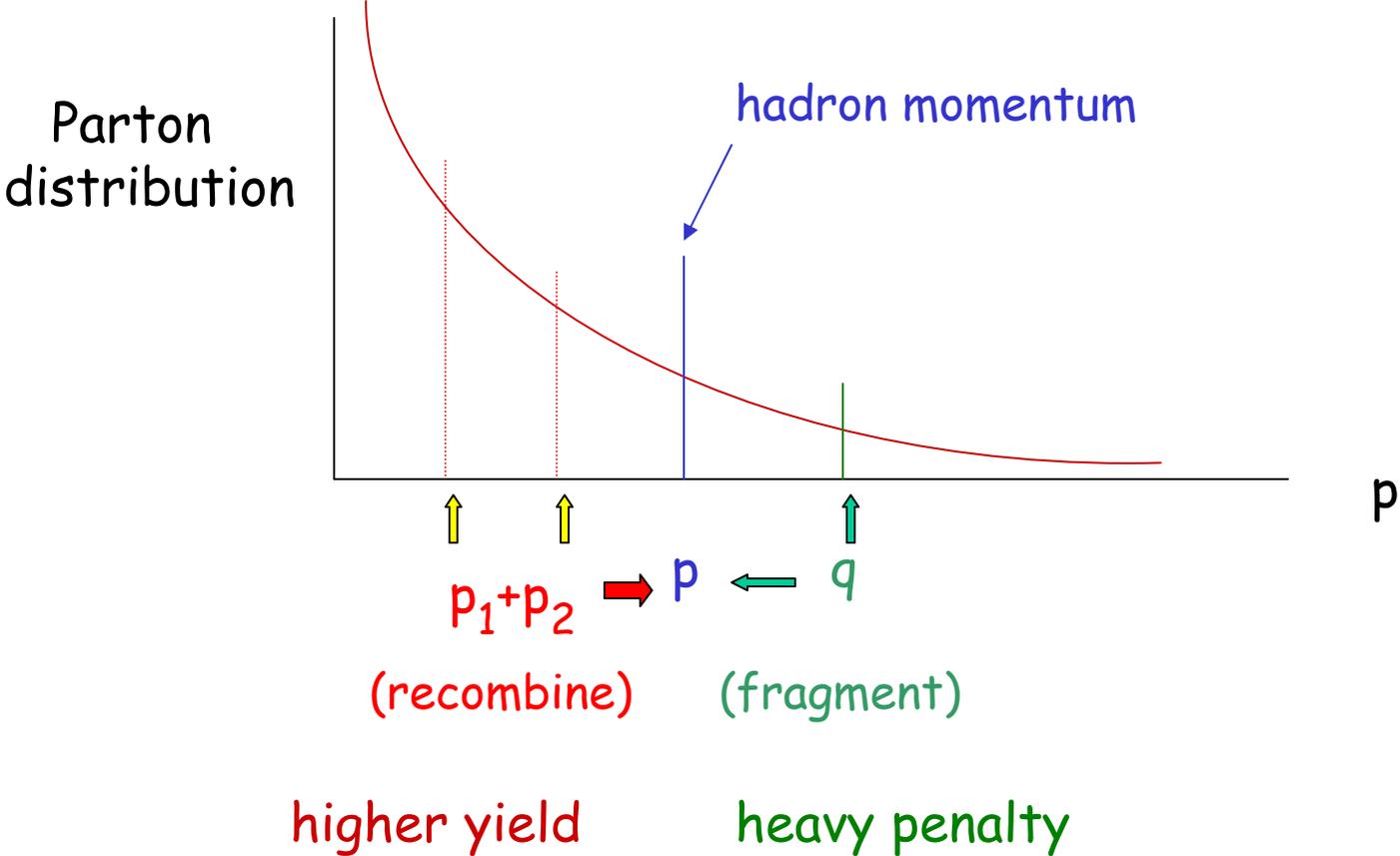
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

The problem:

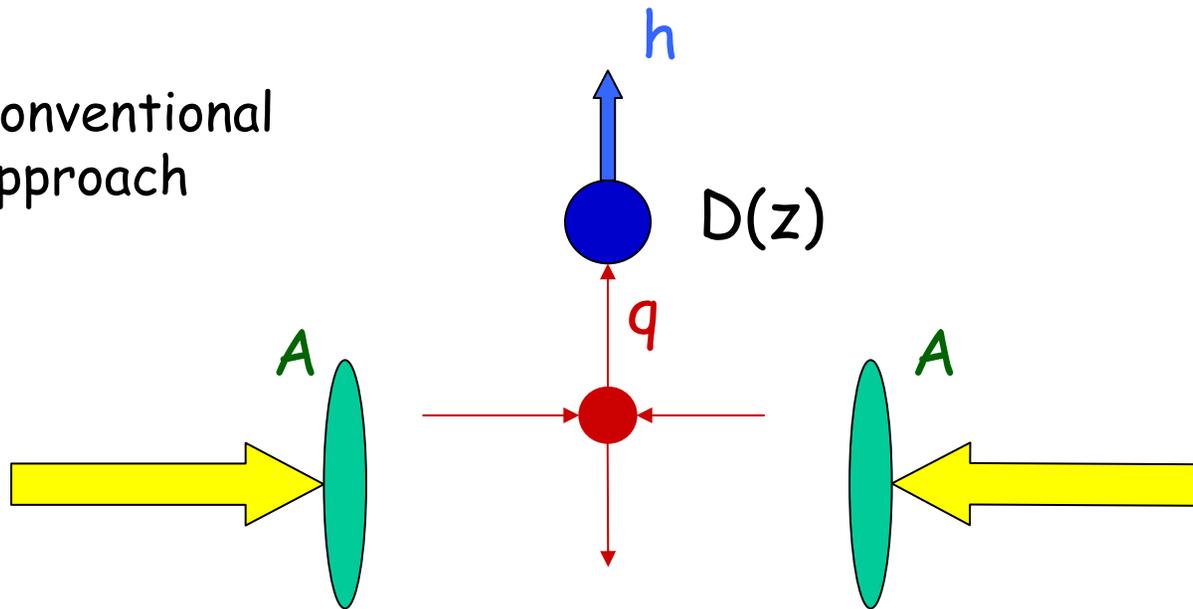
Fragmentation of hard partons
that works so well in leptonic and
hadronic processes

is only one of many components
when the hard parton is in the
environment of many soft partons,
as in heavy-ion collisions.

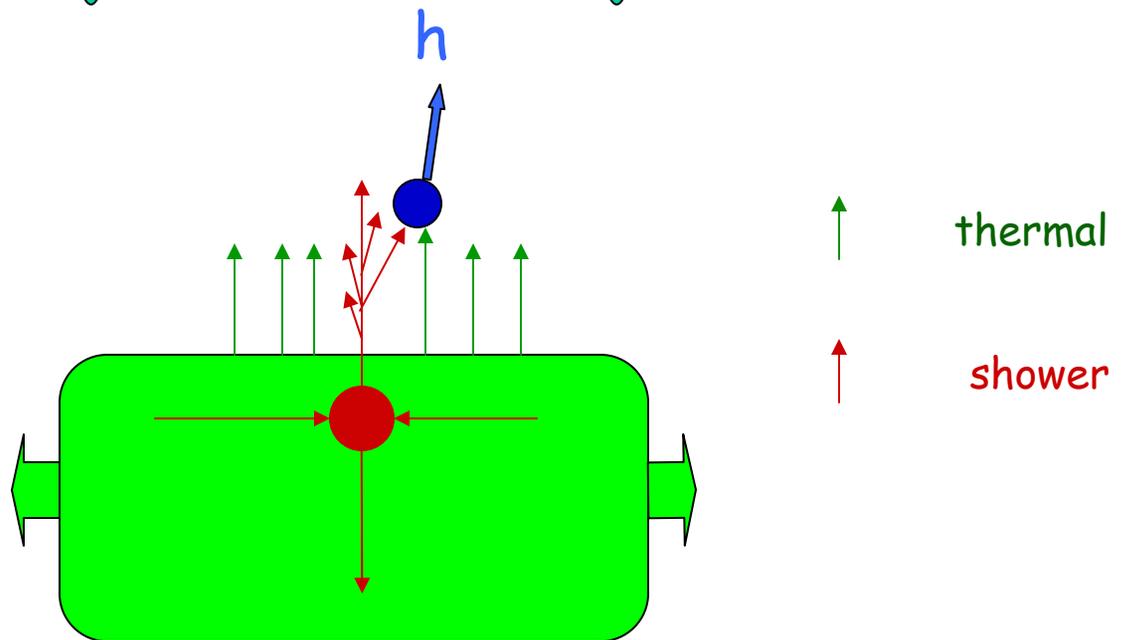
How can recombination solve the puzzles?



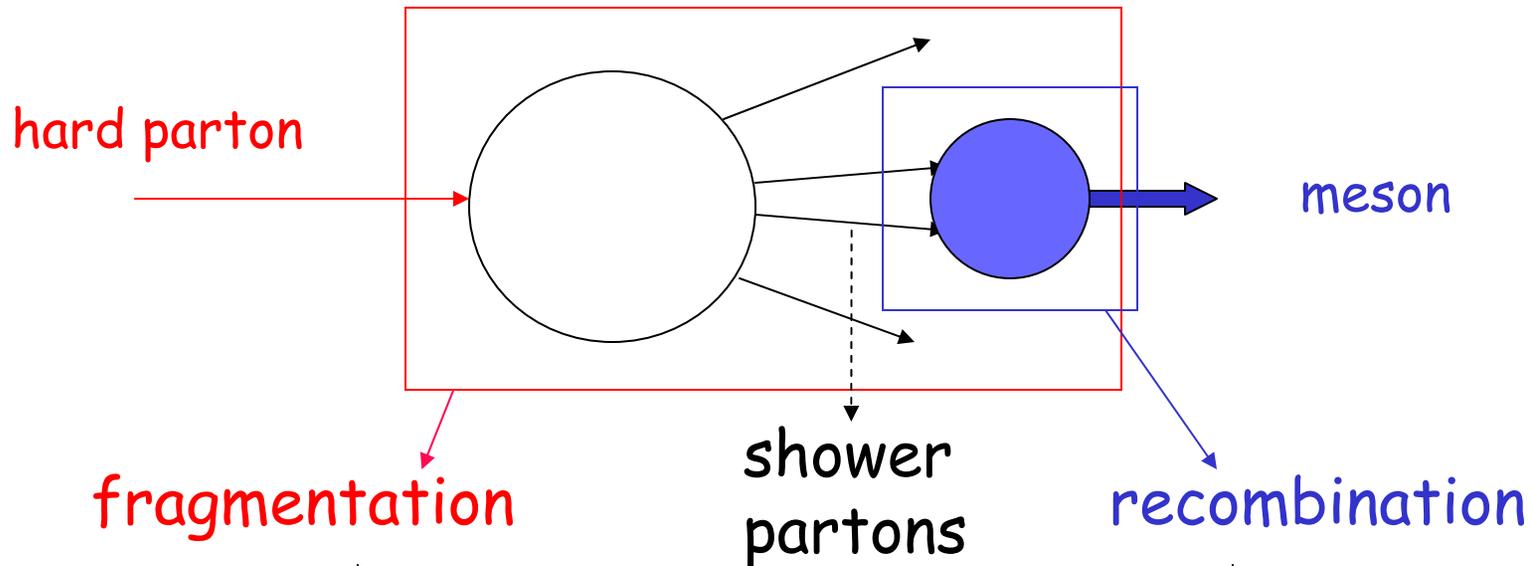
Conventional approach



Thermal-shower recombination



Shower partons in fragmentation process



$$xD_i^h(x) = \int \frac{dx_1}{x_1} \frac{dx_2}{x_2} S_i^j(x_1) S_i^{j'}(x_2) R_{jj'}^h(x_1, x_2, x)$$

known from data (e^+e^- , μp , ...)

can be determined

known from recombination model

Shower parton distributions

$$F_{q\bar{q}'}^{(i)}(x_1, x_2) = S_i^q(x_1) S_i^{\bar{q}'} \left(\frac{x_2}{1-x_1} \right)$$

$$S_i^j = \begin{matrix} & \begin{matrix} u & d & s \end{matrix} \\ \begin{pmatrix} K & L & L_s \\ L & K & L_s \\ L & L & K_s \\ G & G & G_s \end{pmatrix} & \begin{matrix} u \\ d \\ s \\ g \end{matrix} \end{matrix}$$

$$K = \overset{\text{valence}}{\downarrow} K_{NS} + L \longleftarrow \text{sea}$$

$$S_u^{d, \bar{d}, \bar{u}, u(\text{sea})} = L$$

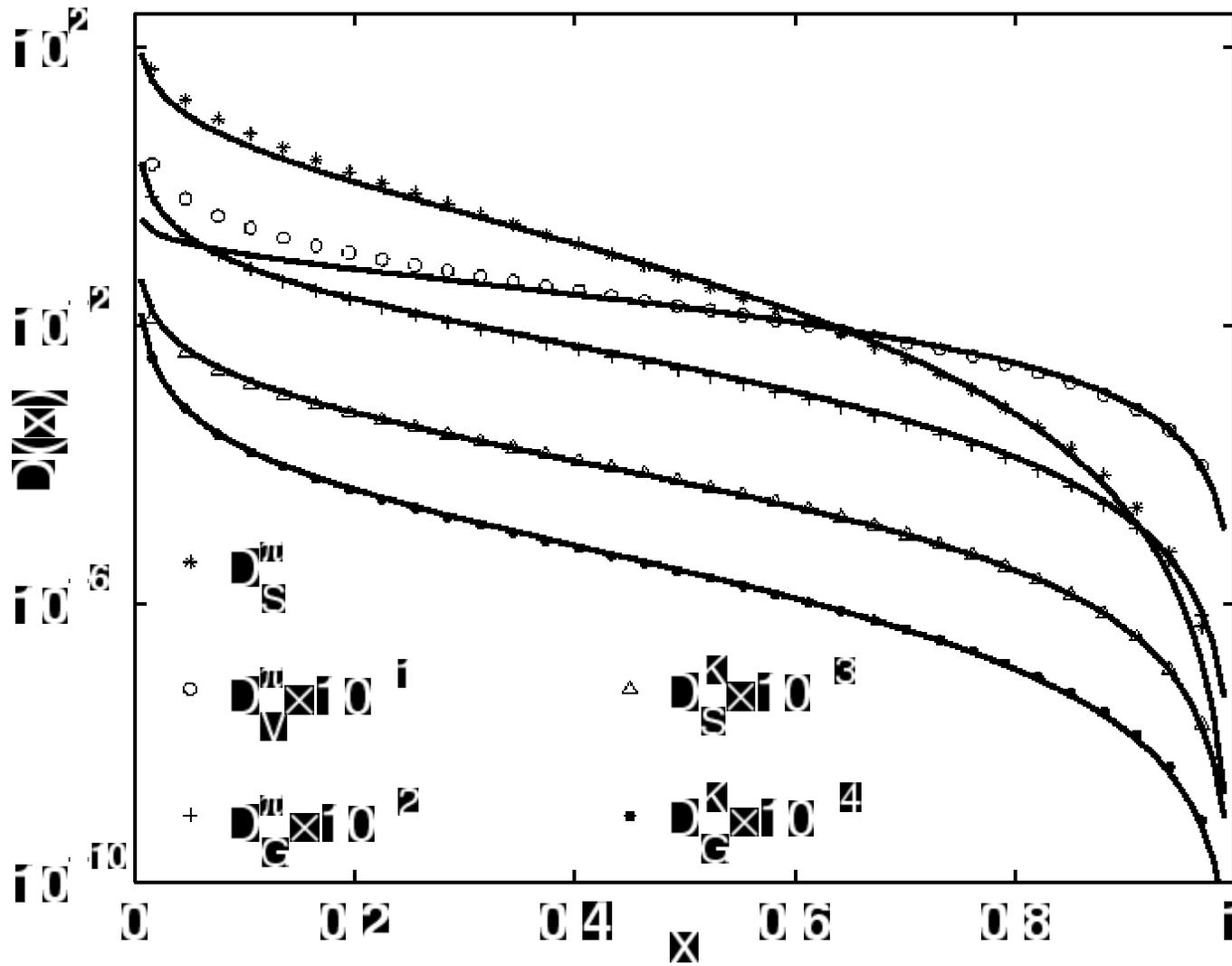
$$K_s = K_{NS} + L_s$$

5 SPDs are determined from 5 FFs.

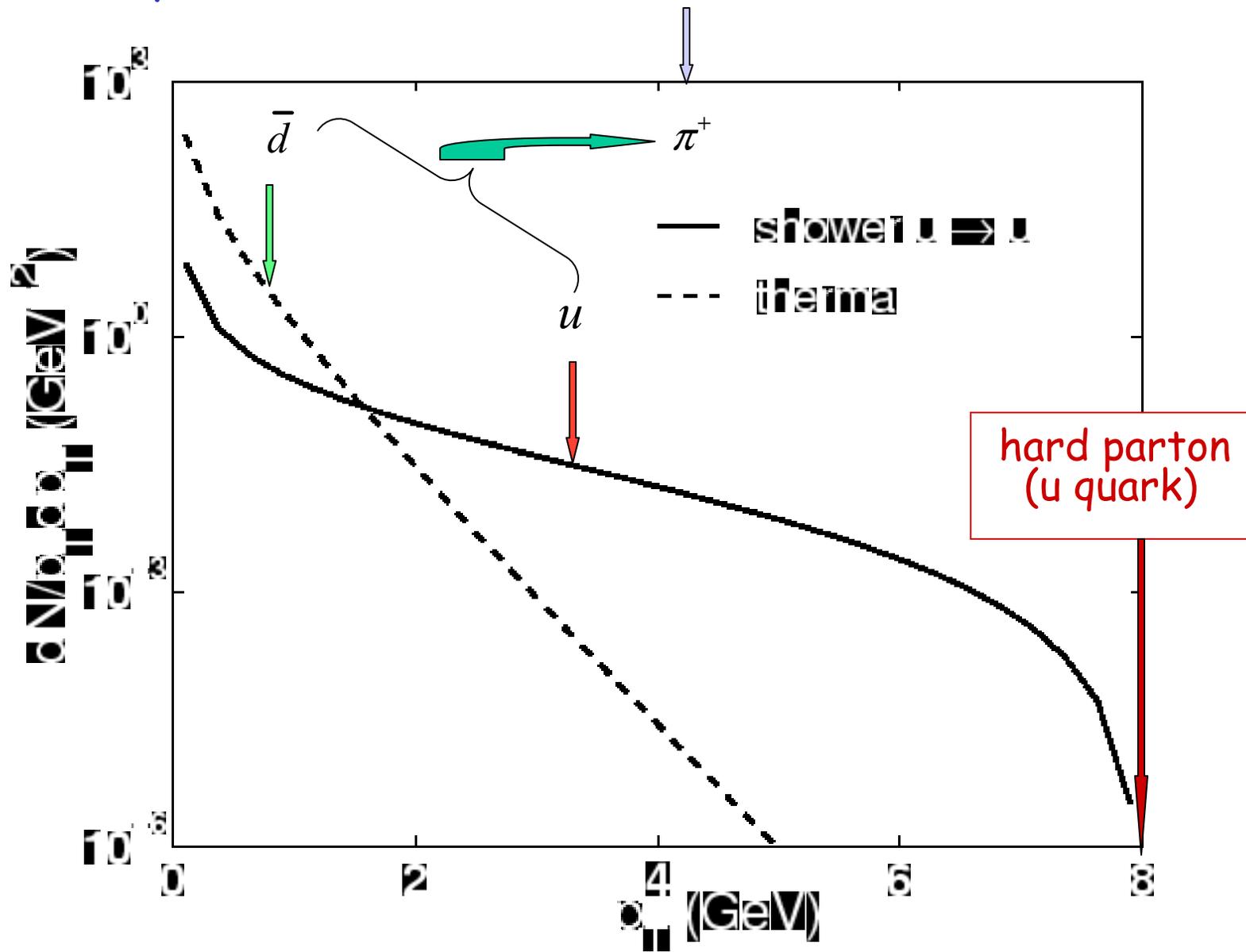
$$\left. \begin{matrix} LL \\ K_{NS} L \\ GG \end{matrix} \right\} R^\pi \begin{matrix} \rightarrow D_{\text{Sea}}^\pi \\ \rightarrow D_V^\pi \\ \rightarrow D_G^\pi \end{matrix}$$

$$\left. \begin{matrix} LL_s \\ GG_s \end{matrix} \right\} R^K \begin{matrix} \rightarrow D_{\text{Sea}}^K \\ \rightarrow D_G^K \end{matrix}$$

BKK fragmentation functions



An example



Inclusive distribution of pions in any direction

$$p \frac{dN_\pi}{dp} = \int \frac{dp_1}{p_1} \frac{dp_2}{p_2} F_{q\bar{q}}(p_1, p_2) R_\pi(p_1, p_2, p)$$

Determine T by fitting dN_π/dp_T
at low p_T ($< 2\text{GeV}/c$)

$$\frac{p_1 p_2}{p} \delta(p_1 + p_2 - p)$$

Pion formation: $q\bar{q}$

$F_{q\bar{q}} = \text{TT} \rightarrow$ soft pions at low p_T

T thermal

S shower

Proton formation: uud distribution

$$F_{uud} = \text{TTT} + \text{TTS} + \text{T(SS)}_1 + \boxed{\text{(SSS)}_1} \leftarrow \text{fragmentation}$$

Thermal partons

$$T(p_1) = p_1 \frac{dN_q^{th}}{dp_1} = Cp_1 \exp(-p_1 / T)$$

Determine C and T by fitting low- p_T data on π production

Shower partons

$$S(p_2) = \xi \sum_i \int dk k f_i(k) S_i^j(p_2 / k)$$

fraction of hard partons that get out of medium to produce shower

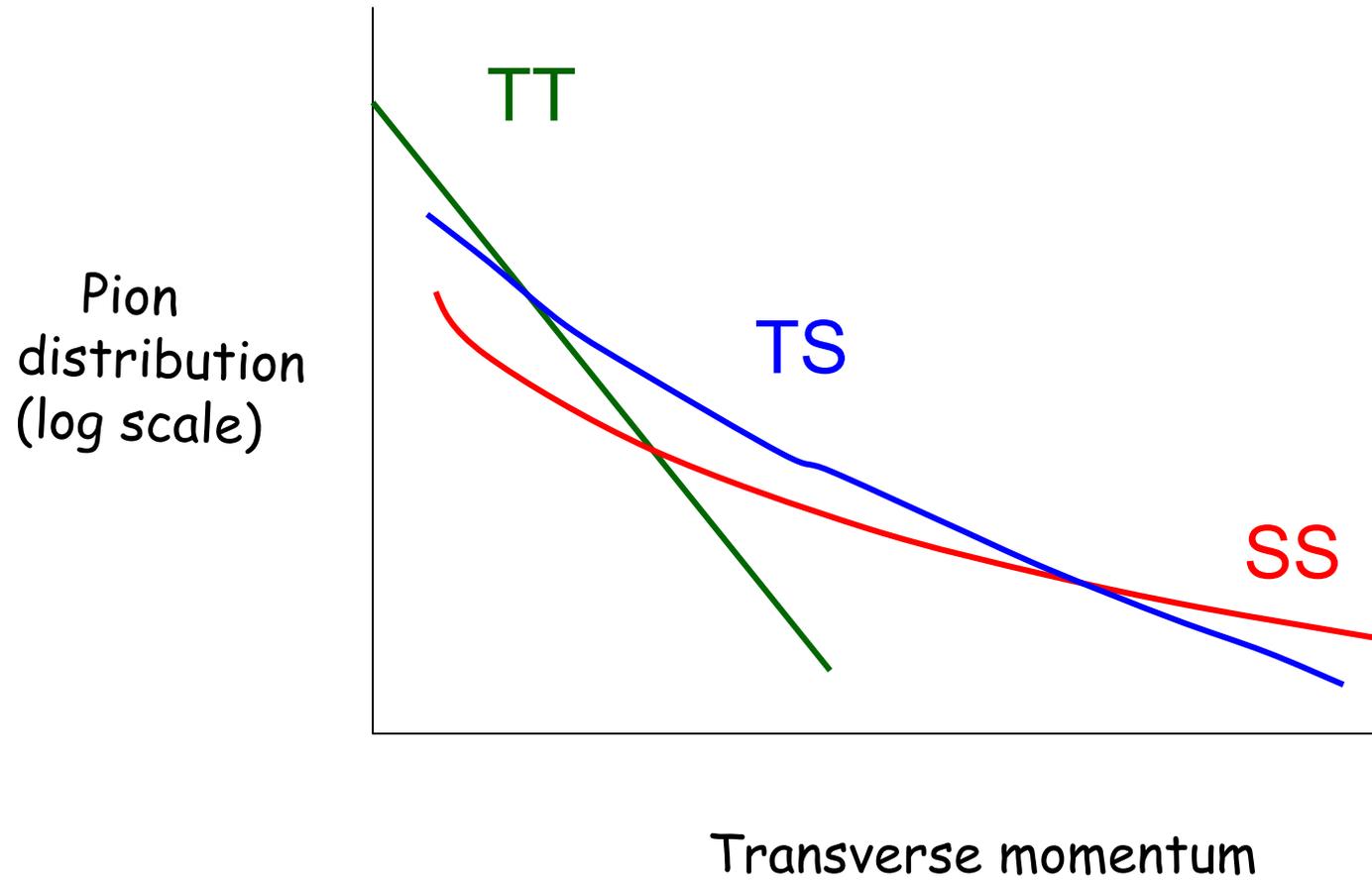
hard parton momentum

distribution of hard parton i in AuAu collisions

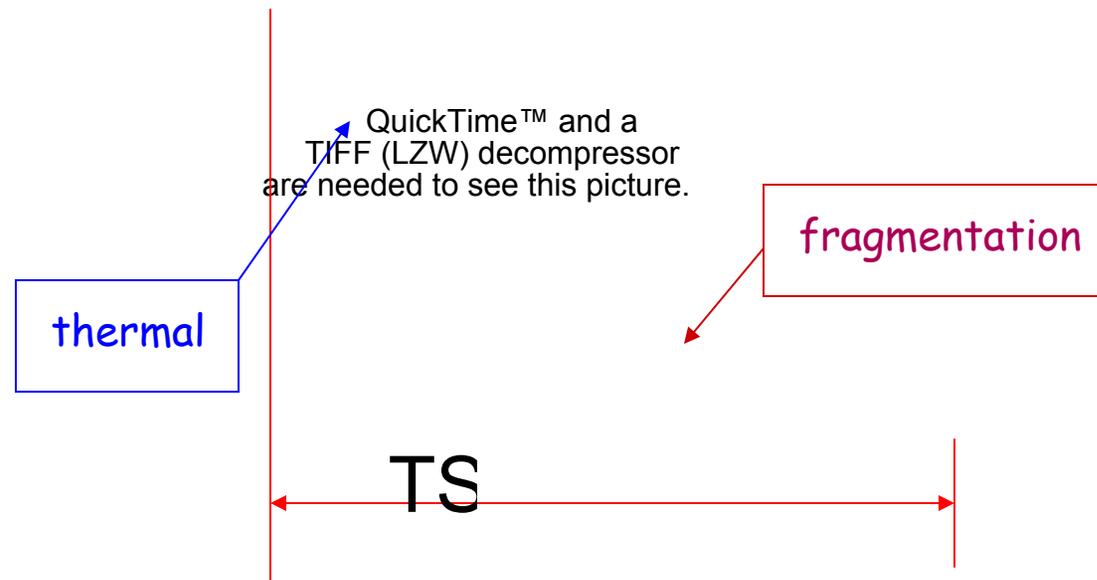
SPD of parton j in shower of hard parton i

$$\xi = 0.07$$

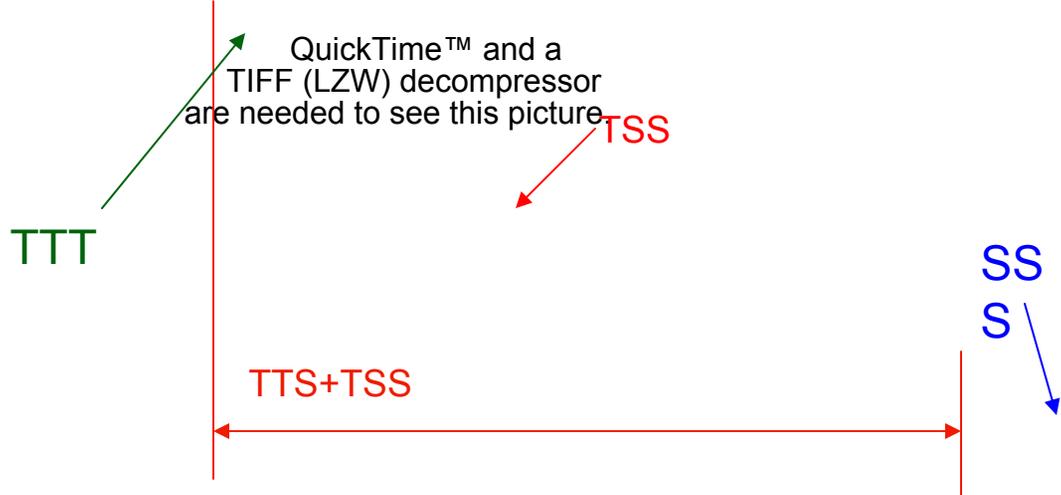
calculable



π production in AuAu central collision at 200 GeV



Proton production in AuAu collisions



Puzzle #1

Proton/pion ratio

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

resolved

Hwa & CB Yang, PRC70, 024905 (2004)

Compilation of $R_{p/\pi}$ by R. Seto (UCR)

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture

All in recombination/
coalescence model

Duke: Fries, Mueller, Nonaka, Bass, PRL90,202303(2003);PRC68,044902(2003).
TAM: Greco, Ko, Levai, PRL,90,202302(2003); PRC68,034904(2003).

d+Au collisions

Pions

$\xi = 1$

soft-soft



QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

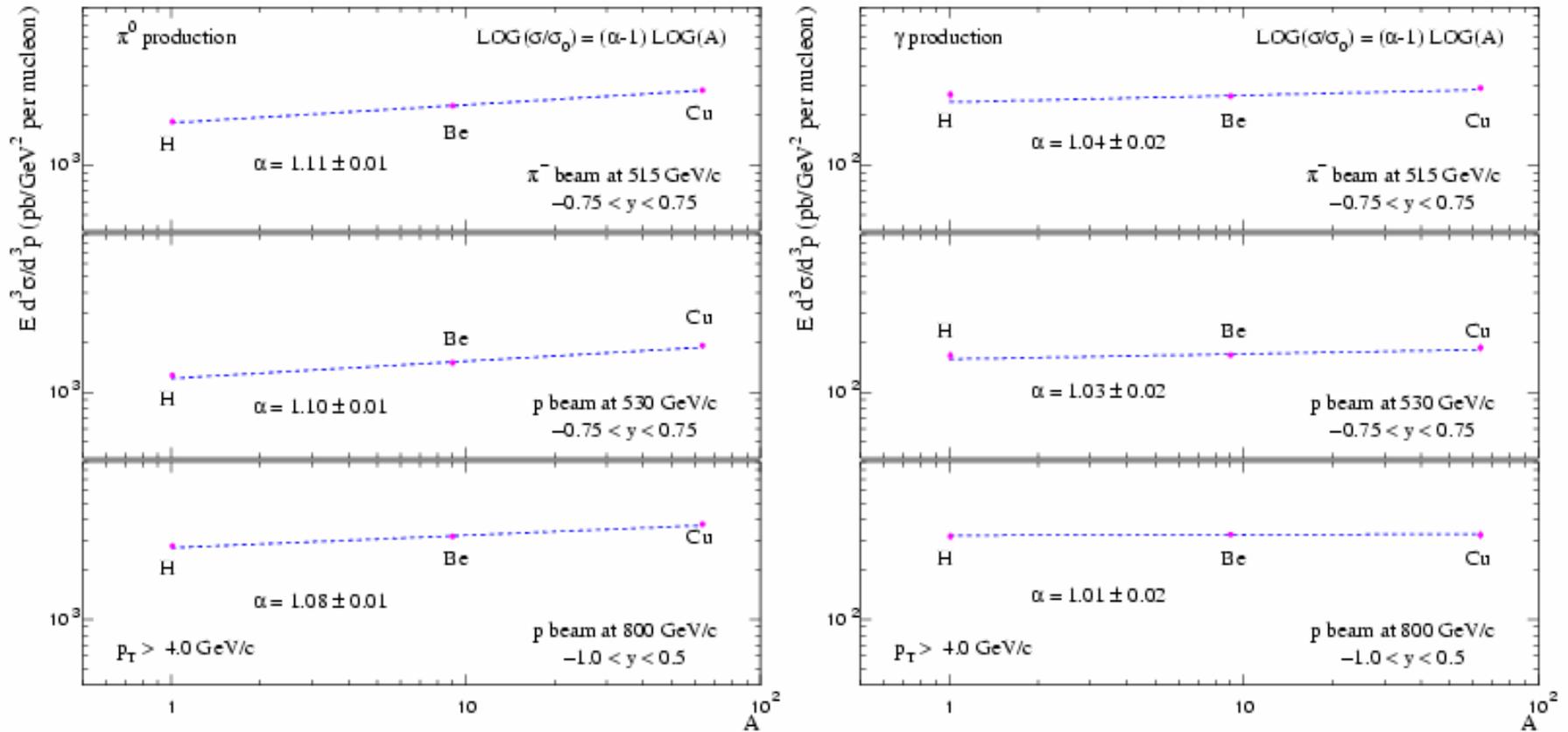
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

- ★ No p_T broadening by multiple scattering in the initial state.

Medium effect is due to thermal (soft)-shower recombination in the final state.

Hwa & CB Yang, PRL 93, 082302 (2004)

π^0 , γ production in p+A collisions



$\pi^0: \alpha = 1.1$

$\gamma: \alpha = 1.0$

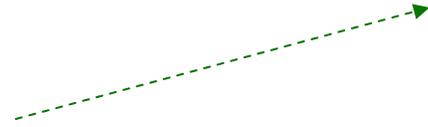
Proton

Thermal-shower
recombination is
NOT negligible.



QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Thermal-shower
recombination is
negligible.



Nuclear Modification Factor

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Puzzle #2
resolved

$$R_{CP}^p > R_{CP}^\pi$$

because $3q \rightarrow p, 2q \rightarrow \pi$

Jet Structure

Exhibit #3 Jet structure in Au+Au
different from that in p+p collisions

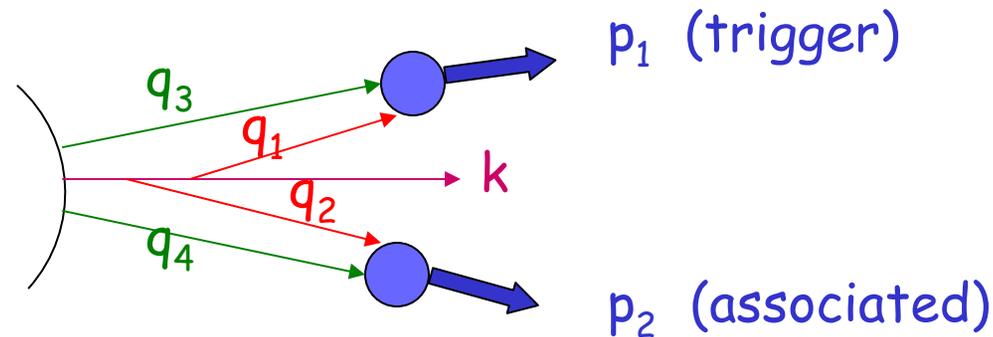
A simple consequence of TS recombination
being more important in Au+Au collisions
than in p+p collisions.

Consider dihadron correlation in the same jet
on the near side.

Trigger at $4 < p_T < 6 \text{ GeV}/c$

p+p: mainly SS \longleftrightarrow fragmentation
 Au+Au: mainly TS

Associated particle



$$\xi \sum_i \int dk k f_i(k) \begin{pmatrix} S(q_1) T(q_3) R(q_1, q_3, p_1) \\ S(q_2) T(q_4) R(q_2, q_4, p_2) \end{pmatrix} \begin{matrix} \longrightarrow & \text{trigger} \\ \longrightarrow & \text{associated} \end{matrix}$$

There are other contributions as well.

trigger

associated

(TS) + [TS]

(TS) + [SS]

(SS) + [TS]

(SS) + [SS]



too small for Au+Au

but the only term for p+p

Associated particle distribution

(TS)[TS] →

} π^+

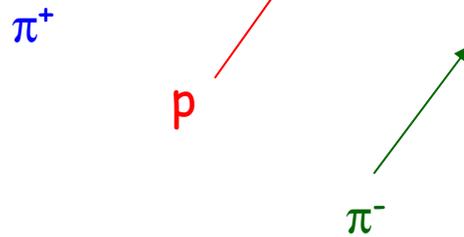
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

(TS)[SS]+(SS)[TS]

(SS)[SS] for p+p collisions much lower

Associated particle distribution in central Au+Au collisions

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.



Centrality dependence of associated-particle dist.

Distributions

Ratios

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

p+p

Au+Au central

d+Au peripheral

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

d+Au central

d+Au peripheral

Exhibit #3 Jets in Au+Au and in p+p are very different.

Dihadron correlation with soft parton correlation

Fries, Bass, Mueller, nucl-th/0407102

- (1) soft-soft recombination X soft-soft recom
- (2) fragmentation X fragmentation
- (3) soft-hard recom X fragmentation

No shower partons

(1) + (2) with 8% scaled soft parton correlation

Comparable associated particle yields for meson and baryon triggers

Meson trigger

Baryon trigger

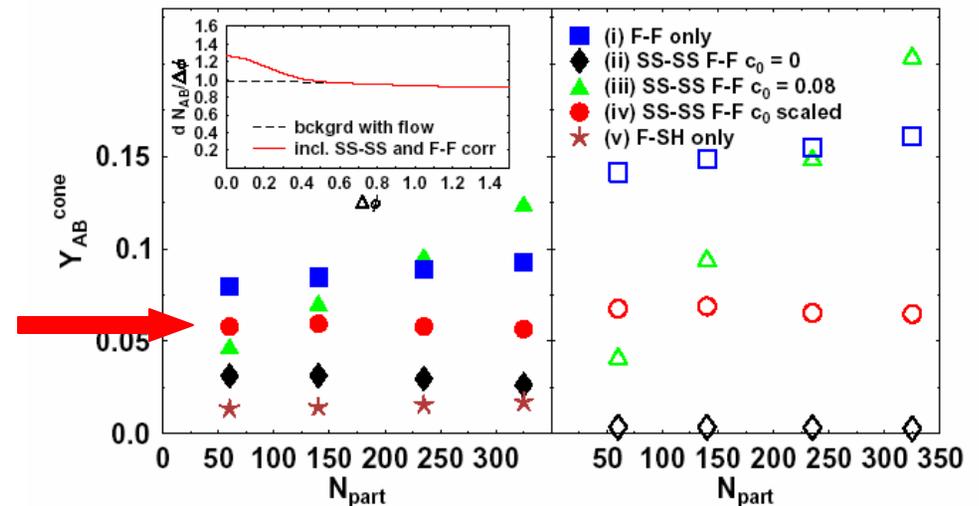
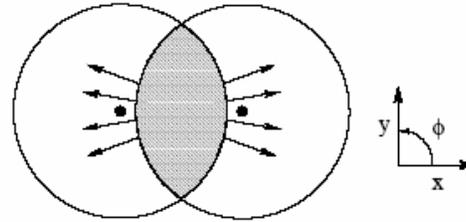
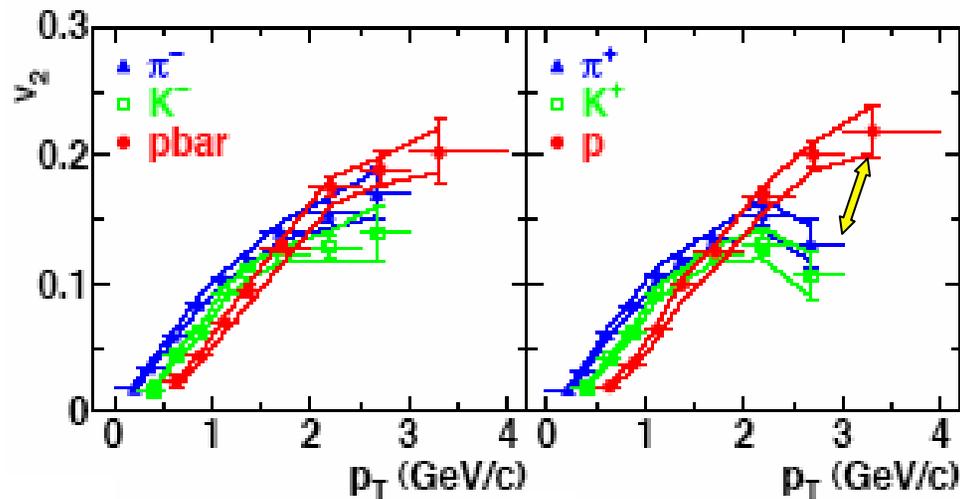


Exhibit #4

Azimuthal anisotropy



v_2 : coeff. of 2nd harmonic of ϕ distribution



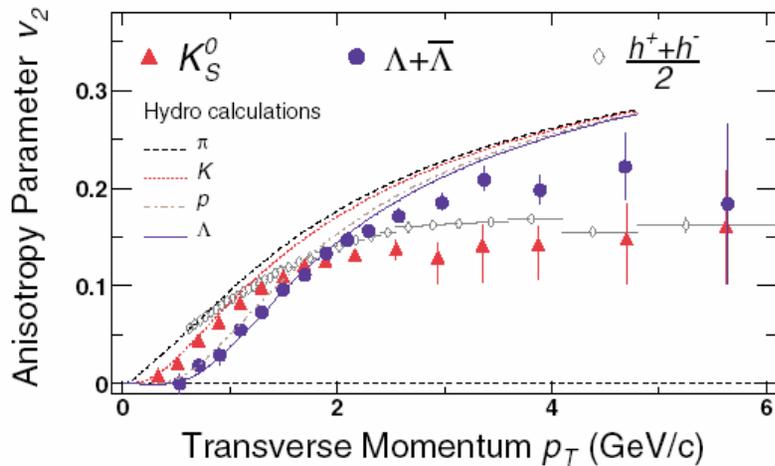
PHENIX,
PRL 91 (2003)

$$v_2(p) > v_2(\pi) \quad \text{at } p_T > 2.5 \text{ GeV}/c$$

Molnar and Voloshin, PRL 91, 092301 (2003).

Parton coalescence implies that $v_2(p_T)$ scales with the number of constituents

STAR data



VIEW LETTERS

week ending
6 FEBRUARY 2004

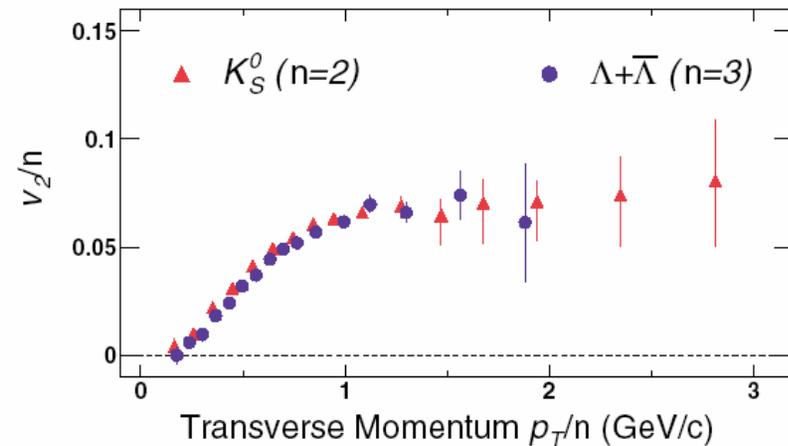


Exhibit #5 Forward-backward asymmetry at intermed. p_T in d+Au collisions

backward
forward

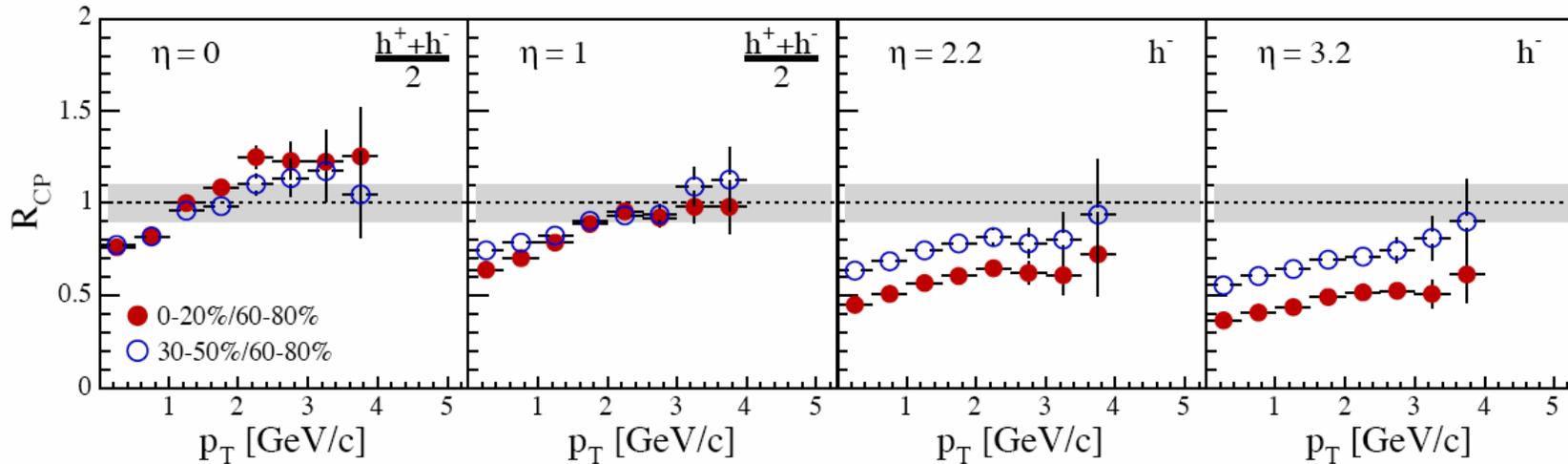
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

STAR preliminary data

Rapidity dependence of R_{CP} in d+Au collisions

BRAHMS

nucl-ex/0403005



$R_{CP} < 1$ at $\eta = 3.2$ \longrightarrow

Central more suppressed
than peripheral collisions

This can be understood in terms of parton recombination.

Forward production in d+Au collisions

Same hadronization process as at $\eta=0$.

Soft component: $T(p_1) = p_1 \frac{dN_q^{soft}}{dp_1} = Cp_1 \exp(-p_1/T)$

Pion distribution $\frac{dN_\pi^{soft}}{p_T dp_T d\eta} \Big|_{\eta=0} = \frac{C^2}{6} \exp(-p_T/T)$

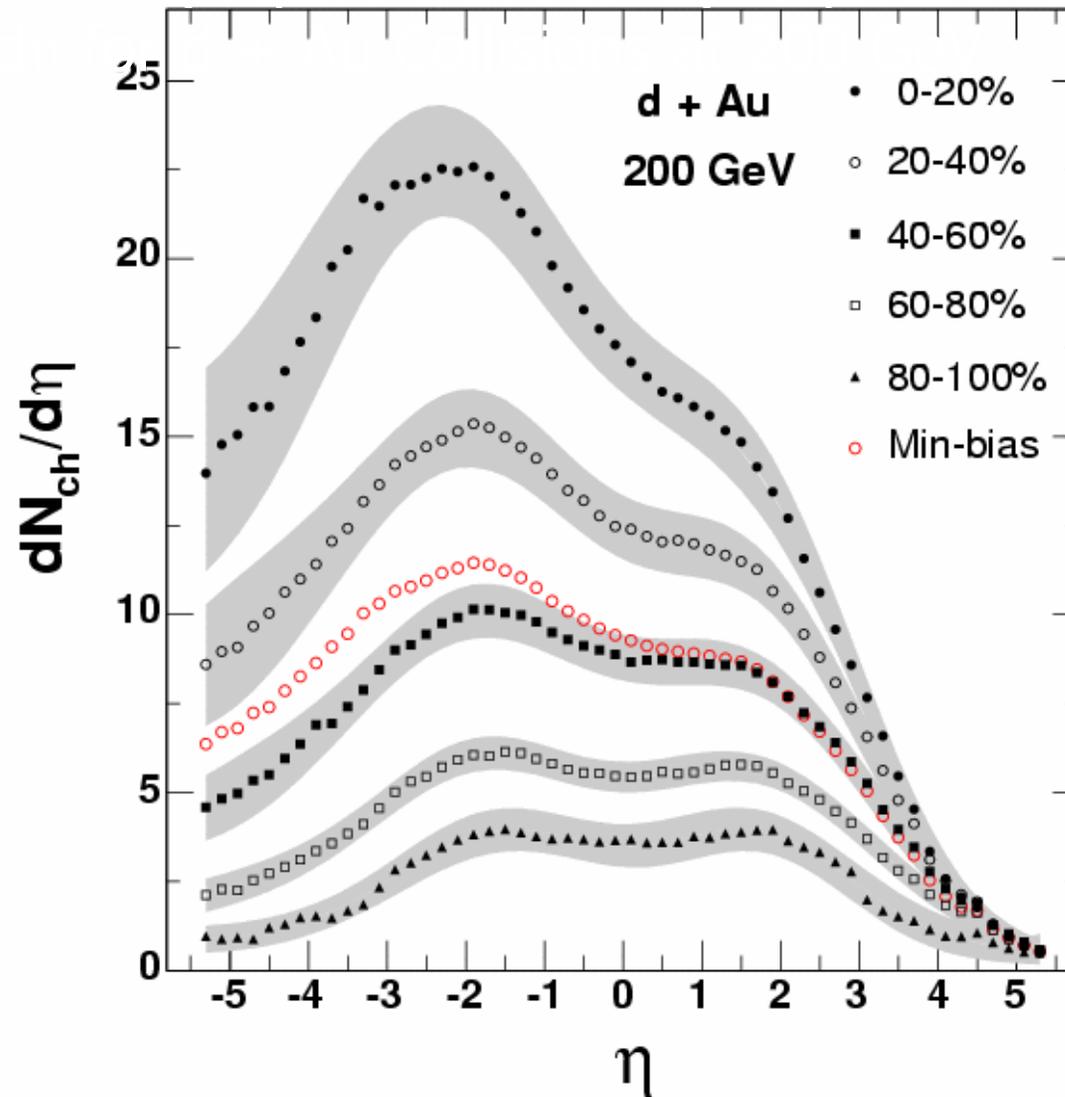
$C(\beta, \eta)$

Notation:

Centrality 0-20% \longrightarrow $\beta = 0.1$
 60-80% \longrightarrow 0.7

$C(\beta, \eta)$
 for all β and η

$$C(\beta, \eta) = C(\beta, 0) \left[\frac{dN_{ch} / d\eta(\beta)}{dN_{ch} / d\eta|_{\eta=0}(\beta)} \right]^{1/2}$$



$\therefore C(\beta, \eta)$ can be calculated.

Our prediction on the π^+ spectra for $0 \leq \eta \leq 3.2$, $0.1 \leq \beta \leq 0.7$

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Hwa, Yang, and Fries,
nucl-th/0410111

T=constant, no free parameter

No new physics, e.g. gluon saturation, explicitly.

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Include β and η dependences of T (one free parameter) and momentum degradation



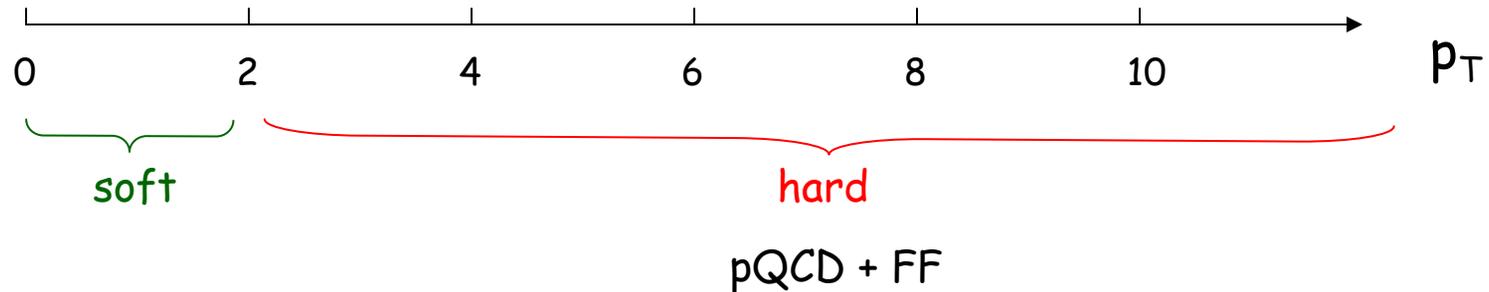
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Without any more free parameters

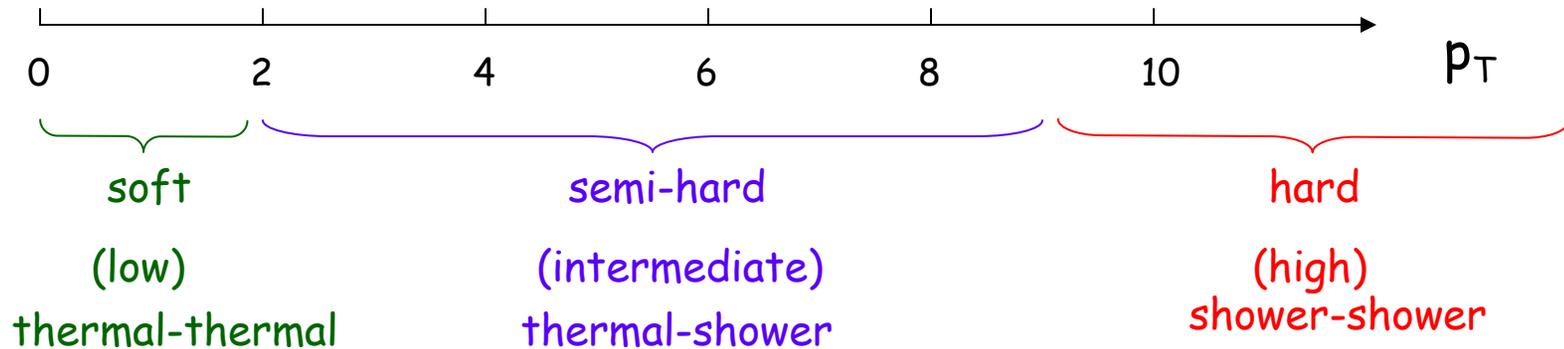
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Summary

Traditional picture



More realistic picture



All anomalies at intermediate p_T can be understood in terms of recombination of

thermal and shower partons

Recombination is the hadronization process.

At $p_T > 9 \text{ GeV}/c$ fragmentation dominates, but it can still be expressed as shower-shower recombination.

Hence, parton recombination at all p_T .

Conclusion in one sentence

Due to the soft parton environment
in heavy-ion collisions at RHIC,

hadronization by thermal-shower
recombination

dominates all phenomena on particle
production at intermediate p_T

for all centralities and all rapidities.

All of the work presented
was done in collaboration with

Chunbin Yang

Hua-zhong Normal University

Wuhan, China