

# HARD DIFFRACTION FROM PARTON RESCATTERING IN QCD

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- Introduction
- Soft Colour Interaction model  $\longrightarrow$  soft gaps in  $ep$  and  $p\bar{p}$
- Parton rescattering in QCD  $\longrightarrow$  rapidity gaps !  
Brodsky, Enberg, Hoyer, GI, hep-ph/0409119
- Summary & Outlook

# Diffractive hard scattering

**Idea:**

Ingelman-Schlein 1984

- hard scale probes parton level
- $\mathbb{P}$  flux
- $\mathbb{P}$  structure function

Predicted:

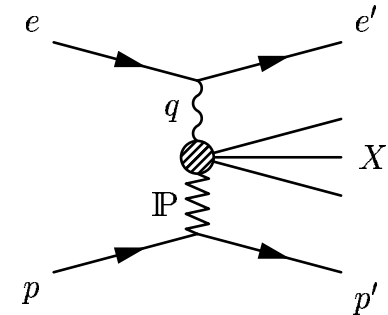
- jets etc. in single diffraction
- diffractive deep inelastic scattering

**Discovery:** UA8 at  $S_{p\bar{p}S}$  1988

- jets in single diffraction  $\simeq$  model
- hard gluons  $xg(x) \sim x(1-x)$  in  $\mathbb{P}$
- 'superhard'  $\delta(1-x)$  component

**More exp's:** HERA  $ep$ , Tevatron  $p\bar{p}$

Diffractive DIS



$$\frac{d\sigma}{dx dQ^2 dx_{\mathbb{P}} dt} = \frac{2\pi\alpha^2}{xQ^4} (1 + (1-y)^2) F_2^{D(4)}$$

$$F_2^{D(4)}(x, Q^2, x_{\mathbb{P}}, t) =$$

$$\underbrace{f(x_{\mathbb{P}}, t)}_{\mathbb{P} \text{ flux}} \underbrace{F_2^{\mathbb{P}}(\beta, Q^2)}_{\mathbb{P} \text{ structure}}$$

Fits HERA rapidity gap data

## Pomeron problems:

- $\mathbb{P}$  model fitted to HERA data  
→ fails for Tevatron data  
 $\sigma(\text{hard diff})$  factor 6–100 too large  
→ need ‘damping’ at high energies,  
*e.g.*  $\mathbb{P}$  flux ‘renormalisation’
- $\mathbb{P}$  flux & structure not universal  
ill-defined for virtual  $\mathbb{P}$
- Factorisation broken in diffractive  $p\bar{p}$   
– coherent interactions
- Improper with  $\mathbb{P}$  ‘emitted’ from  $p$   
soft, long space-time-scale interaction  
→  $\mathbb{P}$ - $p$  cross-talk

## Soft Colour Interactions:

- hard pQCD left unchanged  
– not affect by soft interactions
- non-pQCD below  $Q_0^2 \sim 1 \text{ GeV}^2$
- $\alpha_s$  large  $\Rightarrow$  large interaction probability  
*e.g.* unity for hadronisation!
- no proper theory → models
- colour exchange modifies colour/string topology → different final state
- single model describing all final states  
– diffractive  $\leftrightarrow$  nondiffractive

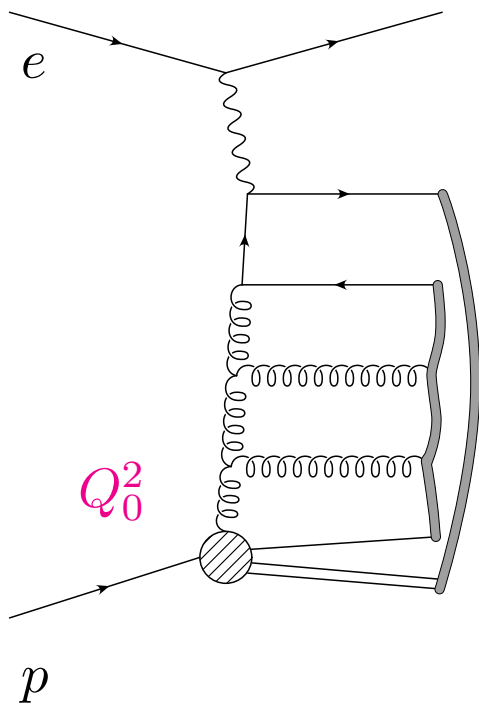
# Soft Colour Interaction model (SCI)

Soft interactions among partons & remnants ( $\leftrightarrow$  proton colour field) below  $Q_0^2 \sim 1 \text{ GeV}^2$

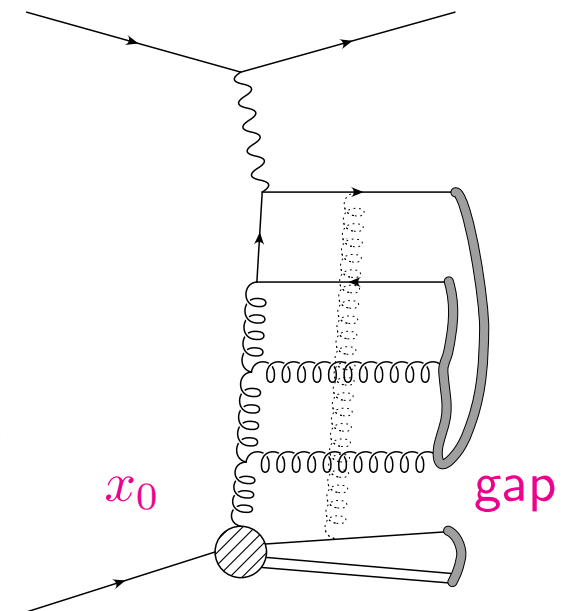
Add-on to Lund Monte Carlo's LEPTO ( $ep$ ) and PYTHIA ( $p\bar{p}$ )

ME + DGLAP PS  $> Q_0^2$   $\rightarrow$  SCI model  $\rightarrow$  String hadronisation  $\sim \Lambda$   
 colour ordered parton state  $\rightarrow$  rearranged colour order  $\rightarrow$  modified final state

Single model describing all final states: diffractive  $\leftrightarrow$  nondiffractive



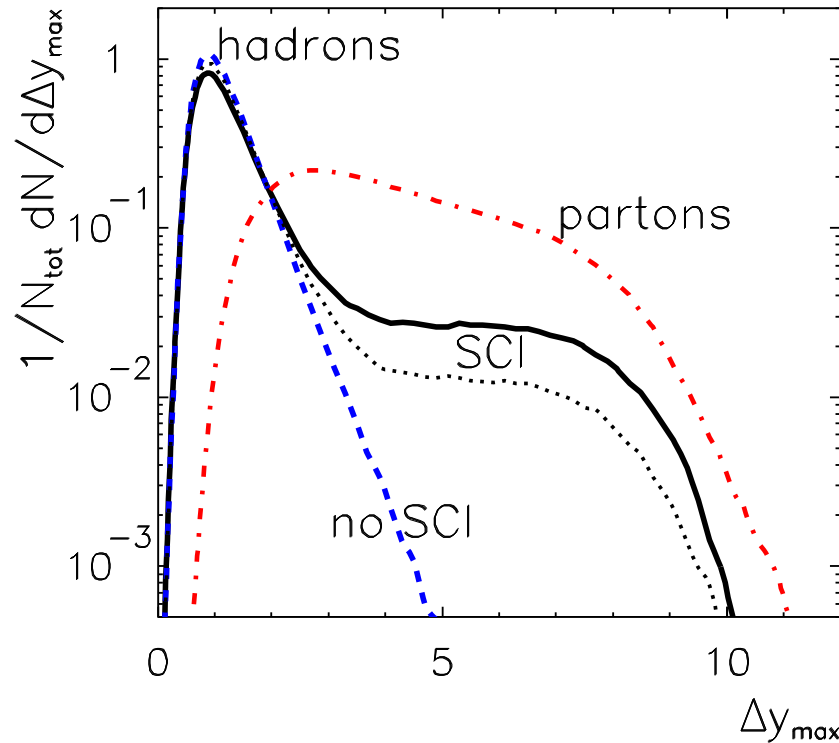
Single parameter  $P = \text{const} \approx 0.5$   
 gives probability for soft ( $p \simeq 0$ )  
 colour-anticolour (gluon) exchange  
 between parton pairs,  
 determined from HERA rap-gap data



Proton remnant with  $(1 - x_0)$  important for large gaps

# Gap-size is infrared sensitive observable !

Size  $\Delta y_{max}$  of largest gap in DIS events



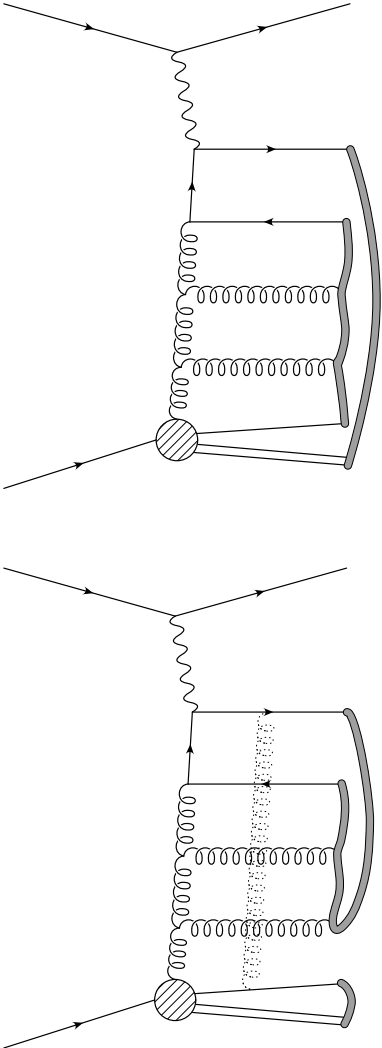
SCI  $\Rightarrow$  plateau in  $\Delta y_{max}$   
characteristic for diffraction

Small parameter sensitivity

—  $P = 0.5$

...  $P = 0.1$

Large gaps at parton level  
normally string across  $\rightarrow$  hadrons fill up  
SCI  $\rightarrow$  new string topologies, some with gaps



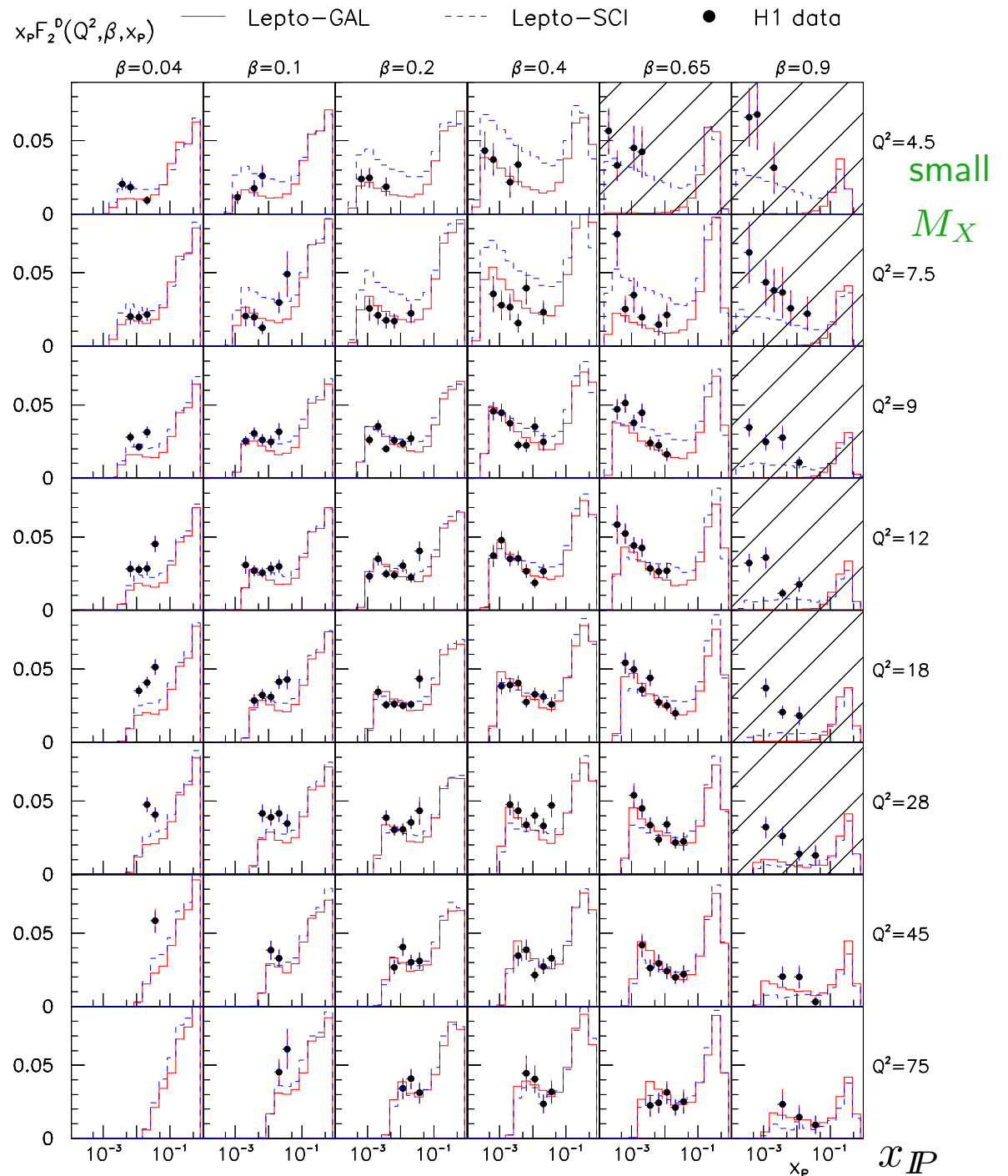
# Diffractive structure function in DIS

$$x_{\mathcal{I}P} F_2^{D(3)}(Q^2, \beta, x_{\mathcal{I}P})$$

SCI model describes  
main features of HERA  
rapidity gap data

Not bad for a  
one-parameter model !

GAL is an alternative  
formulation of the model,  
based on string interactions



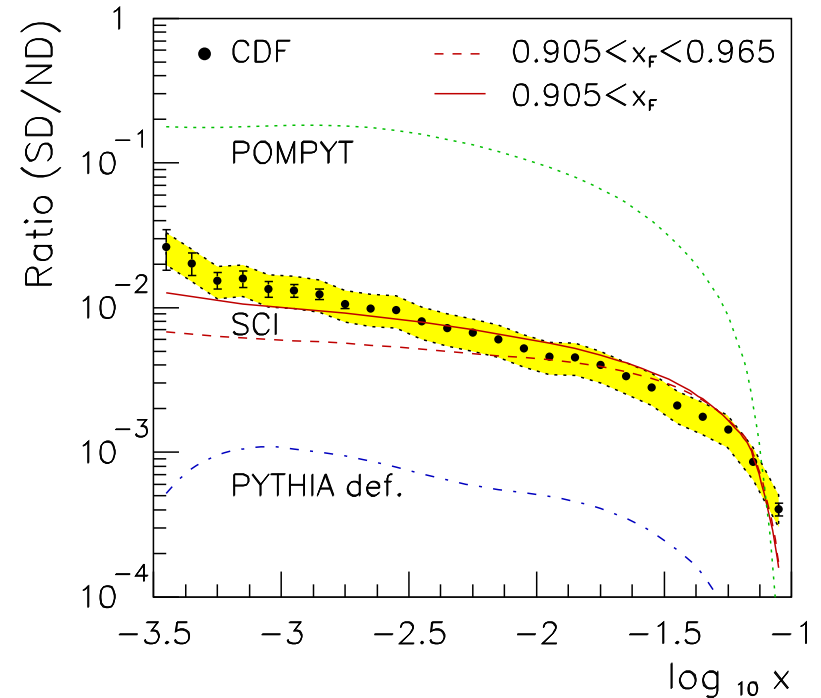
# Single diffractive jets, $W$ , $Z$ , $b\bar{b}$ , $J/\psi$ at the Tevatron

$$R_{\text{hard}} = \frac{1}{\sigma_{\text{hard}}^{\text{tot}}} \int_{x_{F\text{min}}}^1 dx_F \frac{d\sigma_{\text{hard}}}{dx_F}$$

$R_{\text{hard}}[\%]$	Exp.	observed	SCI
dijets	CDF	$0.75 \pm 0.10$	0.7
W	CDF	$1.15 \pm 0.55$	1.2
W	DØ	$1.08^{+0.21}_{-0.19}$	1.2
$b\bar{b}$	CDF	$0.62 \pm 0.25$	0.7
Z	DØ	$1.44^{+0.62}_{-0.54}$	1.0
$J/\psi$	CDF	$1.45 \pm 0.25$	1.4

↑
predictions

$R_{\text{dijets}}$  vs  $x$  of parton in  $\bar{p}$

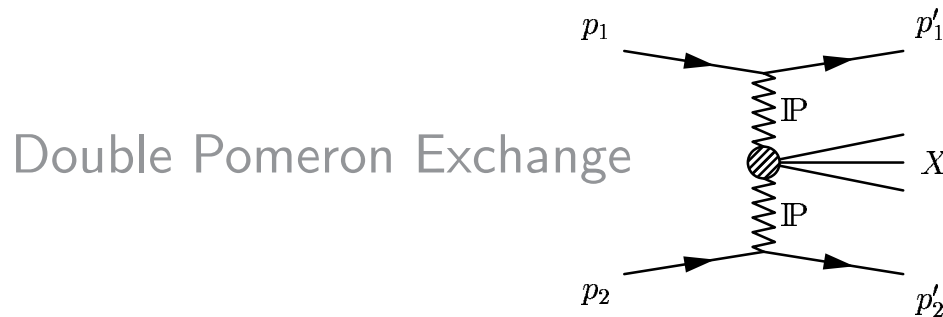


SCI  $\rightarrow$  gap &  $c\bar{c}$  colour octet  $\rightarrow$  singlet  $\rightarrow J/\psi$

SCI model OK, Pomeron model too high, default PYTHIA too low  
 SCI also correctly describes two-gap events (Double Pomeron Exchange)

$\Rightarrow$  Basis for predictions of diffractive Higgs production at Tevatron & LHC  
 Idea: easier to reconstruct Higgs in gap-event with less hadronic activity

# DPE: “Double leading Proton Events”

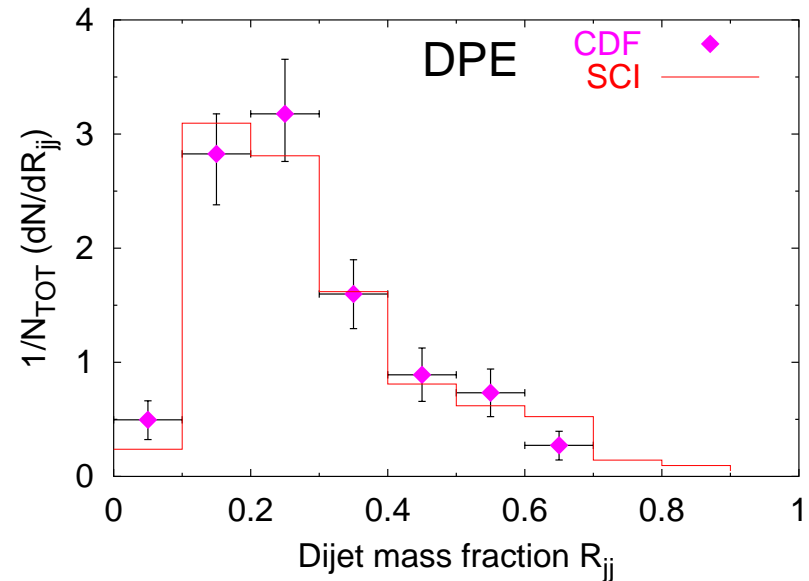


Tevatron  $p\bar{p} \rightarrow$  dijets in DPE  
 defined by leading  $\bar{p}$  and gap on  $p$  side

	CDF	SCI
$\tilde{R}_{SD}^{DPE}$ [%]	$0.80 \pm 0.26$	0.54
$\sigma^{DPE}$ [nb]	$43.6 \pm 22.0$	5–25

↑ ↑  
 proton gap

Strong dependence on remnant treatment  
 Good ratios and cross sections



Measured dijet mass fraction  
 $M_{jj}/\sqrt{\hat{s}}$  well reproduced



# New theoretical basis for SCI model

Brodsky, Enberg, Hoyer, GI, hep-ph/0409119

## Parton rescattering in QCD:

Leading twist gluon exchange between fast outgoing partons and target 'spectators'

Instantaneous 'Coulomb' gluons (light-front/Breit frame)

→ soft rescattering on 'frozen' target

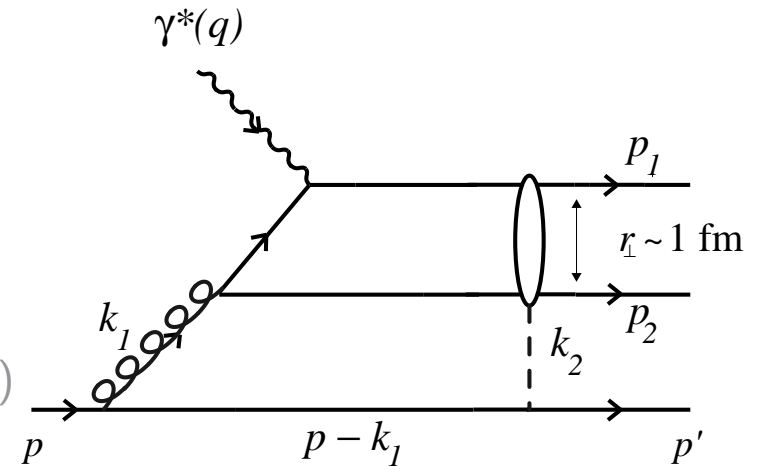
Gluon attached *after* photon vertex ⇒ no pre-existing pomeron in proton

Shadowing and diffraction are rescattering effects

Colour exchange → modified colour (string) field topology

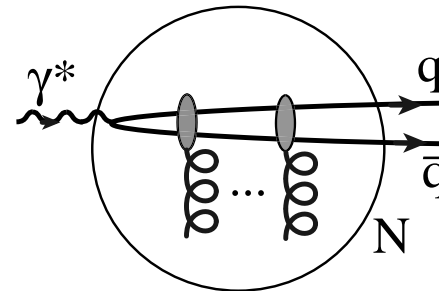
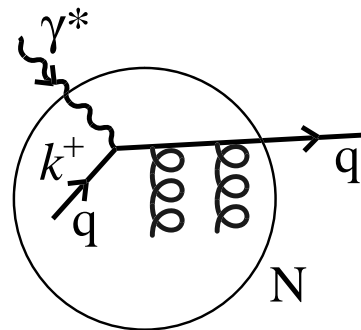
→ modified hadronic final state → gaps may arise

- gluon  $k_2$  quickly after  $k_1$  ⇒ can screen colour giving singlet exchange
- even/odd # gluons → pomeron/odderon exchange
- path-ordered exponential of gluon field (Wilson line)
- on-shell intermediate state ⇒ imaginary amplitude ↔ diffraction  
 $k_2$  is 'soft' → small kinematic effect
- $F_2^D \sim g_p(x_{\mathbb{P}}) q_g(\beta)$ , i.e. gluon distr. ↔  $\mathbb{P}$  and  $g \rightarrow q\bar{q}$  gives  $(\beta^2 + (1 - \beta)^2)$



# DIS on the 'Light Front'

Light Front (LF, or Light Cone LC) formalism – invariant under boosts along  $z$   
 → proton rest frame: interpretation depends on  $q_z$  (LF time  $x^+ = t + z$ )



'parton model' frame:

$$q = (\nu, 0, 0, -\sqrt{\nu^2 + Q^2})$$

$$q^+ \simeq -m_p x_B, \quad q^- \simeq 2\nu$$

Photon probes target at an  
 instant in LF time  $x^+$

'dipole model' frame:

$$q = (\nu, 0, 0, +\sqrt{\nu^2 + Q^2})$$

$$q^+ \simeq 2\nu, \quad q^- \simeq -m_p x_B$$

Dipole scatters on target  
 within finite LF time  $x^+$

The gluons are rescatterings within coherence (loffe) length  $x^- \sim 1/m_p x_B$

## QCD factorisation and rescattering

QCD factorization theorem  $\Rightarrow$  separation of hard and soft

Quark parton density function (PDF) is given by

$$f_{q/N} \sim \int dx^- e^{-ix_B p^+ x^- / 2} \langle N(p) | \bar{\psi}(x^-) \gamma^+ W[x^-; 0] \psi(0) | N(p) \rangle_{x^+=0}$$

where  $W[x^-; 0] = \text{P exp} \left[ ig \int_0^{x^-} dw^- A_a^+(0, w^-, 0_\perp) t_a \right]$  is the **Wilson line**

Expanding the exponential we have

$$W[x^-; 0] \sim 1 + g \int \frac{dk_1^+}{2\pi} \frac{\tilde{A}^+(k_1^+)}{k_1^+ - i\varepsilon} + g^2 \int \frac{dk_1^+ dk_2^+}{(2\pi)^2} \frac{\tilde{A}^+(k_1^+) \tilde{A}^+(k_2^+)}{(k_1^+ + k_2^+ - i\varepsilon)(k_2^+ - i\varepsilon)} + \dots$$

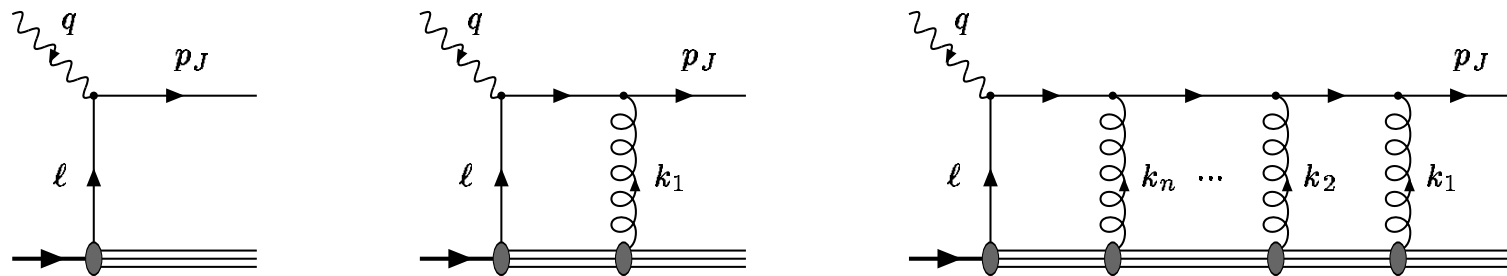
corresponding to the diagrams

for **rescattering of the struck quark via longitudinal ( $A^+$ ) instantaneous (in  $x^+$ ) gluon exchange**. No  $A^\perp$  within Ioffe coherence length  $x^- \sim 1/m_p x_B$

Brodsky, Hoyer, Marchal, Peigné, Sannino, Phys. Rev. D65 (2002) 114025

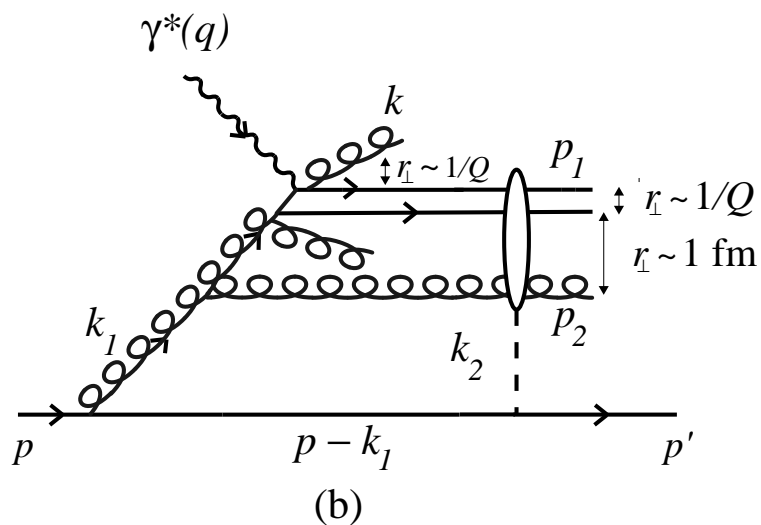
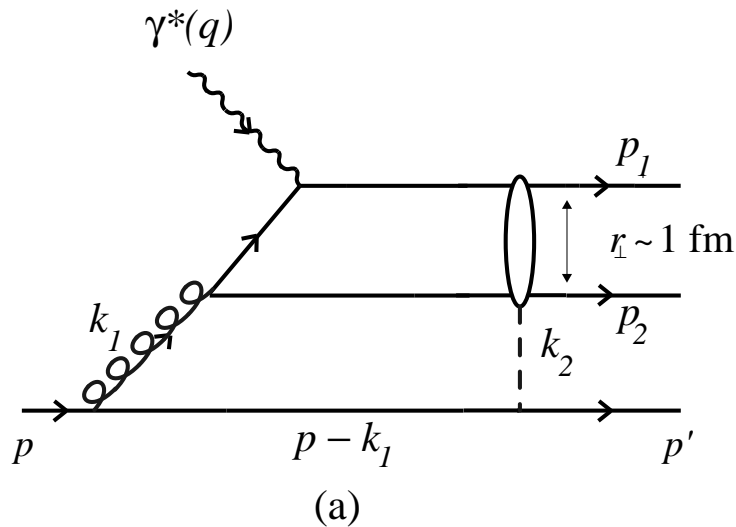
# QCD rescattering and factorisation in DIS

The Wilson line means that DIS looks something like this:



- Rescattering compatible with factorization theorems *by construction*  
Wilson line part of definition of PDF  $\Rightarrow$  rescattering part of PDF
- Experimentally measured PDF's *include* rescattering
- Diffractive PDF's included in inclusive PDF's
- Rescattering can give on-shell intermediate states and *imaginary amplitudes*  $\Rightarrow$  diffraction

## Mechanism for diffraction in DIS



The 'pomeron remnant' is soft  
— first splitting is in the sea

1. A soft gluon with  $k_1^+ / p^+ \sim x_{\mathbb{P}}$  splits into large  $q\bar{q}$  pair
2. The  $\gamma^*$  scatters and a large size  $q - \bar{q}$  'dipole' is formed
3. **Instantaneous gluon exchanges** may make dipole color singlet

- At large  $M_X$ ,  $q\bar{q} - g$  dipole dominates
  - soft gluon splits into  $gg$  pair
  - compact  $q\bar{q}$  pair not resolved

- **Higher order emissions do not destroy the rapidity gap!**

$$k_\perp \gg p_{2\perp} \text{ and } k^+ \lesssim p_2^+$$

$\Rightarrow$  rapidity of  $k$  larger than that of  $p_2$

**Radiation is close to  $\gamma^*$  vertex and is not resolved by the rescattering**

# Consequences for diffractive DIS

- Same hard sub-process in diffractive and non-diffractive events
- Same  $Q^2$  dependence in diffractive and inclusive DIS
- Same energy ( $W$  or  $x_B$ ) dependence in diffractive and inclusive DIS

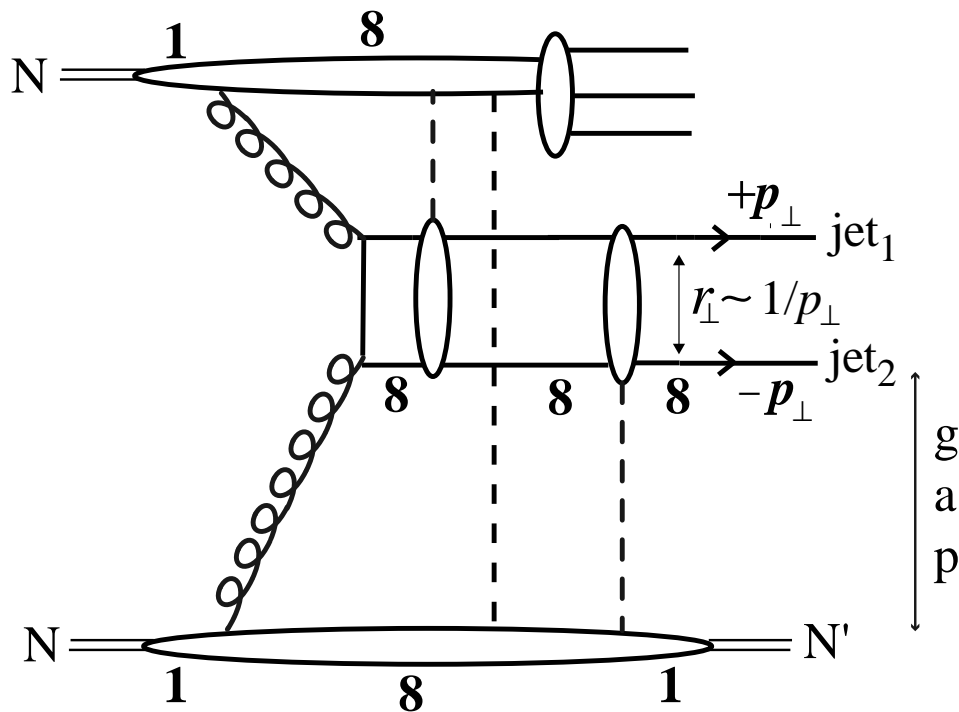
$\Rightarrow \sigma_{\text{diff}}/\sigma_{\text{tot}}$  independent of  $x_B$  and  $Q^2$ , as observed in data

- Amplitudes from rescattering dominantly imaginary, as expected for diffraction
- Rescattering gluons have small momenta  $\Rightarrow \beta$  dependence of diffractive PDF's from underlying (non-perturbative)  $g \rightarrow q\bar{q}$  and  $g \rightarrow gg$  processes

• Effective  $\mathbb{P}$  flux:  $f_{\mathbb{P}/p}(x_{\mathbb{P}}) \propto g(x_{\mathbb{P}}, Q_0^2)$

Effective  $\mathbb{P}$  structure function:  $f_{q/\mathbb{P}}(\beta, Q_0^2) \propto \beta^2 + (1 - \beta)^2$

# Hard diffraction in hadron-hadron collisions



- Diffractive factorization theorem does not hold
- Data shows  $\sim 1\%$  diffraction instead of  $\sim 10\%$  in DIS
- Both target and projectile coloured
  - different rescattering
  - lower probability for colour neutralization
- DPE possible
- SCI model reproduces data

## Related models

- **Colour evaporation model** (Halzen *et al.*)  
many exchanges  $\rightarrow$  random colour  
Diffractive DIS (Buchmüller *et al.*)  
 $\gamma g \rightarrow q\bar{q}$  octet  $\rightarrow$  singlet  $q\bar{q}$  (prob=1/9)  
which decouples from remnant  $\rightarrow$  gap event
- **Interactions with background colour field** (Nachtmann *et al.*)  
*e.g.* quark traversing colour field gives 'synchrotron' radiation  
of soft  $\gamma$ 's
- **Semiclassical model** (Buchmüller, Hebecker *et al.*)  
Diffractive DIS:  $\gamma \rightarrow q\bar{q}$  which traverses proton colour field  
interaction  $\sim$  Wilson loop averaging over colour field
- **String reconnection models** (Lund group)  
*e.g.* applied to  $e^+e^- \rightarrow WW \rightarrow q\bar{q} q\bar{q}$   
gives shift in  $W$  mass determination



# Summary

SCI model OK with data:

- gap events in DIS
- leading protons/neutrons in DIS
- diffractive jets,  $W$ ,  $Z$ ,  $b\bar{b}$ ,  $J/\psi$  at Tevatron
- high- $p_{\perp}$   $J/\psi$ ,  $\psi'$ ,  $\Upsilon$  at Tevatron
- $J/\psi$ ,  $\psi'$  in fixed target  $\pi A$  and  $pA$

Not bad for simple (one-parameter) model !

$\exists$  alternative/related models

Parton rescattering theory in QCD:

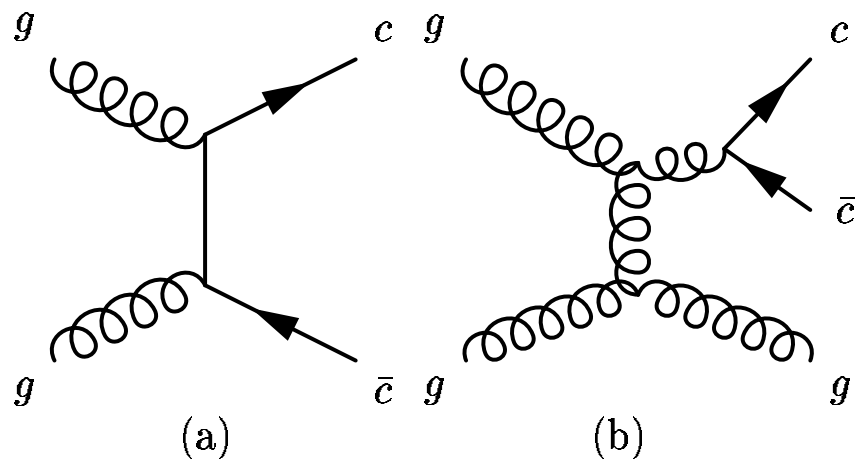
- Hard sub-process universal  
not affected by soft interaction
- PDF  $\sim$  LF wave fcn  $\otimes$  soft rescattering
- Rescattering via soft gluons,  
instantaneous in LF time  
May shield colour  $\rightarrow$  rapidity gap
- Soft rescattering does not resolve  
hard emissions
- Different colour environments  
(particles/collisions)  
 $\Rightarrow$  diffraction/gaps process dependent

# Outlook

- Non-perturbative QCD major unsolved problem
  - QCD rescattering theory
    - theory for diffraction/gaps
    - basis for successful SCI model and implies developments  
*e.g.* colour coherence from dynamical colour exchange probability
  - Parton interaction with colour background field
    - 'old' idea with new aspects
    - different theoretical approaches
- ⇒ Better understanding of non-perturbative QCD

Additional material . . .

# SCI $\rightarrow J/\psi$ in fixed target



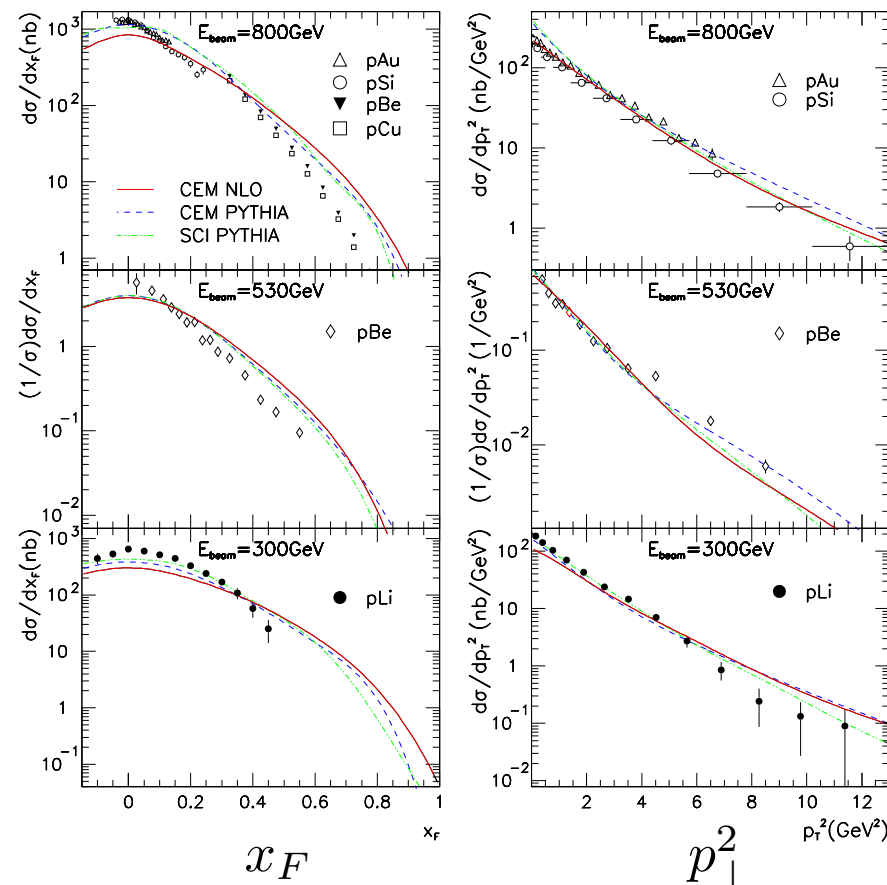
$pA$  @ 800, 530, 300 GeV

pQCD (NLO or LO+PS)  $\rightarrow c\bar{c}$  pair

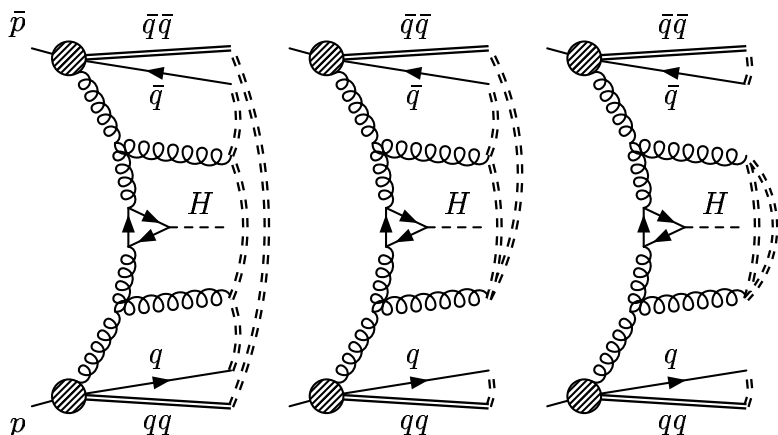
Colour octet  $c\bar{c}$  turned into singlet  $c\bar{c}$   
 $m_{c\bar{c}} < 2m_D$  mapped on charmonium states  
 with spin statistics (+ soft smearing)

Reproduces:

- high- $p_{\perp}$   $J/\psi$ ,  $\psi'$ ,  $\Upsilon$  at Tevatron
- $x_F$ ,  $p_{\perp}$  of  $J/\psi$ ,  $\psi'$  in  $\pi A$ ,  $pA$   $\implies$



# SCI $\rightarrow$ diffractive Higgs at Tevatron/LHC



Idea: diffractive events cleaner  
 $\rightarrow$  easier Higgs reconstruction  
 $\rightarrow$  Higgs discovery channel ?

$m_H = 115 \text{ GeV}$		Tevatron $\mathcal{L} = 20 \text{ fb}^{-1}$	LHC $\mathcal{L} = 30 \text{ fb}^{-1}$
	$\sigma [\text{fb}]$ Higgs-total	600	27000
SD	$\sigma [\text{fb}]$ leading-p	1.2	190
	$\sigma [\text{fb}]$ gap	2.4	27
	$\#$ H + leading-p	24	5700
	$\hookrightarrow \#$ H $\rightarrow \gamma\gamma$	0.05	13
DPE	$\sigma [\text{fb}]$ leading-p's	$1.2 \cdot 10^{-4}$	0.19
	$\sigma [\text{fb}]$ gaps	$2.4 \cdot 10^{-3}$	$2.7 \cdot 10^{-4}$
	$\#$ H + leading-p's	0.0024	6

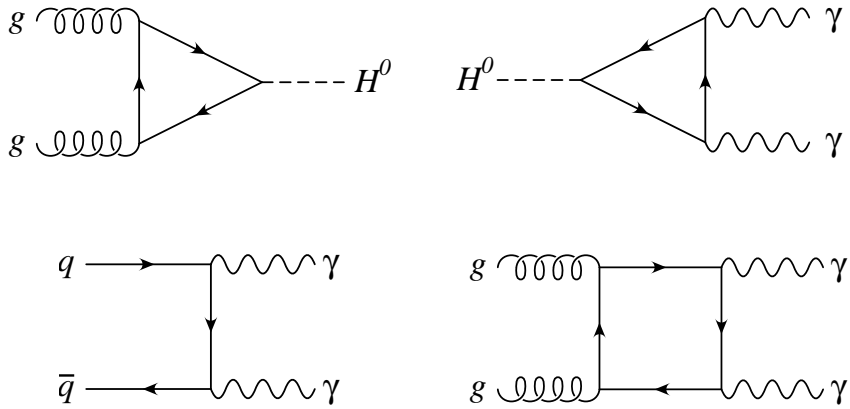
large  $m_H$  + large  $x_F$  proton  $\rightarrow$  kinematical conflict  $\rightarrow$  reduced cross-section

Tevatron: only few SD Higgs  $\Rightarrow H \rightarrow b\bar{b}$  background/reconstruction problems

LHC: Higgs observable in SD (also  $H \rightarrow \gamma\gamma$ ) and DPE

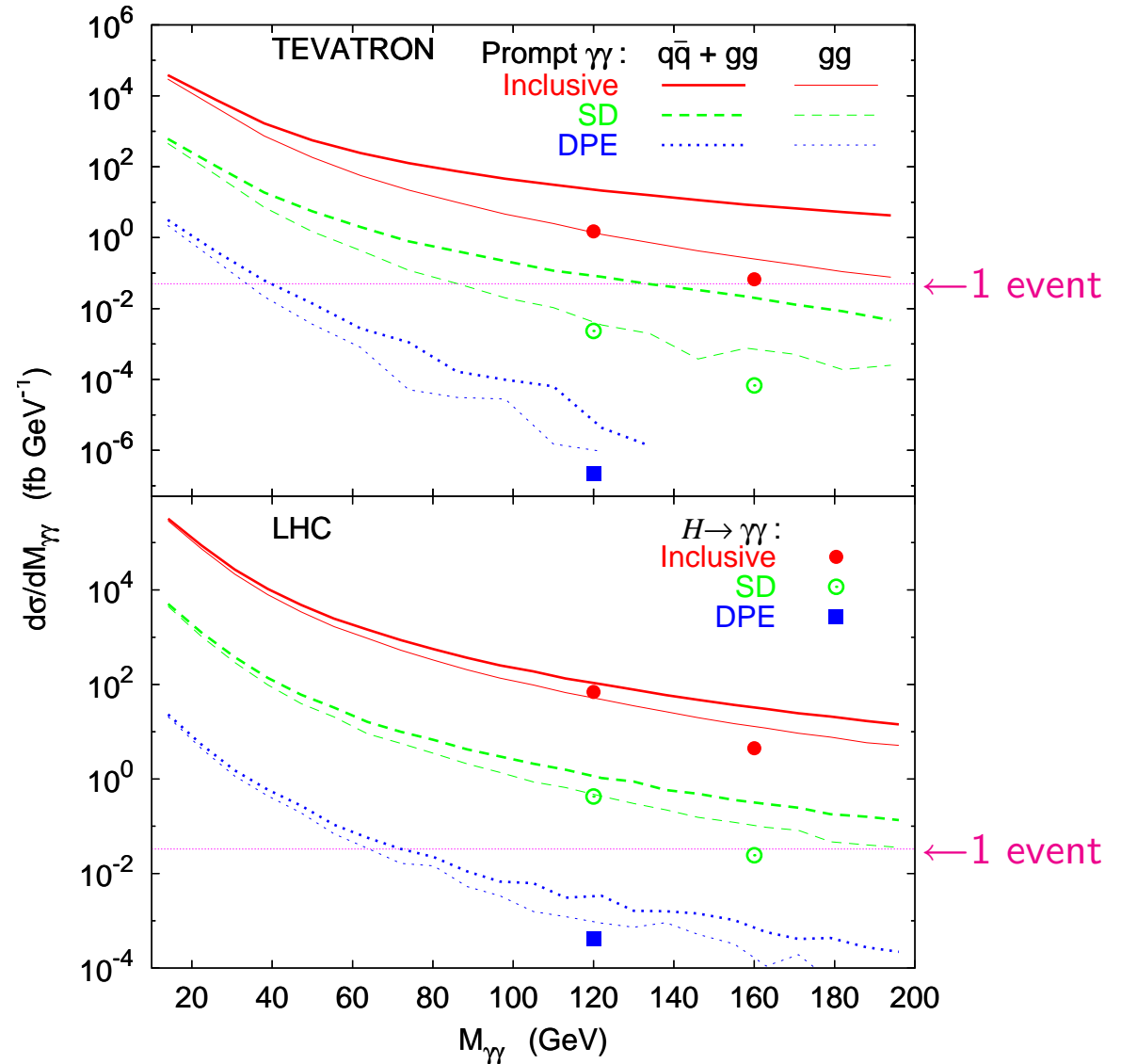
but not so clean since particles in 'gap' region  $\rightarrow$  'pots' better

# Diffractive $H \rightarrow \gamma\gamma$ and prompt $\gamma\gamma$



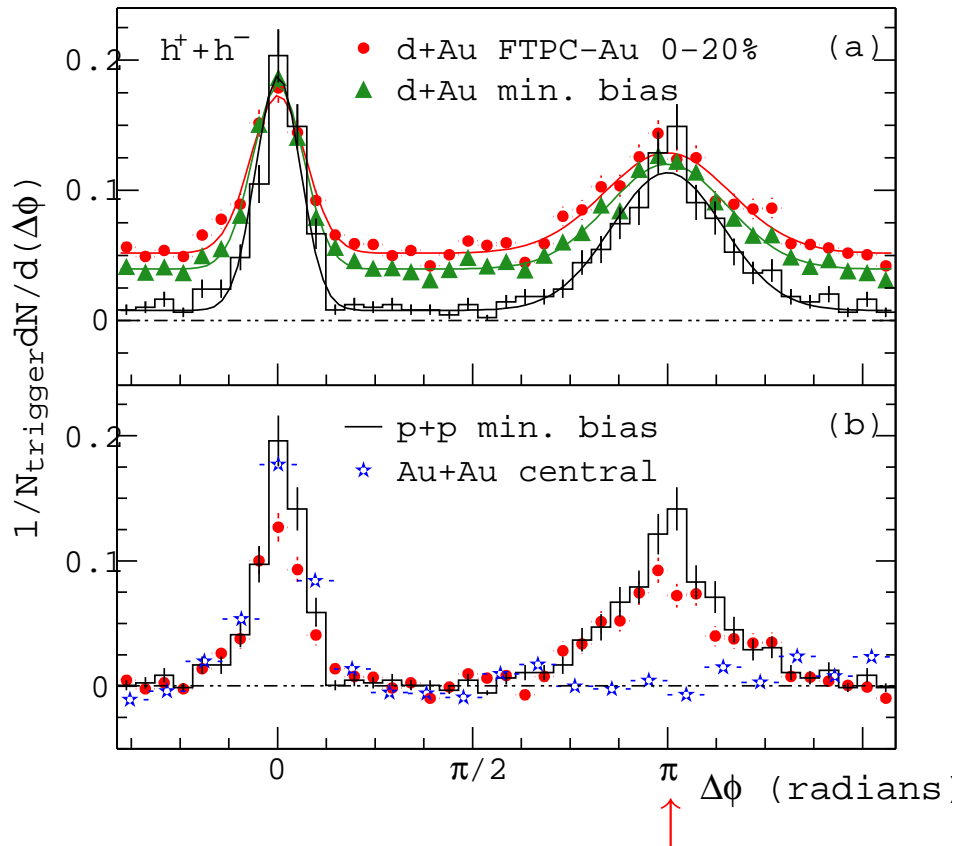
Prompt  $\gamma\gamma$  irreducible background  
 $> \sigma(H) \cdot \text{BR}(H \rightarrow \gamma\gamma)$

$\gamma\gamma$  tests  $gg$  fusion via quark loop  
 $\Rightarrow$  safer Higgs prediction



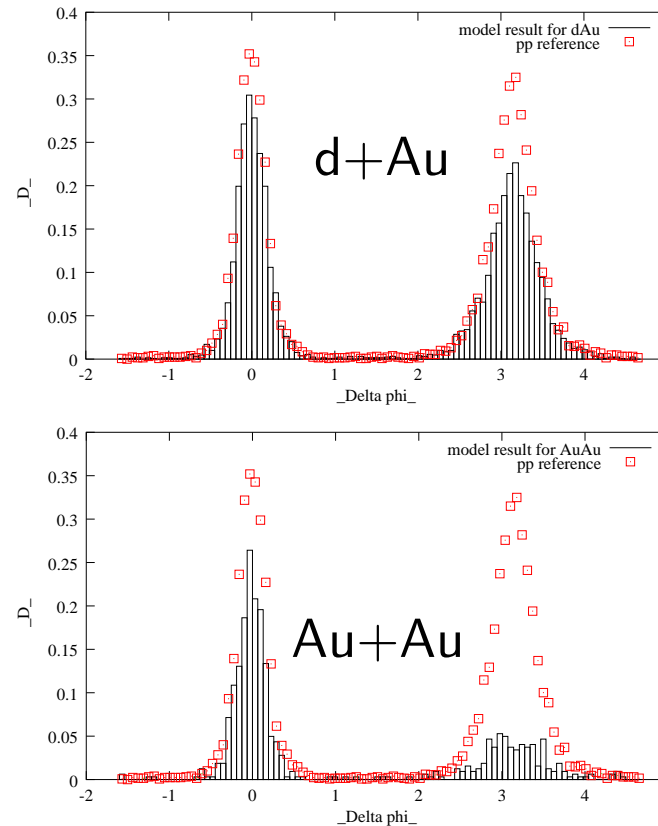
# Jet quenching in quark-gluon plasma

Hadron flow in azimuth  $\Delta\phi$   
relative to trigger particle/jet



Disappearance of back-to-back jet  
in central gold-gold collisions

MC model for soft colour interactions  
of hard-scattered parton with  
background quark-gluon plasma



Results from  
K. Zapp

Uppsala-  
Heidelberg  
collaboration

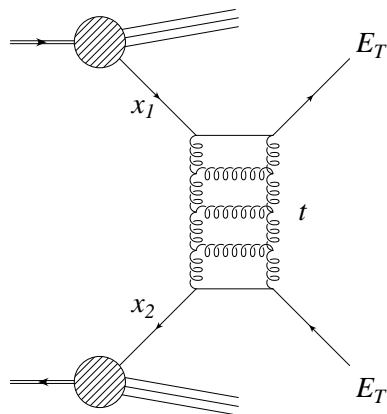
Reproduces effects qualitatively

# Soft and hard rapidity gaps

Soft/hard gap  $\Leftrightarrow$  small/large momentum transfer across gap  
 non-perturbative/perturbative QCD description

Hard gap between jets:

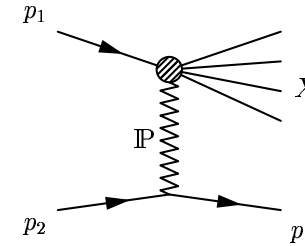
- observed at Tevatron
- understood/described by exchange of colour singlet BFKL gluon ladder with non-leading corrections



Enberg, GI, Motyka, PL B524 (2002) 273

Soft gap  $\leftrightarrow$  leading proton:

- 'soft-soft'  $\Rightarrow$  hadron basis  $\rightarrow$  Regge



- 'hard-soft'  $\leftrightarrow$  hard int'n – gap – proton  
 DIS/jets/W + gap  
 $\rightarrow$  pQCD partons + hadronisation  
 $\Rightarrow$  mixed basis: parton & hadron

Models:

- Regge  $\rightarrow$  pomeron structure works, but problematic
- Soft Colour Interactions . . .



# Gap between jets $\rightarrow$ hard QCD exchange

Rapidity gap **between** a pair of jets

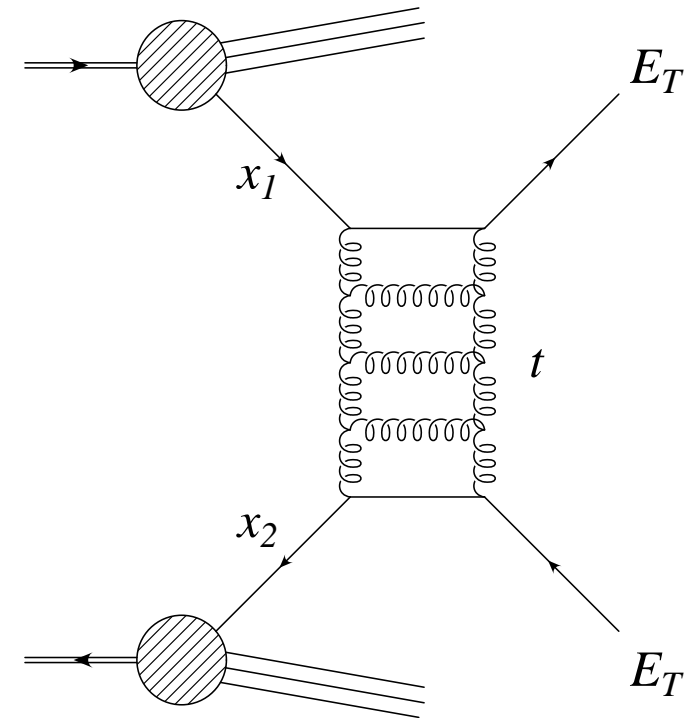
- $\rightarrow |t|$  across gap is large
- $\rightarrow$  different from “normal” diffraction

Elastic parton-parton scattering by hard colour singlet exchange (hard pomeron)

High energy limit  $s/|t| \gg 1 \Rightarrow$  amplitude dominated by terms  $\sim [\alpha_s \ln(s/|t|)]^n$

**BFKL** equation resums these terms, including

- virtual corrections
- reggeization of the exchanged gluons



Numerical solution of BFKL eqn. with non-leading corrections

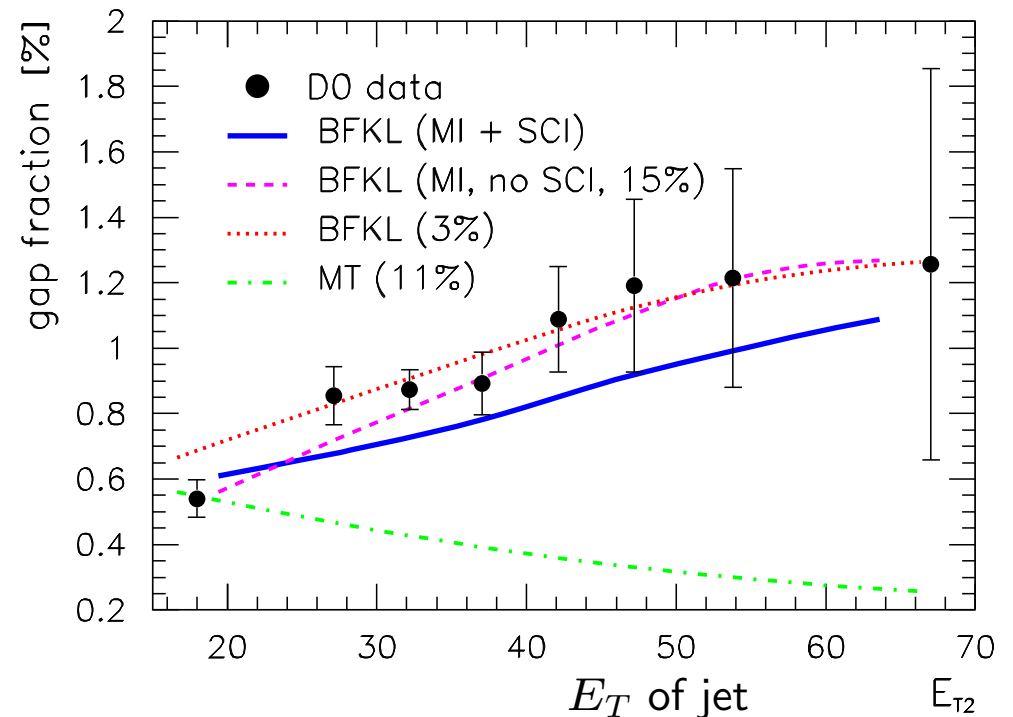
consistency constraint & running coupling  $\leftrightarrow$

# Comparison to data – gap survival

This subprocess implemented in PYTHIA → account for gap survival depending on HO parton emissions, multiple scatterings, hadronization

SCI destroys gaps by re-arranging strings across gaps

- Good agreement with DØ data
  - $E_T$  and  $\Delta\eta$  dependences OK
  - absolute normalisation OK
  - i.e.* correct gap survival probability
  - too large gap survival prob. without SCI
- Previous asymptotic result (Mueller-Tang) ⇒ wrong  $E_T$ -dependence



First clear evidence of BFKL dynamics ⇒ better understanding of data and QCD