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## Higgs production through WBF and gluon fusion at the LHC

in collaboration with the ALPGEN team (M.L. Mangano, M. Moretti, R. Pittau, A.D. Polosa) V. Del Duca and D. Zeppenfeld

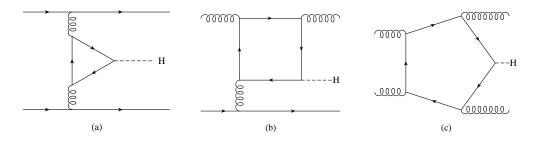
TeV4LHC, 20 October, Fermilab, 2005

Weak Boson Fusion will be an important channel for Higgs boson search at LHC and even more for the determination of its couplings to fermions and gauge bosons

Typical signature: H + 2 forward jets



Another possible contribution is Higgs production through gluon fusion with two additional jets, which is at one loop perturbative order but important because of  $\mathcal{O}(\alpha_s)^4$  against  $\mathcal{O}(\alpha_w^2)$  of WBF



V. Del Duca et al., Phys. Rev. Lett. 87 (2001) 122001

Since QCD H + 2 jets is a background to WBF for Higgs coupling studies, it is important to study the features of the two different

processes

What is available in the literature:

• WBF at QCD NLO

T. Han, G. Valencia and S. Willenbrock, *Phys. Rev. Lett.* 69 (1992) 3274
 T. Figy, C. Oleari and D. Zeppenfeld, *Phys. Rev.* D68 (2003) 073005

• QCD production of H + 2 jets at LO in the limit  $m_t \to \infty$  and keeping the complete  $m_t$  dependence

V. Del Duca et al., Phys. Rev. Lett. 87 (2001) 122001
V. Del Duca et al., Nucl. Phys. B616 (2001) 367
V. Del Duca et al., Phys. Rev. D67 (2003) 073003

• LO calculation of  $pp \to H+3$  jets in the limit  $m_t \to \infty$ 

V. Del Duca, A. Frizzo and F. Maltoni, JHEP 0405 (2004) 064

• very recent improvement on calculation of virtual corrections to H + 4 parton processes, essential ingredient for the calculation of NLO corrections to  $pp \rightarrow H + 2$  jets

K. Ellis, W. Giele and G. Zanderighi, hep-ph/0506196

The parton level analysis carried out so far show particular features (such as a correlation in  $\Delta \phi_{jj}$  between the tagging jets, see below) allowing to distinguish between WBF and gluon fusion production

On the other hand the picture could be changed once the higher order QCD radiation effects are taken into account. This is what has been found in the analysis by

K. Odagiri, JHEP **03030** (2003) 009

where the final state H + 2 jets has been generated starting from the kernel process  $gg \rightarrow H$  and adding QCD radiation with HERWIG Parton Shower, i.e. neglecting exact matrix elements for hard QCD radiation

The situation can be improved if the combined information of matrix elements and Parton Shower can be used. This can be done for instance with the help of the ALPGEN event generator

M.L. Mangano, M. Moretti, F. P., R. Pittau, A.D. Polosa, JHEP 0307 (2003) 001

## In ALPGEN v2.0 the effective coupling ggH in the limit $m_t \to \infty$ has been implemented at the Lagrangian level

M.A. Shifman, A.I. Vainshtein, M.B. Voloshin and V.I. Zakharov, Sov. J. Nucl. Phys 30 (1979) 711J. Ellis, M.K. Gaillard and D.V. Nanopoulos, Nucl. Phys. B106 (1976) 292

This approximation has been shown to be very good for  $m_H$  and  $p_j^T < m_t$ 

V. Del Duca et al., Phys. Rev. Lett. 87 (2001) 122001
V. Del Duca et al., Phys. Rev. D67 (2003) 073003

The process  $pp \to H + N$  jets (N < 5) has been introduced (and interfaced to the Parton Shower) with exact LO QCD matrix elements for up to N additional partons

Our aim is to study the effects of higher order QCD radiation (and eventually hadronisation) with the HERWIG Parton Shower on top of parton level events of the form  $pp \rightarrow H + 2$  jets and  $pp \rightarrow H + 3$ generated with ALPGEN  $m_H = 120 \text{ GeV}$ 

Event selection:

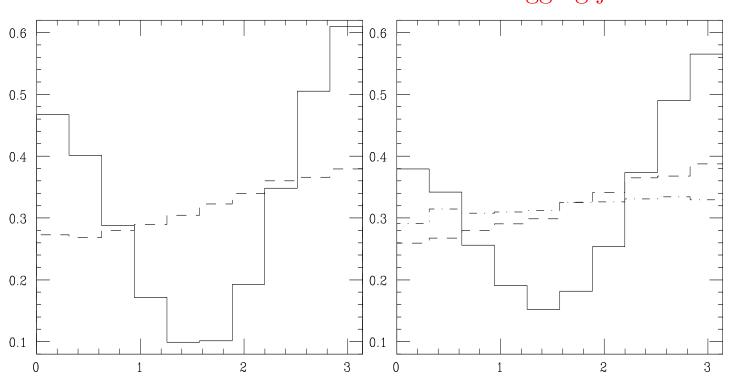
$$p_{\perp}^{j} \ge 20 \text{ GeV}, \quad |\eta_{j}| \le 5, \quad R_{jj} \ge 0.6$$
  
 $|\eta_{j1} - \eta_{j2}| \ge 4.2, \quad \eta_{j1} \cdot \eta_{j2} \le 0, \quad M_{jj} \ge 600 \text{ GeV}$ 

## PDF set: CTEQ5L

Scales (affecting much more QCD Higgs productions than WBF):

$$\alpha_s^{2+N \text{jets}}(\mu_R) \rightarrow \alpha_s^2(M_H) \Pi_i \alpha_s(p_i^T)$$
$$\mu_F = (\Pi_i p_i^T)^{(1/N)}$$

Jets defined according to the routine GETJET, which uses a simplified version of the UA1 jet algorithm. Parameters for the jets as for the partons

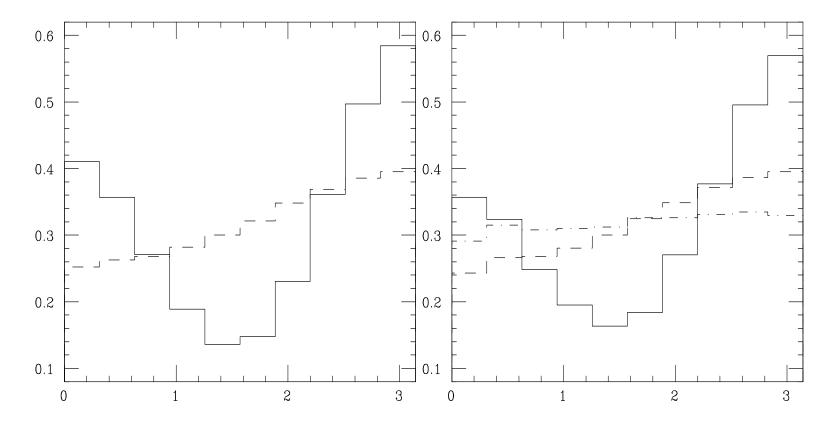


Azimuthal correlation between the tagging jets

Left:  $\frac{1}{\sigma} d\sigma / d\phi_{jj}$  distribution at LO partonic level for the process  $pp \to H + 2$  jets. Solid line: QCD Higgs production; dashed line: WBF

Right:  $\frac{1}{\sigma} d\sigma / d\phi_{j_1 j_2}$  distribution with Parton Shower  $(j_1 \text{ and } j_2 \text{ are the leading } p_T \text{ jets})$  on top of  $pp \to H + 2$  jets generated events

Right: Dot-dashed line has been obtained generating  $gg \to H$  with <code>HERWIG</code> and taking all jets from shower

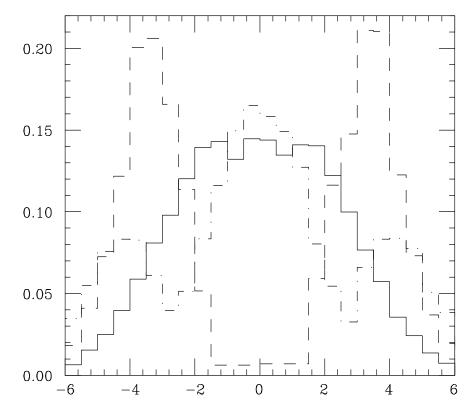


Left:  $\frac{1}{\sigma} d\sigma / d\phi_{j_1 j_2}$  distribution at LO partonic level for the process  $pp \to H + 3$  jets. Solid line: QCD Higgs production; dashed line: WBF

Right:  $\frac{1}{\sigma} d\sigma / d\phi_{j_1 j_2}$  distribution with Parton Shower  $(j_1 \text{ and } j_2 \text{ are the leading } p_T \text{ jets})$  on top of  $pp \to H + 3$  jets generated events

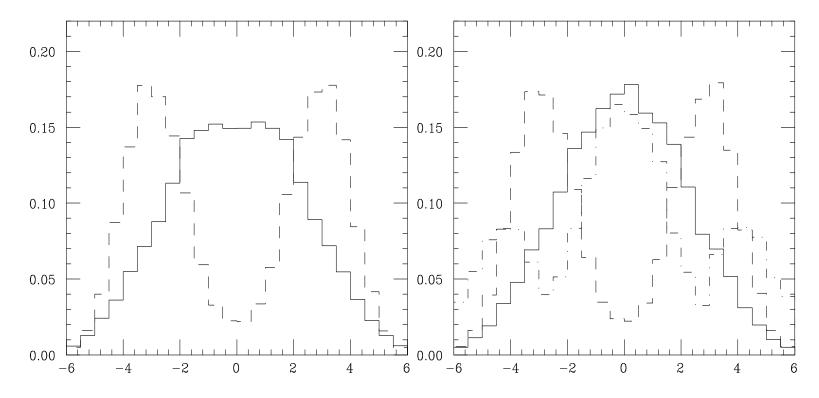
The presence of additional hard radiation doesn't change the  $\Delta\phi_{jj}$  azimuthal correlation





 $\frac{1}{\sigma}d\sigma/dy_{\rm rel}$  distribution with Parton Shower on top of  $pp \to H + 2$  jets generated events; dashed WBF, solid QCD Higgs production, dot-dashed Parton Shower on top of  $pp \to H$  generated events

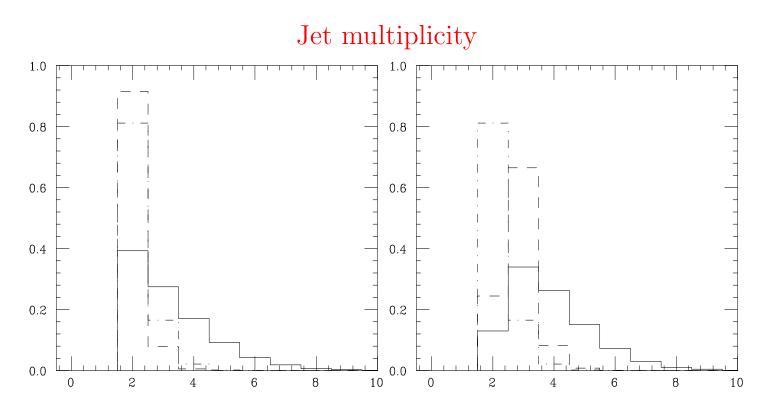
$$y_{\rm rel} = y_3 - (y_1 + y_2)/2$$



Left:  $\frac{1}{\sigma} d\sigma/dy_{\rm rel}$  distribution at LO partonic level for the process  $pp \to H + 3$  jets. Solid line: QCD Higgs production; dashed line: WBF

Right:  $\frac{1}{\sigma}d\sigma/dy_{\rm rel}$  distribution with Parton Shower on top of  $pp \to H + 3$  jets generated events

In WBF the emission of hard radiation starts to populate the central rapidity region but the shape of the parton level simulation is the same as after showering. For gluon fusion the third jet after showering is more likely to be emitted central in rapidity w.r.t. pure parton level



Left:  $\frac{1}{\sigma} d\sigma/dn_{\text{jets}}$  distribution with Parton Shower on top of  $pp \to H + 2$  jets generated events. Solid line: QCD Higgs production; dashed line: WBF; dot-dashed line: Parton Shower on top of  $pp \to H$ 

Right:  $\frac{1}{\sigma} d\sigma / dn_{\text{jets}}$  distribution with Parton Shower on top of  $pp \to H + 3$  jets generated events

While in WBF the number of jets is peaked at the number of final-state partons, a large part of the jets which pass the cuts in gg fusion originates from the parton shower

$$A_{\phi} = \frac{\sigma(\Delta\phi < \pi/4) - \sigma(\pi/4 < \Delta\phi < 3\pi/4) + \sigma(\Delta\phi > 3\pi/4)}{\sigma(\Delta\phi < \pi/4) + \sigma(\pi/4 < \Delta\phi < 3\pi/4) + \sigma(\Delta\phi > 3\pi/4)}$$

where  $\Delta\phi$  is the azimuthal distance between the two tagging jets

| $A_{\phi}$ for x partons + H f.s. | parton level | shower level |
|-----------------------------------|--------------|--------------|
| $A_{\phi}(0j - ggH)$              |              | 0.006        |
| $A_{\phi}(2j - ggH)$              | 0.50         | 0.33         |
| $A_{\phi}(3j - ggH)$              | 0.38         | 0.31         |
| $A_{\phi}(2j - WBF)$              | 0.015        | 0.017        |
| $A_{\phi}(3j - WBF)$              | 0.013        | 0.014        |

Work in progress to perform an inclusive analysis adding together samples with different jet multiplicities, without double counting and using proper reweighting of  $\alpha_s$ , with MLM prescription for jet matching (now available in ALPGEN v2.0)

| σ                  | fixed scale $\alpha_s(m_H)$ | modified*            | $lpha_s^{ m rew}$   |
|--------------------|-----------------------------|----------------------|---------------------|
| $\sigma(2j - ggH)$ | 0.50(1)  pb                 | $0.73(1) { m ~pb}$   | $0.66(1) {\rm ~pb}$ |
| $\sigma(3j - ggH)$ | 0.36(1)  pb                 | 0.44(1)  pb          | 0.49(1)             |
| $\sigma(2j - WBF)$ | 1.36(1)  pb                 | 1.41(3)  pb          |                     |
| $\sigma(3j - WBF)$ | 0.22(1)  pb                 | $0.26(1) {\rm \ pb}$ | 0.20(1)             |

\*)

$$\alpha_s^{2+N \text{jets}}(\mu_R) \rightarrow \alpha_s^2(M_H) \Pi_i \alpha_s(p_i^T)$$
$$\mu_F = (\Pi_i p_i^T)^{(1/N)}$$

according to

V. Del Duca, A. Frizzo and F. Maltoni, JHEP 0405 (2004) 064

Possible check vs NLO calculation for WBF, for ggH will be possible in the future

## Summary

- We proved the importance of the exact LO matrix element calculations for *H*+ jets final states in order to exploit all the correlations allowing to disentangle WBF from QCD Higgs production
- ALPGEN v2.0 includes such processes with up to five jets in the final state
- While the use of the Parton Shower on the process gg → H can give unreliable results, once the complete matrix element for H+ jets is used, the shower doesn't alter the picture
- This has still to be proved when including hadronization (work in progress)
- A good integral quantity  $A_{\phi}$  has been identified to distinguish between gluon fusion and WBF
- The use of exact matrix elements for H+ jets final state calls for a consistent matching between matrix elements and parton shower (work in progress)
- The availability of a NLO calculation would be useful in the study of the optimal scales