

Exploring New Physics with B Physics

Workshop on Flavour in the Era of the LHC
CERN – November 7, 2005

Yossi Nir (*Weizmann Institute of Science*)

Thanks to:

- Guy Raz
- Stephane Monteil *et al.* (CKMfitter)
- Luca Silvestrini (UTfit)

Plan of Talk

1. Recent Era: Excluding alternatives to the KM mechanism

- (a) Is the KM mechanism at work?
- (b) Is δ_{KM} the only source of CPV in meson decays?
- (c) Is CPV in $K \rightarrow \pi\pi$ small because of flavor suppression?
- (d) Is there direct CPV?
- (e) Is there New Physics in $B^0 - \bar{B}^0$ mixing?
- (f) Is there New Physics in $b \rightarrow s$ transitions?

2. Future: Looking for corrections to the KM mechanism

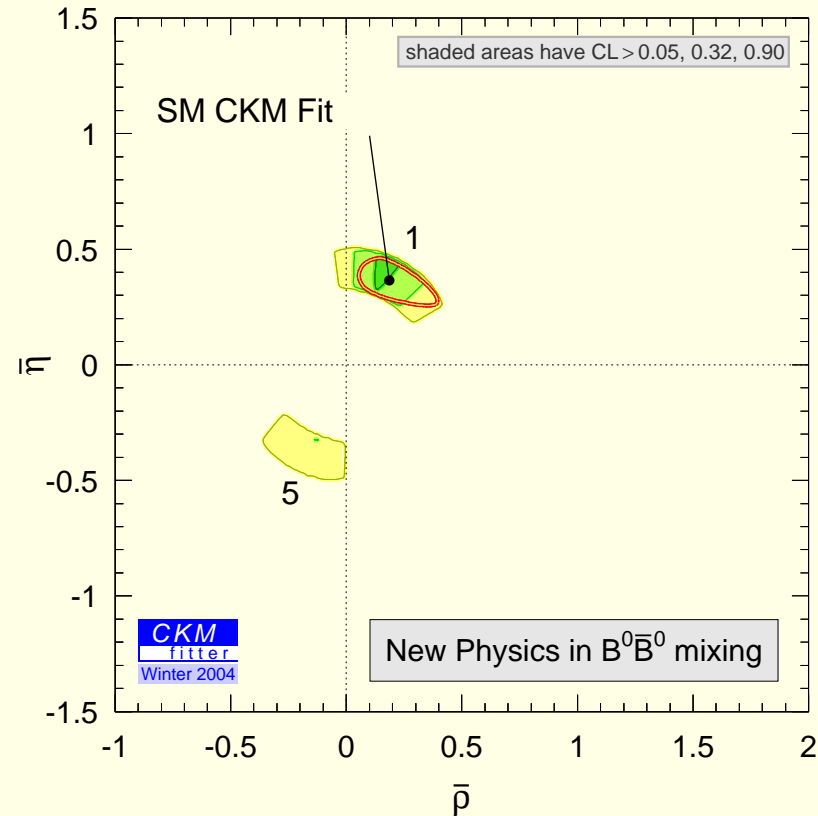
$S_{\pi^0 K_S}$: SM and NP

- (a) Factorization related methods
- (b) SU(3) based methods
- (c) Supersymmetry as an example

Is the KM mechanism at work?

- Assume: New Physics in tree decays - negligible
- Define $r_d^2 \exp(2i\theta_d) = \langle B^0 | \mathcal{H}^{\text{full}} | \bar{B}^0 \rangle / \langle B^0 | \mathcal{H}^{\text{SM}} | \bar{B}^0 \rangle$
- Use $|V_{ub}/V_{cb}|, \mathcal{A}_{DK}, S_{\psi K}, S_{\rho\rho}, \Delta m_{B_d}, \mathcal{A}_{\text{SL}}$
- Fit to $\boxed{\eta}, \rho, r_d, \theta_d$
- Find whether $\eta = 0$ is allowed

The KM mechanism is at work!



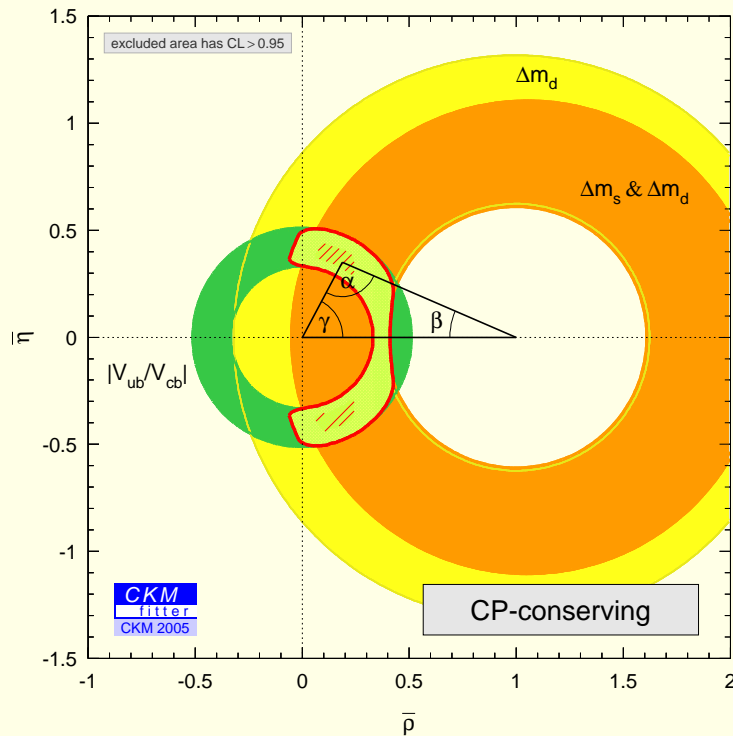
SM tree level + NP in $B^0 - \bar{B}^0$ mixing ($\Delta m_B, S_f, A_{SL}$)

$$\delta_{\text{KM}} \neq 0$$

CKMFitter, hep-ph/0406184

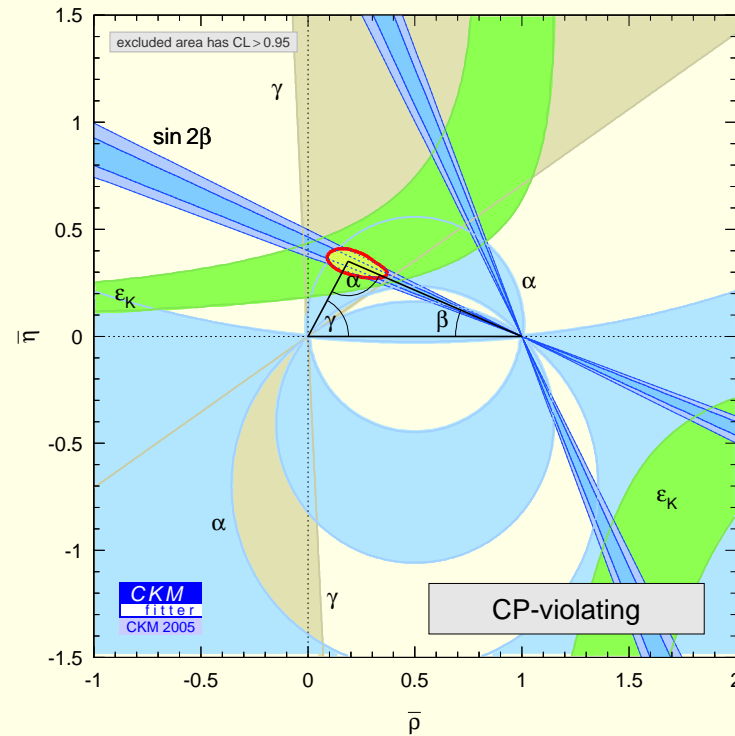
β, α, γ

Is δ_{KM} the only source of CPV in meson decays?



Tree level + CPC observables

$$\Delta m_B, \Delta m_{B_s}$$



Tree level + CPV observables

$$\varepsilon, S_{\psi K}, \alpha, \gamma$$

Using CKMFitter package (Höcker *et al.*, Eur. Phys. J. C21, 225 (01))

Very likely, δ_{KM} is dominant!

$$S_{\psi_K} = +0.69 \pm 0.03 \Leftrightarrow \sin 2\beta(\text{CKM fit}) = +0.74_{-0.03}^{+0.07}$$

The KM mechanism successfully passed
its first precision test

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$$\alpha(\pi\pi, \pi\rho, \rho\rho) = [101_{-9}^{+16}]^\circ \Leftrightarrow \alpha(\text{CKM fit}) = 96 \pm 16^\circ$$

The KM mechanism successfully passed its second precision test

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$$\gamma(DK) = [63_{-13}^{+15}]^\circ \Leftrightarrow \gamma(\text{CKM fit}) = [57_{-14}^{+7}]^\circ$$

The KM mechanism successfully passed its third precision test

Is CPV in $K \rightarrow \pi\pi$ small because of flavor?

SM:

- $\epsilon \sim 10^{-3}$, $\epsilon' \sim 10^{-5}$ because of flavor suppression
- Some CP violating phases are order one

Approximate CP:

- All CPV phases are small
- All CP asymmetries are small

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- All CP asymmetries are small

B Physics:

- $S_{\psi K} \sim 0.7$
- \implies Some CP violating phases are indeed $\mathcal{O}(1)$

$$\mathcal{A}_{K \mp \pi \pm}, \mathcal{A}_{\rho \pi}^{-+}$$

Is CP violated in $\Delta B = 1$ processes?

SM:

- Indirect (M_{12}) and direct (A_f) CP violations are both large

Superweak:

- There is no direct (A_f) CP violation

K Physics:

- $\epsilon'/\epsilon = (1.72 \pm 0.18) \times 10^{-3}$

\implies CP is violated in $\Delta S = 1$ processes ($s \rightarrow u\bar{u}d$)

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B Physics:

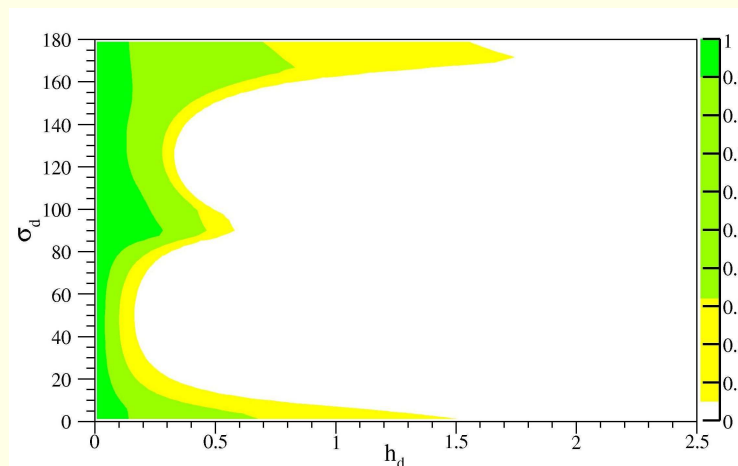
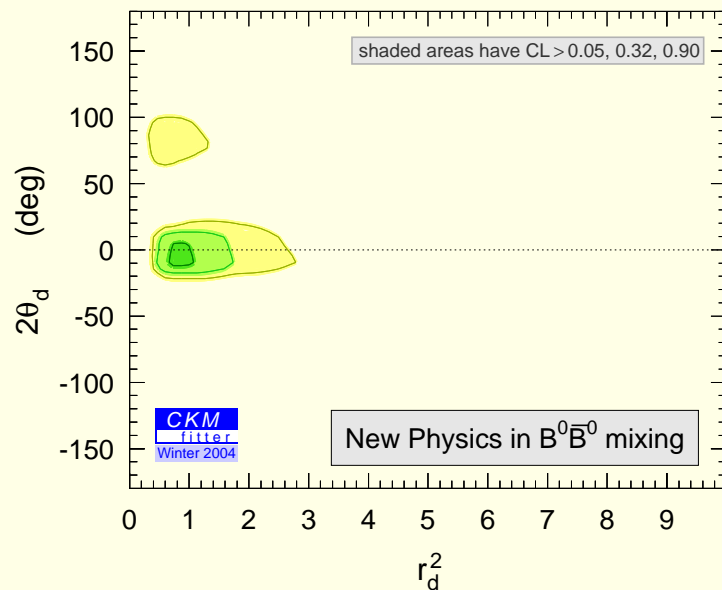
- $\mathcal{A}_{K\mp\pi\pm} = -0.115 \pm 0.018$, $\mathcal{A}_{\rho\pi}^{-+} = -0.48 \pm 0.14$

\implies CP is violated in $\Delta B = 1$ processes ($b \rightarrow u\bar{u}s$, $b \rightarrow u\bar{u}d$)

Is there NP in $B^0 - \bar{B}^0$ mixing?

- Assume: New Physics in tree decays - negligible
- Define $r_d^2 \exp(2i\theta_d) \equiv 1 + h e^{i\sigma} \equiv \langle B^0 | \mathcal{H}^{\text{full}} | \bar{B}^0 \rangle / \langle B^0 | \mathcal{H}^{\text{SM}} | \bar{B}^0 \rangle$
- Use $|V_{ub}/V_{cb}|, \mathcal{A}_{DK}, S_{\psi K}, S_{\rho\rho}, \Delta m_{B_d}, \mathcal{A}_{\text{SL}}$
- Fit to $\eta, \rho, \boxed{r_d}, \theta_d$ (or \boxed{h}, σ)
- Find whether $h \neq 0$ ($r_d \neq 1$) is allowed

Very likely, the KM mechanism dominates



CKMFitter

Agashe et al., hep-ph/0509117

For arbitrary phase, $h = |A_{NP}/A_{SM}| \sim 0.2 \pm 0.2$

$$B \rightarrow X_s \gamma, \quad B \rightarrow X_s \ell^+ \ell^-$$

Is there NP in $b \rightarrow s$ transitions?

- Rare $b \rightarrow s$ processes consistent with the SM predictions

\implies New Physics contributions to certain operators are strongly constrained (Z -penguin, magnetic)

\implies New physics contributions to other operators are still very weakly constrained (chromomagnetic, dim-6)

Is there NP in $b \rightarrow s$ transitions?

$$\text{SM: } \Delta S \equiv -\eta_{\text{CP}} S - S_{\psi K} \approx 0, \quad C \approx 0$$

f_{CP}	ΔS	C
ϕK_S	-0.22 ± 0.19	-0.09 ± 0.15
$\eta' K_S$	$-0.19 \pm 0.09^\dagger$	-0.07 ± 0.07
$f_0 K_S$	$+0.06 \pm 0.24$	$+0.06 \pm 0.21$
$\pi^0 K_S$	-0.38 ± 0.26	-0.02 ± 0.13
ωK_S	-0.06 ± 0.30	-0.44 ± 0.24
$K_S K_S K_S$	-0.09 ± 0.23	-0.31 ± 0.17
$K^+ K^- K_S^\ddagger$	-0.16 ± 0.17	$+0.09 \pm 0.10$

\dagger Belle and Babar not quite consistent ($\Rightarrow -0.19 \pm 0.13$) \ddagger Not a CP eigenstate

\Rightarrow How good is \approx ?

Basics

- Formalism:

- Effective \mathcal{H} for $b \rightarrow sq\bar{q}$ decays:

$$\mathcal{H} = \frac{G_F}{\sqrt{2}} \sum_{p=u,c} V_{ps}^* V_{pb} \left(\sum_{i=1}^2 C_i O_i^p + \sum_{i=3}^{10} C_i O_i + C_{7\gamma} O_{7\gamma} + C_{8g} O_{8g} \right)$$

- Decay amplitudes:

$$A_f = \langle f | \mathcal{H} | B^0 \rangle, \quad \bar{A}_f = \langle f | \mathcal{H} | \bar{B}^0 \rangle, \quad \lambda_f = e^{-i\phi_B} (\bar{A}_f / A_f)$$

- CP asymmetries:

$$S_f = 2\mathcal{I}m(\lambda_f) / (1 + |\lambda_f|^2), \quad C_f = (1 - |\lambda_f|^2) / (1 + |\lambda_f|^2)$$

- SM:

- $A_f^{\text{SM}} = A_f^c + A_f^u$ with $A_f^c \propto V_{cb}^* V_{cs}$ and $A_f^u \propto V_{ub}^* V_{us}$

- $A_f^{\text{SM}} = A_f^c (1 + a_f^u e^{i\gamma})$

The Problem

- Simple NP:

- $C_i(m_W) = C_i^{\text{SM}}(m_W) + x_i \varepsilon e^{i\theta}$, x_i known

- $A_f^{\text{SM}} = A_f^c \left[1 + a_f^u e^{i\gamma} + b_f^c \varepsilon e^{i\theta} \right]$

- CP asymmetries:

- $\Delta S_f = +2 \cos 2\beta_{\text{eff}} \left[\mathcal{R}e(b_f^c) \varepsilon \sin \theta + \mathcal{R}e(a_f^u) \sin \gamma \right]$

- $C_f = -2\mathcal{I}m(b_f^c) \varepsilon \sin \theta - 2\mathcal{I}m(a_f^u) \sin \gamma$

- Problem:

- To be convinced that $\varepsilon \neq 0$, we need to know a_f^u

- $a_f^u = |(V_{ub}V_{us})/(V_{cb}V_{cs})| \times \text{hadronic parameters}$

Factorization-related methods

Example: $B \rightarrow K^0 \pi^0$:

$$A_{K^0 \pi^0}^c \approx iV_{cb}^* V_{cs} \frac{G_F}{2} f_K F^{B \rightarrow \pi}(m_K^2)(m_B^2 - m_\pi^2)(a_4 + r_\chi a_6),$$

$$A_{K^0 \pi^0}^u \approx iV_{ub}^* V_{us} \frac{G_F}{2} [f_K F^{B \rightarrow \pi}(m_K^2)(m_B^2 - m_\pi^2)(a_4 + r_\chi a_6) \\ - f_\pi F^{B \rightarrow K}(m_\pi^2)(m_B^2 - m_K^2)a_2]$$

where $r_\chi = 2m_K^2/[m_b(m_s + m_d)]$, $a_i \equiv C_i + C_{i\pm 1}/N_c$



$$a_{K\pi}^u \approx \lambda^2 R_u \left(1 - \frac{f_\pi}{f_K} \frac{F^{B \rightarrow K}}{F^{B \rightarrow \pi}} \frac{a_2}{a_4 + r_\chi a_6} \right) \approx 2.75 \lambda^2 R_u \approx 0.052$$

$$\implies \Delta S_{\pi K_S} \approx 0.06, \quad C_{\pi K_S} \approx 0$$

$\Delta S_{\pi^0 K_S}$ in Factorization-related methods

$\Delta S_{\pi^0 K_S}$	Method	hep-ph/	Authors
$+0.06 \pm 0.04$	NF	0503151	Buchalla, Hiller, Nir, Raz
$+0.04 \pm 0.03$	NF+model	0502235	Cheng, Chua, Soni
$+0.07 \pm 0.04$	QCDF	0505075	Beneke
$+0.06 \pm 0.03$	PQCD	0508041	Li, Mishima, Sanda
$+0.08 \pm 0.16$	SCET+SU(3)	0510241	Bauer, Rothstein, Stewart

SU(3)-based methods

- For $b \rightarrow s$, define: $d_f^q = A_f^q / (V_{qb}^* V_{qs})$
 $A_f = V_{cb}^* V_{cs} d_f^c + V_{ub} V_{us}^* d_f^u$
- For $b \rightarrow d$, define: $h_{f'}^q = A_{f'}^q / (V_{qb}^* V_{qd})$
 $A_{f'} = V_{cb}^* V_{cd} h_{f'}^c + V_{ub} V_{ud}^* h_{f'}^u$

The approximate SU(3) symmetry of the strong interactions

$$\implies d_f^q = \sum_{f'} X_{f'} h_{f'}^q$$

where $X_{f'}$ are known (CG) coefficients

$$\boxed{|a_f^u| = \frac{|V_{ub} V_{us} d_f^u|}{|V_{cb} V_{cs} d_f^c|} \leq \left| \frac{V_{us}}{V_{ud}} \right| \frac{\sum_{f'} |X_{f'}| \sqrt{\mathcal{B}(B \rightarrow f')}}{\sqrt{\mathcal{B}(B \rightarrow f)}}$$

SU(3)-based methods

⊖ SU(3) breaking effects of $\mathcal{O}(0.3)$

⇒ The bounds are only approximate

⊖ Adding (conservatively) the amplitudes coherently

+ Dependence on measured \mathcal{B} 's

⇒ Bounds often much weaker than actual estimates

⊕ Hadronic model independence

⇒ Complimentary to the factorization-related methods

SU(3)-based methods

Example: $B \rightarrow K^0 \pi^0$:

$$d_{B_d^0 \rightarrow K^0 \pi^0}^q = \frac{1}{\sqrt{2}} h_{B_d^0 \rightarrow K^+ K^-}^q - h_{B_d^0 \rightarrow \pi^0 \pi^0}^q$$

$$|a_{K^0 \pi^0}^u| \leq \left| \frac{V_{us}}{V_{ud}} \right| \frac{\frac{1}{\sqrt{2}} \sqrt{\mathcal{B}(B \rightarrow K^+ K^-)} + \sqrt{\mathcal{B}(B \rightarrow \pi^0 \pi^0)}}{\sqrt{\mathcal{B}(B \rightarrow K^0 \pi^0)}}$$



$$|a_{K\pi}^u| \leq 0.15$$

$$\implies [\Delta S_{\pi K_S} / \cos 2\beta]^2 + C_{\pi K_S}^2 \leq [0.30 \sin \gamma]^2$$

hep-ph/0310020 (Gronau, Grossman, Rosner), hep-ph/0509125 (Raz)

Estimating $\Delta S_f \equiv -\eta_f S_f - S_{\psi K_S}$

f_{CP}	EXP	NF*	QCDF**	SU(3)***
ϕK_S	-0.22 ± 0.19	$+0.02 \pm 0.01$	$+0.02 \pm 0.01$	$(K^* K)^\dagger$
$\eta' K_S$	-0.19 ± 0.09	$+0.01 \pm 0.02$	$+0.01 \pm 0.01$	0.31
$f_0 K_S$	$+0.06 \pm 0.24$			
$\pi^0 K_S$	-0.38 ± 0.26	$+0.06 \pm 0.04$	$+0.07 \pm 0.04$	0.18
ωK_S	-0.06 ± 0.30	$+0.19_{-0.14}^{+0.06}$	0.13 ± 0.08	$(K^* K)^\dagger$
$K_S K_S K_S$	-0.08 ± 0.23			1(0.37)
$K^+ K^- K_S^\dagger$	-0.16 ± 0.17			0.94

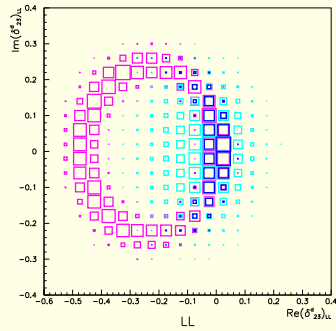
* Buchalla, Hiller, Nir, Raz, hep-ph/0503151

** Beneke, hep-ph/0505075

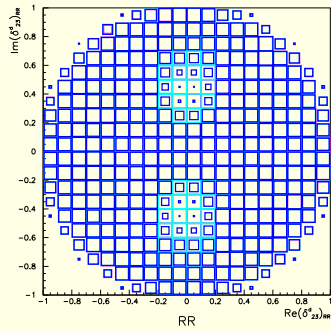
*** Grossman, Ligeti, Quinn, Nir, hep-ph/0303171; Raz *et al.*, hep-ph/0505195,0508046,0509125

† Available once $\mathcal{B}(K^* K)$ is measured

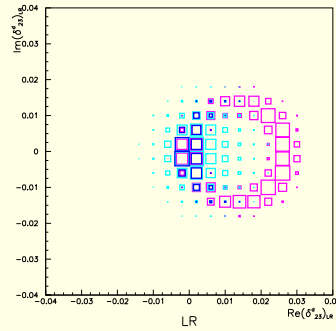
Supersymmetry: constraints and predictions



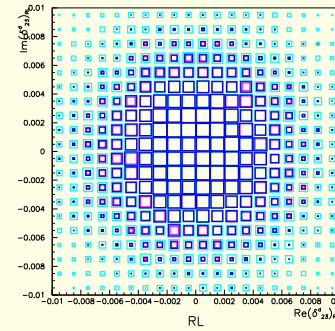
$(\delta_{23}^d)_{LL}$



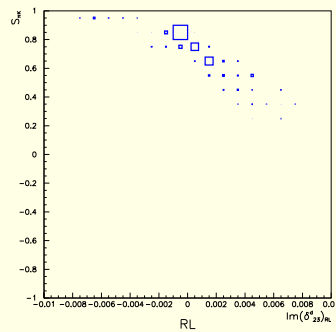
$(\delta_{23}^d)_{RR}$



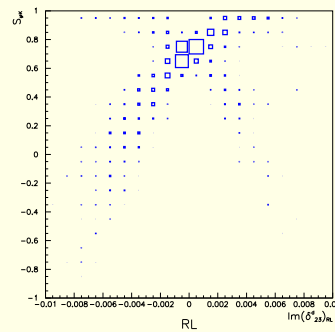
$(\delta_{23}^d)_{LR}$



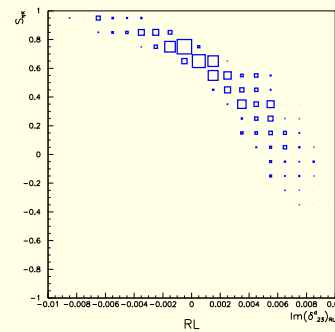
$(\delta_{23}^d)_{RL}$



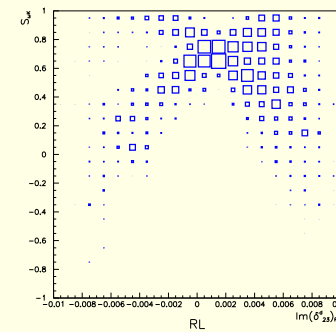
$S_{\pi K}$



$S_{\phi K}$



$S_{\eta' K}$



$S_{\omega K}$

Silvestrini, hep-ph/0510077

Supersymmetry: constraints and predictions

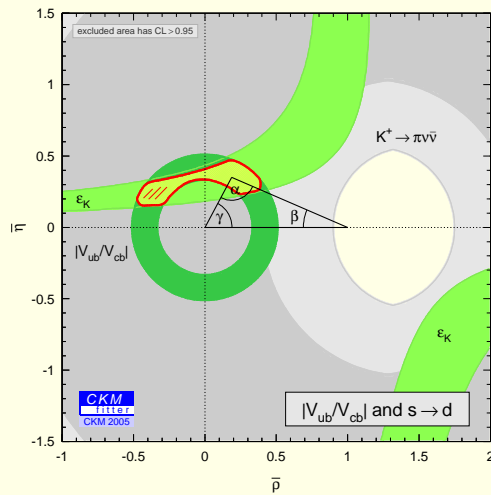
$(\delta_{23}^d)_{MN}$	Upper Bound*	$ \Delta S_{\phi K_S} \sim 0.1^{**}$	Alignment***
$(\delta_{23}^d)_{LL}$	$\lambda^2(\mathcal{R}e) - \lambda(\mathcal{I}m)$	λ	λ^2
$(\delta_{23}^d)_{RR}$	1	λ	$\lambda^4 - \lambda^2$
$(\delta_{23}^d)_{LR}$	$\lambda^4(\mathcal{R}e) - \lambda^3(\mathcal{I}m)$	λ^4	$\lambda^2(m_b/\tilde{m})$
$(\delta_{23}^d)_{RL}$	λ^3	λ^4	$\lambda^4(m_b/\tilde{m})$

* Ciuchini *et al.*, hep-ph/0407073

** Silvestrini, hep-ph/0510077

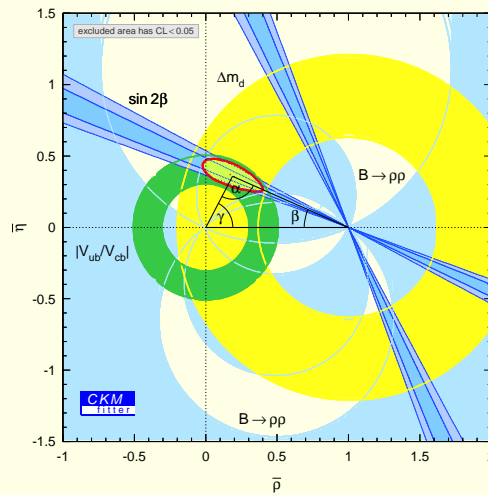
* * * Nir, Raz, hep-ph/0206064

Unitarity Triangles 2005



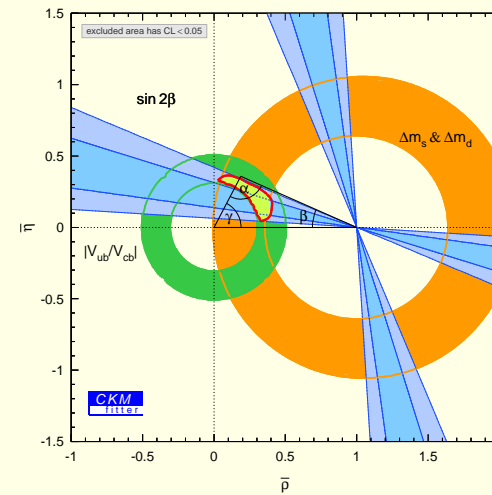
$s \rightarrow d$

$\varepsilon, \mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$



$b \rightarrow d$

$\Delta m_{B_d}, S_{\psi K}, S_{\rho\rho}$



$b \rightarrow s$

$\Delta m_{B_s}, S_{b \rightarrow s \bar{s}}$

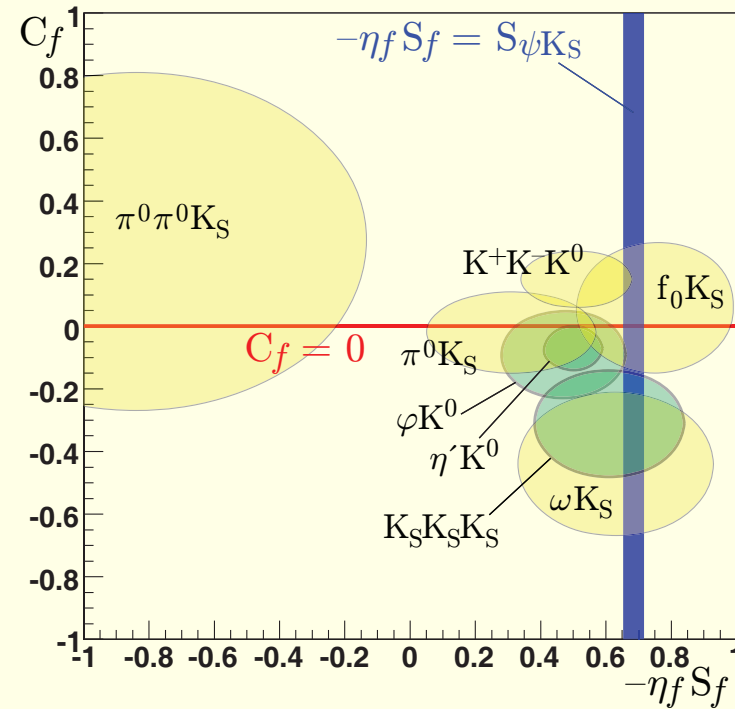
There is still a lot to be learned from future measurements

Conclusions

- The KM phase is different from zero (SM violates CP)
- The KM mechanism is, very likely, the dominant source of the CP violation observed in meson decays
- The size and the phase of NP contributions to $B^0 - \bar{B}^0$ mixing are severely constrained
- Complete alternatives to the KM mechanism are excluded (Superweak, Approximate CP)
- Corrections to KM are possible, particularly for $b \rightarrow s$; No evidence for such corrections at present
- There is still a lot to be learned from flavor/CP physics

$$S_{\phi K_S}, S_{\eta' K_S}, S_{\pi^0 K_S} \dots$$

Is there NP in $b \rightarrow s$ transitions?



Kirkby and Nir, PDG

No evidence at present

Experimental status of CP asymmetries

f_{CP}	$-\eta_{\text{CP}}S$	C
$\psi\pi^0$	$+0.69 \pm 0.25$	-0.11 ± 0.20
D^+D^-	$+0.29 \pm 0.63$	$+0.11 \pm 0.35$
$D^{*+}D^{*-}$	$+0.75 \pm 0.23$	-0.04 ± 0.14
$\pi^+\pi^-$	$+0.50 \pm 0.12(0.18)$	$-0.37 \pm 0.10(0.23)$
$\pi^0\pi^0$		-0.28 ± 0.39
$\rho^+\rho^-$	$+0.22 \pm 0.22$	-0.02 ± 0.17

The NP CP/Flavor Problem

- $m_H^2 \sim (m_H^2)_{\text{tree}} + \frac{1}{16\pi^2} \Lambda_{\text{NP}}^2$

To avoid fine-tuning of the Higgs mass,

$$\Lambda_{\text{NP}} \lesssim 4\pi m_W \sim 1 \text{ TeV}.$$

- $\mathcal{L}_{\text{NP}} \sim \frac{1}{\Lambda_{\text{NP}}^2} s\bar{d}s\bar{d}$

To avoid too large contributions to ε_K and to $\Delta m_{K,D,B}$,

$$\Lambda_{\text{NP}} \gtrsim 10^{3-4} \text{ TeV}.$$

New Physics at the TeV scale must have a very non-generic flavor and CP structure

SU(3) Relations

$$A_f = V_{cb}^* V_{cs} a_f^c + V_{ub}^* V_{us} a_f^u, \quad \xi_f \equiv |V_{ub} V_{us} / V_{cb} V_{cs}| (a_f^c / a_f^u)$$

$$-\eta_f S_f - S_{\psi K} = 2 \cos 2\phi_1 \sin \phi_3 \operatorname{Re}(\xi_f)$$

$$C_f = -2 \sin \phi_3 \operatorname{Im}(\xi_f)$$

Grossman, Ligeti, Nir, Quinn (03)
Engelhard, Nir, Raz (05)

- Example: $|\xi_{\eta' K_S}| \leq \sqrt{\frac{3\mathcal{B}(\eta'\eta)}{2\mathcal{B}(\eta' K_S)}} + \sqrt{\frac{\mathcal{B}(\eta'\pi^0)}{2\mathcal{B}(\eta' K_S)}}$

mode	$\eta' K_S$	$\pi^0 K_S$	$K^- \pi^+$	$\eta' K^+$	ϕK^+	$K_S K_S K_S$
$ \xi <$	0.25	0.18	0.23	0.07	0.22	0.31 [†]

† Extra (mild) dynamical assumptions