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# **Detectors and electronics for Super-LHC: IRRADIATION RESULTS**

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# OUTLINE

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The presented results are from experiments performed by:

- SMART** Collaboration (Bari, Firenze, Pisa, Padova, Perugia, Trieste, Trento)  
(V. Khomenkov, A. Litovchenko, A. Candelori)
- Department of Information Engineering (**DEI**), University of Padova  
(A. Paccagnella, A. Cester, S. Gerardin)

**Measurements are in progress ... some new results after the Trento RD50 Meeting**

## **1) Nuclear reactor irradiation in Lubjana** (V. Cindro):

- Epitaxial** n-type detectors (**SMART**)
- Pre-irradiated** n-type detectors (**Kiev, IRST, Padova**)

## **2) 24 GeV proton irradiation at CERN** (M. Glaser, M. Moll, F. Ravotti):

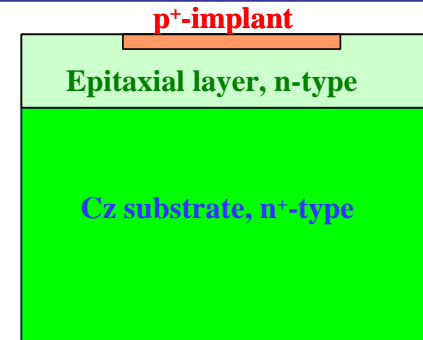
- Epitaxial** n-type detectors (**SMART**)
  - FZ thinned** n-type detectors (**SMART**)
  - **0.13  $\mu\text{m}$  CMOS** technology (**DEI**)
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# DETECTORS

## -**Epitaxial** detectors:

**Epitaxial layer:** n-type, 50  $\mu\text{m}$  thick,  $\rho = 50 \Omega\cdot\text{cm}$ ,  
grown by ITME on CZ substrate

**Detector processing:** IRST.



## -**Pre-irradiated** detectors:

**Substrate:** FZ, n-type, 300  $\mu\text{m}$  thick,  $\rho = 3\text{-}4 \text{ k}\Omega\cdot\text{cm}$ .

**Pre-irradiation:** irradiation by fast nuclear reactor neutrons up to  $\approx 10^{16} \text{ n/cm}^2$ ;  
maximum annealing temperature: 850  $^{\circ}\text{C}$ ;  
total annealing time with  $\text{RT} \leq T \leq 850 \text{ }^{\circ}\text{C}$ : 2 hours.

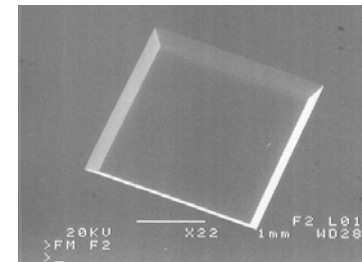
**Detector processing:** IRST.

## -**FZ thinned** detectors:

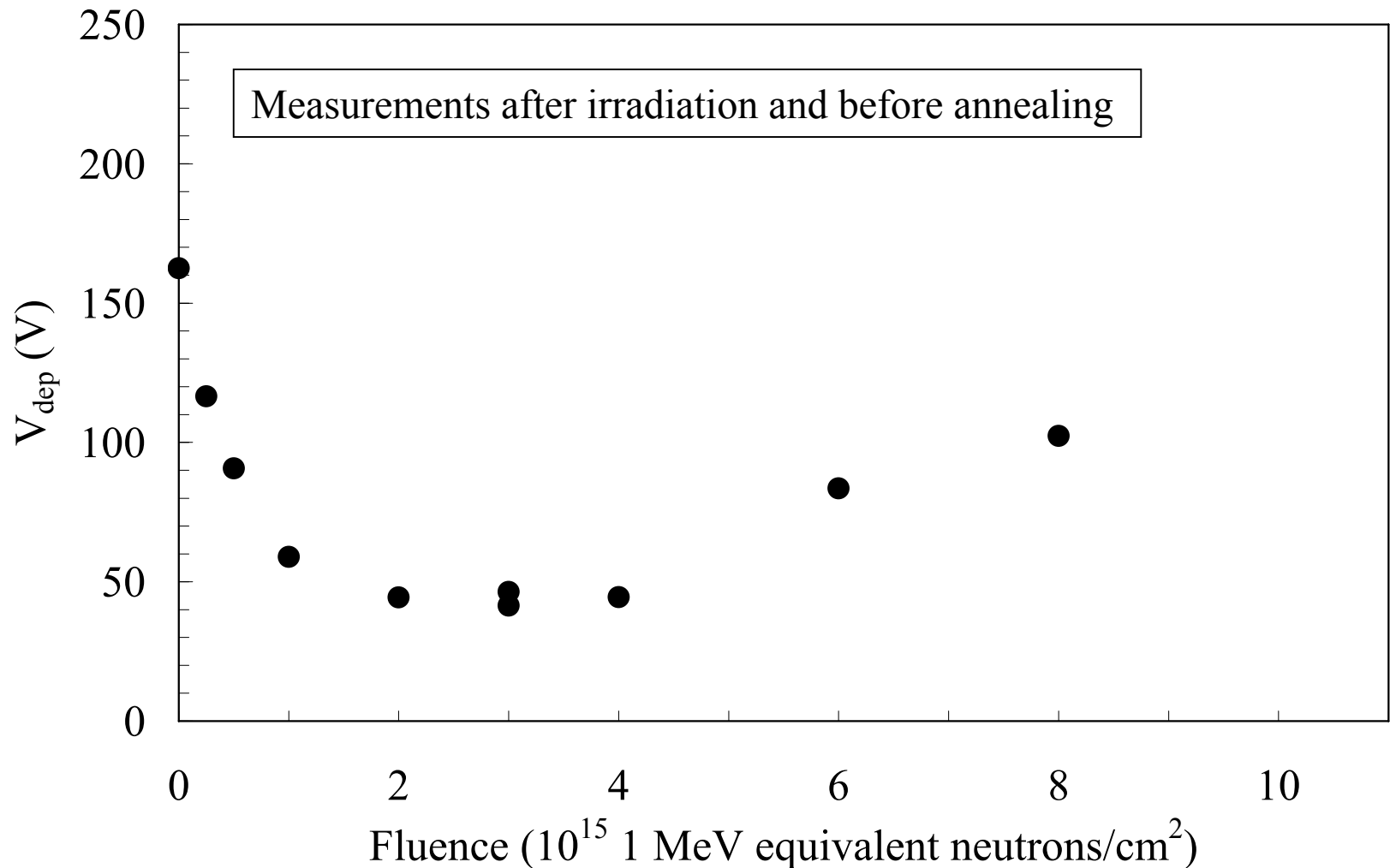
**Substrate:** FZ, n-type, 300  $\mu\text{m}$  thick,  $\rho = 6 \text{ k}\Omega\cdot\text{cm}$ .

**Thinning:** down to 50-100  $\mu\text{m}$  by Tetra Methyl Ammonium Hydroxide (TMAH) etching from backside: IRST.

**Detector processing:** IRST.

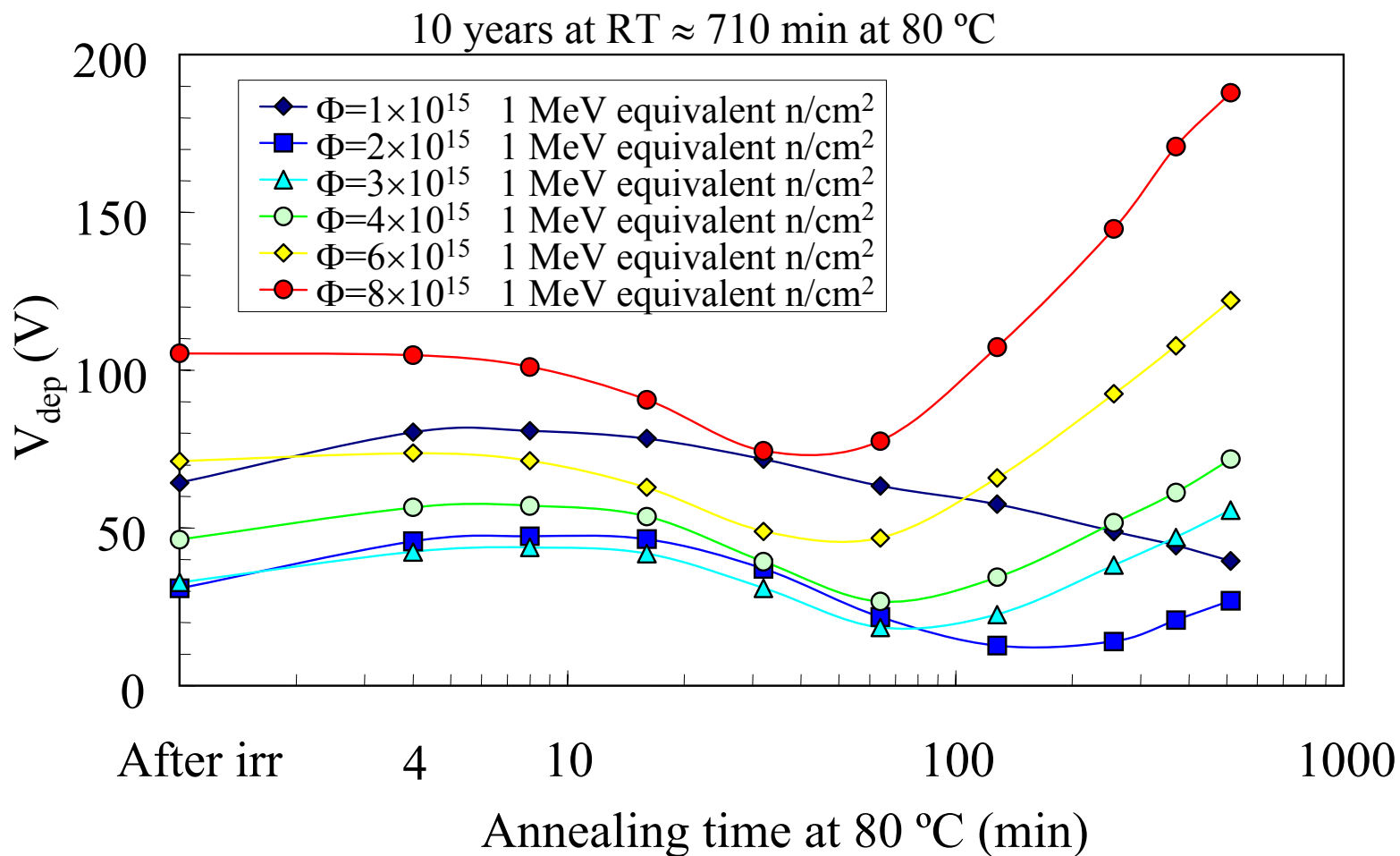


# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: neutron irradiation



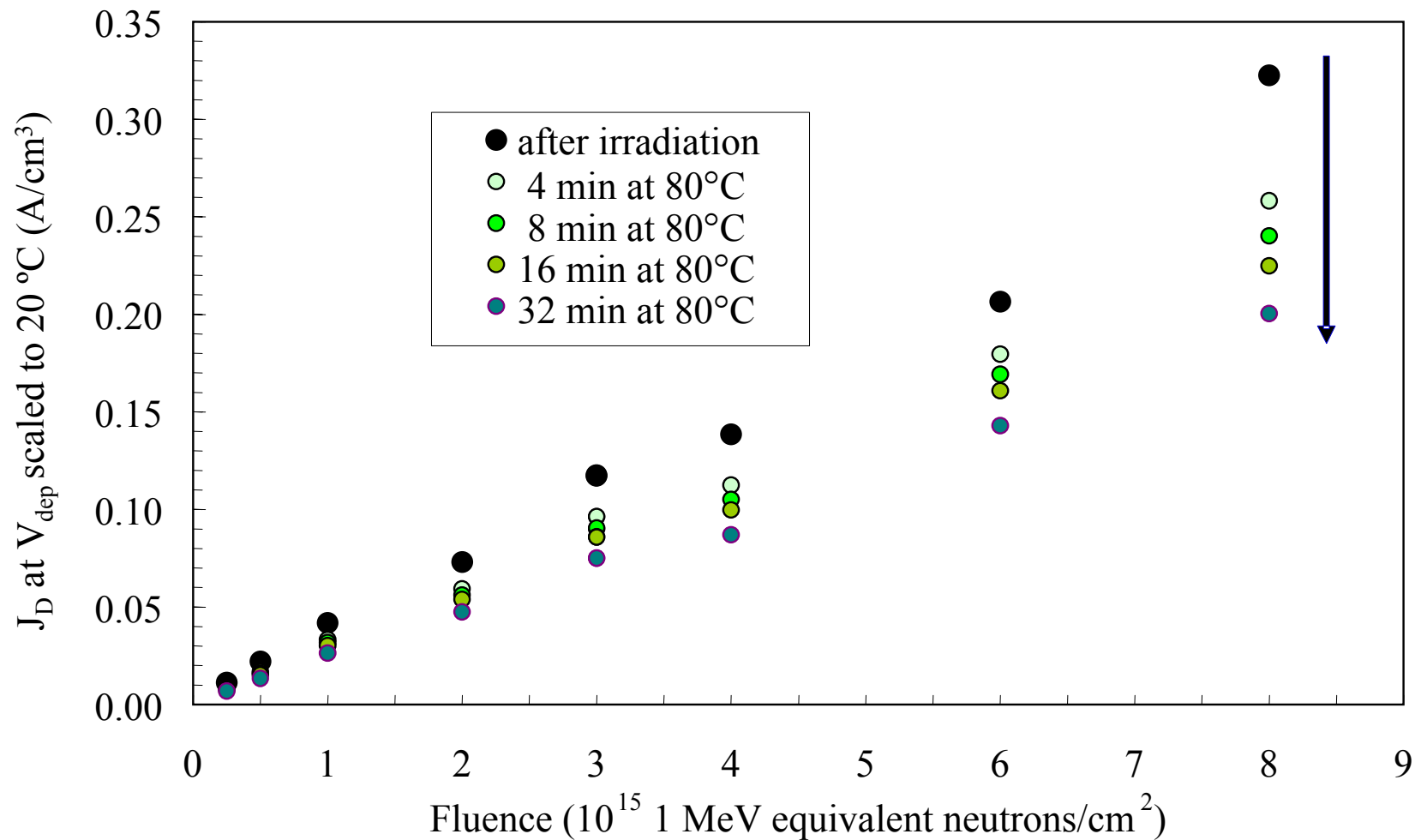
- The minimum of  $V_{\text{dep}} \approx 40\text{-}50$  V is reached at  $2\text{-}4 \times 10^{15}$  n/cm<sup>2</sup>.
- $V_{\text{dep}} < V_{\text{dep},0}$  at the maximum fluence ( $8 \times 10^{15}$  n/cm<sup>2</sup>).

# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: neutron irradiation



- Initial  $V_{\text{dep}}$  increase as a function of the annealing time: less evidence by increasing  $\Phi$ .
- Low  $\Phi$ .  $V_{\text{dep}}$   $\uparrow$ ,  $\downarrow$ , no SCSI during long-term annealing.
- Medium  $\Phi$ .  $V_{\text{dep}}$   $\uparrow$ ,  $\downarrow$ ,  $\uparrow$ , SCSI during long-term annealing.
- High  $\Phi$ .  $V_{\text{dep}}$   $\uparrow$  less evident,  $\downarrow$ ,  $\uparrow$ : already SCSI after irradiation or SCSI during long-term annealing?

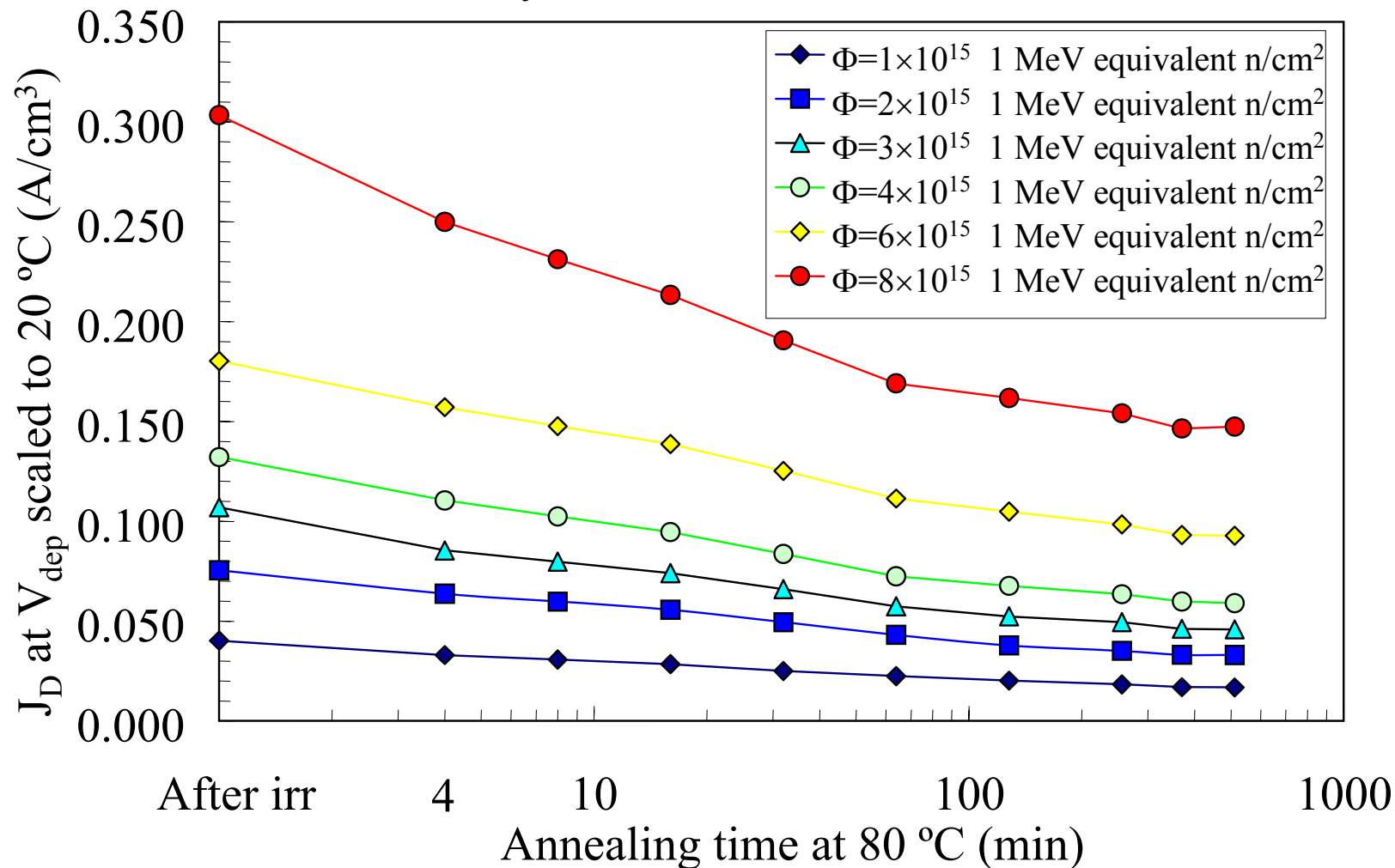
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: neutron irradiation



- $J_D$  linearly increases with fluence.
- $J_D$  decreases with annealing time at 80 °C.

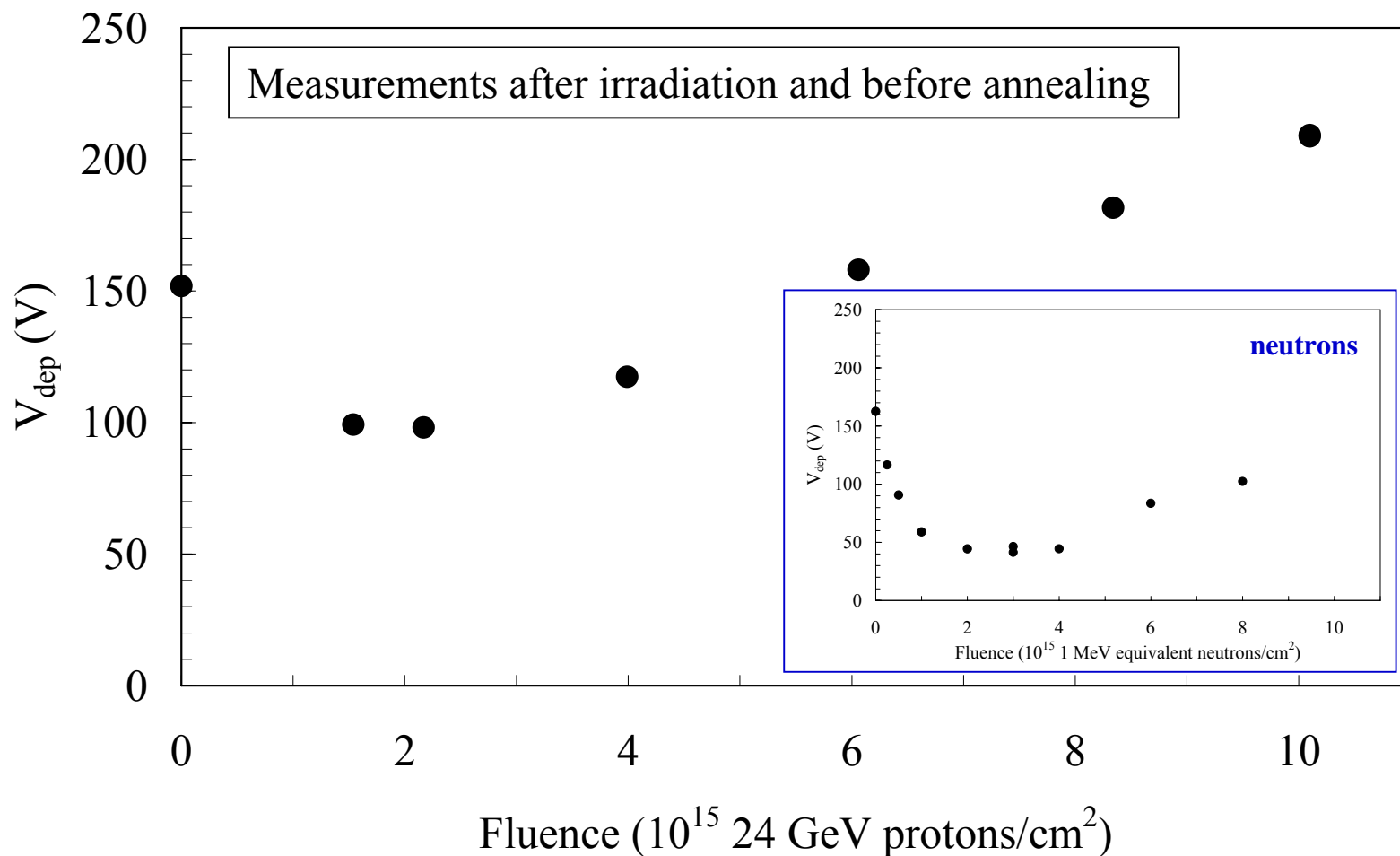
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: neutron irradiation

10 years at RT=710 min at 80 °C



•  $J_D$  decreases with annealing time at 80 °C.

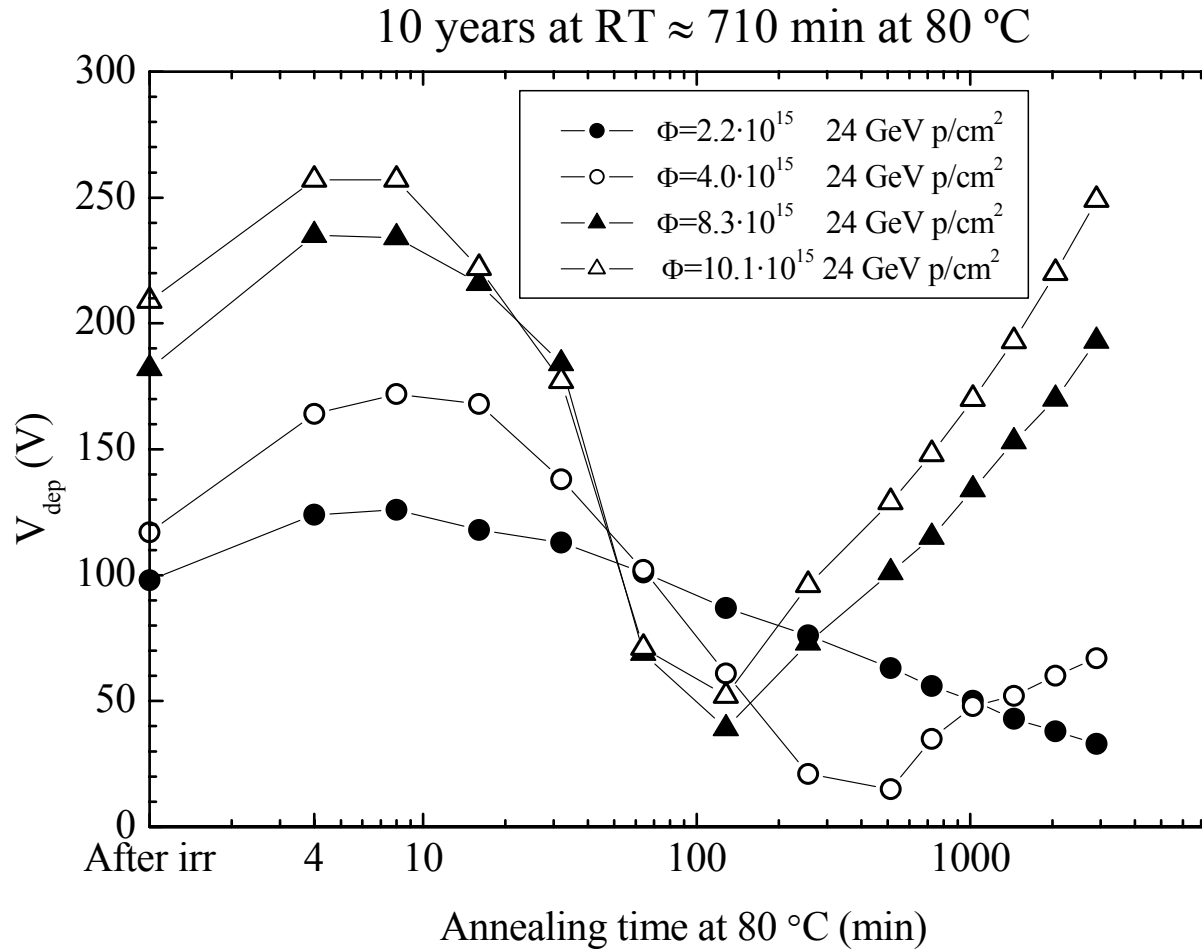
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation



- The minimum of  $V_{\text{dep}}$  for protons (90-100 V) is higher than for neutrons (40-50 V).
- $V_{\text{dep}} > V_{\text{dep},0}$  at the maximum irradiation fluence ( $10^{16}$  p/cm<sup>2</sup>).
- Radiation effects induced by neutrons and protons are different.

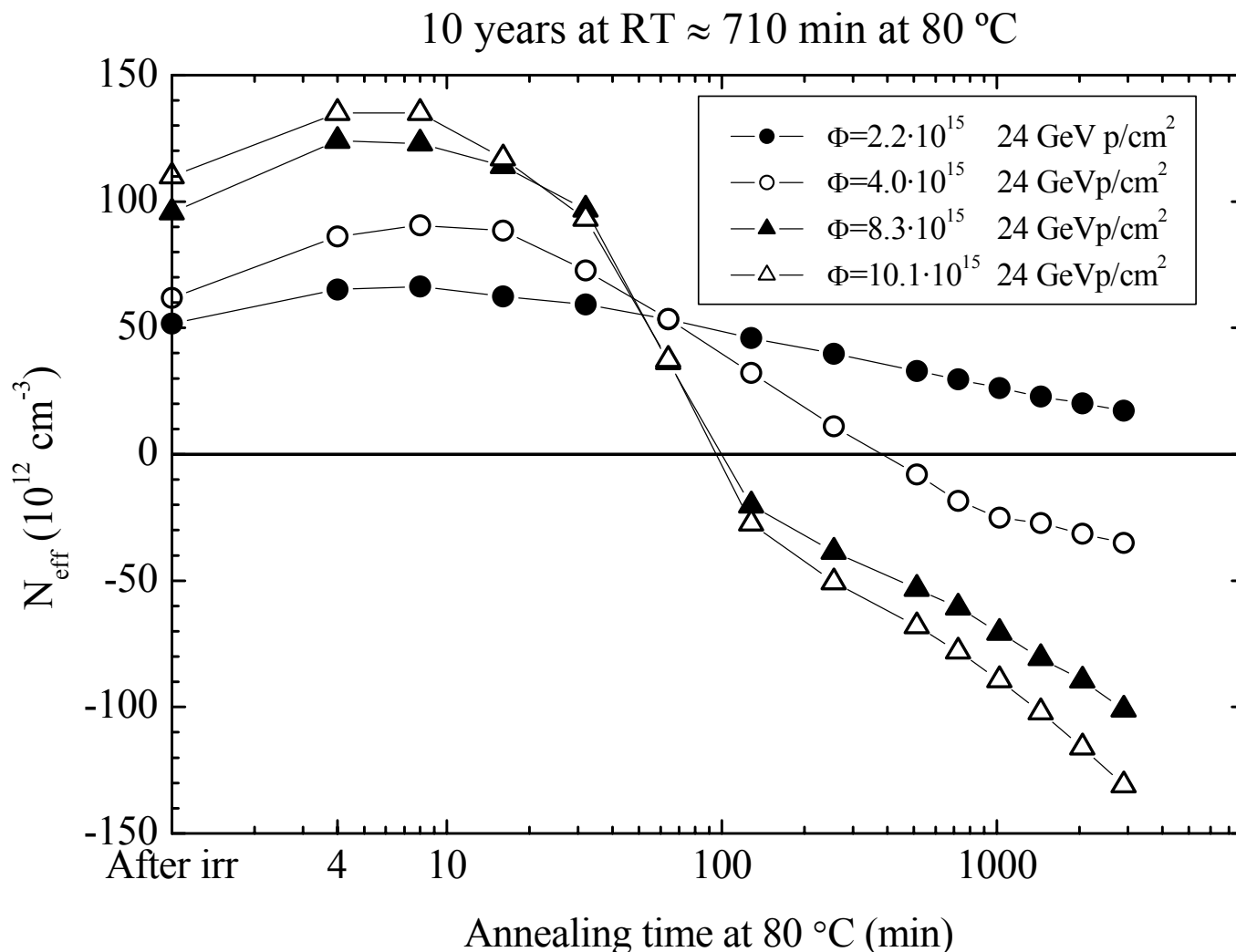


# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation



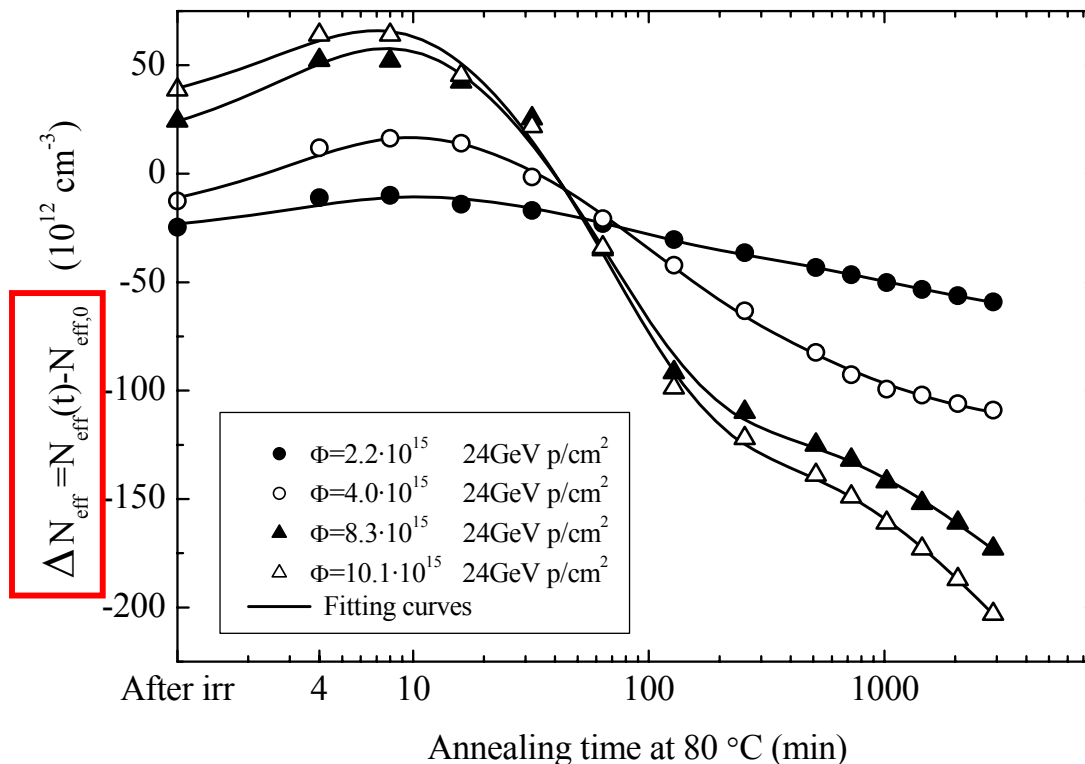
- $V_{\text{dep}}$  decreases for annealing times at 80 °C higher than 8 minutes.
- If this effect is due to deep acceptor generation, devices are before SCSI after irradiation.
- For devices irradiated at **low fluences**  $V_{\text{dep}} \uparrow, \downarrow$ : **no SCSI during long-term annealing**;
- For devices irradiated at **high fluences**  $V_{\text{dep}} \uparrow, \downarrow, \uparrow$ : **SCSI during long-term annealing**.
- **Radiation effects induced by neutrons and protons are different.**

# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation



$N_{\text{eff}}$  as a function of the annealing time, assuming that the  $V_{\text{dep}}$  increase, at high annealing times and for devices irradiated at high fluences, is due to SCSI.

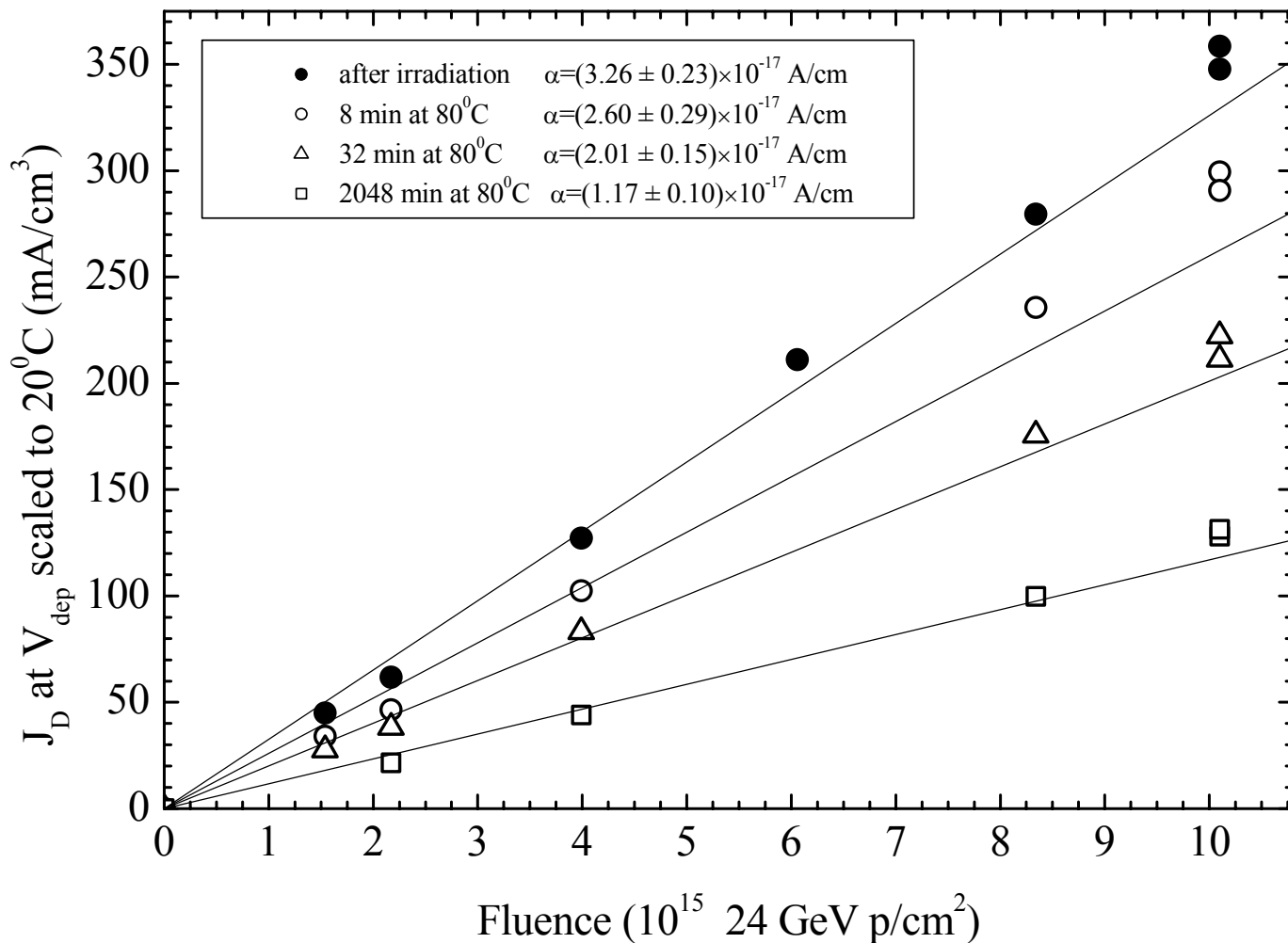
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation



$$\Delta N_{\text{eff}} = N_{\text{eff}}(t) - N_{\text{eff},0} = -N_0 - N_1 \times \exp(-t/\tau_1) - N_2 \times (1 - \exp(-t/\tau_2)) - N_3 \times (1 - 1/(1 + t/\tau_3))$$

	$\Phi$ ( $10^{15}$ 24 GeV p/cm <sup>2</sup> )			
	2.17	3.99	8.34	10.1
$N_0$ ( $10^{13} \text{ cm}^{-3}$ )	$0.4 \pm 0.4$	$-3.3 \pm 0.5$	$-9.5 \pm 0.6$	$-10.3 \pm 0.6$
$N_1$ ( $10^{13} \text{ cm}^{-3}$ )	$1.9 \pm 0.6$	$4.4 \pm 0.6$	$7.2 \pm 0.7$	$6.5 \pm 0.7$
$\tau_1$ (min)		$5.0 \pm 0.7$		
$N_2$ ( $10^{13} \text{ cm}^{-3}$ )	$2.6 \pm 0.9$	$5.7 \pm 2.1$	$20.2 \pm 0.8$	$22.0 \pm 0.7$
$\tau_2$ (min)		$65 \pm 4$		
$N_3$ ( $10^{13} \text{ cm}^{-3}$ )	$3.9 \pm 0.8$	$9.5 \pm 1.5$	$13.8 \pm 3.9$	$18.9 \pm 3.6$
$\tau_3$ ( $10^3$ min)	$1.0 \pm 1.0$	$0.3 \pm 0.2$	$3.1 \pm 1.9$	$2.9 \pm 1.2$

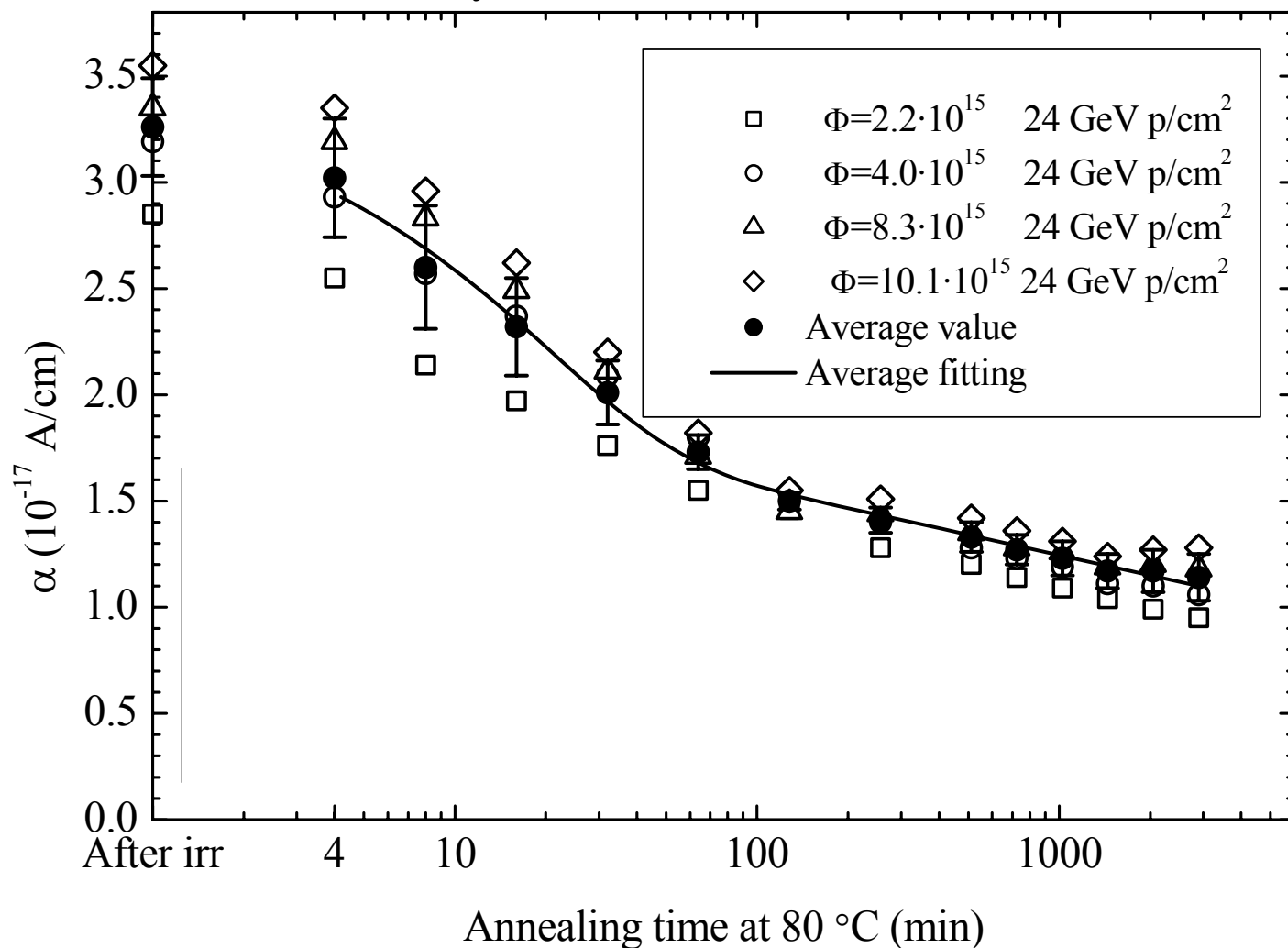
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation



- $J_D$  linearly increases with fluence.
- Average  $\alpha=J_D/\Phi$  value after **8 min at 80 °C** is  **$(2.60 \pm 0.29) \times 10^{-17}$  A/cm**, in agreement within **10%** with the value measured by Hamburg group for epitaxial diodes ( $2.43 \times 10^{-17}$  A/cm).

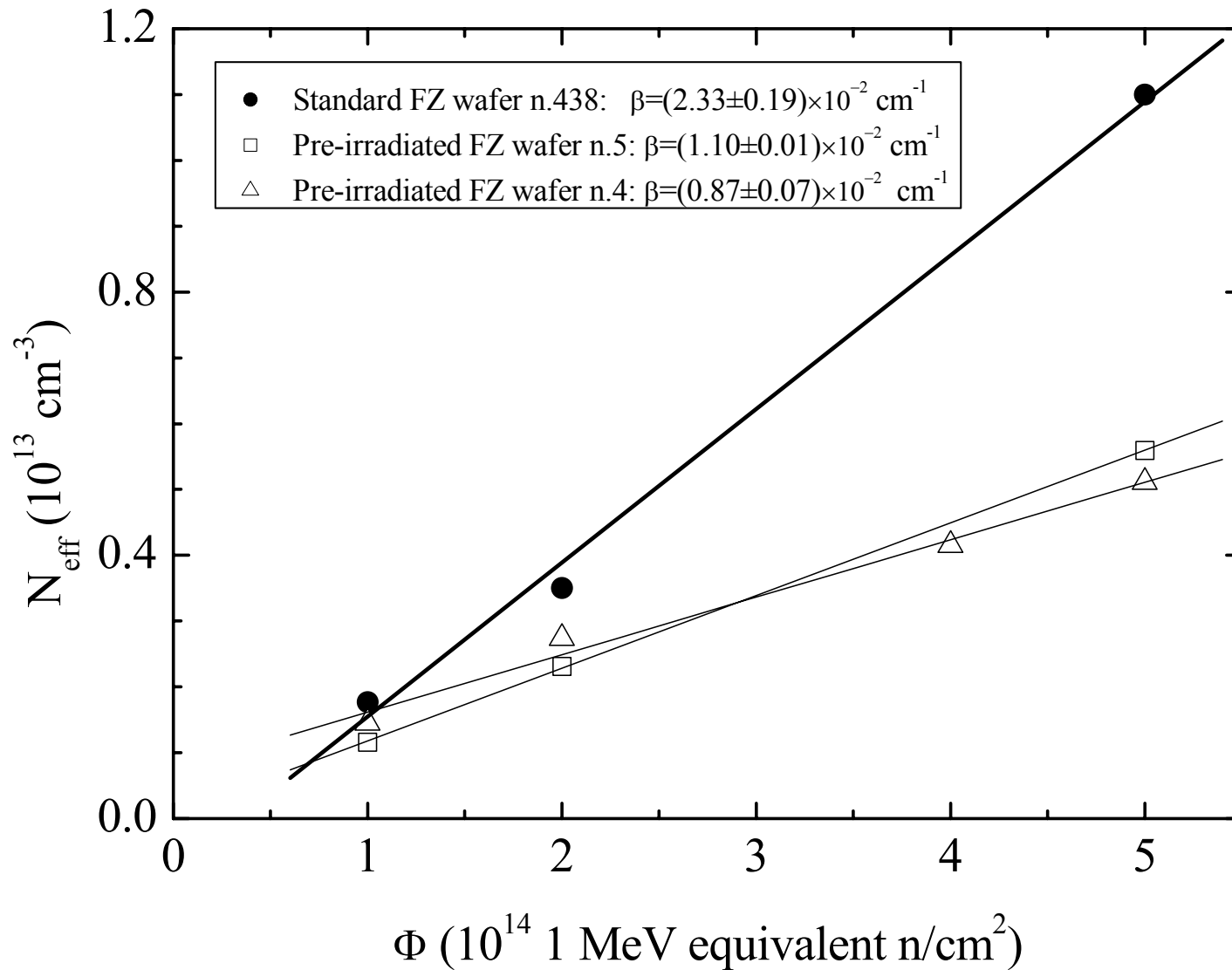
# Epitaxial n-type 50 $\mu\text{m}$ thick detectors: proton irradiation

10 years at RT  $\approx$  710 min at 80  $^{\circ}\text{C}$



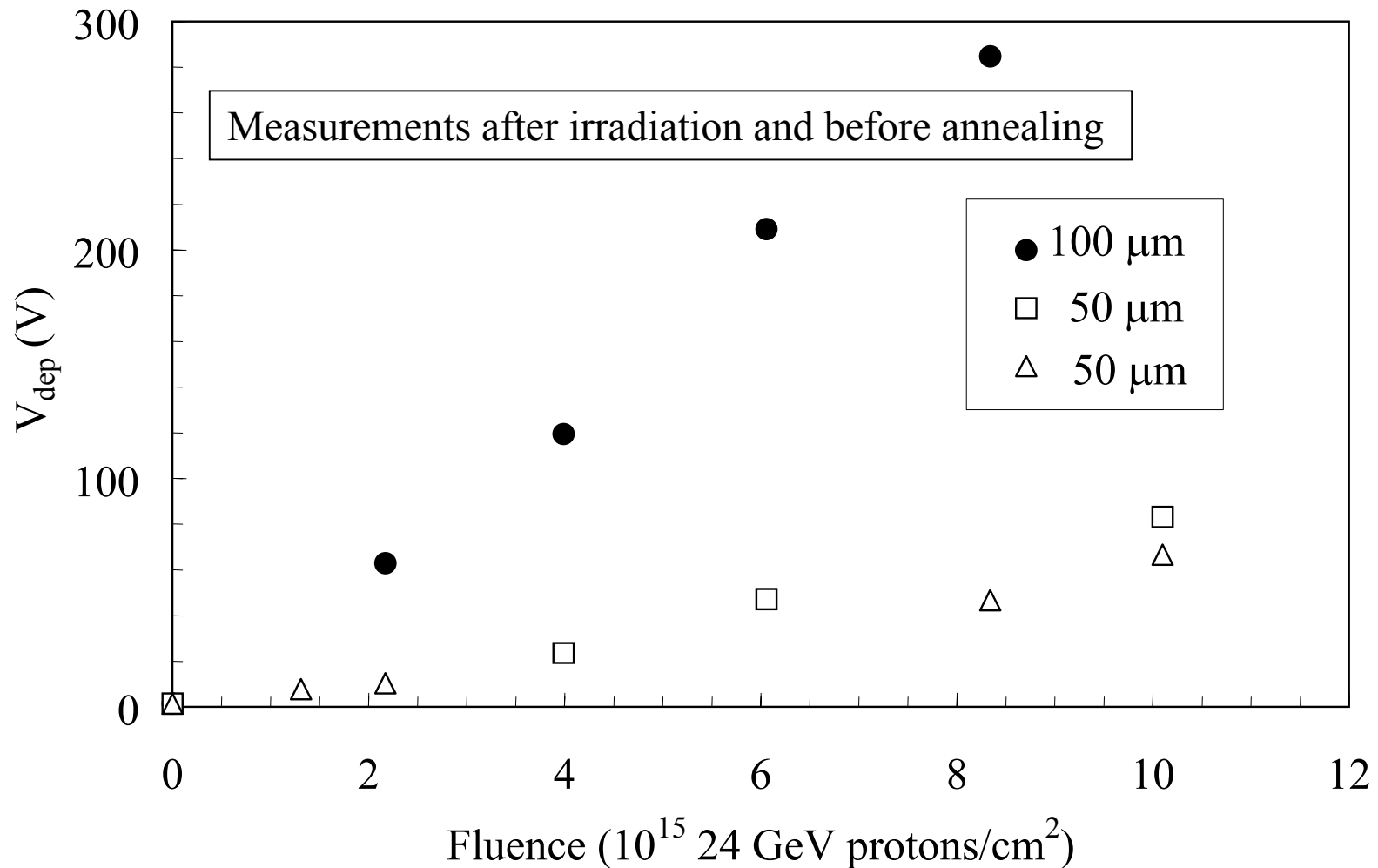
$$\alpha(t) = \alpha_0 + \alpha_1 \times \exp(-t / \tau_1) + \alpha_2 \times \ln(t / \tau_2)$$

# Pre-irradiated n-type detectors: neutron irradiation



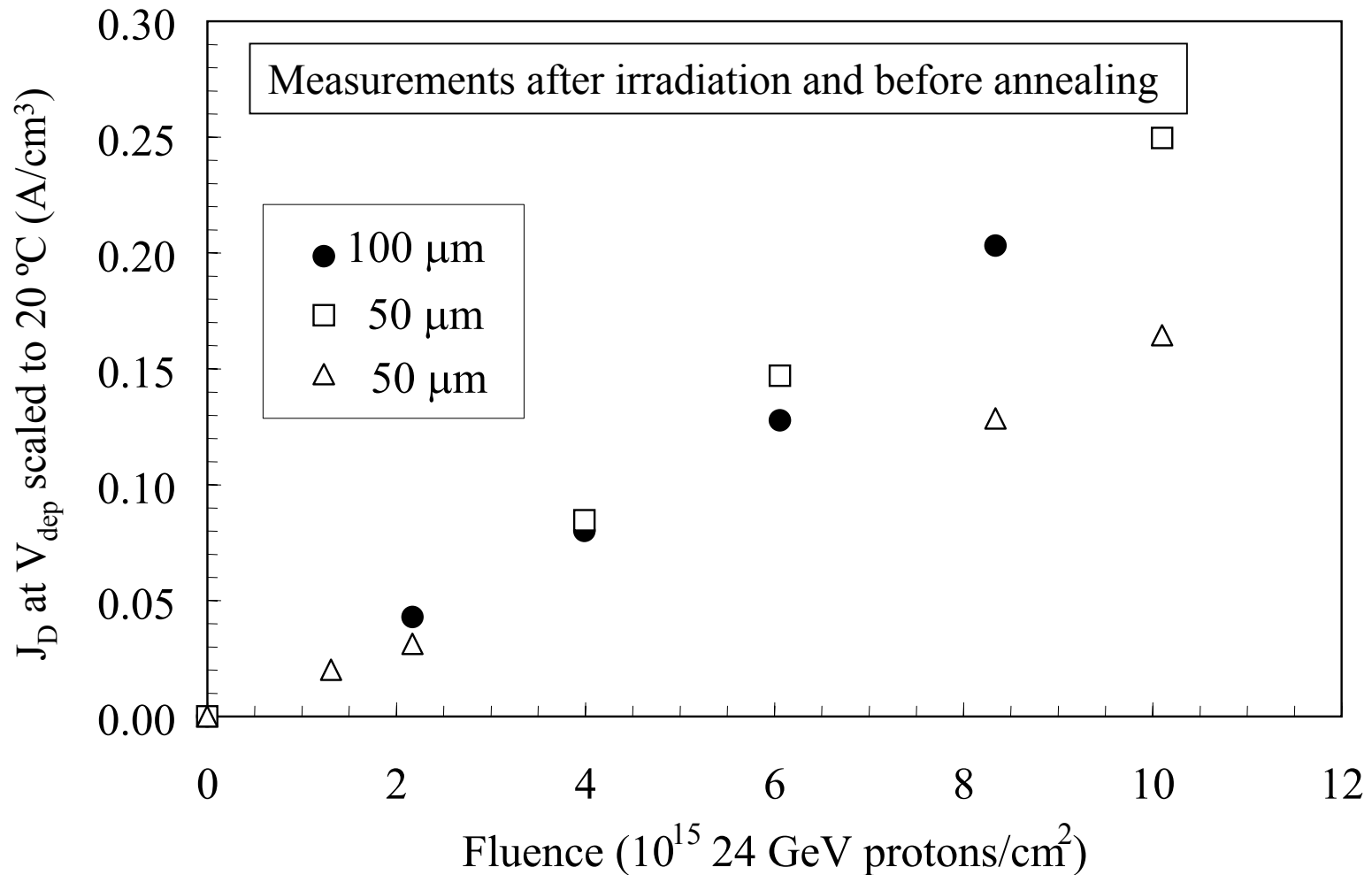
The deep acceptor introduction rate  $\beta$  is lower for pre-irradiated than standard FZ devices.

# FZ thinned n-type detectors: proton irradiation



Device thinning limits the  $V_{\text{dep}}$  increase after irradiation:  
-  $V_{\text{dep}} \approx 290$  V for 100  $\mu\text{m}$  thick detectors at  $8 \times 10^{15}$  p/cm<sup>2</sup>;  
-  $V_{\text{dep}} \approx 70$  V for 50  $\mu\text{m}$  thick detectors at  $10^{16}$  p/cm<sup>2</sup>.

# FZ thinned n-type detectors: proton irradiation



- $J_D$  linearly increases with fluence, independently on the thickness (some data dispersion is present).

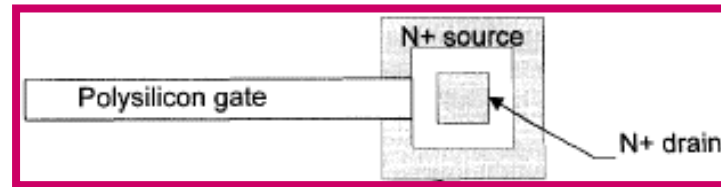


# ELECTRONICS

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-**Electronics for LHC** (10 Mrad(Si) and  $10^{15}$  fast hadrons/cm<sup>2</sup>):

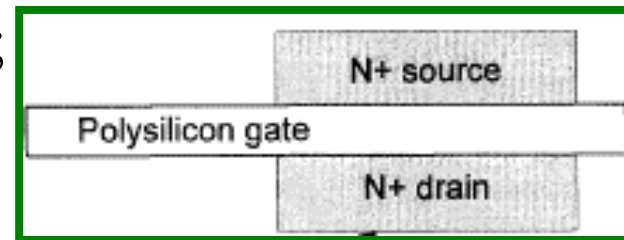
- **0.25  $\mu\text{m}$  CMOS** technology from **IBM**;
- **5.5 nm** oxide thickness;
- radiation hardened by design: **enclosed geometry** transistors.



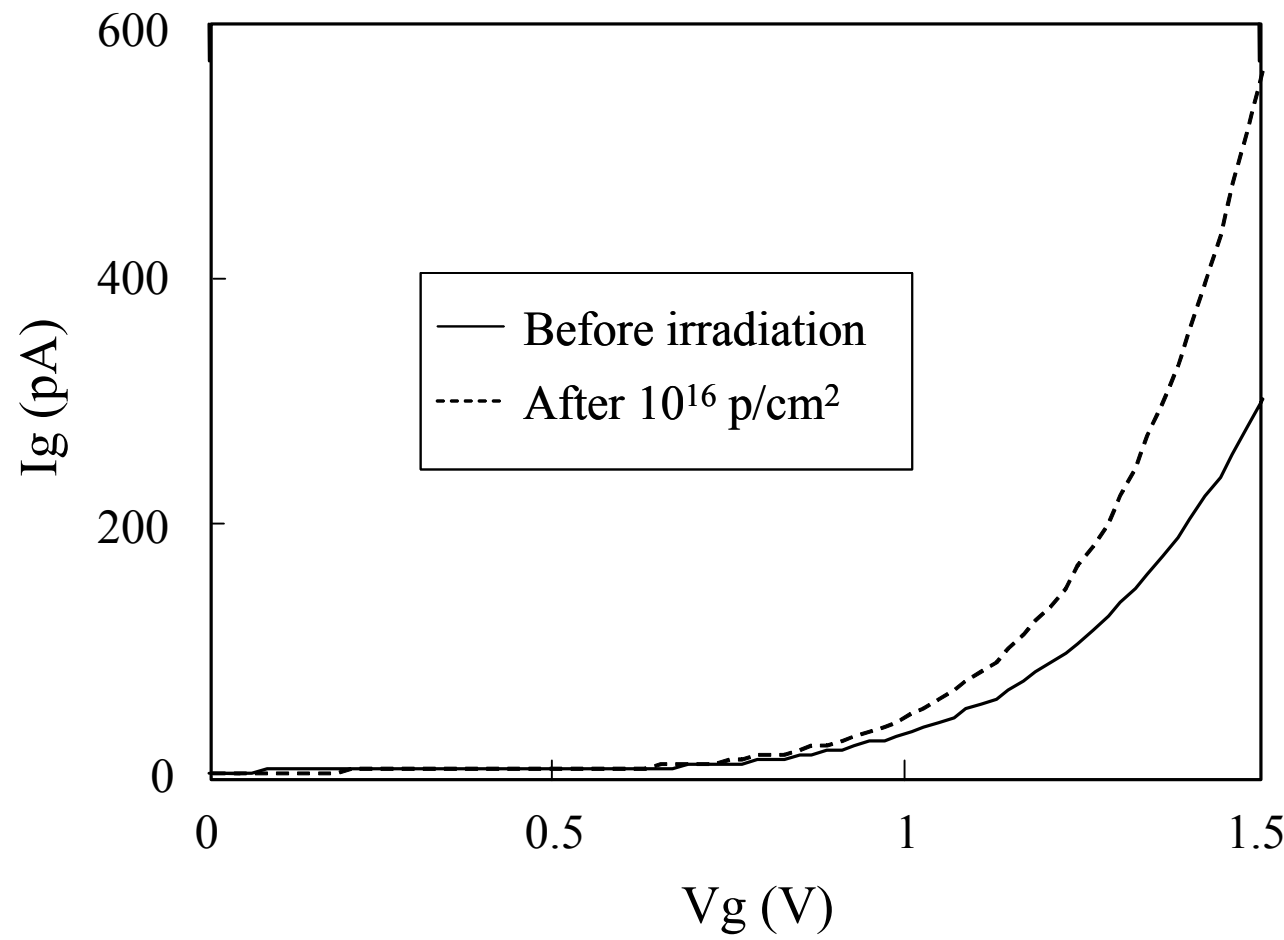
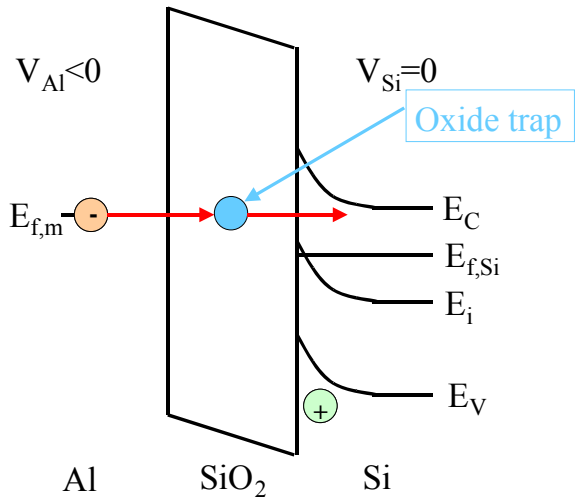
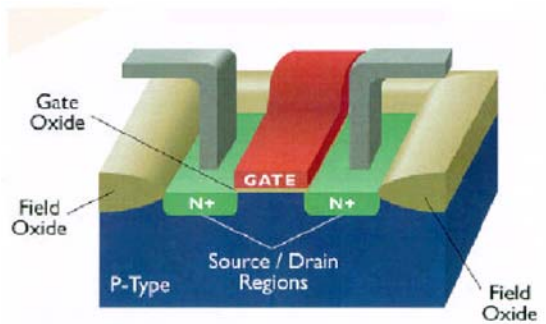
**The 0.25  $\mu\text{m}$  CMOS technology will not be probably any more available for Super-LHC**

-**Electronics for Super-LHC** (100 Mrad(Si) and  $10^{16}$  fast hadrons/cm<sup>2</sup>):  
following the scaling down of the commercial CMOS technologies?

- **0.13  $\mu\text{m}$**  minimum channel length;
- **2.5 nm** oxide thickness;
- no radiation hardened by process;
- no radiation hardened by design (i.e., **no enclosed geometry**);
- from **ST-Microelectronics**.

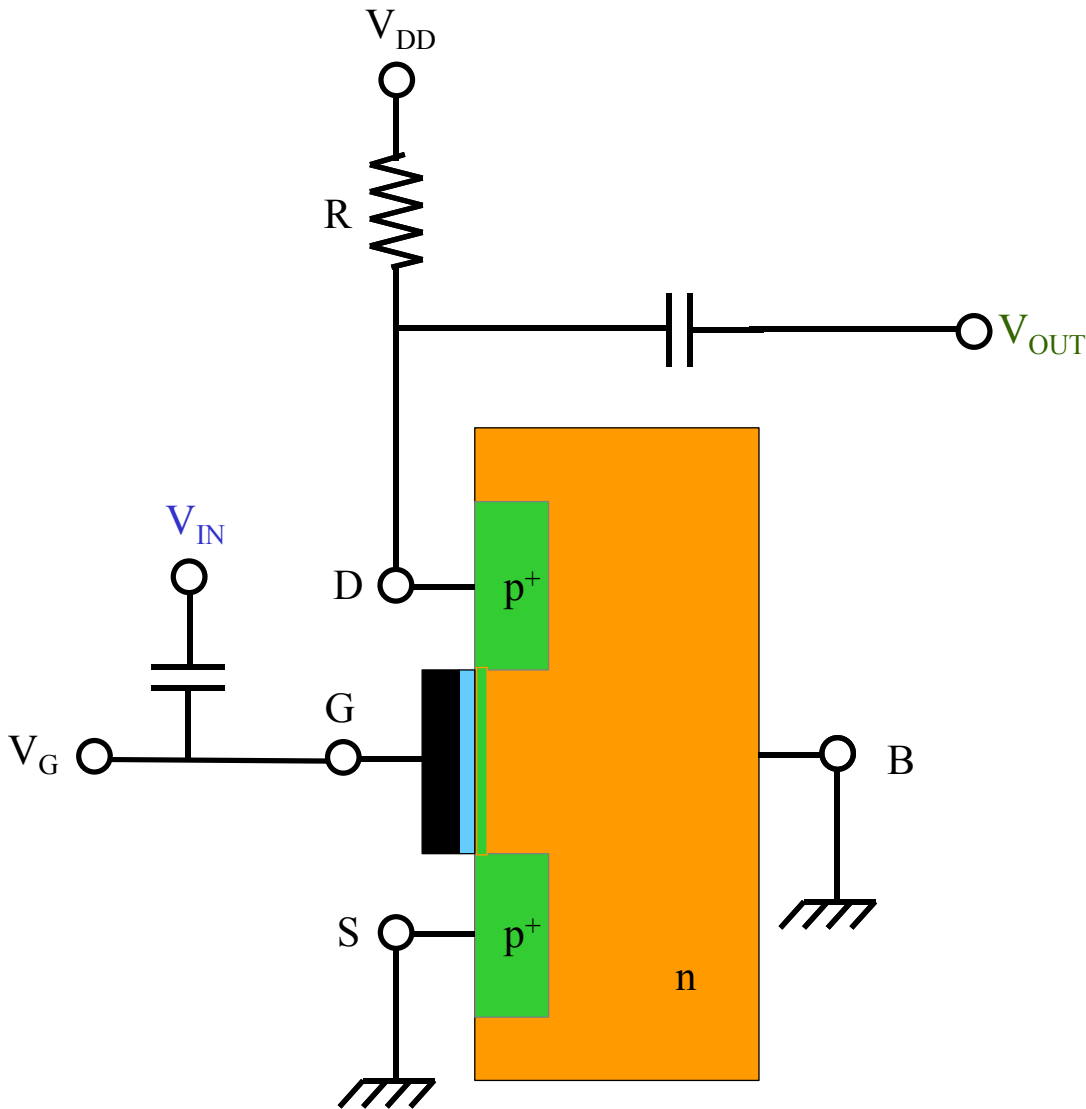


# 0.13 $\mu\text{m}$ CMOS technology: proton irradiation



- The gate leakage current is almost doubled at  $\Phi=10^{16}$  24 GeV p/cm<sup>2</sup>.
- This is due to defect generation in the gate oxide: RILC (already observed after  $\gamma$ -ray, X-ray and ion irradiation).

# The MOSFET as an amplifier



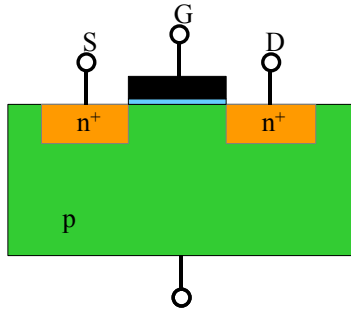
$$V_{OUT} = V_{IN} \times g_M \times R$$

$$g_M = \Delta I_{DS} / \Delta V_{GS}$$

Trans-conductance

# 0.13 $\mu\text{m}$ CMOS technology: proton irradiation

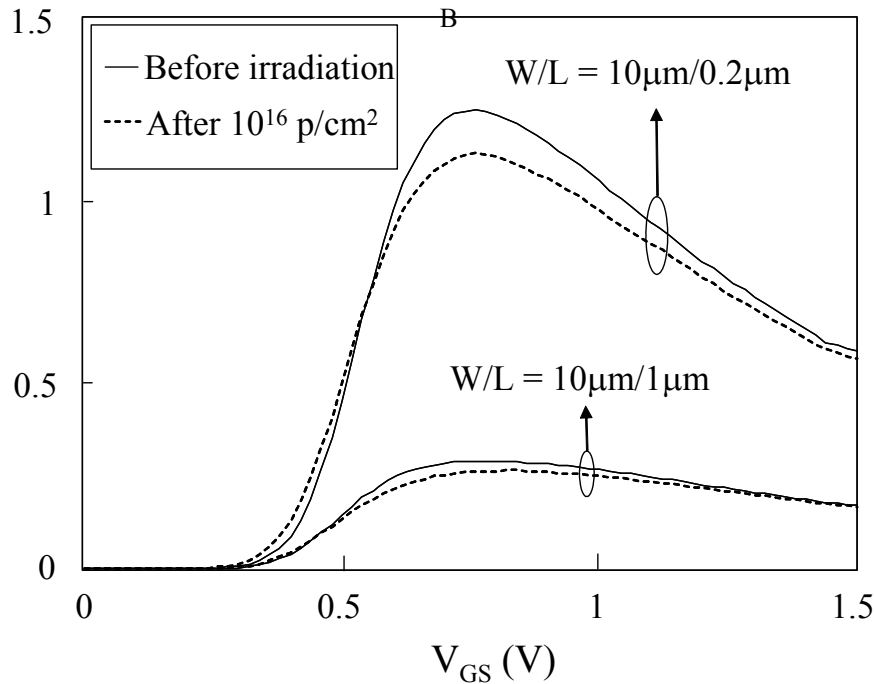
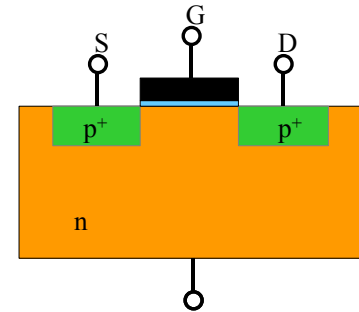
n-channel MOSFET



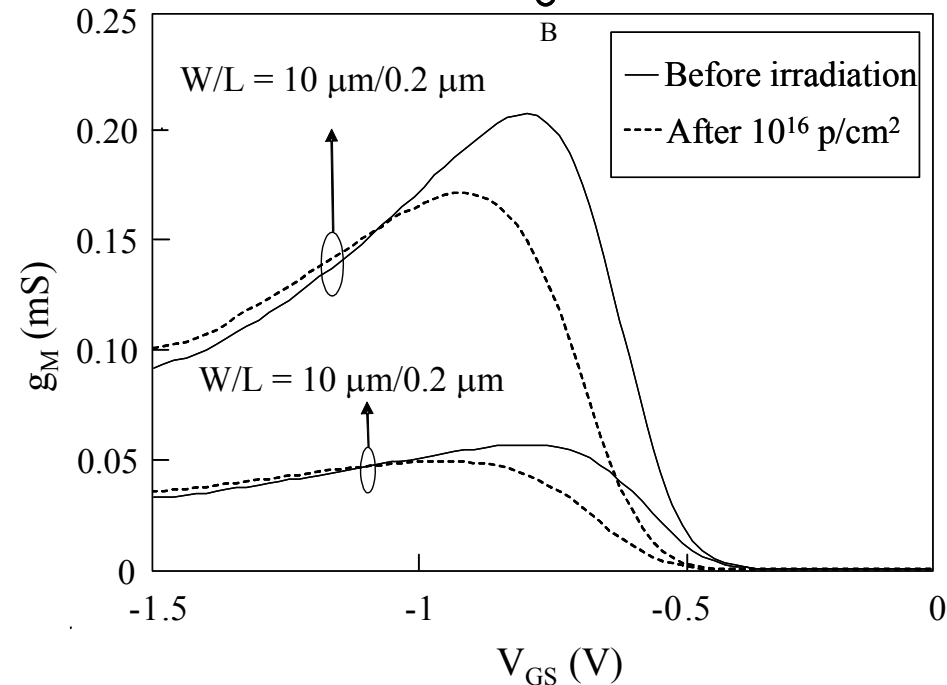
$$g_M = \Delta I_{DS} / \Delta V_{GS}$$

Trans-conductance

p-channel MOSFET



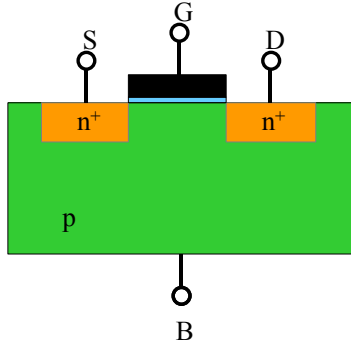
- $g_M$  peak decrease 15%.
- $g_M$  peak shift negligible.



- $g_M$  peak decrease higher than in n-MOSFET.
- $g_M$  peak shift of 200 mV.

# 0.13 $\mu\text{m}$ CMOS technology: proton irradiation

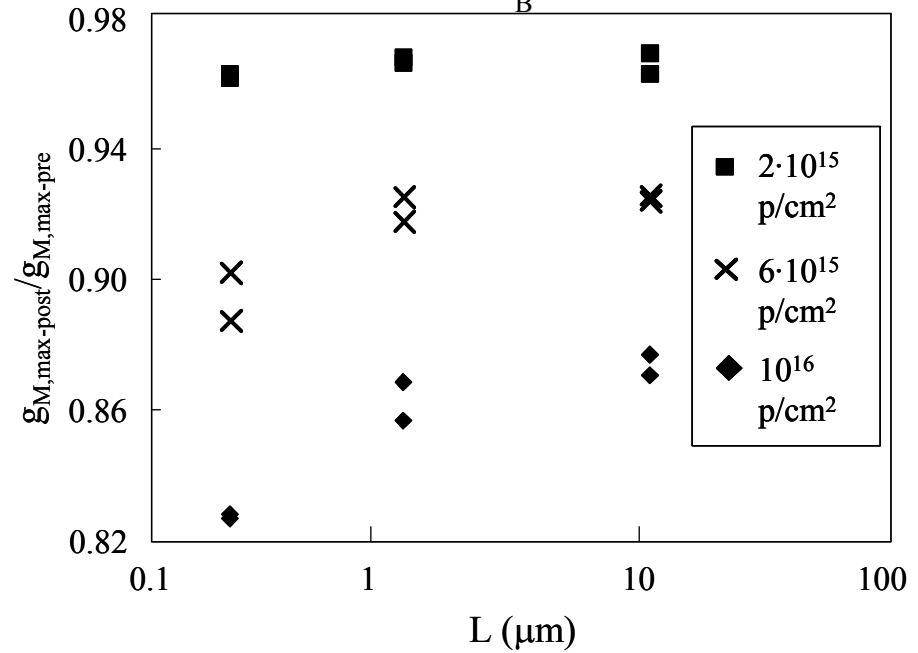
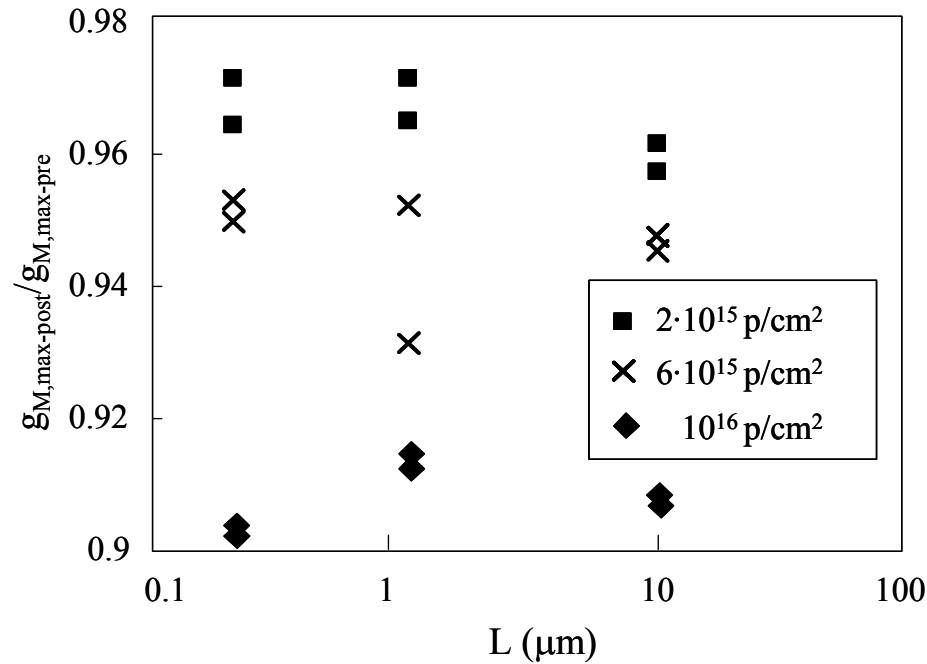
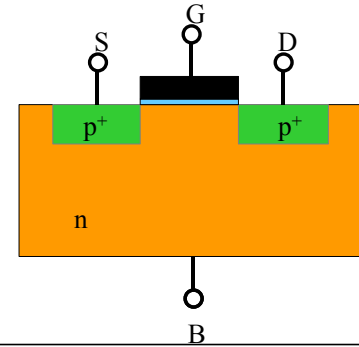
n-channel MOSFET



$$g_M = \Delta I_{DS} / \Delta V_{GS}$$

Trans-conductance

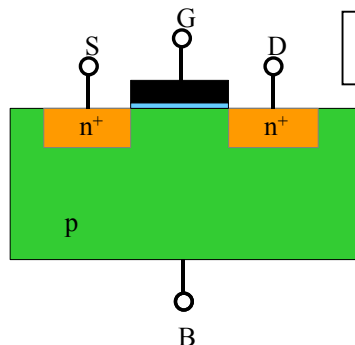
p-channel MOSFET



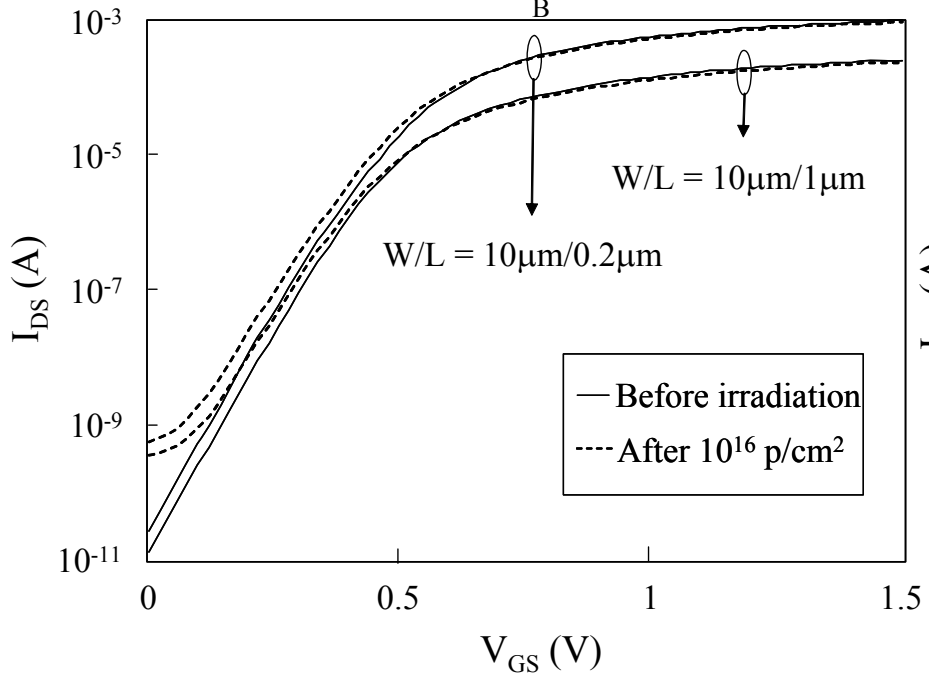
- The  $g_M$  peak degradation increases with fluence.

# 0.13 μm CMOS Technology: proton irradiation

n-channel MOSFET

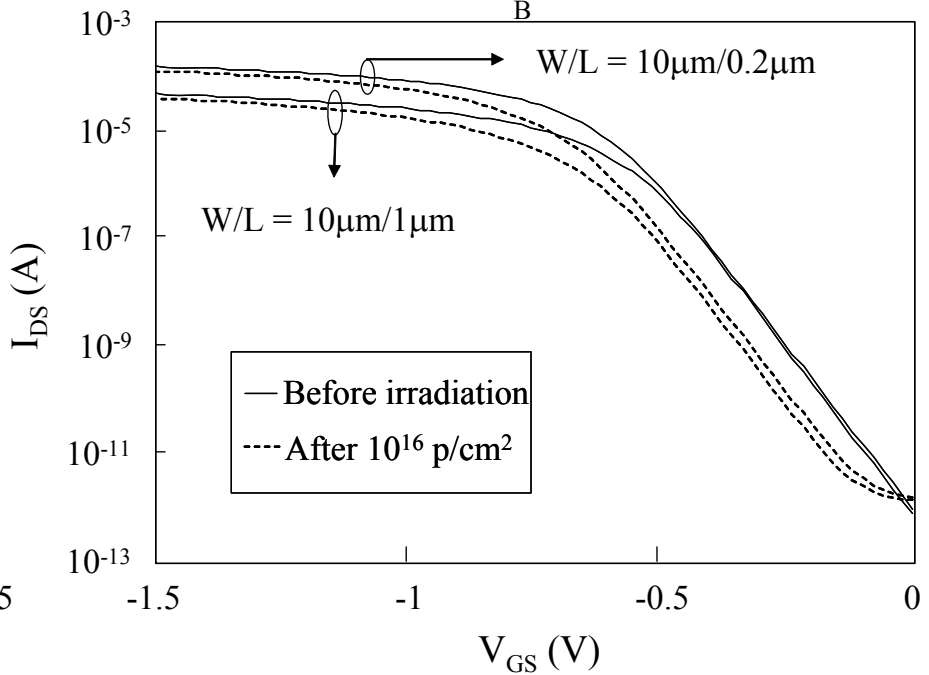
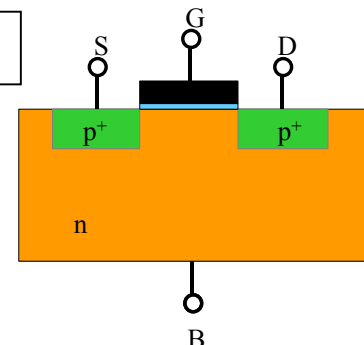


$$S = \ln 10 \times (\Delta V_{GS} / \Delta(\ln(I_{DS})))$$



- Subthreshold swing change negligible
- Subthreshold shift lower than in p-MOSFET

p-channel MOSFET



- Subthreshold swing change negligible
- Subthreshold shift higher than in n-MOSFET

# Conclusion

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## -Epitaxial detectors

- $V_{\text{dep}}$  vs  $\Phi$ :  $V_{\text{dep}}$  **minimum lower for neutrons** than protons.
- $V_{\text{dep}}$  vs annealing time for **24 GeV protons**:
  - SCSI can not be reached** during **proton irradiation**.
  - SCSI can be reached** during **annealing** at 80 °C.
- $V_{\text{dep}}$  vs annealing time for **nuclear reactor neutrons**:
  - SCSI can be reached during neutron irradiation?
  - SCSI can be reached** during **annealing** at 80 °C.

## -Pre-irradiated detectors

- The deep acceptor introduction rate  $\beta$  is **lower by a factor 2-3** for pre-irradiated **than standard FZ** devices **irradiated by neutrons**.

## -TMAH-thinned detectors

- **Device thinning limits** the  $V_{\text{dep}}$  **increase** after irradiation:
  - $V_{\text{dep}} \approx 290$  V for 100  $\mu\text{m}$  thick detectors at  $8 \times 10^{15}$  24 GeV p/cm<sup>2</sup>.
  - $V_{\text{dep}} \approx 70$  V for 50  $\mu\text{m}$  thick detectors at  $10^{16}$  24 GeV p/cm<sup>2</sup>.

## - 0.13 $\mu\text{m}$ CMOS technology

- The degradation has been evaluated up to  $10^{16}$  24 GeV p/cm<sup>2</sup> on MOSFET test structures.
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