Heavy Ion Physics at the LHC Santa Fe, October 23 Santa Fe, October 23rd, 2005

Heavy Quarks Heavy Quarks in view of the LHC in view of the LHC

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Layout

High- $\rho_{\sf T}$ (heavy-flavour) particle production in hadronic collisions –from pp to AA– in factorized QCD \bullet **pp baseline: pQCD calculations** nuclear effects, in initial and final state

What have we learnt at Tevatron and RHIC?

- What will be new at the LHC?
- What do we expect?
	- What should we measure?

PARTON

ENERG

LOSS

Heavy-flavour production: pp

proton-proton collisions: factorized pQCD approach

 \boldsymbol{calc} ulable as perturbative series of strong coupling $\alpha_{\rm s}(\mu_{\mathsf{R}})$

pp [→] *QQ: pQCD calculations vs data*

Calculation of partonic cross section $\hat{\sigma}$

standard: Fixed Order (NLO) Massive (e.g. MNR code)

$$
\frac{d\sigma}{dp_T} = A(m)\alpha_s^2 + B(m)\alpha_s^3 + O(\alpha_s^4) \qquad B(m) = \beta(m) + \gamma(m)\log(\mu/m)
$$

state-of-the-art: Fixed Order Next-to-Leading Log (FONLL) $\mu \approx p_{\rm T}$

$$
\frac{d\sigma}{dp_T} = A(m)\alpha_s^2 + B(m)\alpha_s^3 + G(m, p_T) \left[\alpha_s^2 \sum_{i=2}^{\infty} a_i [\alpha_s \log(\mu/m)]^i + \alpha_s^3 \sum_{i=1}^{\infty} b_i [\alpha_s \log(\mu/m)]^i\right]
$$

B production at Tevatron (1.96 TeV) Charm production slightly underpredicted a 10⁻⁴ Form production slightly underpredicted at Tevatron superinter of the Corporation strategy of the Charm production slightly under production superinter of tevatron (1.96 Te

Cacciari, Frixione, Mangano, Nason and Ridolfi, JHEP0407 (2004) 033

 \rightarrow more accurate at high $\rho_{\rm T}$

Heavy-flavour production: AA

Calculating Parton Energy Loss

 $\overline{\mathcal{L}}$ ⎨ \int ≥ $\propto \alpha_{\rm c} C_{\rm b} \sqrt{\omega_{\rm c}/\omega}$ for ω < *c c c c c c c c c* $\int_S C_R$ *I* ω/ω for $\omega \geq \omega$ ω/ω for $\omega < \omega$ α ω ω $(\omega_{c}/\omega)^{2}$ for / ω for d d 2 *BDMPS-Z formalism*Radiated-gluon energy distrib.: λ 2 ˆ*T k* $\hat{q} = \frac{\sqrt{q}}{q}$ transport coefficient (BDMPS case)

$$
C_R
$$

$$
\omega_c = \hat{q}L^2/2
$$

$$
R = \omega_c L
$$

sets the scale of the radiated energy related to constraint $k_{\text{T}} < \omega_\text{v}$ controls shape at ω *<<* ^ω*c*Casimir coupling factor: 4/3 for q, 3 for g

Baier, Dokshitzer, Mueller, Peigne', Schiff, NPB 483 (1997) 291.

Zakharov, JTEPL 63 (1996) 952.

Salgado, Wiedemann, PRD 68(2003) 014008.

Model vs RHIC "light" data

◆ Model: pQCD + E loss probability + detailed collision geometry

Density $(\hat{q}$) "tuned" to match R_{AA} in central Au-Au at 200 GeV

Small effect of mass for charm (~50% for D, ~30% for e) at low $\rho_{\!\scriptscriptstyle{T}}$ [large uncertainties!] **Basically no effect in "safe"** p_T **-region**

Armesto, Dainese, Salgado, Wiedemann, PRD 71 (2005) 054027.

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c + b (?) decay e[±] *RAA at RHIC*

FONLL: Electron spectrum may be \sim 50% charm + \sim 50% beauty for 3 $<\rho_{\rm T}$ $<$ 8 GeV

Due to larger mass of b quark electron R_{AA} incresed to ~0.4 (mass uncertainty also studied)

Heavy-flavour data in Au-Au 200 GeV

 $R_{\rm AA}$ down to 0.3 for $p_{\rm T}$ > 4 GeV/c! Heavy-quark quenching. Similar to that of light! Small room for mass effect …

Data lower than energy loss calculations, but charm fraction may be higher than predicted by FONLL

Armesto, Dainese, Salgado, Wiedemann, PRD 71 (2005) 054027 + w/Cacciari, in preparation

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Layout

- High- $\rho_{\sf T}$ (heavy-flavour) particle production in hadronic collisions –from pp to AA– in factorized QCD \bullet pp baseline: pQCD calculations
	- nuclear effects, in initial and final state

What have we learnt at Tevatron and RHIC?

- **What will be new at the LHC?**
- ◆ What do we expect?

Novelties at LHC (1): Hard Probes

- LHC: large hard cross sections!
- Our set of "tools" (probes) becomes richer

• **heavy quarks**

$$
\sigma_{LHC}^{c\bar{c}} \approx 10 \times \sigma_{RHIC}^{c\bar{c}}
$$

$$
\sigma_{LHC}^{b\bar{b}} \approx 100 \times \sigma_{RHIC}^{b\bar{b}}
$$

- \bullet γ + jet correlations
- \bullet Z⁰ + jet correlations

Heavy-quark production at the LHC

- ◆ pp: Important test of pQCD in a new energy domain
- Remember the "15-years saga of b production at Tevatron"*
- Baseline predictions: NLO (MNR code) in pp + binary scaling (shadowing included for PDFs in the Pb)
- ALICE baseline for charm / beauty:

Theoretical uncertainty of a factor 2—3 (next slide)

* M.Mangano

MNR code: Mangano, Nason, Ridolfi, NPB373 (1992) 295.

Theoretical Uncertainties (HERA-LHC Workshop)

Evaluation of theoretical uncertainties

MNR code: Mangano, Nason, Ridolfi, NPB373 (1992) 295.

Model Comparisons (HERA-LHC Workshop)

Compare predictions by several different models

 \Rightarrow Good agreement between collinear factorization based calculations: FO NLO and FONLL

 k_T factorization (CASCADE) higher at large ρ_T

Energy extrapolation via pQCD?

- Different systems (pp, p-Pb, Pb-Pb) will have different √*^s* values
- Results in pp at 14 TeV will have to extrapolated to 5.5 TeV (Pb-Pb energy) to compute, e.g., nuclear modification factors R_{AA}
- pQCD: "there ratio of results at 14 TeV/5.5 TeV has 'small' uncertainty"

Novelties at LHC (2): Small x

Probe unexplored small-*^x* region with HQs at low $\rho_{\sf T}$ and/or forward *y*

 \triangleq down to x ~ 10⁻⁴ with charm already at $y=0$

• Window on the rich phenomenology of high-density PDFs:

gluon saturation / recombination effects

Eskola, Kolhinen, Vogt, PLB582 (2004) 157 • Possible effect: enhancement of charm production at **low** *p***T** w.r.t. to standard DGLAP-based predictions, even in pp!

Gotsmann, Levin, Maor, Naftali, hep-ph/0504040

Probing nuclear initial state PDFs with HQs

- **Shadowing in pA (AA)**
- ◆ CGC in pA (AA)
- Double parton scattering in pA

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Probing nuclear initial state PDFs with HQs

- ◆ Shadowing in pA (AA)
- **CGC in pA (AA)**
- Double parton scattering in pA

Saturation scale $Q_s^2(x) \sim xg(x)A/R_a^2 \sim xg(x)A^{1/3}$

Low Energy

- At LHC for $x \sim 10^{-4}$, $Q_s \sim 1.5$ -2 GeV $> m_c$
- For m_{To} -Q_s, charm prod. CGC-dominated:
- \bullet scales with $\sqrt{\mathsf{N}_{\mathsf{part}}}$ in pA (not $\mathsf{N}_{\mathsf{coll}}$)
- \bullet harder p_T spectra, since typical $\mathsf{k}_\mathsf{T}\text{-}\mathsf{Q}_\mathsf{s}\text{-}$ 1.5 GeV, while
- in standard factorization $k_T \sim \Lambda_{\text{QCD}} \sim 0.2 \text{ GeV}$

Kharzeev, Tuchin, NPA 735 (2004) 248

Probing nuclear initial state PDFs with HQs

- ◆ Shadowing in pA (AA)
- CGC in pA (AA)

Double parton scattering in pA

probe "many-body" PDFs

normal and anomalous: different A dep.

predicted rate: cccc/cc ~ 10% (Treleani et al.)

signature: events with "tagged" DD (can use $D^0 + e^+$ or e^+e^+) and ch. conj. NB: there is a "background" from normal $b\overline{b}$ events, but is can be estimated from measured single inclusive b cross section

Cattaruzza, Del Fabbro, Treleani, PRD 70 (2004) 034022

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Heavy Quark Energy Loss at LHC

- ◆ ~100 cc pairs and ~5 bb pairs per central Pb-Pb collision
- Experiments will measure with good precision $R_{\sf AA}$ for D \sf and B, and for their decay leptons

What can we learn from a comparative quenching study of massive and massless probes at the LHC?

Heavy Flavour RAA at LHC

Baseline: PYTHIA, with EKS98 shadowing, tuned to reproduce c and b $\rho_{\sf T}$ distributions from NLO pQCD (MNR)

(*m/E*)-dep. *E* loss with ${\hat{q}}_{LHC} \approx 7^{*} \times {\hat{q}}_{RHIC} = 25 \div 100 \, \text{GeV}^{2}\text{/fm}$

MNR: Mangano, Nason, Ridolfi, NPB 373 (1992) 295.

NPB 570 (2000) 379.

Azimuthal asymmetry (v₂)

- The azimuthal asymmetry --v $_{2}$ or $\mathsf{R}_{\mathsf{AA}}(\phi)$ -- of D and B mesons in *non-central* collisions tests:
	- at low/moderate $\rho_{\sf T}$: recombination scenario, $\mathsf{v}_{\sf 2}$ of c/b quarks, hence degree of thermalization of medium
	- \bullet at higher p_T : path-length dependence of E loss (almond-shaped medium => $v₂$ ~ 0.05--0.10)

Conclusions

- The (open) heavy-flavour era of heavy-ion physics has begun
- Observation of heavy-quarks quenching at RHIC !
- The LHC will be a `hard probes and heavy quarks machine' and quenching studies will play a central role

Promising observables at the LHC:

- R_{AA} of D and B mesons
- ◆ Heavy-to-light ratios as probes of E loss...
	- $m = 1$... color-charge dependence $(R_{D/h})$
	- L … parton-mass dependence $(R_{B/h})$
- Azimuthal anysotropy for D, B mesons (or their decay leptons): interplay between reco/flow and E loss
- + study of dense QCD initial-state phenomenology

EXTRA SLIDES

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B.Muller's QM Theory Summary

Emergence of away-side jet will make determination of q^ easier.

Medium Expansion

• The density of scattering centers is time-dependent:

$$
\hat{q}(\tau) = \hat{q}_0 \times \left(\frac{\tau_0}{\tau}\right)^{\alpha}
$$

◆ Dynamical scaling law: *same* spectrum obtained *for equivalent static trasport coefficient*

$$
\overline{\hat{q}} = \frac{2}{L^2} \int_{\tau_0}^{\tau_0+L} d\tau (\tau - \tau_0) \hat{q}(\tau)
$$

Calculations for a static medium apply to expanding systems

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 α = 1.5, 1.0, 0.5, 0

Application: Parton Quenching Model

- QW + Glauber-based medium geometry and density profile + PYTHIA for parton generation and fragmentation
- The procedure in short:
	- 1) generate parton (q or g) with PYTHIA (or back-to-back pair)
	- 2) calculate its L and average $\;\hat{q}\;$ along the path $[\hat{q}(\vec{s}) \propto T_A T_B(\vec{s})]$
	- 3) use quenching weights to get energy loss
	- 4) quench parton and then hadronize it (independent fragm.)

Calculating Parton Energy Loss

$$
\langle \Delta E \rangle \approx \int_0^{\omega_c} \mathrm{d}\omega \, \omega \frac{\mathrm{d}I}{\mathrm{d}\omega} \propto \alpha_s \, C_R \, \omega_c \propto \alpha_s \, C_R \, \hat{q} \, L^2
$$

 $\propto \hat{q} \propto \frac{$ gluons volume-density and interaction cross section

Probe the medium

Finite parton energy (qualitatively) $\langle E \rangle \approx \int_{0}^{E} d\omega \, \omega \frac{dI}{dt} \propto \alpha_{s} C_{R} \sqrt{E \omega_{c}} \propto \alpha_{s} C_{R} \sqrt{E} \sqrt{\hat{q}} L$ $\langle \Delta E \rangle \approx \int_0^E {\rm d}\omega \, \omega \frac{{\rm d} I}{{\rm d}\omega} \! \propto \alpha_s \, C_{_R} \, \sqrt{E \omega_c} \, \propto \alpha_s \, C_{_R} \, \sqrt{E} \, \sqrt{\hat q}$ $\hat{q} \rightarrow \hat{q}^{1/2}$: smaller sensitivity to density $\triangleright L^2 \rightarrow L$ \cdot If E < ω_c (e.g. small $p_{\rm T}$ parton with large L): ¾ **dependence on parton energy**

 ΔE \rangle \propto

Model vs RHIC data Centrality dependence of RAA

Centrality evolution according to Glauber-model collision geometry $\hat{q}(\vec{s};b) = k \times T_A T_B(\vec{s};b)$

Dainese, Loizides, Paic, EPJC 38 (2005) 461.

Model vs RHIC data Disappearence of the away-side jet

STAR Coll., PRL 90 (2003) 082302. STAR Coll., nucl-ex/0501016.

Extrapolation in c.m.s. energy Intermediate RHIC energy √*s = 62 GeV*

Extrapolation in √*s*: assuming $\;\hat{q}$ ∝ N_{gluons}/volume ∝ (√*s*)^{0.6} (EKRT saturation model)

EKRT: Eskola, Kajantie, Ruuskanen, Tuominen, NPB 570 (2000) 379. PHENIX Coll., JPG 31 (2005) S473.

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Model prediction for LHC

◆ Extrapolation to LHC according to saturation model gives: $\hat{q}_\textrm{5.5 TeV} \leq 7 \times \hat{q}_\textrm{200 GeV} \approx 100 \, \textrm{GeV}^2/\textrm{fm}$

Why R_{AA} is flat ↔ Surface effect

- Long path lengths exploited only by high energy partons \triangleright effectively, $\langle \Delta E \rangle \propto E^{\alpha}$, with $0.5 < \alpha < 1$
	- \triangleright R_{AA} doesn't increase at high p_T
	- $▶$ Exercise: $L \equiv \langle L \rangle \Rightarrow R_{AA}$ increases with p_{T}

Open points (1): limited sensitivity of RAA

- Strong suppression requires very large density
- Surface emission scenario
- $R^{\,}_{\rm AA}$ determined by geometry rather than by density itself

Eskola, Honkanen, Salgado, Wiedemann, NPA 747 (2005) 511.

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?

Open points (2): the opacity problem

- **◆ Can we really** *probe* the medium?
- Need to relate extracted $\;\hat{q}\;$ to an energy density $\bm{\mathcal{E}}$
- QCD estimate for ideal QGP: $\,\hat{q}(\tau) \approx 2 \!\times\! \bm{\mathcal{E}}^{3/4}(\tau) \qquad \hbox{(Baier)}$
- A recent analysis* of RHIC data, similar to that presented, extracts energy density ×5 larger than that estimated from produced transverse energy dE_T/dy (Bjorken estimate)
- **Opacity problem: the interaction of the hard parton with the medium is much stronger than expected**

Baier, NPA 715 (2003) 209.

* Eskola, Honkanen, Salgado, Wiedemann, NPA 747 (2005) 511.

Charm Energy Loss at RHIC

- Compute $R_{\sf AA}$ for c, D, e (from D)
- Baseline: *PYTHIA*, with *EKS98 shadowing*, tuned to match p_T -shape of D cross section measured in d-Au by STAR
- c-quark *E* loss as for light-flv hadrons, but using (m_c/E_c) -dependent QW
- $\widetilde{q}\;$ extracted from π^0,h^\pm $R_{\sf AA}$ (central) ˆ
- **Thermalize charms that lose all energy d/V/d** m_T **∝** m_T **exp(-** m_T **/7), т = 300 MeV**
- Use a range in $|q|$ to visualize limited sensitivity of R_{AA} to $\stackrel{_}{q}$ itself ˆˆ

Armesto, Dainese, Salgado, Wiedemann, PRD 71 (2005) 054027. EKS: Eskola, Kolhinen, Salgado, EPJC 9 (1999) 61.

The heavy-to-light ratio: $\chi_{h}(p_{T}) \equiv \frac{R_{AA}^{D}(p_{T})}{R_{A}^{b}}$ $D/h(P_T) \equiv \frac{P_{AA} \cdot P}{R_{AA}^h(p)}$ $R_{D/h}(p_T) \equiv \frac{R_{AA}^D(p)}{R_{D/h}}$

What enters the game:

- (c) quark vs gluon (Casimir factor) $\rightarrow R_{D/h} > 1$
- 2) harder charm $\rho_{\sf T}$ distribution: $\qquad \rightarrow \;\; R_{\sf D/h} \;$ up
- 3) harder charm fragm.

 $R_{D/h}$ down

 (p_T)

h

D

AA

T

T

