

A modified Balitsky-Kovchegov equation

Eugene Levin, Tel Aviv University



LHC-HERA WS, November 2004

The Goal:

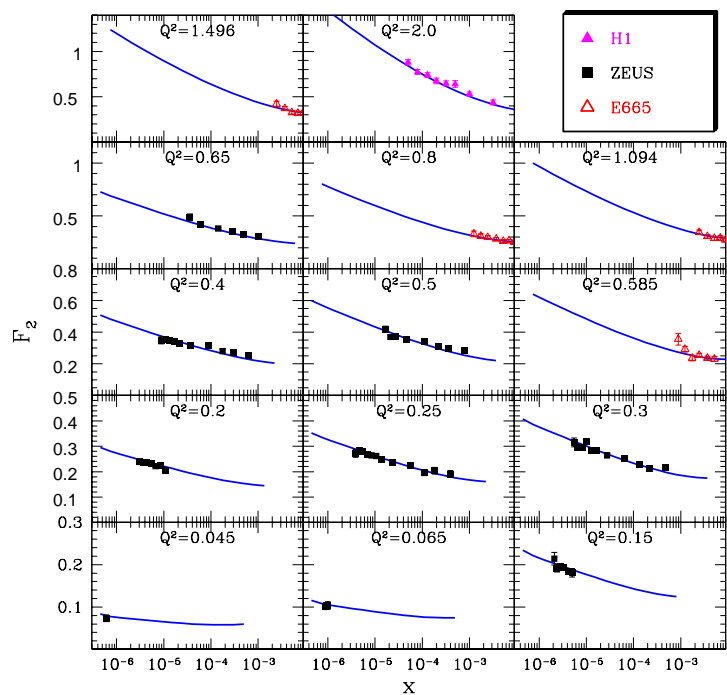
Reliable calculations for parton densities at the LHC energies based on the high parton density QCD theory

Sources:

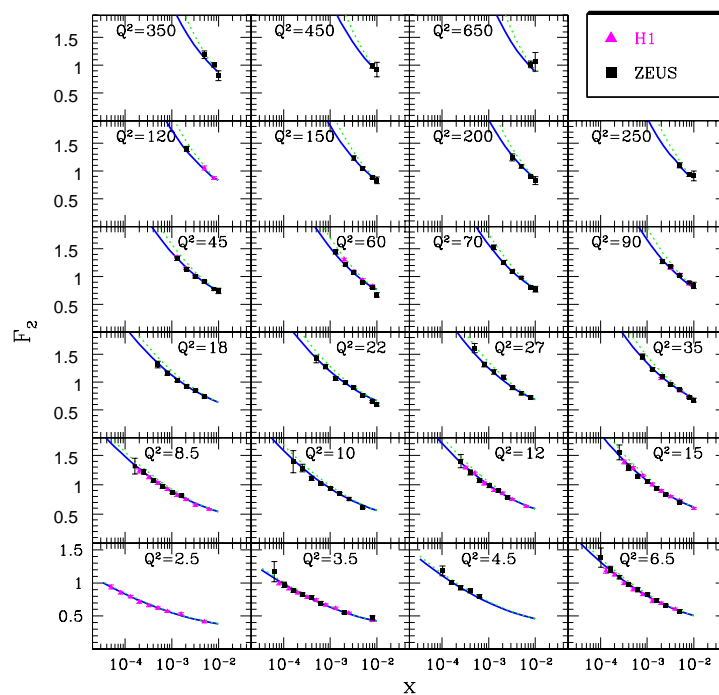
- E. Gotsman, E. Levin, M. Lublinsky and U. Maor: “Towards a new global QCD analysis: Low x DIS data from non-linear evolution,” *Eur. Phys. J. C* 27 (2003) 411 [arXiv:hep-ph/0209074];
- Parameterizations for unintegrated parton densities are available at www.desy.de/~lublinm;
- E. Gotsman, M. Kozlov, E. Levin, U. Maor and E. Naftali : “Towards a new global QCD analysis: Solution to the non-linear equation at arbitrary impact parameter,” *Nucl. Phys. A* 742 (2004) 55 [arXiv:hep-ph/0401021];
- E. Gotsman, E. Levin, U. Maor and E. Naftali : “A modified Balitsky - Kovchegov equation”, (in preparation);

F_2 - HERA data :

Low Q^2 :

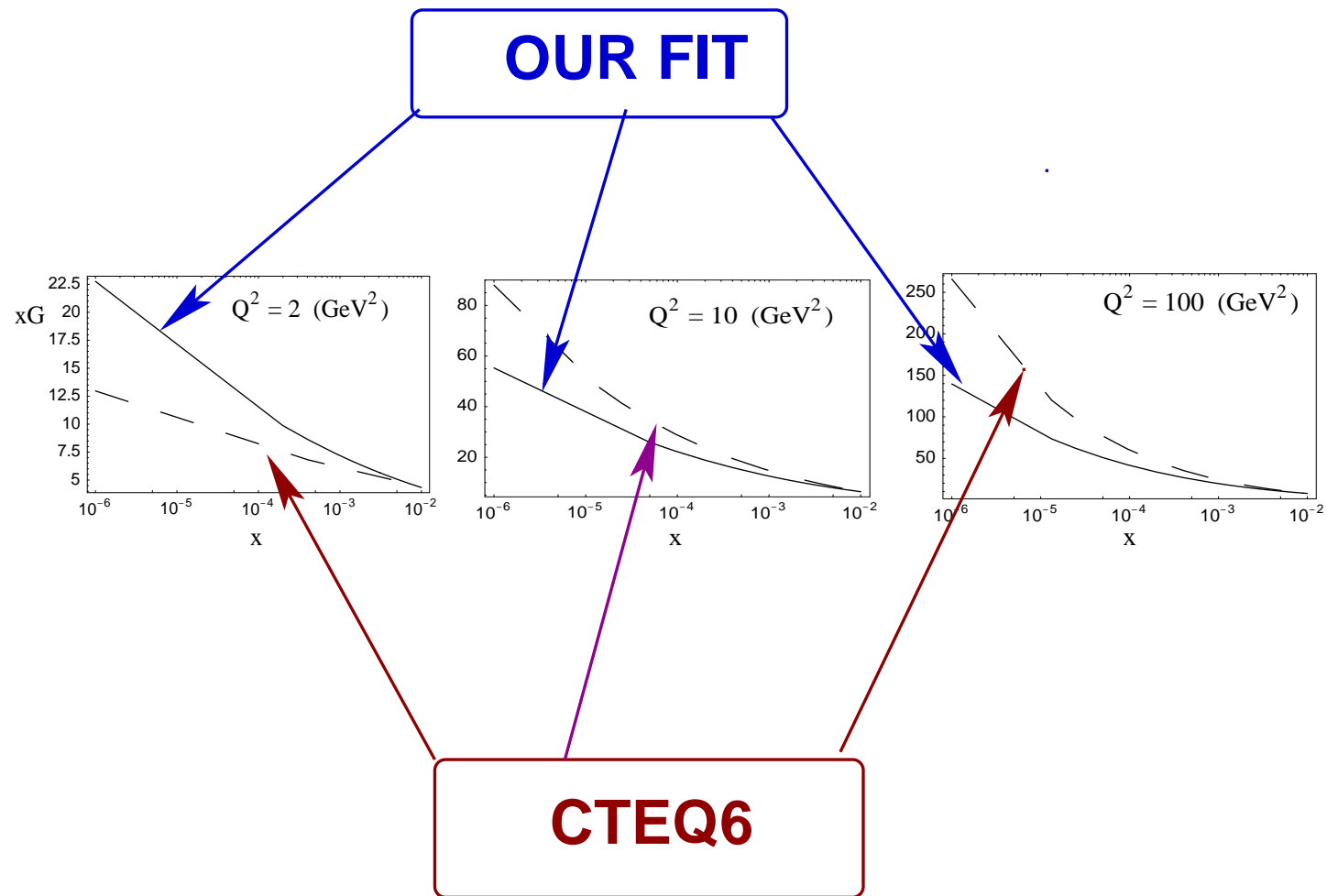


Large Q^2 :



LHC predictions :

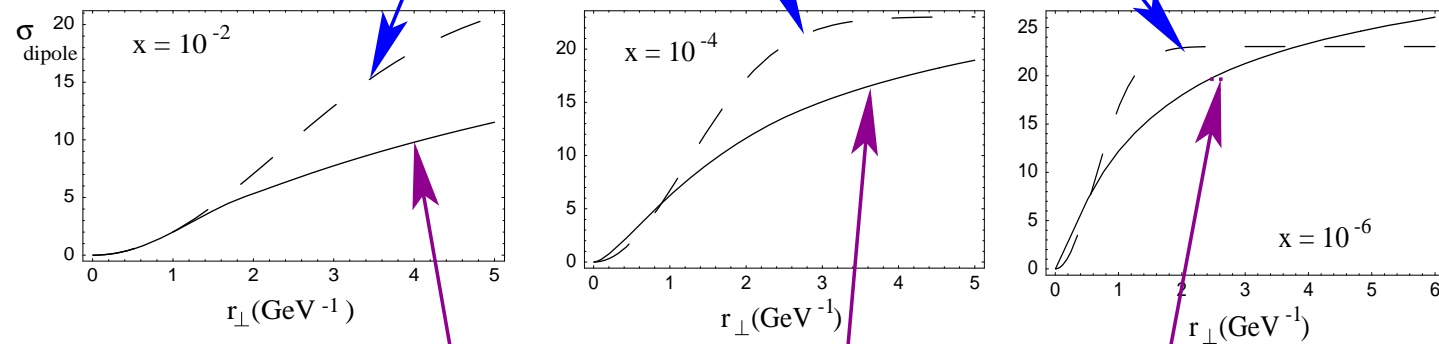
$$xG(x, Q^2):$$



LHC predictions :

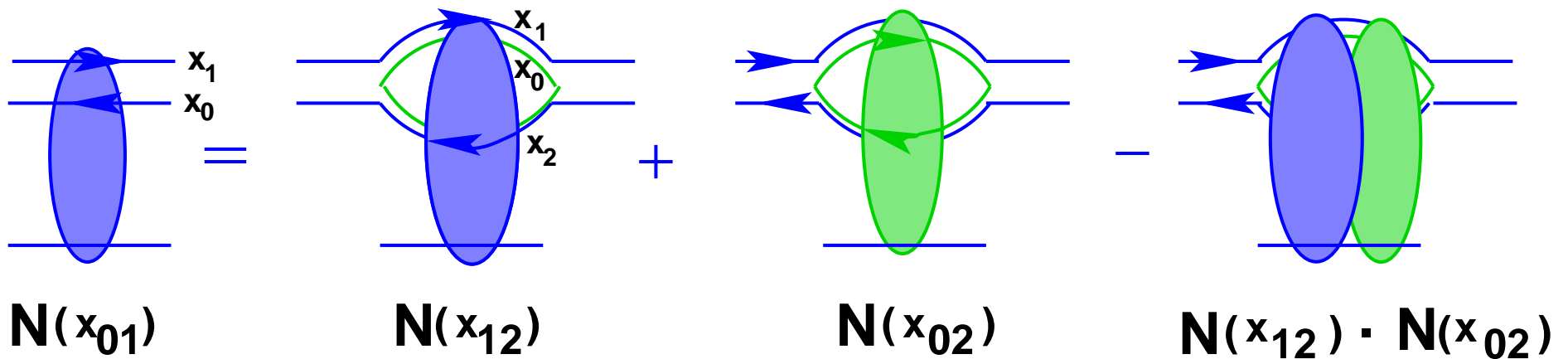
σ -dipole:

GBW MODEL



OUR PREDICTIONS

B-K non-linear equation:



$$\begin{aligned}
 \frac{\partial N(y, \vec{x}_{01}, \vec{b})}{\partial y} = & \frac{C_F \alpha_S}{2\pi^2} \int d^2 x_2 \frac{x_{01}^2}{x_{02}^2 x_{12}^2} \left(2N(y, \vec{x}_{12}, \vec{b} - \frac{1}{2}\vec{x}_{02}) \right. \\
 & \left. - N(y, \vec{x}_{01}, \vec{b}) - N(y, \vec{x}_{12}, \vec{b} - \frac{1}{2}vecx_{02}) N(y, \vec{x}_{02}, \vec{b} - \frac{1}{2}\vec{x}_{12}) \right)
 \end{aligned}$$

Deficiencies of B-K equation:

- Correct only in LLA approximation of pQCD with BFKL kernel in LO;
- The mean field approximation to the JIMWLK equation;
- It is not correct in the saturation region;
- The region where we can neglect the non-linear corrections should be specified by conditions beyond the BK equation;

B-K equation versus NLO BFKL:

$$N_{\text{non-linear term}} \propto \alpha_S^4 s^{2\Delta_{BFKL}}; \quad N_{\text{linear term}} \propto \alpha_S^2 s^{\Delta_{BFKL}};$$

with $\Delta_{BFKL} = \alpha_S \chi_{LO BFKL} + \alpha_S^2 \chi_{NLO BFKL}$

Correct strategy (theory point of view):

For $1/\alpha_S > y = \ln s > 1$

$$N_{\text{linear term}}^{LO BFKL}$$

For $y = \ln s > (2/\alpha_S) \ln(1/\alpha_S)$

$$N_{\text{linear term}}^{LO BFKL} + N_{\text{n-l term}}$$

For $y = \ln s > (1/\alpha_S^2)$

$$N_{\text{linear term}}^{NLO BFKL} + N_{\text{n-l term}}$$

NLO BFKL and saturation scale $Q_s(Y)$:

- The NLO BFKL kernel is known (Lipatov and Fadin , Camisi and Ciafaloni 1998);
- The needed re-summation has been performed (Salam, Ciafaloni et al. 1998 - 2000);
- Our procedure is based on the observation of the Durham group:

$$\bar{\alpha}_S \chi_{NLO}(\gamma) = \left(\frac{1 + \omega A_1(\omega)}{\gamma} - \frac{1}{\gamma} + \frac{1 + \omega A_1(\omega)}{1 - \gamma + \omega} - \frac{1}{1 - \gamma} \right) - \omega \chi_{LO}^{HT}(\gamma)$$

- The pole $\gamma = 0$ corresponds to the normal twist-2 DGLAP contribution with

$$Q > \dots k_{i,t} > k_{t,i+1} > \dots > Q_0$$

where Q_0 is the typical virtuality of the target;

- The pole at $\gamma = 1$ corresponds to inverse k_t ordering

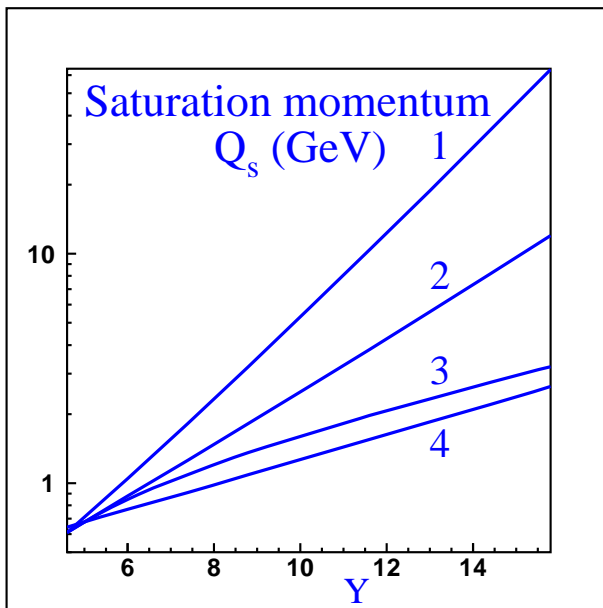
$$k_{i,t} < k_{t,i+1} < \dots < Q_0$$

- The other poles, at $\gamma = -1, -2, \dots$ ($\gamma = 2, 3, \dots$), are the higher twists contributions due to the gluon reggeization.

Our suggestion: (Ellis, Kunszt and Levin 1994)

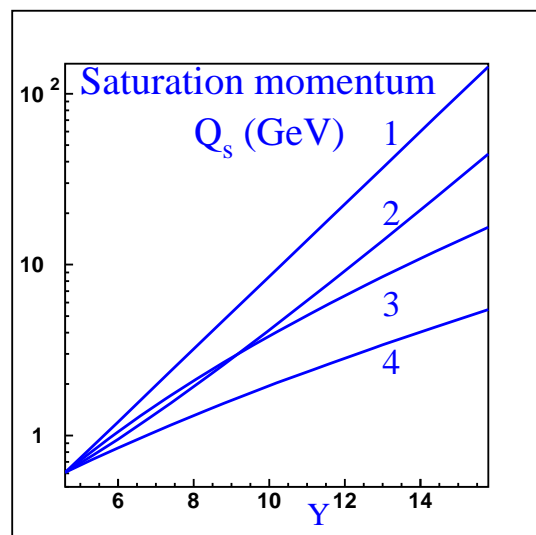
$$\bar{\alpha}_S \chi_{NLO}(\gamma) = -\omega \bar{\alpha}_S \chi_{LO}(\gamma) ;$$

$$\omega(\gamma) = \bar{\alpha}_S (1 - \omega) \bar{\alpha}_S \chi_{LO}(\gamma) ; \quad \gamma^{DGLAP} = \bar{\alpha}_S \left(\frac{1}{\omega} - 1 \right) ;$$



- 1 - LO BFKL
- 2 - Our kernel;
- 3 - NLO BFKL (Durham);
- 4 - BGW model;

Energy dependence of Q_s :



- 1** - High energy behaviour (fixed α_S);
- 2** - Low energy corrections (fixed α_S);
- 3** - High energy behaviour (running α_S);
- 4** - Low energy corrections (running α_S);

$$Q_s^2(Y) = Q_s^2(Y_0) \exp \left(\frac{\omega(\gamma_{cr})}{1 - \gamma_{cr}} (Y - Y_0) - \frac{3}{2(1 - \gamma_{cr})} \ln(Y/Y_0) - \frac{3}{(1 - \gamma_{cr})^2} \sqrt{\frac{2\pi}{\omega''(\gamma_{cr})}} \left(\frac{1}{\sqrt{Y}} - \frac{1}{\sqrt{Y_0}} \right) \right)$$

Modified B-K equation:

$$\bullet \frac{\partial N(r, Y; b)}{\partial Y} = \frac{C_F \alpha_S}{\pi^2} \int \frac{d^2 r' r^2}{(\vec{r} - \vec{r}')^2 r'^2}$$

$$\left(2N\left(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')\right) - N\left(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')\right) N\left(\vec{r} - \vec{r}', Y; b - \frac{1}{2}\vec{r}'\right) \right)_{B-K \text{ term}}$$

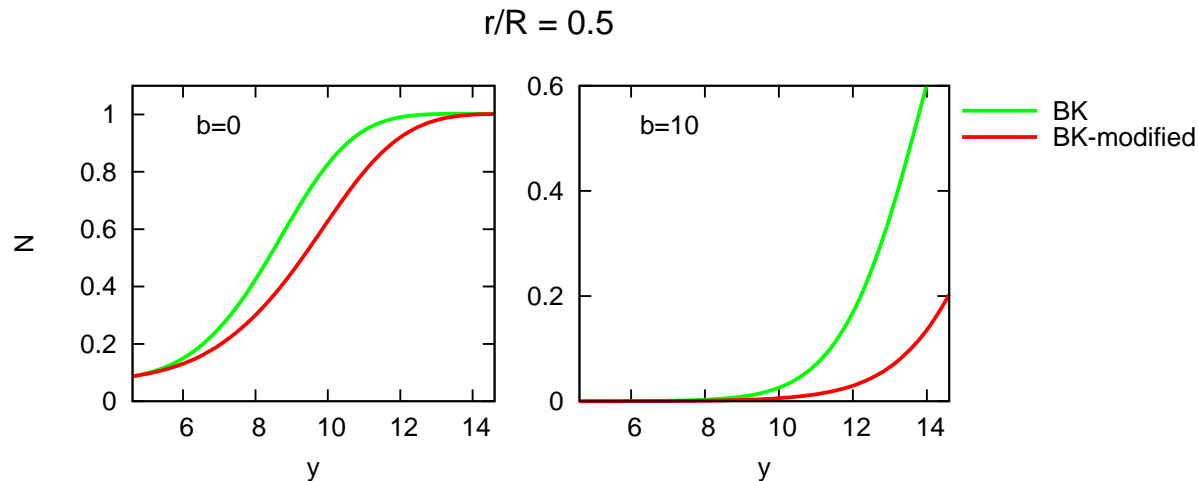
$$- \frac{\partial}{\partial Y} \left(2N\left(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')\right) - N\left(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')\right) N\left(\vec{r} - \vec{r}', Y; b - \frac{1}{2}\vec{r}'\right) \right)_{\text{new}}$$

$\bar{\alpha}_S \omega \chi_{LO}(\gamma)$ has the following form in Y, r representation

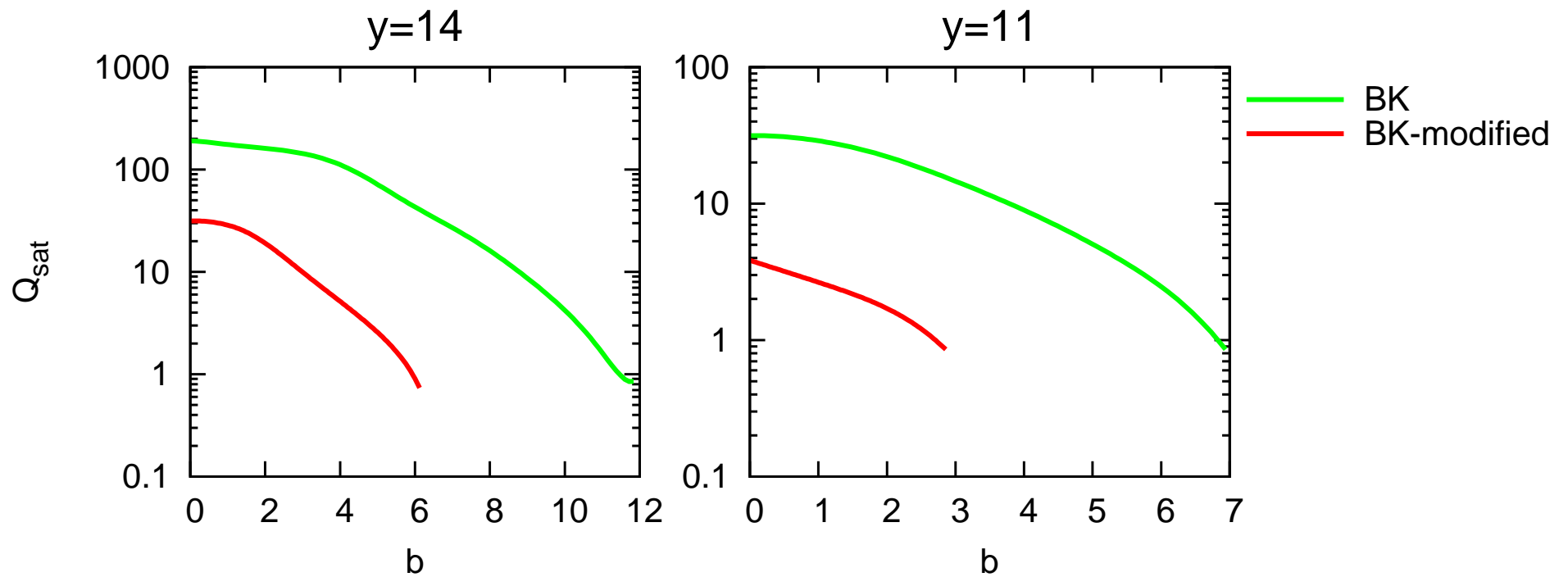
$$\bar{\alpha}_S \omega \chi_{LO}(\gamma) \rightarrow \bar{\alpha}_S \int K_{LO}(Y, r') d^2 r' \frac{\partial N(Y, r')}{\partial Y}$$

Modified B-K equation that we have solved:

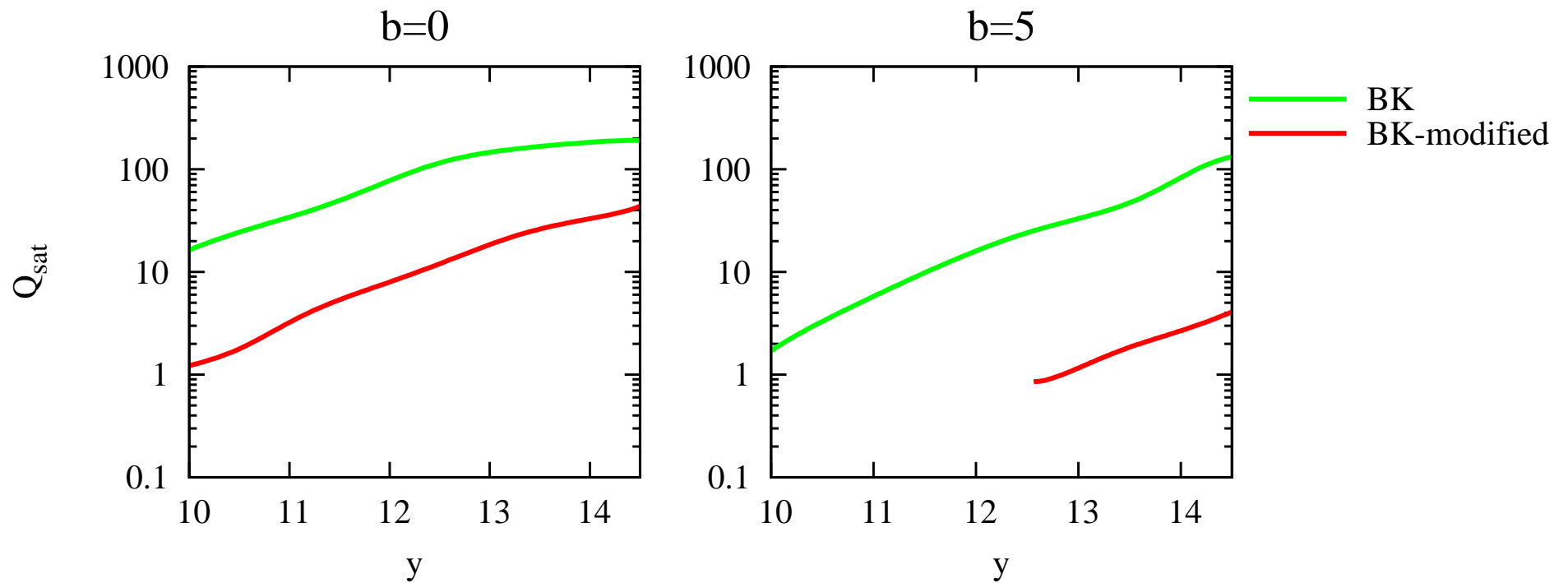
$$\bullet \frac{\partial N(r, Y; b)}{\partial Y} = \frac{C_F \alpha_S}{\pi^2} \left\{ - 2r^2 \int_r^R \frac{d^2 r'}{r'^4} \frac{\partial N(r', Y; b)}{\partial Y} \Big|_{DGLAP} + \int \frac{d^2 r' r^2}{(\vec{r} - \vec{r}')^2 r'^2} \Theta(R - r') \Theta(R - |\vec{r} - \vec{r}'|) \times \left[2N(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')) - N(r', Y; \vec{b} - \frac{1}{2}(\vec{r} - \vec{r}')) N(\vec{r} - \vec{r}', Y; b - \frac{1}{2}r') \right] \right\} .$$



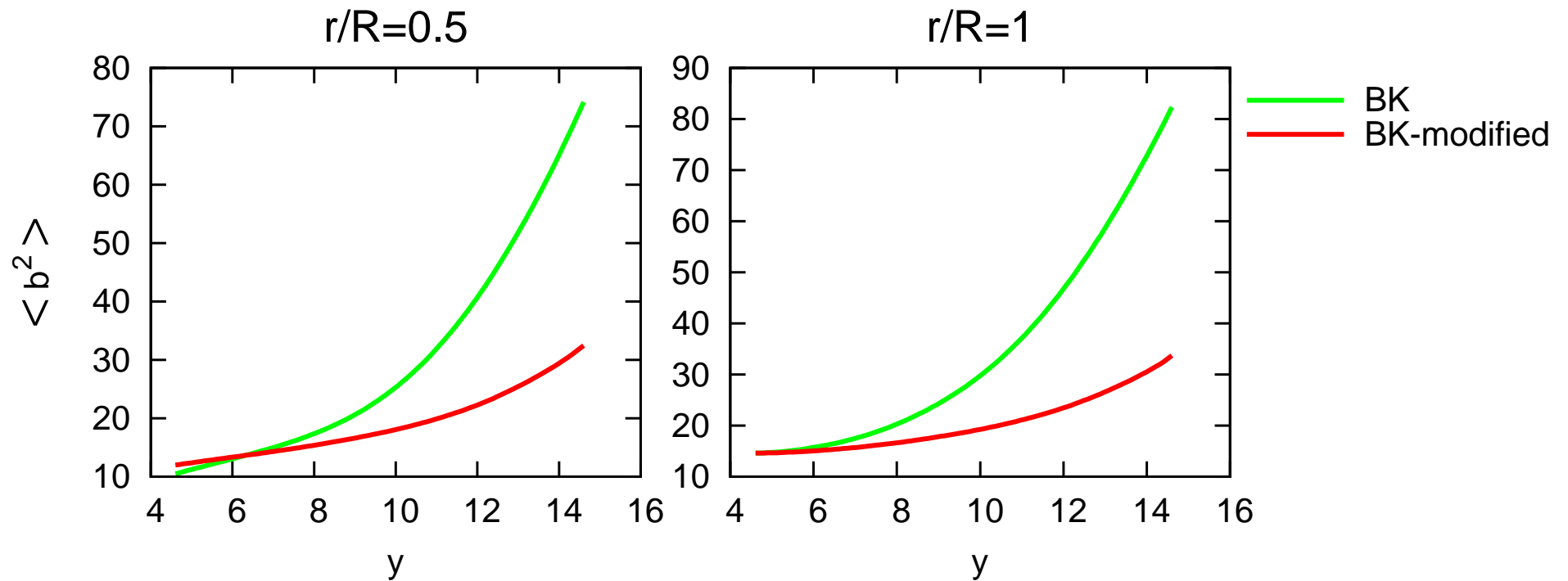
b -dependence of Q_s :



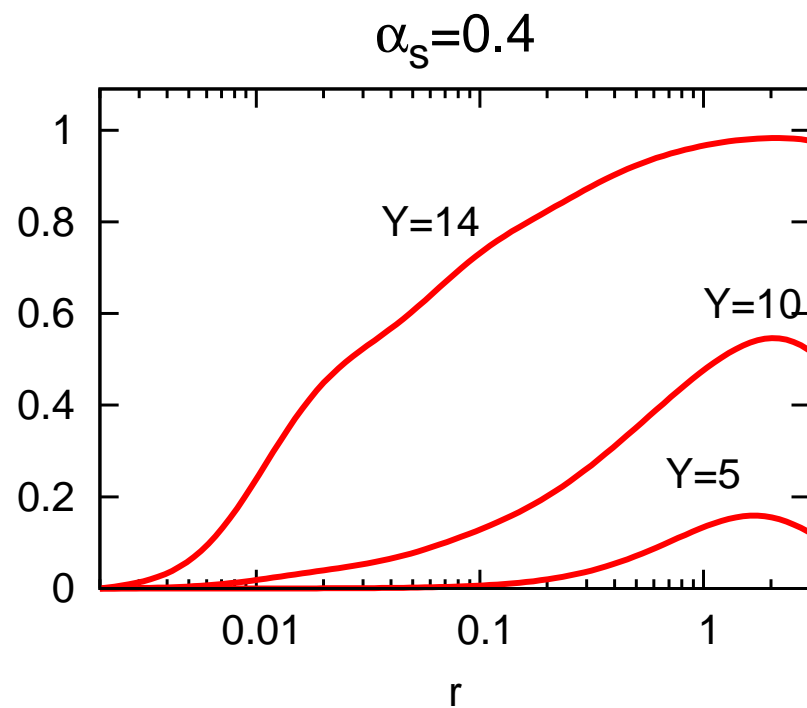
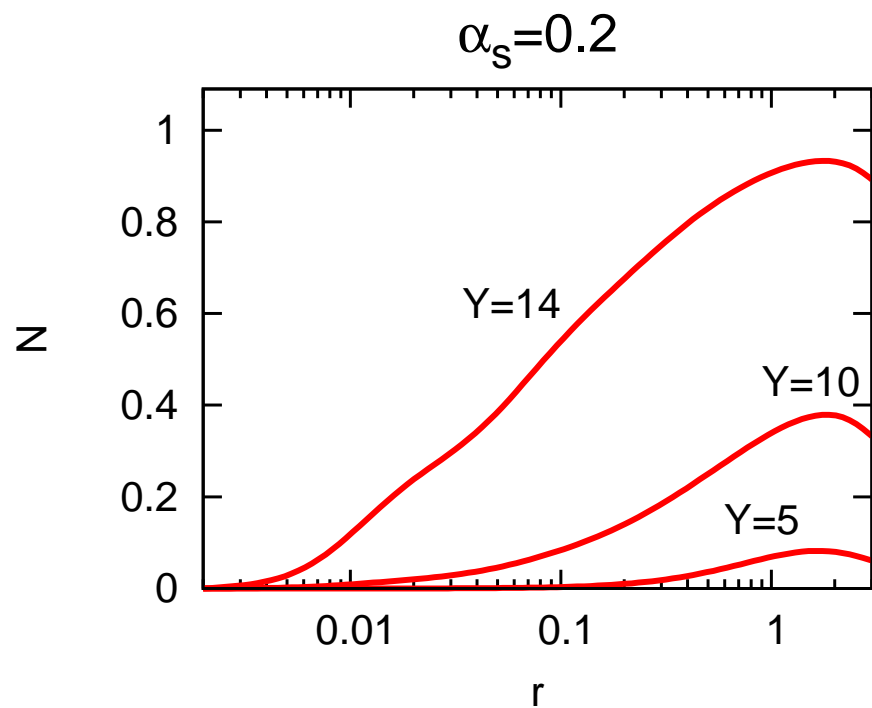
Y -dependence of Q_s :



Energy behaviour of $\langle b^2 \rangle$:



α_S dependence of the dipole amplitude :



Conclusions:

- The influence of the preasymptotic corrections, related to the full anomalous dimension of the DGLAP equation, **is rather large**;
- These corrections **slow down** the energy behaviour **as well as** the value of Q_s ;
- The BK equation without any modification **is not able to provide** reliable predictions for the LHC energies;
- The modified BK equation **has a chance** to do this job;

Problems to be solved:

- Solution to the full modified equation ?
- Running α_S ?
- Global fit of the experimental data?
- Reliable predictions for unintegrated parton densities for LHC energies?