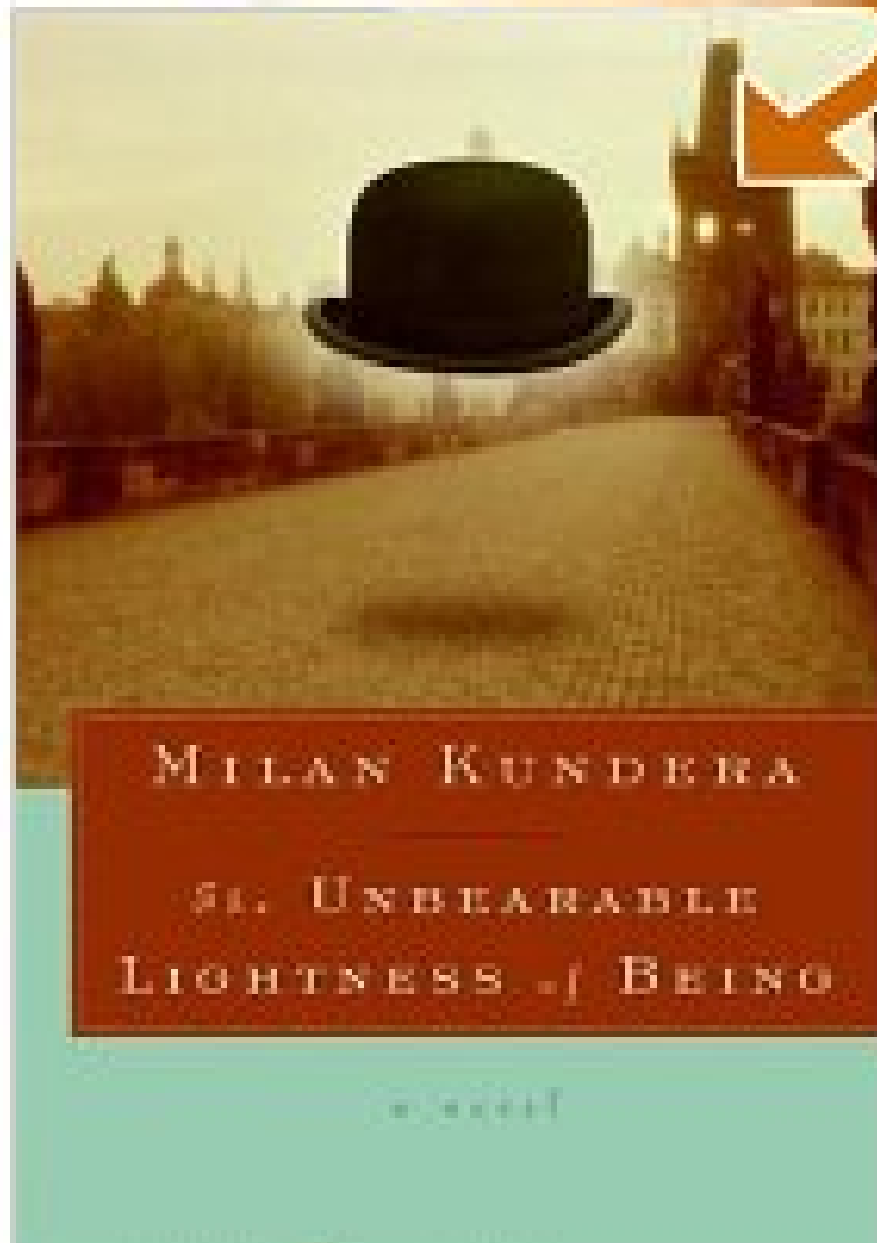


Physics at the proton and neutron drip lines

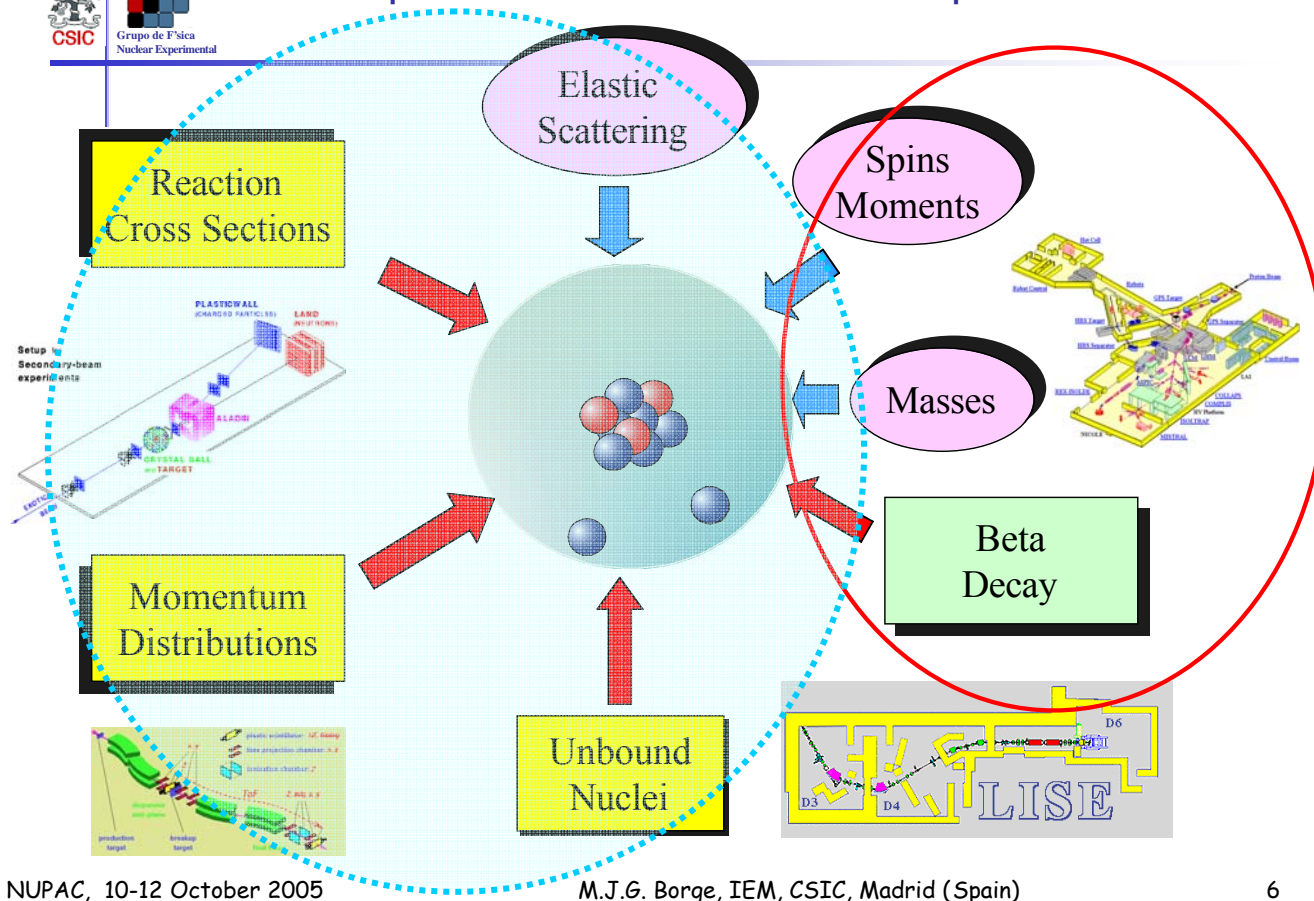
Theoretical perspectives

SEARCH INSIDE!™



-
-
-
-

Experimental Studies at the drip line



Let us see how the “center-left” Physics needs **ISOLDE** and even an upgrade of it.

Unbound
Nuclei

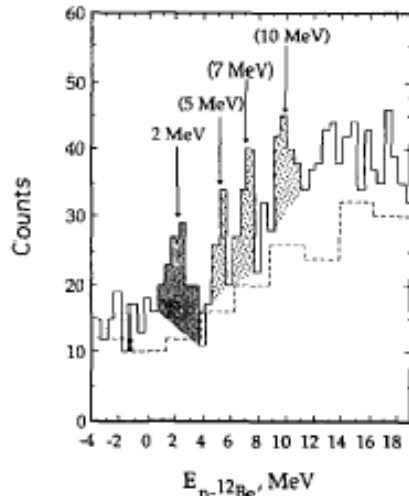
Nuclei *at* the driplines

They are what the reaction mechanism makes them to be...

Interplay between structure and reaction aspects...bound vs. continuum...stronger than in any other part of the nuclear chart.

A perfect balance between short and long range parts of the nuclear interaction.

¹³Be: an example of *creation* by the reaction mechanism



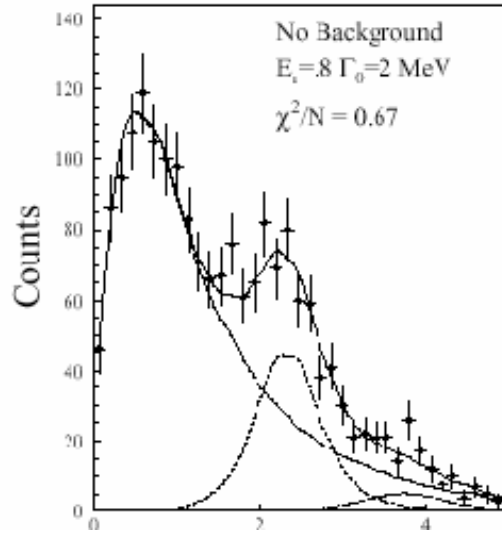
•transfer to the continuum: ¹²Be (d,p) RIKEN (Korshennikov) (1995).

cf. ¹⁰Li at REX-ISOLDE
H.B. Jeppesen, et al.,
NP A748, (2005)374

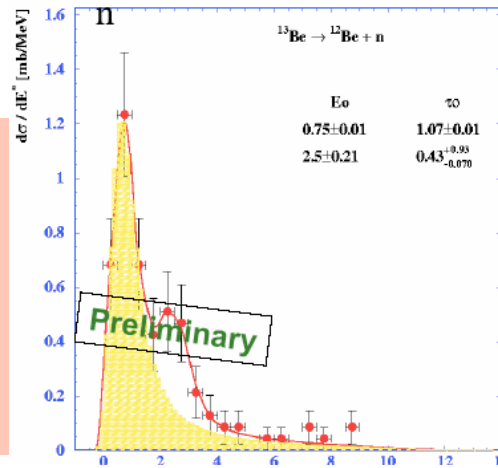
- ¹⁴Be nuclear and Coulomb breakup: GANIL (K. Jones thesis, 2000).
- ¹⁴C + ¹¹B multinucleon transfer: (Berlin Group, 1998).
- ¹⁸O fragmentation MSU (Thoennessen, 2001) n-core relative velocity spectra.

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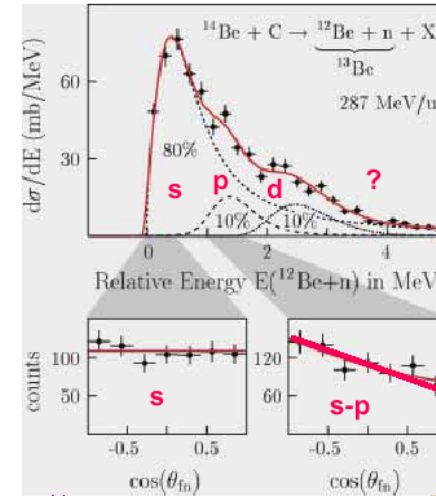
•¹⁴B fragmentation: GANIL (Lecouey, Orr) (2002).



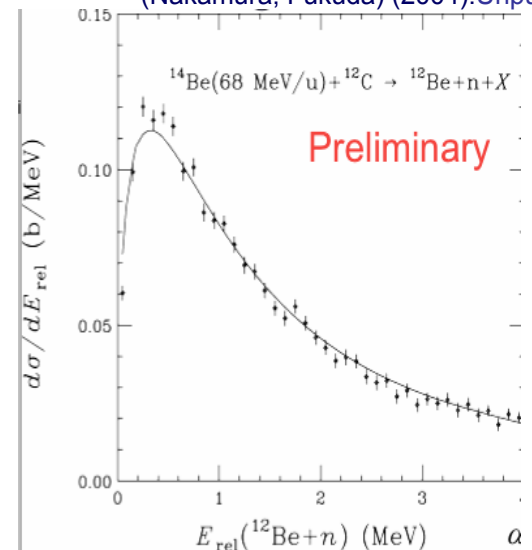
•GSI (U. Datta Pramanik) (2004).Unpublished

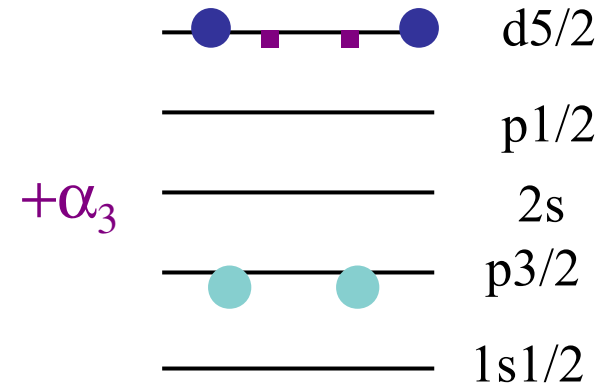
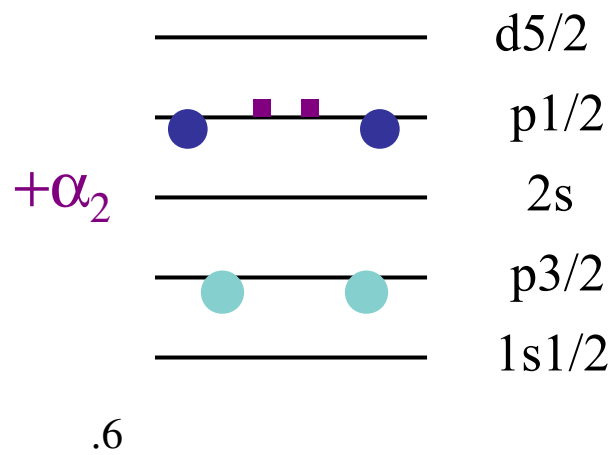
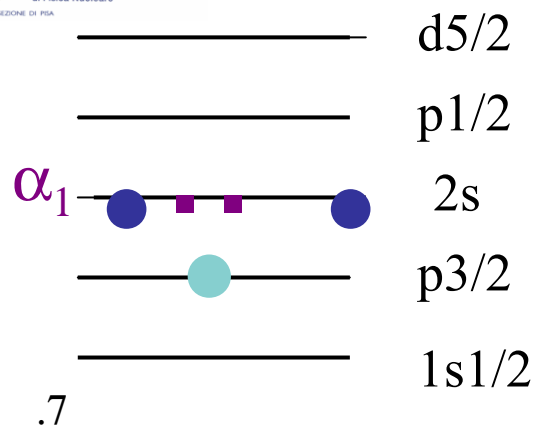


•¹⁴Be nuclear breakup, GSI (Simon), 287MeV, n-core angular correlation



•¹⁴Be nuclear breakup: RIKEN (Nakamura, Fukuda) (2004).Unpublished





$$^{14}\text{B}_{\text{g.s.}} = 2^- = \pi p_{3/2} + \nu 2s$$

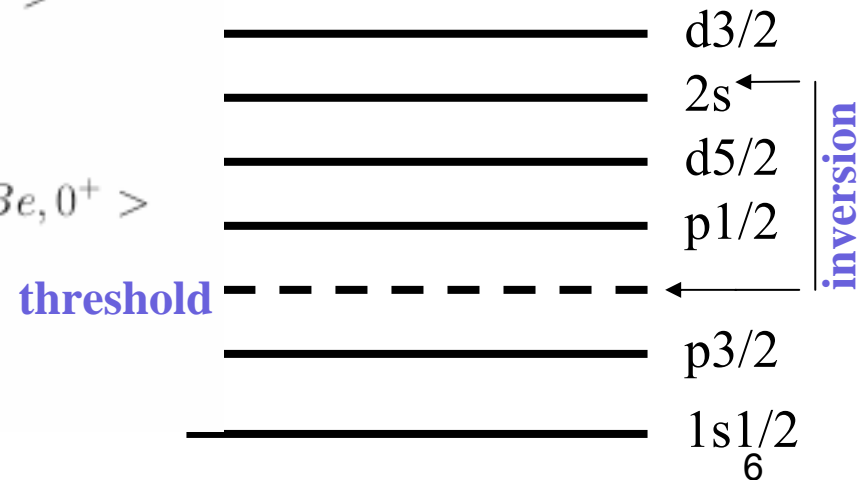
$$^{12}\text{Be}_{\text{g.s.}} = 0^+$$

$$^{14}\text{Be}_{\text{g.s.}} = 0^+ (?)$$

$$|^{14}\text{B}\rangle = [a_1(p_{3/2}, 2s_{1/2}) + a_2(p_{3/2}, d_{5/2})] \otimes |^{12}\text{Be}, 0^+\rangle$$

$$|^{14}\text{Be}\rangle = [b_1(s_{1/2})^2 + b_2(1p_{1/2})^2 + b_3(1d_{5/2})^2] \otimes |^{12}\text{Be}, 0^+\rangle$$

$$c_3(1d_{5/2})^2 \otimes |^{12}\text{Be}^*, 2^+\rangle$$



Potential corrections due to the particle-vibration coupling (N. Vinh Mau and J. C. Pacheco, NPA607 (1996) 163.

also T. Tarutina, I.J. Thompson, J.A. Tostevin NPA733 (2004) 53)

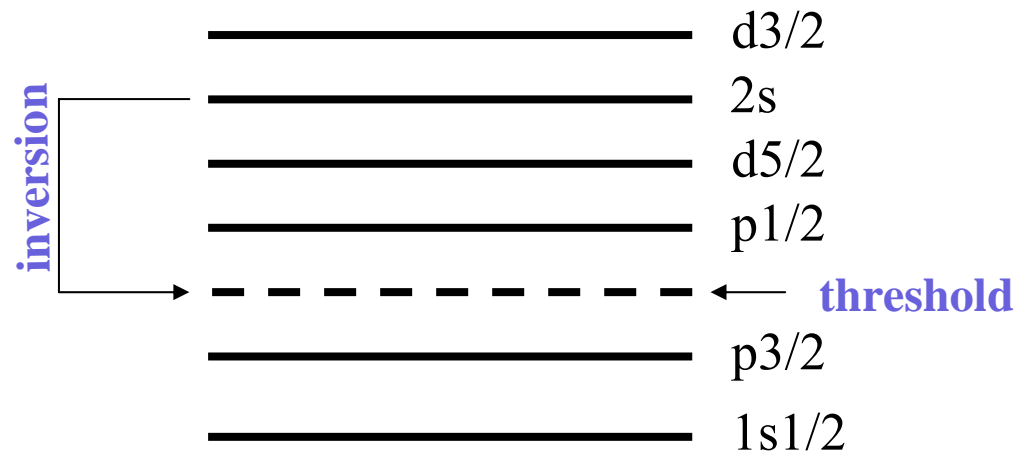
...can be modeled as:

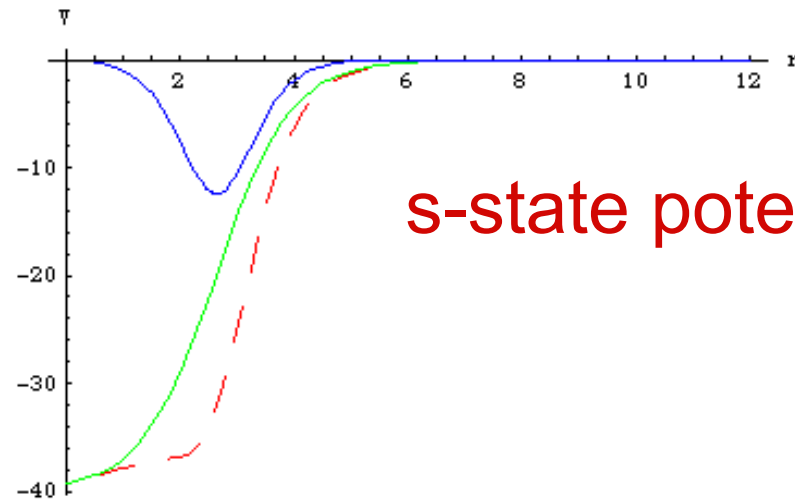
$$U(r) = V_{WS} + V_{so} + \delta V$$

$$\delta V(r) = 16 \alpha e^{(r-R)/a} / (1 + e^{(r-R)/a})^4$$

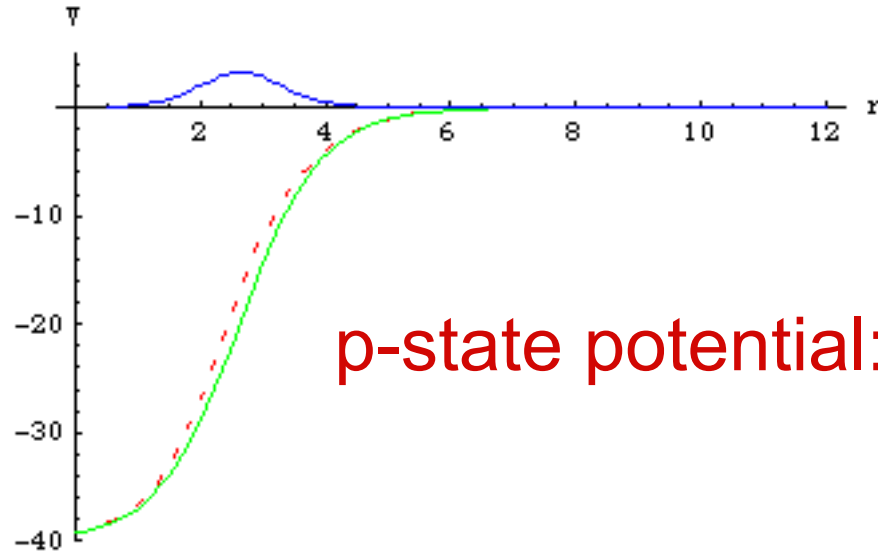
$n+^{12}\text{Be}$:

	ε_{res} (MeV)	Γ (MeV)	α (MeV)
$1p_{1/2}$	0.67	0.28	8.34
$1d_{5/2}$	2.0	0.40	-2.36





s-state potential: long range added

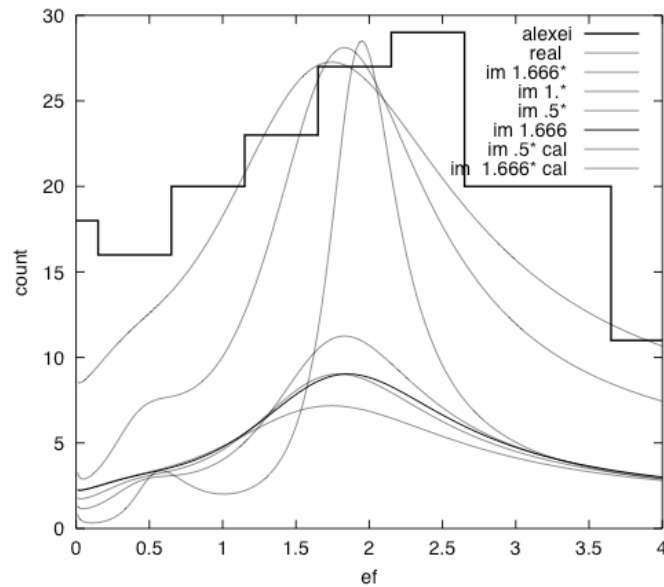


p-state potential: long range subtracted

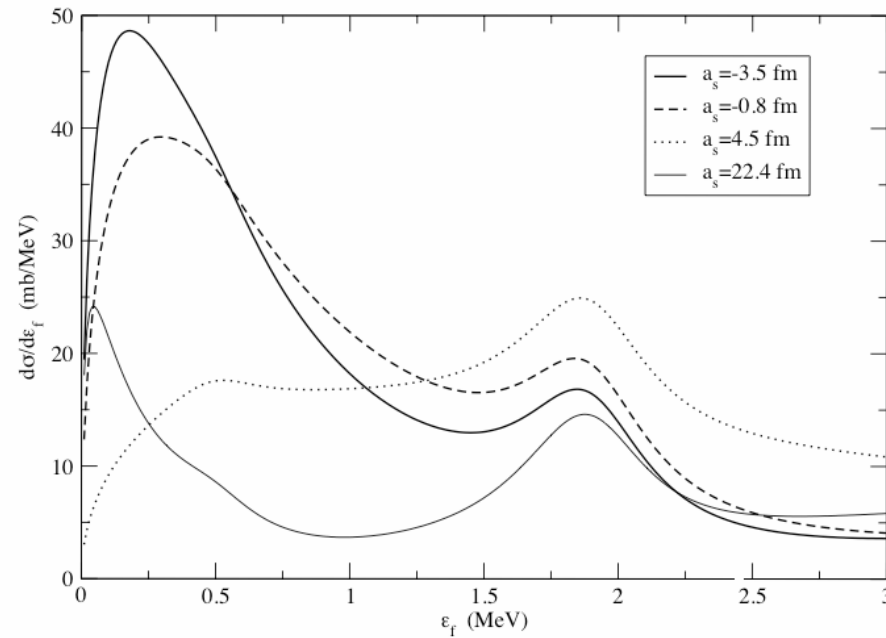
G. Blanchon, A. B. and N. Vinh Mau
Unbound exotic nuclei studied by transfer
to the continuum reactions
Nucl. Phys. A739 (2004) 259.

G. Blanchon, A. B., D. M. Brink and N. Vinh Mau,
Unbound exotic nuclei studied by projectile
fragmentation reactions.
IFUP-TH 29/2004: submitted to NPA

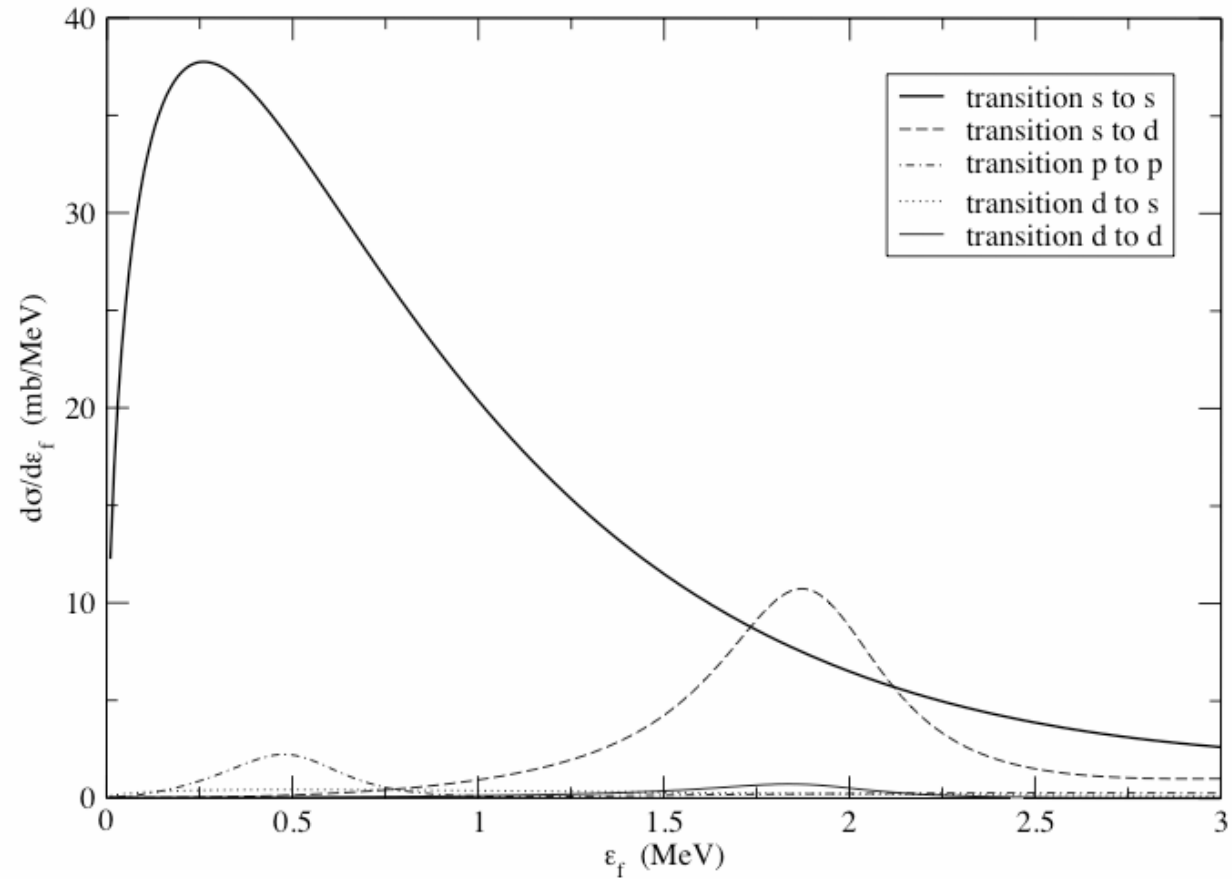
^{12}Be (d,p)



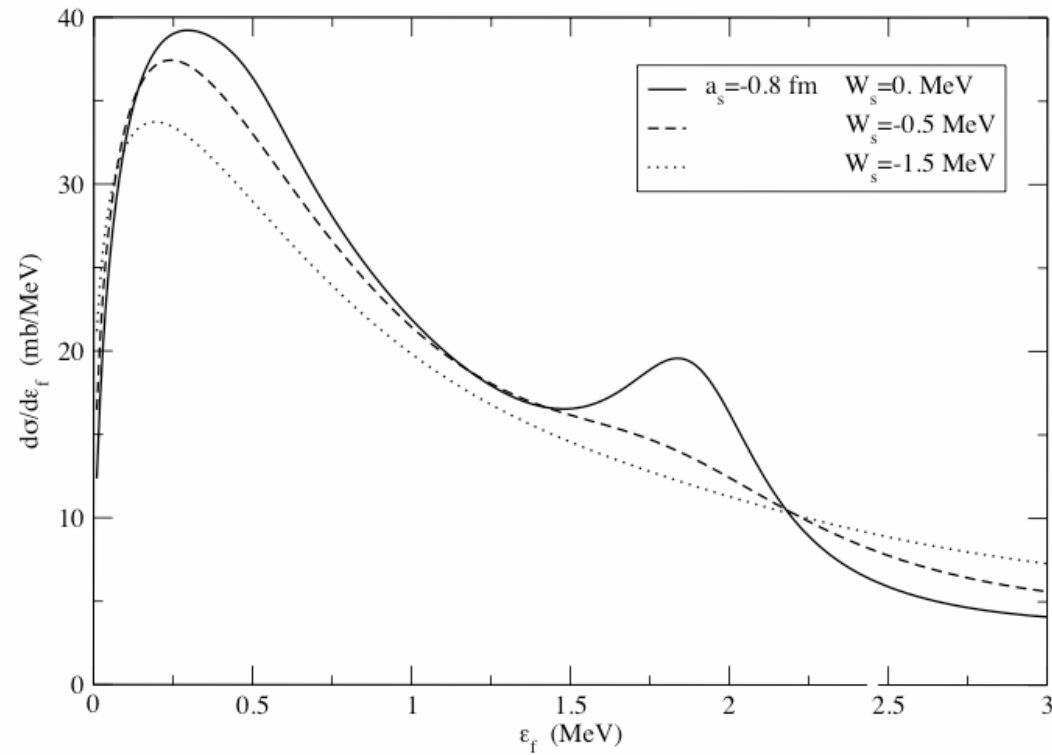
Final s-state bound vs continuum



Bound to unbound transitions



Core excitation simulated by imaginary part of n-core optical potential



effective range theory (1⁰ order)

$$\rightarrow k \cotan \delta = - \frac{1}{a_s} + \frac{1}{2} r_o k^2$$

α	a_s	r_e	$ \epsilon $
(MeV)	(fm)	(fm)	(MeV)
8.0	-0.8	117.0	
4.0	-3.5	17.9	
2.0	-6.6	11.8	
-1.0	-26.1	7.58	
-5.0	22.4	5.9	0.06
-15.0	7.1	3.8	1.34
-35.0	4.5	2.7	6.49

All aspects of the nuclear interaction discussed in structure models: short range vs. long range, real vs. imaginary...volume vs. surface...

amplified at the driplines

J.Enders et al., PHYSICAL REVIEW C **65** 034318

Me
 Dis

ons

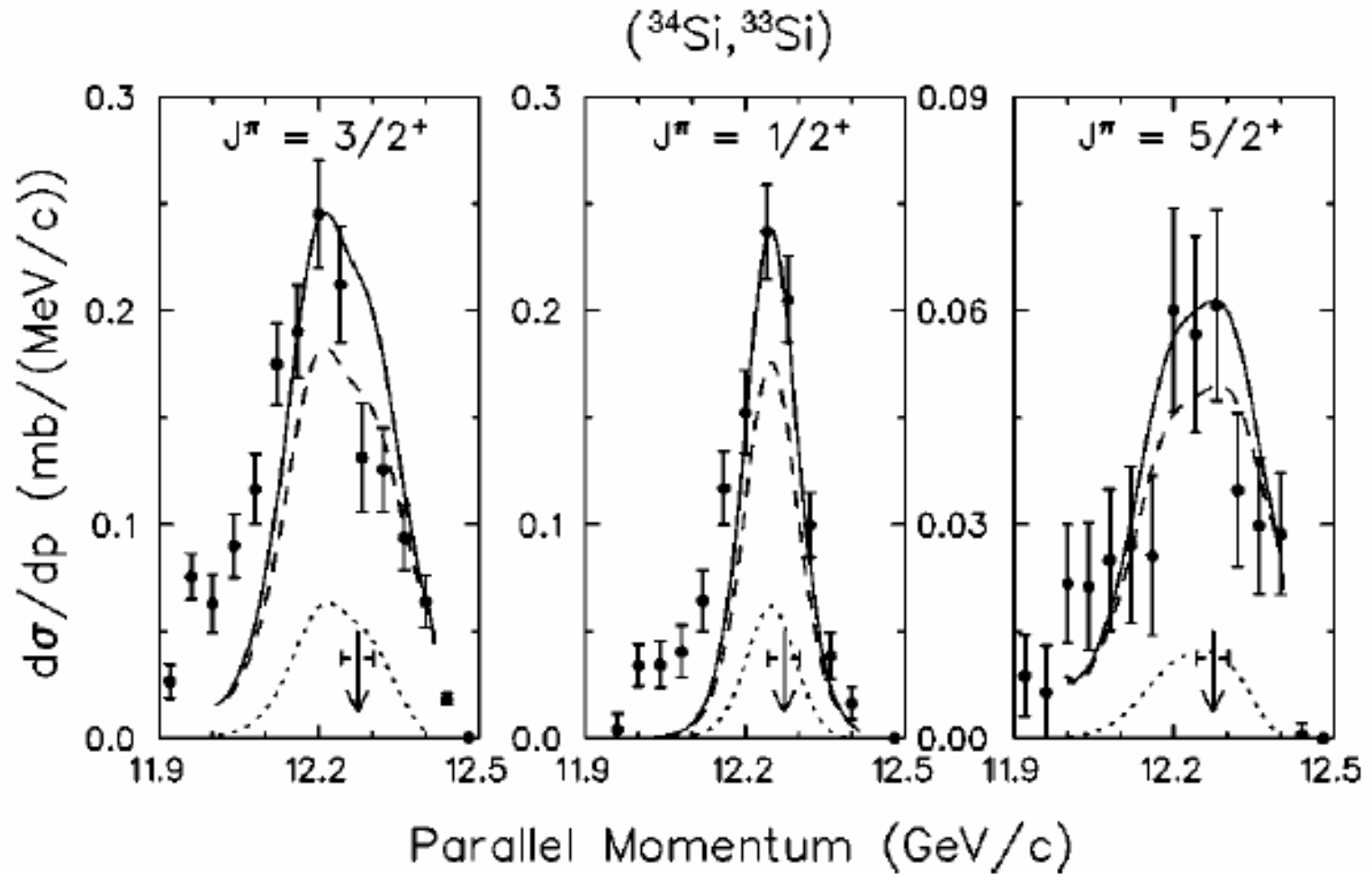
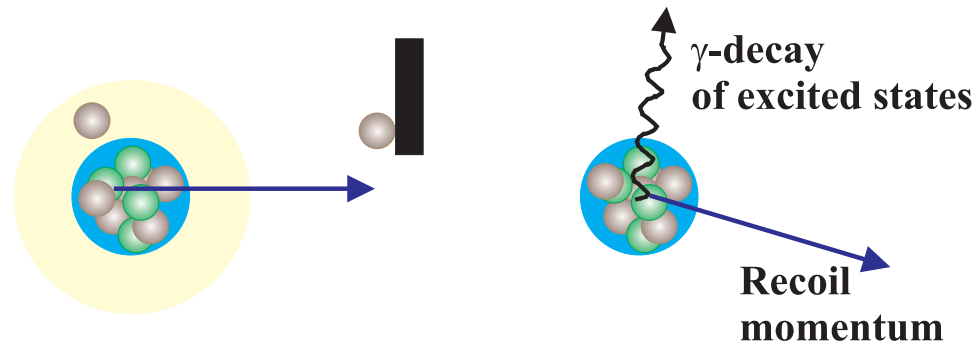


FIG. 4. Measured differential cross-sections for the $(^{34}\text{Si}, ^{33}\text{Si})$ system.

Spectroscopic factors from knockout reactions

(P.G.Hansen et al.)

- Spectroscopic factors from excited states of fragment.



$$\sigma(nI^\pi) = \sum_j C^2 S(j, nI^\pi) \sigma_{sp}(j, B_n)$$

$$\sigma_{sp}(j, B_n) = \sigma_{sp}^{strip}(j, B_n) + \sigma_{sp}^{diffr}(j, B_n)$$

- ℓ of nucleon from momentum distribution of fragment.

Ela
Nuc

Elastic scattering and fusion studies in the reactions $^{10,11}\text{Be} + ^{64}\text{Zn}$.

alo

A. Amorini¹⁾, C. Angulo²⁾, M.J.G. Borge³⁾, A. Di Pietro¹⁾, P. Figuera¹⁾, L.M. Fraile⁴⁾, M. Lattuada¹⁾,
M. Milin⁵⁾, F. Pansini¹⁾, M.G. Pellegriti⁶⁾, R. Raabe⁷⁾, F. Rizzo¹⁾, M. Sawicka⁷⁾, V. Scuderi¹⁾, O.
Tengblad³⁾, M. Zadro⁵⁾

1) INFN- Laboratori Nazionali del Sud and Università di Catania, Catania, Italy

2) CRC- Louvain la Neuve, Belgium

3) Insto. de Estructura de la Materia, CSIC, Madrid, Spain

4) CERN, Geneva, Switzerland

5) Ruđer Bošković Institute, Zagreb, Croatia

6) GANIL, Caen, France

7) Instituut voor Kern- en Stralingsfysica, University of Leuven, Leuven, Belgium

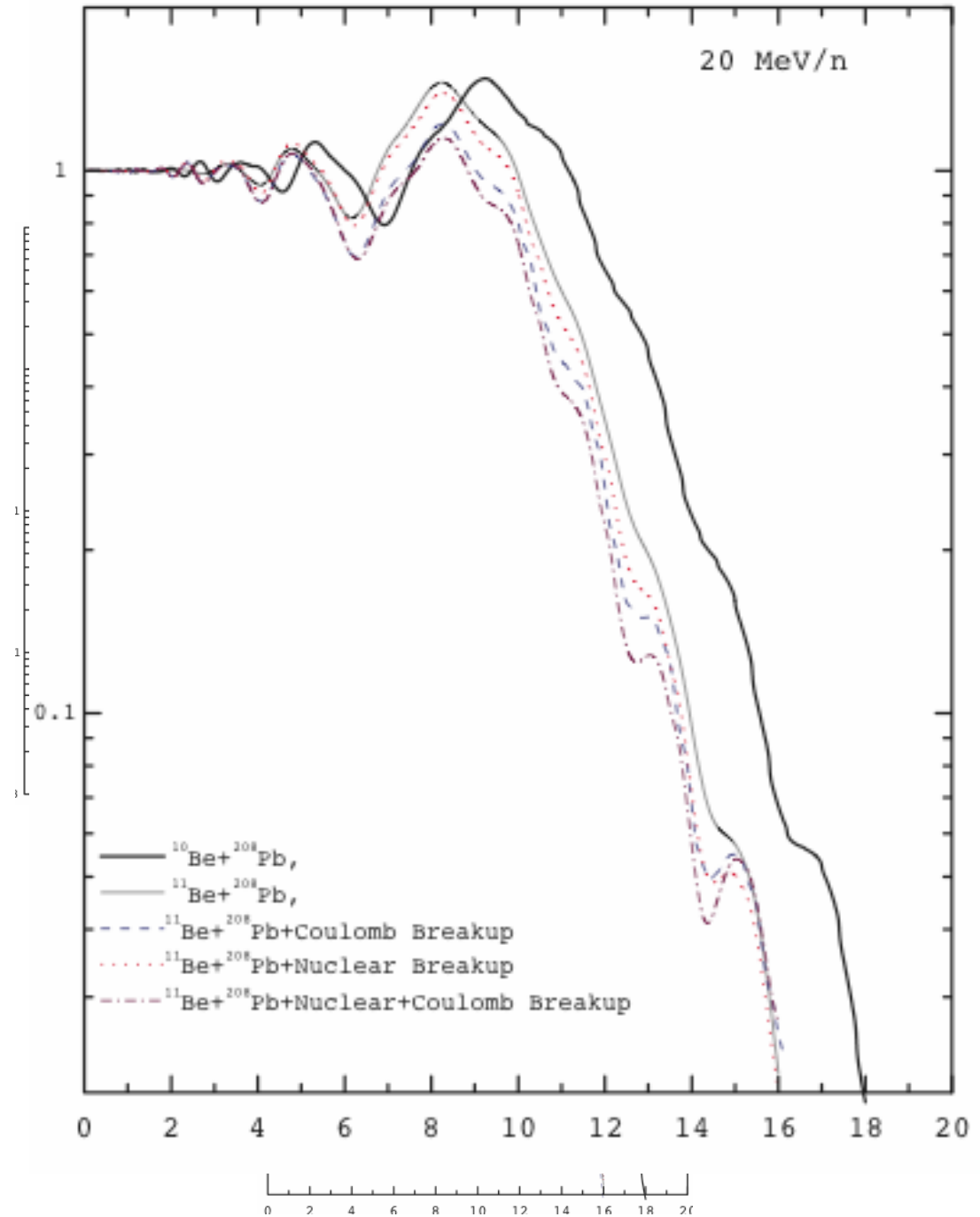
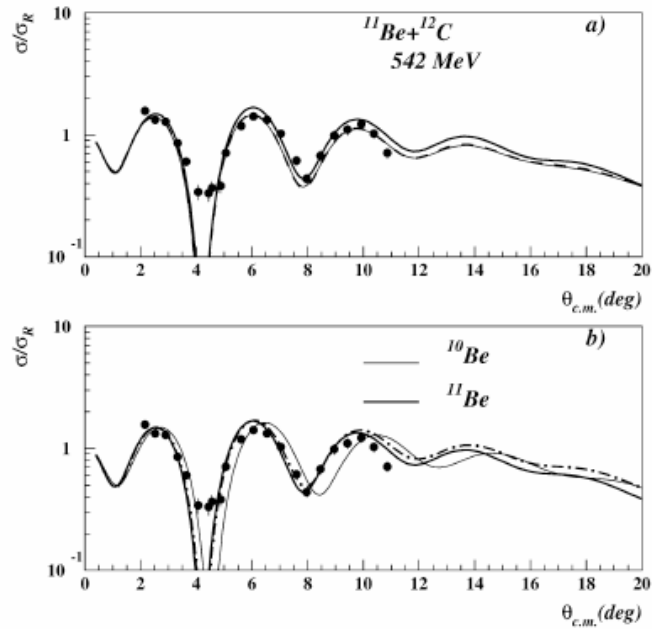
Spokeperson : A. Di Pietro

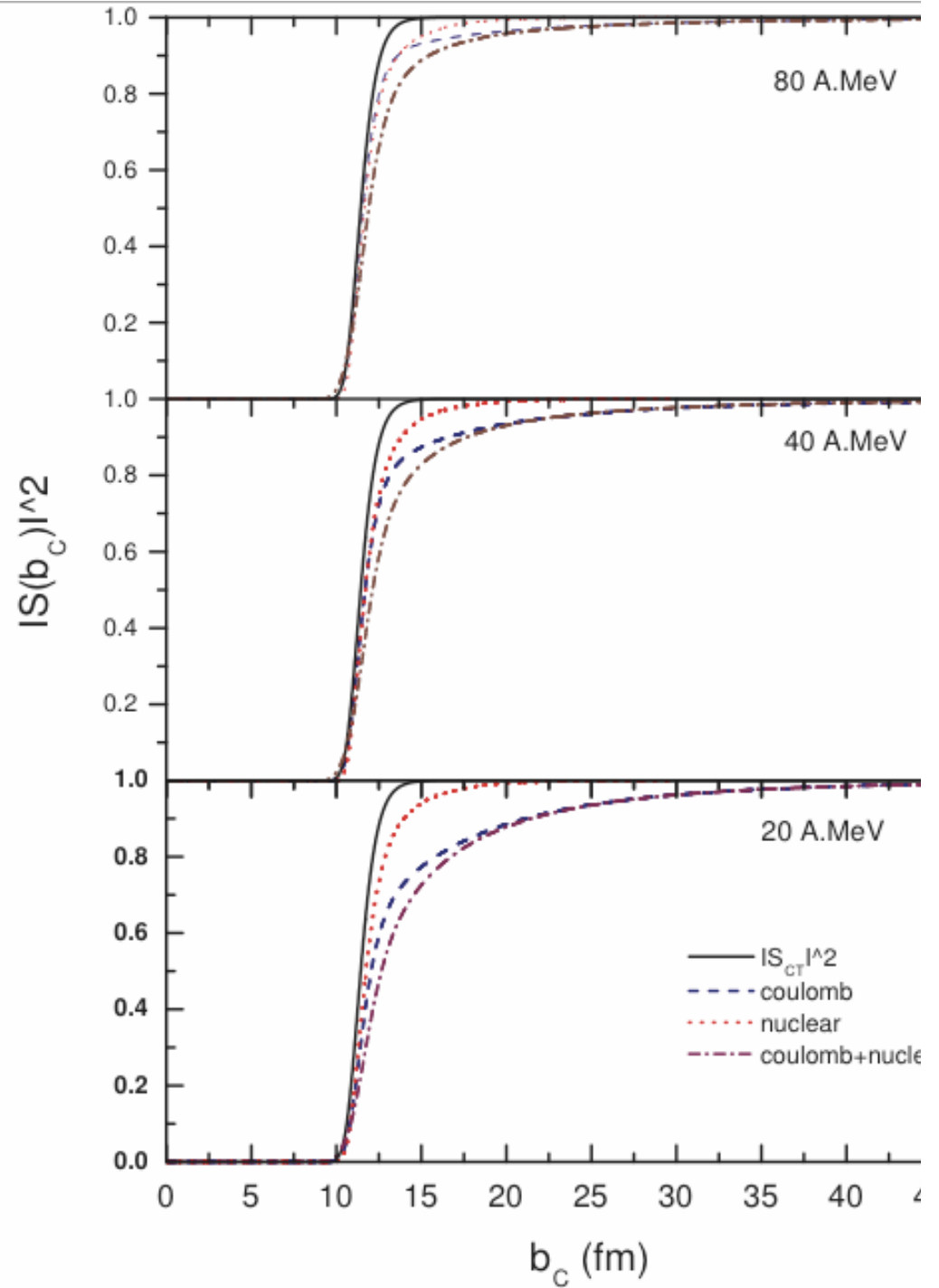
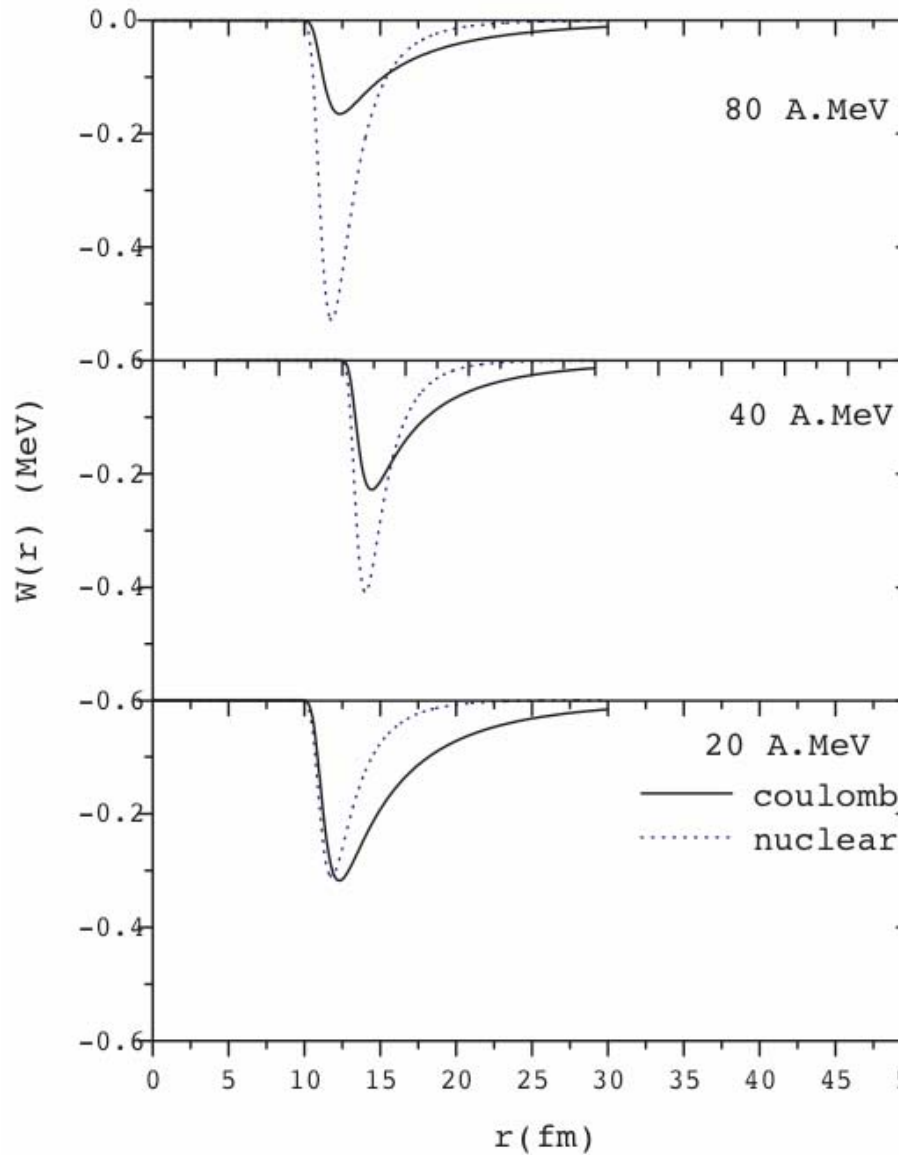
Contactperson: L.M. Fraile

Abstract

We propose to measure elastic scattering and fusion excitation functions for the reactions $^{10,11}\text{Be} + ^{64}\text{Zn}$ at 3.1 MeV/u. The aim of the experiment is to investigate possible effects of the halo structure of the ^{11}Be nucleus on the reaction mechanisms at energy around the Coulomb barrier. For this purpose a comparison with the reaction induced by the ^{10}Be nucleus is required.

A. B. and F. Carstoiu
Optical potentials of halo and
weakly bound nuclei
Nucl. Phys. A706 (2002) 322.





	Nuclear			Coulomb			
E_{inc}	20	40	80	20	40	80	
A_1	1860.38	303.23	339.45	9.672	3.922	1.734	B_1
A_2	19.73	7.91	1.91	1.558	0.6013	0.2549	B_2
-	-	-	-	0.1531	0.0587	0.0248	B_3
α_1	1.1924	1.5334	1.5356	2.7273	2.9438	3.1046	β_1
α_2	2.5403	2.7918	3.4955	6.5189	7.4850	8.2713	β_2

$$W_S^N(r) = -\frac{\hbar v}{2} P_0 \sum_n A_n \exp(-r/\alpha_n) \frac{1}{\sqrt{2\pi\alpha_n r}}$$

$$W_S^C(r) = -\frac{\hbar v}{2} P_0 \sum_n B_n \exp(-r/\beta_n) \frac{1}{\sqrt{2\pi\beta_n r}}$$

$$\alpha_2 \approx 1/2\gamma, \quad \gamma = \sqrt{2\mu\varepsilon_i/\hbar}$$

decay length of initial wave function

$$\beta_3 \approx (\varepsilon_f - \varepsilon_i) / \hbar v$$

adiabaticity parameter of Coulomb excitation theory

...see also a classical reference

**Effect of the breakup channel on ^{11}Li elastic scattering,
J.S. Al-Khalili, Nucl. Phys. A 581 (1995) 315-330**

...and one of the most recent

PHYSICAL REVIEW C **72**, 034606 (2005)

Elastic scattering of the proton drip-line nucleus ^{17}F at 10A.MeV

J. C. Blackmon,¹ F. Carstoiu,^{2,3} L. Trache,² D. W. Bardayan,¹ C. R. Brune,⁴ C. A. Gagliardi,² U. Greife,⁵ C. J. Gross,¹
C. C. Jewett,⁵ R. L. Kozub,⁶ T. A. Lewis,¹ J. F. Liang,¹ B. H. Moazen,⁶ A. M. Mukhamedzhanov,² C. D. Nesaraja,^{1,6}
F. M. Nunes,⁷ P. D. Parker,⁸ L. Sahin,⁹ J. P. Scott,^{1,6} D. Shapira,¹ M. S. Smith,¹ J. S. Thomas,¹⁰ and R. E. Tribble²

...and references therein.

For the *structure* aspect...see also

“Evolution of nuclear structure, shapes, and fission”

**Evolution Of Shell Structure,
Shapes & Collective Modes**

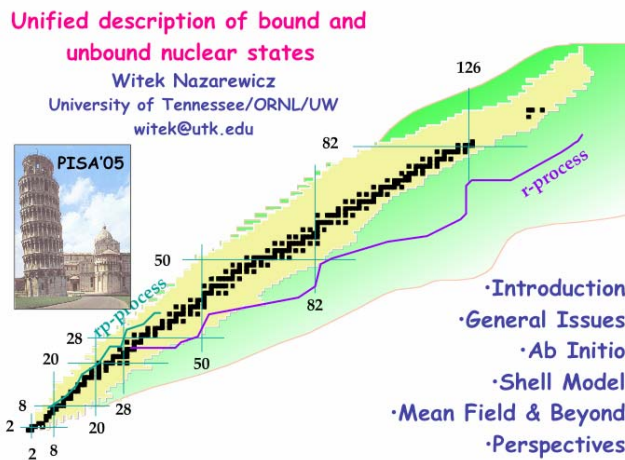
Dario Vretenar
vretenar@phy.hr

Shell model for weakly bound and unbound nuclear states:
Bound and continuum states in one framework

GANIL - ORNL Theory Collaboration

N. Michel, W. Nazarewicz, J. Okolowicz, M. Ploszajczak, J. Rotureau

<http://www.df.unipi.it/~angela/finalprogram.html>



Structure of proton-radioactive nuclei

Lidia S Ferreira and Enrico Maglione 2005 J. Phys. G: Nucl. Part. Phys. 31 S1569-S1572

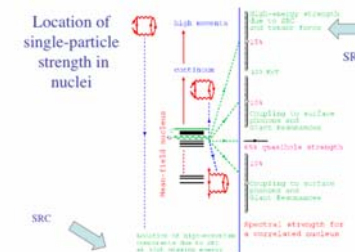
NuPAC 2005

Trento 3/4/04

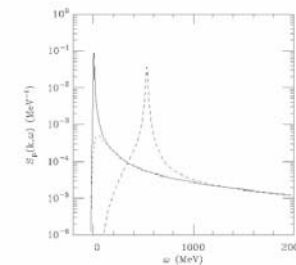
Proton properties in the nucleus
- *quantum physics at all energy scales* -

Wim Dickhoff
Washington University in St. Louis
(where the photon was discovered)

<http://www.df.unipi.it/~angela/talks.html>



Where the depleted strength ends up ...



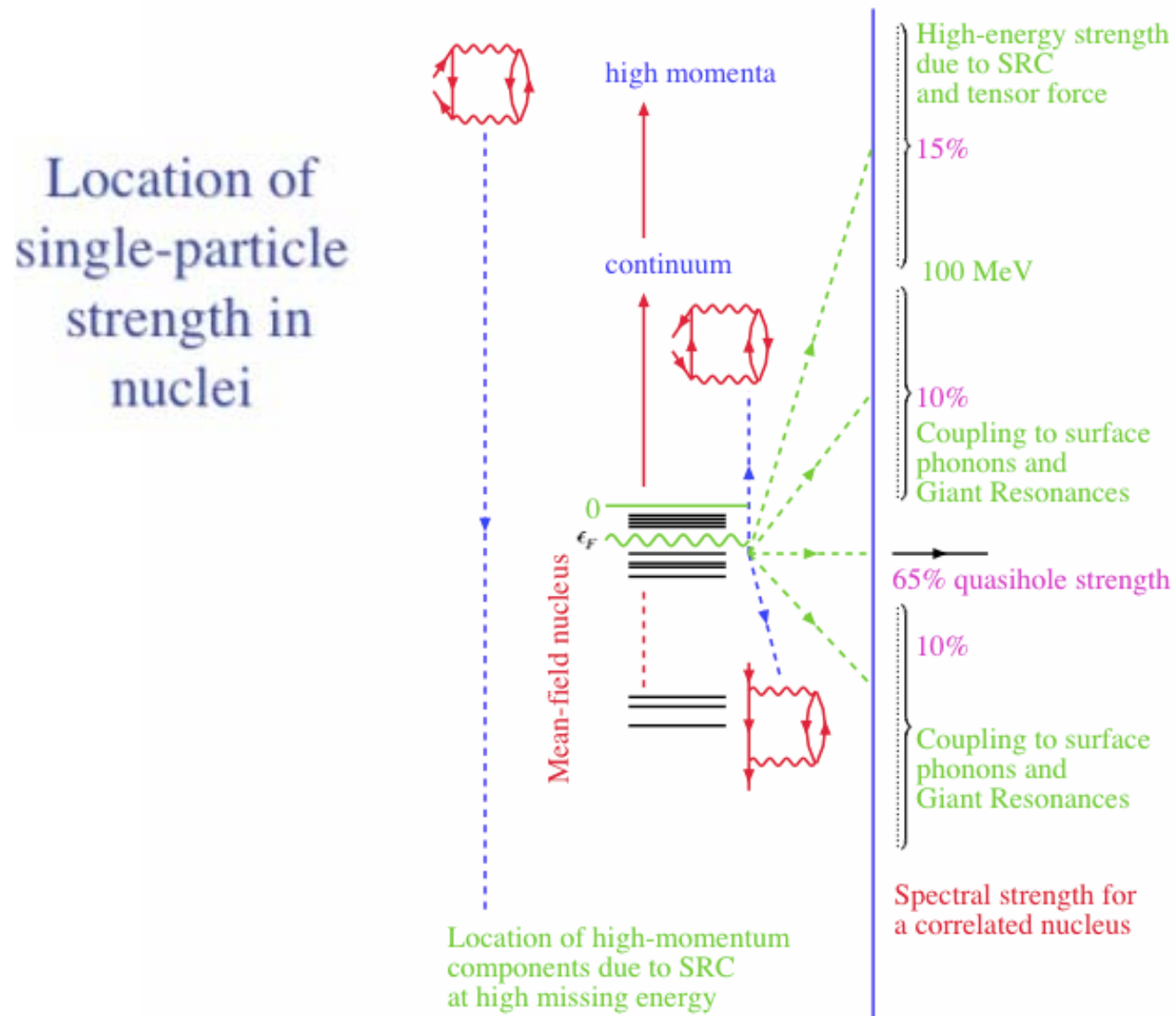


Fig. 55. The distribution of single-particle strength in a nucleus like ^{208}Pb . The present summary is a synthesis of experimental and theoretical work discussed in this review. A slight reduction (from 15% to 10%) of the depletion effect due to SRC must be considered for light nuclei like ^{16}O .

Conclusions and Outlook

- “Exotic” threshold characteristics of dripline nuclei enhanced by low energy experiments at $E_{inc} \leq 20A.MeV$.
- Study of dripline nuclei will lead us to the *best possible phenomenological energy dependent nuclear potential* for *bound* and *scattering* states.

My guess for the future:

- We will have to redefine the *n-n* interaction in terms of a “n-body” interaction (3-body terms are already necessary !!).
- Pairing (*n-p*) will be understood (I hope!!) on a microscopic basis...(*effective field theories ??*).



Hi Gregers, would you agree...?

Appendix I: transfer vs. inelastic

transfer to the continuum

$$\frac{dP_t}{d\varepsilon_f} \approx \sum_{l_f} (2l_f + 1) |1 - S_{l_f}|^2 B_{l_f, l_i},$$

$$B_{l_f, l_i} \approx \frac{e^{-2\eta b_c}}{b_c}$$

projectile fragmentation: inelastic excitation

$$\frac{dP_{in}}{d\varepsilon_f} \approx \sum_{l_f} (2l_f + 1) |1 - \bar{S}_{l_f}|^2 I_{l_f, l_i},$$

$$I_{l_f, l_i} \approx \frac{e^{-2\gamma b_c}}{b_c^3}$$

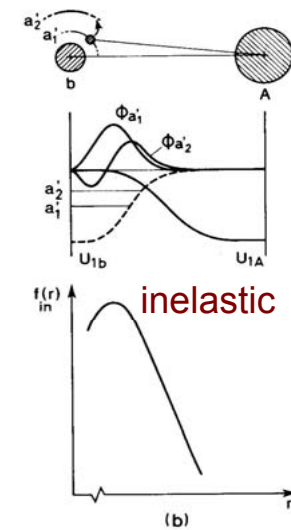
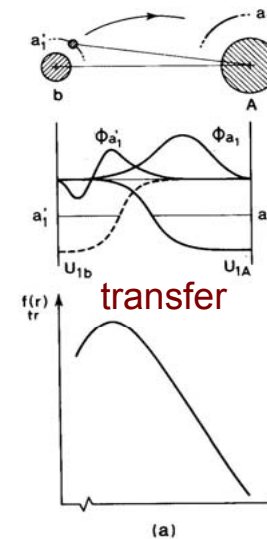
$$\bar{S}_{l_f} = S_{l_f} e^{2i\alpha} = e^{2i(\delta + \alpha)}$$

S and \bar{S} correspond to different phase shifts

check of sudden approximation $\left\{ \begin{array}{l} E_{inc}: \text{independent} \\ \varepsilon_{if} : \text{important} \end{array} \right.$

V Transfer Reactions

Broglia & Winther book



Appendix II: potential from phase shift.

$$|S_{NN}(b)|^2 = e^{-4\delta_I(b)}$$

$$\delta_I(b) = -\frac{1}{2\hbar} \int_{-\infty}^{+\infty} (W_V(\mathbf{r}(t)) + W_S(\mathbf{r}(t))) dt$$

$$\int_{-\infty}^{+\infty} W_S(\mathbf{r}(t)) dt = -\frac{\hbar}{2} P_{b_{up}}$$

$$\mathbf{r}(t) = \mathbf{b}_c + vt$$

$$|S_{NN}|^2 = |S_{CT}|^2 e^{-P_{b_{up}}}$$

$$W_S^N(r) = -\frac{\hbar v}{2} p_{b_{up}}^N(r) \frac{1}{\sqrt{2\pi ar}}$$