

QUARKONIUM FORMATION IN STATISTICAL AND KINETIC MODELS

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WHAT DO HEAVY FLAVORS PROVIDE?

MASS SCALES: $M_c = 1.5 \text{ GeV}$, $M_b = 5.0 \text{ GeV}$

DISTANCE SCALES: 0.07 fm, 0.02 fm

BOUND STATE SCALES: 0.30 fm, 0.13 fm

HARD PROBES: $M_Q \gg \Lambda_{\text{QCD}}$

CALCULATIONS: LO + NLO pQCD

NUCLEAR COLLISIONS:

SHADOWING, SATURATION, k_t
broadening

POINTLIKE PROCESS: SCALES
WITH BINARY COLLISIONS:

$$N_{coll} \propto N_p^{\frac{4}{3}}$$

QUARKONIUM FORMATION

MATSUI-SATZ: $R_{\text{plasma screening}} < R_{\text{quarkonium}}$: SUPPRESSION

KHARZEEV-SATZ: Ionization with deconfined gluons

NA50: Anomalous Suppression

ALTERNATIVES: Dense hadronic medium, comovers

M. Sitta (NA50), hep-ex/0405056

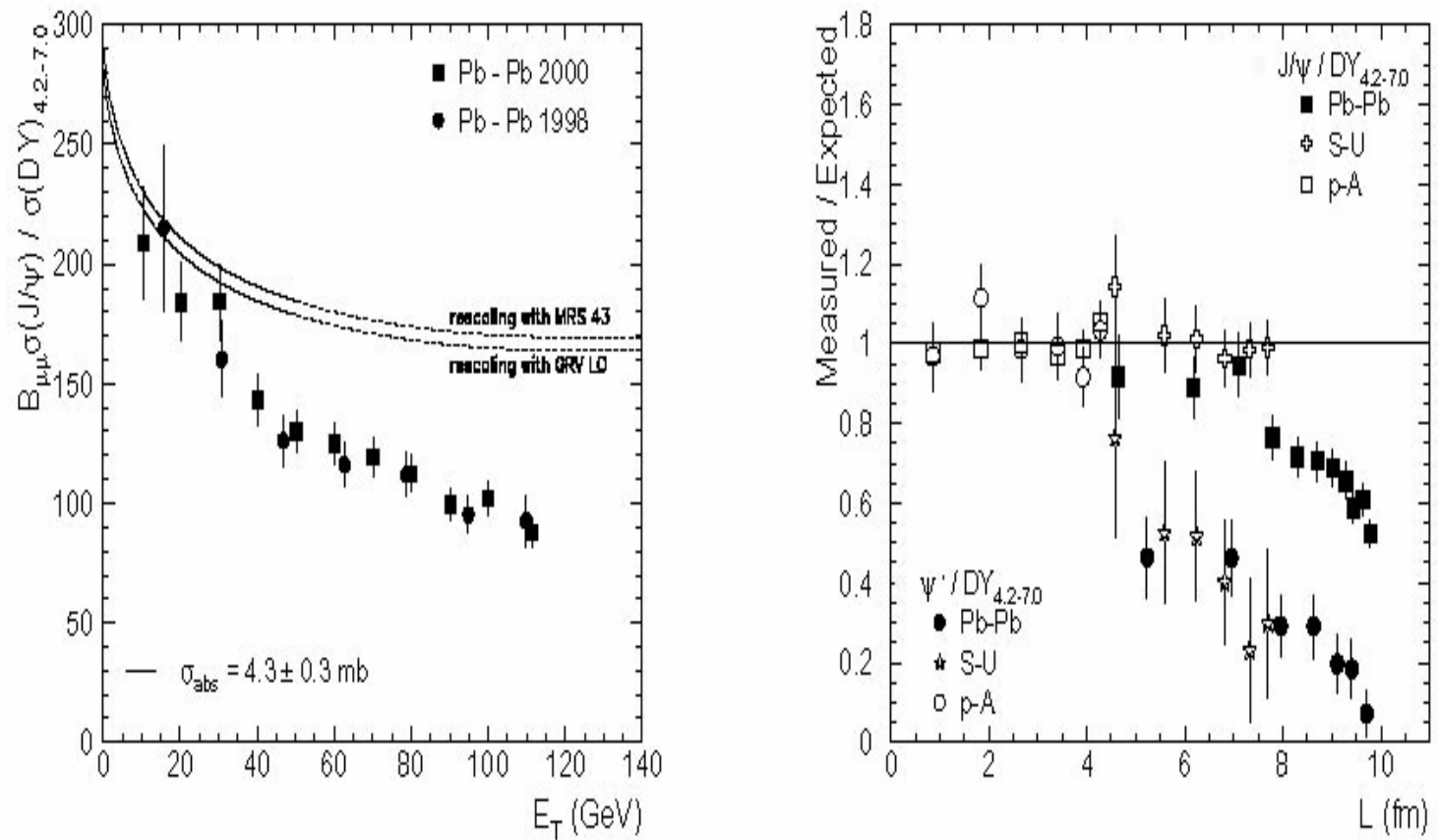


Figure 3: $B_{\mu\mu}\sigma(J/\psi)/\sigma(DY)$ as a function of E_T (left) and the ratio "measured value"/"expected value" for the relative yields $B_{\mu\mu}\sigma(J/\psi)/\sigma(DY)$ and $B_{\mu\mu}\sigma(\psi')/\sigma(DY)$ as a function of L (right).

Multiple $c\bar{c}$ pairs in high energy AA Collisions

$$N_{c\bar{c}}(b=0) \approx 30\sigma_{c\bar{c}}^{pp} (mb)$$

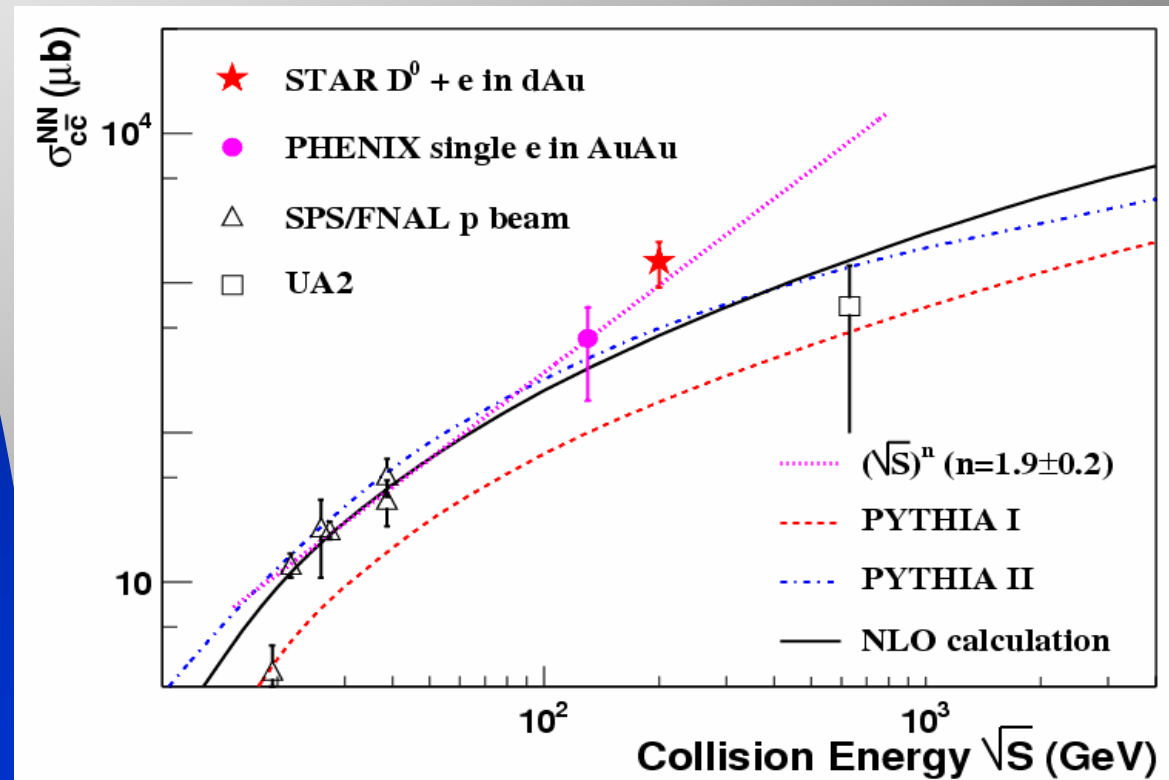
CENTRAL VALUES:

- 10-15 from extrapolation of low energy
- 20 from PHENIX electrons
- 40 from STAR electrons and $K\pi$

$709 \pm 85 \pm \frac{332}{281} \mu b$: PHENIX nucl-ex/0403057 (pp)

$622 \pm 57 \pm 160 \mu b$: PHENIX nucl-ex/0409028 (Au-Au)

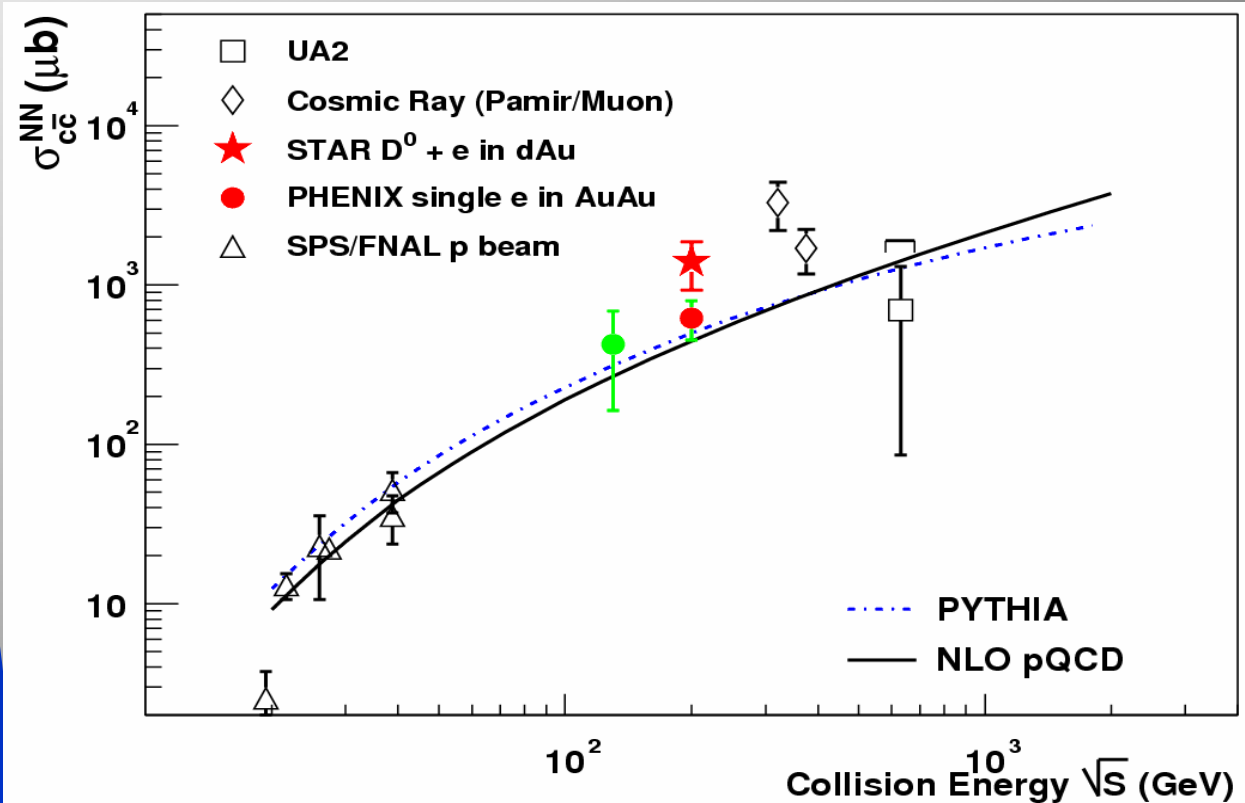
$1.4 \pm 0.2 \pm 0.4 mb$: STAR nucl-ex/0407006 (d-Au)



$709 \pm 85 \pm \frac{332}{281} \mu b$: PHENIX nucl-ex/0403057 (pp)

$622 \pm 57 \pm 160 \mu b$: PHENIX nucl-ex/0409028 (Au-Au)

$1.4 \pm 0.2 \pm 0.4 mb$: STAR nucl-ex/0407006 (d-Au)



PROBE REGION OF COLOR DECONFINEMENT WITH MULTIPLE PAIRS OF HEAVY QUARKS

Avoids Matsui-Satz Condition

Distribute Heavy Flavor at hadronization (statistical)

Form Quarkonium in the medium: inverse suppression

IF THE INCOHERENT RECOMBINATION OF HEAVY QUARKS DETERMINES FINAL HADRONIC ABUNDANCES:

Probability for charm quark to combine with anticharm:

$$\varepsilon = N_c / N_{u,d} \propto N_{c\bar{c}} / N_{ch}$$

Since $\varepsilon \ll 1$, sum for each \bar{c} :

$$N_{quarkonium} \propto N_{c\bar{c}}^2 / N_{ch}$$

Average over fluctuations:

$$\langle J / \psi \rangle = \lambda \langle N_{c\bar{c}} \rangle (\langle N_{c\bar{c}} \rangle + 1) / N_{ch}$$

Centrality dependence in terms of participants N_p :

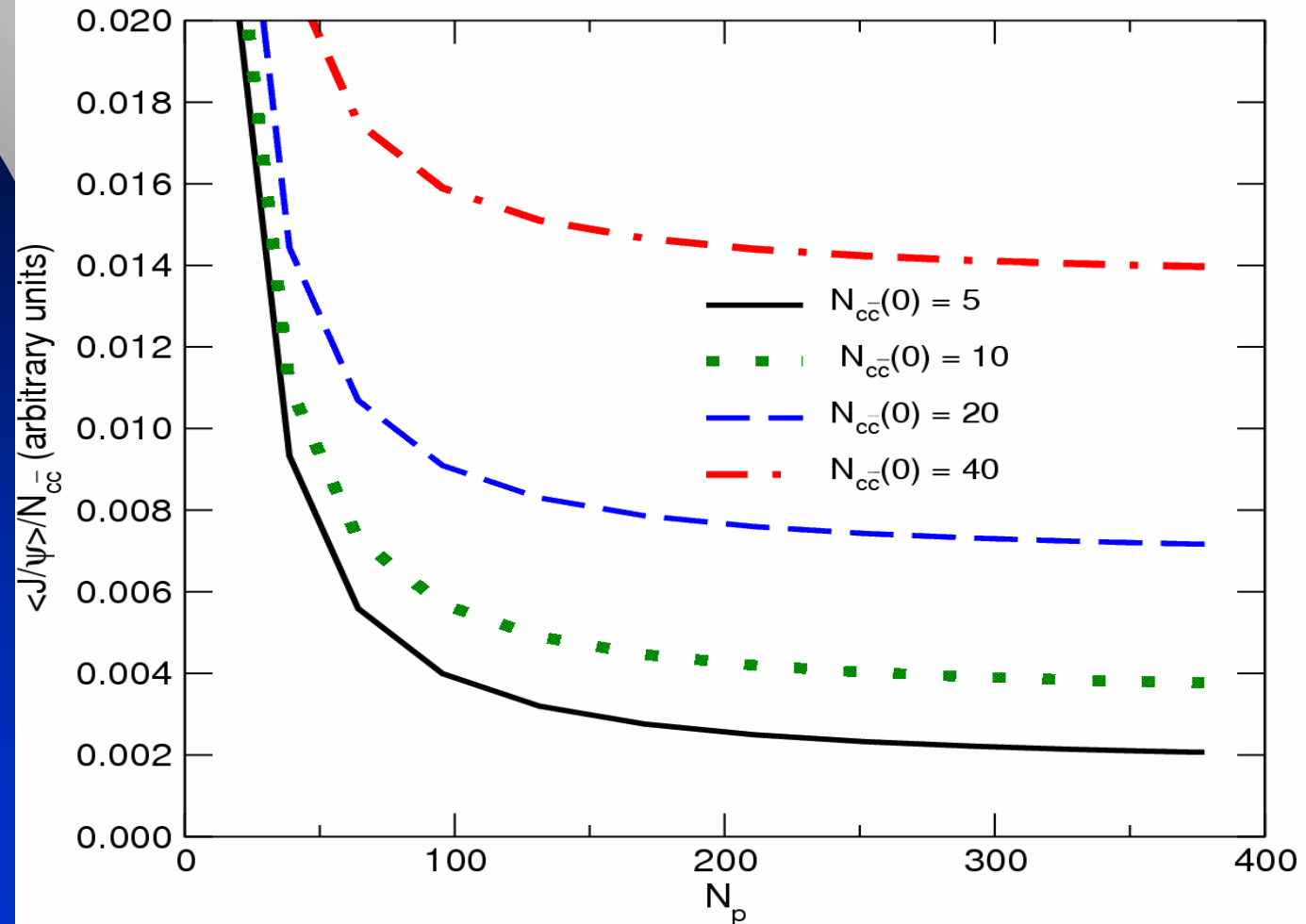
$$\text{Parameterize } N_{ch} \propto N_p^{1+\Delta},$$

$$\langle J / \psi \rangle / \langle N_{c\bar{c}}(binary) \rangle = aN_p^{\frac{1}{3}-\Delta} + bN_p^{-1-\Delta}$$

FORMATION CENTRALITY SIGNATURES

Centrality Dependence of Quarkonium Formation

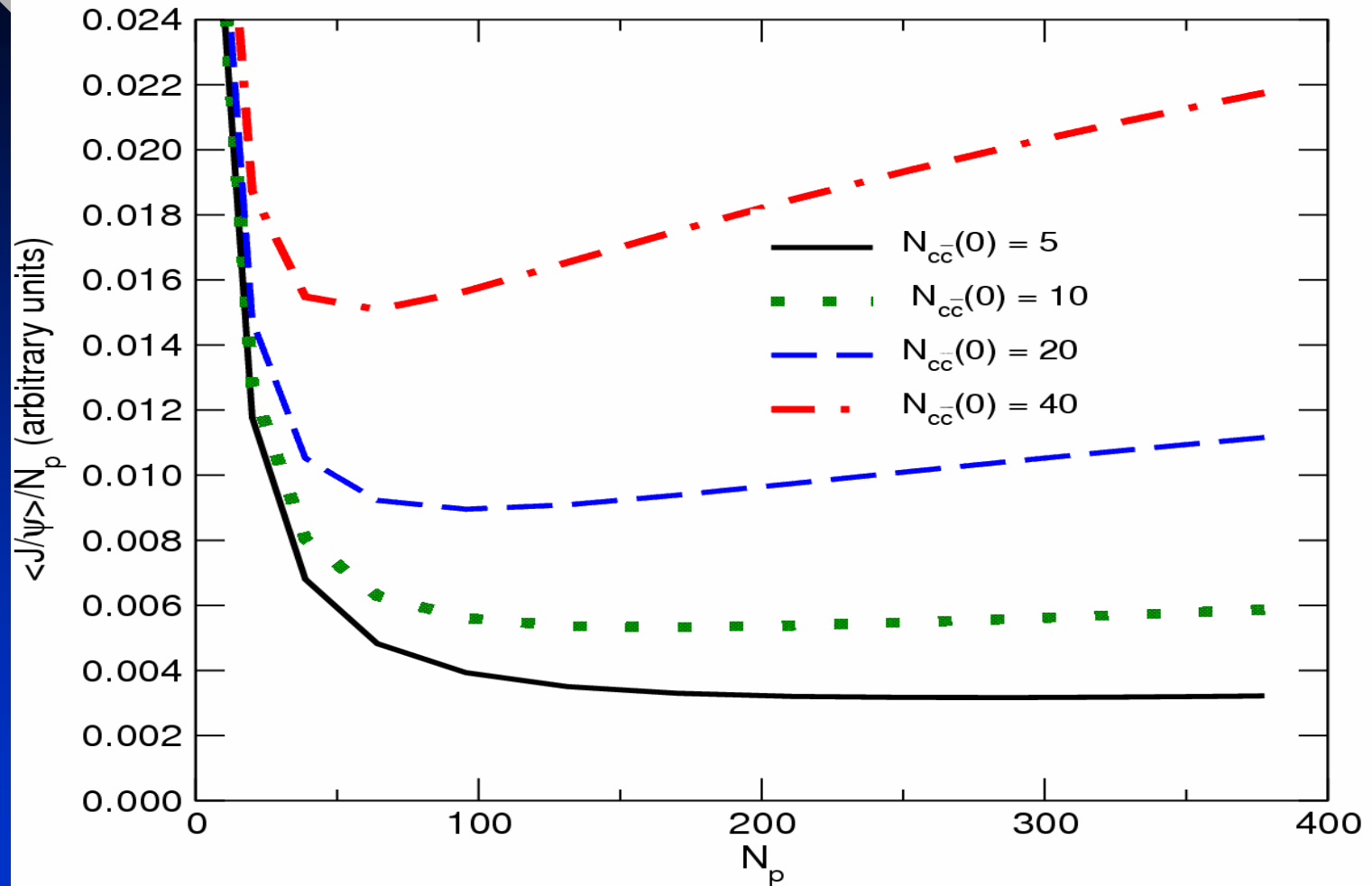
$\Delta = 0.33$



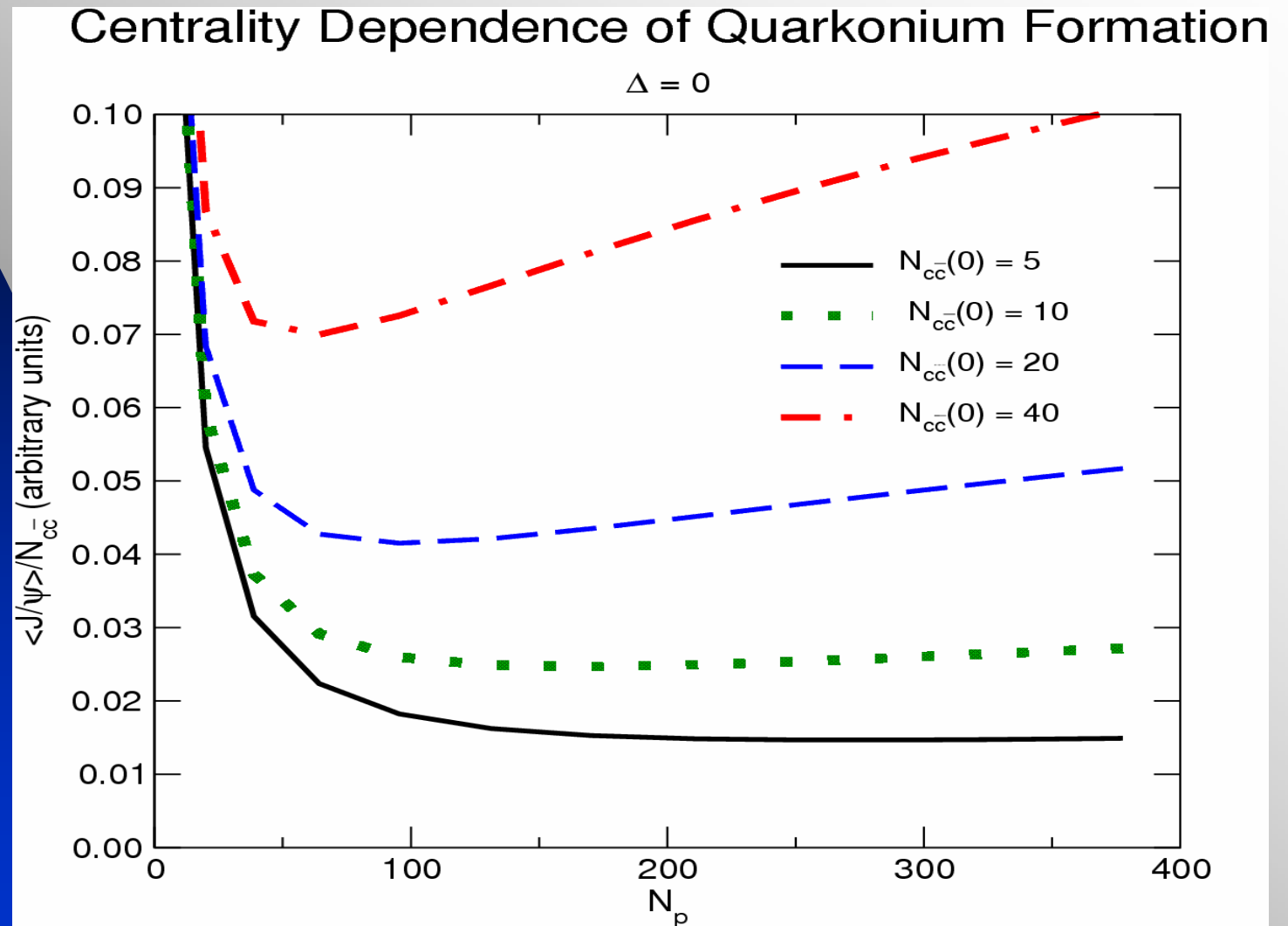
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Centrality Dependence of Quarkonium Formation

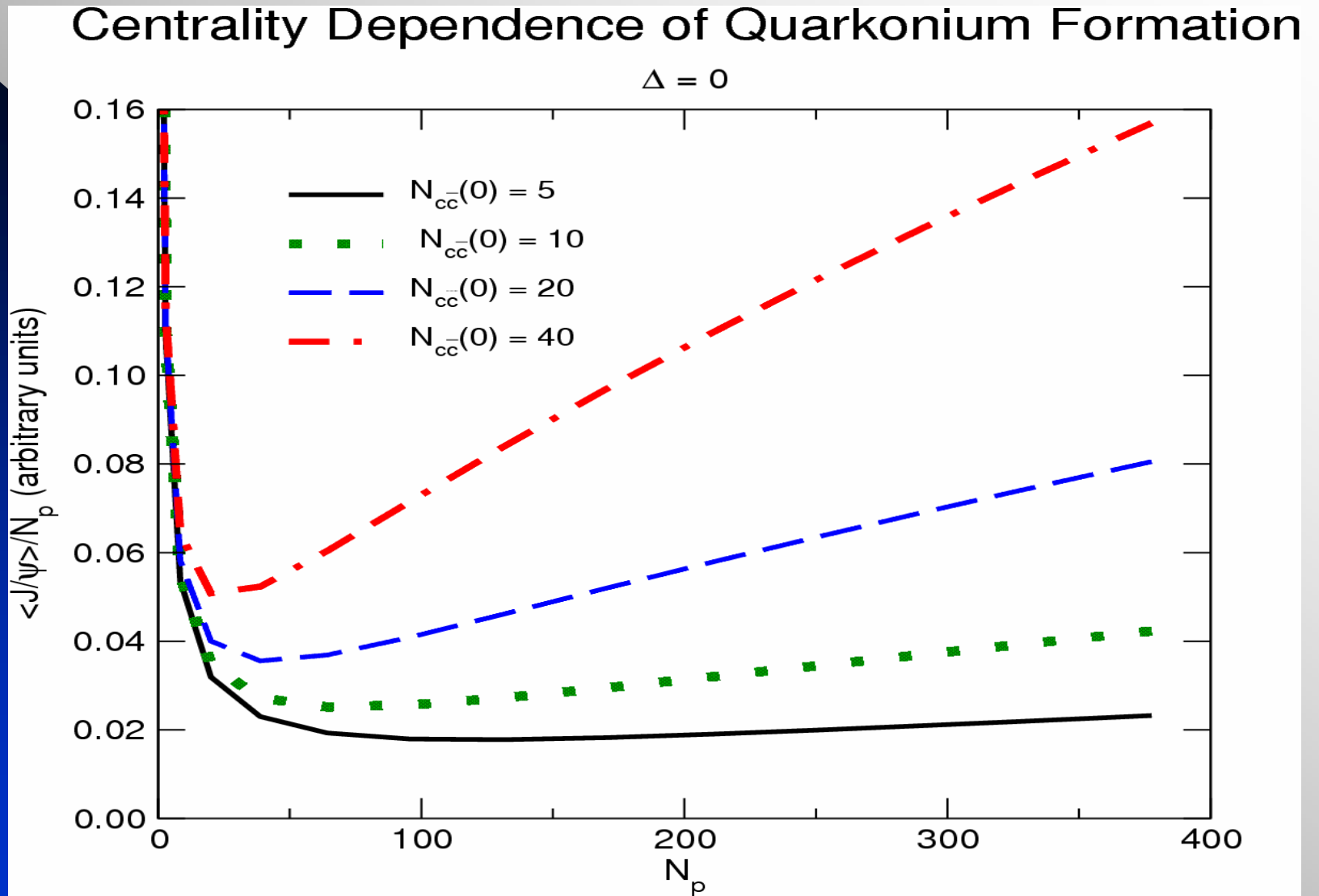
$\Delta = 0.33$



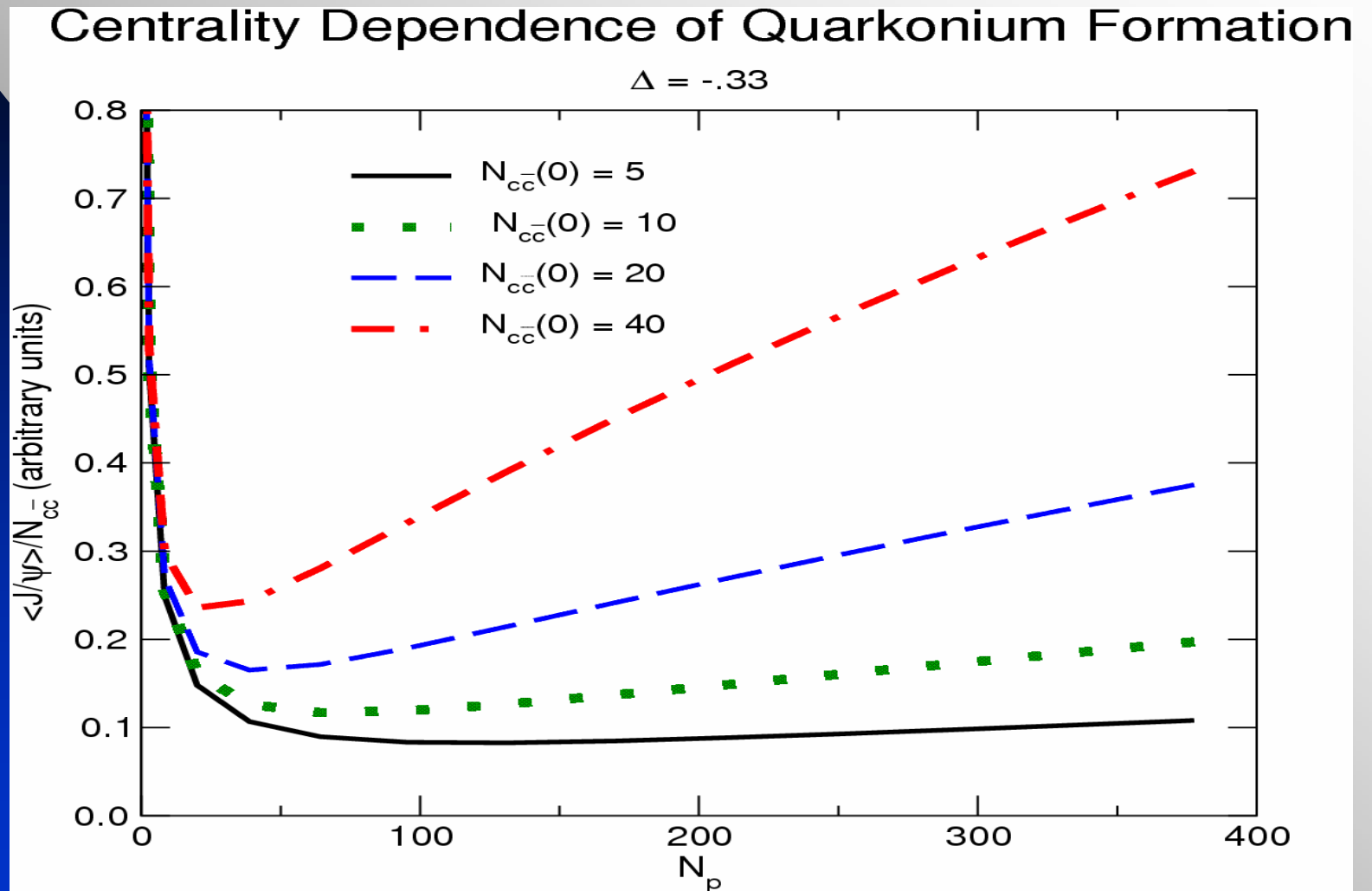
FORMATION CENTRALITY SIGNATURES



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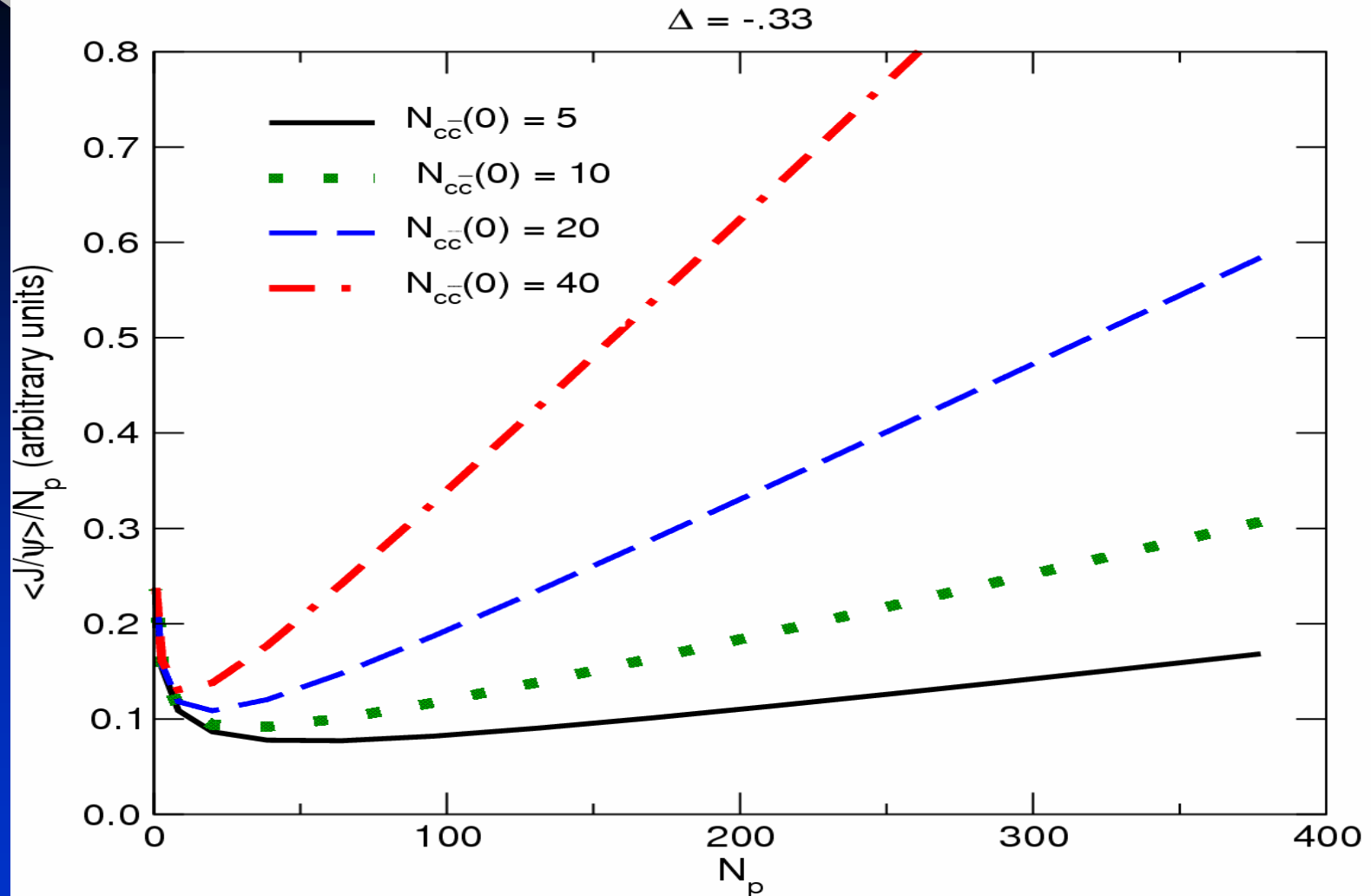


FORMATION CENTRALITY SIGNATURES



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Centrality Dependence of Quarkonium Formation

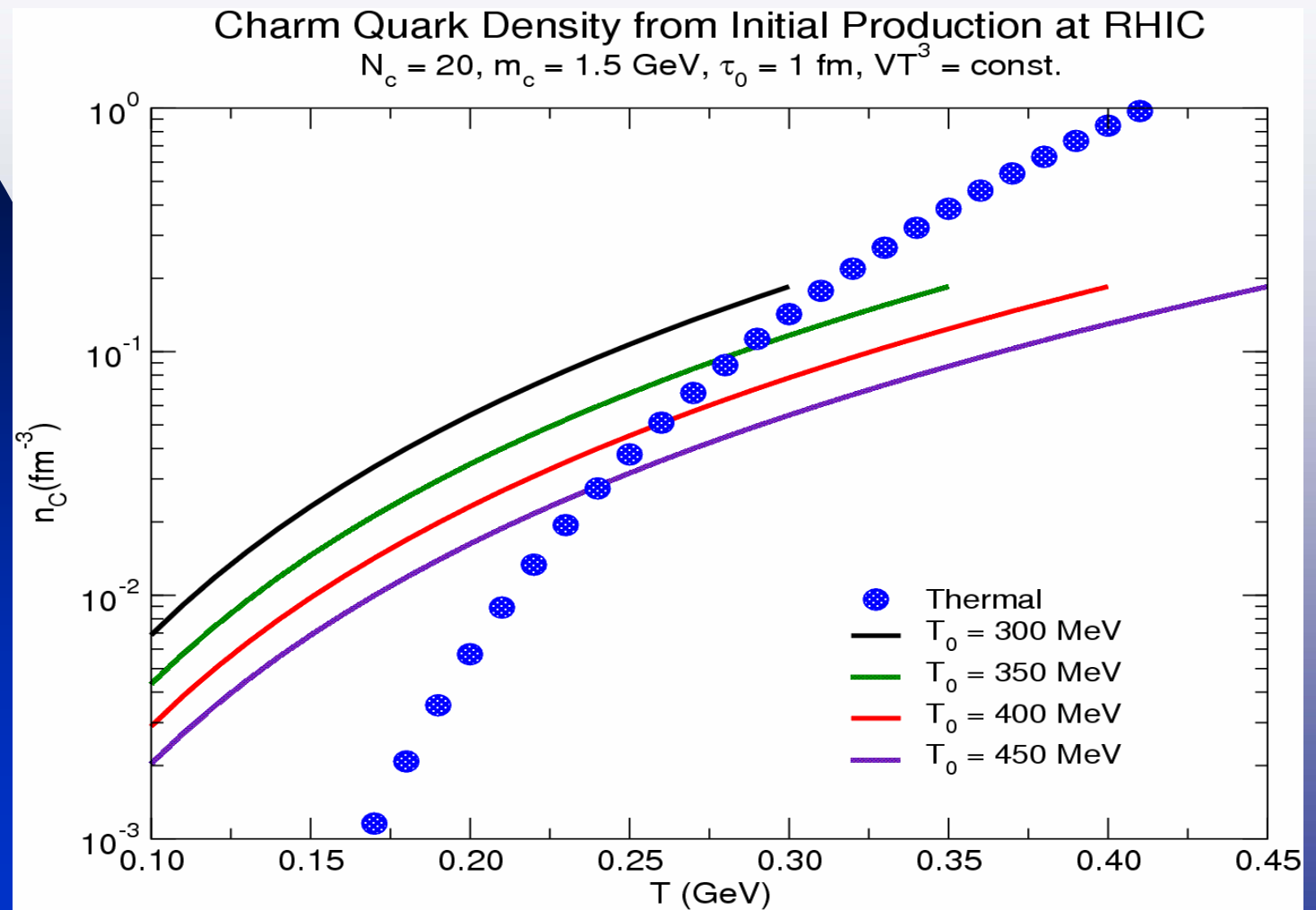


STATISTICAL HADRONIZATION MODEL

P. Braun-Munzinger, J. Stachel, Phys. Lett. B490, 196 (2000)

Chemical abundance of light hadrons fit with statistical model, parameters T , V , μ .
Hadrons with heavy quarks underpredicted.
Initial production of heavy quarks oversaturates chemical equilibrium.
Introduce charm enhancement factor γ_C , which is fixed by conservation of charm.
Then distribute charm quarks into hadrons according to statistical weights.

Charm density from Initial Production exceeds chemical equilibrium density for all $T < T_C$



$$N_{cc} = \frac{1}{2} \gamma_c N_{open}^{thermal} + \gamma_c^2 N_{hidden}^{thermal}$$

$$N_{J/\psi} = \gamma_c^2 N_{J/\psi}^{thermal}$$

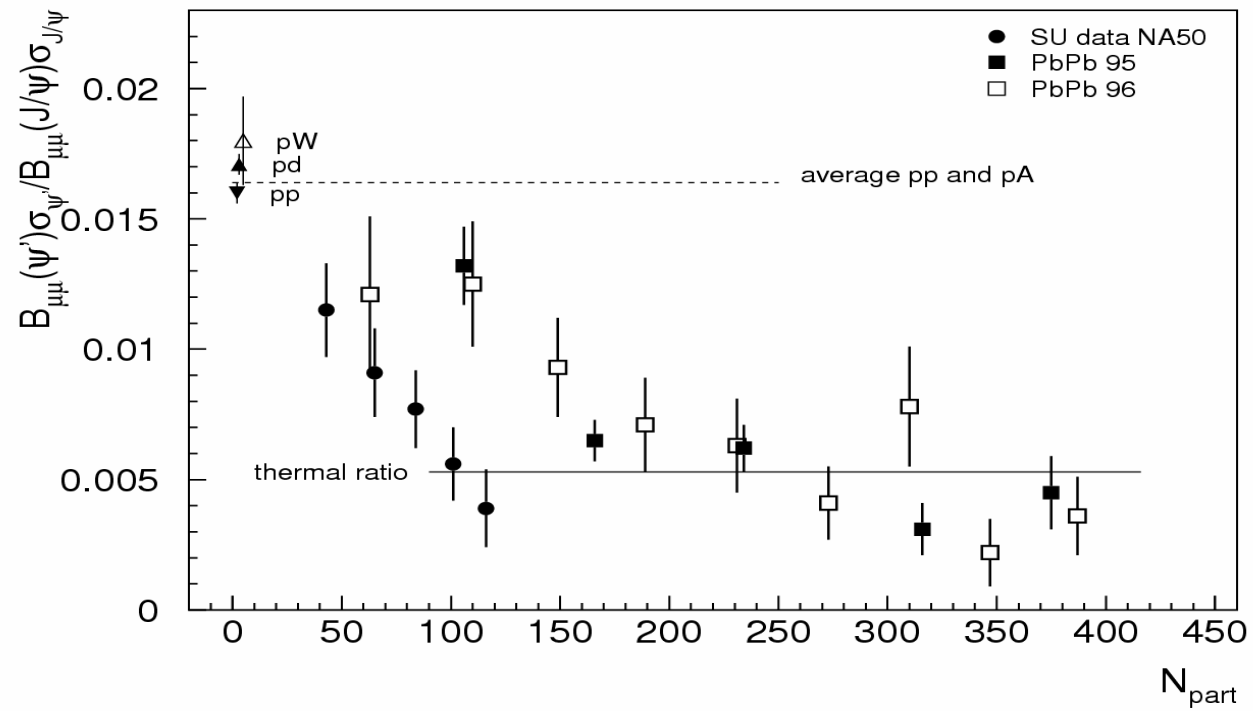
$$N_{canonical} = N_{gc} I_1(N_{gc}) / I_0(N_{gc})$$

$$\rightarrow N_{gc}, N_{gc} \gg 1$$

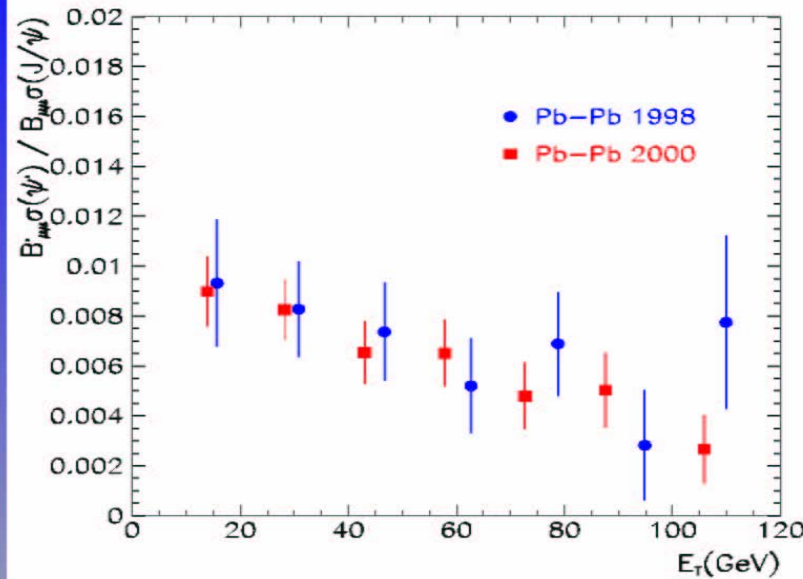
$$\rightarrow \frac{1}{2} N_{gc}^2, N_{gc} \ll 1.$$

$$\gamma_c \rightarrow 2 N_{cc} / N_{open}^{thermal}, \text{ gc limit}$$

$$\rightarrow 2 \sqrt{N_{cc}} / N_{open}^{thermal}, \text{ canonical limit}$$



ψ'/ψ as a function of E_T



The ratio of the **two charmonium states** decreases with centrality by a factor of 2.5 between peripheral and central collisions

The ψ' is more suppressed than the J/ψ in **Pb-Pb collisions**

$$N_{J/\psi} = 4 \left[\frac{n_{ch}^{therm} n_{J/\psi}^{therm}}{(n_{open}^{therm})^2} \right] \frac{N_{cc}^2}{N_{ch}}$$

Hadronization Volume $(N_{ch}/n_{ch})^{therm}$ is physical parameter

Ratio of thermal densities slowly-varying with Temperature

SPS (NA50): Requires charm enhancement approx 3x

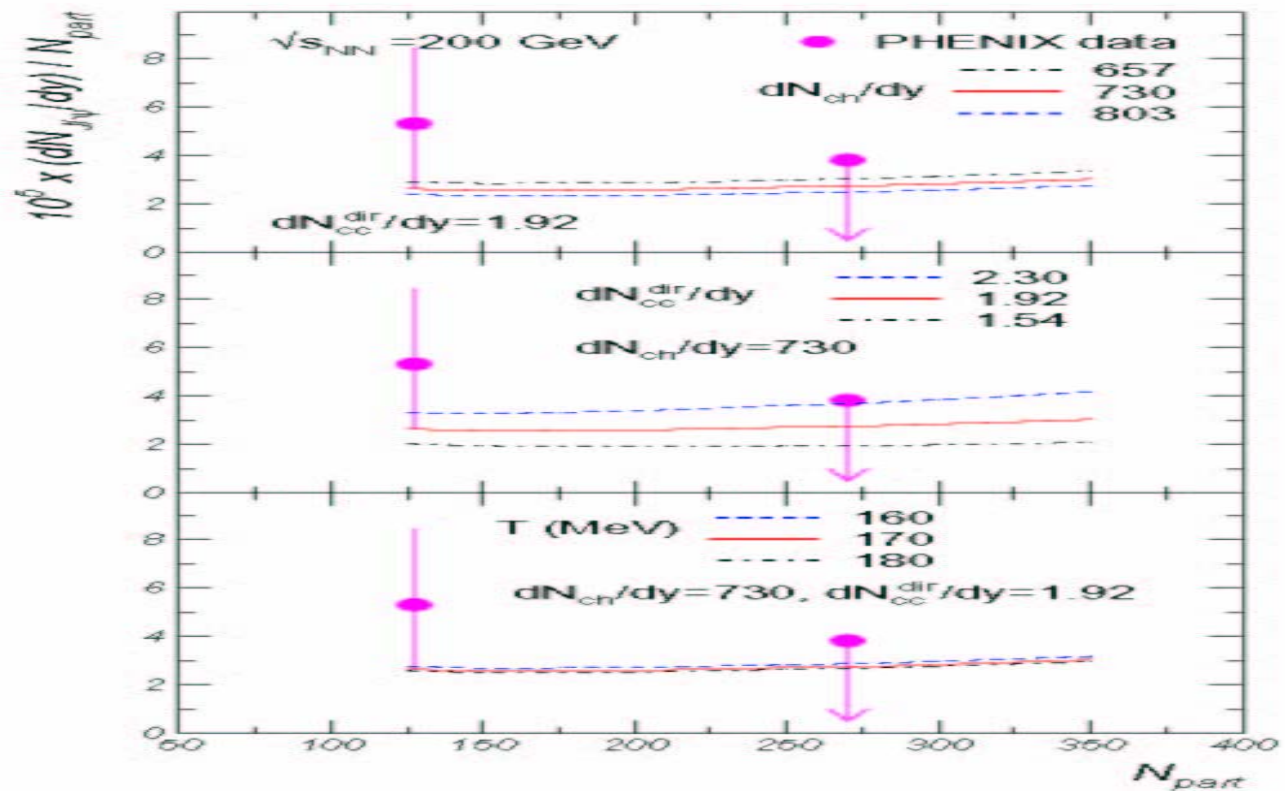
RHIC & LHC: Express in terms of rapidity densities

Additional predictions: Open charm mesons and baryons

A Andronic, P. Braun-Munzinger, K. Redlich, J. Stachel,

Phys. Lett. B571:36-44, 2003

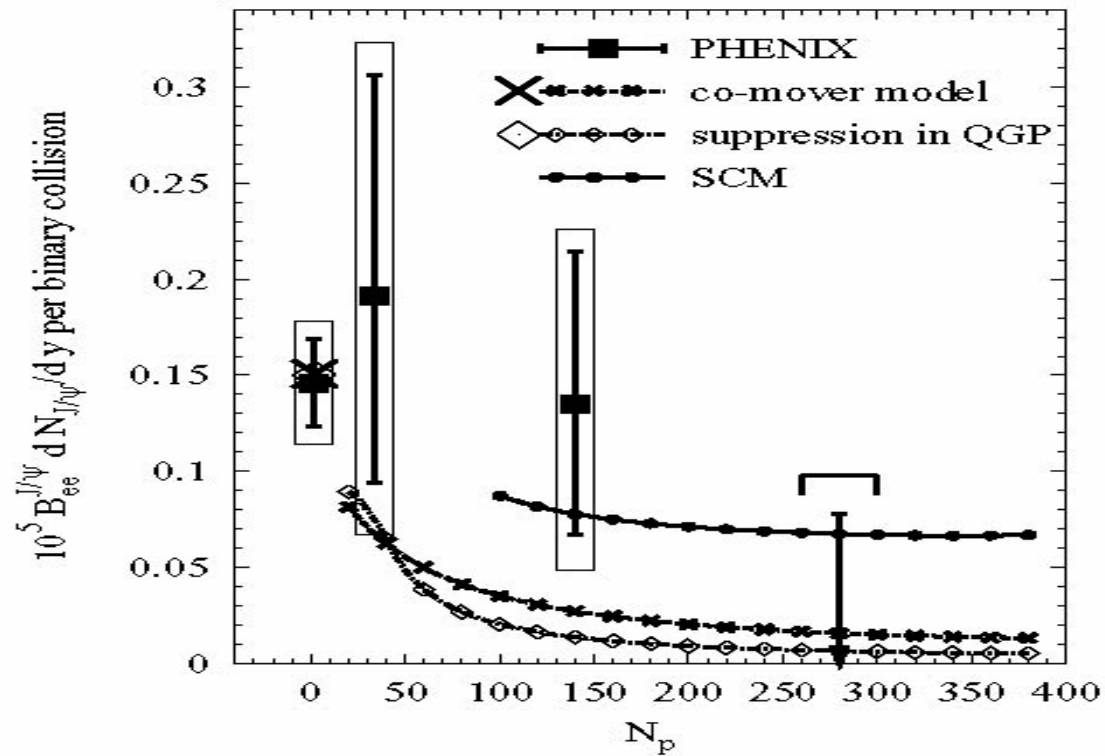
$$\sigma_{cc} = 390 \mu\text{b}$$



A.P. Kostyuk, M.I. Gorenstein, H. Stocker, W. Greiner

Phys. Rev. C68: 041902, 2003

$$\sigma_{cc} = 650 \mu\text{b}$$



FORMATION OF QUARKONIUM IN REGION OF COLOR DECONFINEMENT

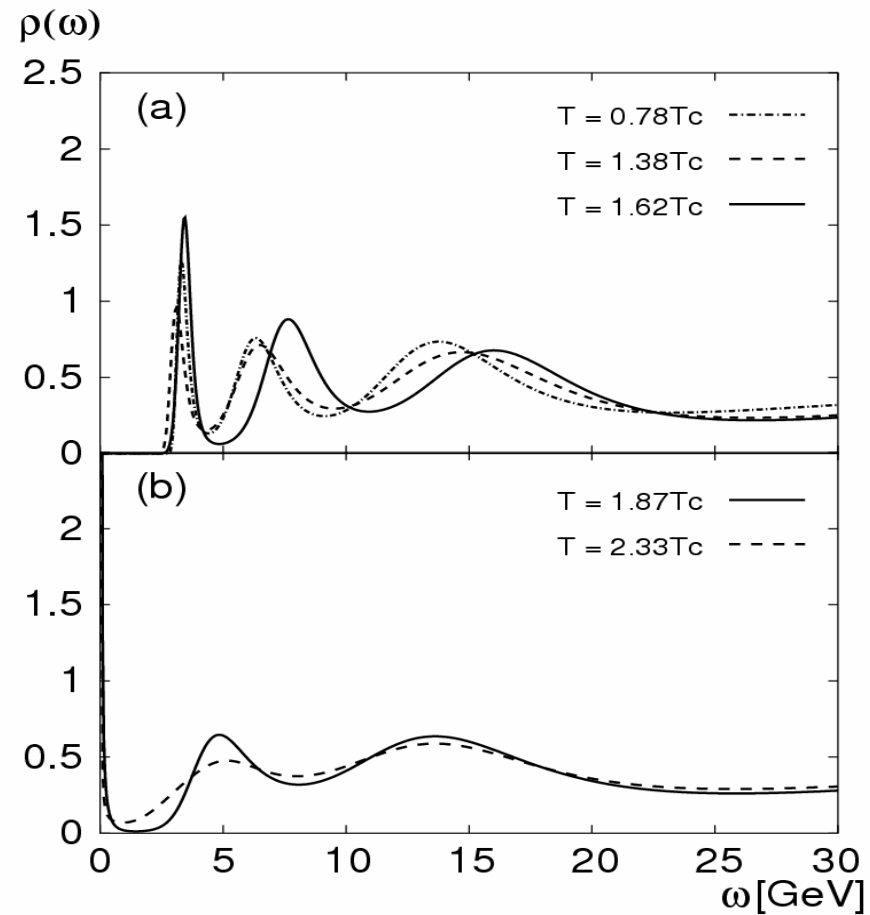
R. L. Thews, M. Schroedter, J. Rafelski, Phys. Rev. C63:054905, 2001

Formation process is **Inverse** of dissociation

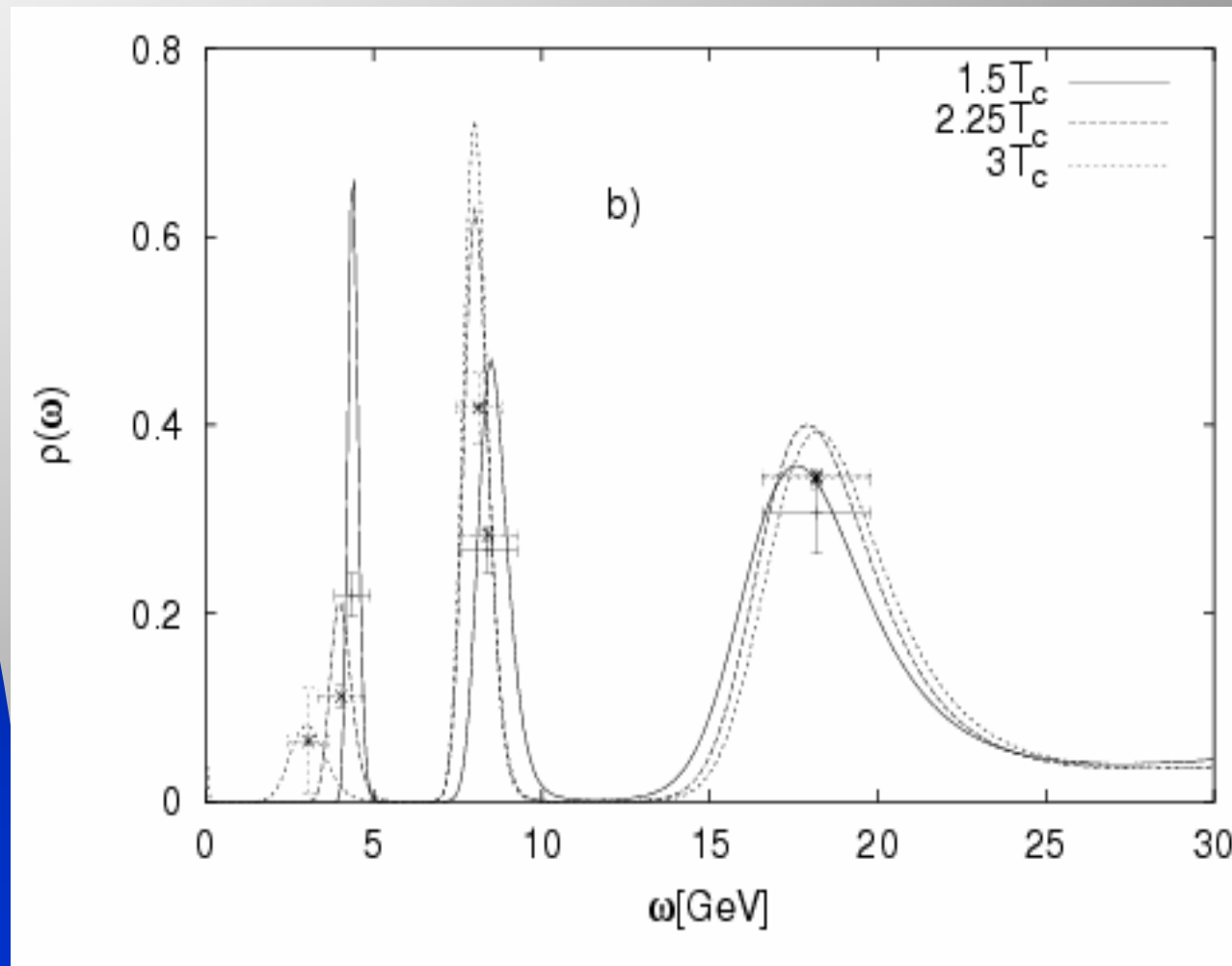
Model evolution of region with initial temperature and isentropic expansion, depends on contours of participant density

Final population determined by competition between formation and dissociation rates

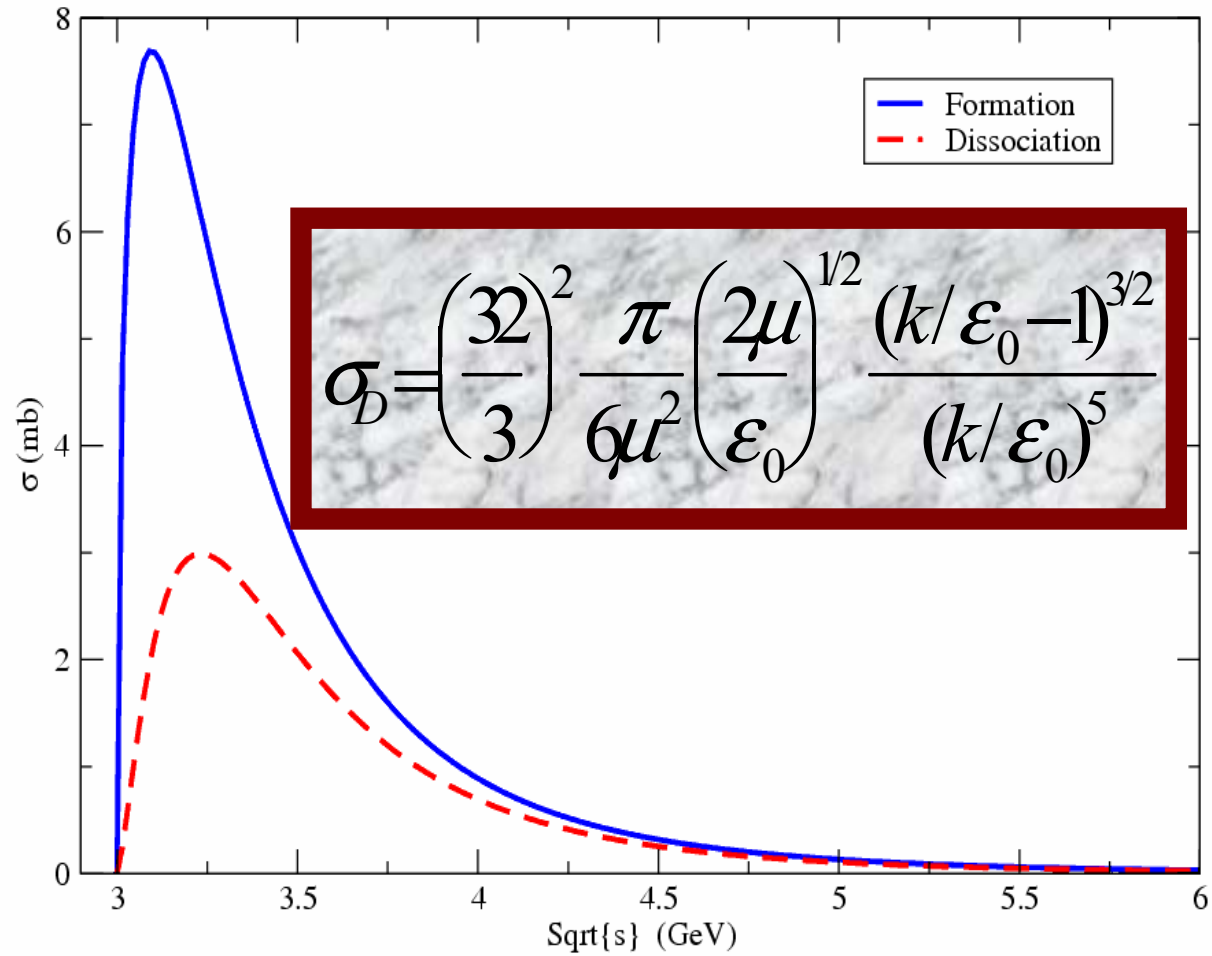
M. Asakawa and T. Hatsuda, Phys. Rev. Lett: 012001 (2004)



S. Datta, F. Karsch, P. Petreczky, I. Wetzorke, Phys. Rev. D69:094507 (2004)

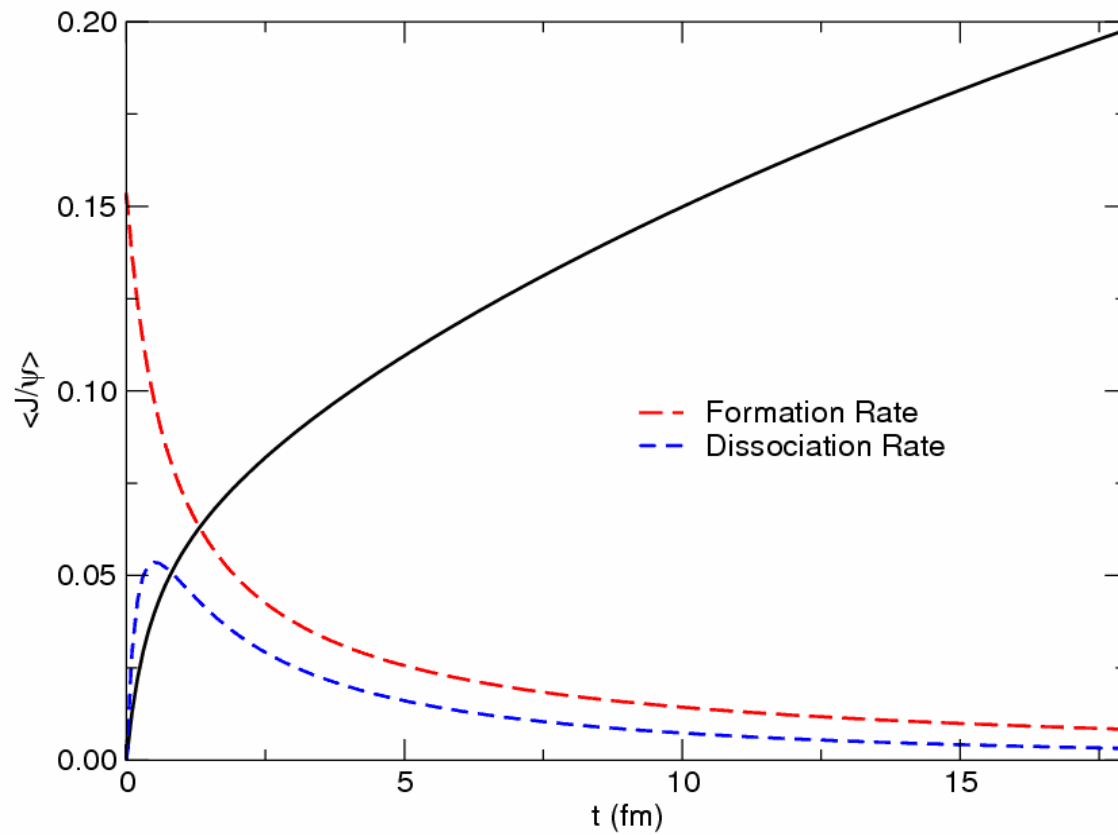


OPE Cross Sections $g + J/\psi \leftrightarrow c + \bar{c}$



$$\frac{dN_{J/\psi}}{d\tau} = \langle v\sigma_F \rangle \rho_c N_c - \langle v\sigma_D \rangle \rho_g N_{J/\psi}$$

Evolution of Charmonium Formation and Dissociation Rates



If $N_{J/\psi} \ll N_{cc}$, solution is:

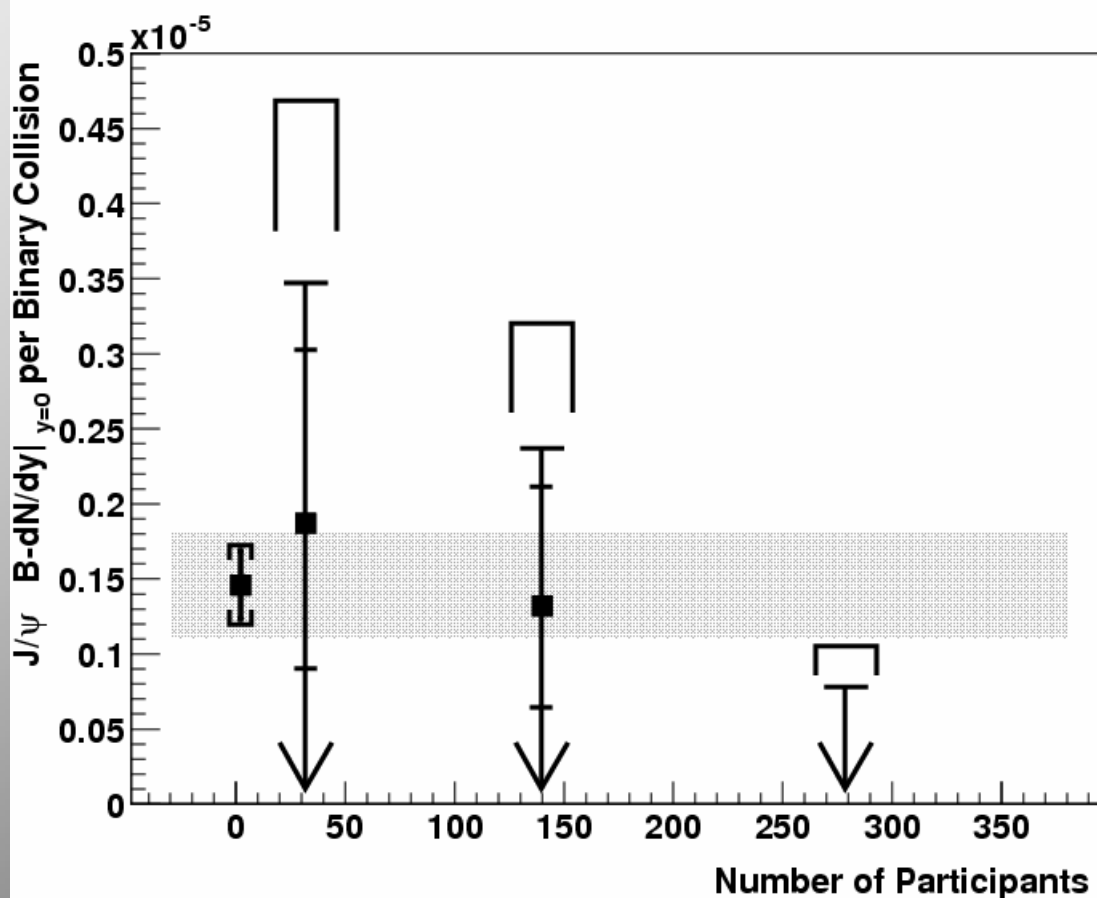
$$N_{J/\psi}(\tau_f) = \varepsilon(\tau_f) \times N_{cc}^2 \int_{\tau_0}^{\tau_f} \lambda_F [V(\tau) \varepsilon(\tau)]^1 d\tau$$

$$\varepsilon(\tau) = \exp\left[-\int_{\tau_0}^{\tau} \lambda_D \rho_g d\tau\right]$$

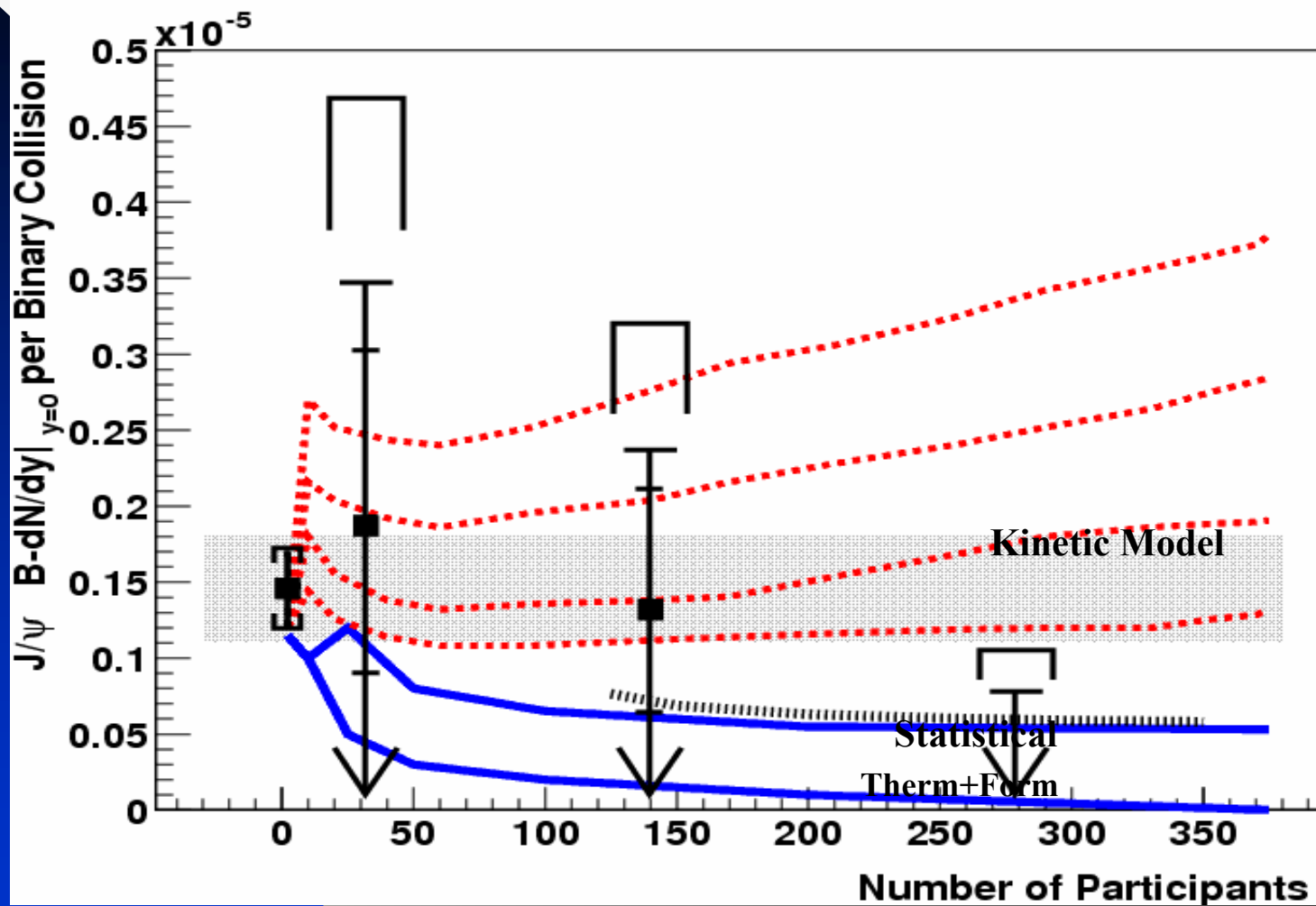
COMPARISON WITH INITIAL PHENIX DATA AT RHIC 200

Rates very sensitive to quark momentum distribution
Centrality signature varies with magnitude of N_{cc}

PHENIX – Phys. Rev. C69, 014901 (2004)

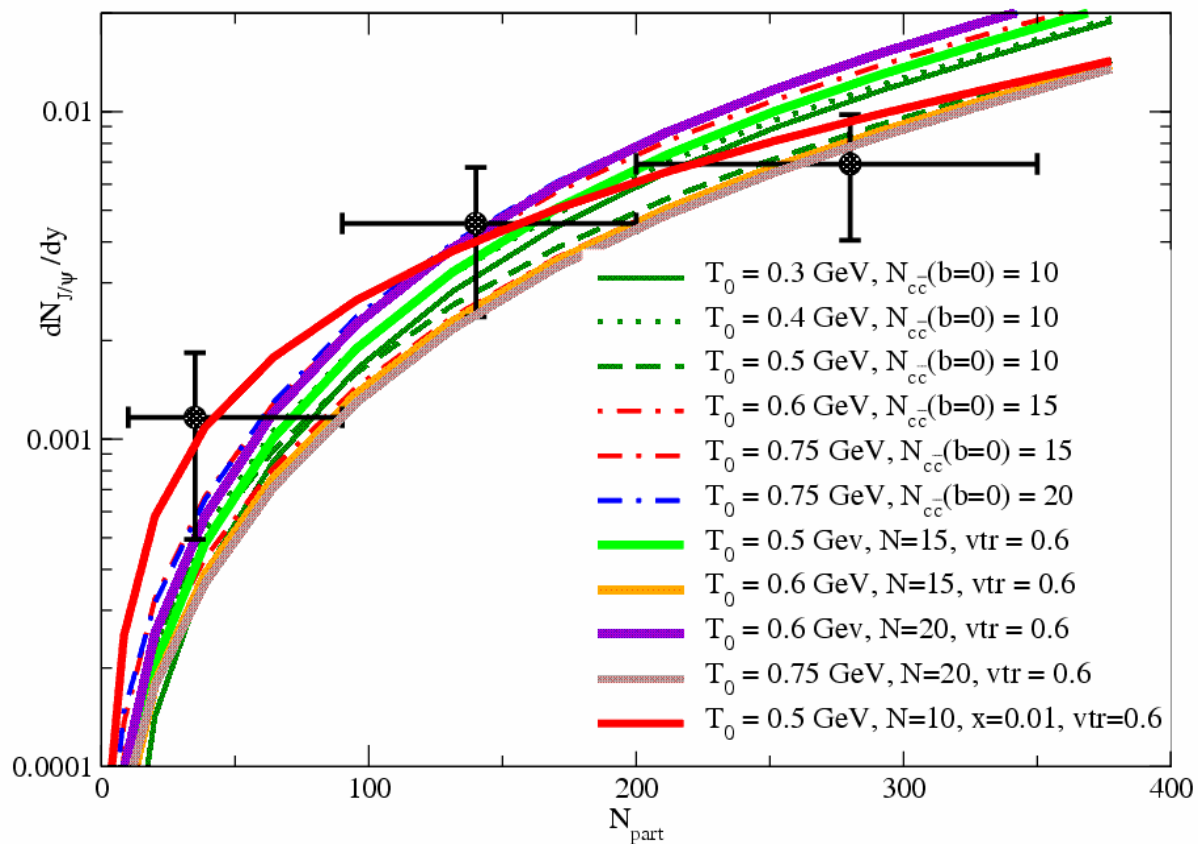


Model predictions very sensitive to N_{cc} and distribution

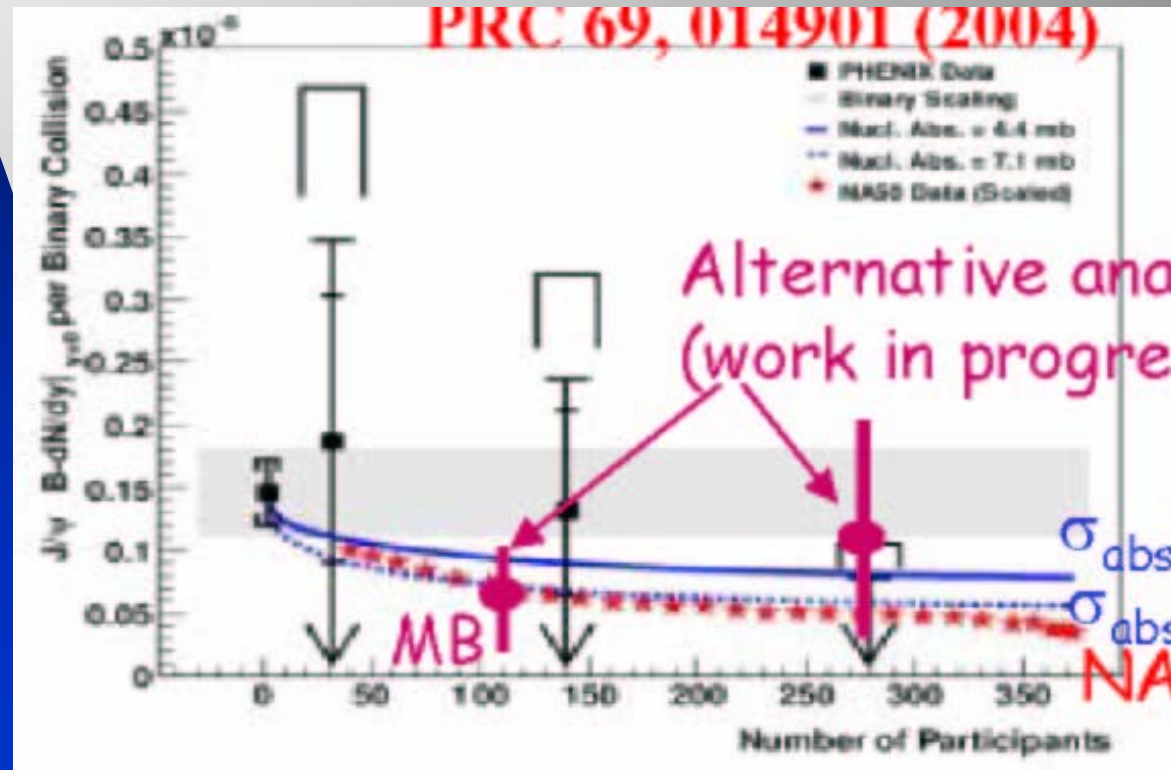


R. L. Thews, J. Phys. G30: S369 (2004)

J/ψ Formation at RHIC, $x=0$



T. Gunji, JPS 2004, PHENIX (prelim)



DO THE Y AND P_T SPECTRA PROVIDE A FORMATION SIGNATURE?

M. Mangano and R. L. Thews (work in progress)

1. Generate sample of $c\bar{c}$ pairs from NLO pQCD (smear LO q_t)
2. Supplement with k_t to simulate initial state and confinement effects
3. Integrate formation rate using these events to define particle distributions

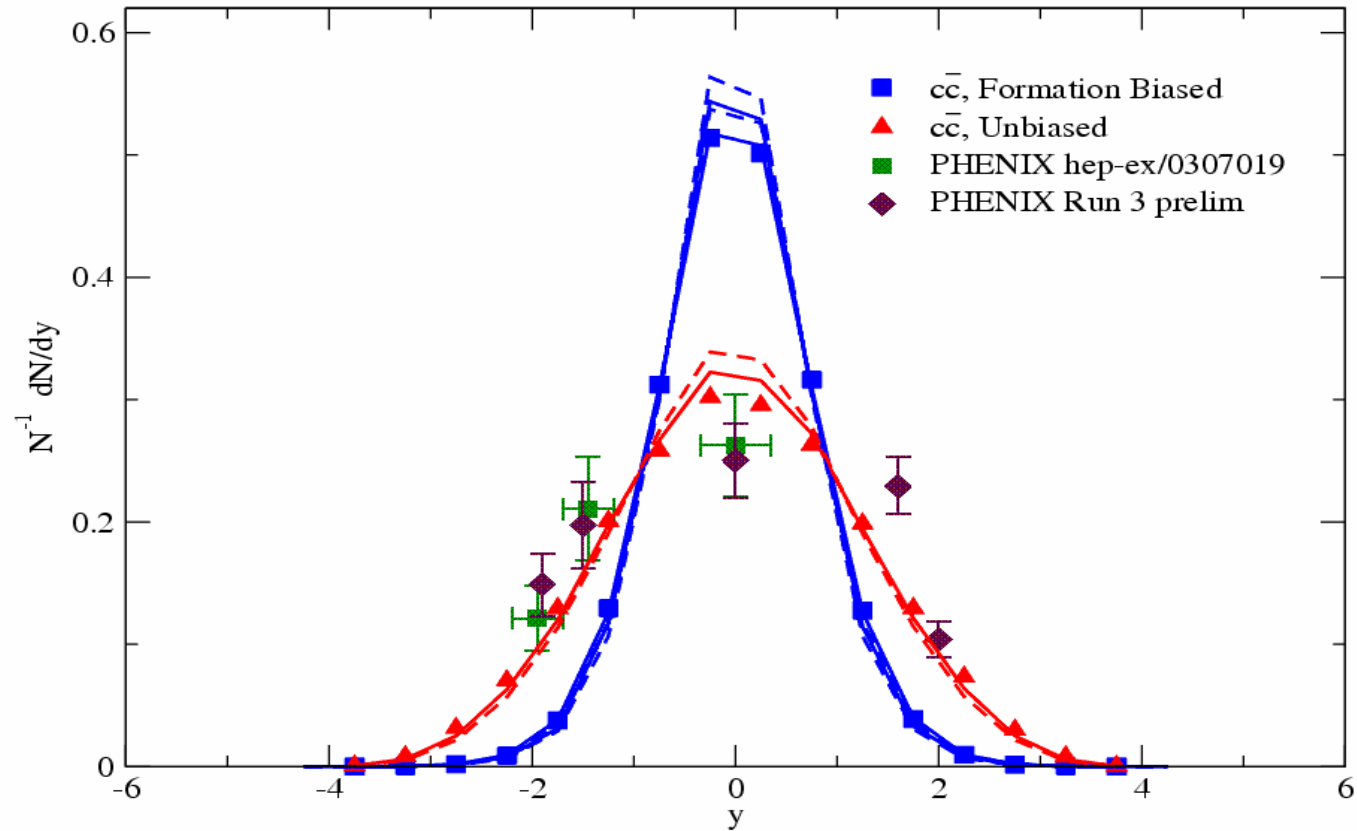
$$\frac{dN_{J/\psi}}{d^3p_{J/\psi}} = \int \frac{dt}{V(t)} \sum_{i=1}^{N_{\bar{c}}} \sum_{j=1}^{N_{\bar{c}}} v_{rel} \frac{d\sigma(p_i + p_j \rightarrow p_{J/\psi} + X)}{d^3p_{J/\psi}}$$

- All combinations of c and cbar contribute
- Prefactor is integrated flux for given pair
- Total has expected $(N_{c\bar{c}})^2 / V$ behavior

p-p data “select” unbiased diagonal c-cbar pairs

J/ ψ Formation in pp Interactions at RHIC200

84K NLO diagonal $c\bar{c}$ pairs, variable k_t broadening

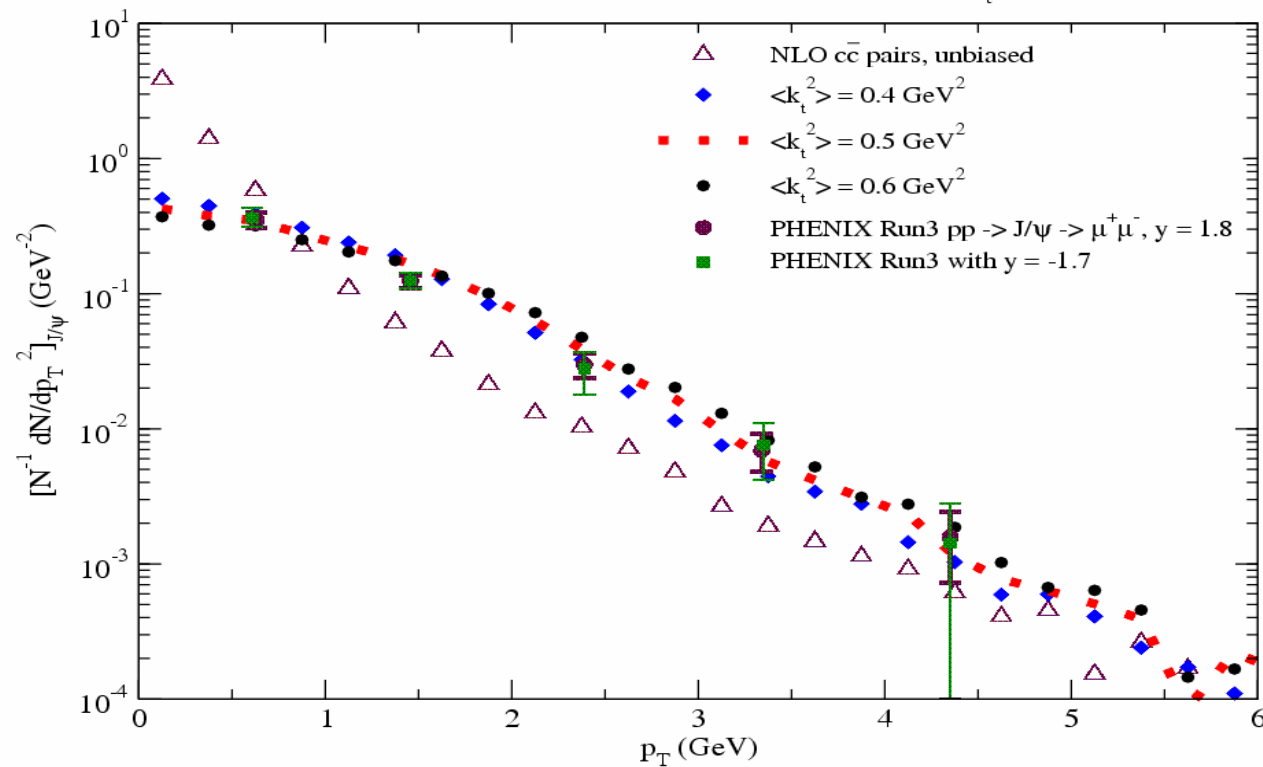


p-p data determine intrinsic k_t

$$\langle k_t^2 \rangle_{c\text{-quarks}} = 0.5 \pm 0.1 \text{ GeV}^2$$

J/ ψ Formation in pp Interactions at RHIC200

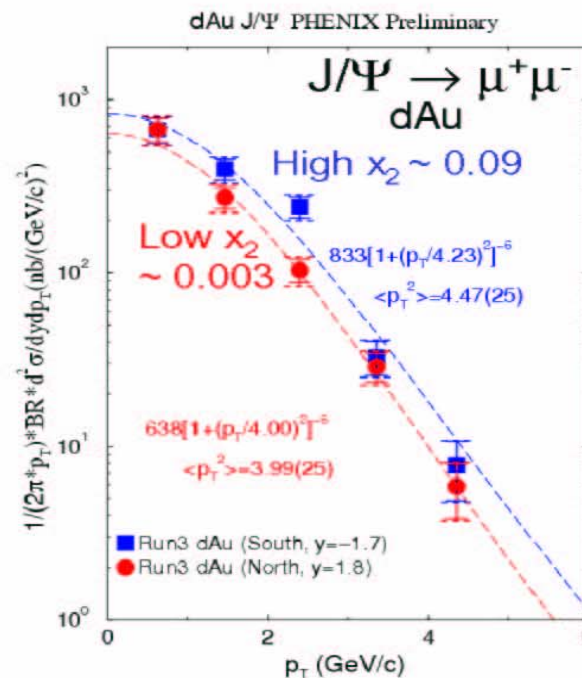
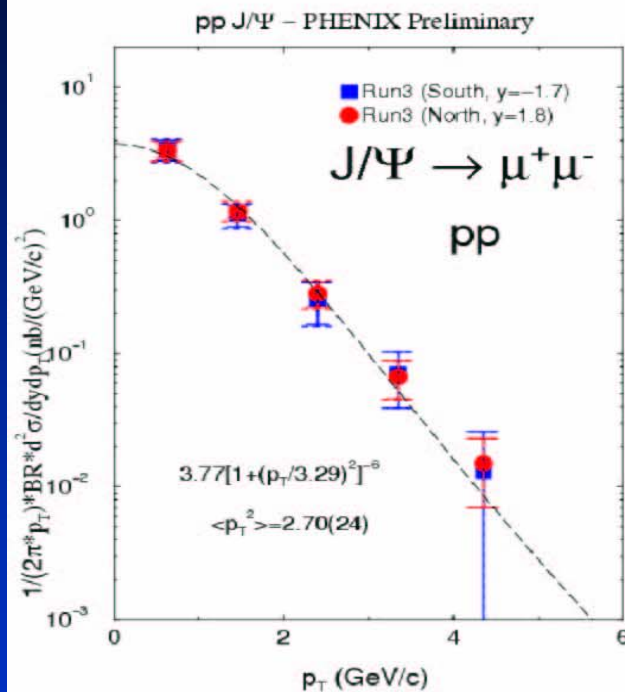
1.2 < y < 2.4, 84K NLO diagonal $c\bar{c}$ pairs, Sensitivity to k_t broadening



Use dAu broadening to determine nuclear k_t

$$\Rightarrow \langle k_t^2 \rangle_{AA} = 1.3 \pm 0.3 \text{ GeV}^2$$

Cross section versus p_T



$$\Delta \langle p_T^2 \rangle =$$

$$\langle p_T^2 \rangle_{dAu} - \langle p_T^2 \rangle_{pp}$$

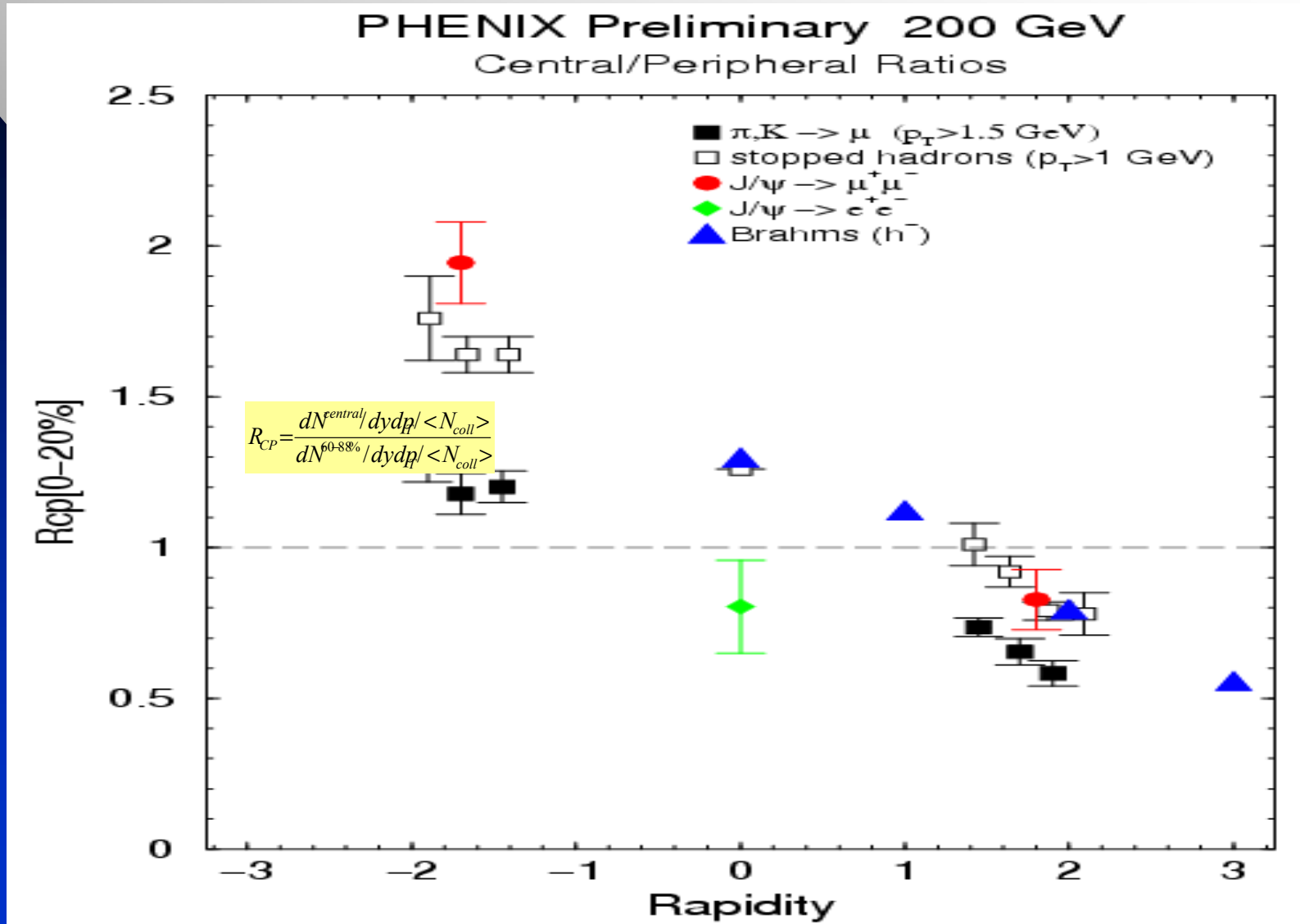
$$1.77 \pm 0.35 \text{ GeV}^2$$

$$1.29 \pm 0.35 \text{ GeV}^2$$

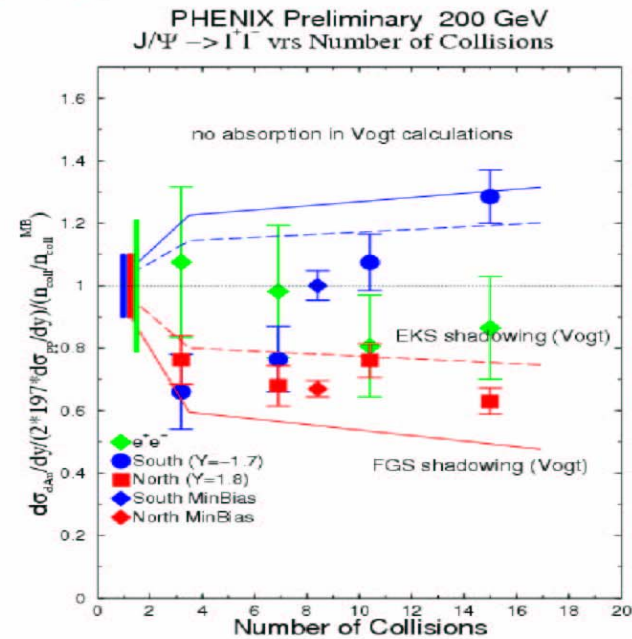
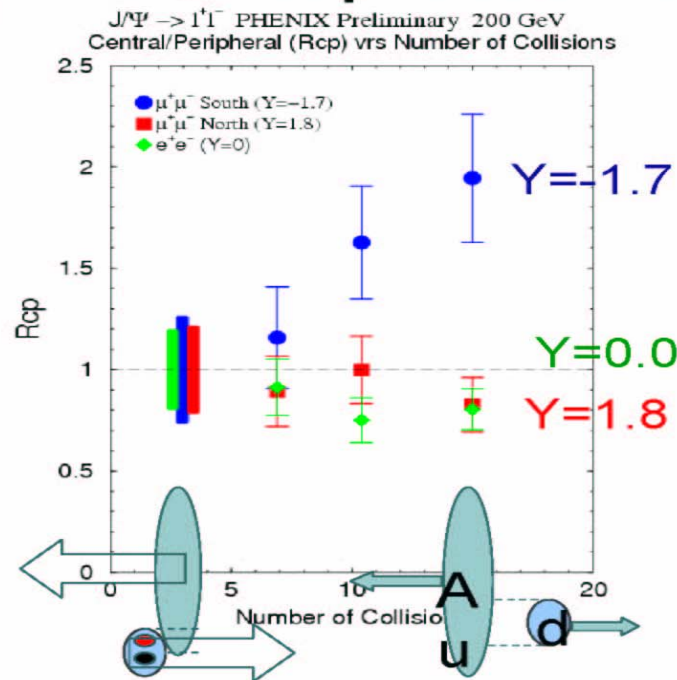
(preliminary)

p_T is broadened for dAu

P. Steinberg, Hot Quark 2004 Workshop, July 2004



Rcp and RdAu

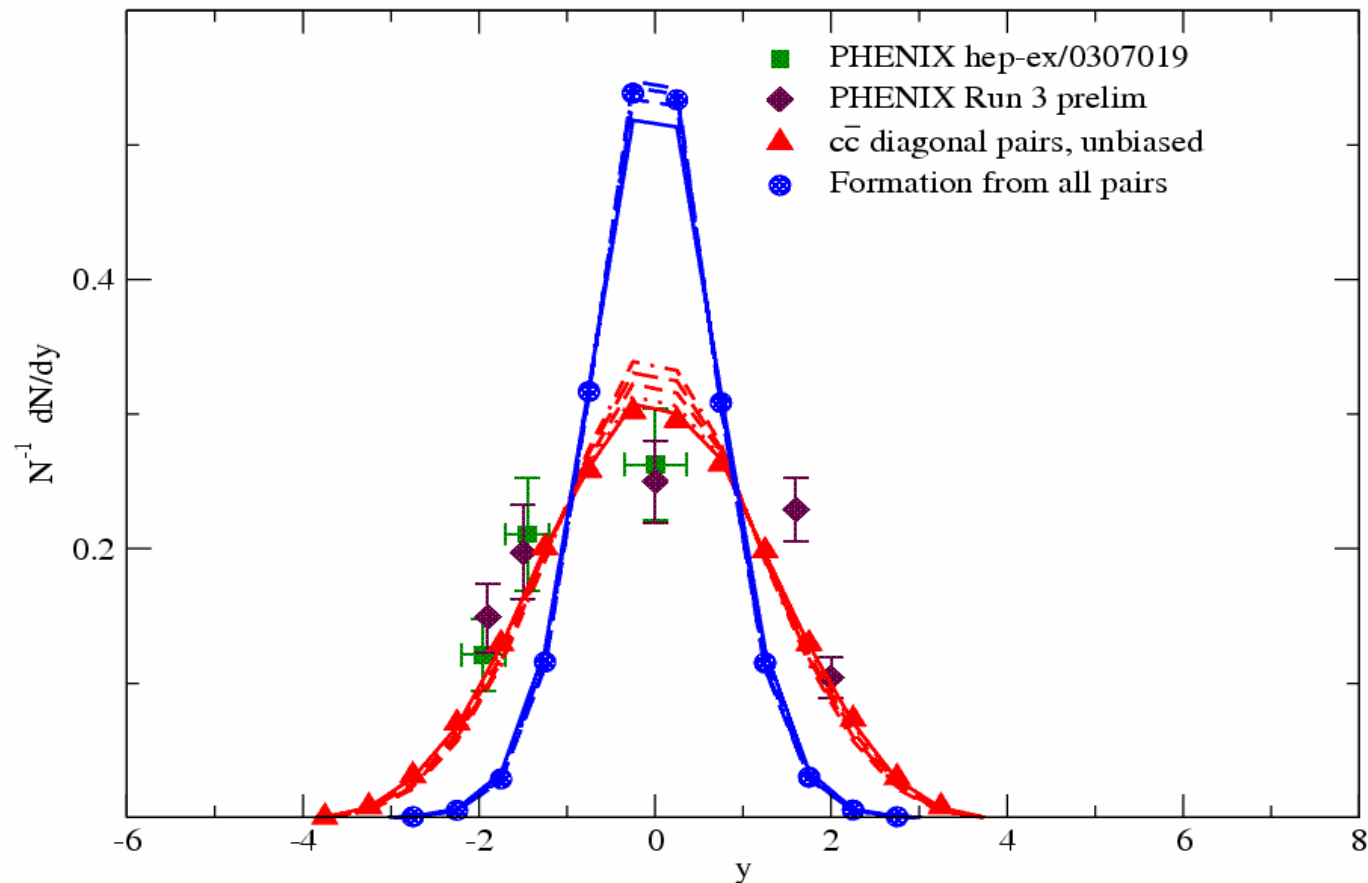


- RdA: average over centrality consistent with minimum bias result.
- Weak nuclear effects at forward rapidities (green points at midrapidity)
- Stronger centrality dependence at backward rapidities

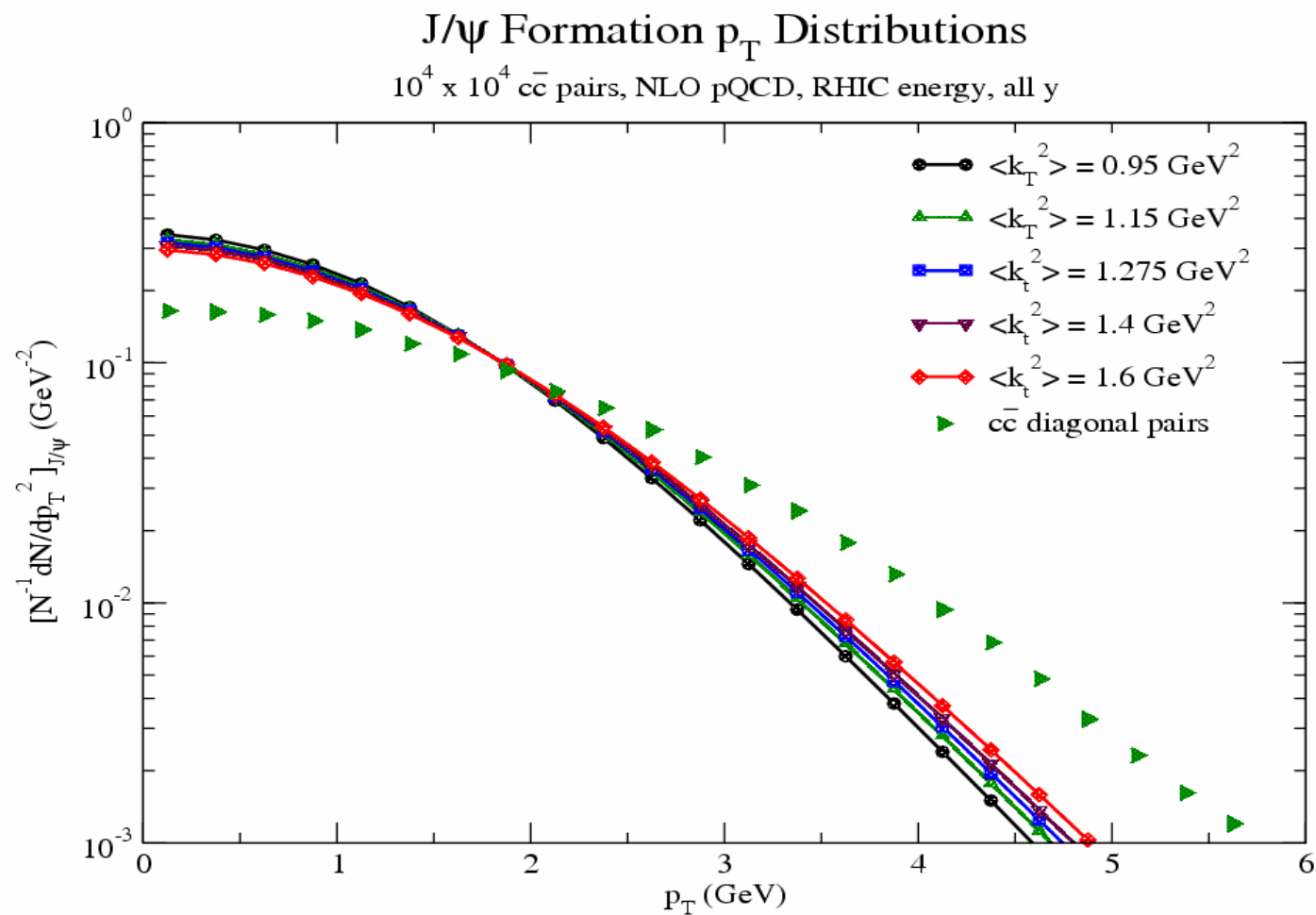
Formation through “off-diagonal” pairs narrows rapidity distribution

J/ψ Formation in AA Interactions at RHIC200

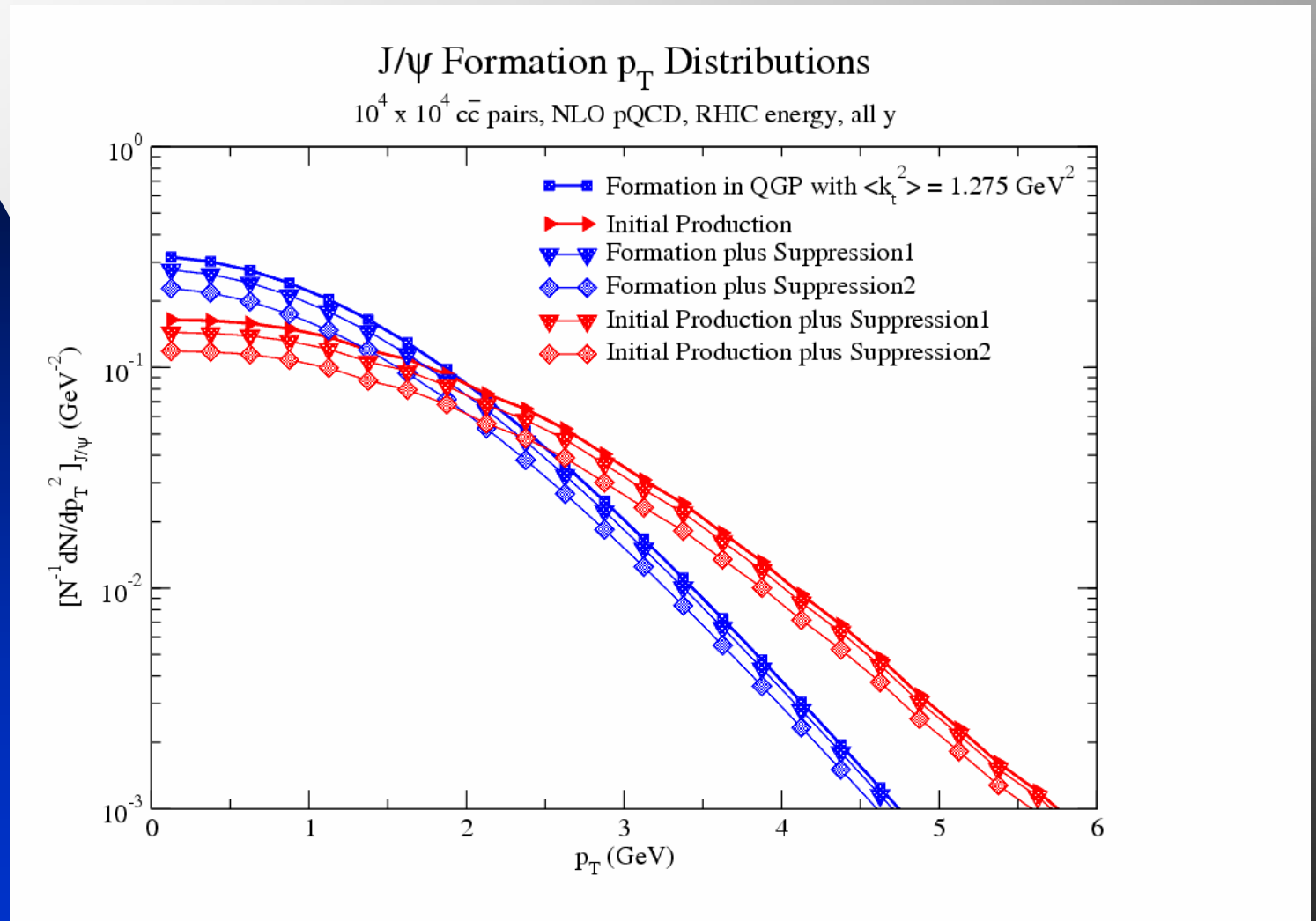
Normalized Rapidity Distributions, $10^4 \times 10^4$ NLO $c\bar{c}$ pairs



Formation through “off-diagonal” pairs narrows p_t distribution



Suppression of formed or initial J/ψ in partonic medium



Comparison with Thermal + Transverse Flow c-Quark Distributions

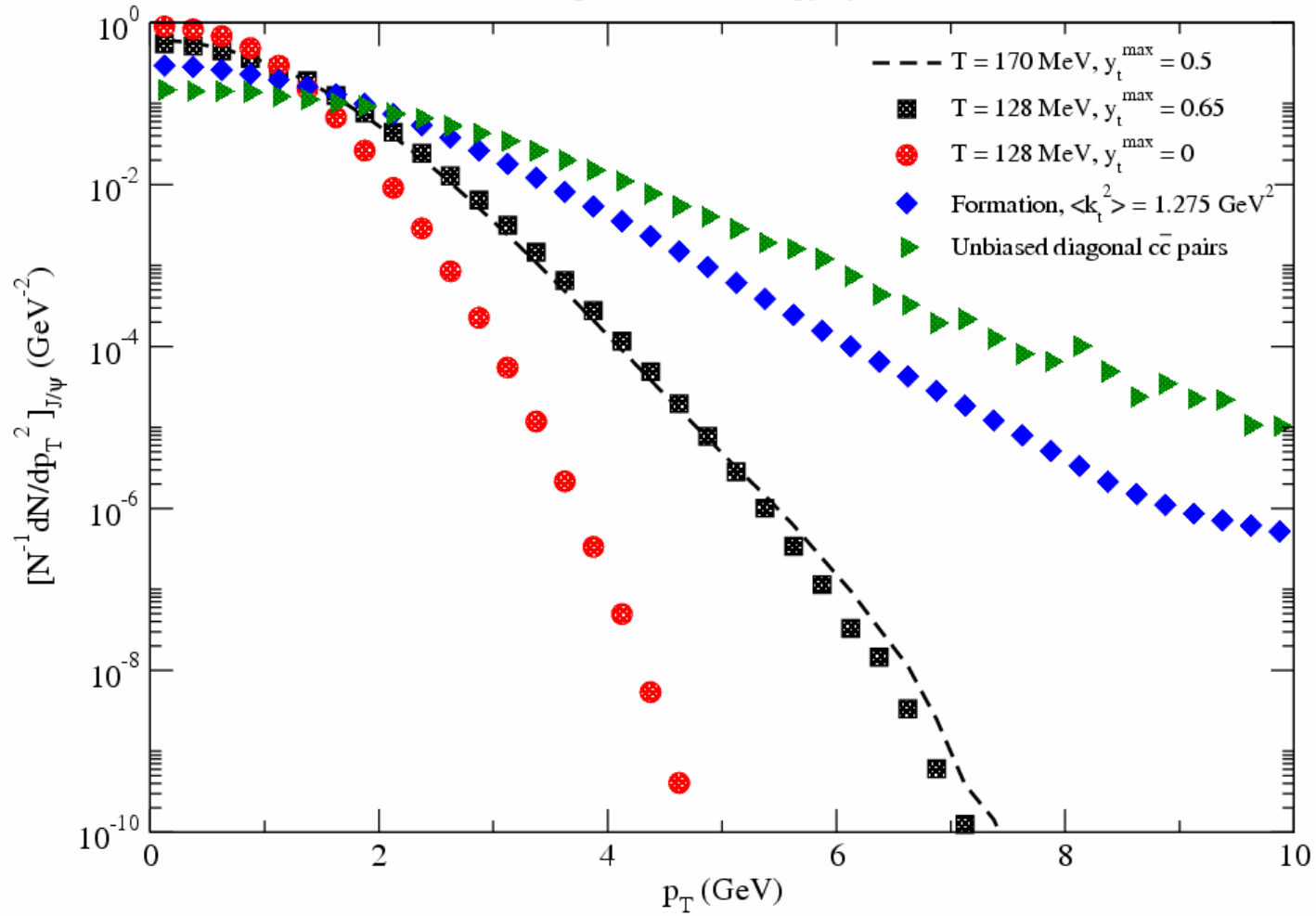
K.A.Bugaev, M. Gazdzicki, M.I.Gorenstein, Phys.Lett.B544,127(2002)

S.Batsouli, S.Kelly, M.Gyulassy, J.L.Nagle, Phys.Lett.B557,26 (2003)

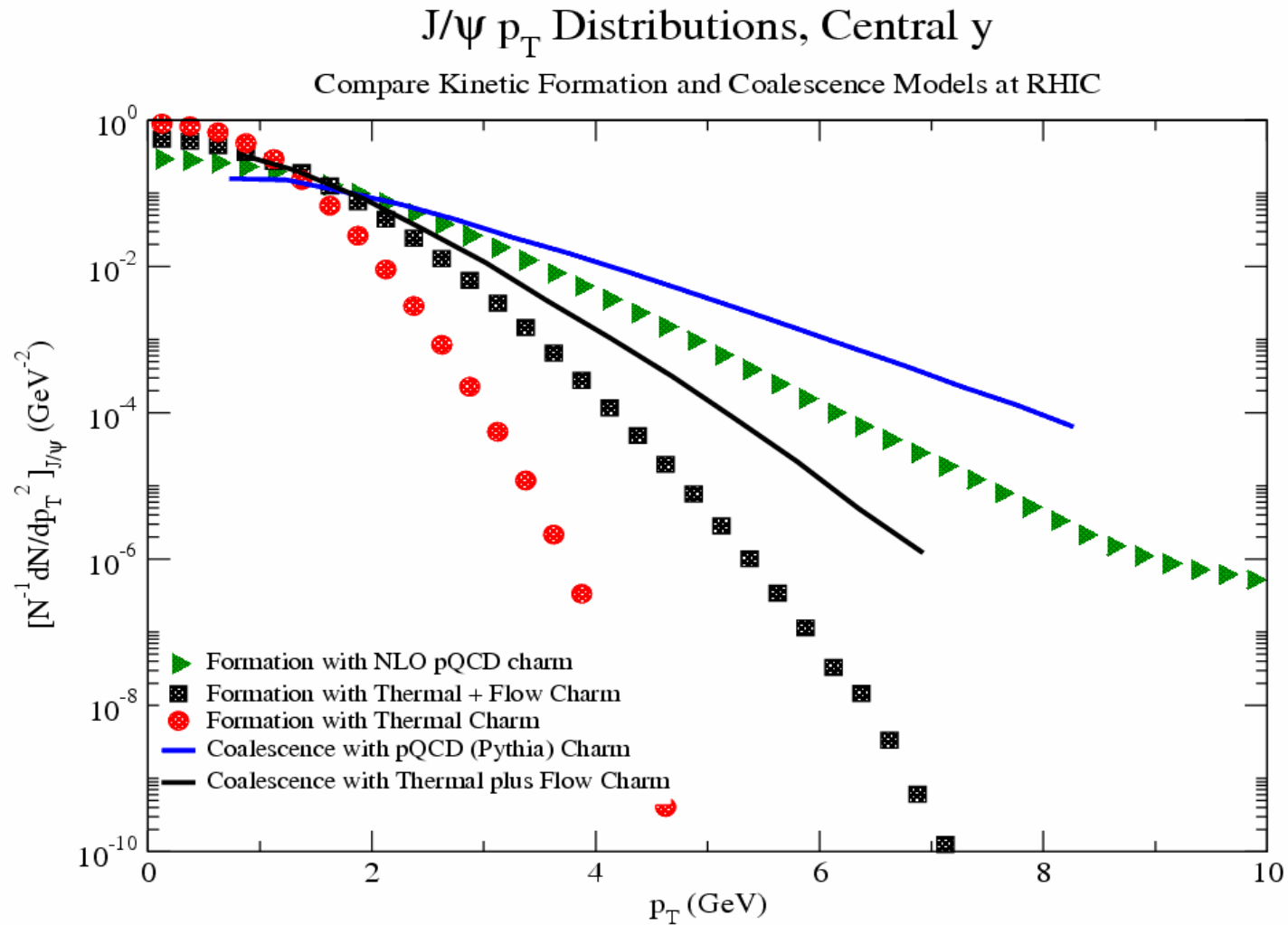
$$\frac{dN}{dp_t^2} \propto m_t \int_0^R r dr I_0\left[\frac{p_t \sinh\left(\frac{r}{R} y_t^{\max}\right)}{T}\right] K_1\left[\frac{m_t \cosh\left(\frac{r}{R} y_t^{\max}\right)}{T}\right]$$

J/ψ Formation p_T Distributions

$10^4 c\bar{c}$ pairs, RHIC energy, $|y| < 0.35$



Comparison with coalescence model: V Greco, C. M. Ko, R. Rapp, Phys. Lett. B595:202 (2004)



SUMMARY

- Absolute magnitude and centrality dependence tests require both open and hidden flavor
- P_t and y signature of kinetic process
- Both mechanisms contribute to total production
- Stat hadronization also predicts open/hidden
- Formation process very sensitive to quark flow

