

Automating radiative corrections in Bhabha scattering

Alejandro Lorca



DESY Zeuthen

Contents

1. Motivation
2. Introduction to one-loop issues
3. Automation with *a*TALC
 - (a) Diagrams \rightarrow Algebra \rightarrow Numerics
4. Results
5. Conclusions & Outlook



I. Motivation

Motivation: ILC and Bhabha physics

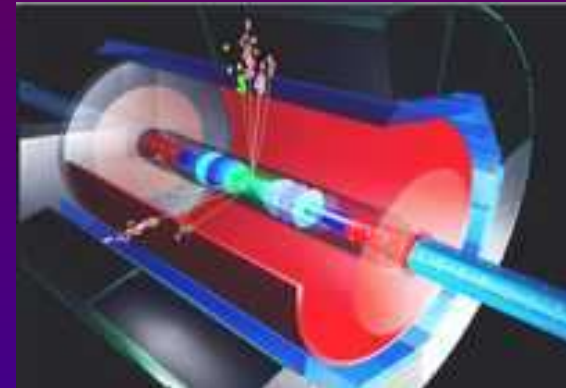
The Standard Model (SM) is very **successfull**, but . . .
present few “experiment vs. theory” $2\text{-}3\sigma$ disagreements
(A_{FB}^b , $\sin \theta_{\text{eff}}$, N_ν, \dots) require **better control** of theoretical uncertainties.

Motivation: ILC and Bhabha physics

The Standard Model (SM) is very **successful**, but . . .
present few “experiment vs. theory” $2-3\sigma$ disagreements
(A_{FB}^b , $\sin\theta_{\text{eff}}$, N_ν, \dots) require **better control** of theoretical uncertainties.

Next International Linear Collider (ILC)
will be a **challenge**

e^+e^- beam, $E_{\text{CM}} \approx 1\text{TeV}$, $\%_0$ precision

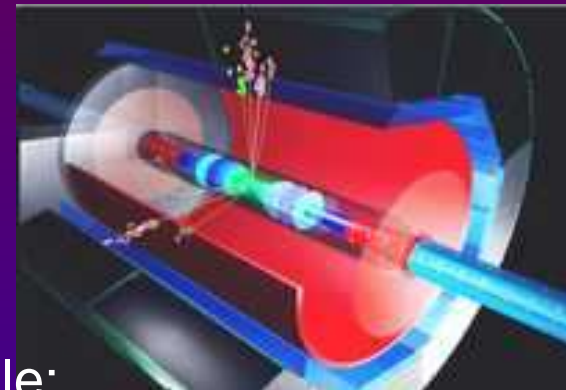


Motivation: ILC and Bhabha physics

The Standard Model (SM) is very **successful**, but . . . present few “experiment vs. theory” $2\text{-}3\sigma$ disagreements (A_{FB}^b , $\sin \theta_{\text{eff}}$, N_ν, \dots) require **better control** of theoretical uncertainties.

Next International Linear Collider (ILC) will be a **challenge**

e^+e^- beam, $E_{\text{CM}} \approx 1\text{TeV}$, $\%_0$ precision



Bhabha scattering ($e^+e^- \rightarrow e^+e^-$) plays a dominant role:

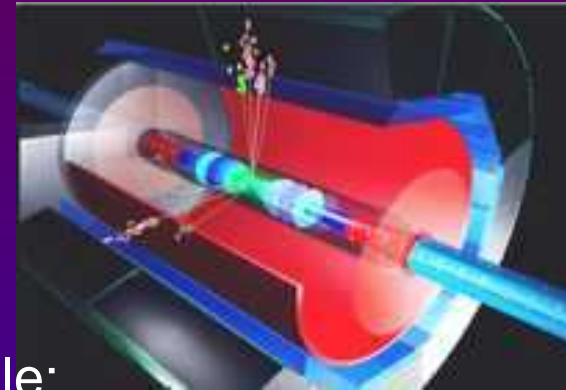
- Luminosity monitoring and precise parameter determination
- Disentangle limits on *New Physics* from SM predictions and backgrounds

Motivation: ILC and Bhabha physics

The Standard Model (SM) is very **successful**, but . . . present few “experiment vs. theory” $2\text{-}3\sigma$ disagreements (A_{FB}^b , $\sin \theta_{\text{eff}}$, N_ν, \dots) require **better control** of theoretical uncertainties.

Next International Linear Collider (ILC) will be a **challenge**

e^+e^- beam, $E_{\text{CM}} \approx 1\text{TeV}$, $\%_0$ precision



Bhabha scattering ($e^+e^- \rightarrow e^+e^-$) plays a dominant role:

- Luminosity monitoring and precise parameter determination
- Disentangle limits on *New Physics* from SM predictions and backgrounds

Massive effects important for heavy (t , b) fermions \rightarrow test of Higgs

Motivation: computing

We need **reliable** and **independent** theoretical predictions, especially for fermion processes, at the next colliders

based on different methods and codes.

Motivation: computing

We need **reliable** and **independent** theoretical predictions, especially for fermion processes, at the next colliders

based on different methods and codes.

The idea is to develop a tool for precise calculations. A must:

- Automatic **beyond** tree-level (1-loop or higher)
- **Free** software (also based upon)
- **Documented** and easy to **install**
- Profit from **available** good packages

Motivation: computing

We need **reliable** and **independent** theoretical predictions, especially for fermion processes, at the next colliders

based on different methods and codes.

The idea is to develop a tool for precise calculations. A must:

- Automatic **beyond** tree-level (1-loop or higher)
- **Free** software (also based upon)
- **Documented** and easy to **install**
- Profit from **available** good packages

☞ Do other tools in the market satisfy these 3-4 points? 😞

Packages in the market to be mentioned:

FEYNARTS, GRACE, SANC, COMPHEP, MADGRAPH ...

Motivation: Bhabha calculations

An incomplete historical summary . . .

Who?	When?	What?
Bhabha	1935	Tree-level QED
Readhead	1953	Leading 1-loop QED
Berends, Gaemers and Gastmans	1974	Complete 1-loop QED with hard γ
Consoli	1979	1-loop Electro-weak (approx. weak boxes)
Böhm, Denner, Hollik and Sommer	1984	Complete 1-loop Electro-weak
Bardin, Hollik and Riemann	1991	Leading weak 2-loop at Z-resonance
Jadach and others	90's	Precise Monte Carlo
Bern, Dixon and Ghinculov	2000	2-loops QED
Fleischer et al	2002	(1-loop) ² massive QED
Penin, others (Bonciani <i>et al</i> , Gluza <i>et al</i>)	2005	leading 2-loops massive QED (ILC precision)

Introduction: electroweak model

Electroweak interactions (Glashow-Salam-Weinberg 60's): local gauge group $G_{\text{local}}^{\text{EW}} = SU(2)_L \otimes U(1)_Y$ broken to $U(1)_{\text{em}}$ invariant \rightarrow QED

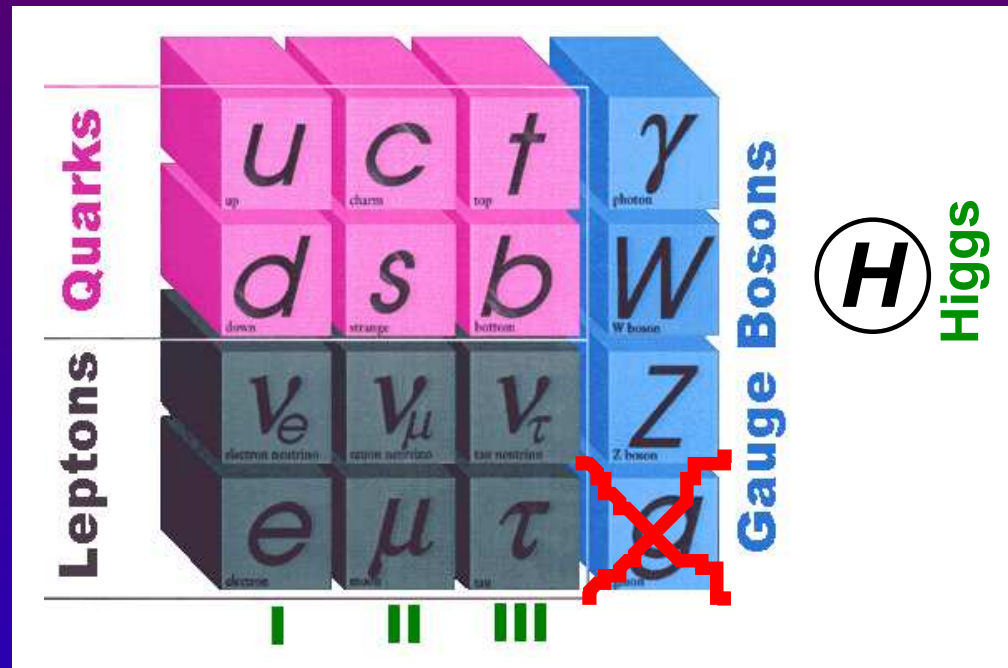
$$\mathcal{L} = \mathcal{L}_{\text{YM}} + \mathcal{L}_{\text{H}} + \mathcal{L}_{\text{F}} + \mathcal{L}_{\text{fix}} + \mathcal{L}_{\text{FP}}$$

Introduction: electroweak model

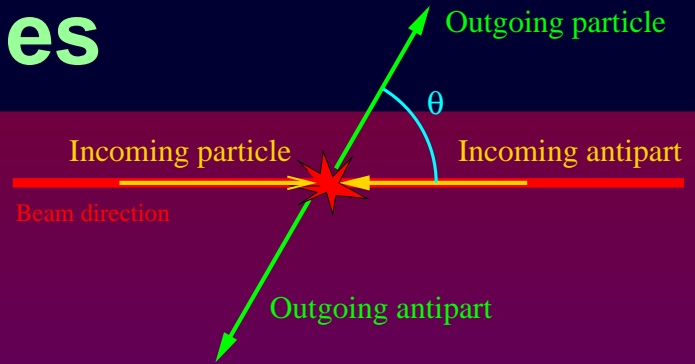
Electroweak interactions (Glashow-Salam-Weinberg 60's): local gauge group $G_{\text{local}}^{\text{EW}} = SU(2)_L \otimes U(1)_Y$ broken to $U(1)_{\text{em}}$ invariant \rightarrow QED

$$\mathcal{L} = \mathcal{L}_{\text{YM}} + \mathcal{L}_{\text{H}} + \mathcal{L}_{\text{F}} + \mathcal{L}_{\text{fix}} + \mathcal{L}_{\text{FP}}$$

Particle content and couplings \rightarrow Feynman Rules ('t Hooft '71)



Introduction: observables



In $2 \rightarrow 2$ fermion processes we calculate

- **Differential Cross Section** (pb) [l^2]

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{32\pi} \frac{\beta_{\text{out}}\beta_{\text{in}}^{-1}}{s} \sum_{\text{conf}} |\mathcal{M}(\cos\theta)|^2$$

- **Total** Cross Section:

$$\sigma_{\text{tot}} = \int_{-1}^1 d\cos\theta \frac{d\sigma}{d\cos\theta}$$

- **Forward–Backward** Asymmetry:

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_{\text{tot}}} = \frac{[\int_0^1 - \int_{-1}^0] d\cos\theta \frac{d\sigma}{d\cos\theta}}{\sigma_{\text{tot}}}$$

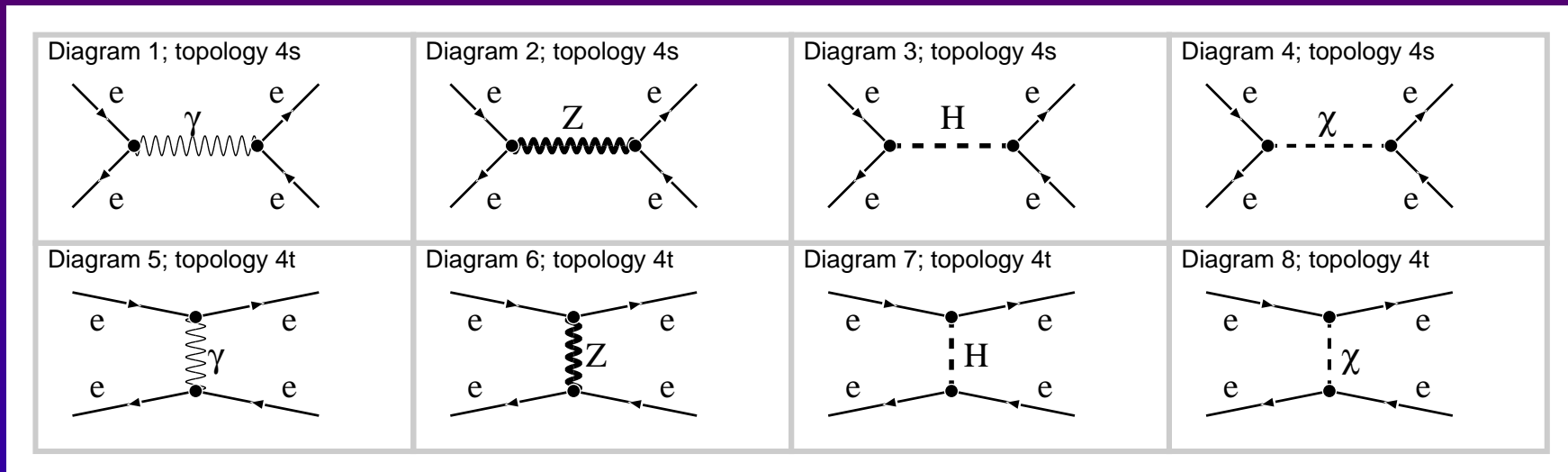
Introduction: perturbation theory

Thanks to the perturbative approach (Dyson, Feynman 50's), we represent the different contributions through Feynman diagrams.

Example: massive Bhabha scattering: $e^- e^+ \rightarrow e^- e^+$

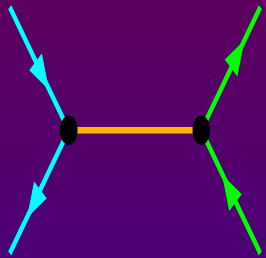
Vector

Scalar

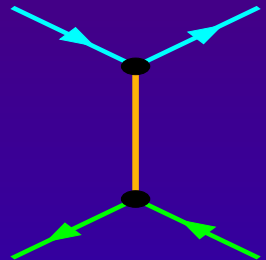


Introduction: perturbation theory

Thanks to the perturbative approach (Dyson, Feynman 50's), we represent the different contributions through Feynman diagrams.



$$\begin{aligned}
 \mathcal{M}_S^{(0)} = & \bar{v}_e \quad ie\gamma^\mu Q_e \quad u_e \quad \frac{-ig_{\mu\nu}}{s} \quad \bar{u}_e \quad ie\gamma^\nu Q_e \quad v_e \\
 & + \bar{v}_e ie\gamma^\mu (V_e - A_e\gamma_5) u_e \quad \frac{-ig_{\mu\nu}}{s-m_Z^2} \quad \bar{u}_e ie\gamma^\nu (V_e - A_e\gamma_5) v_e \\
 & + \bar{v}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \quad u_e \quad \frac{i}{s-m_H^2} \quad \bar{u}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \quad v_e \\
 & + \bar{v}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \gamma_5 \quad u_e \quad \frac{i}{s-m_Z^2} \quad \bar{u}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \gamma_5 \quad v_e
 \end{aligned}$$



$$\begin{aligned}
 \mathcal{M}_T^{(0)} = & \bar{u}_e \quad ie\gamma^\mu Q_e \quad u_e \quad \frac{-ig_{\mu\nu}}{t} \quad \bar{v}_e \quad ie\gamma^\nu Q_e \quad v_e \\
 & + \bar{u}_e ie\gamma^\mu (V_e - A_e\gamma_5) u_e \quad \frac{-ig_{\mu\nu}}{t-m_Z^2} \quad \bar{v}_e ie\gamma^\nu (V_e - A_e\gamma_5) v_e \\
 & + \bar{u}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \quad u_e \quad \frac{i}{t-m_H^2} \quad \bar{v}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \quad v_e \\
 & + \bar{u}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \gamma_5 \quad u_e \quad \frac{i}{t-m_Z^2} \quad \bar{v}_e \quad -i\frac{1}{2s_W} \frac{m_e}{m_W} \gamma_5 \quad v_e
 \end{aligned}$$

Coupling

Propag.

Coupling

Introduction: one-loop decomposition

Going to **next** perturbation **order** ...

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}^{(0)} + \mathcal{M}^{(1)} + \dots|^2 \\ &= \underbrace{\mathcal{M}^{(0)} \mathcal{M}^{(0)*}}_{\text{LO}(\alpha^2)} + \underbrace{2\Re(\mathcal{M}^{(1)} \mathcal{M}^{(0)*})}_{\text{NLO}(\alpha^3)} + \dots \end{aligned}$$

Introduction: one-loop decomposition

Going to **next** perturbation **order** ...

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}^{(0)} + \mathcal{M}^{(1)} + \dots|^2 \\ &= \underbrace{\mathcal{M}^{(0)} \mathcal{M}^{(0)*}}_{\text{LO}(\alpha^2)} + 2\Re(\underbrace{\mathcal{M}^{(1)} \mathcal{M}^{(0)*}}_{\text{NLO}(\alpha^3)}) + \dots \end{aligned}$$

The invariant amplitude \mathcal{M} is decomposed into a sum:

$$\mathcal{M} = \sum_i^{\text{complete set}} \mathbf{M}_i \left(\mathbf{F}_i^{(0)} + \mathbf{F}_i^{(1)} + \dots \right)$$

Introduction: one-loop decomposition

Going to **next** perturbation **order** ...

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}^{(0)} + \mathcal{M}^{(1)} + \dots|^2 \\ &= \underbrace{\mathcal{M}^{(0)} \mathcal{M}^{(0)*}}_{\text{LO}(\alpha^2)} + \underbrace{2\Re(\mathcal{M}^{(1)} \mathcal{M}^{(0)*})}_{\text{NLO}(\alpha^3)} + \dots \end{aligned}$$

The invariant amplitude \mathcal{M} is decomposed into a sum:

$$\mathcal{M} = \sum_i^{\text{complete set}} \mathbf{M}_i (\mathbf{F}_i^{(0)} + \mathbf{F}_i^{(1)} + \dots)$$

❗ Arbitrary basis of **Matrix Elements** (kinematics)

Introduction: one-loop decomposition

Going to **next** perturbation **order** ...

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}^{(0)} + \mathcal{M}^{(1)} + \dots|^2 \\ &= \underbrace{\mathcal{M}^{(0)} \mathcal{M}^{(0)*}}_{\text{LO}(\alpha^2)} + \underbrace{2\Re(\mathcal{M}^{(1)} \mathcal{M}^{(0)*})}_{\text{NLO}(\alpha^3)} + \dots \end{aligned}$$

The invariant amplitude \mathcal{M} is decomposed into a sum:

$$\mathcal{M} = \sum_i^{\text{complete set}} \mathbf{M}_i \left(\mathbf{F}_i^{(0)} + \mathbf{F}_i^{(1)} + \dots \right)$$

❖ Arbitrary basis of **Matrix Elements** (kinematics)

❖ Polarization independent, scalar **Form Factors** (dynamics)

Introduction: matrix elements

We use 9×4 elements for each channel (S , T or U) = 36

Example for the fermionic S -channel:

\mathbf{M}_S	$1,j$	$= \bar{v}_e$	$\{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$2,j$	$= \bar{v}_e$	$\not{p}_a \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$3,j$	$= \bar{v}_e$	$\{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\not{p}_b \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$4,j$	$= \bar{v}_e$	$\not{p}_a \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\not{p}_b \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$5,j$	$= \bar{v}_e$	$\gamma^\mu \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\gamma_\mu \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$6,j$	$= \bar{v}_e$	$\gamma^\mu \not{p}_a \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\gamma_\mu \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$7,j$	$= \bar{v}_e$	$\gamma^\mu \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\gamma_\mu \not{p}_b \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$8,j$	$= \bar{v}_e$	$\gamma^\mu \not{p}_a \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\gamma_\mu \not{p}_b \{\mathbb{1}, \gamma_5\}_j$	v_f
\mathbf{M}_S	$9,j$	$= \bar{v}_e$	$\gamma^\mu \gamma^\nu \{\mathbb{1}, \gamma_5\}_j$	$u_e \otimes \bar{u}_f$	$\gamma_\mu \gamma_\nu \{\mathbb{1}, \gamma_5\}_j$	v_f

Introduction: renormalization

It turns out that some **loop** integrals are **divergent** !

- Ultra-Violet: High energy limit in virtual momenta
- Infra-Red: Massless photon between external legs

How can we still make **predictions**?

Introduction: renormalization

It turns out that some **loop** integrals are **divergent** !

- Ultra-Violet: High energy limit in virtual momenta
- Infra-Red: Massless photon between external legs

How can we still make **predictions**?

DIMENSIONAL REGULARIZATION :

$$\epsilon = \frac{d-4}{2}$$

Shift the Minkowski space-time \rightarrow Convergent integrals

Introduction: renormalization

It turns out that some **loop** integrals are **divergent** !

- Ultra-Violet: High energy limit in virtual momenta
- Infra-Red: Massless photon between external legs

How can we still make **predictions**?

DIMENSIONAL REGULARIZATION :

$$\epsilon = \frac{d-4}{2}$$

Shift the Minkowski space-time \rightarrow Convergent integrals

ON-SHELL RENORMALIZATION :

$$\psi_0 = Z_\psi^{1/2} \psi, \quad \alpha_0 = Z_\alpha \alpha$$

Redefinition of bare parameters and fields, absorbing ϵ

Introduction: renormalization

It turns out that some **loop** integrals are **divergent** !

- Ultra-Violet: High energy limit in virtual momenta
- Infra-Red: Massless photon between external legs

How can we still make **predictions**?

DIMENSIONAL REGULARIZATION :

$$\epsilon = \frac{d-4}{2}$$

Shift the Minkowski space-time \rightarrow Convergent integrals

ON-SHELL RENORMALIZATION :

$$\psi_0 = Z_\psi^{1/2} \psi, \quad \alpha_0 = Z_\alpha \alpha$$

Redefinition of bare parameters and fields, absorbing ϵ

 Only **physical** parameters and fields plus **counter-terms** appear

Introduction: photon emission

X Real detectors cannot observe soft photons

Introduction: photon emission

- ✗ Real detectors cannot observe soft photons
- ☞ Photon emission mixes **incoherently** with pure processes

Introduction: photon emission

- ✗ Real detectors cannot observe soft photons
- ☞ Photon emission mixes **incoherently** with pure processes
- ✓ Cancellation of **InfraRed** singularities from external self-energies, vertices and boxes in one-loop integrals

Introduction: photon emission

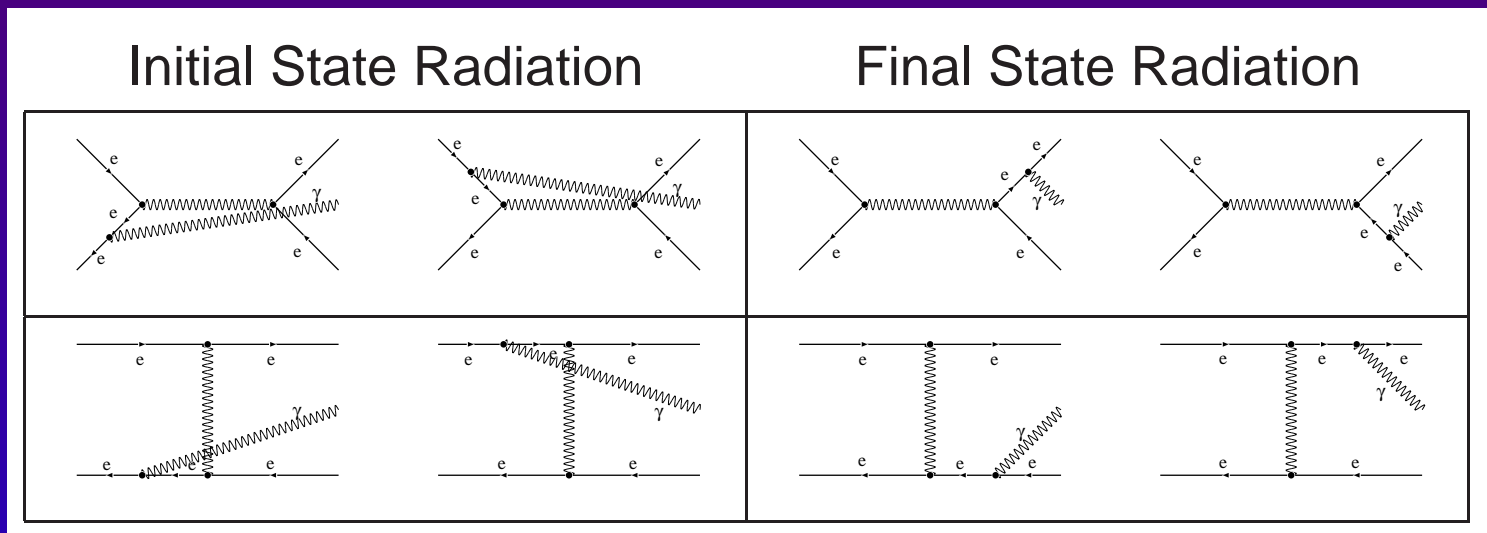
- ✗ Real detectors cannot observe soft photons
- ☞ Photon emission mixes **incoherently** with pure processes
- ✓ Cancellation of **InfraRed** singularities from external self-energies, vertices and boxes in one-loop integrals

SOFT Contribution:
$$\frac{d\sigma}{d\cos\theta} \Big|_{\text{Soft}} = \text{Factor}(\gamma_{\text{soft}}) \frac{d\sigma}{d\cos\theta} \Big|_{\text{Born}}$$

Introduction: photon emission

- ✗ Real detectors cannot observe soft photons
- ☞ Photon emission mixes **incoherently** with pure processes
- ✓ Cancellation of **InfraRed** singularities from external self-energies, vertices and boxes in one-loop integrals

SOFT Contribution: $\frac{d\sigma}{d\cos\theta} \Big|_{\text{Soft}} = \text{Factor}(\gamma_{\text{soft}}) \frac{d\sigma}{d\cos\theta} \Big|_{\text{Born}}$



Introduction: on the Z-peak

Strict one-loop corrections in $2 \rightarrow 2$ fermions lead to

$\sigma(s = m_Z^2) \rightarrow \infty$, coming from lowest order Z-boson propagator

$$\frac{1}{s - m_Z^2}$$

Introduction: on the Z-peak

Strict one-loop corrections in $2 \rightarrow 2$ fermions lead to

$\sigma(s = m_Z^2) \rightarrow \infty$, coming from lowest order Z-boson propagator

$$\frac{1}{s - m_Z^2}$$

In order to have $\mathcal{O}(\alpha)$ accuracy beyond IBA we implemented *fixed width* scheme:

- **Dyson** summation $\frac{1}{s - m_Z^2} \rightarrow \frac{1}{s - m_Z^2 + i\Gamma_Z m_Z}$
- Discard self-energies topologies from 1-loop amplitude
- Ensure **Infrared** finiteness when adding soft-photon emission

Introduction: IR-finiteness on the Z-peak

The infrared cancellation reads

The diagram shows the infrared cancellation in electron-muon scattering at the Z-peak. It is represented as an equation: $IR \left(\text{tree} + \text{radiative} + \text{loop} \right) = 0$. The tree-level diagram shows an electron and muon interacting via a Z boson. The radiative correction term consists of two diagrams: one with a photon radiated from the electron line and another with a photon radiated from the muon line. The loop correction term is a one-loop diagram with a photon and a Z boson in the loop. The entire expression is enclosed in large parentheses and set equal to zero.

Introduction: IR-finiteness on the Z-peak

The infrared cancellation reads

The diagram shows the infrared cancellation equation: $IR \left(\text{tree} + \text{tree} \times \text{radiative} + \text{tree} \times \text{loop} \right) = 0$. The tree-level diagram is $e^+e^- \rightarrow Z \rightarrow \mu^+\mu^-$. The first radiative correction is $e^+e^- \rightarrow Z \rightarrow \mu^+\mu^- + \gamma$ (real emission). The loop correction is a Z -boson loop. The cancellation is shown as the sum of these terms multiplied by the infrared factor IR equals zero.

Minimally we need

- D_0 with complex mass argument
Beenakker and Denner. Nucl. Phys. B338 (1990)
- C_0, B_0, A_0 with arbitrary complex arguments (due to the reduction to masters of projected $D_{i,ij}$)
't Hooft and Veltman. Nucl. Phys. B153 (1979)



II. Automated tool: *a*ĪTALC

Automation: aTALC overview

an Integrated Tool for Automated Loop Calculations

Automation: aTALC overview

an Integrated Tool for Automated Loop Calculations

- Restricted to automated $2 \rightarrow 2$ fermions (EWSM and QED)
- GNU/LINUX tool, GPL licensed, **free available** since Oct'04
- <http://www-zeuthen.desy.de/theory/aitalc>
- To be accepted CPC: A.L and T.Riemann. [hep-ph/0412047](http://arxiv.org/abs/hep-ph/0412047)

Automation: aTALC overview

an Integrated Tool for Automated Loop Calculations

- Restricted to automated $2 \rightarrow 2$ fermions (EWSM and QED)
- GNU/LINUX tool, GPL licensed, **free available** since Oct'04
- <http://www-zeuthen.desy.de/theory/aitalc>
- To be accepted CPC: A.L and T.Riemann. [hep-ph/0412047](http://arxiv.org/abs/hep-ph/0412047)

Three structural blocks:

Diagram
generation

DIANA 2.35
(QGRAF)

Algebra
simplification

FORM 3.1

Numerical
evaluation

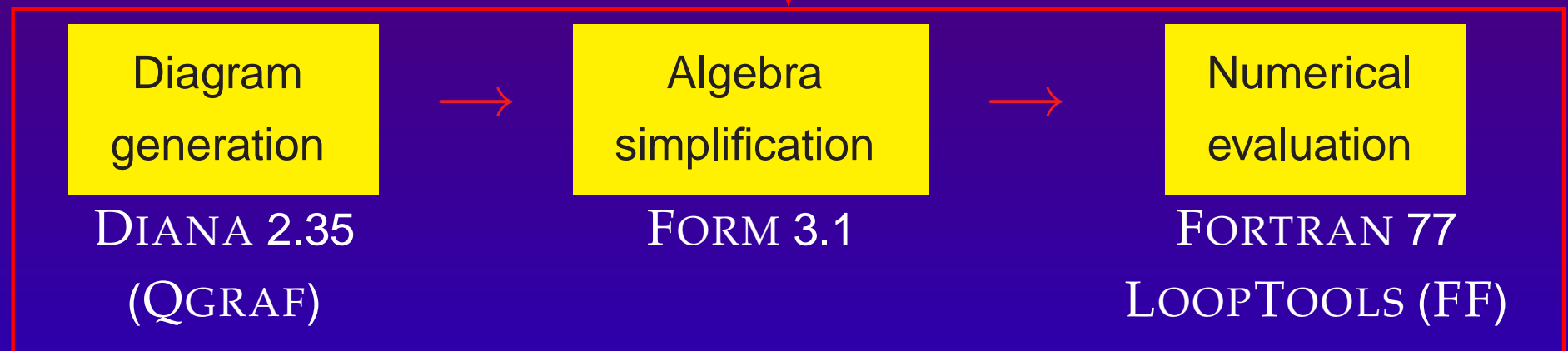
FORTRAN 77
LOOPTOOLS (FF)

Automation: aTALC overview

an Integrated Tool for Automated Loop Calculations

- Restricted to automated $2 \rightarrow 2$ fermions (EWSM and QED)
- GNU/LINUX tool, GPL licensed, **free available** since Oct'04
- <http://www-zeuthen.desy.de/theory/aitalc>
- To be accepted CPC: A.L and T.Riemann. [hep-ph/0412047](http://arxiv.org/abs/hep-ph/0412047)

Three structural blocks: all running under **MAKE** environment



Automation: Feynman Diagram Analyzer DIAN

Developed at U.Bielefeld 1997-2004 ([Fleischer and Tentyukov](#))

Automation: Feynman Diagram Analyzer DIAN

Developed at U.Bielefeld 1997-2004 (Fleischer and Tentyukov)

- C program, based on Nogueira's FORTRAN generator QGRAF2
- Command line: requires a driver file and model file
- High portability, running in many UNIX systems
- Front-end topology editor (tedi) included for GNU/LINUX

<http://www.physik.uni-bielefeld.de/~tentukov/diana.html>

Automation: Feynman Diagram Analyzer DIAN

Developed at U.Bielefeld 1997-2004 (Fleischer and Tentyukov)

What do we ask?

```
SET _processname = Bhabha
```

```
\Begin(model,EWSM.model)
```

```
\Begin(process)
```

```
ingoing le(;p1),Le(;p4);
```

```
outgoing le(-p2),Le(-p3);
```

```
loops = 1;
```

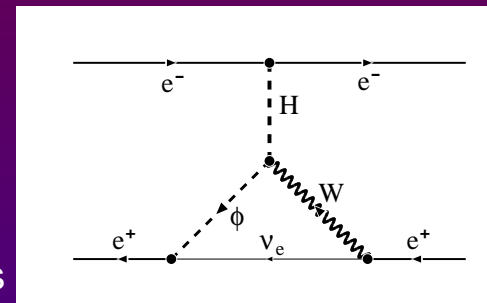
```
options = onshell,notadp;
```

```
* \excludevertex(Le,le,H)
```

```
SET MakeEps = "!"
```

```
...
```

What does Diana answer?



Bhabha626.eps

G Amplitude =

$$(-1)*F(1,1,1,0,0)*(-i_-)*e/2/sw*Mle/MW*F(2,2,1,-1,0)*$$

$$(-i_-)*e/2/sqrt2/sw*Mle/MW*FF(3,2,+q,Mne)*i_-*$$

$$F(3,2,mu1,1,-1,1)*(+i_-)*e/2/sqrt2/sw*SS(4,0)*i_-*$$

$$SS(1,2)*i_-*VV(2,mu2,mu1,-q-k2,2)*i_-*$$

$$V(4,mu2,+p1+p2-(+q+k1),1)*(-i_-)*e/2/sw;$$

```
#define COUNTER "626" #define LINE "4"
```

```
#define LOOPTYPE "c" ...
```

Automation: aTALC algebra



DIANA
(symbolic level)

??? 1-Loop Library ???



FORTRAN
(numeric level)

Written in FORM

```
#call feynmanrules()
...
#call tracefermiloops()
#call integration()
#call chisholm()
#call dimensionfour()
#call gammaalgebra()
#call onshell()
#call diracequation()
#call massiveformfactors()
.end
```

These general procedures perform all algebra simplifications

✓ Write **automatically** FORTRAN subroutines from DIANA output

Automation: numerics with aTALC

- For numerical evaluation language FORTRAN 77 is used

Automation: numerics with aTALC

For numerical evaluation language FORTRAN 77 is used

The code is decomposed into different routines

- **Local:** Process-dependent automatically generated (me, ff)
- **Global:** Fixed coming with the distribution (renorm.)
- **External:** LOOPTOOLS package (evaluation of loop integrals)

Automation: numerics with aTALC

For numerical evaluation language FORTRAN 77 is used

The code is decomposed into different routines

- **Local:** Process-dependent automatically generated (me, ff)
- **Global:** Fixed coming with the distribution (renorm.)
- **External:** LOOPTOOLS package (evaluation of loop integrals)

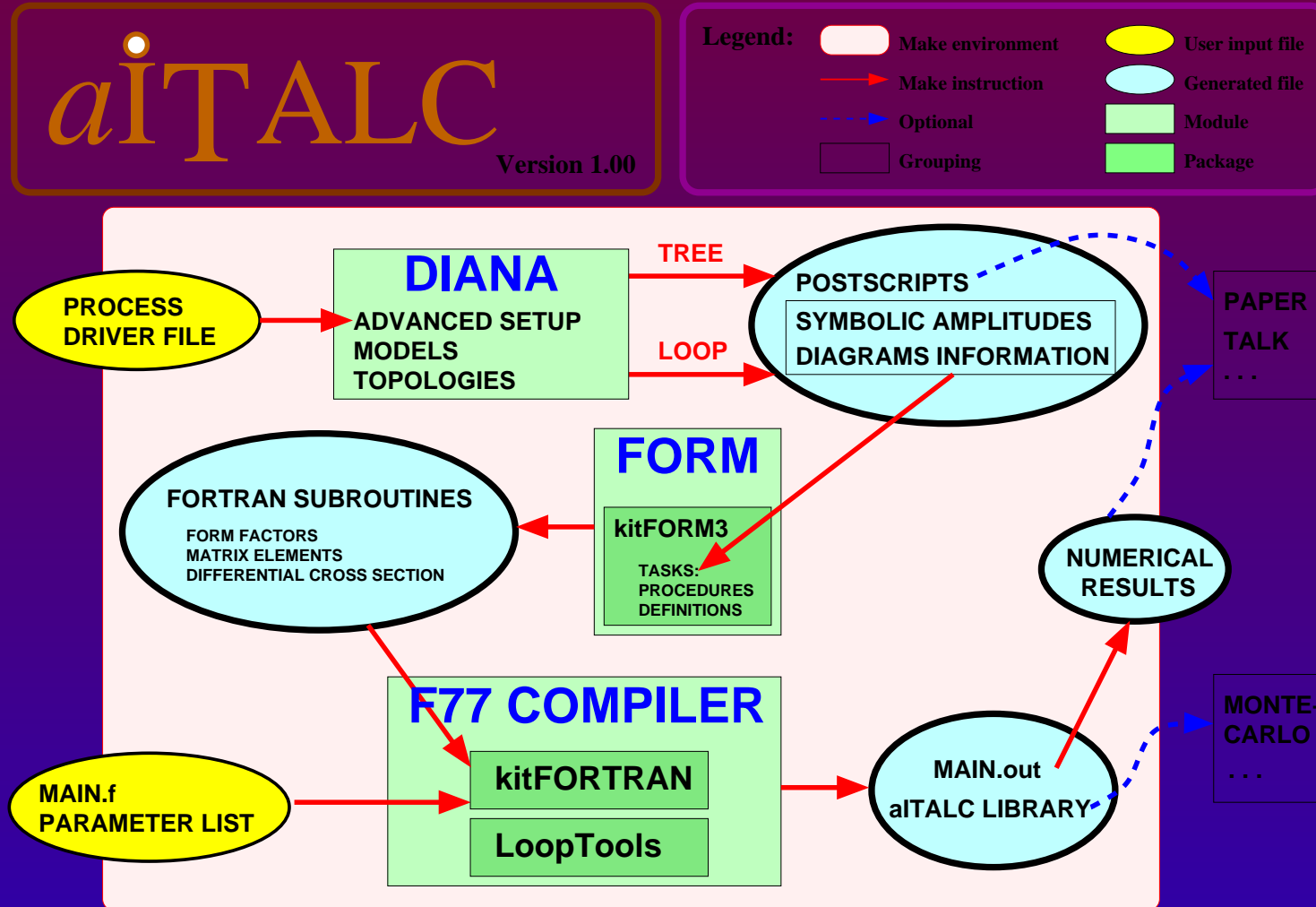
Executable file `main.out`

Input → parameter list, control flags.

Output ← tables for **differential** and **integrated** cross sections and **forward-backward** asymmetries

Tests ✓ ultraviolet and infrared finiteness against parameter variation.
Quadruple precision

Automation: aITALC Flowchart

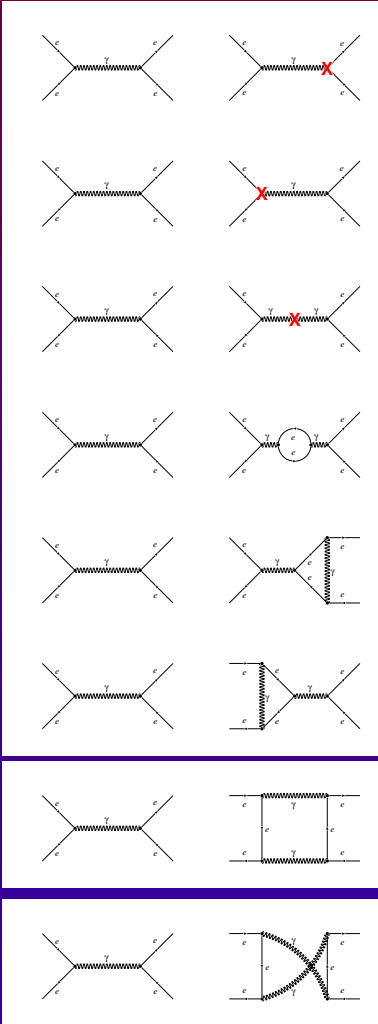




III. Results

Results: Bhabha analytical

$a\dot{\Gamma}$ ALC helped in getting analytical results: massless QED



$$e^4 \frac{t^2+u^2}{s^2} \left(2\delta Z_e + \delta Z_{AA} + 2\delta Z_e^{f,V} \right),$$

$$e^4 \frac{t^2+u^2}{s^2} \left(2\delta Z_e + \delta Z_{AA} + 2\delta Z_e^{f,V} \right),$$

$$e^4 \frac{t^2+u^2}{s^2} \left(-2\delta Z_{AA} \right),$$

$$\frac{e^6}{\pi^2} \frac{t^2+u^2}{s^2} \left(+\frac{1}{18} - \frac{1}{6} B_s^e \right),$$

$$\frac{e^6}{\pi^2} \frac{t^2+u^2}{s^2} \left(-\frac{1}{4} + \frac{1}{2} B_e - \frac{3}{8} B_s^e - \frac{1}{4} \tilde{C}_s^e \right),$$

$$\frac{e^6}{\pi^2} \frac{t^2+u^2}{s^2} \left(-\frac{1}{4} + \frac{1}{2} B_e - \frac{3}{8} B_s^e - \frac{1}{4} \tilde{C}_s^e \right),$$

$$\frac{e^6}{\pi^2} \left(-\frac{1}{4} \frac{u}{s} (B_s^\gamma - B_t^e) - \frac{1}{4} \frac{t-u}{s} \tilde{C}_t^e + \frac{1}{4} \frac{3t^2+u^2}{s^2} \tilde{C}_s^\gamma - \frac{1}{8} \frac{s^2+u^2}{st} \tilde{D}_{ts} \right),$$

$$\frac{e^6}{\pi^2} \left(\frac{1}{4} \frac{t}{s} (B_s^\gamma - B_u^e) - \frac{1}{4} \frac{t-u}{s} \tilde{C}_u^\gamma - \frac{1}{8} \frac{t^2+3u^2}{s^2} (2\tilde{C}_s^\gamma - \tilde{D}_{su}) \right),$$

Results: Bhabha comparisons

$$e^-e^+ \rightarrow e^-e^+ (\gamma) \text{ at ILC: } \sqrt{s} = 500 \text{ GeV}, \quad E_{\max}(\gamma_{\text{soft}}) = \frac{\sqrt{s}}{10}$$

$\cos\theta$	$\left[\frac{d\sigma}{d\cos\theta}\right]_{\text{Born}}$ (pb)	$\left[\frac{d\sigma}{d\cos\theta}\right]_{\mathcal{O}(\alpha^3)=\text{Born+QED+weak+soft}}$ (pb)	Tool
-0.9	0.21699 88288 10920	0.19344 50785 26862 70315 89...	<i>a</i> ITALC
-0.9	0.21699 88288 10920	0.19344 50785 26862	FEYNARTS
-0.9	0.21699 88288 41513	0.19344 50785 62638	$m_e = 0$
+0.0	0.59814 23072 50331	0.54667 71794 69423 03528 77...	<i>a</i> ITALC
+0.0	0.59814 23072 50329	0.54667 71794 69422	FEYNARTS
+0.0	0.59814 23072 88584	0.54667 71794 99961	$m_e = 0$
+0.9	0.18916 03223 32271 · 10 ³	0.17292 83490 66508 29307 47... · 10 ³	<i>a</i> ITALC
+0.9	0.18916 03223 32271 · 10 ³	0.17292 83490 66508 · 10 ³	FEYNARTS
+0.9	0.18916 03223 31849 · 10³	0.17292 83490 61347 · 10³	$m_e = 0$
+0.9999	0.20842 90676 46391 · 10 ⁹	0.19140 17861 11883 04292 09... · 10 ⁹	<i>a</i> ITALC
+0.9999	0.20842 90676 464 36 · 10⁹	0.19140 17861 11 979 · 10⁹	FEYNARTS

Great independent agreement saturating **limit in double precision**

Thanks to T. Hahn for supplying FEYNARTS' numbers

Results: Bhabha within different models

Forward region, SABH (Small Angle Bhabha scattering)

$$e^-e^+ \rightarrow e^-e^+ (\gamma) \text{ at ILC: } \sqrt{s} = 500 \text{ GeV, } E_{\text{max}}(\gamma_{\text{soft}}) = \frac{\sqrt{s}}{10}$$

rad	$\cos \theta$	Born EWSM	Born QED	$\mathcal{O}(\alpha)$ EWSM	$\mathcal{O}(\alpha)$ QED $N_f=9$	$\mathcal{O}(\alpha)$ QED $N_f=1$
0.451	+0.9000	$1.89160 \cdot 10^2$	-0.0222%	-8.58%	-6.71%	-15.60%
0.142	+0.9900	$2.06556 \cdot 10^4$	-0.0840%	-7.72%	-7.16%	-14.06%
0.045	+0.9990	$2.08236 \cdot 10^6$	0.0031%	-7.98%	-7.53%	-12.50%
0.014	+0.9999	$2.08429 \cdot 10^8$	0.0005%	-8.17%	-7.75%	-11.02%

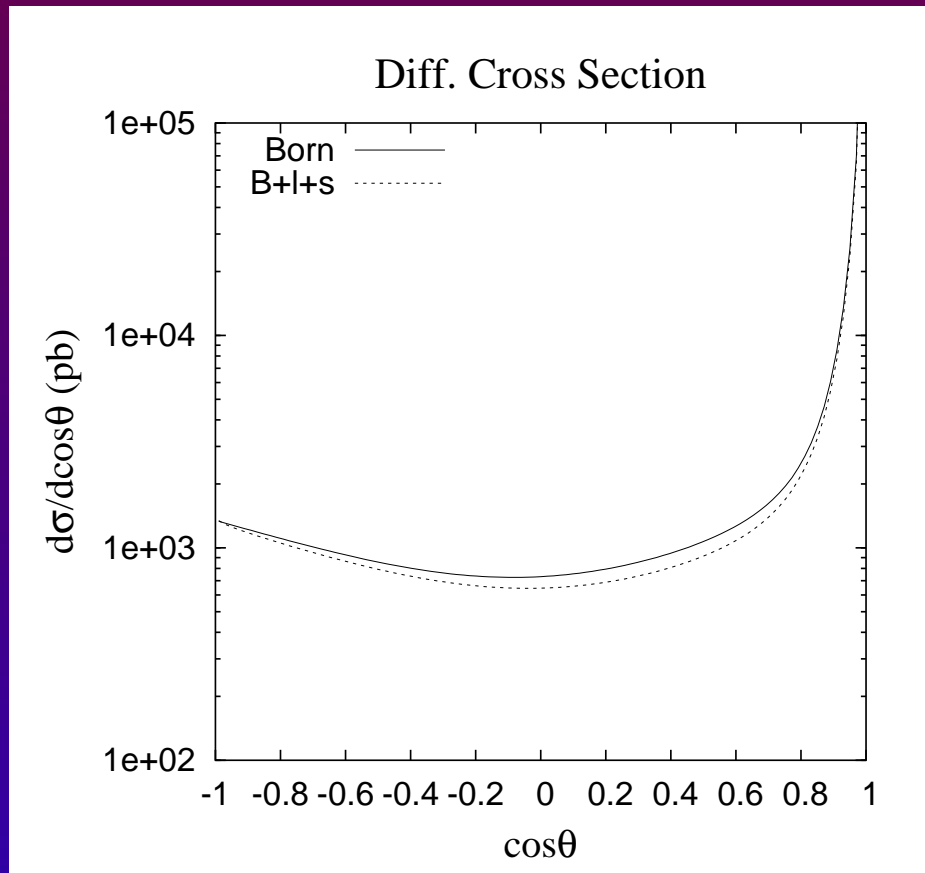
Precision required at ILC for SABH is achieved incorporating

fully flavour QED two-loops corrections to the one-loop EWSM ones.

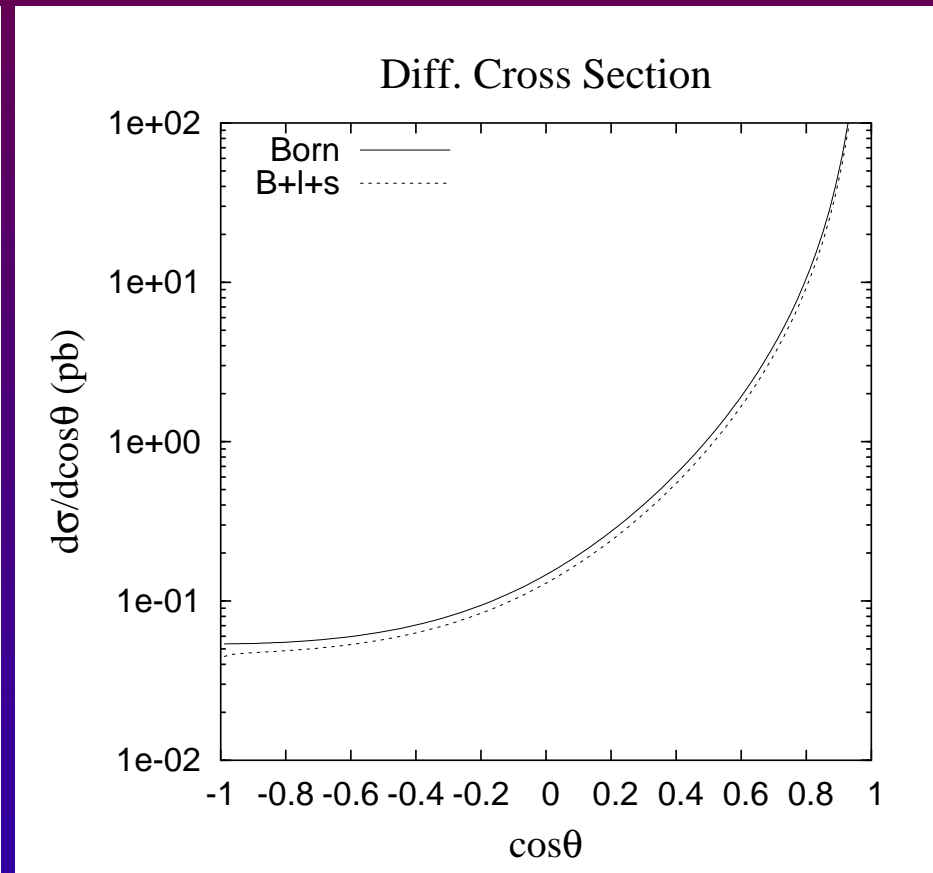
Results: some plots

Differential cross sections

(a) $\sqrt{s} = m_Z$



(b) $\sqrt{s} = 1000$ GeV



Conclusions & Outlook

- Complete $\mathcal{O}(\alpha)$ electroweak corrections to Bhabha scattering (also other different $2 \rightarrow 2$ fermion processes not shown)
 - ▶ Resonances and masses included
 - ▶ Other contributions still required (hard γ , QCD, kin. cuts...)

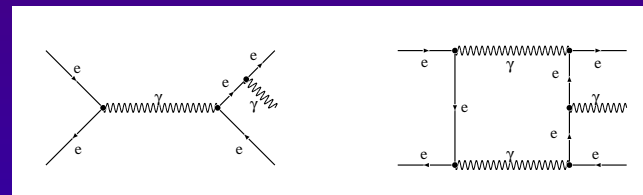
Conclusions & Outlook

- Complete $\mathcal{O}(\alpha)$ electroweak corrections to Bhabha scattering (also other different $2 \rightarrow 2$ fermion processes not shown)
 - ▶ Resonances and masses included
 - ▶ Other contributions still required (hard γ , QCD, kin. cuts...)
- ✓ We show the capability of $a^{\circ}\text{TALC}$
 - ▶ Fully automated, Free available and Tested
 - ▶ Potential extension to other models (QCD or SUSY), to *parallelization* or to higher loops through DIANA

Conclusions & Outlook

- Complete $\mathcal{O}(\alpha)$ electroweak corrections to Bhabha scattering (also other different $2 \rightarrow 2$ fermion processes not shown)
 - ▶ Resonances and masses included
 - ▶ Other contributions still required (hard γ , QCD, kin. cuts...)
- ✓ We show the capability of $a\dot{I}TALC$
 - ▶ Fully automated, Free available and Tested
 - ▶ Potential extension to other models (QCD or SUSY), to *parallelization* or to higher loops through DIANA
- ♣ Ongoing project related to 2-loops Bhabha:

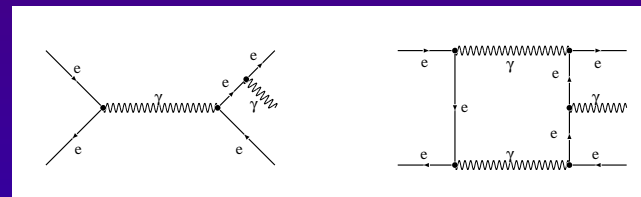
Radiative Bhabha scattering:



Conclusions & Outlook

- Complete $\mathcal{O}(\alpha)$ electroweak corrections to Bhabha scattering (also other different $2 \rightarrow 2$ fermion processes not shown)
 - ▶ Resonances and masses included
 - ▶ Other contributions still required (hard γ , QCD, kin. cuts...)
- ✓ We show the capability of $a\dot{I}TALC$
 - ▶ Fully automated, Free available and Tested
 - ▶ Potential extension to other models (QCD or SUSY), to *parallelization* or to higher loops through DIANA
- ♣ Ongoing project related to 2-loops Bhabha:

Radiative Bhabha scattering:



? Should we join into a project in loop calculations for colliders?