

# Standard model Physics

The double-simplex

QuickTime™ and a  
Animation decompressor  
are needed to see this picture.

[Quigg,04]

at the LHC  
Fabio Maltoni

## Motivations

SM Physics we know but we need to know better

SM Physics we know is there but not seen yet

SM Physics we do not know if it is there or not!

Outlook

## Examples

QCD,  
 $\sigma(W,Z), m_W, m_t$

Single-top, Rare dec  
Vector Boson scatter

Higgs

# The SM Lagrangian

$$\mathcal{L}_{eff} = -\frac{1}{4}G^{A\mu\nu}G_{\mu\nu}^A - \frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{4}W^{A\mu\nu}W_{\mu\nu}^A$$

Gaug

$$+i\bar{Q}_L^i \not{D} Q_L^i + i\bar{u}_R^i \not{D} u_R^i + i\bar{d}_R^i \not{D} d_R^i + i\bar{L}_L^i \not{D} L_L^i + i\bar{e}_R^i \not{D} e_R^i$$

Matt

$$+M_W^2 W^{+\mu} W_{\mu}^{-} + \frac{1}{2} \frac{M_W^2}{\cos^2 \theta_W} Z^{\mu} Z_{\mu}$$

Mass

$$-m_u^{ij} \bar{u}_L^i u_R^j$$

\*Non linearly realized. Unitary gauge.

\*\* Except for neutrino's masses and dark matter

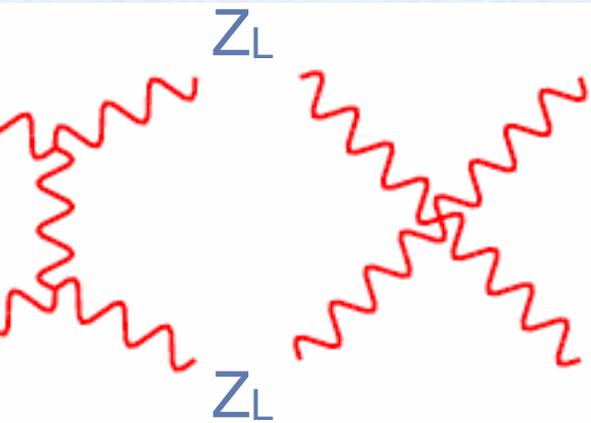
Perfectly defined gauge-invariant\* Lagrangian!

It describes all\*\* we know!

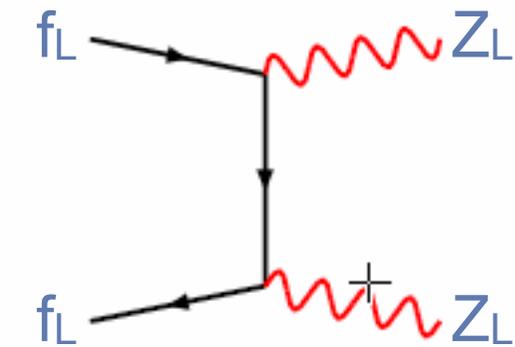
# Effective Lagrangian

[Chanowitz, Gallard.1985]

[Appelquist, Chanowitz, 1985]



$$a_0 \sim \frac{s}{v^2}$$



$$a_0 \sim \frac{\sqrt{sm_f}}{v^2}$$

Inelastic tree-level amplitude for longitudinal W and Z fermions violate unitarity at scale:

$$\Lambda_{EWSB} = \sqrt{8\pi}v$$

Our effective description contains information on when is going to fail.

Only case we know of where unknown physics has to appear below 1 TeV.

# Chiral Lagrangian

$$\mathcal{L}_M = -\frac{1}{4} G^{A\mu\nu} G_{\mu\nu}^A - \frac{1}{4} W^{A\mu\nu} W_{\mu\nu}^A - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} \quad \text{Gaug}$$

SU(3)                      SU(2)<sub>L</sub>                      U(1)<sub>Y</sub>

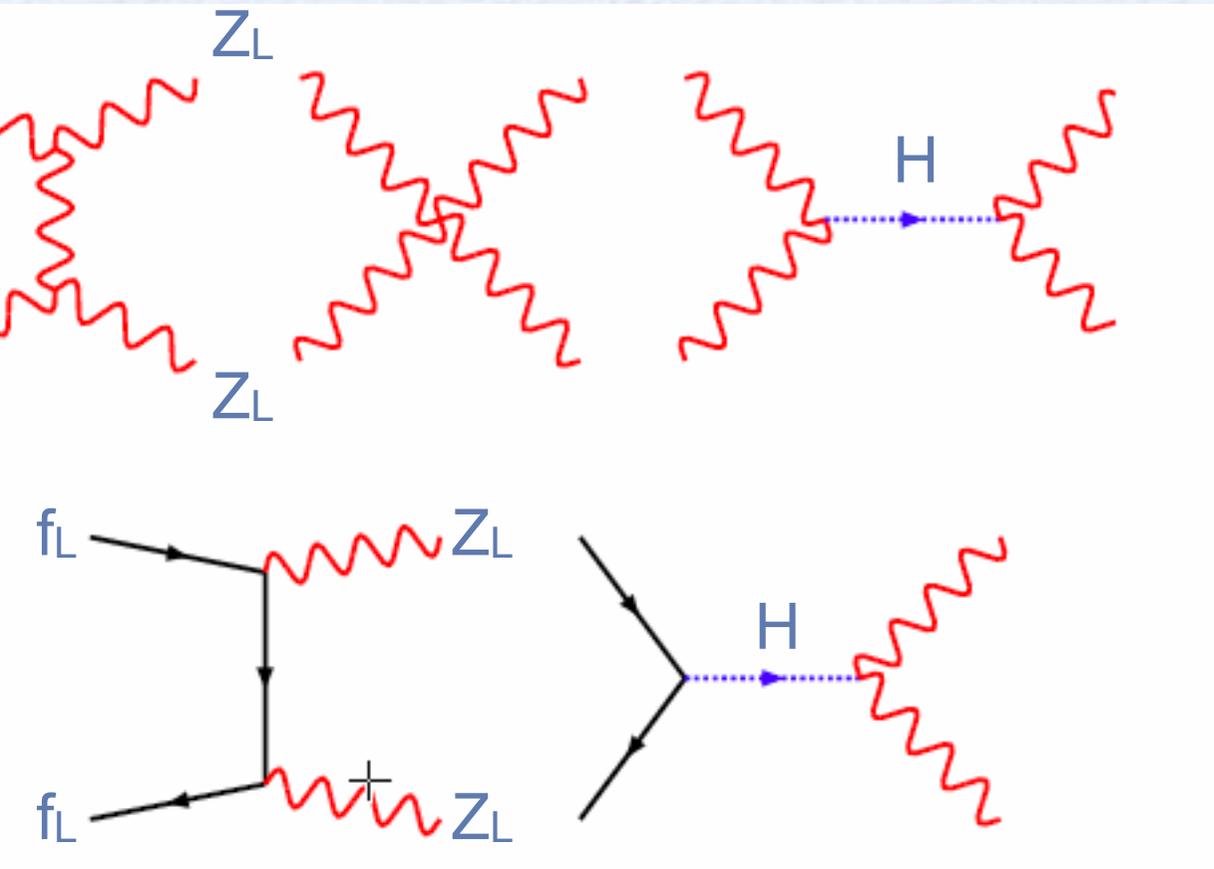
$$+ i\bar{Q}_L^i \not{D} Q_L^i + i\bar{u}_R^i \not{D} u_R^i + i\bar{d}_R^i \not{D} d_R^i + i\bar{L}_L^i \not{D} L_L^i + i\bar{e}_R^i \not{D} e_R^i \quad \text{Matt}$$

$$- \Gamma_u^{ij} \bar{Q}_L^i \epsilon \phi^* u_R^j - \Gamma_d^{ij} \bar{Q}_L^i \phi d_R^j - \Gamma_e^{ij} \bar{L}_L^i \phi e_R^j + h.c. \quad \text{Yuka}$$

$$+ (D^\mu \phi)^\dagger (D^\mu \phi) + \mu^2 \phi^\dagger \phi - \lambda (\phi^\dagger \phi)^2 \quad \text{Higg}$$

Spontaneous symmetry breaking  $\Rightarrow \phi = \frac{H + i\eta}{\sqrt{2}}$

[Chanowitz, Gallard.1985]  
 [Appelquist, Chanowitz, 19



$$a_0 \sim \frac{s}{v^2} - \frac{s}{v^2} \sim \frac{m_H^2}{v^2}$$

$$a_0 \sim \frac{\sqrt{sm_f}}{v^2} - \frac{\sqrt{sm_f}}{v^2} \sim$$

SM is a linear gauge theory valid up to arbitrary scales,  
 with  $m_H < 1$  TeV.

$$\begin{aligned}
\mathcal{L}_{SM} = & -\frac{1}{4}G^{A\mu\nu}G_{\mu\nu}^A - \frac{1}{4}W^{A\mu\nu}W_{\mu\nu}^A - \frac{1}{4}B^{\mu\nu}B_{\mu\nu} && \text{Gauge} \\
& +i\bar{Q}_L^i \not{D} Q_L^i + i\bar{u}_R^i \not{D} u_R^i + i\bar{d}_R^i \not{D} d_R^i + i\bar{L}_L^i \not{D} L_L^i + i\bar{e}_R^i \not{D} e_R^i && \text{Matter} \\
& -\Gamma_u^{ij} \bar{Q}_L^i \epsilon \phi^* u_R^j - \Gamma_d^{ij} \bar{Q}_L^i \phi d_R^j - \Gamma_e^{ij} \bar{L}_L^i \phi e_R^j + h.c. && \text{Yukawa} \\
& +(D^\mu \phi)^\dagger (D^\mu \phi) + \mu^2 \phi^\dagger \phi - \lambda(\phi^\dagger \phi)^2 && \text{Higgs}
\end{aligned}$$

Perform precision measurements to establish **direct** evidence for Higgs and BSM physics.

Use the heaviest particles (W,Z, top) as **direct** probes of the EWSB sector and BSM physics.

# Motivations

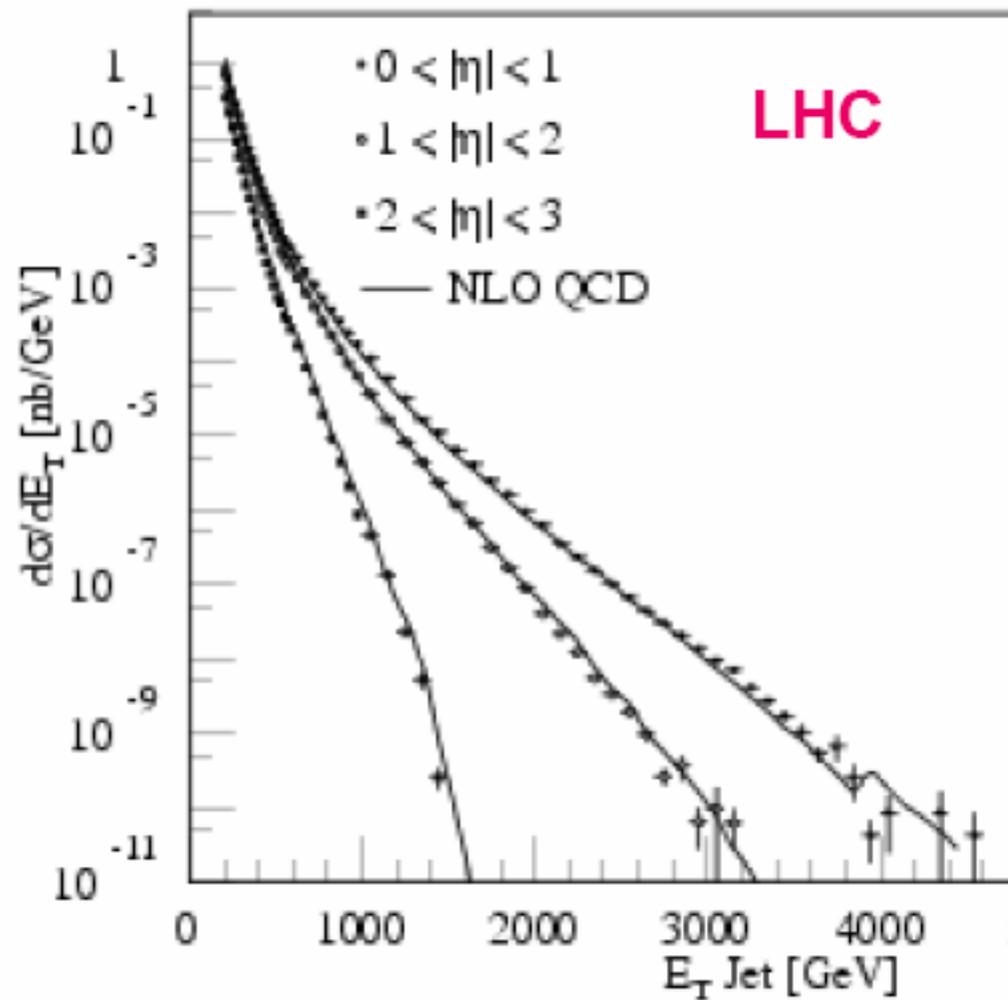
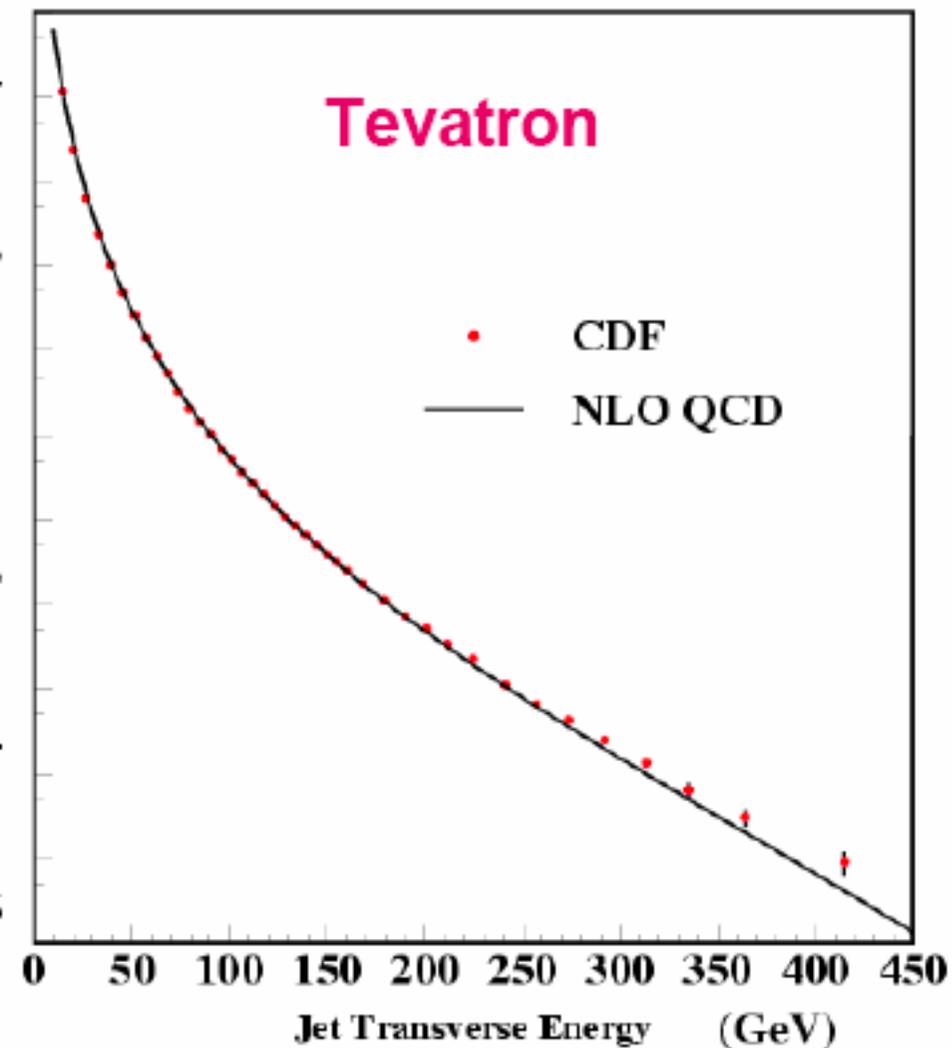
SM Physics we know but we need to know better

QCD,  
 $\sigma(W,Z), m_W, m_t$

SM Physics we know is there but not seen yet

SM Physics we do not know if it is there or not!

Outlook



ressive agreement over 9 orders  
 magnitude! At high  $E_T$  statistically  
 ted. Theoretical uncertainties

Enormous rates ( $10^3$  events/s  
 $E_T > 100$  GeV).  
 How to calibrate jet energy scale?

increased by order of magnitudes wrt  
on.

s with vectors bosons, tops and heavy and  
s with rates >1 Hz.

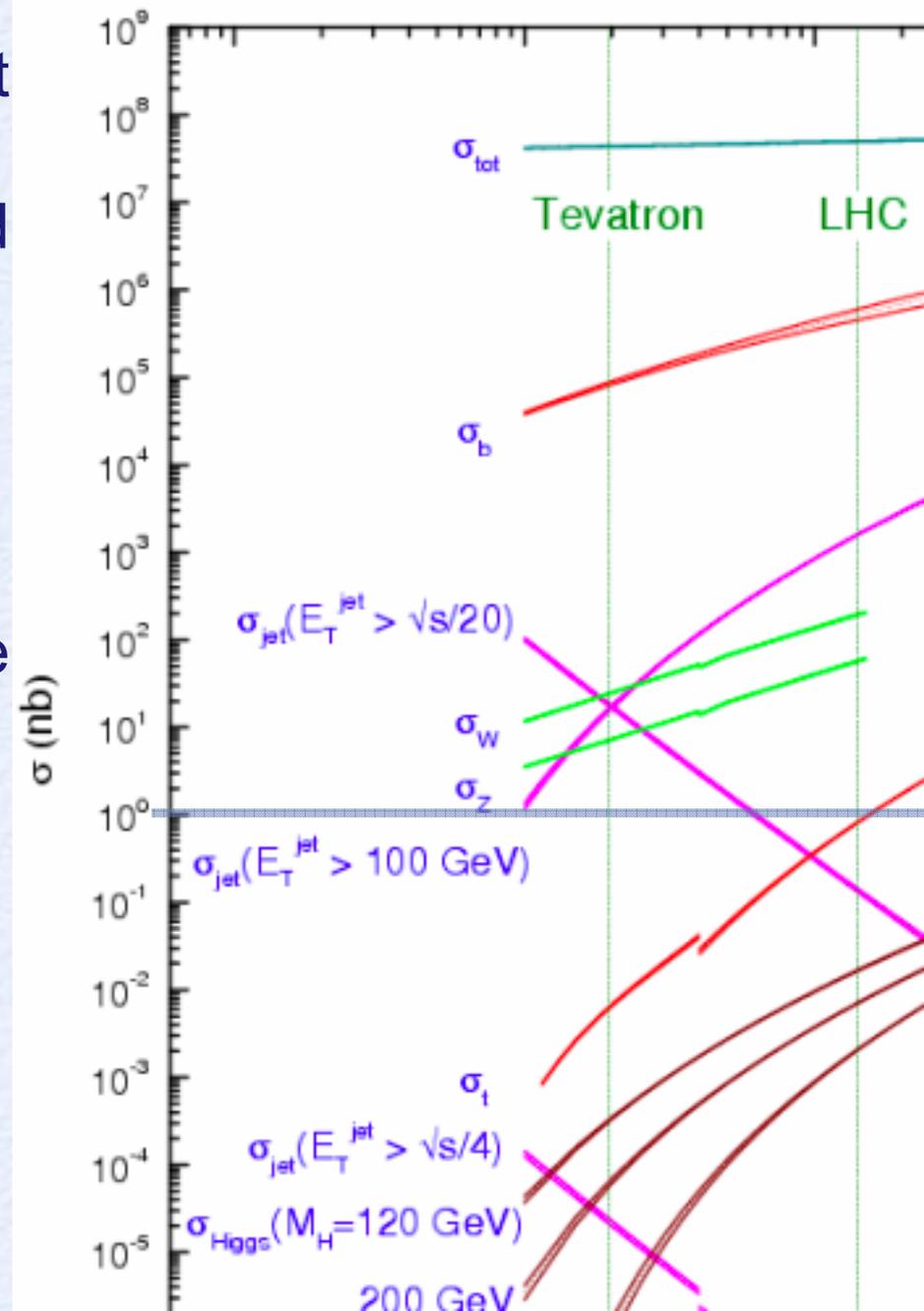
physics down order of magnitudes.  
to understand QCD backgrounds well!

factorization theorem for short-distance  
e processes:

$$= \sum_{a,b} \int_0^1 dx_1 dx_2 f_a(x_1, \mu_F^2) f_b(x_2, \mu_F^2) \\ \times \hat{\sigma}_{ab \rightarrow X}(x_1, x_2, \alpha_S(\mu_R^2), \frac{Q^2}{\mu_F^2}, \frac{Q^2}{\mu_R^2})$$

ingredients necessary:  
on Distribution functions (from exp).

proton - (anti)proton cross section



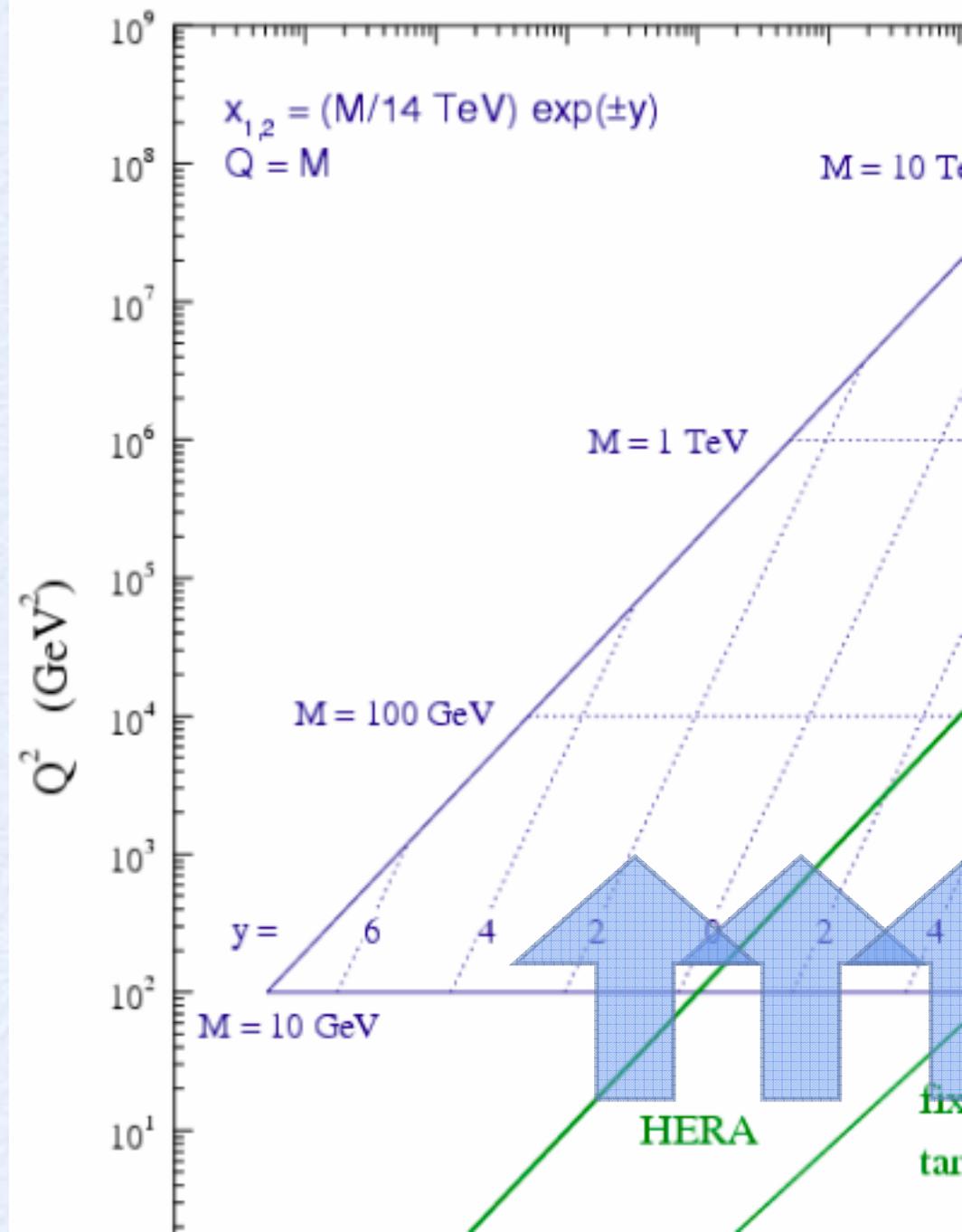
measured at HERA and fixed-target experiments. x dependence from data. Q dependence from DGLAP evolution.

Recently:

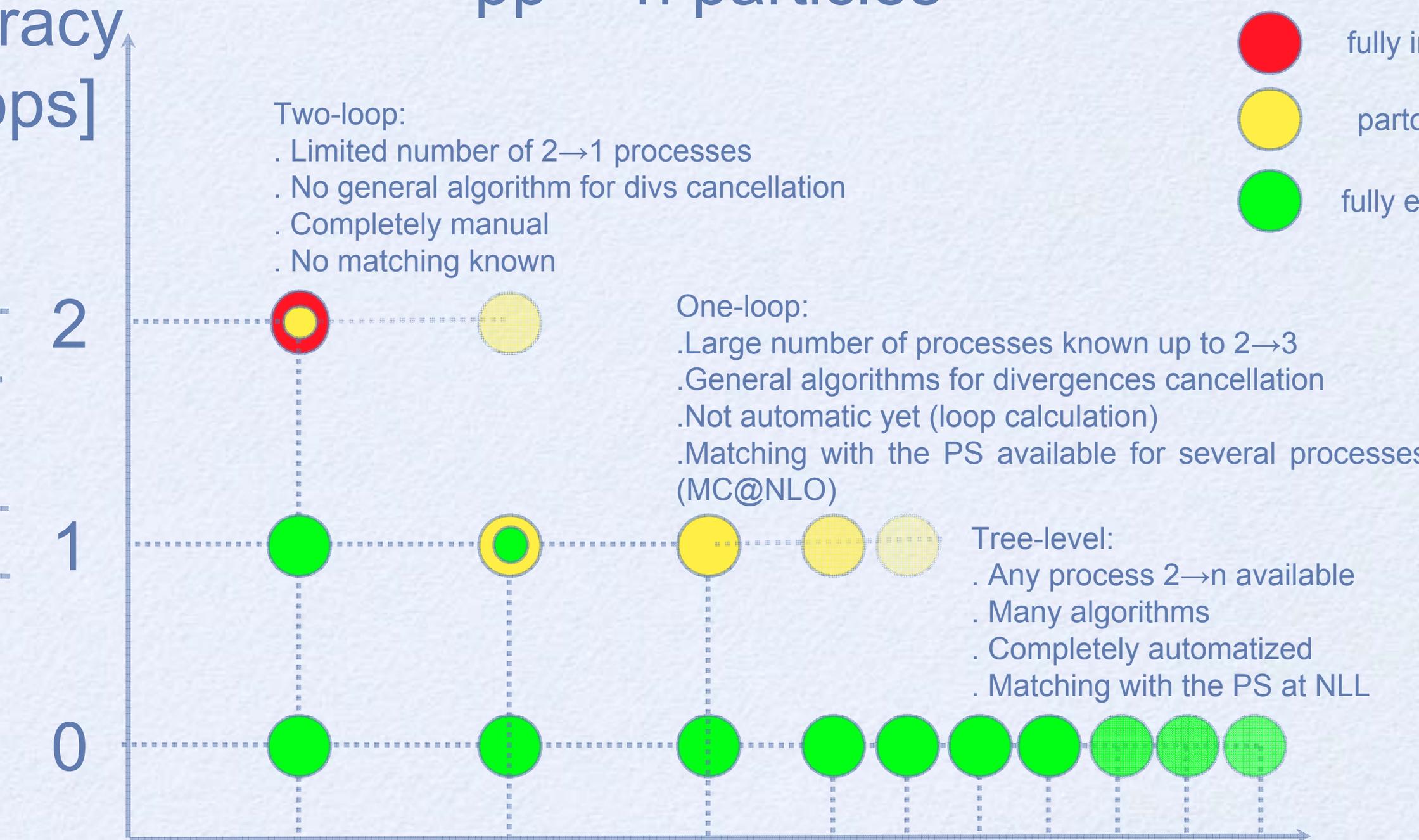
calculation of the 3-loop splitting kernels (hardest calculation in QCD") [Carter, Vermaseren, Vogt. 2004]

together with short distance NNLO calculation (first sets of NNLO PDF sets. [MRST and Alekhin,

with errors: Various "traditional PDF sets", [CTEQ and MRST, 2003]. Also new approaches, the functional space [Giele, Keller, ...]



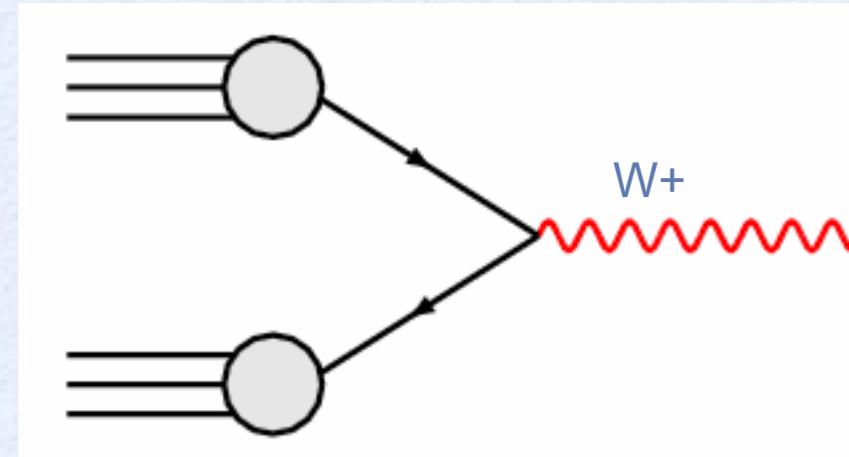
# pp → n particles



# W PRODUCTION

$$\sigma^h(W) = \sum_{ab} \mathcal{P}_{ab} \otimes \hat{\sigma}_{ab}(W)$$

$$\sigma^{exp}(W) = \frac{1}{BR(W \rightarrow \ell\nu)} \frac{1}{\int \mathcal{L} dt} \frac{N^{sig}}{A_W}$$



Theory Status:

measuring  $\sigma(W)$ , one needs to estimate the acceptance  $A_W$  theory and the luminosity from an independent source. A exclusive description of the final state is needed.

if  $\sigma(W)$  is accurate enough, one can use  $\sigma(W)$  to:

precisely measure  $\Gamma_W$

$R = \sigma(W) BR(W \rightarrow \ell\nu) / \sigma(Z) BR(Z \rightarrow \ell\ell)$

obtain direct information on the PDF

measure the collider luminosity

obtain parton-parton luminosity (=luminosity+PDF)

Best QCD predictions at present:

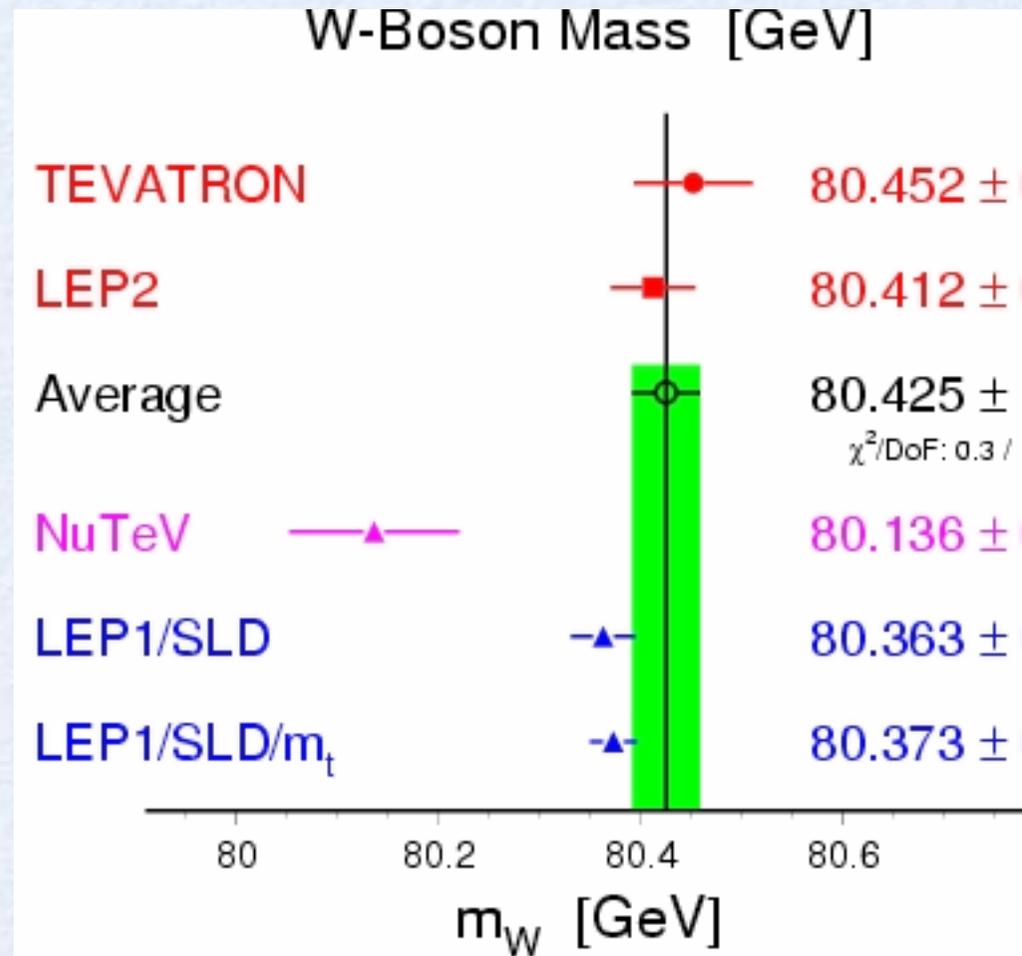
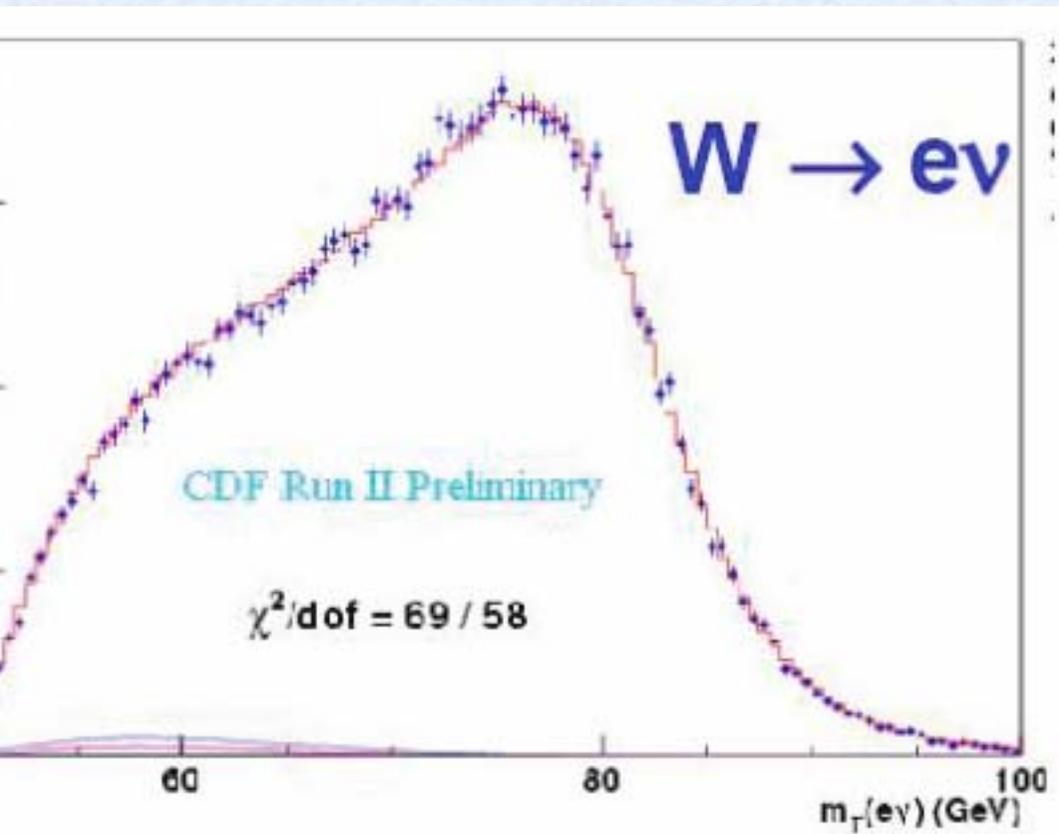
- > Fully exclusive (PS interfaced) at NLO+NLL [Frixione, Webber, 2002]
- > Partly exclusive prediction at NLO [Anastasiou, Dixon, Melnikov, 2003]
- > Resummed pt distribution at NLO [Balazs, Qiu, Yuan, 1995]
- > 1-loop EW corrections [Baur, Wackerroth, 2000]

Spin correlations are important!

[Frixione, Mangano, 2004]

Exclusive NNLO calculation WITH

# W mass



Run II expectation:

improve on LEP2 result:  $\delta m_W = 40$  MeV  
 $2\text{fb}^{-1}$  per lepton channel per  
 experiment.

HL-LHC expectation:

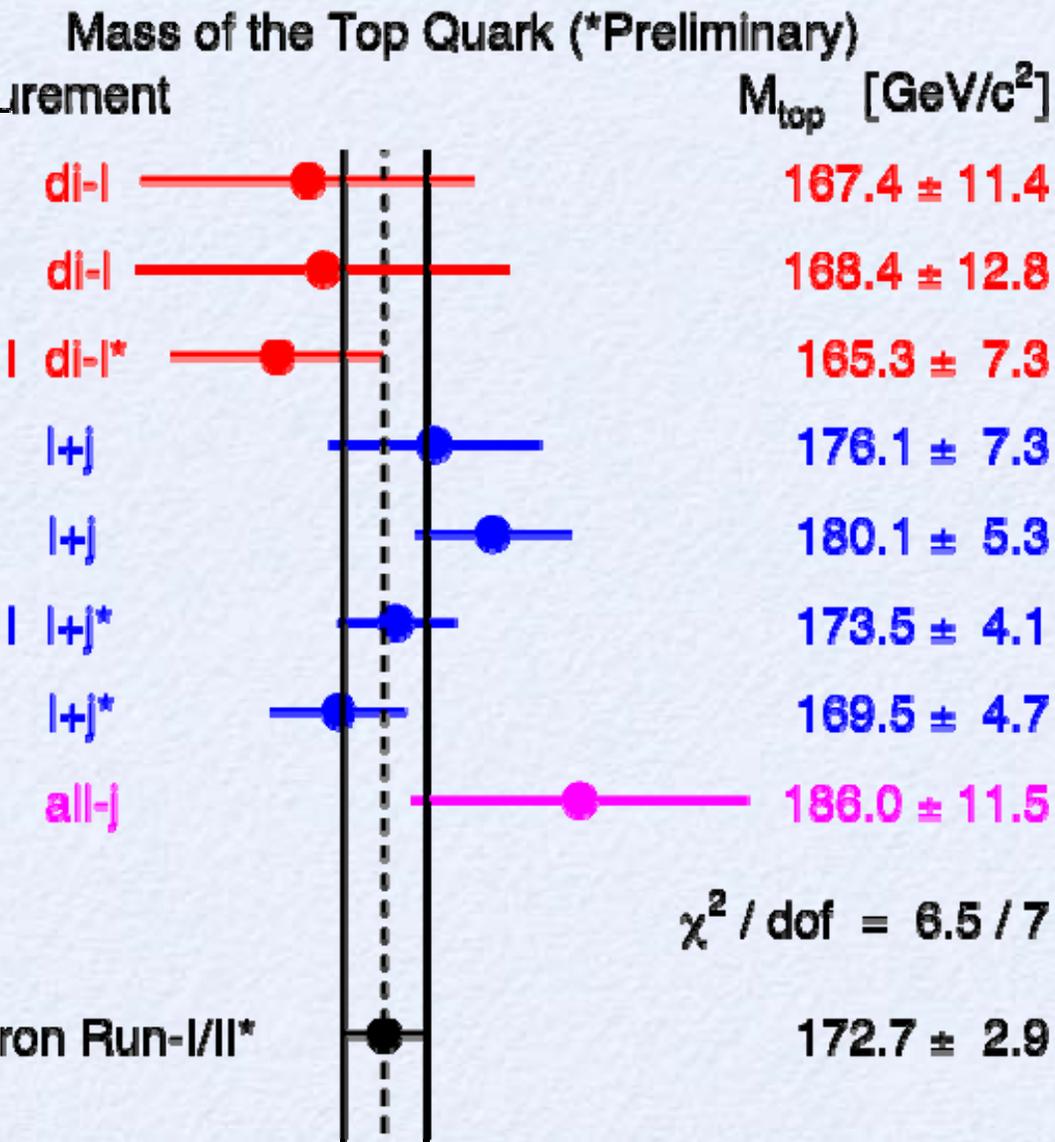
$\delta m_W = 15$  MeV from transverse mass

Need:

$$\delta m_W \sim 7 \times 10^{-3} \delta m_t$$

In the SM, it is the **only** quark with a “natural mass”:

$$m_{\text{top}} = y_t v / \sqrt{2} \approx 174 \text{ GeV} \Rightarrow y_t \approx 1$$



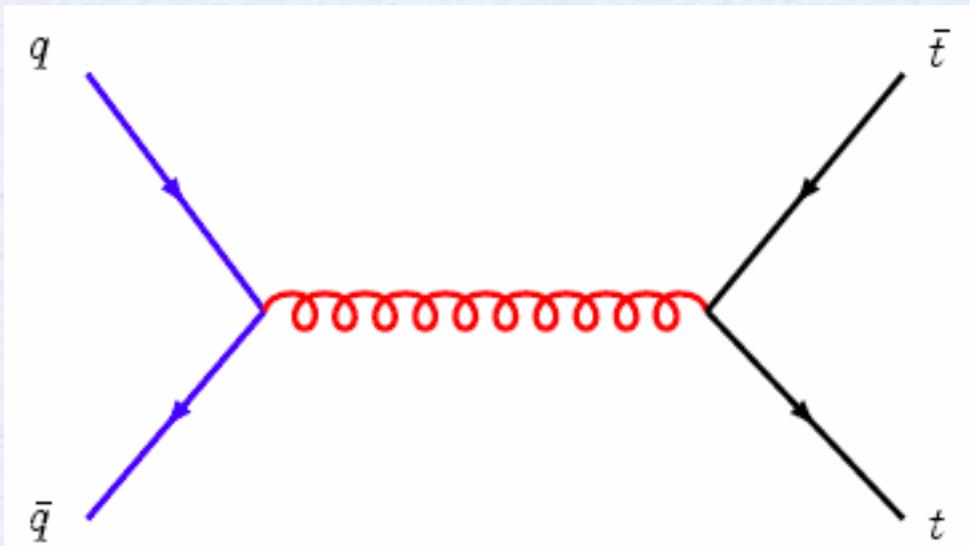
Many methods. Two important ones:

I. **Template mass**: global kinematic fit based on likelihood of matching the masses of the different objects in the chain.

II. **Matrix element** (also called DLM): “maximal” use of the information contained in the ME (both signal and background) to build a likelihood of an event has a top mass.

Both methods suffer from jet ES and ISR/FSR uncertainties.

## Tevatron



85% of the total cross section

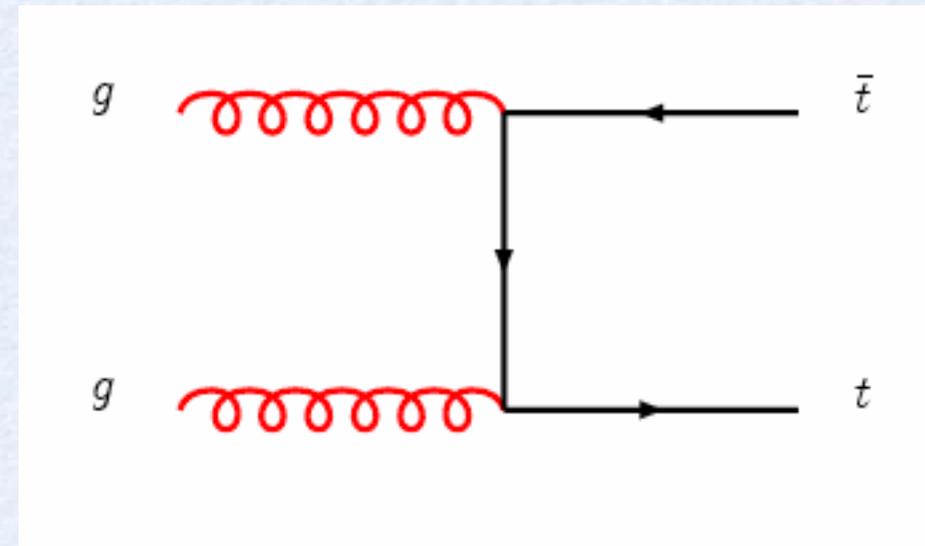
10 tt pairs per day

90% of the time there is extra radiation  
that  $p_t(t\bar{t}) > 15$  GeV.

tt are produced closed to threshold, in a  
 $^1[8]$  state. Same spin directions. 100%  
correlated in the off-diagonal basis.

Sorry because of the backgrounds:

## LHC



90% of the total cross section

1 tt pair per second

Almost 70% of the time there is  
radiation so that  $p_t(t\bar{t}) > 30$  GeV.

tt can be easily produced away  
threshold. On threshold they are  
in a  $^3[0]$  state, with opposite spin directions.  
100% correlation.

## Motivations

SM Physics we know but we need to know better

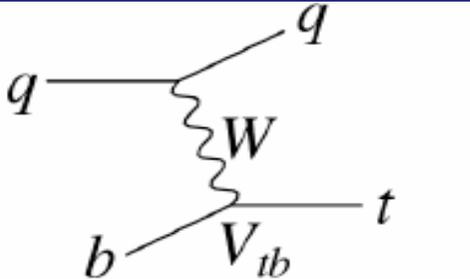
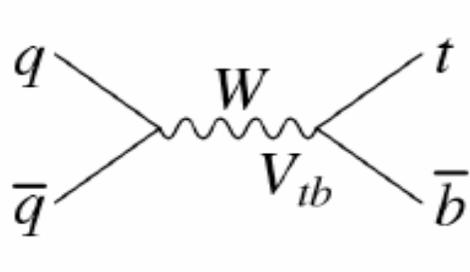
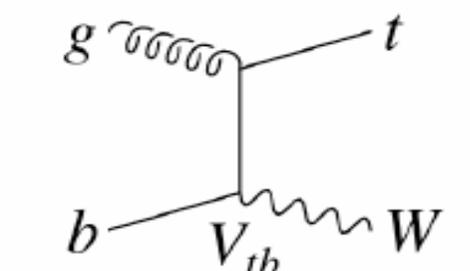
SM Physics we know is there but not seen yet

Single-top

SM Physics we do not know if it is there or not!

Outlook

# Single top

Process	Diagram	Accuracy	CTEQ6M, $m_t=178$ GeV, $\theta_{\text{ren}}=0$ $\sigma$ (pb)	
			TeV II	LHC
s-channel		<p><b>NLO</b></p> <p>[Stelzer, Sullivan, Willenbrock. 1997]</p>	1.85	239
t-channel		<p><b>(N)NLO</b></p> <p>[Smith, Willenbrock. 1996 Chetyrkin, Steinhauser. 2001]</p>	0.82	9.8
tW		<p><b>NLO</b></p> <p>[Campbell, Tramontano. '05]</p>	0.129	64

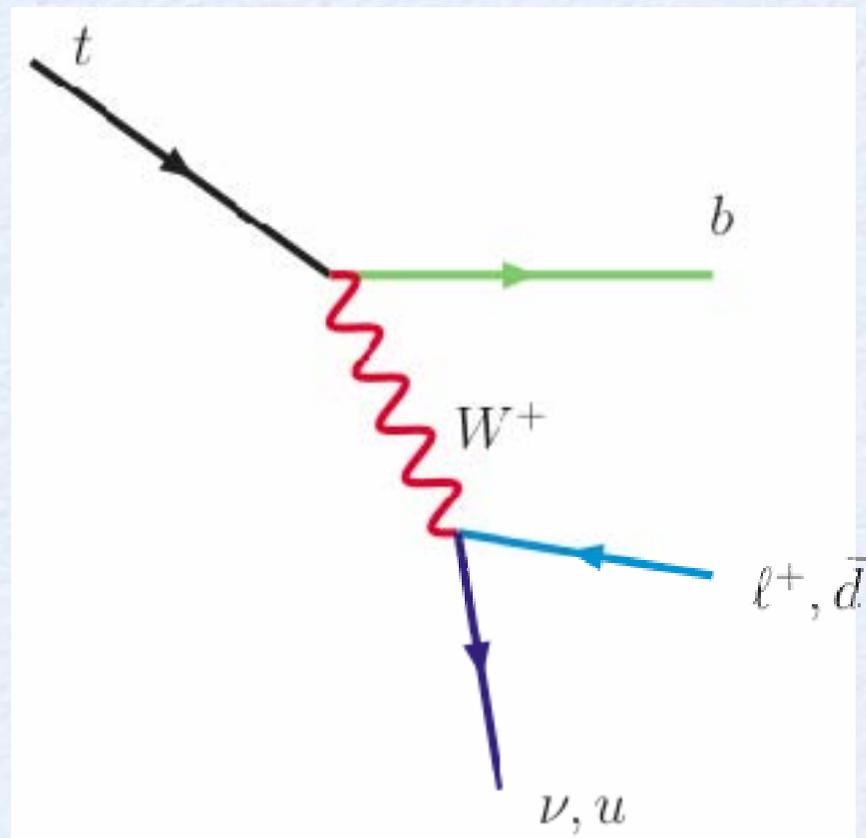
signals available in MCFM (Campbell, Ellis) and t- and s-channel soon in MC@NL (Alfonso, Webber). Most of the backgrounds are also known at NLO. However, analysis

# Top decay channels

Top can decay into a real  $W \Rightarrow$

$\Gamma \approx GF m_t^3 |V_{tb}|^2 \gg \Lambda_{\text{QCD}} \Rightarrow$

Very short life. Top is the only quark that does not feel non perturbative QCD effects! No top-hadrons, no top-spectroscopy but a "clean" quark.



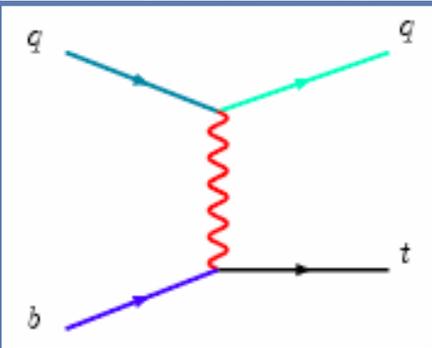
In an experiment one is sensitive not to the total width but to the branching ratio:

$$R = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wq)} = \frac{|V_{tb}|^2}{|V_{td}|^2 + |V_{ts}|^2 + \dots}$$

CDF has performed such a measurement

$R=0.94$  does only tell us that  $V_{tb} \gg V_{td}$

# t channel



## SM info

highest rate, dominant at the LHC, where 62% top, 38% anti-top.

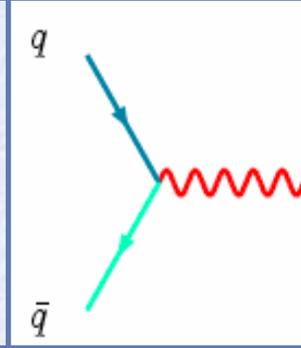
$\sigma \propto |V_{tb}|^2$ .  
 Hard jet in final state, top central, sometimes one extra forward bottom. FB asymmetric at the Tevatron. Main background s+jet (and tt at the LHC).

Top is polarized along spectator jet (most of times) in the 2→2 configuration.

## BSM window

Sensitive to new production modes,

# s channel



## SM info

Smallest at the LHC, where 63% top, 37% anti-top.

$\sigma \propto |V_{tb}|^2$ . Very well known. DY might be used for normalization.

Central high-pt b-jet. Main backgrounds s+jet and W+Q's+jets.

Top is polarized along beam axis at Tevatron.

## BSM window

Sensitive to vector (extra Z) and scalar (toponium) resonances.

## Motivations

SM Physics we know but we need to know better

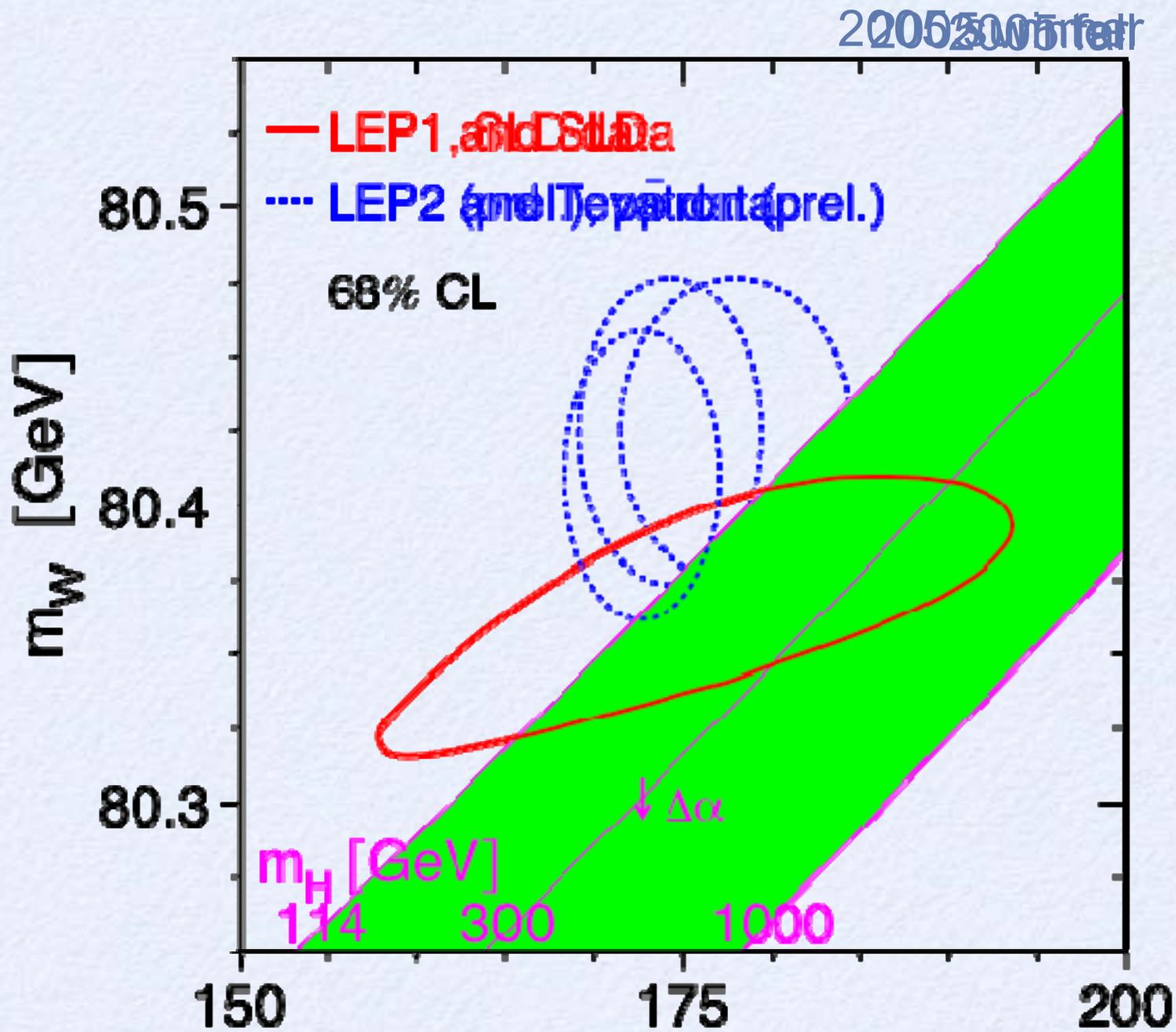
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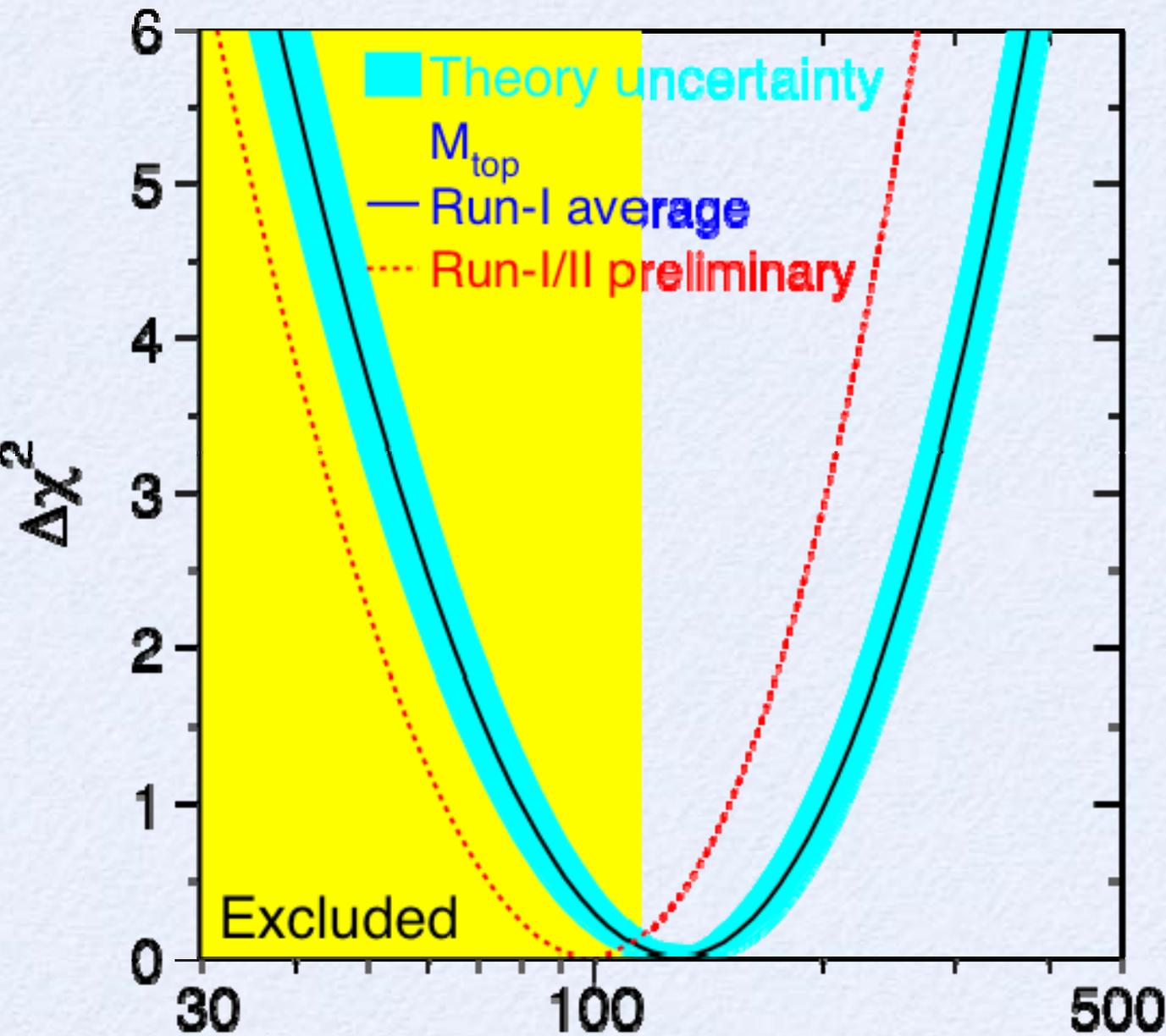
Higgs

## Outlook

# $m_{top} \propto m_W \propto m_H$



# $m_{top}$ & $m_{H^0}$ / $m_{A^0}$

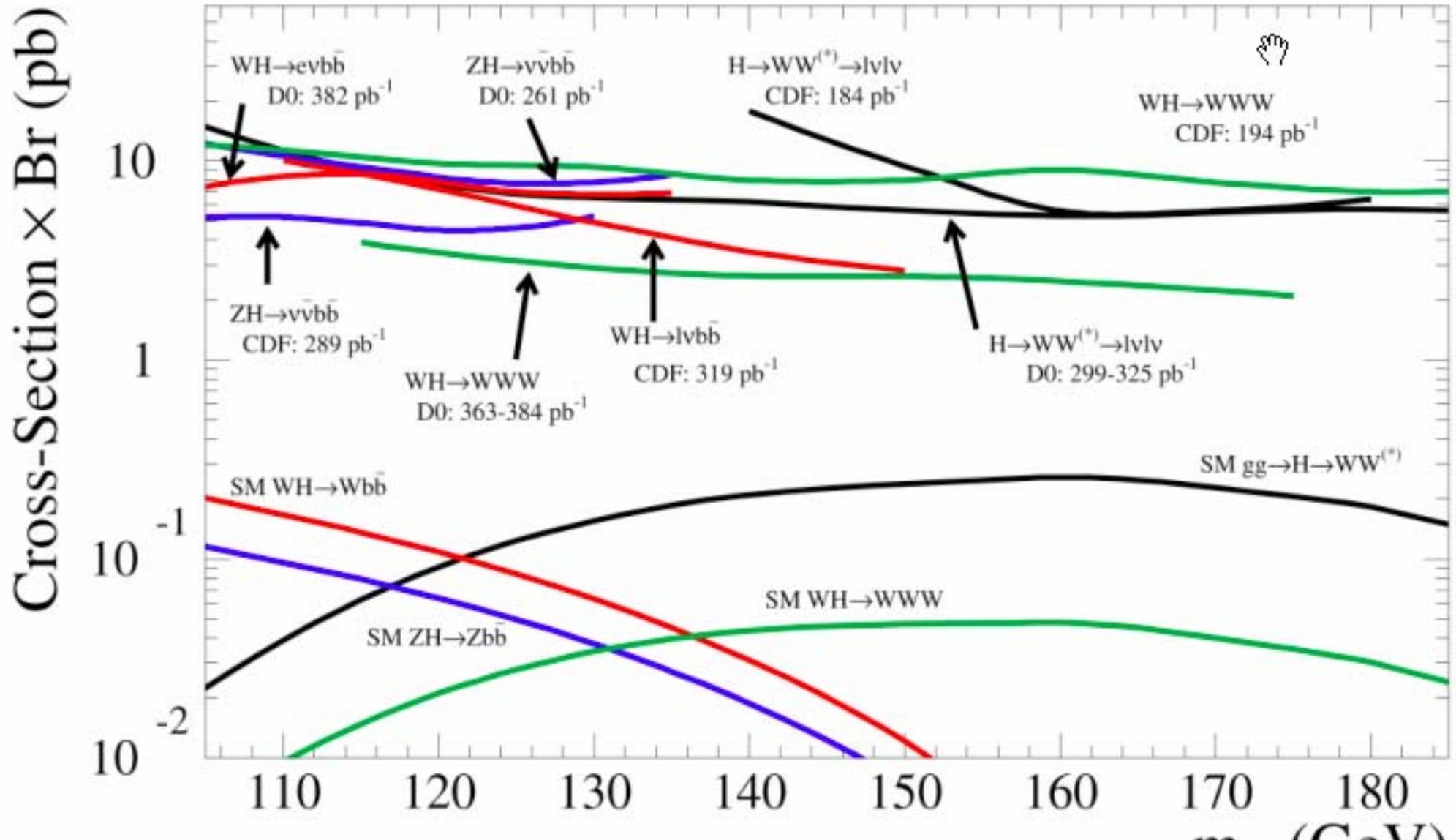


$m_H > 114$  GeV from L

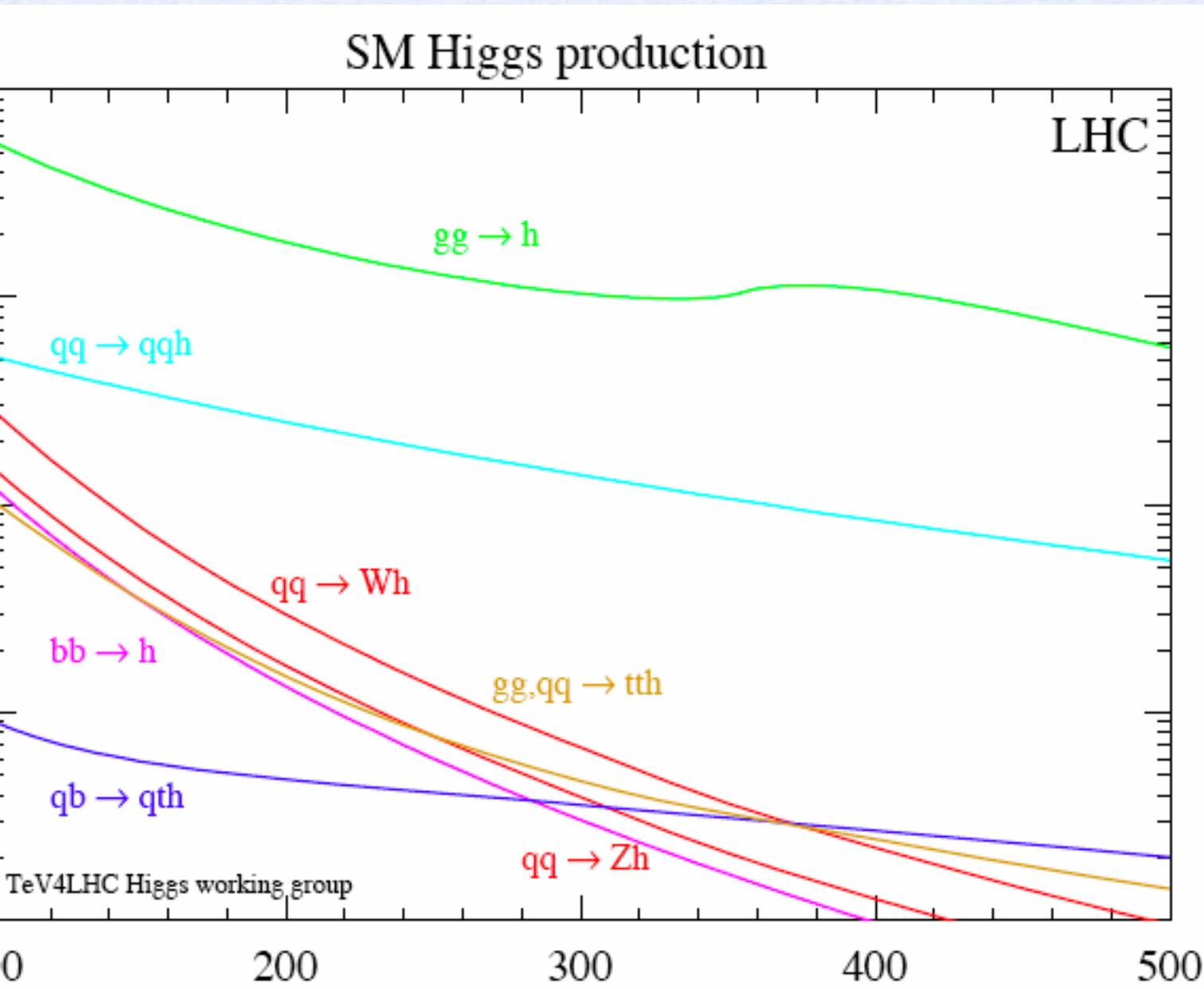
$m_H < 208$  GeV @ 95%

# STAND

## Tevatron Run II Preliminary



# LHC

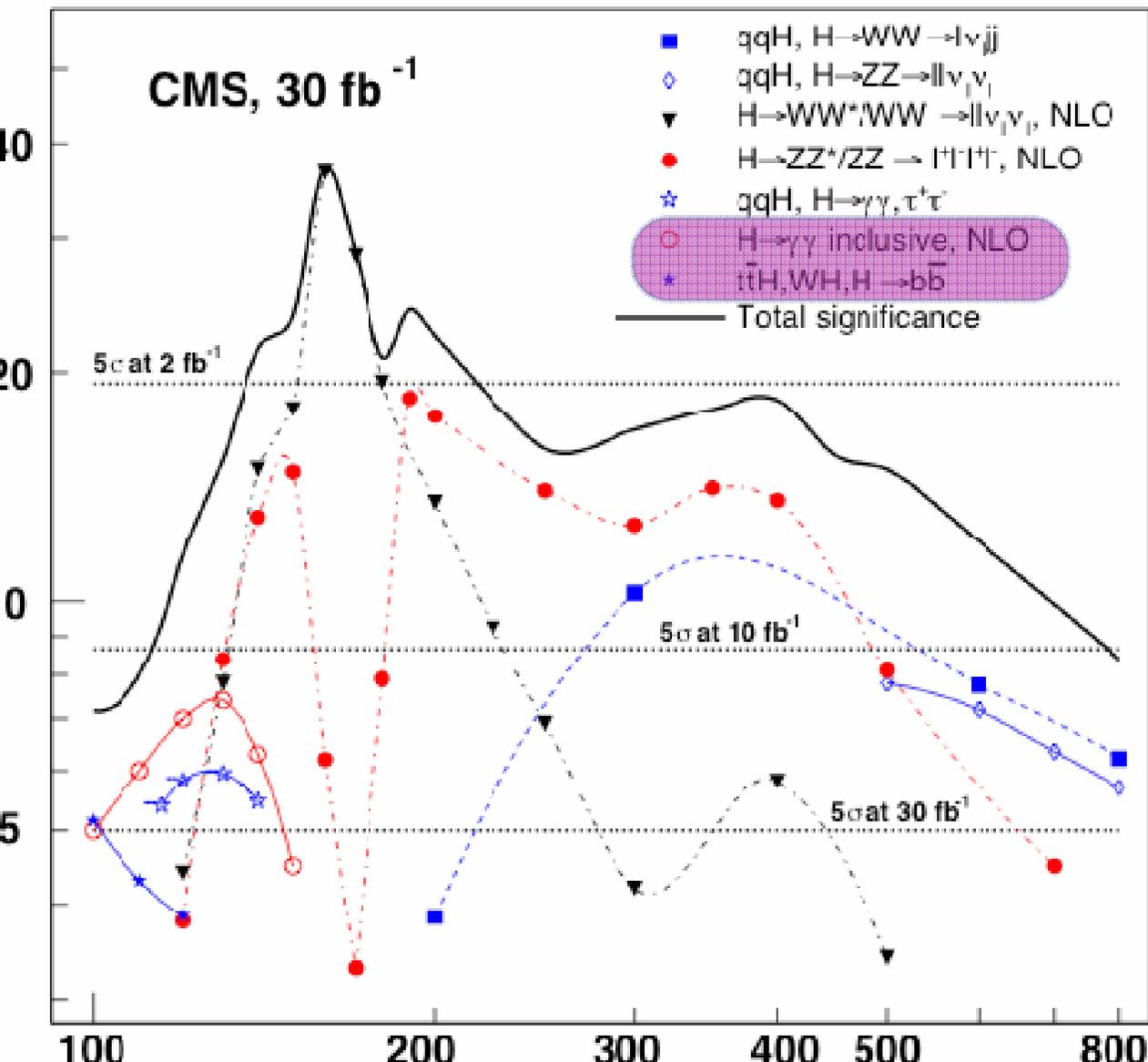


All relevant signals known  
at NLO in QCD.

Various channels/decay  
are complementary:  
Inclusive/associated pro  
with rare/common decay m

Very different strategies and  
systematics.

# LHC



All relevant signals known  
at NLO in QCD.

Various channels/decay  
are complementary:  
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Very different strategies an  
systematics.

Two Examples for light Hig

1.  $gg \rightarrow H \rightarrow \gamma\gamma$
2.  $gg \rightarrow ttH \rightarrow ttbb$

CMS, 100 fb<sup>-1</sup>

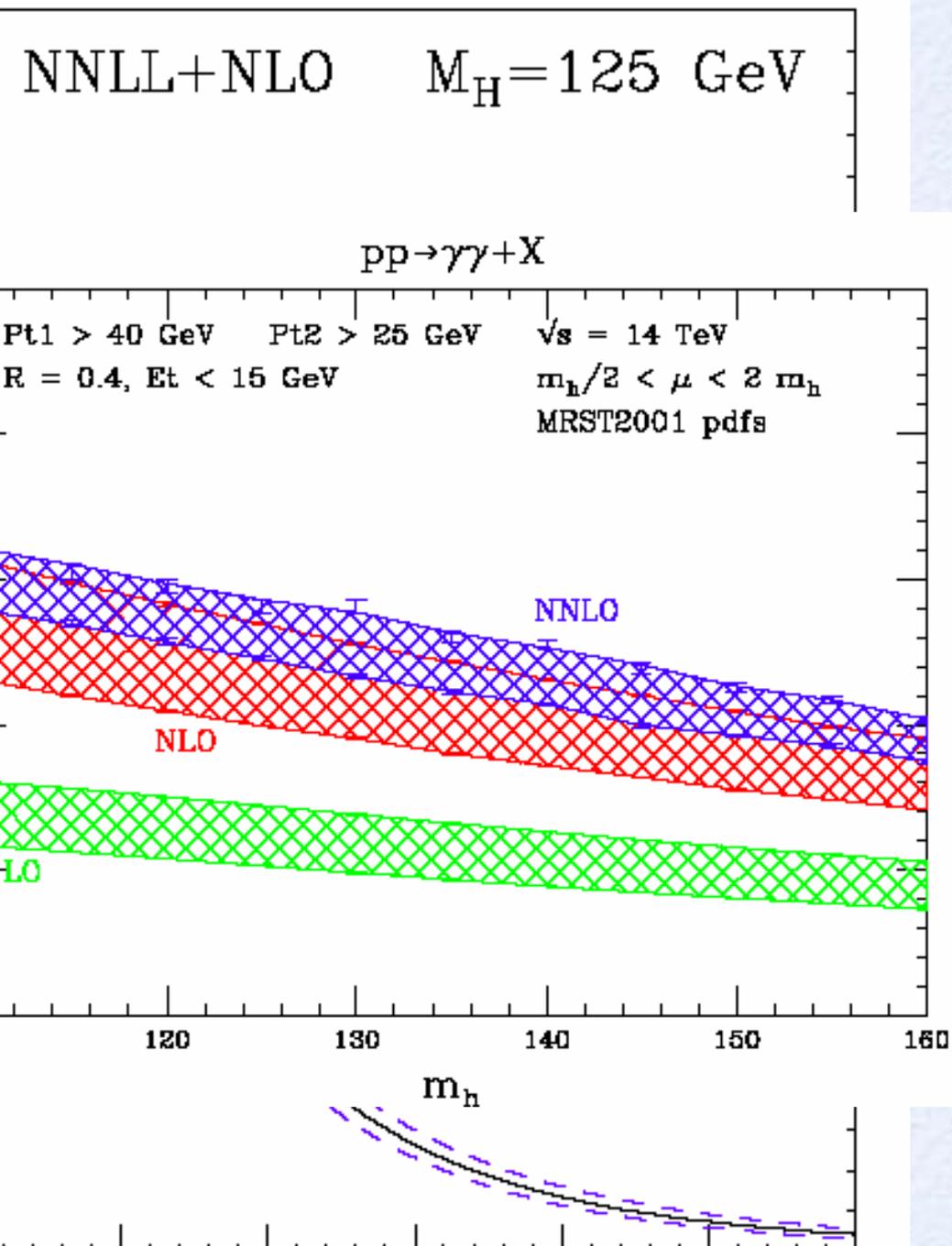


Huge background from QCD.

$qq \rightarrow \gamma\gamma$  known at NLO (DIPHOX) including fragmentation contributions  
[Binoth, Guillet, Pilon, Werlen. 2000]

$gg \rightarrow \gamma\gamma$  direct known at NLO (two-loop)  
[Bern, Dixon, Schmidt. 2002]

On the other hand this is an example of a discovery that does not need an accurate theoretical prediction for the background. Data models suffice.



Dominant production mechanism at hadron colliders.  
 The story of the most accurate prediction in QCD:

QCD corrections:

[Daswon.1991] [Djouadi, Graudenz, Spira, Zerwas. 1997]  
 [Kramer, Laenen, Spira.1998] [Catani, De Florian, Grazzini. 2001]  
 [Harlander, Kilgore.2001,2002] [Anastasiou, Melnikov.2002]  
 [Ravindran,Smith, Van Neerven. 2003]  
 [Catani, De Florian, Grazzini, Nason.2003]

Two-loop EW corrections:

[Djouadi, Gambino, Kniehl. 1998]  
 [Aglietti, Bonciani, Degrandi, Vicini. 2004]  
 [Degrandi, FM. 2004]

PDF evolution at NNLO (“Guinness of QCD”):

[Moch, Vogt, Vermaseren, 2004]

Best QCD predictions at present:

- > Fully exclusive (PS interfaced) prediction at NLO+NLL [Frixione, Webber, 2003]
- > Fully exclusive prediction at NNLO (first) [Anastasiou, Melnikov, Petriello. 2004]



QCD corrections:

[Djouadi, Spira, van der Bij, Z

[Daswon, Kaufmann. 1993]

[Fleischer, Tarasov<sup>2</sup>. 2004]

Two-loop EW corrections:

[Djouadi, Gambino, Kniehl. 1

[Aglietti, Bonciani, Degrassi,

[Degrassi, FM. 2005]

Very small corrections

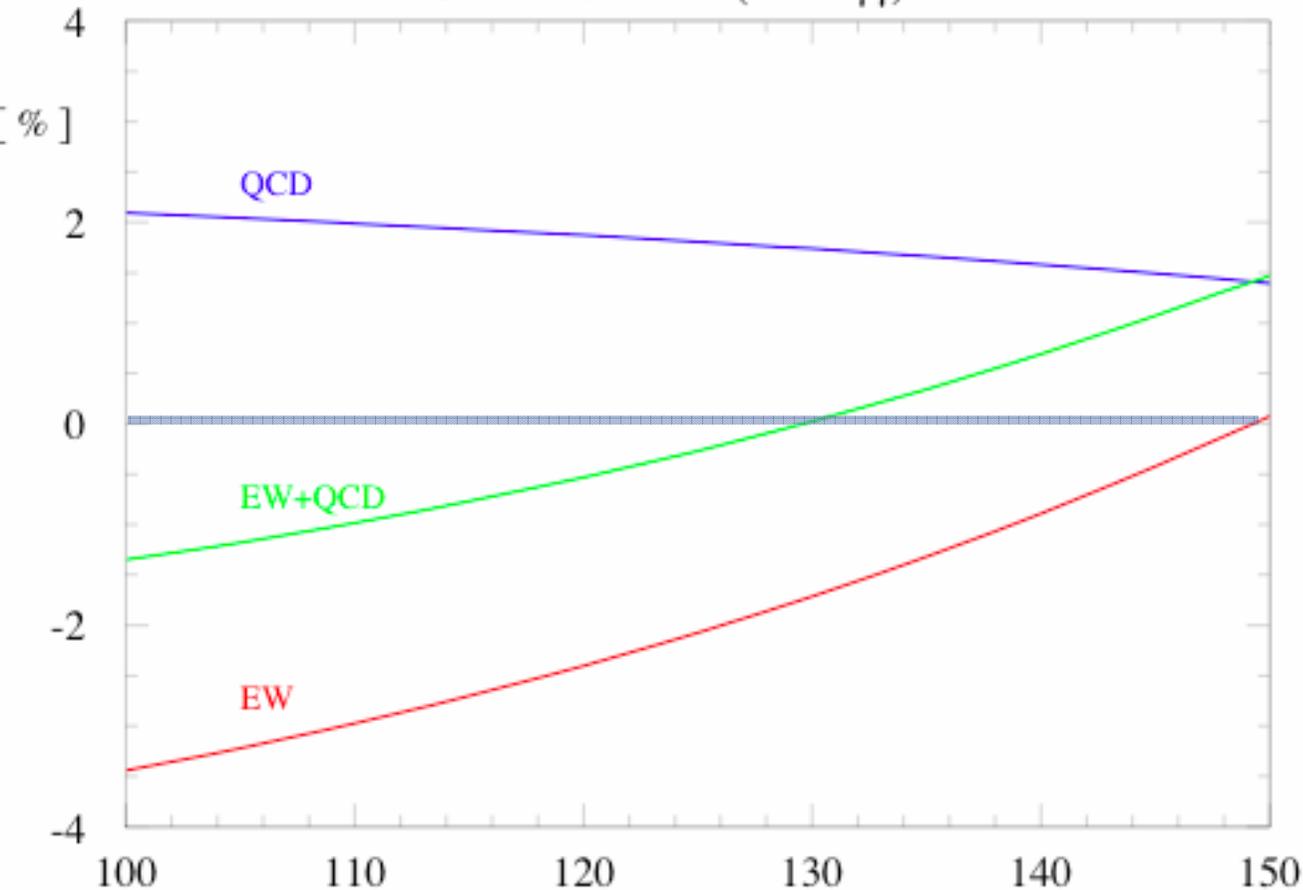
Overall,  $gg \rightarrow H \rightarrow \gamma\gamma$  is

of the most

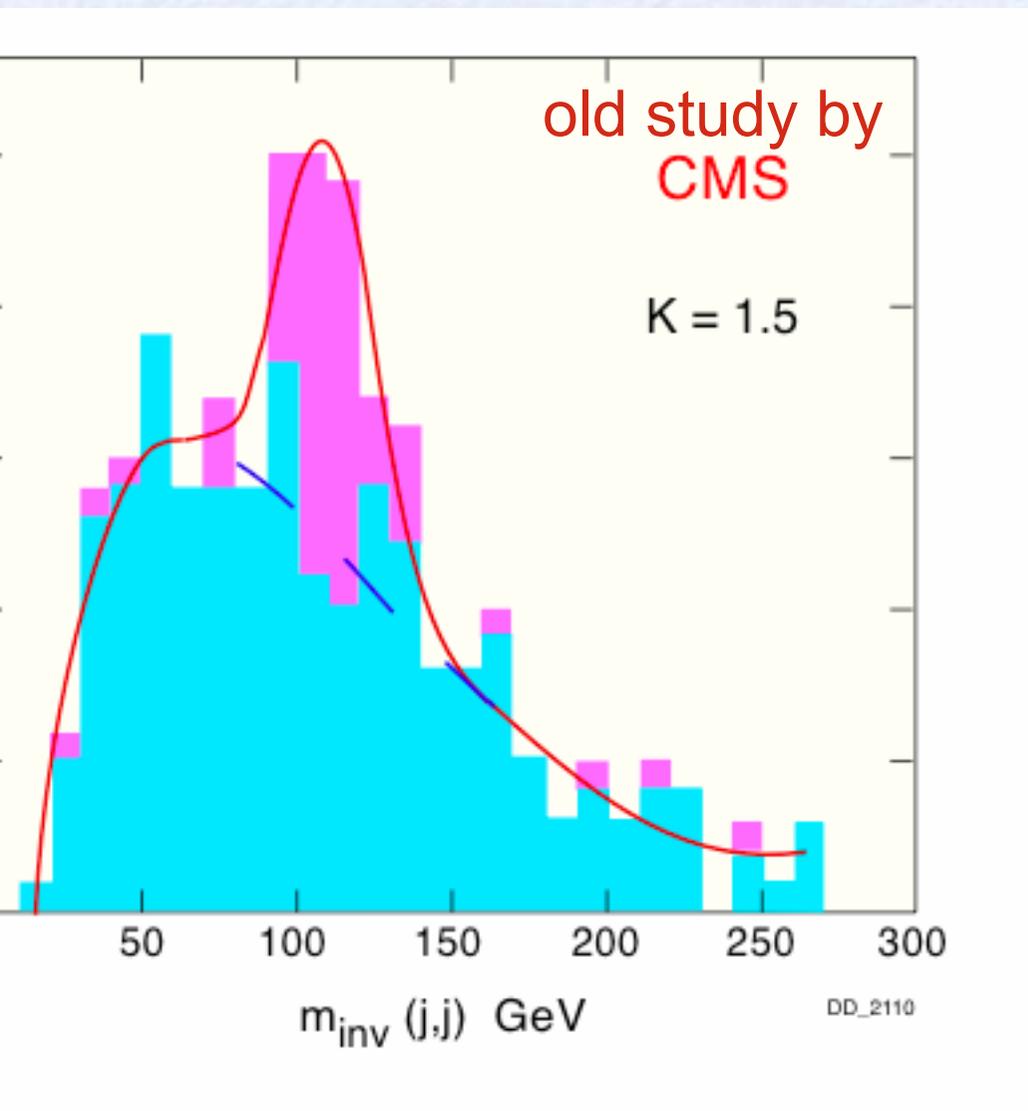
accurate predictions

in hadron collisions!

Corrections to  $\Gamma(H \rightarrow \gamma\gamma)$



# tt production



Signal cross section of  $\sim 1$  pb, known at NLO.  
K-factor of order 1.

[Beenakker, Dittmaier, Kramer, Plumper, Spira, Zanderighi, 2002] [Dawson, Jackson, Orr, Reina, Wackerom]

Typical signature  $4b+2j+l+mEt$ : very difficult!

Key issues:

1. Combinatorics
2. b-tagging
3. Invariant mass resolution
4. Background modeling:  $t\bar{t}b\bar{b}, t\bar{t}j\bar{j}$  are known  
LO  $\Rightarrow$  normalization very uncertain.

Extremely good knowledge of the cross section is necessary.

# Higgs couplings

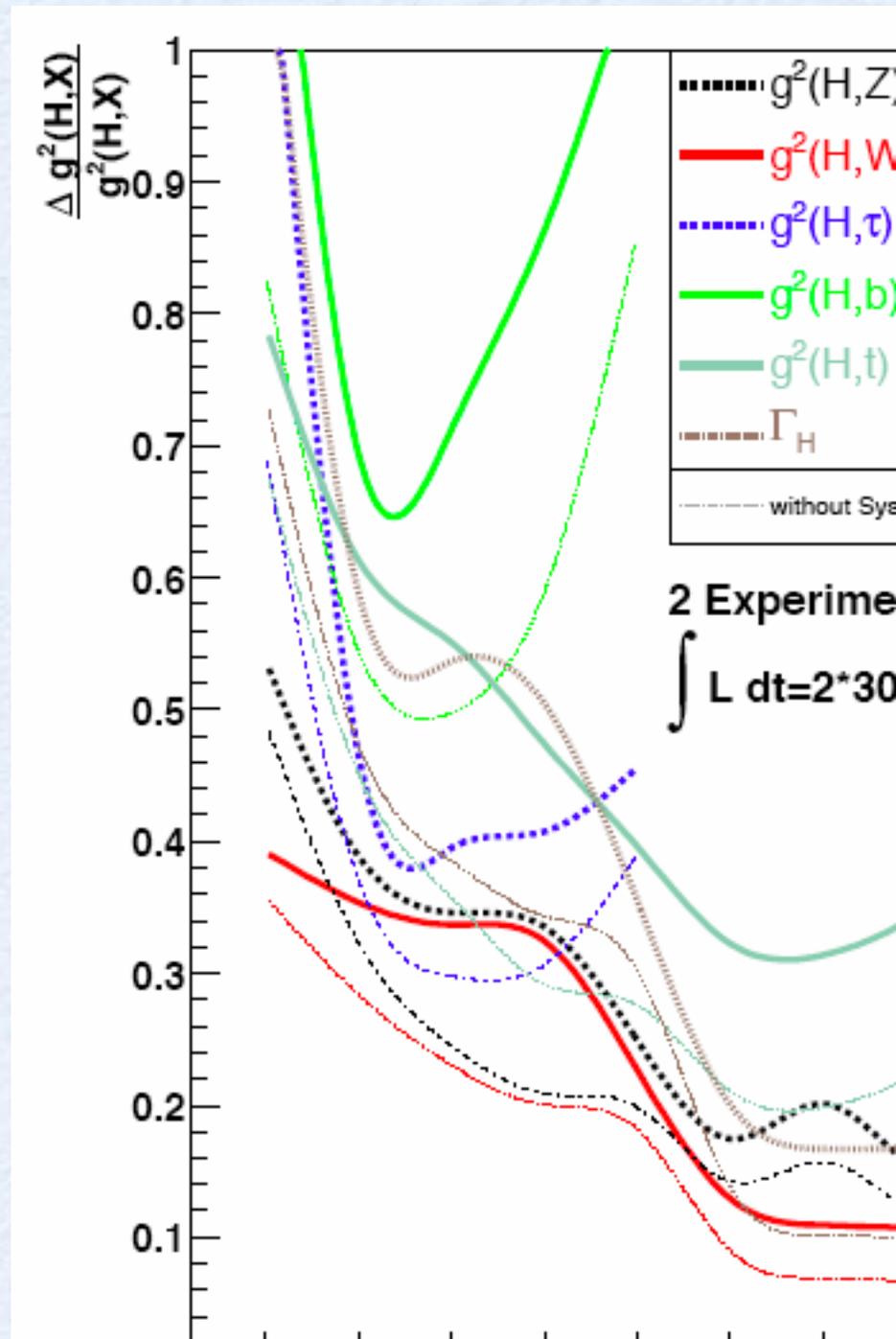
Once a “Higgs” has been found, its SM nature can be tested by measuring its couplings to fermions, bosons as well as the Higgs self-coupling.

Work on all the relevant channels for each mass (e.g.  $gg \rightarrow H$ ;  $qq \rightarrow Hqq$ ;  $qq \rightarrow ZH, WH$ ;  $gg \rightarrow ttH$ ):

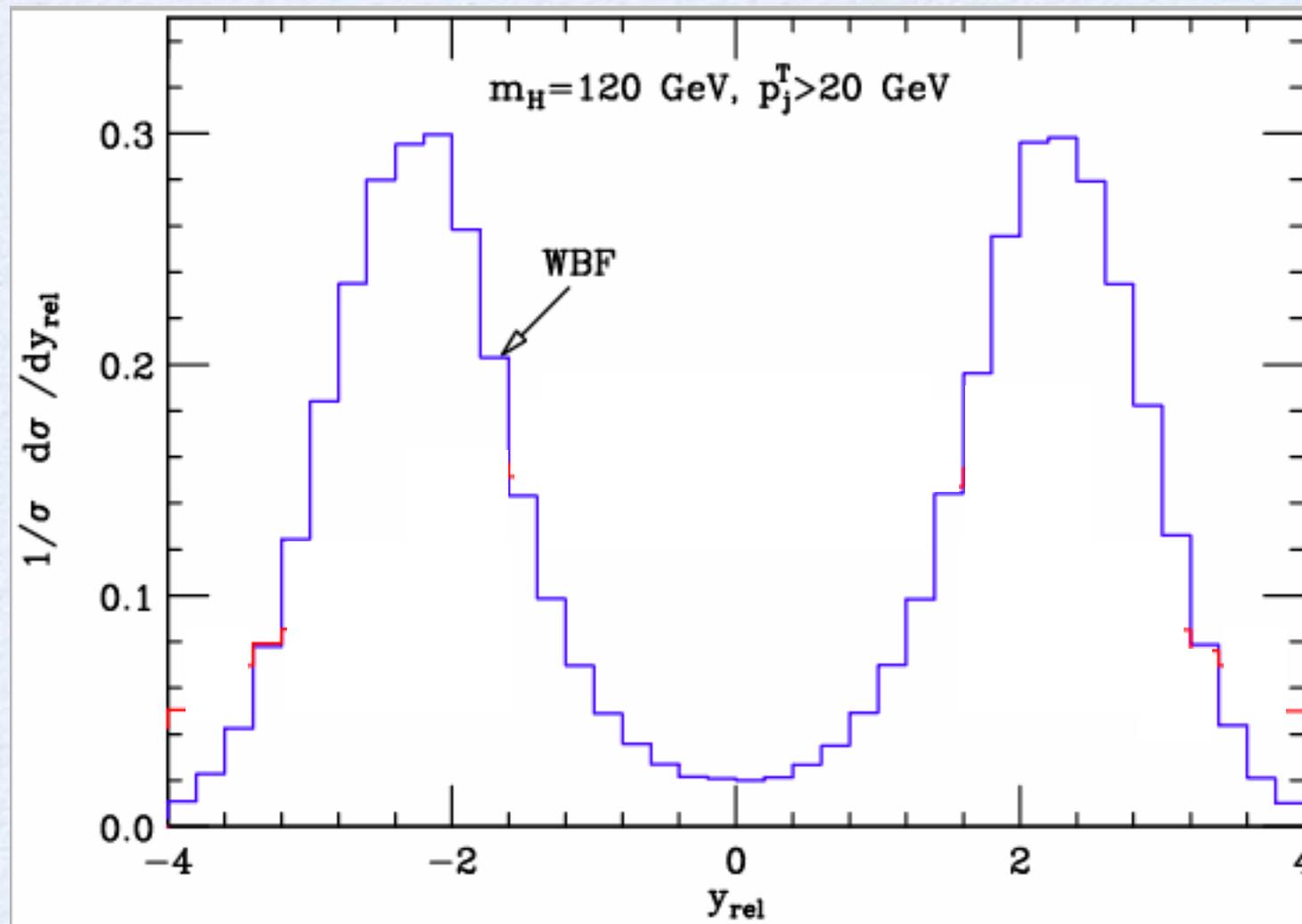
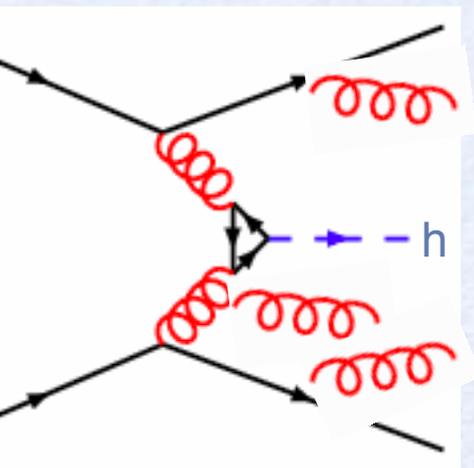
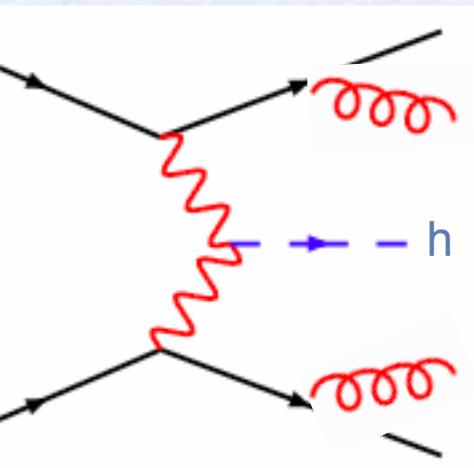
$$\sigma_p \times BR(H \rightarrow xx) = \frac{\sigma_p^{SM}(H)}{\Gamma_p^{SM}} \times \frac{\Gamma_p \Gamma_x}{\Gamma}$$

Accuracy of the couplings of 10-30% can be achieved.

Control of systematic uncertainty from  $gg \rightarrow H$  production at NNLO [Anastasiou, Melnikov, Petriello. 2005]



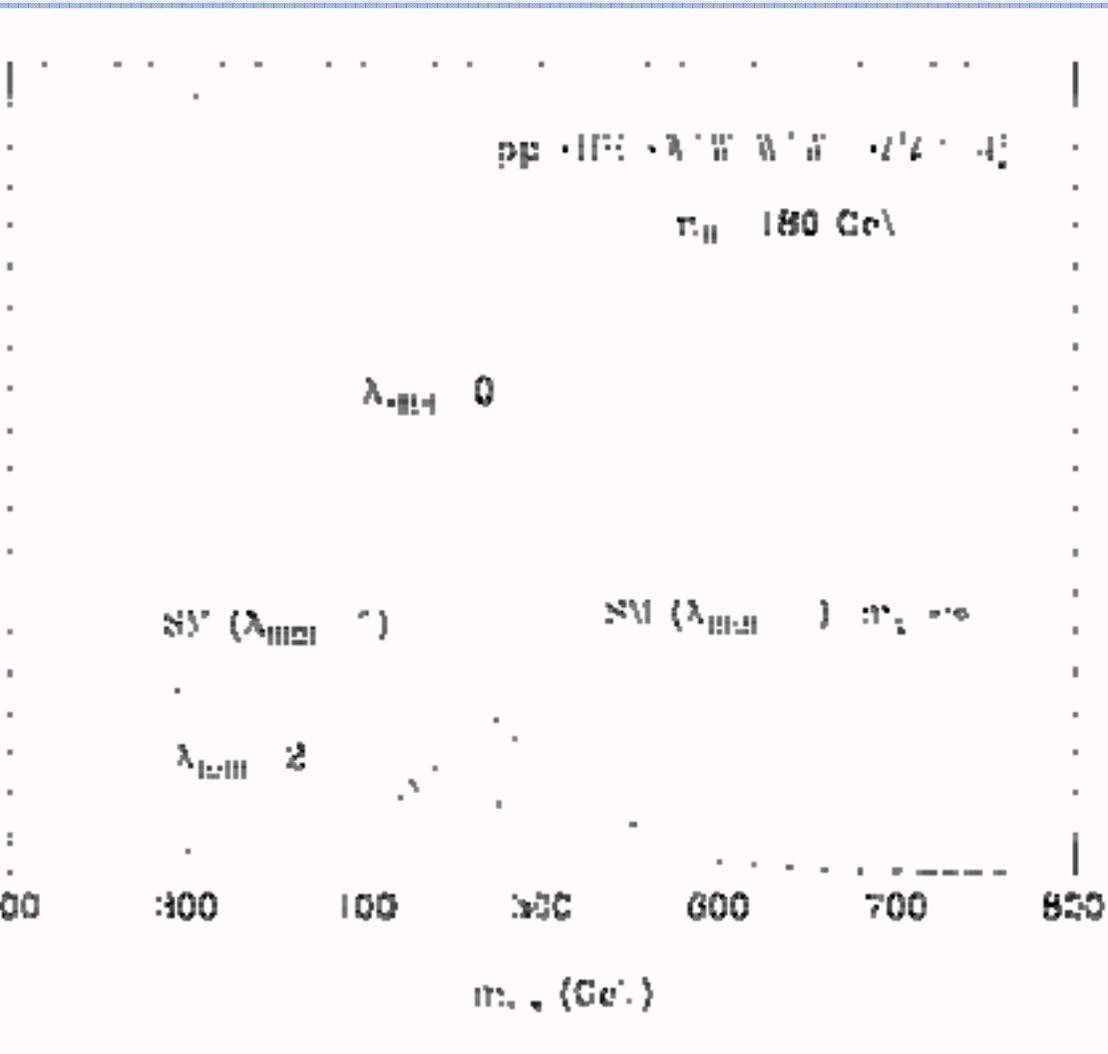
Vector boson fusion will play a crucial role in studying the Higgs properties, in many decay channels (ZZ, WW,  $\pi\pi$ ,  $\Upsilon\Upsilon$ ). Typical signature is two forward jets and a “rapidity gap”. Central jet veto will be essential to select not only signal from background, but also VBF from QCD production.



Central jet veto will be essential to select not only signal from background, but also VBF from QCD production.

# Higgs couplings

$$\mathcal{L}_H = \frac{1}{2}(\partial^\mu H)^2 - \lambda v^2 H^2 - \lambda v H^3 - \frac{\lambda}{4} H^4$$



Higgs potential completely determined by knowledge of the Higgs mass.

$gg \rightarrow HH$ , two diagrams with negative interference

Large top mass limit not a good approximation

NLO corrections known only in the large top limit...

$0 < \lambda < 3$  @ 95% with  $300 \text{ fb}^{-1}$   
with  $160 < m_H < 200$

For smaller Higgs masses  $< 140$ ,  $HH \rightarrow \gamma\gamma b\bar{b}$



Quartic couplings need  $HH$  production and cannot be determined at the LHC

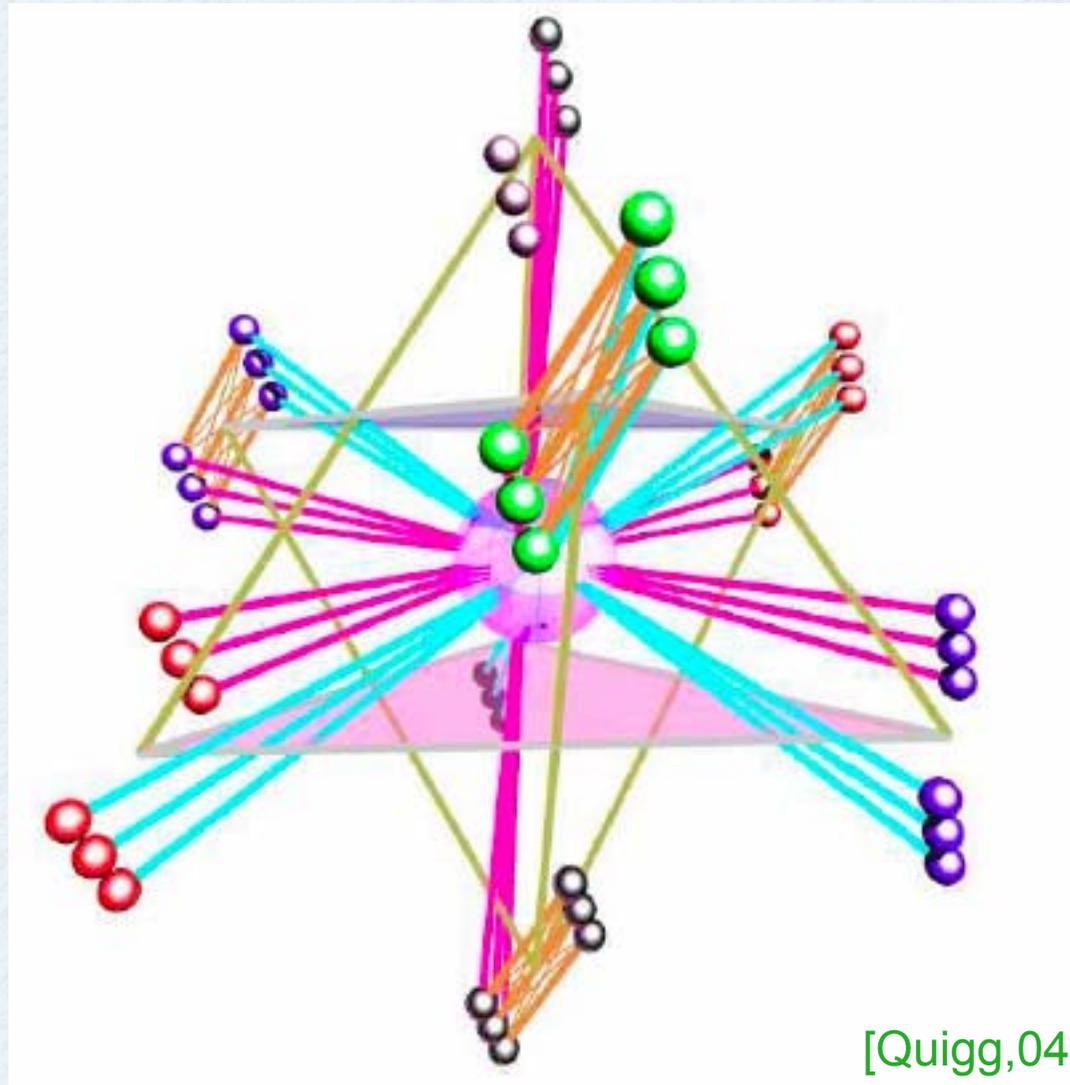
# Outlook

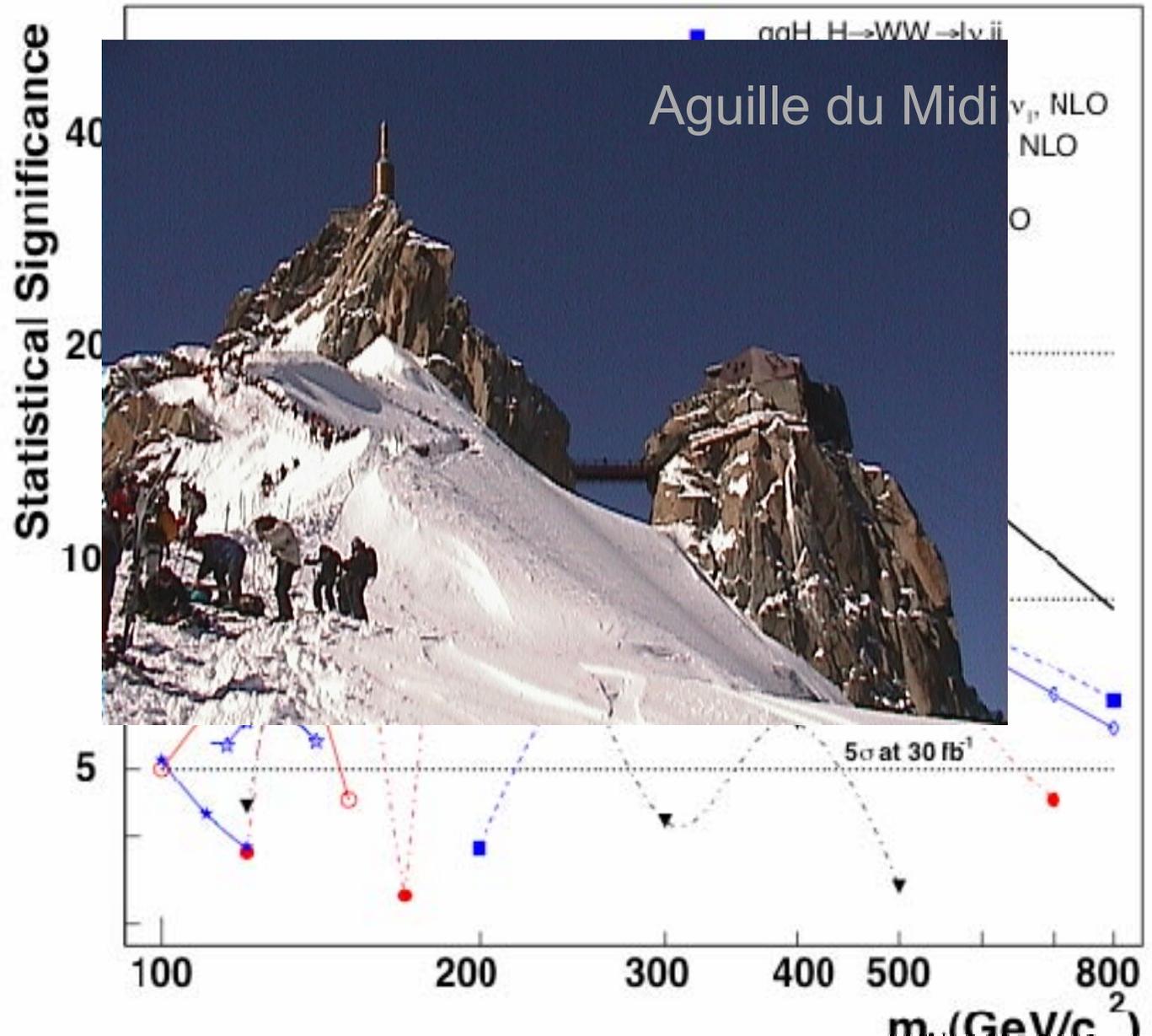
LHC will uncover the mechanism of EWSB and possibly new physics at the TeV scale.

Most probably discoveries will rely not only on our understanding of the detectors but also on our detailed knowledge of SM processes

Many developments and on-going efforts:

- Predictions for signals (and backgrounds) available at an unprecedented level of accuracy  $\Rightarrow$  precision physics possible;
- New analysis strategies developed;
- More data and models to be challenged with.





$$\frac{\Delta g^2(H, X)}{g^2(H, X)}$$



# Tevatron Run II Preliminary

Cross-Section  $\times$  Br (pb)

