

Graph 5

Graph 6

Graph 7

Graph 8

# NLO-QCD calculation in GRACE

- GRACE status -

Graph 9

Graph 10

Graph 11

Graph 12

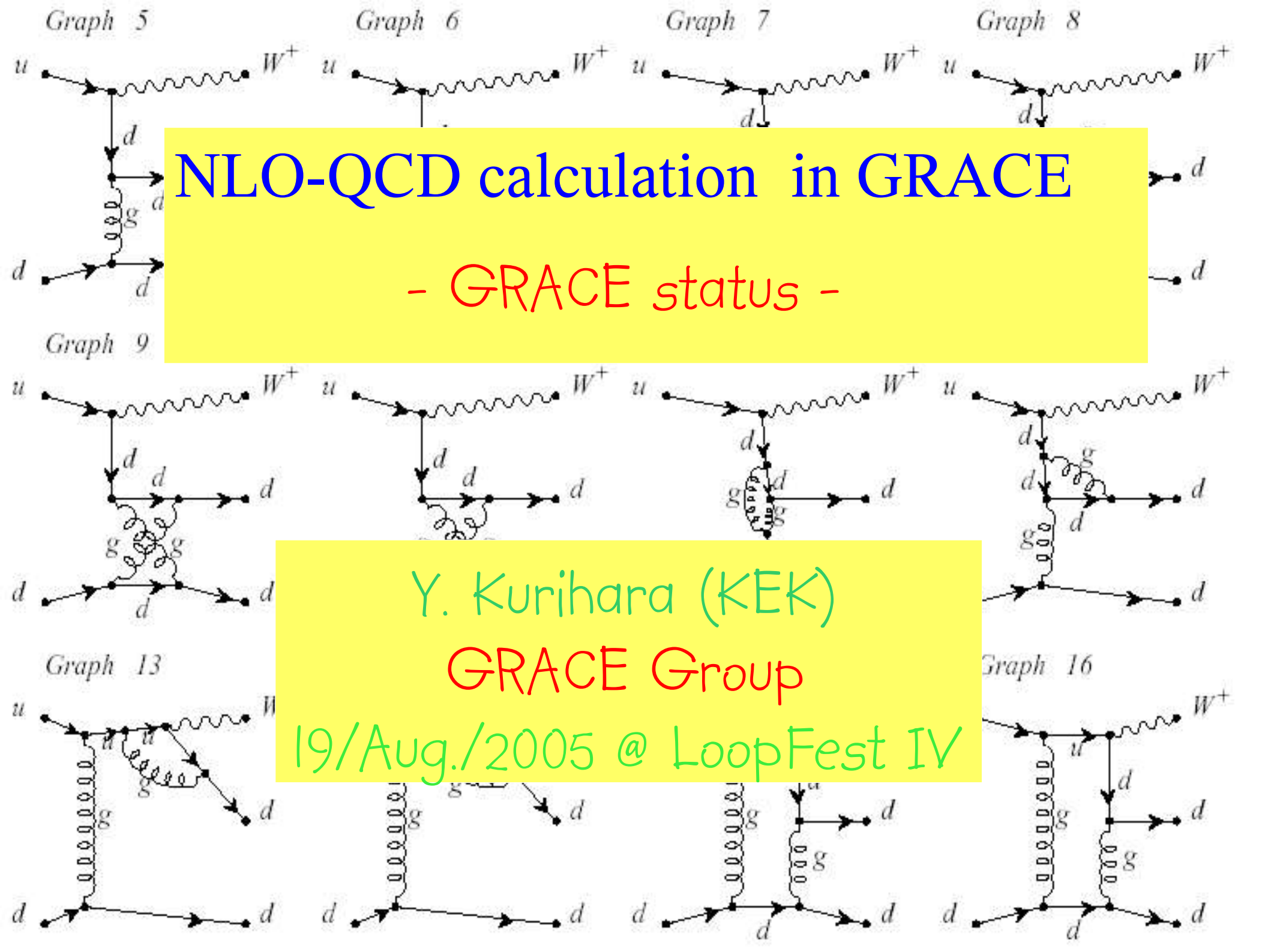
Y. Kurihara (KEK)

GRACE Group

19/Aug./2005 @ LoopFest IV

Graph 13

Graph 16

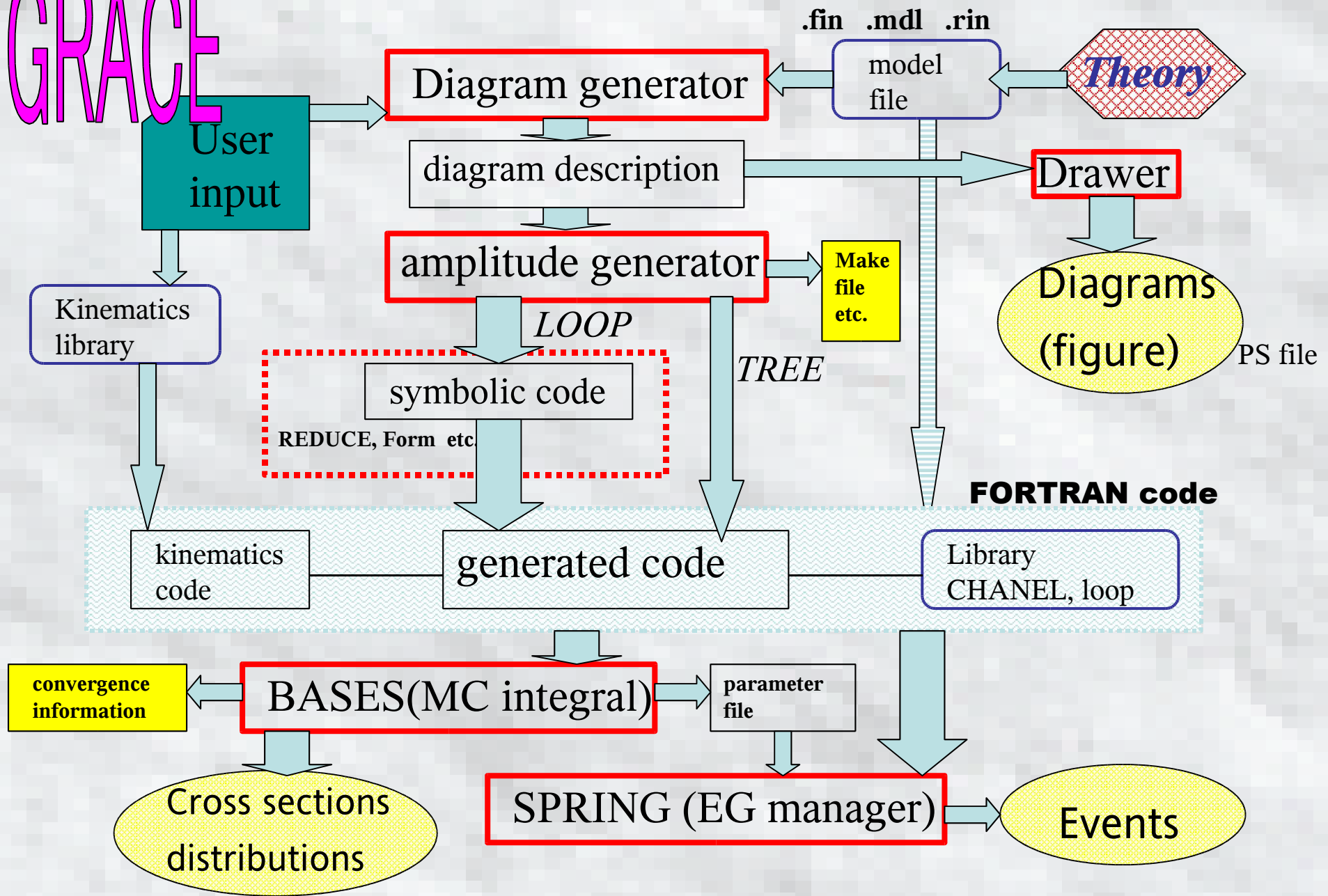


# GRACE Author list

A map of Japan is shown in the background. The main islands are colored in shades of green and yellow. A red dot is placed on the main island of Honshu, specifically in the Kanto region, indicating the location of the authors.

**J. Fujimoto, T. Ishikawa,  
M. Jimbo, T. Kaneko,  
K. Kato, S. Kawabata,  
T. Kon, Y. Kurihara,  
M. Kuroda, N. Nakazawa,  
Y. Shimizu, H. Tanaka,  
Y. Yuasa**

# GRACE





# Physics in LHC

LHC Experimental requirement

New Particle Search/Precision Measurements

LO-QCD Event generator+K-factor

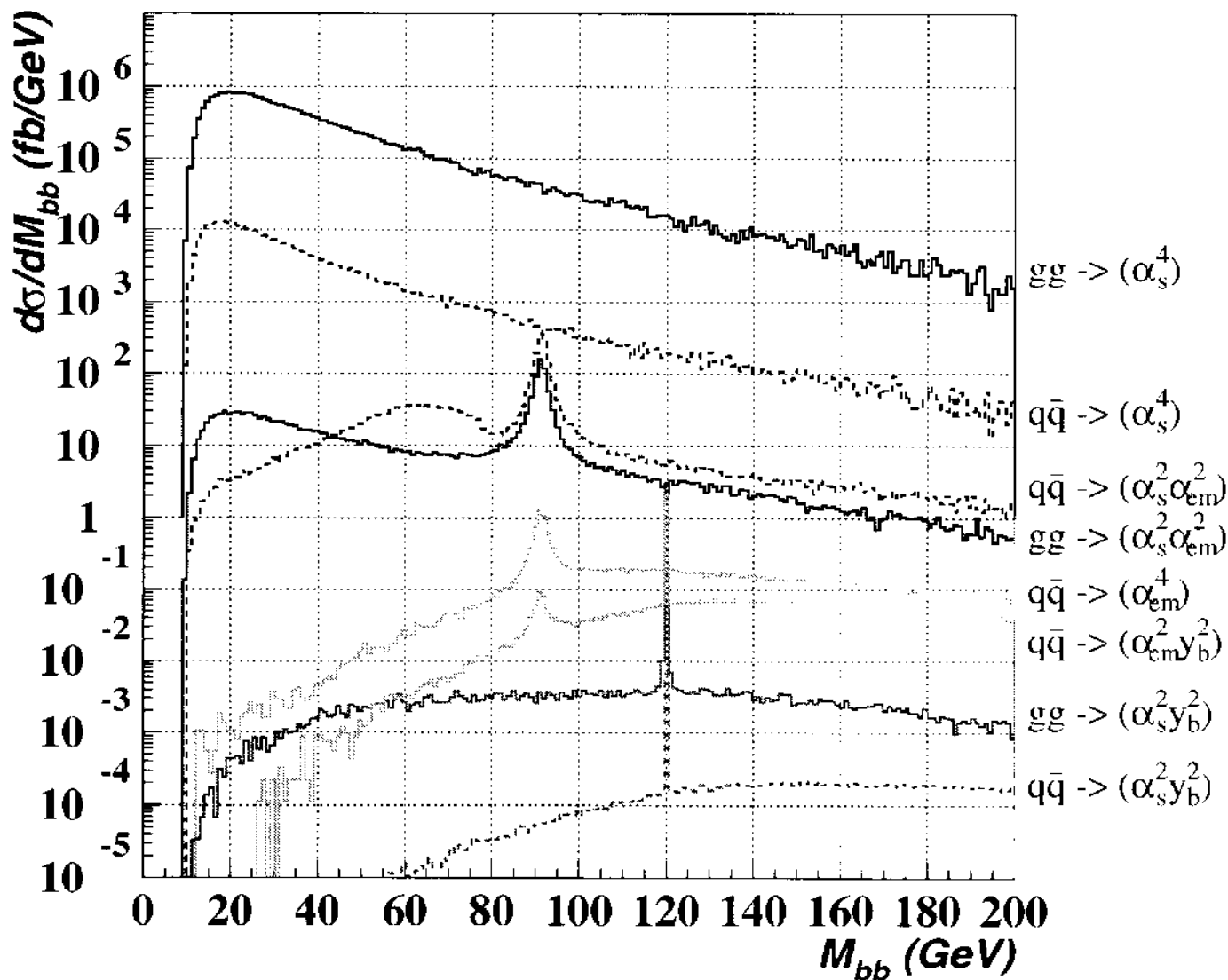


Obviously not enough!

We need  
NLO Event generator!

# QCD-event generator @ Tree-level

$pp \rightarrow bbbb$ , TEVATRON/LHC, GR@PPA\_4b, S. Tsuno



pp → many, TEVATRON/LHC, GR@PPA\_ALL,  
S. Tsuno

- W + jets (up to 4 jets) with the subsequent W decay to a fermion pair,
- Z + jets (up to 4 jets) with the subsequent Z decay to a fermion pair,
- Four bottom quarks via Z and Higgs-boson mediated processes as well as those from pure QCD interactions (same as GR@PPA\_4b),
- top-quark pair with the subsequent decay to W and b, and the W decay to a fermion pair,
- di-boson (WW, WZ and ZZ) with the subsequent W/Z decay to a fermion pair.

# Loop Calc. by GRACE-loop in ELWK proc. (Higgs production)

## ★ *Single Higgs production*

→  $e^+e^- \rightarrow ZH$  (full number of graphs = 341)

→  $e^+e^- \rightarrow \nu \nu H$  (1,350) *Phys.Lett. B559 (2003) 252-262*  
*A.Denner et.al. PLB 560(2003)196, NPB 660(2003)289*

→  $e^+e^- \rightarrow e^+e^- H$  (4,470) *Phys.Lett.B600 (2004) 65-76*

## ★ *top Yukawa*

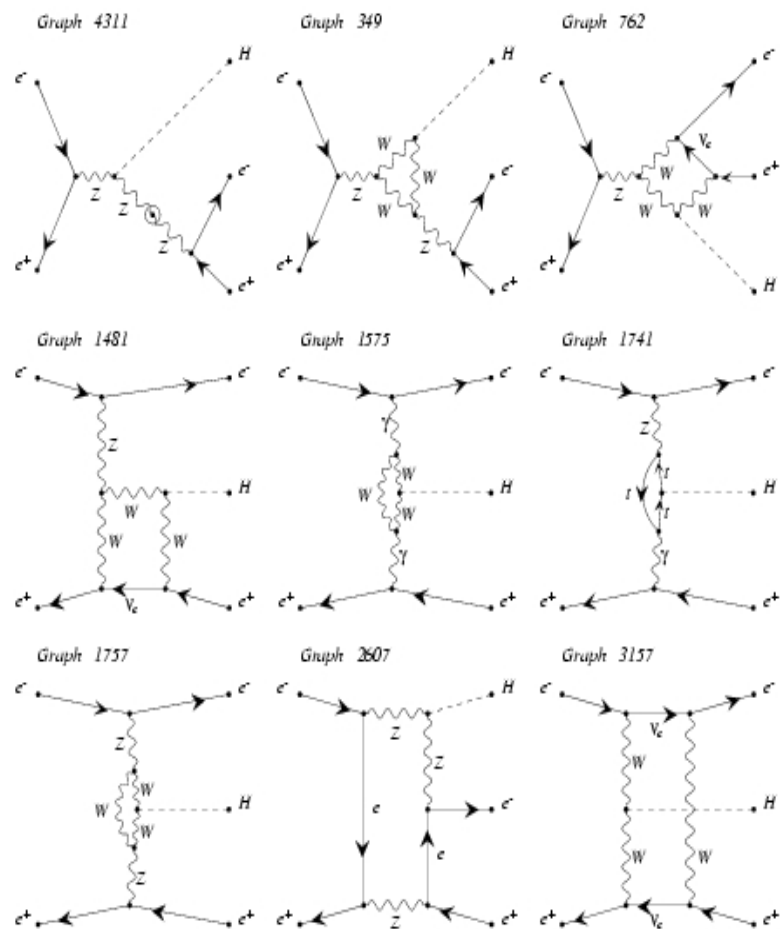
→  $e^+e^- \rightarrow ttH$  (2,327) *Phys.Lett. B571 (2003) 163-172*  
*Y.You et.al.PLB 571(2003)85*  
*A.Denner et.al. PLB 575(2003)290, NPB 680 (2004)85*

## ★ *Multi Higgs production*

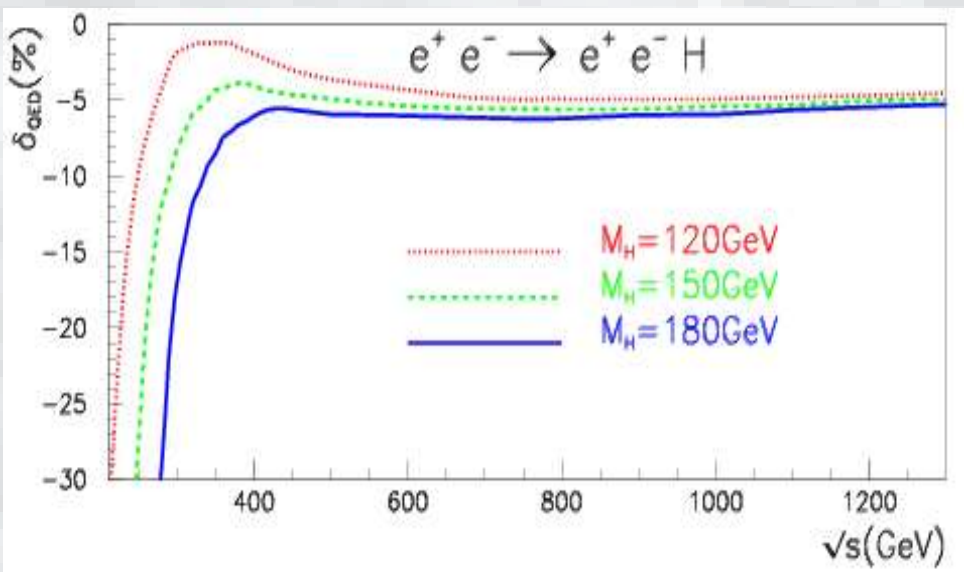
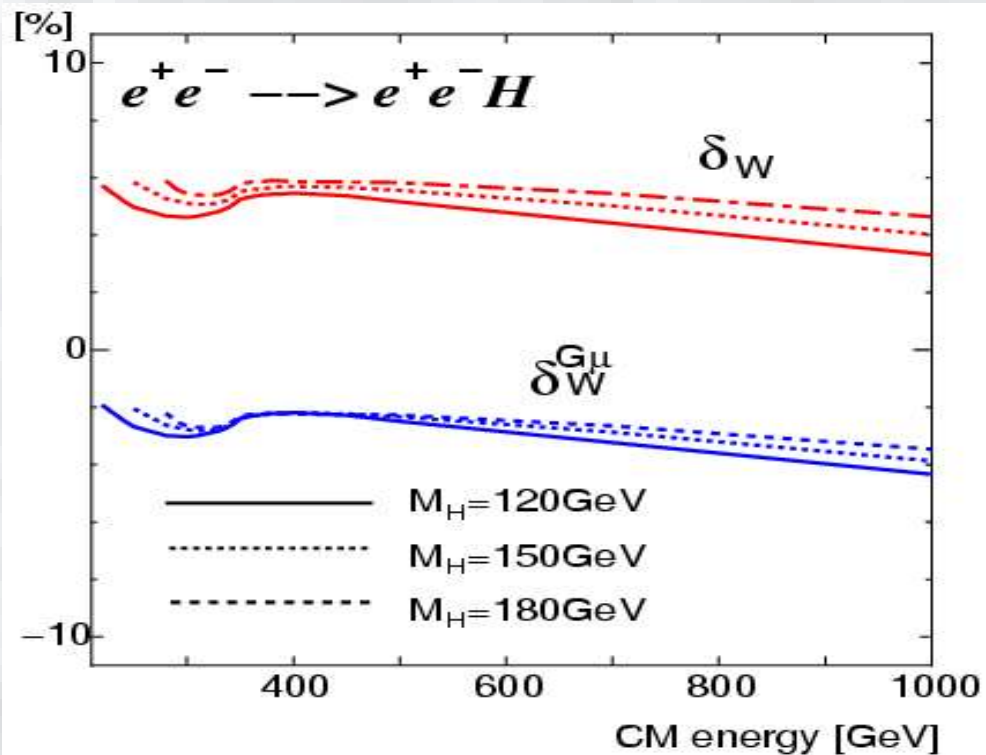
→  $e^+e^- \rightarrow ZHH$  (5,417) *Phys.Lett. B576 (2003) 152-164*  
*R.Zhang et.al.PLB(2004)349*

→  $e^+e^- \rightarrow \nu_e \nu_e HH$  (19,638) ⇒ *Preliminary*

$$e^+e^- \rightarrow e^+e^-H$$



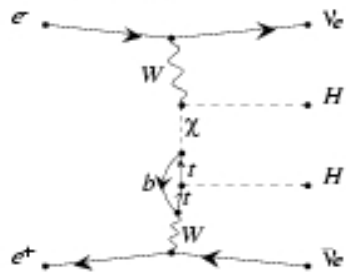
produced by GRACEFIG



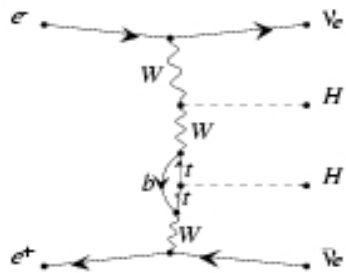


# $e^+e^- \rightarrow \nu_e \bar{\nu}_e HH$ (final 4-body process)

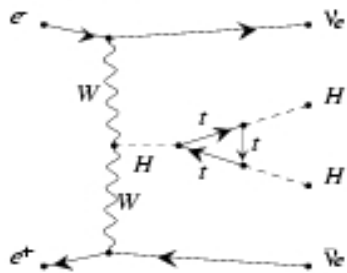
Graph 4403



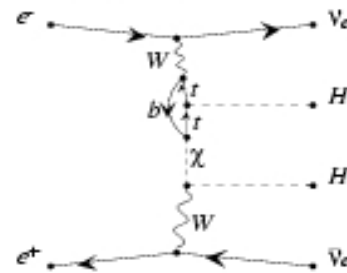
Graph 4471



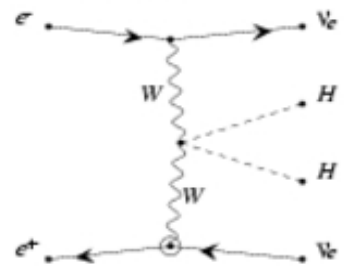
Graph 5384



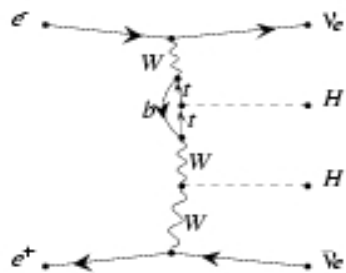
Graph 6652



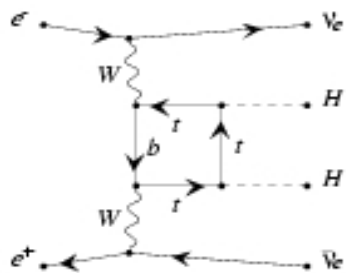
Graph 19205



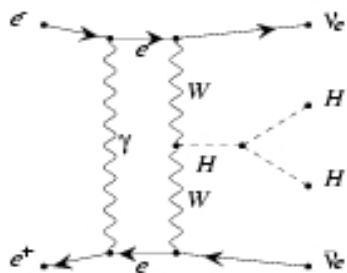
Graph 6686



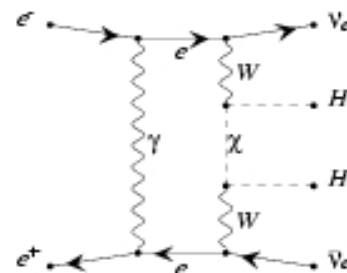
Graph 7205



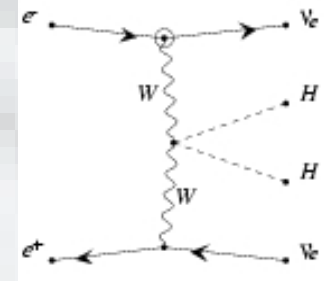
Graph 11477



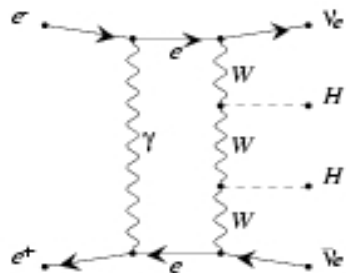
Graph 11496



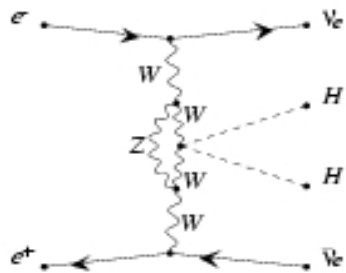
Graph 19208



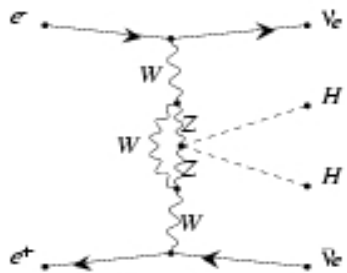
Graph 11497



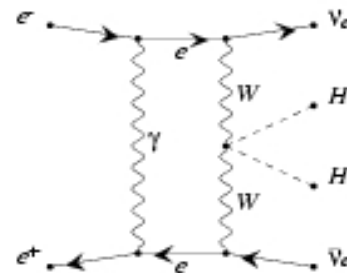
Graph 17308



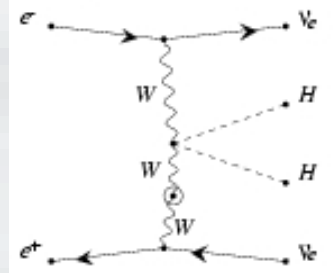
Graph 17310



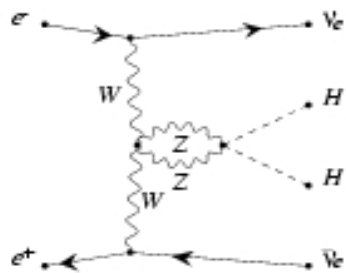
Graph 18283



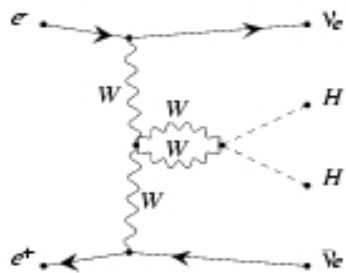
Graph 19631



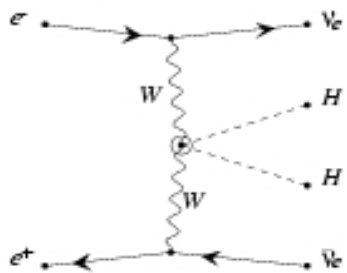
Graph 18857



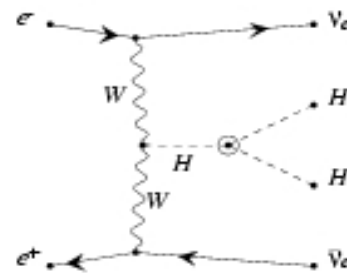
Graph 18858



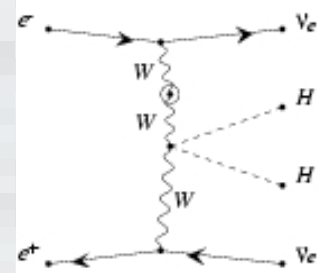
Graph 18873



Graph 18970

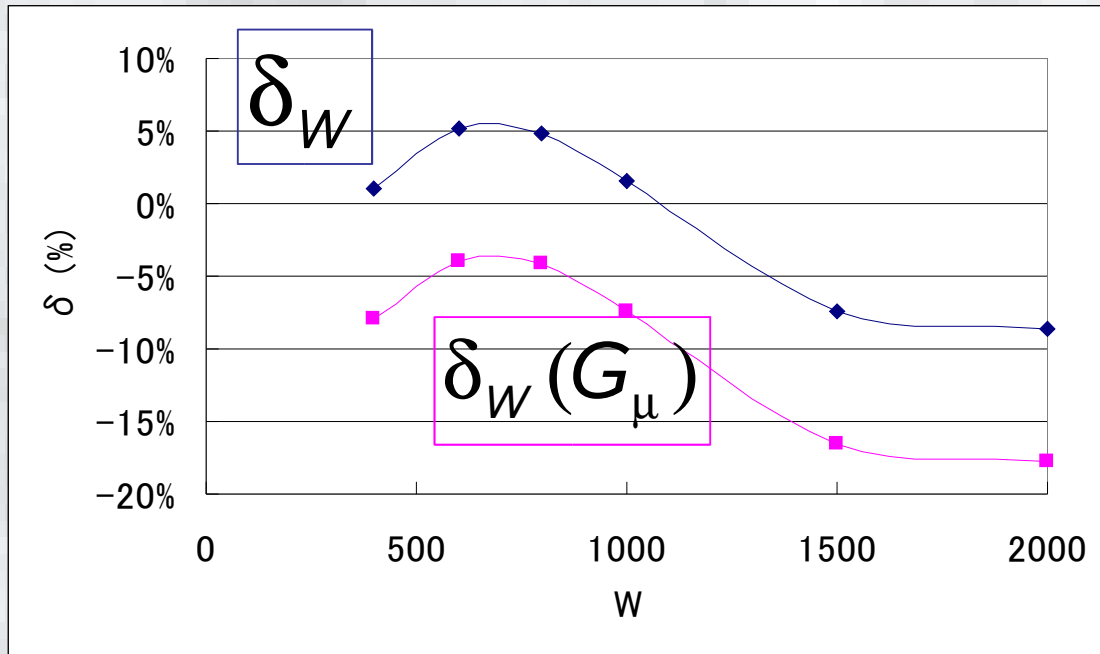


Graph 19635



$$e^+e^- \rightarrow \nu_\mu \bar{\nu}_\mu HH$$

$$\delta = \sigma(O(\alpha)) / \sigma(\text{tree}) - 1$$



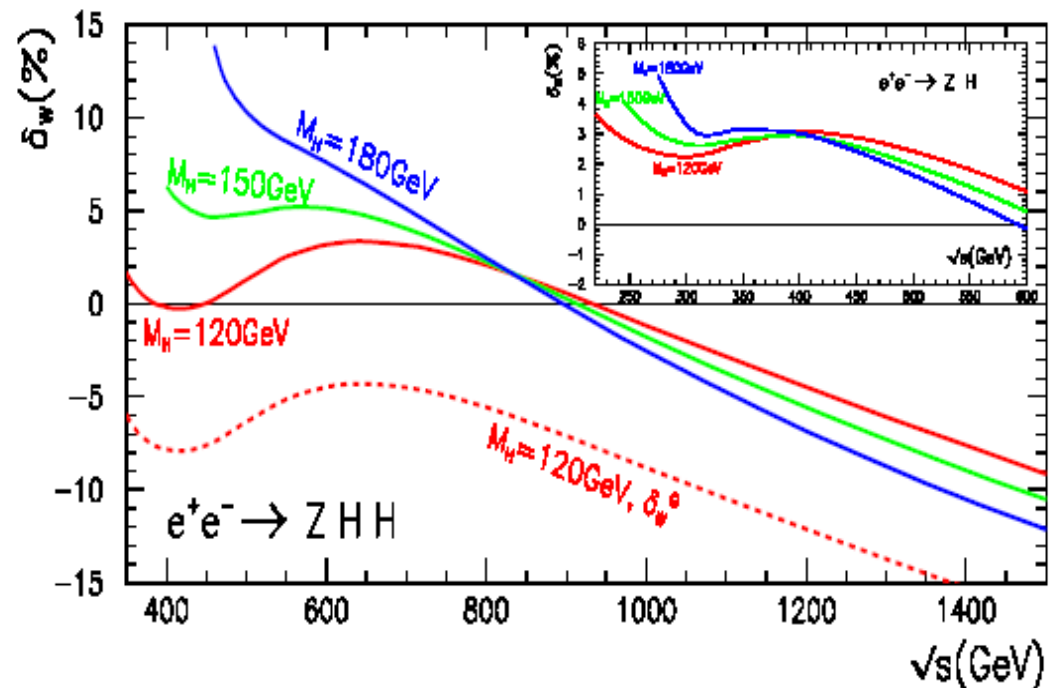
$M_H = 120 \text{ GeV}$

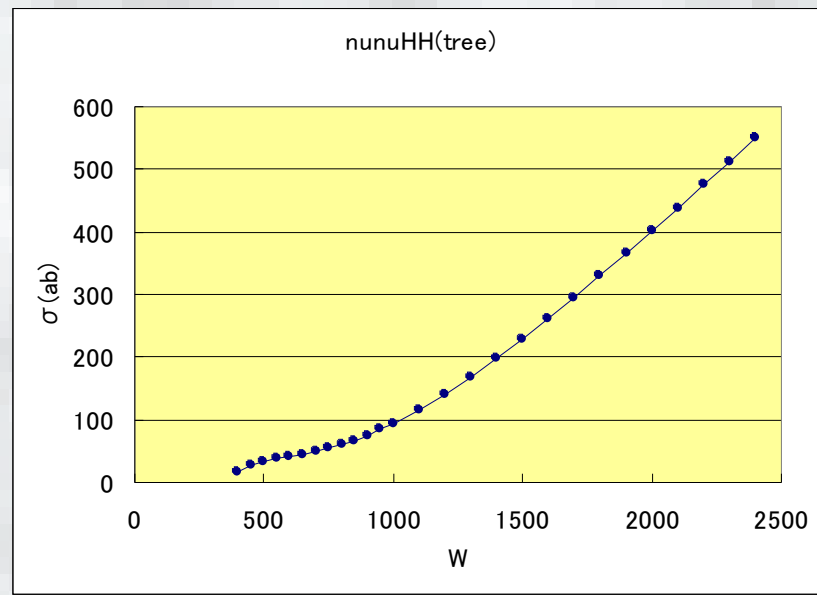
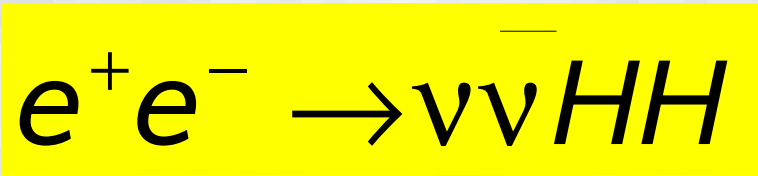
$m_t = 180 \text{ GeV}$

“s-channel”  

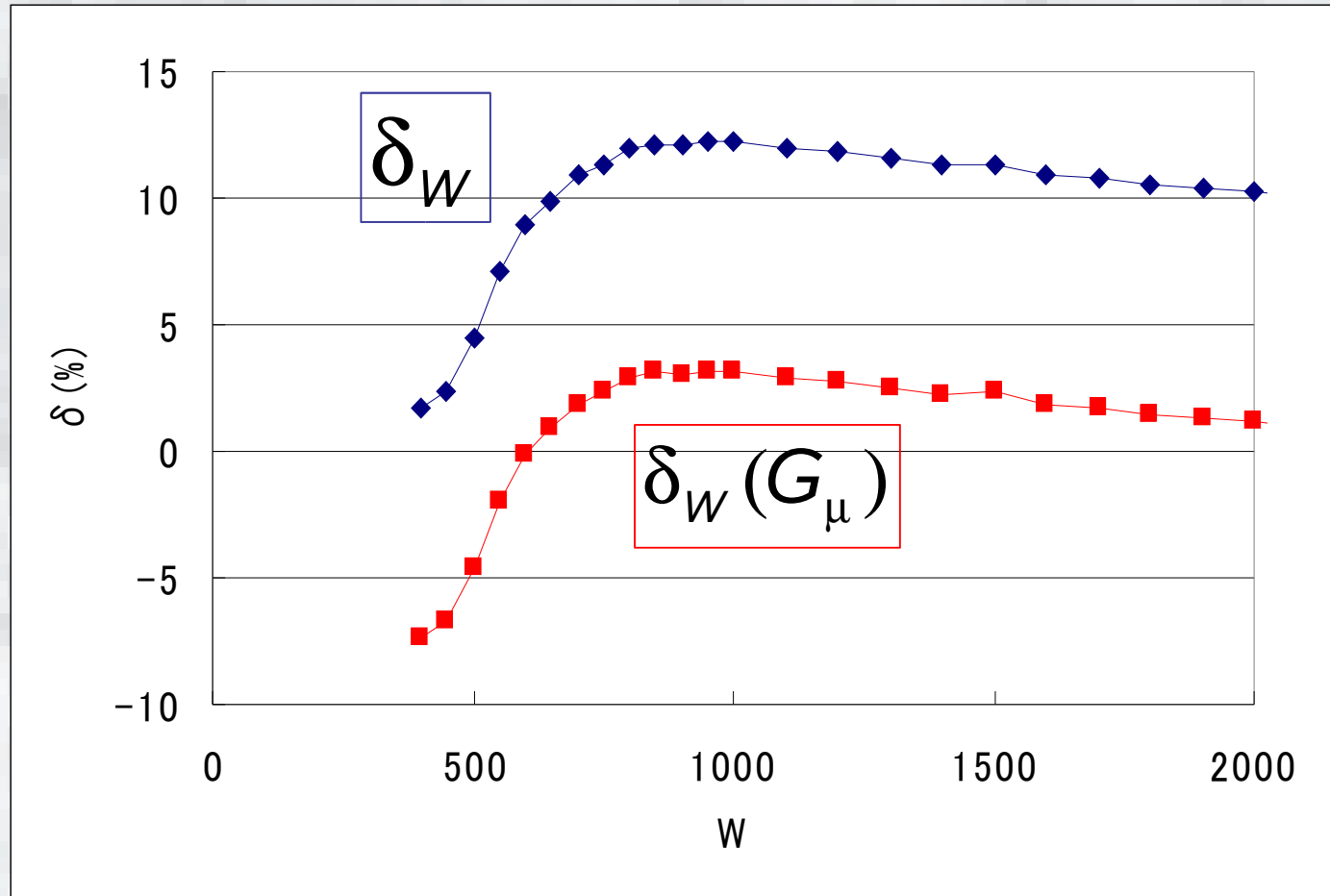
$$e^+e^- \rightarrow ZHH$$

Phys.Lett.B 576(2003)152.  
 $m_t = 174 \text{ GeV}$





$M_H=120\text{GeV}$



# Internal consistency check

For EW part:

System passes successfully the usual checks:

- ultraviolet finiteness (better than 20 digits)
- infrared finiteness (better than 20 digits)
- NLG dependence (better than 20 digits)
- $k_c$  dependence consistent with MonteCarlo statistical error (0.02%)

1. One random phase space point
2. Full set of diagrams
3. Quadruple precision

# • New feature of NLO-QCD issues in GRACE

QCD-tree : OK

ELWK 1-loop : OK

- PDF/PS  $\leftrightarrow$  Real emission Double counting

→ Virtuality ordering/LL-subtraction

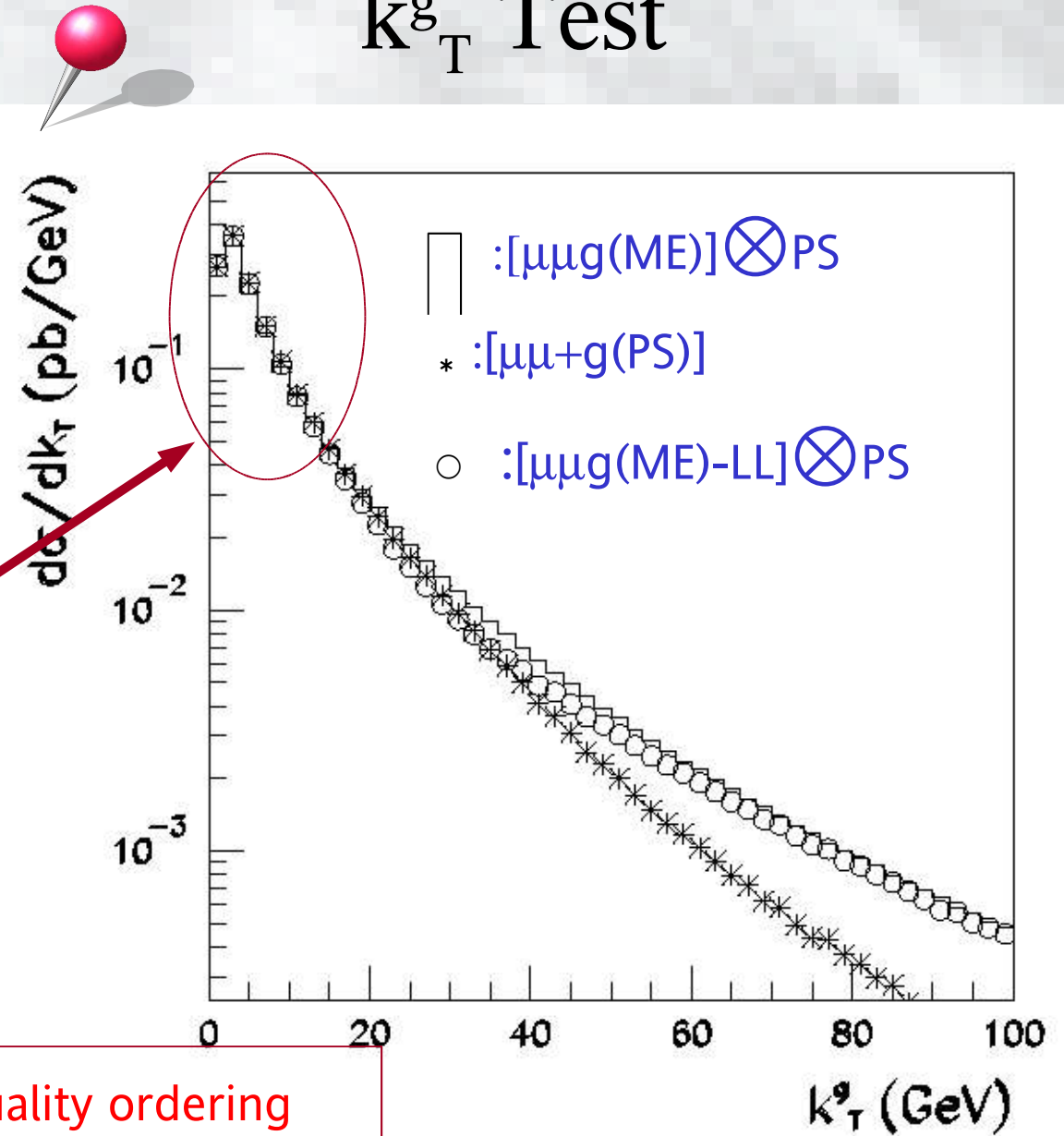
(YK et al. Nucl. Phys. B654 (2003) 301)

- Dimensional regularization in loop integrals for IR  
(fictitious mass in photon in ELWK)
- IR (soft/collinear) approximation terms

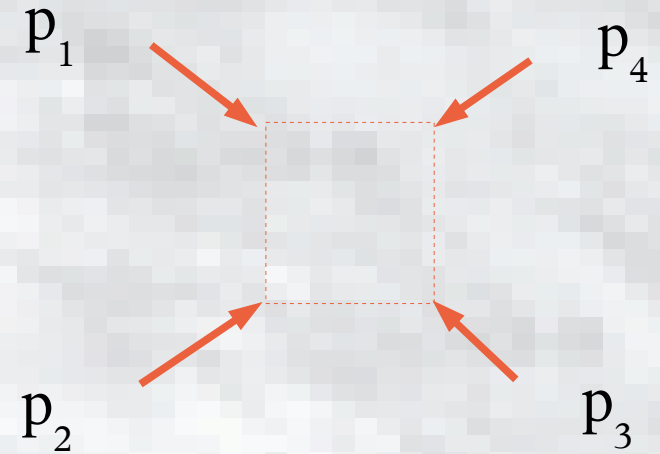
# Drell-Yan process

$k_T^g$  Test

- Process :  
 $u\bar{u} \rightarrow \mu^+\mu^- (+\text{gluon})$   
in  $p\bar{p}$  collision
- Cuts:  
 $\sqrt{s_{\mu\mu}} > 40 \text{ GeV}$   
 $k_T^g > 1 \text{ GeV}$



# Box Integral



$$J_{(4)}(s, t; p_1^2, p_2^2, p_3^2, p_4^2; n_x, n_y, n_z) = \frac{\Gamma(2 - \epsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\epsilon_{IR}}} \int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{D^{2-\epsilon_{IR}}},$$

$$D = -s xz - t yw - p_1^2 xy - p_2^2 yz - p_3^2 zw - p_4^2 xw - i0,$$

$$w = 1 - x - y - z,$$

$$s = (p_1 + p_2)^2,$$

$$t = (p_1 + p_4)^2.$$

# All on-shell (massless) external legs

$$\begin{aligned}
 J_4(s, t; 0, 0, 0, 0; n_x, n_y, n_z) &= \frac{1}{(4\pi)^2 s t} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) n_x! \Gamma(\varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR}) \\
 &\times \left[ \left( \frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \left( \frac{-t}{s} \right)^{n_x} \frac{B(1 + n_z, n_x + n_y + \varepsilon_{IR})}{\Gamma(n_x + \varepsilon_{IR})} \right. \\
 &\times {}_2F_1 \left( 1 + n_x, n_x + n_y + \varepsilon_{IR}, 1 + n_x + n_y + n_z + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \\
 &+ \left( \frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \sum_{l=0}^{n_x} \left( \frac{-s}{t} \right)^l \frac{(-1)^l}{\Gamma(l + \varepsilon_{IR}) (n_x - l)!} B(1 + n_y, l + n_z + \varepsilon_{IR}) \\
 &\left. \times {}_2F_1 \left( 1 + l, l + n_z + \varepsilon_{IR}, 1 + l + n_y + n_z + \varepsilon_{IR}, -\frac{\tilde{u}}{t} \right) \right],
 \end{aligned}$$

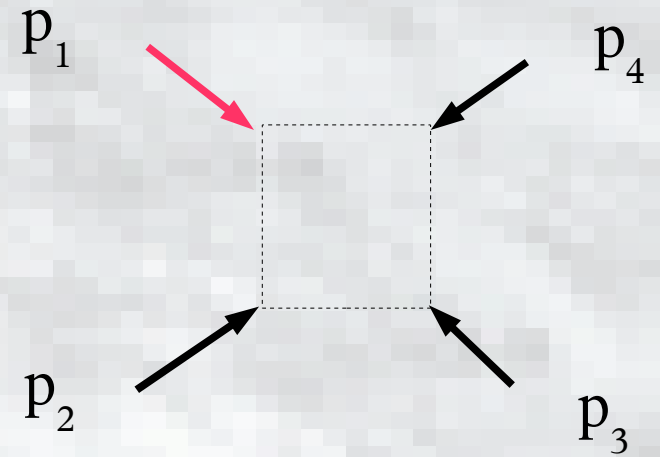
## Scalar Integral

$$\begin{aligned}
 J_{(4)}(s, t; 0, 0, 0, 0; 0, 0, 0) &= \frac{1}{(4\pi)^2 s t} \frac{B(\varepsilon_{IR}, \varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR})}{\varepsilon_{IR}} \\
 &\times \left[ \left( \frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left( 1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) + \left( \frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left( 1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \right]
 \end{aligned}$$

This result is compared with G. Duplanić, B. Nižić, Eur. Phys. J. C **20**, 357 (2001)



# One off-shell box integral



$$J_4(s, t; p_1^2, 0, 0, 0; n_x, n_y, n_z) = \frac{\Gamma(2 - \varepsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\varepsilon_{IR}}} \int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{(-xzs - y(1-x-y-z)t - p_1^2 xy - i0)^{2-\varepsilon_{IR}}}$$

$$= \frac{1}{(4\pi)^2 s t} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) n_x! \Gamma(\varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR})$$

$$\times \left[ \left( \frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \left( \frac{-t}{s} \right)^{n_x} \frac{B(1 + n_z, n_x + n_y + \varepsilon_{IR}) \mathcal{I}^{(1)}}{\Gamma(n_x + \varepsilon_{IR})} \right.$$

$$\left. + \left( \frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \sum_{l=0}^{n_x} \frac{(-1)^l B(1 + n_y, l + n_z + \varepsilon_{IR}) \mathcal{I}_l^{(2)}}{\Gamma(l + \varepsilon_{IR}) (n_x - l)!} \right]$$

$$\mathcal{I}^{(1)} = B(1 + n_z, n_x + n_y + \varepsilon_{IR}) {}_2F_1 \left( 1 + n_x, n_x + n_y + \varepsilon_{IR}, 1 + n_x + n_y + n_z + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right)$$

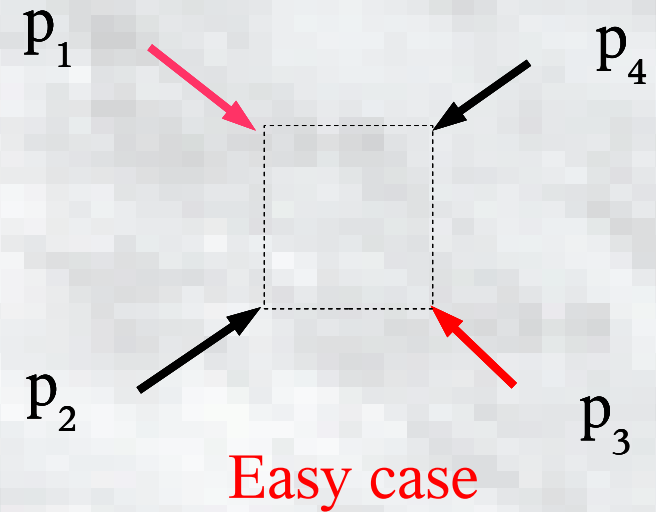
$$\begin{aligned} \mathcal{I}_l^{(2)} &= \sum_{k_1=0}^{n_z} n_z C_{k_1} \left( \frac{s}{p_1^2 - s} \right)^{n_y + k_1} \sum_{k_2=0}^{n_y + k_1} n_{y+k_1} C_k (-1)^{n_y + k_2} \left( \frac{-t}{s} \right) \\ &\times \int_0^1 dw \left( 1 + \frac{\tilde{u}}{\tilde{s}} w \right)^{-(l+1)} \left( 1 + \frac{\tilde{t} + \tilde{u}}{\tilde{s}} w \right)^{k_2 + l - 1 + \varepsilon_{IR}} \\ &= \sum_{k_1=0}^{n_z} \sum_{k_2=0}^{n_y + k_1} n_z C_{k_1} n_{y+k_1} C_k (-1)^{k_1 + k_2} \left( \frac{s}{p_1^2 - s} \right)^{n_y + k_1} \frac{1}{l + k_2 + \varepsilon_{IR}} \left( 1 + \frac{u}{t} \right)^l \\ &\times \left[ {}_2F_1 \left( 1 + l, l + k_2 + \varepsilon_{IR}, 1 + l + k_2 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) \right. \\ &\left. - \left( \frac{\tilde{p}_1^2}{\tilde{s}} \right)^{l + k_2 + \varepsilon_{IR}} {}_2F_1 \left( 1 + l, l + k_2 + \varepsilon_{IR}, 1 + l + k_2 + \varepsilon_{IR}, -\frac{\tilde{u}\tilde{p}_1^2}{\tilde{t}\tilde{s}} \right) \right], \end{aligned}$$

## Scalar Integral

$$\begin{aligned}
 J_4(s, t; p_1^2, 0, 0, 0; 0, 0, 0) &= \frac{1}{(4\pi)^2 s t} \frac{B(\varepsilon_{IR}, \varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR})}{\varepsilon_{IR}} \\
 \times &\left[ \left( \frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left( 1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) + \left( \frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left( 1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \right. \\
 - &\left. \left( \frac{-\tilde{p}_1^2}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left( 1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}\tilde{p}_1^2}{\tilde{t}\tilde{s}} \right) \right],
 \end{aligned}$$

This result is compared with G. Duplanić, B. Nižić, Eur. Phys. J. C **20**, 357 (2001)

# Two off-shell box integral



$$J_4(s, t; p_1^2, 0, p_3^2, 0; n_x, n_y, n_z) = \frac{\Gamma(2 - \epsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\epsilon_{IR}}}$$

$$\int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{(-xzs - y(1-x-y-z)t - p_1^2 xy - p_3^2 z(1-x-y-z) - i0)^{2-\epsilon_{IR}}}$$

$$= \frac{1}{(4\pi)^2 (s - p_3^2)(t - p_3^2)} B(n_x + \epsilon_{IR}, n_y + n_z + \epsilon_{IR}) n_x! \Gamma(\epsilon_{IR}) \Gamma(1 - \epsilon_{IR})$$

$$\times \left[ \left( \frac{\bar{t} - p_3^2}{4\pi\mu_R^2} \right)^{\epsilon_{IR}} \left( \frac{t - p_3^2}{s - p_3^2} \right)^{n_x} \frac{1}{\Gamma(n_x + \epsilon_{IR})} \mathcal{I}^{(1)} + \left( \frac{-\bar{s}}{4\pi\mu_R^2} \right)^{\epsilon_{IR}} \sum_{l=0}^{n_x} \frac{(-1)^l}{\Gamma(l + \epsilon_{IR})(n_x - l)!} \mathcal{I}_l^{(2)} \right]$$

$$\begin{aligned}
\mathcal{I}^{(1)} &= \frac{1}{n_x + \varepsilon_{IR}} \sum_{k_1=0}^{n_x} \sum_{k_2=0}^{n_y+k_1} n_x C_{k_1} n_{y+k_1} C_{k_2} (-1)^{k_1+k_2} \left( \frac{p_3^2 - s}{u} \right)^{n_y+k_1} (1-\alpha)^{k_2-n_x-1} \\
&\times \left[ \left( 1 + \frac{\bar{p}_3^2}{\bar{t} - \bar{p}_3^2} \right)^{n_x + \varepsilon_{IR}} {}_2F_1 \left( 1 + n_x - k_2, n_x + \varepsilon_{IR}, 1 + n_x + \varepsilon_{IR}, \frac{\bar{u}/(\bar{s} - \bar{p}_3^2) + \alpha}{\alpha - 1} \right) \right. \\
&\left. - \left( \frac{\bar{p}_3^2}{\bar{t} - \bar{p}_3^2} \right)^{n_x + \varepsilon_{IR}} {}_2F_1 \left( 1 + n_x - k_2, n_x + \varepsilon_{IR}, 1 + n_x + \varepsilon_{IR}, \frac{\alpha}{\alpha - 1} \right) \right]
\end{aligned}$$

$$\begin{aligned}
\mathcal{I}_l^{(2)} &= \sum_{k_1=0}^{n_x} \sum_{k_2=0}^{n_y+k_1} n_x C_{k_1} n_{y+k_1} C_{k_2} (-1)^{k_1+k_2} \left( \frac{s}{s - p_1^2} \right)^{n_y+k_1} \frac{1}{l + k_2 + \varepsilon_{IR}} \left( \frac{1}{1 - \beta} \right)^{l+1} \left( \frac{t - p_3^2}{t + u - p_3^2} \right) \left( \frac{s}{s - p_3^2} \right)^l \\
&\times \left[ {}_2F_1 \left( 1 + l, l + k_2 + \varepsilon_{IR}, 1 + l + k_2 + \varepsilon_{IR}, \frac{\beta}{\beta - 1} \right) \right. \\
&\left. - \left( \frac{\bar{p}_1^2}{\bar{s}} \right)^{l+k_2+\varepsilon_{IR}} {}_2F_1 \left( 1 + l, l + k_2 + \varepsilon_{IR}, 1 + l + k_2 + \varepsilon_{IR}, \frac{\beta}{\beta - 1} \frac{\bar{p}_1^2}{\bar{s}} \right) \right]
\end{aligned}$$

$$\alpha = \frac{\bar{p}_3^2}{\bar{t} - \bar{p}_3^2} \frac{\bar{u}}{\bar{s} - \bar{p}_3^2}, \quad \beta = \frac{\bar{u}}{\bar{s} - \bar{p}_3^2} \frac{\bar{s}}{\bar{t} + \bar{u} - \bar{p}_3^2}$$

# Numerical calculation (IR finite case)

For  $l, m, n \in \mathcal{N}$ ,

$$\begin{aligned} & {}_2F_1(l, m+1, n+m+2; z) \\ &= \frac{1}{B(m+1, n+1)} \int_0^1 \tau^m (1-\tau)^n (1-z\tau)^{-l} d\tau \\ &= \sum_{k_1=0}^n \sum_{k_2=0}^{m+k_1} (-1)^{k_1+k_2} \frac{n C_{k_1} m+k_1 C_{k_2}}{B(m+1, n+1) z^{m+k_1}} \int_0^1 (1-z\tau)^{-l+k_2} d\tau. \end{aligned}$$

$$\int_0^1 (1-z\tau)^{-l+k_2} d\tau = \begin{cases} -\frac{\ln(1-z)}{z} & k_2 - l + 1 = 0, \\ \frac{1}{k_2 - l + 1} \frac{(1-z)^{k_2 - l + 1} - 1}{-z} & k_2 - l + 1 \neq 0. \end{cases}$$

Numerical check: mathematica  $\leftrightarrow$  our FORTRAN program  
more than ten digit agreement

# Numerical calculation (IR divergent case)

$$\begin{aligned}
 \mathcal{I}_{l,m,n} &\equiv \int_0^1 \tau^{l+n-1+\varepsilon_{IR}} (1-\tau)^m (1-z\tau)^{-(l+1)} d\tau, \\
 &= B(1+m, l+n+\varepsilon_{IR}) \underline{{}_2F_1(1+l, l+n+\varepsilon_{IR}, 1+l+n+m+\varepsilon_{IR}, z)} \\
 &= \sum_{j=-1}^{\infty} \mathcal{F}_{l,m,n}^{(j)}(z) \varepsilon_{IR}^j,
 \end{aligned}$$

Expansion w.r.t.  $\varepsilon_{IR}$

$$\tilde{F}_{j_1, j_2}^{(n)}(z) \equiv \frac{(-1)^n}{n!} \int_0^1 d\tau \tau^{j_1-1} (1-z\tau)^{-(j_2+1)} \ln^n \tau,$$

When  $n=1$

$$\tilde{F}_{1,0}^{(1)}(z) = \frac{\text{Li}_2(z)}{z},$$

$$\tilde{F}_{1,1}^{(1)}(z) = -\frac{\ln(1-z)}{z}$$

$$\tilde{F}_{1,j_2+1}^{(1)}(z) = \frac{j_2}{j_2+1} \tilde{F}_{1,j_2}^{(1)} + \frac{(1-z)^{-j_2} - 1}{j_2(j_2+1)z}$$

$$\tilde{F}_{j_1,j_2}^{(1)}(z) = \frac{1}{z^{j_1-1}} \sum_{k=0}^{j_1-1} (-1)^k {}_{j_1-1}C_k \tilde{F}_{1,j_2-k}^{(1)}(z)$$



When  $n=2$

$$\tilde{F}_{1,0}^{(2)}(z) = \frac{\text{Li}_3(z)}{z}$$

$$\tilde{F}_{1,1}^{(2)}(z) = \frac{\text{Li}_2(z)}{z}$$

$$(j_2 + 1)\tilde{F}_{1,j_2+1}^{(2)}(z) - j_2\tilde{F}_{1,j_2}^{(2)}(z) - \tilde{F}_{1,j_2}^{(1)}(z) = 0$$

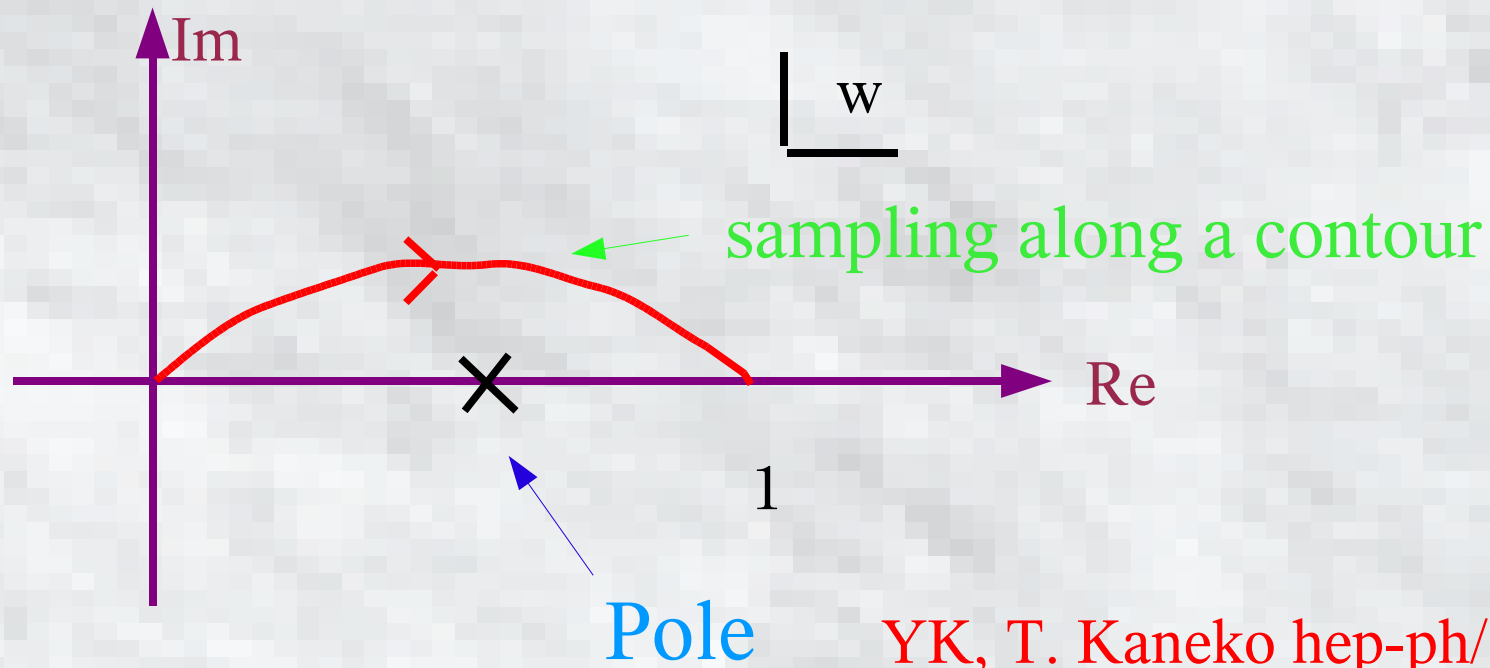
$$\tilde{F}_{j_1,j_2}^{(2)}(z) = \frac{1}{z^{j_1-1}} \sum_{k=0}^{j_1-1} (-1)^k {}_{j_1-1}C_k \tilde{F}_{1,j_2-k}^{(2)}(z)$$



# Numerical test for Infrared finite case

$$\begin{aligned} J_4 &= (4\pi\mu^2)^{-\varepsilon_{IR}} \int_0^1 dr r^{-1+n_x+\varepsilon_{IR}} (1-r)^{-1+n_y+n_z+\varepsilon_{IR}} \\ &\times \int_0^1 dv \int_0^1 dw \frac{w^{n_y} (1-w)^{n_z} v^{n_x}}{(-s v(1-w) - t (1-v)w - i0)^{2-\varepsilon_{IR}}} \\ &= (4\pi\mu^2)^{-\varepsilon_{IR}} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) \\ &\times \int_0^1 dv \int_0^1 dw \frac{w^{n_y} (1-w)^{n_z} v^{n_x}}{(-s v(1-w) - t (1-v)w - i0)^{2-\varepsilon_{IR}}} \end{aligned}$$

Numerical contour integral



$$J_4(s, t; 0, 0, 0, 0; n_x, n_y, n_z)$$

$n_x$	$n_y$	$n_z$	real/imag.	analytic	NCI
1	2	3	real	$-2.15298 \times 10^{-9}$	$-2.15297 \times 10^{-9}$
			imag.	$-2.78647 \times 10^{-9}$	$-2.78650 \times 10^{-9}$
2	0	2	real	$9.74570 \times 10^{-9}$	$9.74572 \times 10^{-9}$
			imag.	$-3.22229 \times 10^{-8}$	$-3.22230 \times 10^{-8}$

$$J_4(s, t; p_1^2, 0, 0, 0; n_x, n_y, n_z)$$

$n_x$	$n_y$	$n_z$	real/imag.	analytic	NCI
1	2	3	real	$-7.88683 \times 10^{-10}$	$-7.88689 \times 10^{-10}$
			imag.	$-1.95176 \times 10^{-9}$	$-1.95176 \times 10^{-9}$
2	0	2	real	$1.48133 \times 10^{-8}$	$1.48133 \times 10^{-8}$
			imag.	$-2.04318 \times 10^{-8}$	$-2.04318 \times 10^{-8}$

$$s=123, t=-200, p_1^2=80$$

# Numerical test for IR divergent case

A sector decomposition method can be used.

$$J_4(s, t; 0, 0, 0, 0; 0, 0, 0)$$

	real/imag.	analytic	SD
$1/\epsilon_{IR}^2$	real	$-1.029686826 \times 10^{-6}$	$-1.029686826 \times 10^{-6}$
	imag.	0	$\mathcal{O}(10^{-16})$
$1/\epsilon_{IR}$	real	$-5.205325212 \times 10^{-6}$	$-5.205325212 \times 10^{-6}$
	imag.	$1.617428283 \times 10^{-6}$	$1.617428283 \times 10^{-6}$
$\epsilon_{IR}^0$	real	$-9.739160873 \times 10^{-6}$	$-9.739160872 \times 10^{-6}$
	imag.	$8.569648363 \times 10^{-6}$	$8.569648363 \times 10^{-6}$

$$s=123, t=-200$$

# IR cancellation test

ex. Prompt photon production

@ one phase point

$$\delta = a_2/\epsilon_{\text{IR}}^2 + a_1/\epsilon_{\text{IR}} + a_0$$

**BOX**

$$a_1 = -18601.9993715016$$

$$a_2 = -3793.95539013131$$

**S/C+vertex-(terms included in PDF)**

$$a_1 = 18601.9993714494$$

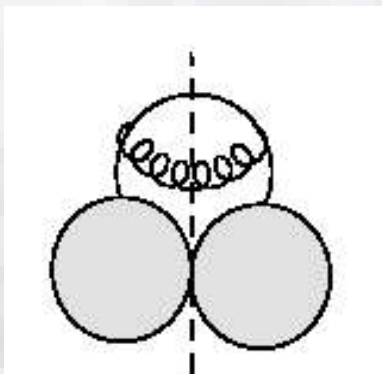
$$a_2 = 3793.95539013130$$

$1/\epsilon_{\text{IR}}^2$  : Soft/Coll. + Loop  $\sim O(10^{-15})$

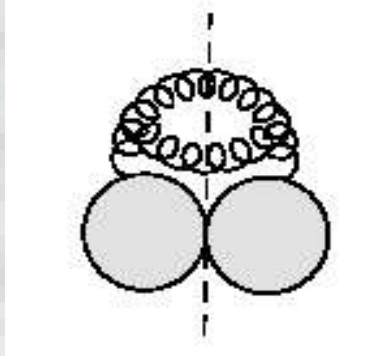
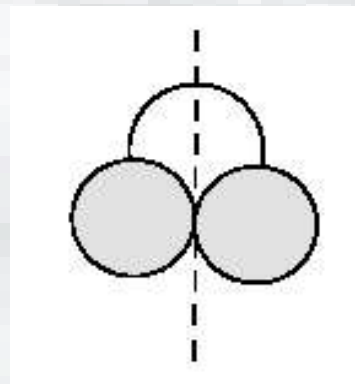
$1/\epsilon_{\text{IR}}$  : Soft/Coll. + Loop  $\sim O(10^{-12})$

# Soft/Collinear Approximation

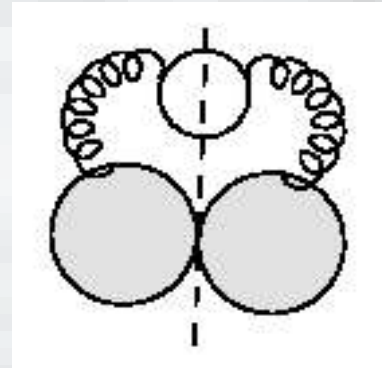
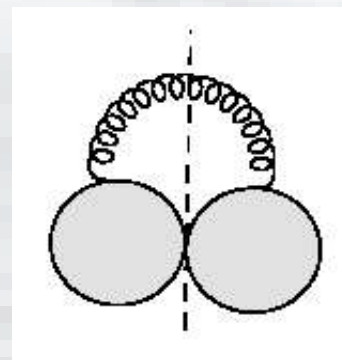
In axial gauge,



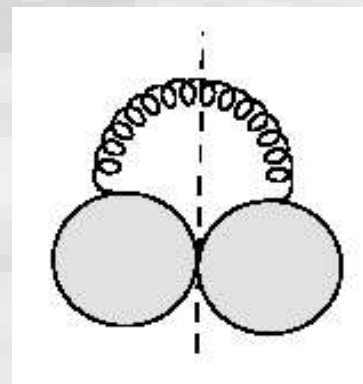
$$= f_{q \rightarrow qg} \times$$



$$= f_{g \rightarrow gg} \times$$

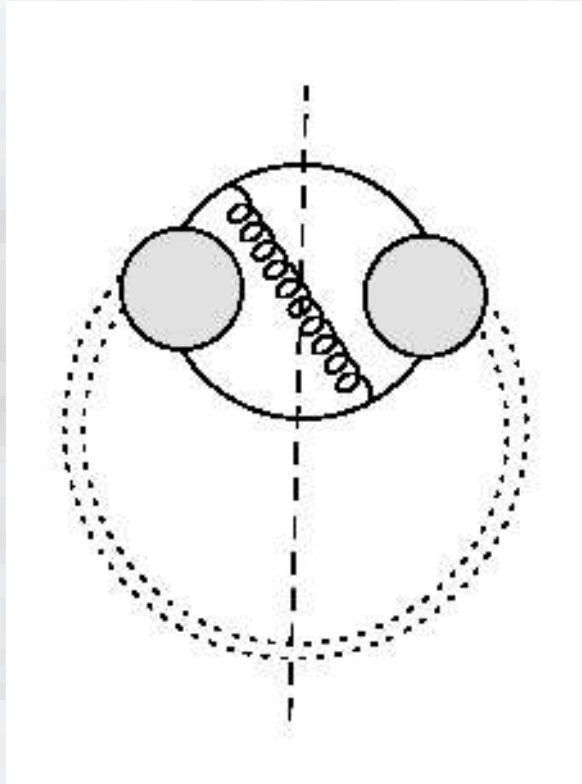


$$= f_{g \rightarrow qq} \times$$

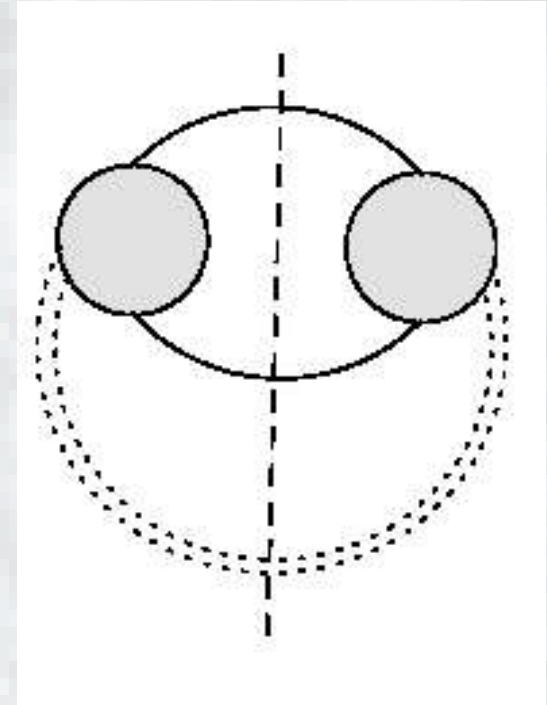


# Soft/Collinear Approximation

In axial gauge,



$$= f^{int}_{q \rightarrow qg} \times$$



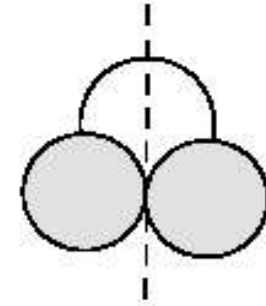
$$\begin{aligned} f^{int}_{q \rightarrow qg} &= 0 && \text{in collinear region} \\ &\neq 0 && \text{in soft region} \end{aligned}$$

All  $f$ 's will be implemented in the system soon.

# Soft/Collinear Approximation

$$f_{q_{in} \rightarrow q_{in} g_{out}}$$

$$\sigma_{coll} = \frac{1}{(2p_1^0)(2p_2^0)v_{rel}} \int_{\Omega_{full}} d\Phi_{N+1}^{(d)} f_{q \rightarrow qg} \times$$



$$= \left(\frac{s}{4\pi\mu^2}\right)^{\varepsilon_{IR}} \frac{B(\varepsilon_{IR}, \varepsilon_{IR})}{2\Gamma(1 + \varepsilon_{IR})} f_c \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) P(x) \left(\frac{1-x}{x}\right)^{2\varepsilon_{IR}}$$

$$= \sigma_0(s) \frac{\alpha_s}{2\pi} f_c \left[ \frac{1}{\varepsilon_{IR}^2} + \frac{2L-3}{2\varepsilon_{IR}} - \frac{\pi^2}{4} + \frac{L^2}{2} \right]$$

$$+ \int_0^1 dx \sigma_0(xs) \phi(x, \varepsilon_{IR})$$

$$+ f_c \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) \left[ L \frac{1+x^2}{(1-x)_+} + \frac{(1+x^2) \ln(1-x)}{(1-x)_+} - \frac{1+x^2}{1-x} \ln x \right]$$

$$\phi(x, \varepsilon_{IR}) = \frac{1}{\varepsilon_{IR}} f_c \frac{\alpha_s}{2\pi} P(x) = \frac{1}{\varepsilon_{IR}} f_c \frac{\alpha_s}{2\pi} \frac{1+x^2}{(1-x)_+}$$

$$L = \ln(s/\mu^2).$$



# Status of GRACE NLO Generator



## Already tested

- + Drell-Yan process
- + W production
- + Prompt photon



## Under development

- +V+1 jet



## Future plan

- +V+2 jets,.....
- +VV+jet,2jets,.....

# Summary

- (1) Automatic loop calculation in ELWK is established.
- (2) QCD-NLO Matrix Elements
  - Automatic generation by GRACE
- (3) Loop integral
  - Numerical loop-integration library
- (4) Soft/Collinear treatment
  - Basic tools of NLL-PS is available (under implementation)
  - LL-subtraction method
  - General calculation recipe
- (5) Application
  - Drell-Yan process/w production
  - prompt photon
  - $V+1\text{jet}, 2\text{jet}, VV+1\text{ jet}$