

ECM'08
Electron Cloud Mitigation 2008 (CARE-HHH Mini-Workshop)
20 – 21 November 2008, CERN

MEST

Multipactor Effect Simulation Tool

Work done under ESA AO 4025:
Surface Treatment and Coating for the Reduction of Multipactor
and Passive Intermodulation (PIM) Effects in RF Components

by



J. de Lara, F. Pérez, M. Alfonseda, [L. Galán](#), (UAM)

I. Montero, E. Román, (CSIC)

D. Raboso (ESA)



Overview

For detecting multipactor in infinite parallel plate geometry

Determines initial multipactor susceptibility

For studying influence of secondary electron emission properties

Outline

Simplifications and limitations

Complete detailed simulation of electron cloud

Simple detailed realistic Secondary Electron Emission model

Results

Final comments

Electromagnetic field extremely simplified

Infinite parallel plate

Only electric field

$$F_z(t) = e \frac{V_o}{d} \sin(\omega t)$$

Homogeneous every where, only dependent of time

No relativistic effects

No space charge effects

Each electron moves alone in the unperturbed electric field

Each electron trajectory calculated independently

Exact analytical trajectory: emission event \Rightarrow impact event

event: time, position, velocity

No need to calculate trajectories, series of (time, position)

Initial stage or tendency (susceptibility) of multipactor discharge

All the electrons in the cloud are simulated individually

Detailed discrete branching process of multiplications and absorptions

Electron cloud: impact event queue ordered by time to be processed

Electron: object defined by impact event, can relate to emission event

At impact event, new electrons are generated by Secondary Emission

The simulations discrete event approach

loops along the queue

advance simulation time

generate new events

placing new event orderly in the queue

No resonant conditions are imposed

Simulation procedure

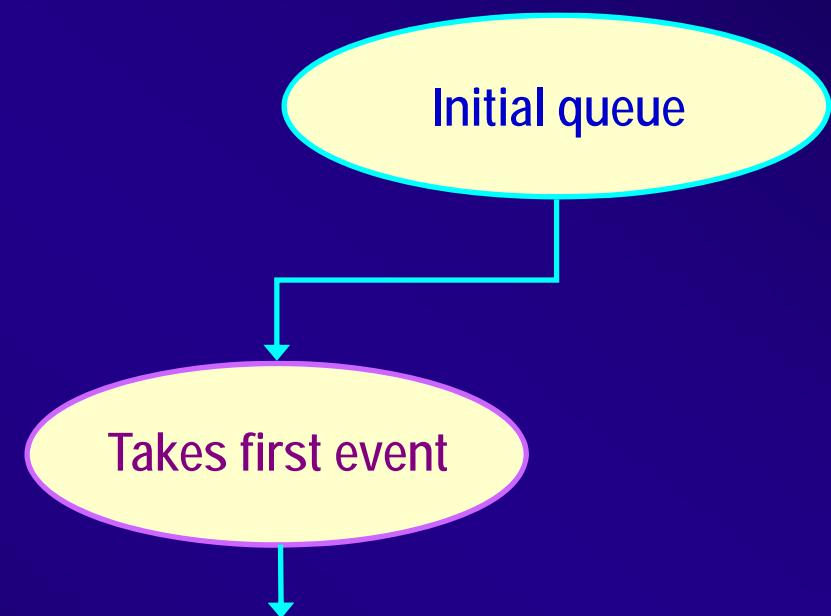
Initial queue: initial electrons, seeding electrons

Random Distribution (time, position, velocity) number \leftarrow user

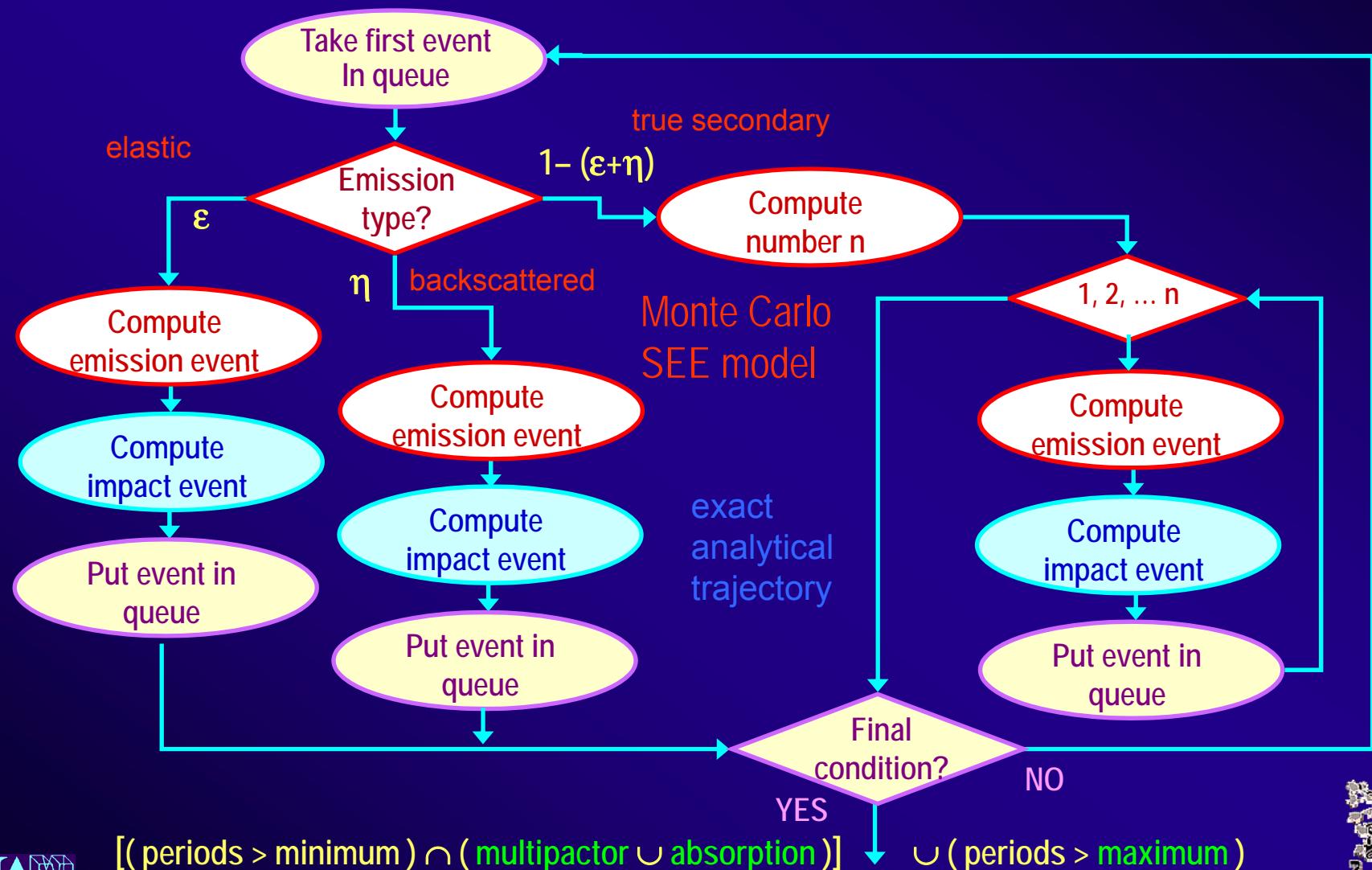
time $\in \mathcal{D}\text{Uniform}(\text{first period})$

position $\in \mathcal{D}\text{Uniform}(\text{first plate})$

velocity $\in \mathcal{D}\text{Normal}(\text{initial velocity})$



Simulation procedure



Monte Carlo Secondary Electron Emission Model

Impact event = (time, position, velocity)

velocity $\Leftrightarrow (E_p, \varphi)$

emission event = (time, position, velocity)

velocity $\Leftrightarrow (E, \theta, \zeta)$

TYPE OF EMISSION EVENT

Elastically reflected electron

probability \Leftarrow SEE yield functions

$$P_e(E_p, \varphi) = \varepsilon(E_p, \varphi)$$

$$P_b(E_p, \varphi) = \eta(E_p, \varphi)$$

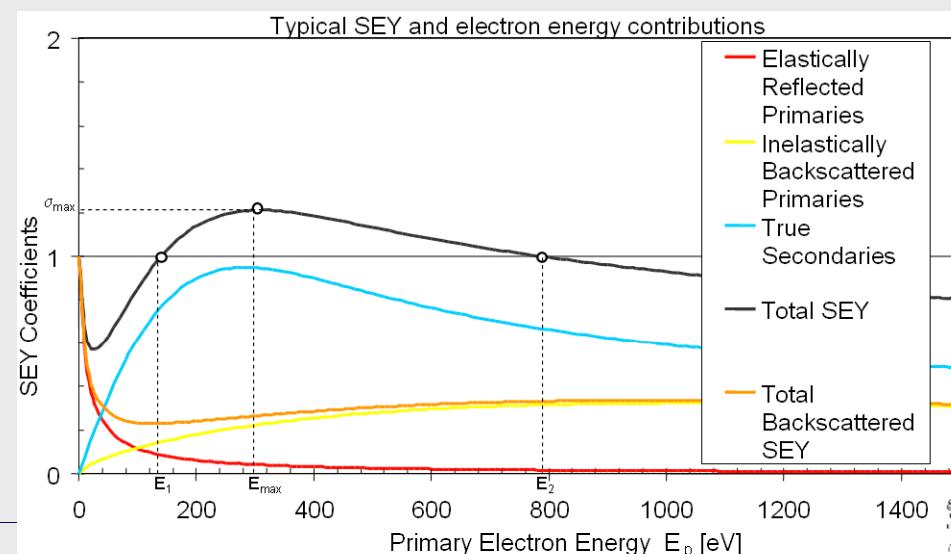
$$P_s(E_p, \varphi) = 1 - P_e(E_p, \varphi) - P_b(E_p, \varphi)$$

Inelastically backscattered electron

True secondary emission of n electrons

$n \in \mathcal{D}$ Poisson(impact velocity)

$$\langle n \rangle = \delta(E_p, \varphi) / P_s(E_p, \varphi)$$



Monte Carlo Secondary Electron Emission Model

SEE YIELD FUNCTIONS empirical

$$\varepsilon(E_p, \varphi) = \varepsilon(E_p)^{\cos(\varphi)} \cdot C_2^{1-\cos(\varphi)}$$

$$C_2 = \chi \cdot \frac{\varepsilon(E_p)}{\eta(E_p) + \varepsilon(E_p)}$$

$$\varepsilon(E_p) = \frac{\varepsilon_1}{1 + \frac{E_p}{E_{e_1}}} + \frac{\varepsilon_2}{1 + \frac{E_p}{E_{e_2}}}$$

$$\varepsilon(E_p \approx 0) = 0$$

$$\eta(E_p, \varphi) = \eta(E_p)^{\cos(\varphi)} \cdot C_1^{1-\cos(\varphi)}$$

$$C_1 = \chi \cdot \frac{\eta(E_p)}{\eta(E_p) + \varepsilon(E_p)}$$

$$\eta(E_p) = a \cdot (1 - b \cdot E_p) \cdot E_p^\gamma \cdot \exp\left(-\left(\frac{E_p}{E_b}\right)^\mu\right)$$

$$\delta(E_p, \varphi) = \delta(E_p) \frac{k+1}{k + \cos(\varphi)}$$

$$\delta(E_p) = \delta_m \cdot \frac{s \cdot \frac{E_p}{E_m}}{s - 1 + \left(\frac{E_p}{E_m}\right)^s}$$

constants \Leftarrow Material properties

Monte Carlo Secondary Electron Emission Model

EMISSION EVENT: EMISSION ENERGY AND ANGLES

Elastically reflected electron

$$E = E_p, \theta = \varphi, \zeta = 0$$

Inelastically backscattered electron

$$E = E_p \cdot G_b(u) \quad \text{where } u = \in \mathcal{D}\text{Uniform}[0, 1]$$

$$G_b(u) = \alpha^{\frac{1}{n_b}} \cdot (\arccos(1 - \beta \cdot u))^{\frac{1}{n_b}} \quad \beta = 1 - \cos(\alpha) \quad \alpha = \pi \cdot X_{cb}^{n_b}$$

inverse cumulative probability function, empirical

$$X = G(u) \in \text{Distribution(probability } f(X)) [0, 1]$$

$$\text{where } f(X) = \frac{dF}{dX} \quad F \equiv G^{-1}$$

$$\theta = \varphi, \zeta = 0$$

constants \Leftarrow Material properties

Monte Carlo Secondary Electron Emission Model

EMISSION EVENT: EMISSION ENERGY AND ANGLES

True secondary emission of n electrons

$$E = E_{remain} \cdot G_b(u) \quad G_s(u) = \left(\frac{2}{\pi} \cdot \arctan \left(\sqrt{\tan\left(\frac{\pi}{2} \cdot X_{cs}\right) \cdot \tan\left(\frac{\pi}{2} \cdot u\right)} \right) \right)^{\frac{1}{n_s}}$$

$$E_{remain}(\text{initial}) = E_p$$

$$E_{remain}(\text{next}) = E_{remain} - E \quad \text{ensures conservation of energy, realistic}$$

$$X_{cs}(E_{remain}, \text{material})$$

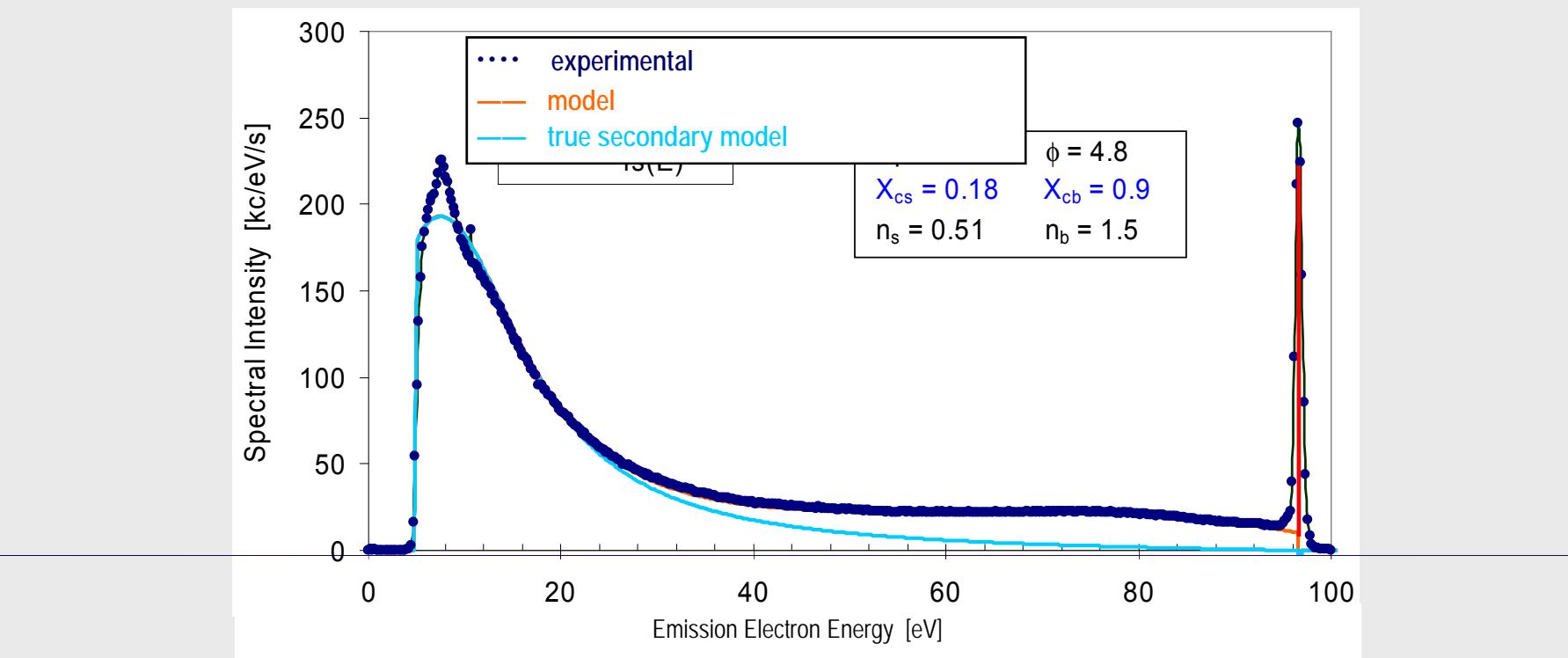
$$\theta = \arcsin\left(\sqrt{x^2 + y^2}\right) \quad \zeta = \arctan(y/x) \quad (x, y) \in \mathcal{D}\text{Uniform}(\text{Circle } x^2 + y^2 \leq 1)$$

constants \Leftarrow Material properties

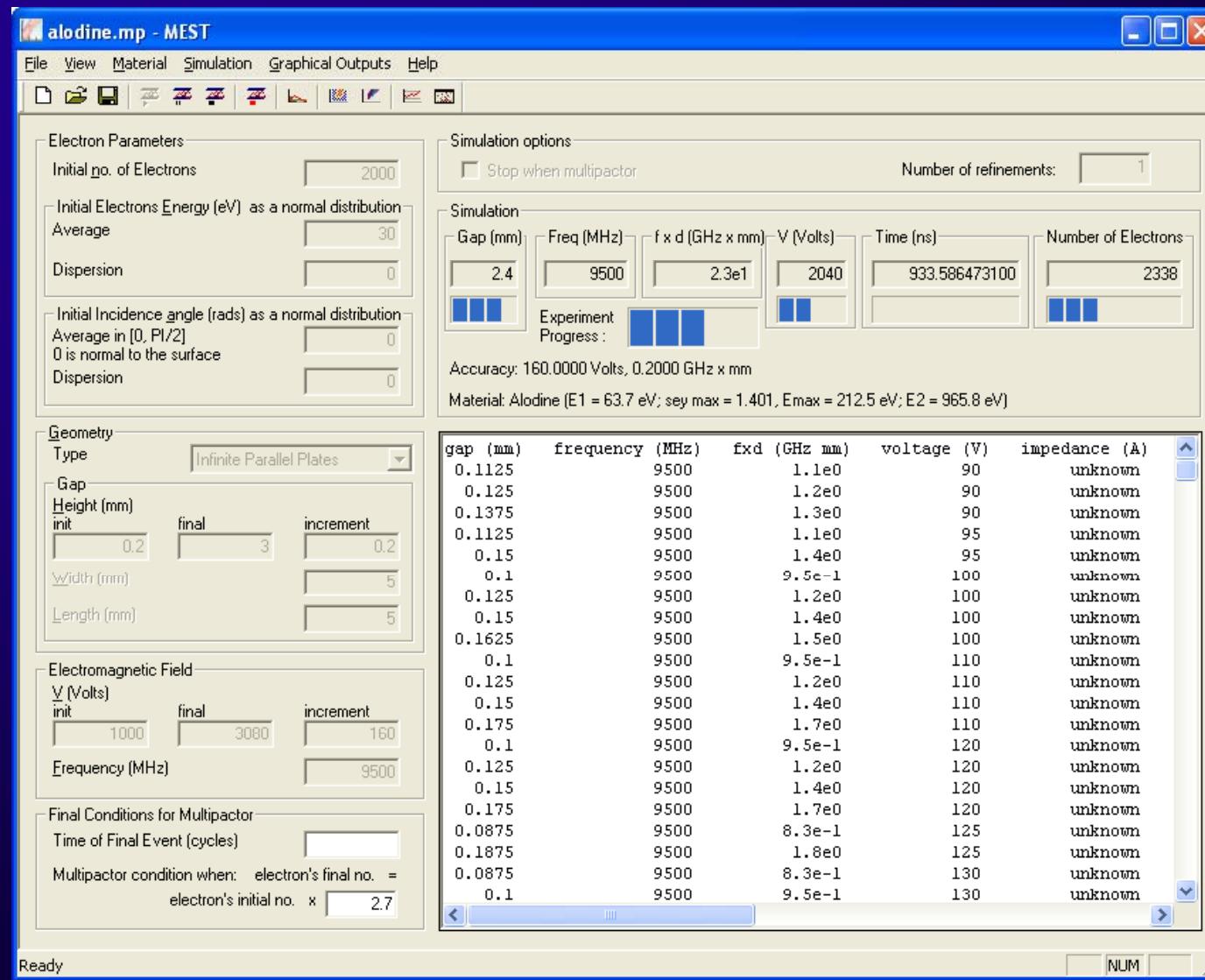
Monte Carlo Secondary Electron Emission Model

EMISSION PROBABILITY FUNCTIONS AND EXPERIMENTAL EDC

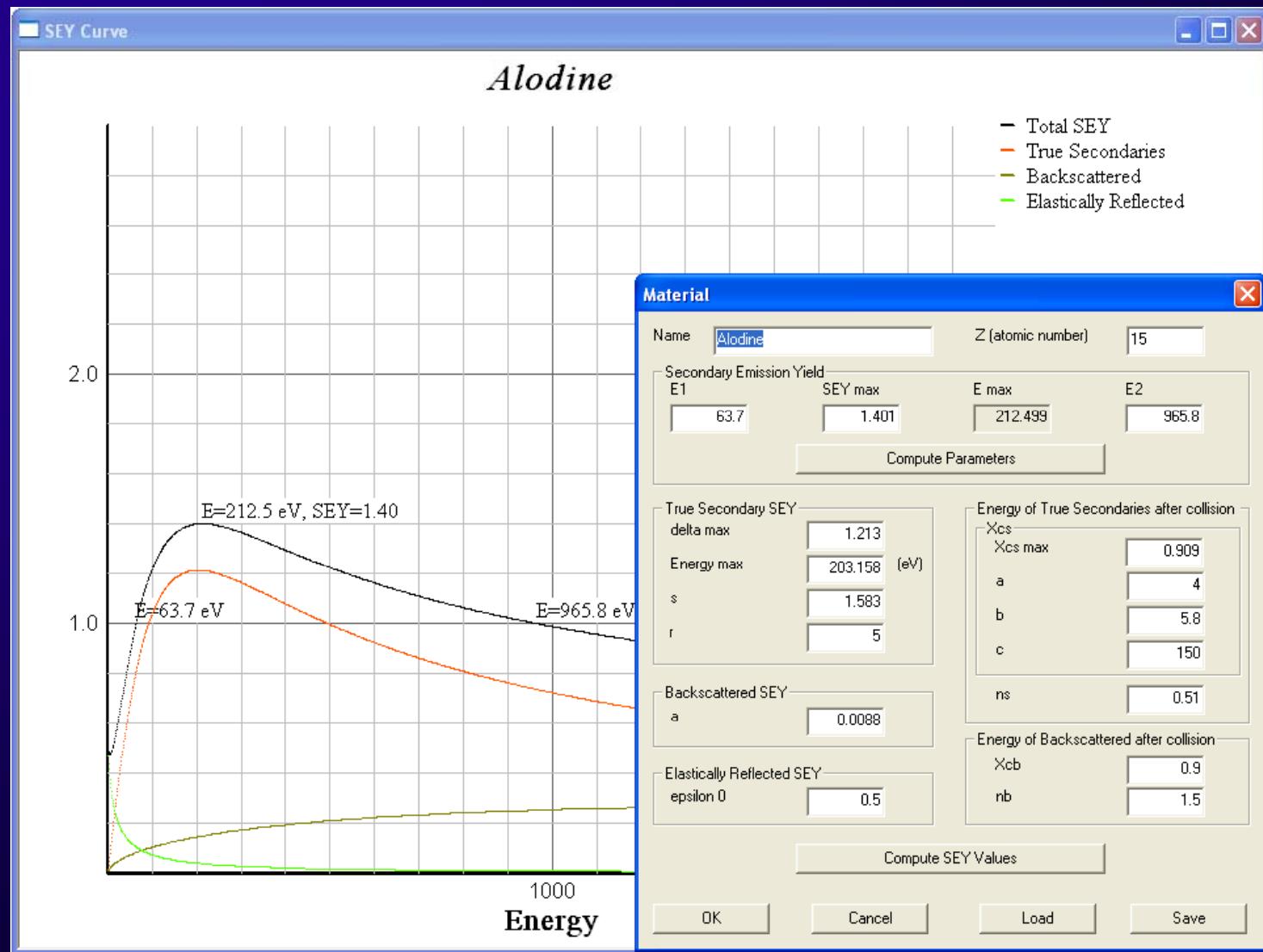
Clean Cu



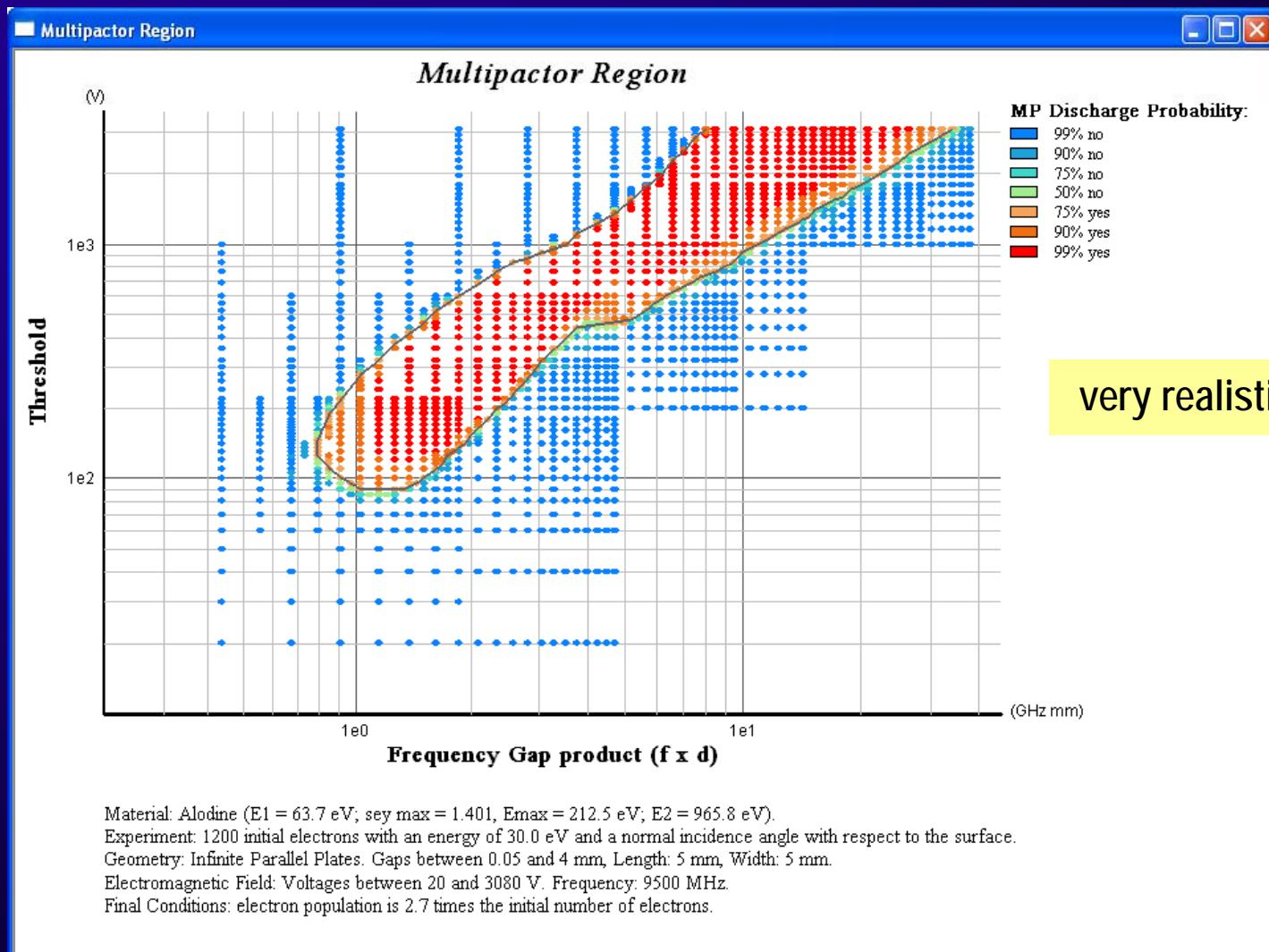
MEST main window



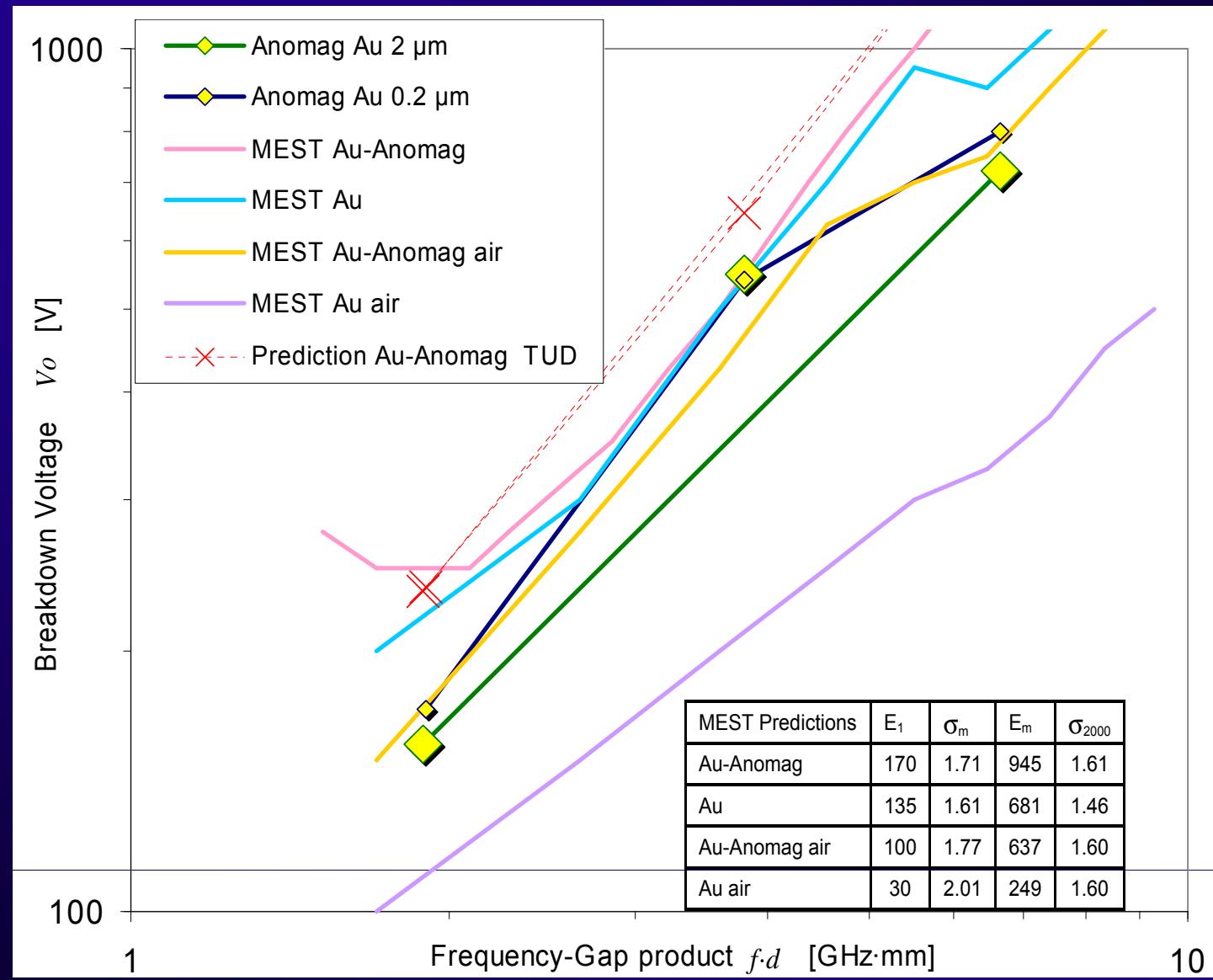
MEST material SEE properties



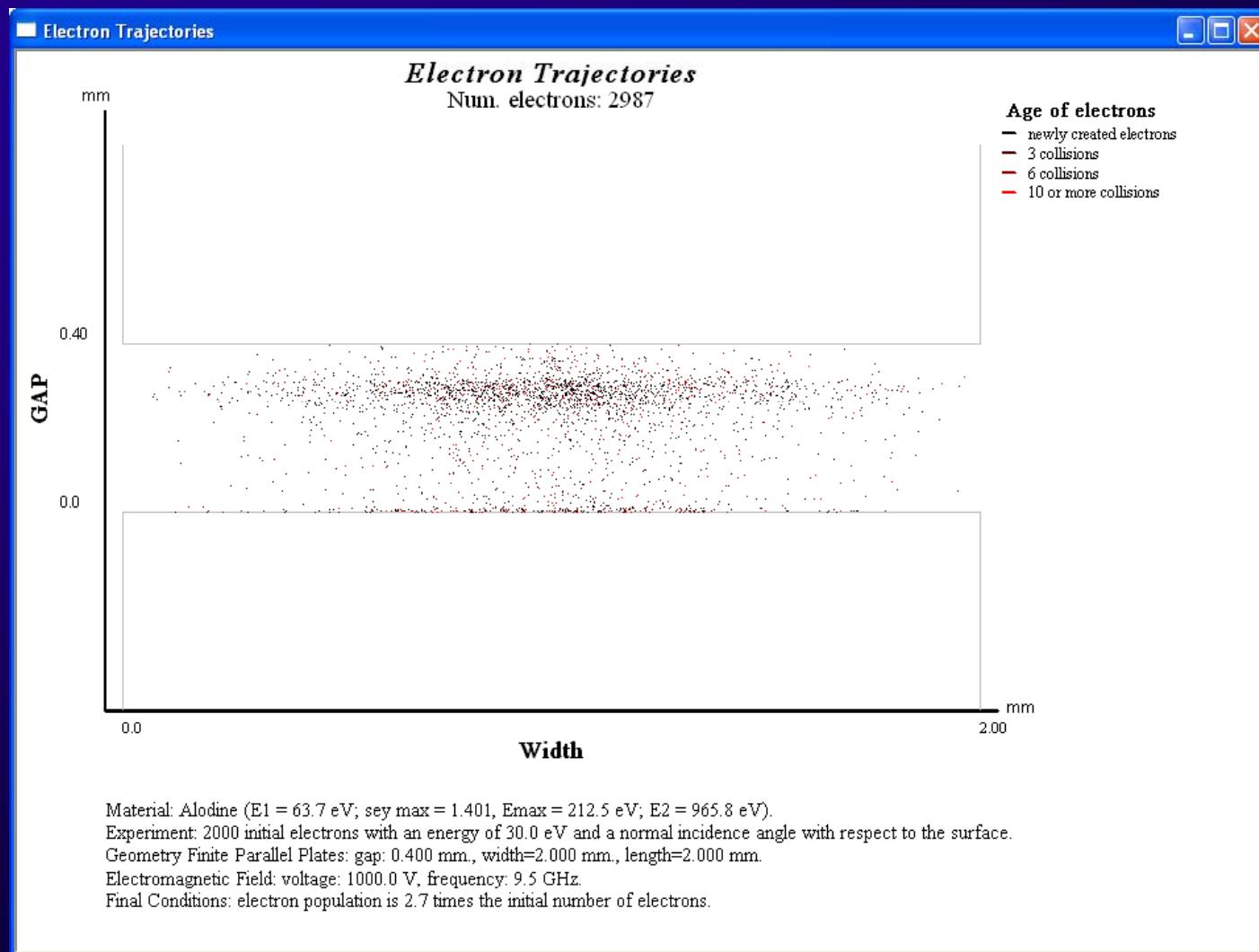
MEST results: multipactor region



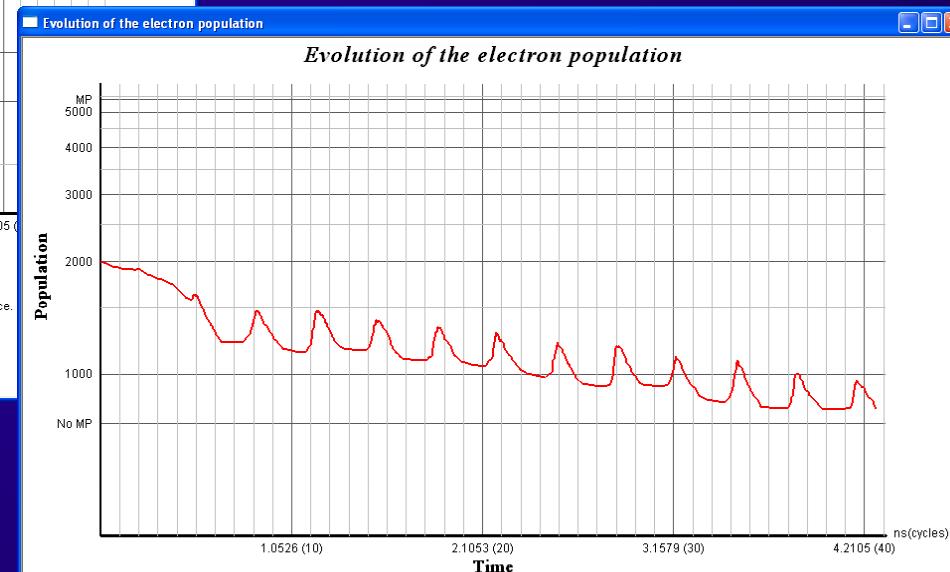
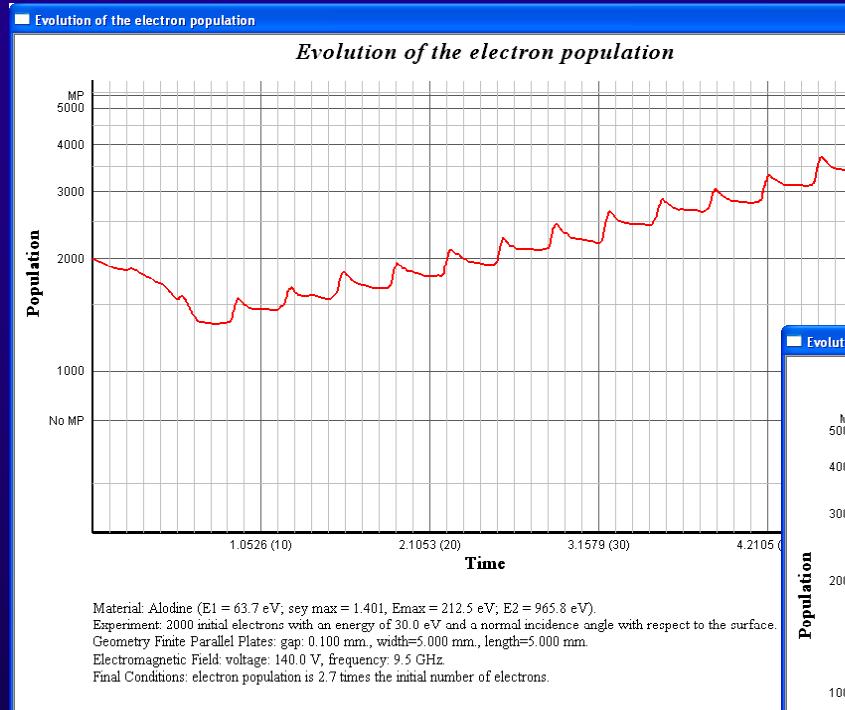
MEST results: multipactor region



MEST results: electron cloud evolution

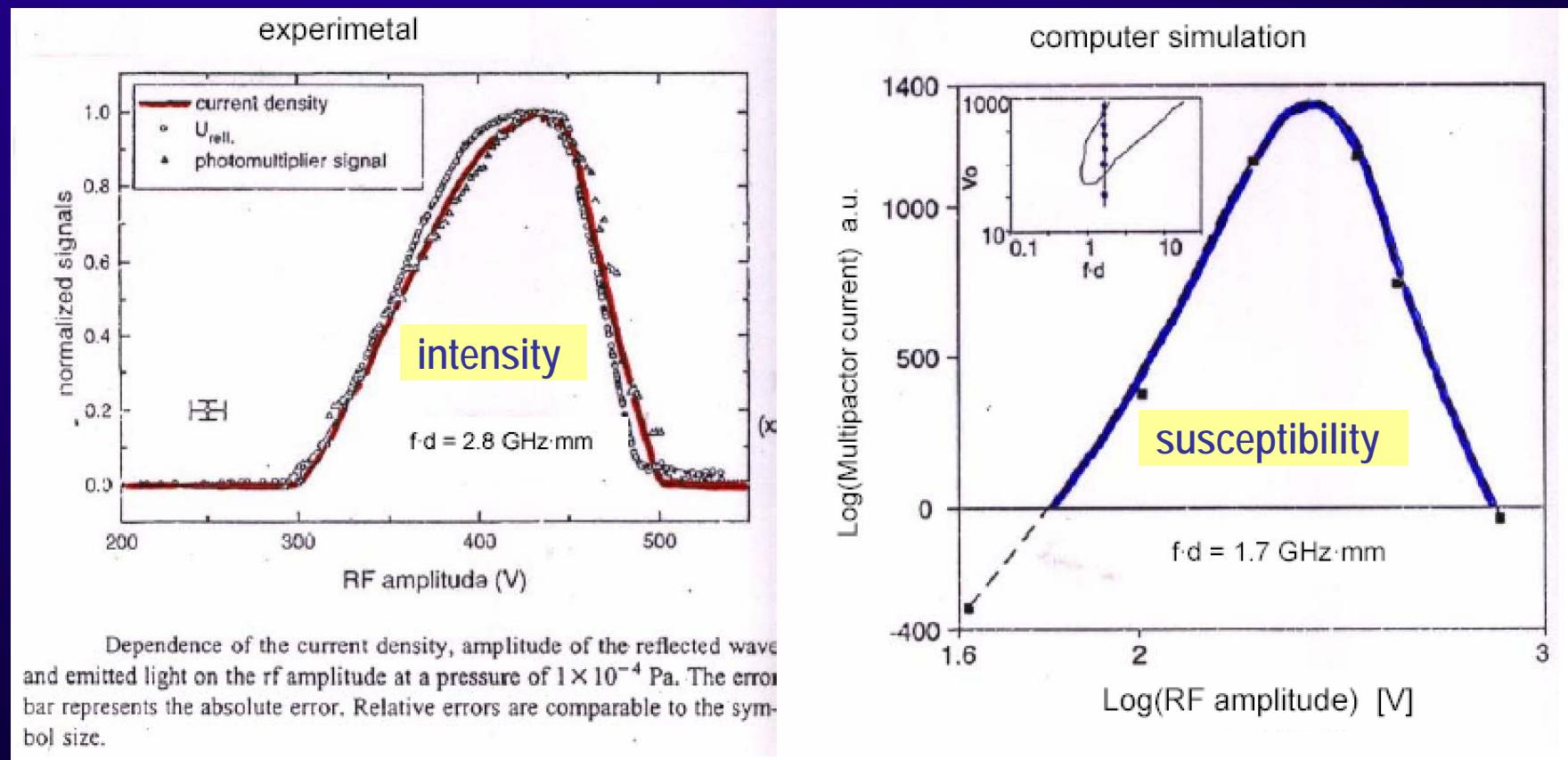


MEST results: electron population evolution



MEST results: electron cloud internal parameters

Precursor of MEST



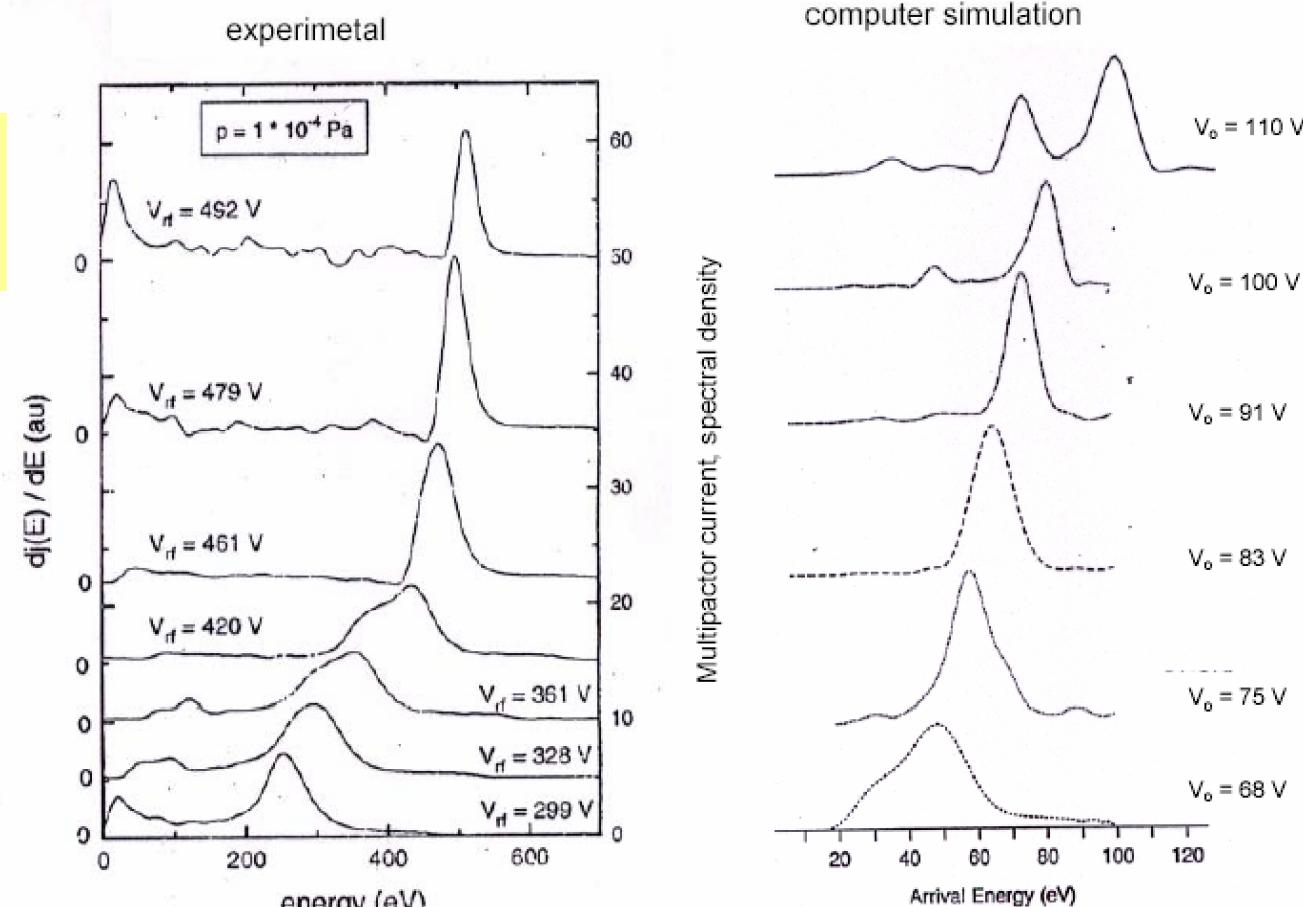
F. Höhn et al, *The transition of a multipactor to a low pressure gas discharge.*

Phys. Of Plasmas 4, (10), pp. 940947 (1997)

MEST results: electron cloud internal parameters

Precursor of MEST

Bunching of
electron
energies



Electron energy distribution as a function of the rf amplitude at pressure of 1×10^{-4} Pa.

F. Höhn

Phys. Of Plasmas 4, (10), pp. 940947 (1997)

FINAL COMMENTS

- Multipactor predictions very realistic in spite of strong simplifications in RF field
- Encouragement for using as a tool for testing SEE models and the influence of SEE parameters
- It is still a question how precise should be the SEE model for multipactor predictions
- Are experimental BSE curves important for multipactor prediction?
- Many SEE material parameters are not well known
 - which are “universal”?
 - which are only dependent on Z ?
 - which need experimental measurement?
- Is there a need for a materials SEE data base with more than SEY curves?
- Can SEE models explain abnormal SEY curves of rough surfaces?