

LHC luminosity upgrade: detector challenges

Lecture 2: Semiconductor detectors

D. Bortoletto

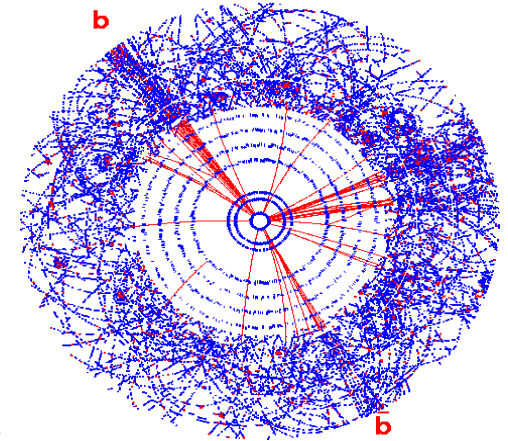
Outline

- Tracking requirements for the LHC
- Principles of operation of semiconductor detectors (4 slides)
- Radiation damage (6 slides)
- What did learn from the LHC detectors for the SLHC challenges (lower power, less mass, higher radiation tolerance and higher speed performance)
- Strategies to improve semiconductor detectors:
 - New materials
 - Defect engineering
 - New structures
 - Optimization of operation conditions
- Electronics and integration issues
- Detector specific R&D and a snapshot to the future



Tracking requirements for the LHC

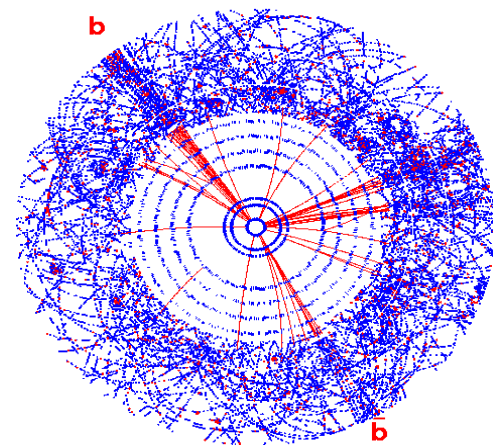
- Bunch spacing of 25ns \Rightarrow fast detector response to resolve bunch crossings
- High luminosity: $10^{33}\text{-}10^{34}\text{ cm}^{-2}\text{ s}^{-1}$ up to 20 minimum bias/bunch crossing
 - high detector granularity to keep occupancy low and resolve nearby tracks
- Excellent momentum resolution for low and high p_T tracks
- High track reconstruction efficiency
- Ability to tag b-jets and identify B-hadrons and τ 's
 - Excellent impact parameter resolution
- Unprecedented irradiation level \Rightarrow radiation hardness
 - Operating $T \approx -10^\circ\text{ C}$ to minimize radiation damage
 - Dose at 4 cm after 10 years: $500\text{ fb}^{-1} \sim 3 \times 10^{15}\text{ cm}^{-2}$
 - Dose at 22 cm after 10 years: $500\text{ fb}^{-1} \sim 1.5 \times 10^{14}$
- Small amount of material in front of electromagnetic calorimeter
- Risk of failure (preference for known industrial technologies) and cost



$H \rightarrow b\bar{b}$ LHC high lumi

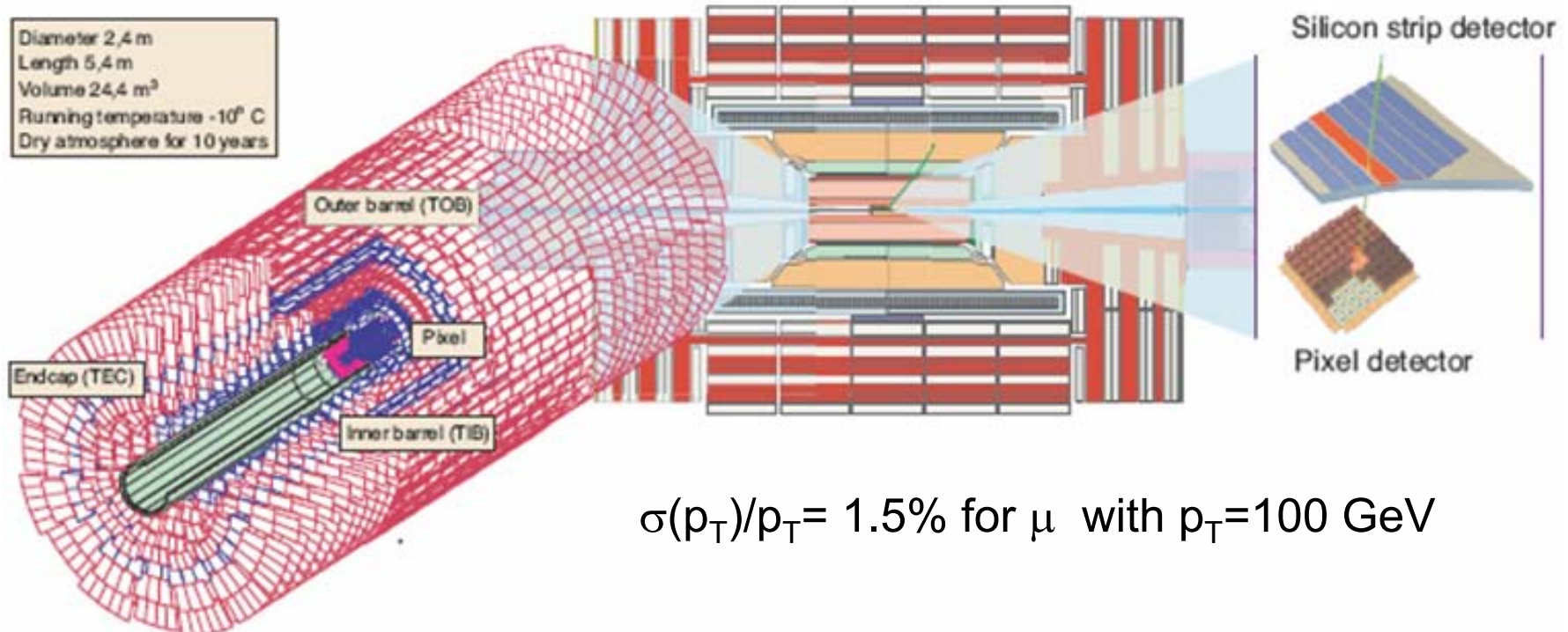
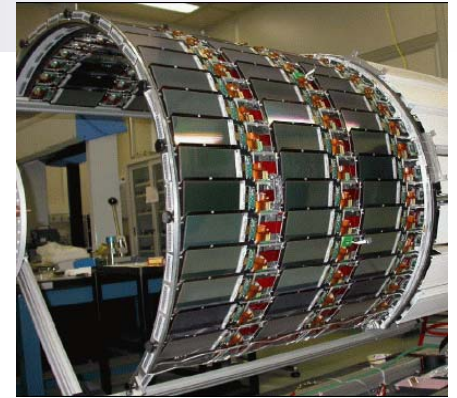
Tracking requirements for the SLHC

- Bunch spacing of 25ns (12.5 ns?) \Rightarrow fast detector response to resolve bunch crossings
- High luminosity: 10^{33} - 10^{34} $\text{cm}^{-2} \text{s}^{-1}$ (10^{35} $\text{cm}^{-2} \text{s}^{-1}$) up to 20 minimum bias/bunch crossing (100 minimum bias/bunch crossings)
 - high detector granularity to keep occupancy low and resolve nearby tracks
- Excellent momentum resolution for low and high p_T tracks
- High track reconstruction efficiency
- Ability to tag b-jets and identify B-hadrons and τ 's
 - Excellent impact parameter resolution
- Unprecedented irradiation level \Rightarrow radiation hardness
 - Operating $T \approx -10^\circ \text{C}$ to minimize radiation damage
 - Dose at 4 cm after 10 years: $500 \text{ fb}^{-1} \sim 3 \times 10^{15} \text{ cm}^{-2}$ ($3000 \text{ fb}^{-1} \sim 1.8 \times 10^{16} \text{ cm}^{-2}$)
 - Dose at 22 cm after 10 years: $500 \text{ fb}^{-1} \sim 1.5 \times 10^{14} \text{ cm}^{-2}$ ($3000 \text{ fb}^{-1} \sim 9 \times 10^{15} \text{ cm}^{-2}$)
- Small amount of material in front of electromagnetic calorimeter
- Risk of failure (preference for known industrial technologies), cost



CMS tracker

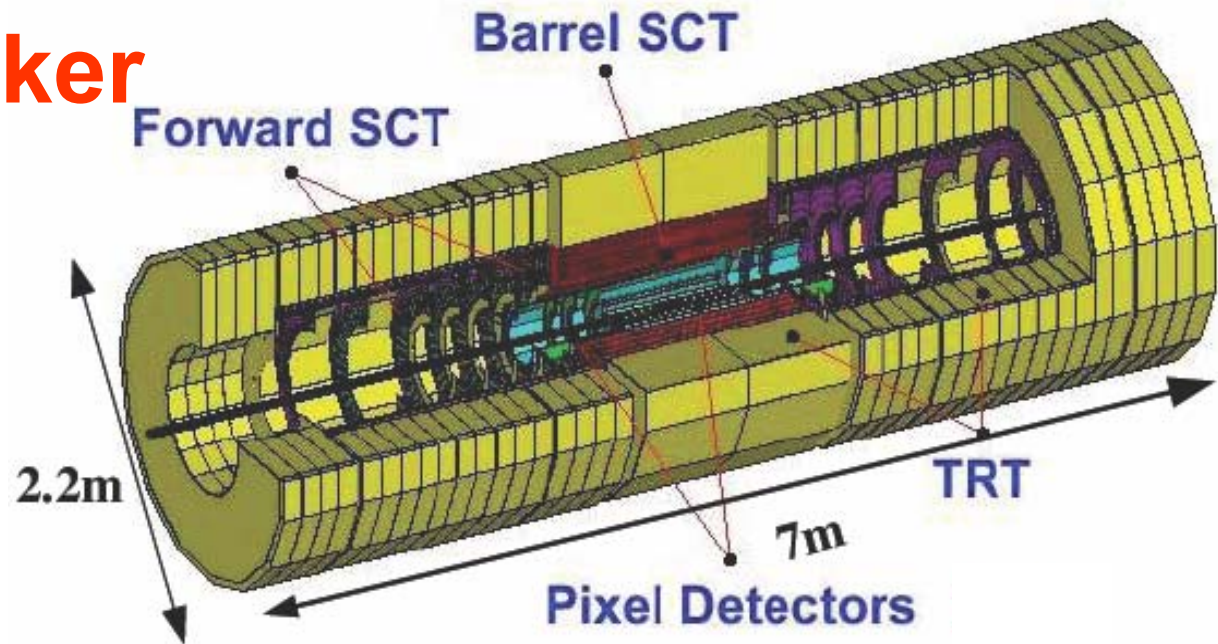
- About 200m² of active silicon area
- 1440 pixel modules with 66 million pixels
 - Pixel size: 100μm (r-φ) x 150μm (r-z),
 - Charge sharing due to large Lorentz angle (23°) + analog readout
 - spatial resolution ~10μm in r-φ, ~20μm in r-z
- 15148 silicon strip modules ~10 million strips



ATLAS tracker

■ Pixel:

- 3 barrel layers, 2 x 3 disks \Rightarrow three 3d space points for 98% of tracks
- Spatial resolution $\sim 12\mu\text{m}$ in $r-\phi$, $\sim 60\mu\text{m}$ in $r-z$



■ Semiconductor tracker (SCT):

- 4 barrel layers, 2 x 9 disks; 4088 modules, 61m²
- All modules are double sided (2.3°)

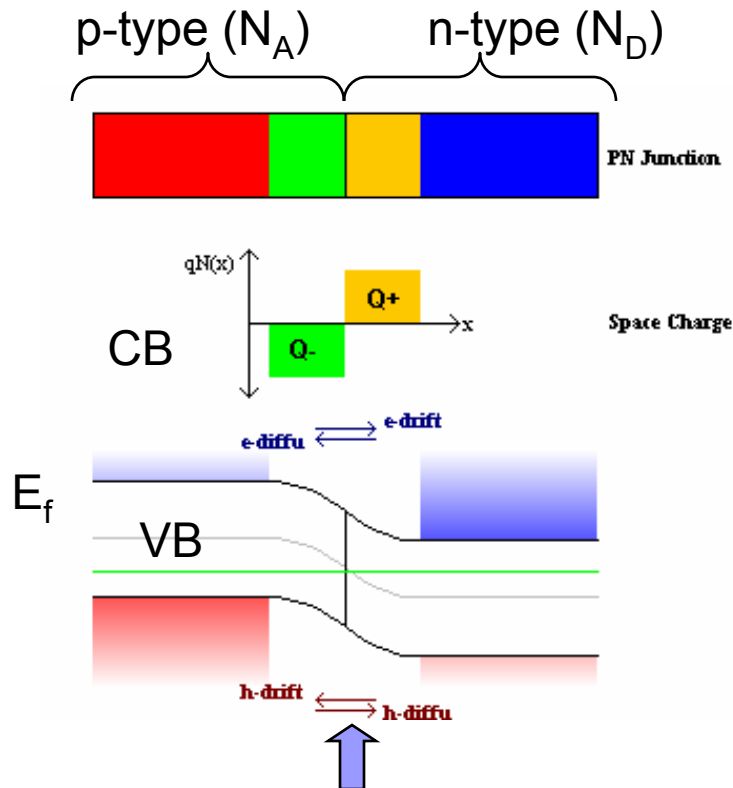
$$B=2T$$

$$\sigma(p_T)/p_T \sim 2 [\sigma(p_T)/p_T]_{\text{CMS}}$$

■ Transition radiation tracker (TRT):

- 370 000 drift tubes; spatial res. from drift time: $170\mu\text{m}$ per straw
- Continuous tracking (> 30 hits per track), low cost, less material per point
- Electron/pion separation
- Concerns: occupancy, speed (maximal drift time: 40ns)

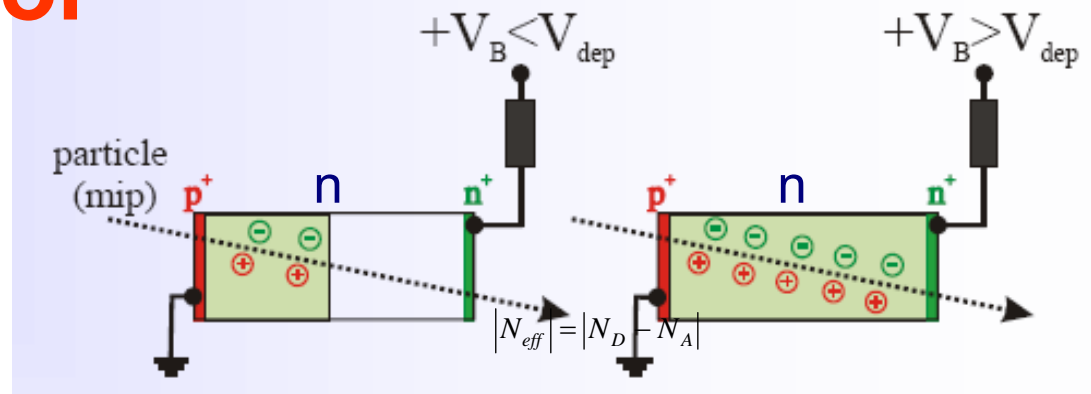
A silicon detector



Electron and holes combine

Solve Poisson eq.

$$-\frac{d^2}{dx^2} \phi(x) = \frac{q_0}{\epsilon\epsilon_0} N_{eff}$$



- Reversed biased p-n junction to establish region with no mobile carriers

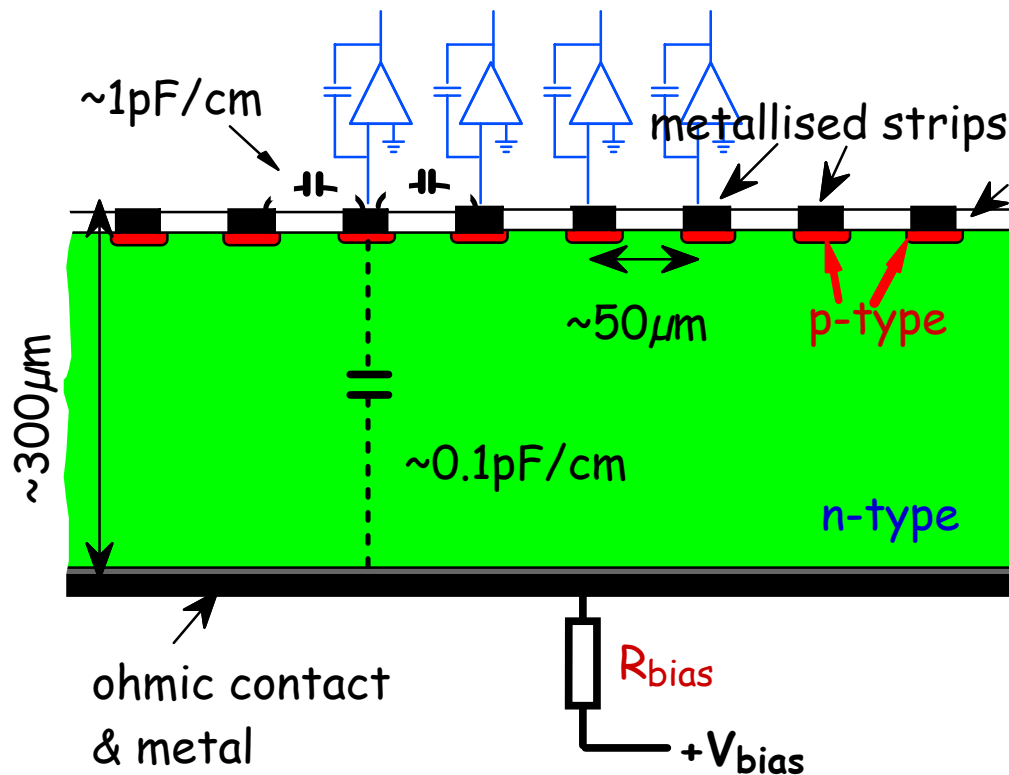
$$V_{dep} = \frac{q_0}{\epsilon\epsilon_0} |N_{eff}| d^2 \quad |N_{eff}| = |N_D - N_A|$$

- Increase external reverse bias
 - Increase E field \Rightarrow e^- and h drift to electrodes
 - Increase depletion region size
 - Reduce capacitance $\approx \epsilon\epsilon_0 A/d$ (Measurement of C yields full depletion voltage)
 - Small current flow

Single sided detectors



- Make several p-n junctions by segmenting the p+ layer into strips
- Connect the strips to individual readout channels

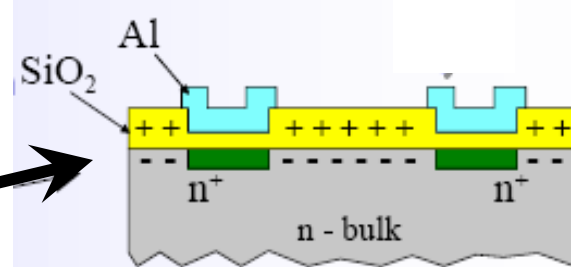
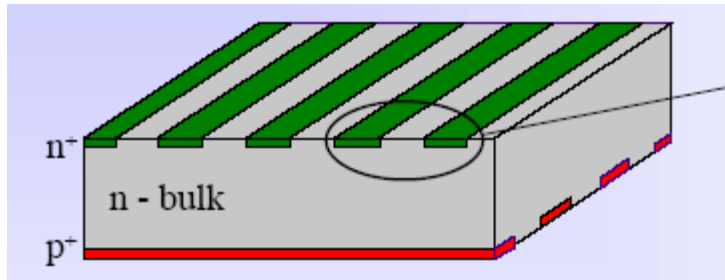


- **Detector size**
 - limited by wafer size < 15cm diameter
- **Signal speed**
 - $\langle E \rangle \approx 150\text{V}/300\mu\text{m}$
 - p-type strips collect holes $v_{\text{hole}} \approx 15 \mu\text{m}/\text{ns}$
- **Connect amplifier to each strip**
 - can also use inter-strip capacitance \Rightarrow reduce number of amplifiers to share charge over strips
- **Spatial measurement precision**
 - defined by strip dimensions and readout method $\sigma = p/\sqrt{12}$
 - ultimately limited by charge diffusion $\sigma \sim 5\text{-}10\mu\text{m}$

300 μm n-type silicon,
 $\rho = 2\text{K}\Omega\cdot\text{cm}$ ($N_D \approx 2.2 \times 10^{12} \text{ cm}^{-2}$) $\Rightarrow V_d = 150\text{V}$
 $p = 50 \mu\text{m} \Rightarrow \sigma \sim 14.4 \mu\text{m}$

Double sided detectors

- Segment both n and p-side \Rightarrow 2D

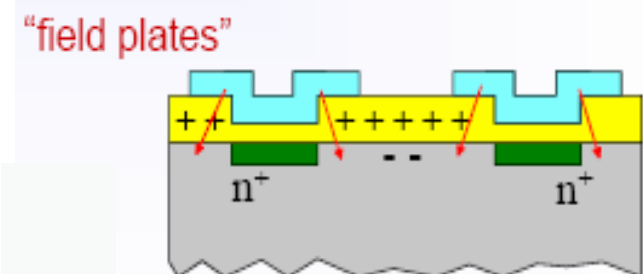
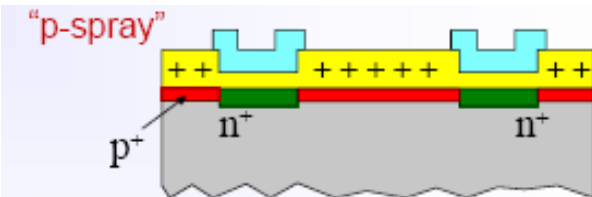
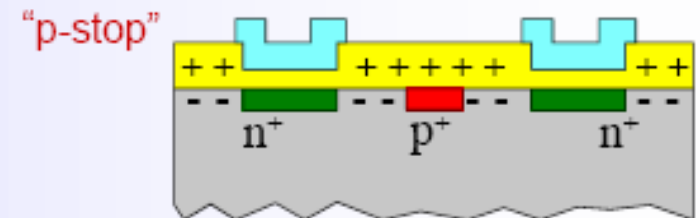


- Problem electron accumulation layer
 - n^+ strips are not isolated because of an accumulation layer at Si-SiO₂ interface due to positive charges in the SiO₂ layer

- Solution:

- p -strips between n -strips (**p-stops**)
- Moderate p -implantation over the all surface (**p-spray**)
- Metal at negative potential over n^+ strips to repel electrons (**field plates**)

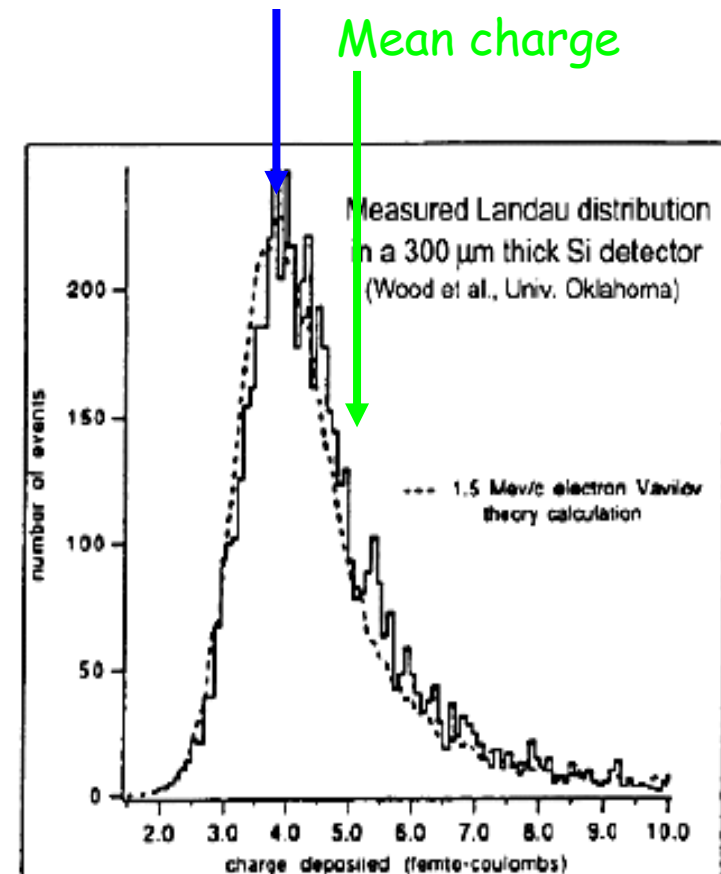
- These isolation techniques are also used in n^+ on n detectors . Advantage $v = \mu E$:
 $\mu_e = 1350 \text{ cm}^2 / \text{V}\cdot\text{s} \gg \mu_h = 450 \text{ cm}^2 / \text{V}\cdot\text{s}$



S/N before irradiation

- **Collected charge usually given for Minimum Ionizing Particle (MIP):**
 - Mean $dE/dx_{Si} = 3.88 \text{ MeV/cm} \Rightarrow 116 \text{ keV}$ for $300 \mu\text{m}$ thick Si
 - Most probable loss = 81 keV for $300 \mu\text{m}$ Si
 - Since 3.6eV needed to make e-h pair \Rightarrow charge in $300 \mu\text{m} = 22500 e^-$ ($=3.6 \text{ fC}$)
- **Landau distribution has a low energy tail which broadens because of noise.**
 - **Noise sources:**
 - Capacitance $\text{ENC} \propto C_d$
 - Leakage Current $\text{ENC} \propto \sqrt{I}$
 - Thermal Noise $\text{ENC} \propto \sqrt{(kT/R)}$
- **Typical Values S/N > 10-15**

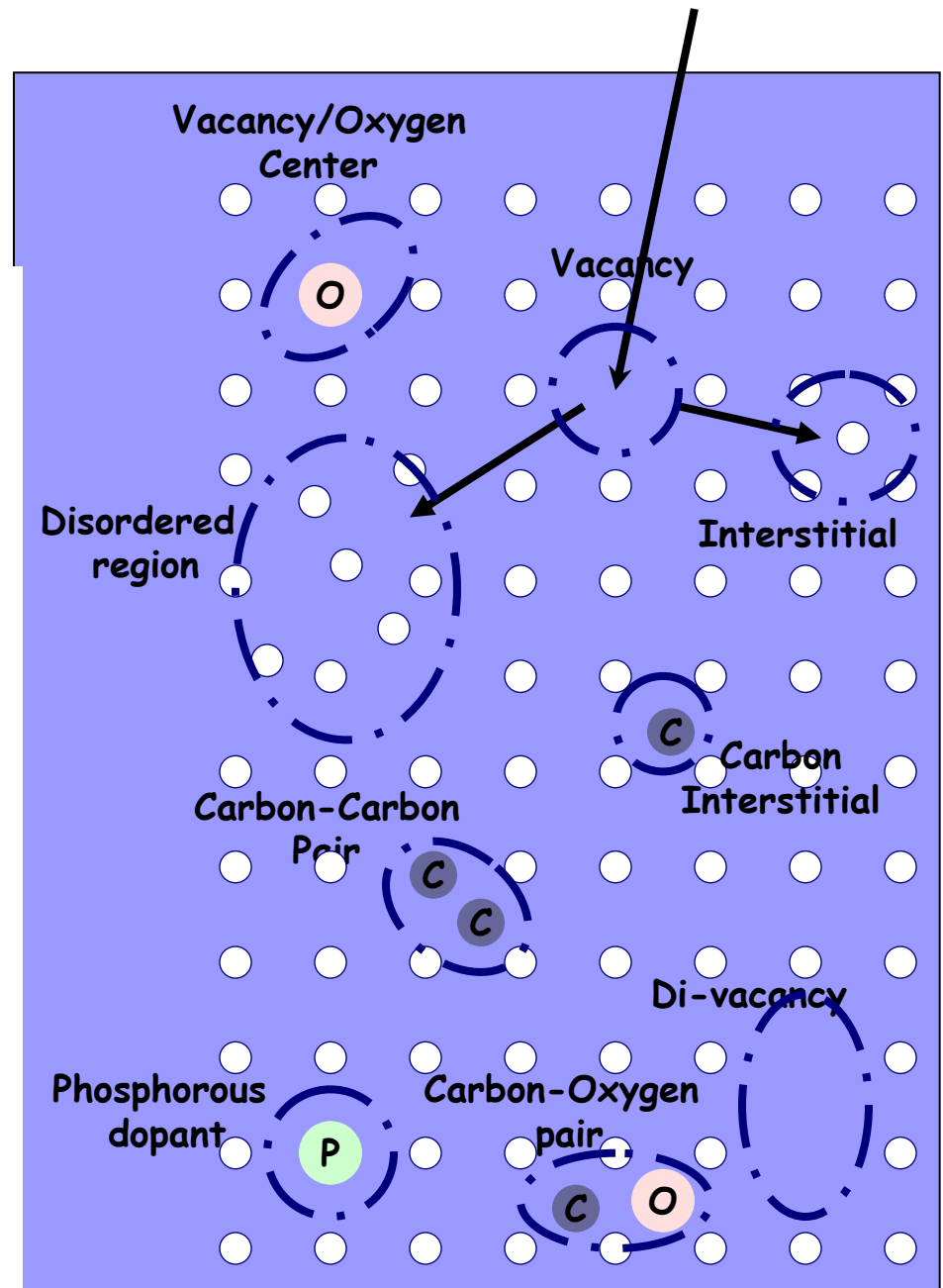
Most probable charge $\approx 0.7 \times$ mean



Bulk Damage- Microscopic view



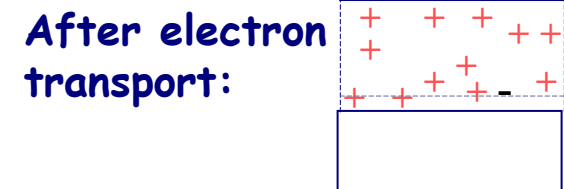
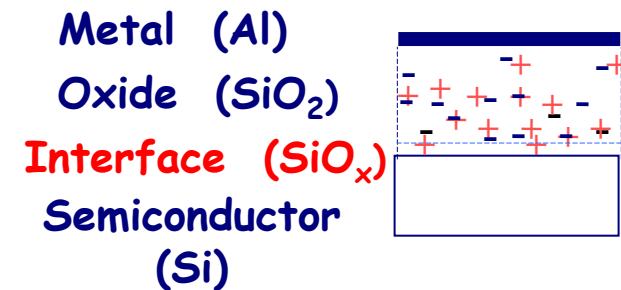
- Bulk damage is mainly from **hadrons displacing primary lattice atoms**:
 - Vacancies, silicon interstitials, and large disordered regions
 - 1 MeV neutron transfers $\approx 60-70$ keV to recoiling silicon atom, which in turn can displace clusters
- Defects can recombine or migrate through the lattice to form more complex and stable defects
 - Annealing can be beneficial
 - Defects can be stable or unstable
- Defects add levels in the band gap affecting macroscopic properties: N_{eff} , trapping and leakage current



Surface Damage



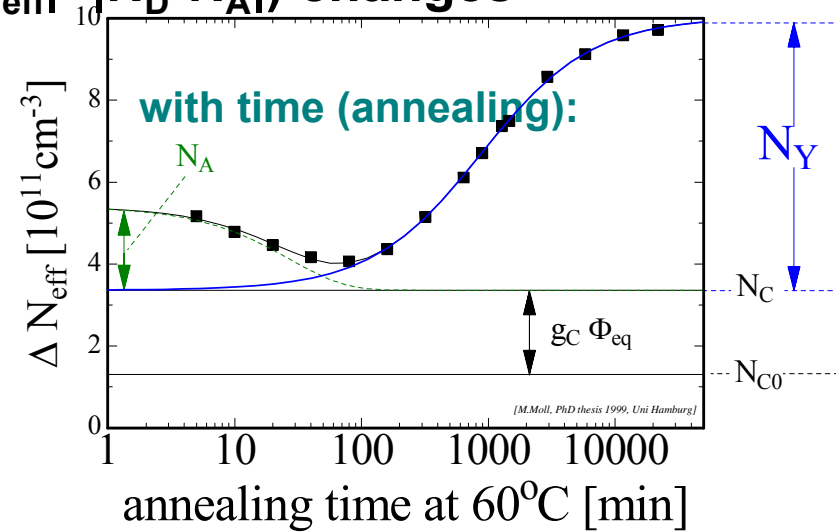
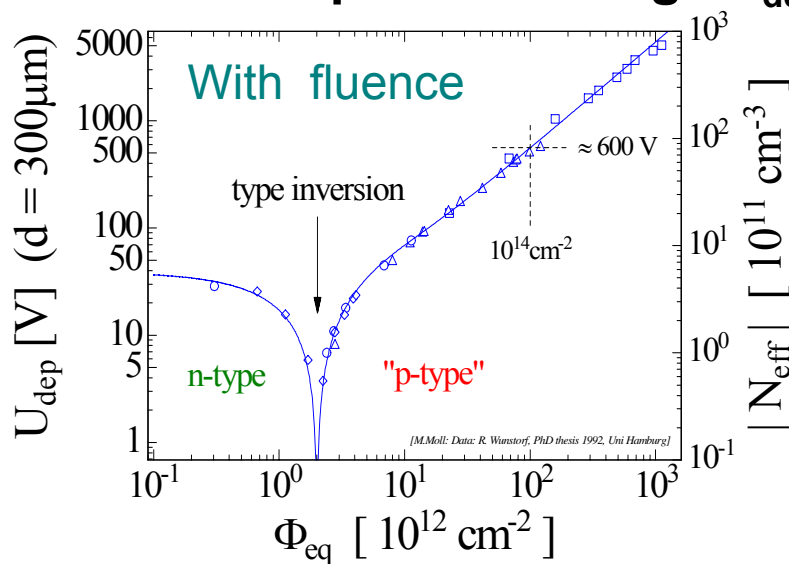
- Surface damage generation:
 - Ionizing radiation creates e-h pairs in SiO_2
 - Many recombine, electrons migrate quickly, holes slowly migrate to Si/SiO_2 interface since $\mu_{\text{hole}} \ll \mu_{\text{electron}}$
 - Some holes 'stick' in the boundary layer
- Surface damage results in
 - Increased interface trapped charge
 - Increased fixed oxide charges
 - Surface generation centers
- Macroscopic effects
 - Increase in the sensors capacitance
 - Decrease in the interstrip resistance
 - Surface current
 - Risk to readout electronics
 - threshold shifts
 - noise and gain deterioration



Effective doping concentration



■ The depletion Voltage V_{dep} ($|N_{eff}| = |N_D - N_A|$) changes



Most irradiation damages in the Si behave like acceptors ⇒
“Hamburg” model

$$|N_{eff}(\Phi_{eq}, t, T)| = |N_D - N_A| =$$

$N_a(\Phi_{eq}, t, T)$	Annealing
$+ N_c(\Phi_{eq})$	Constant
$+ N_y(\Phi_{eq}, t, T)$	Anti-annealing

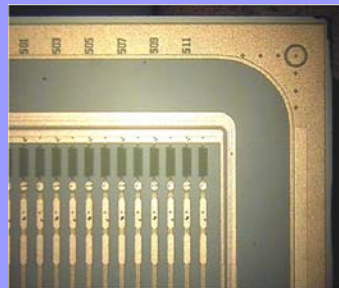
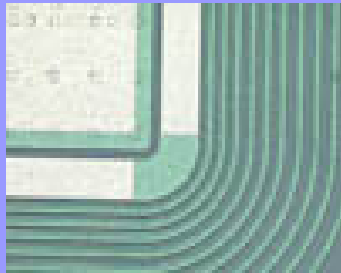
- Short term: “Beneficial annealing”
- Long term: “Reverse annealing”
 - time constant depends on temperature:
 - ~ 500 years (-10°C)
 - ~ 500 days (20°C)
 - ~ 21 hours (60°C)

Detectors must be cooled even when the experiment is not running!

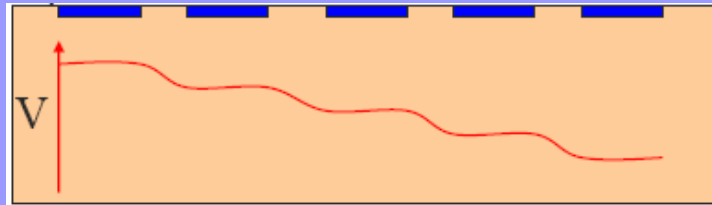
What did we learn for LHC?

Operation of devices at higher bias Voltages

- Bias ring design

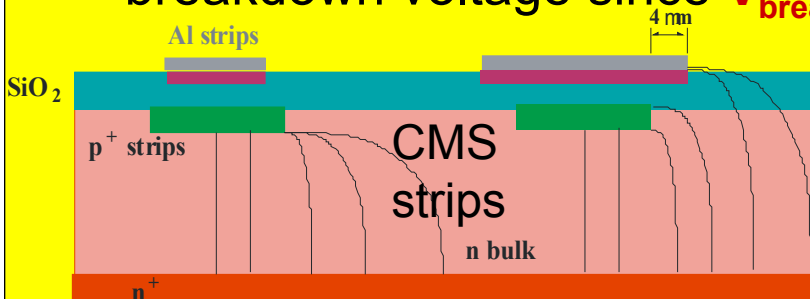


Multi-guard ring
CMS and Atlas Pixels

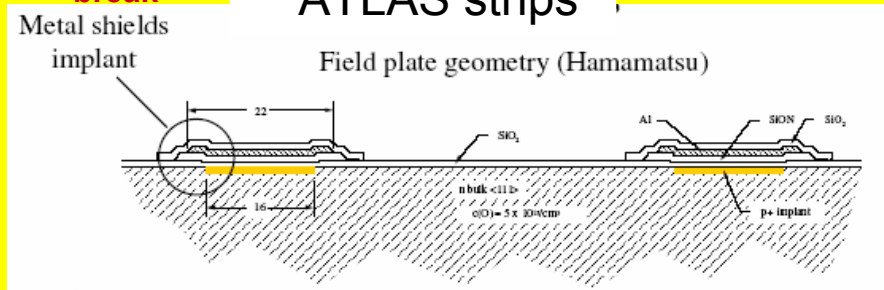


CMS strips: one guard ring, the metal overhang as “continuous” multi-guard-structure.

- **CMS:** Metal overhang design Al strips overhang implants \rightarrow high breakdown voltage since $V_{break}(SiO_2) > V_{break}(Si)$



ATLAS strips

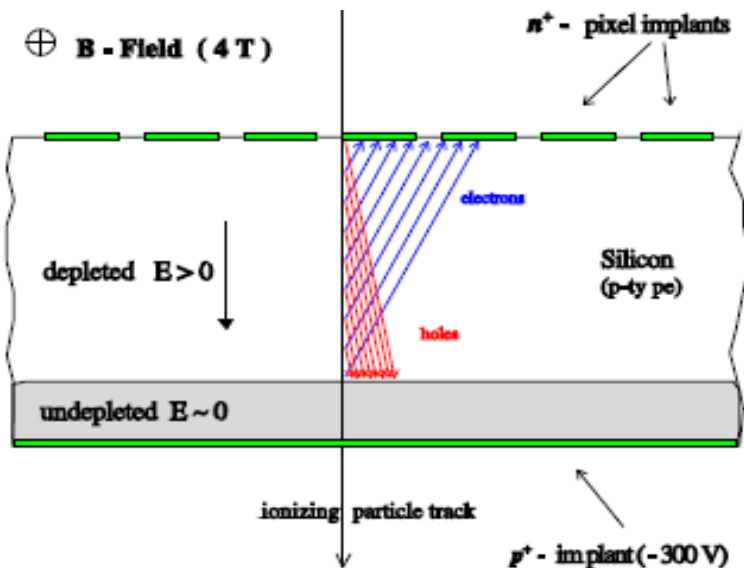
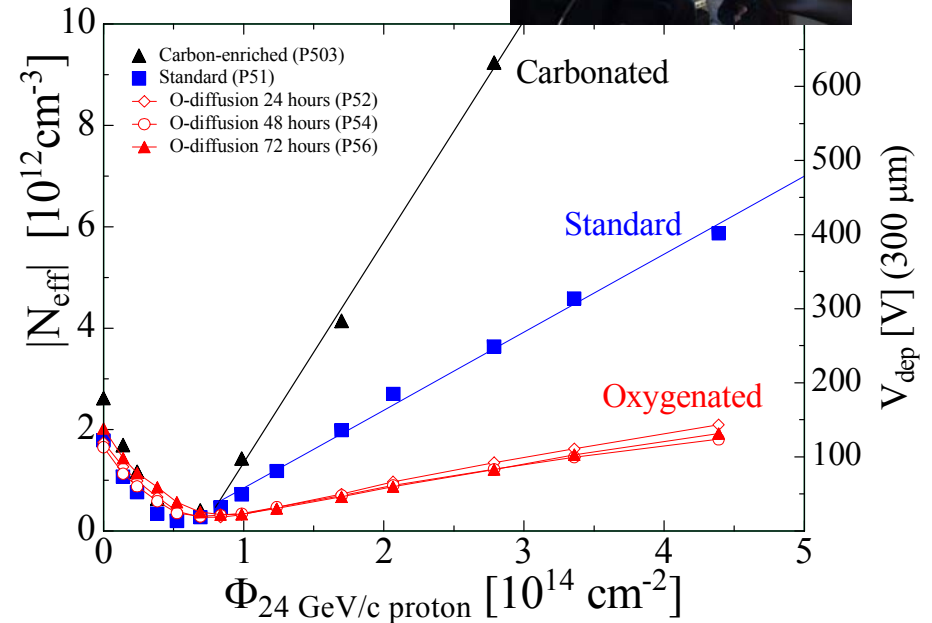


ATLAS: field plate geometry for the Hamamatsu detectors (86%), non field plate CiS (14%). CiS detectors show sensitivity to humidity

What did we learn for LHC



- Oxygen is good!!! (RD48 was formed at CERN to develop radiation hard sensors for the LHC.)
 - DOFZ silicon has less variation in V_{fd} with radiation compared to FZ \Rightarrow more *radiation hard*
 - Atlas and CMS barrel pixel detectors use oxygenated silicon

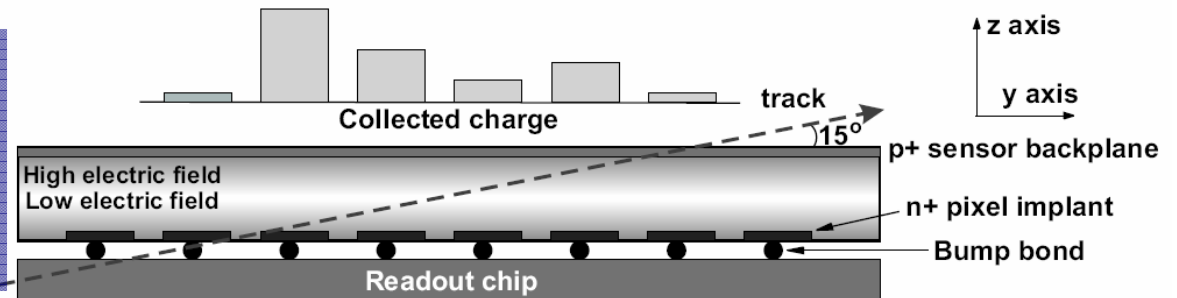


Space Charge Sign Inversion (SCSI)

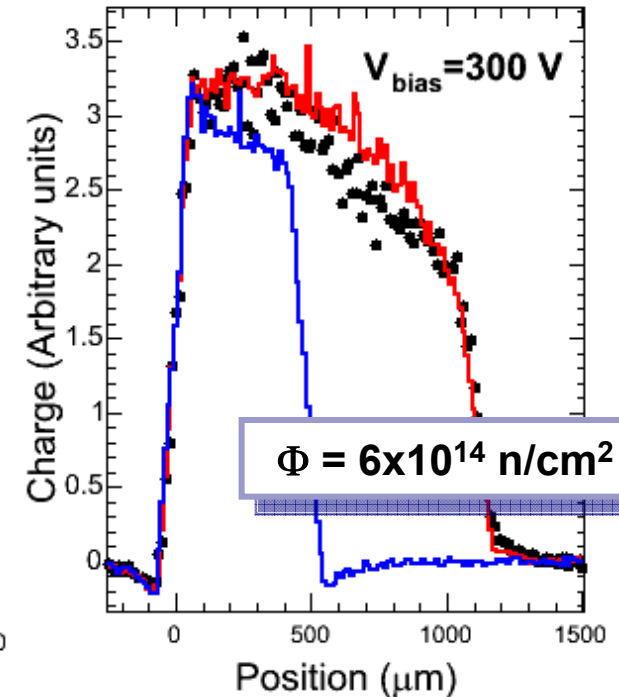
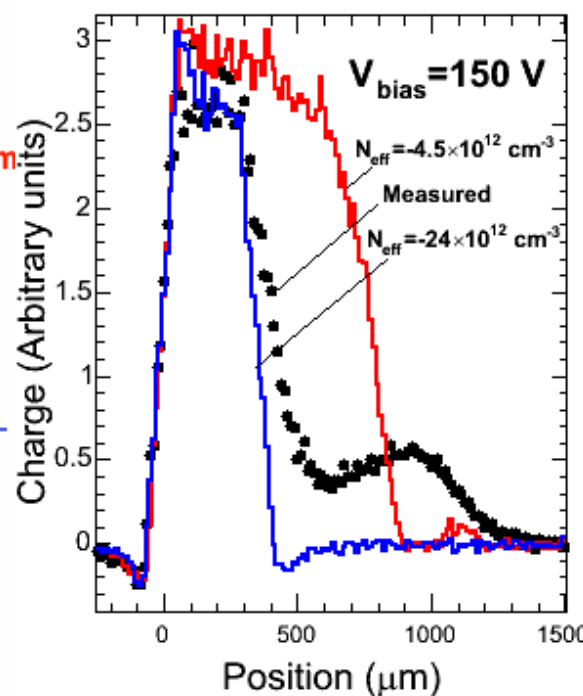
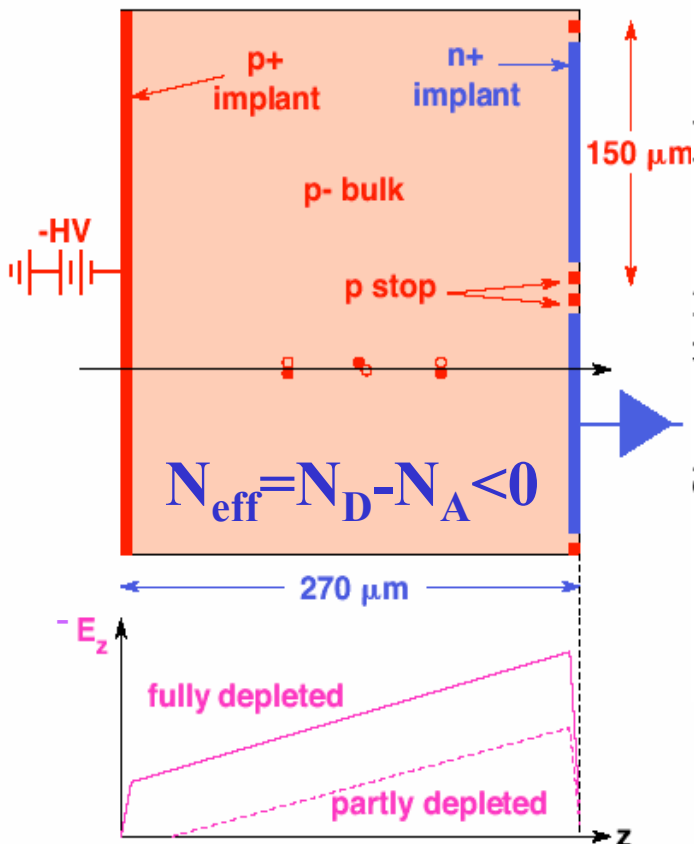
- n-on-n sensors allows operation with undepleted detectors.
- Faster charge collection
- Option chosen for Atlas and CMS pixels

Models with constant N_{eff}

Charge collection measured using cluster profiles in a row of pixels illuminated by a 15° beam and no magnetic field

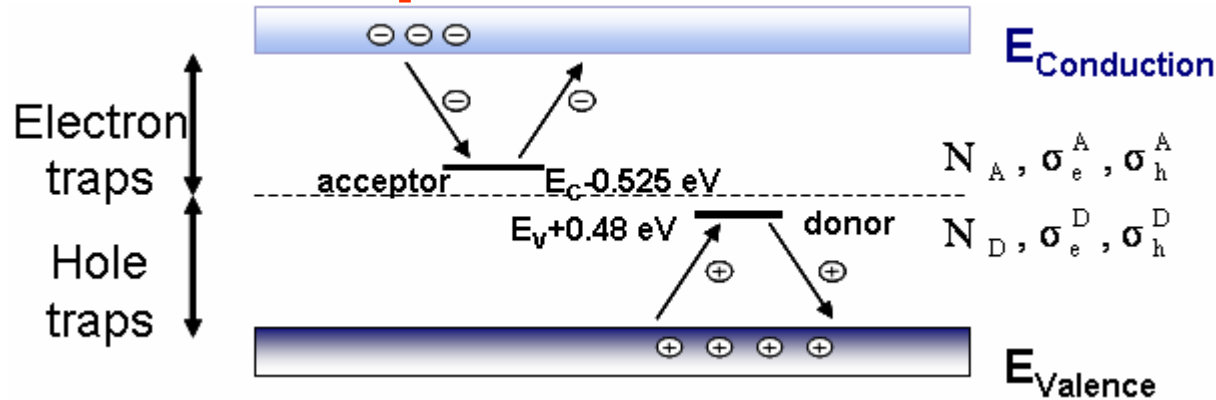


after type inversion

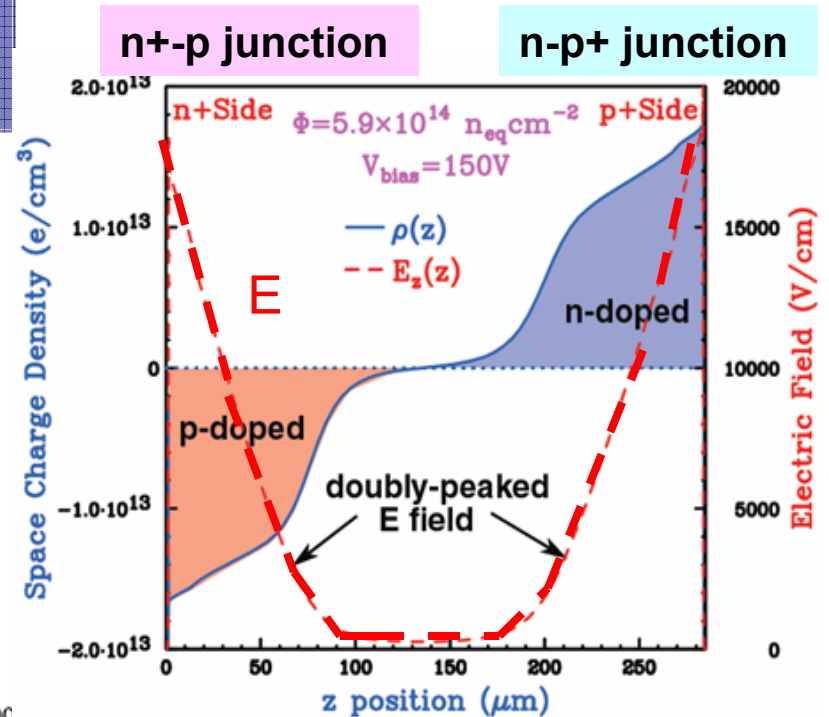
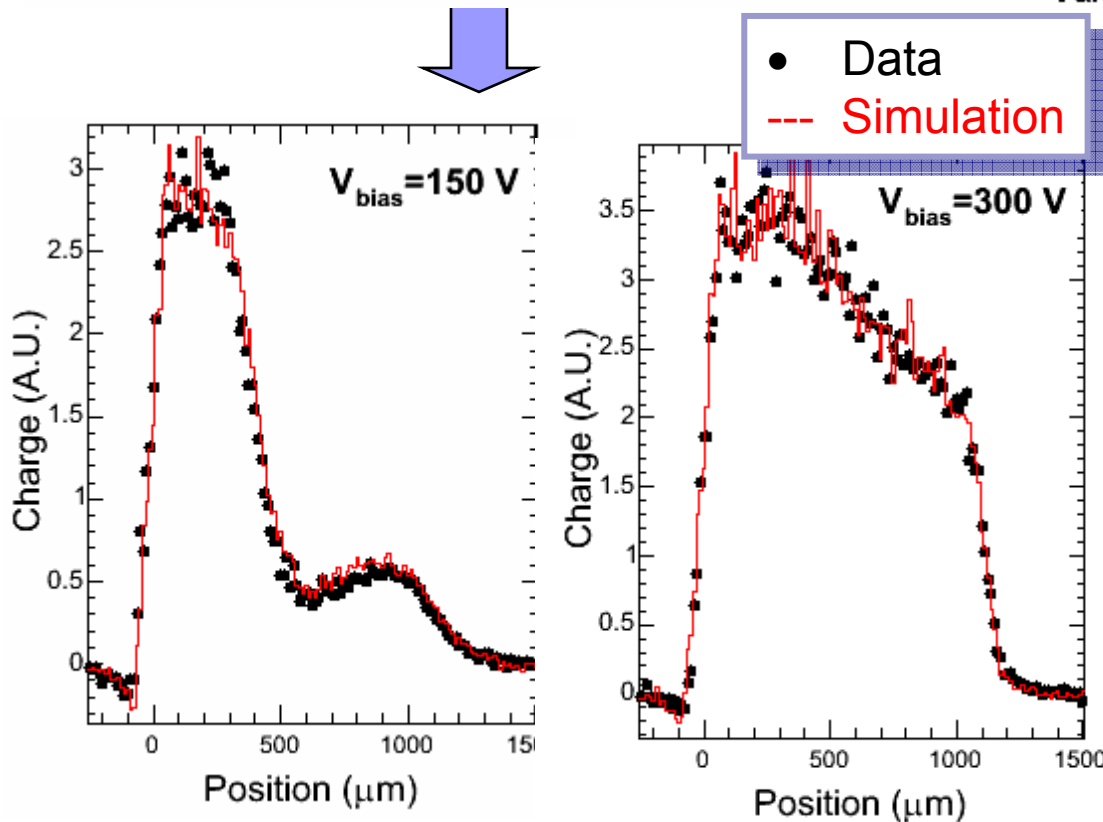


Model with **constant space charge density** does not describe the measured charge collection

Two trap model and double peak E field



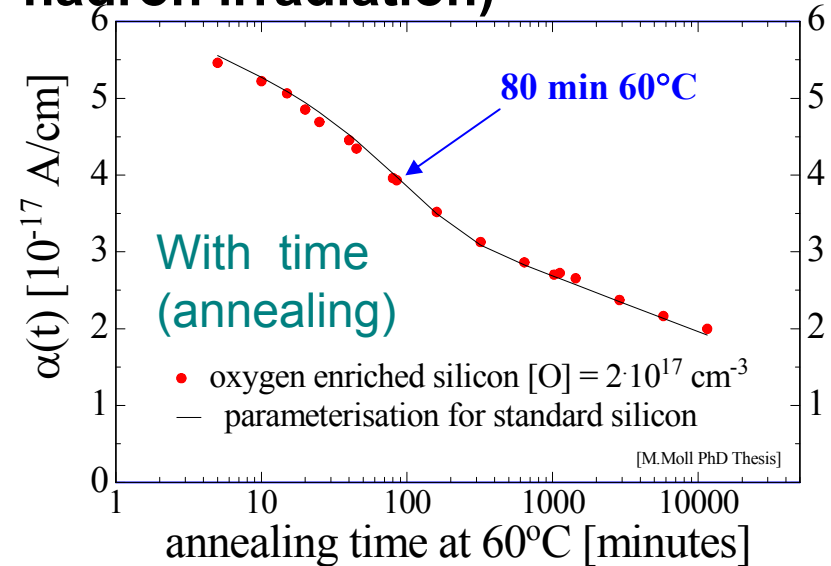
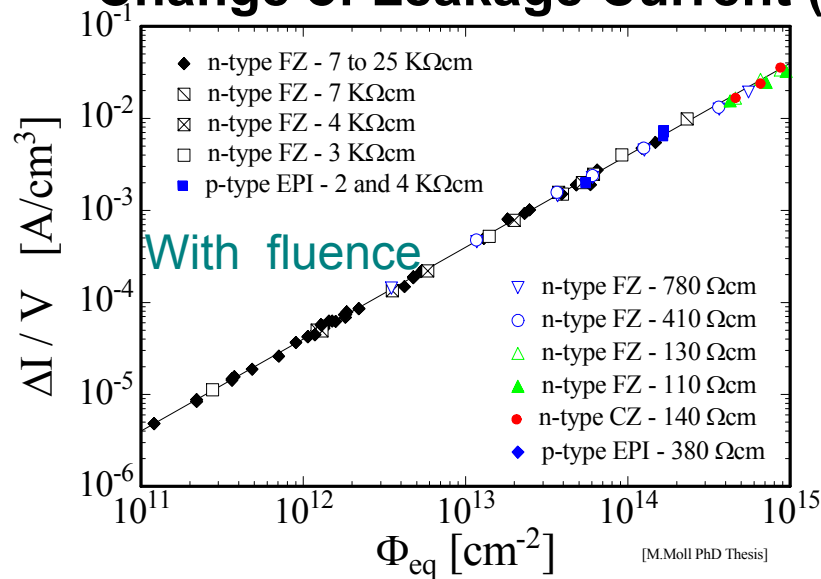
$E_{A/D}$ = trap energy level fixed
 $N_{A/D}$ = trap densities from fit
 $\sigma_{e/h}$ = trapping cross sections from fit
 $\Phi_1 = 6 \times 10^{14} \text{ n/cm}^2$,
 $N_A/N_D = 0.40, \sigma_h/\sigma_e = 0.25$



Leakage Current



Change of Leakage Current (after hadron irradiation)



- Damage parameter α (slope in figure)

$$\alpha = \frac{\Delta I}{V \cdot \Phi_{eq}}$$

Leakage current per unit volume and particle fluence

- α is constant over several orders of fluence and independent of impurity concentration in Si \Rightarrow can be used for fluence measurement

- Leakage current decreasing in time (depending on temperature)
- Strong temperature dependence

$$I \propto \exp\left(-\frac{E_s}{2k_B T}\right)$$

\Rightarrow Cool detectors during operation!
Ex: $I(-10^\circ\text{C}) \sim 1/16 I(20^\circ\text{C})$



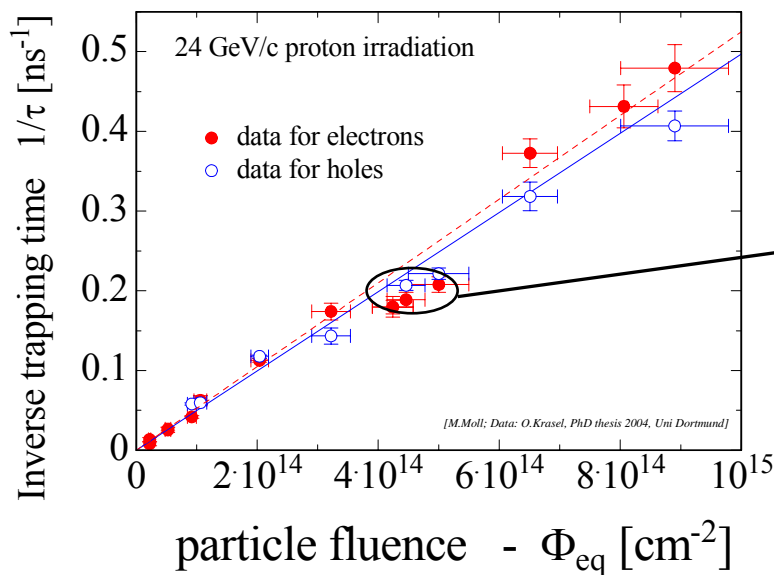
Trapping

- Deterioration of Charge Collection Efficiency (CCE) by trapping

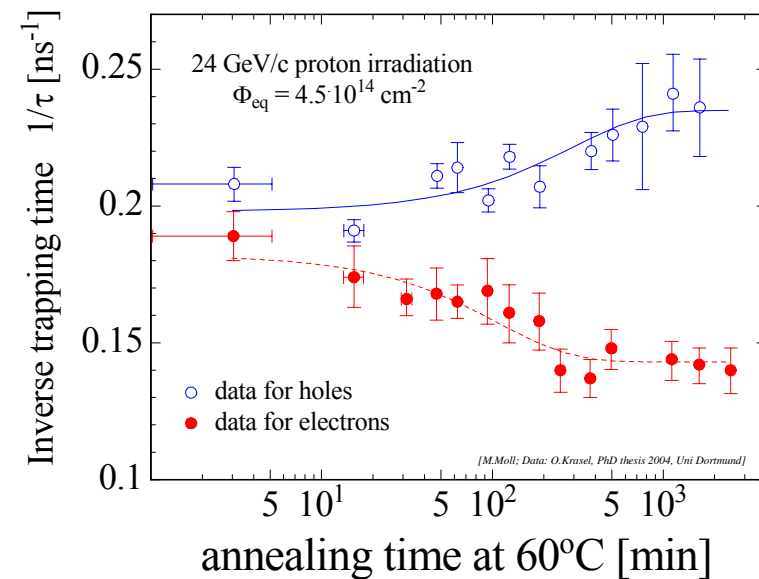
Trapping is characterized by an effective trapping time τ_{eff} for e^- and h^+ :

$$Q_{e,h}(t) = Q_{0e,h} \exp\left(-\frac{1}{\tau_{eff\ e,h}} \cdot t\right) \quad \text{where} \quad \frac{1}{\tau_{eff\ e,h}} \propto N_{defects} \propto \text{fluence}$$

Increase of $1/\tau$ with fluence



$1/\tau$ changes with annealing

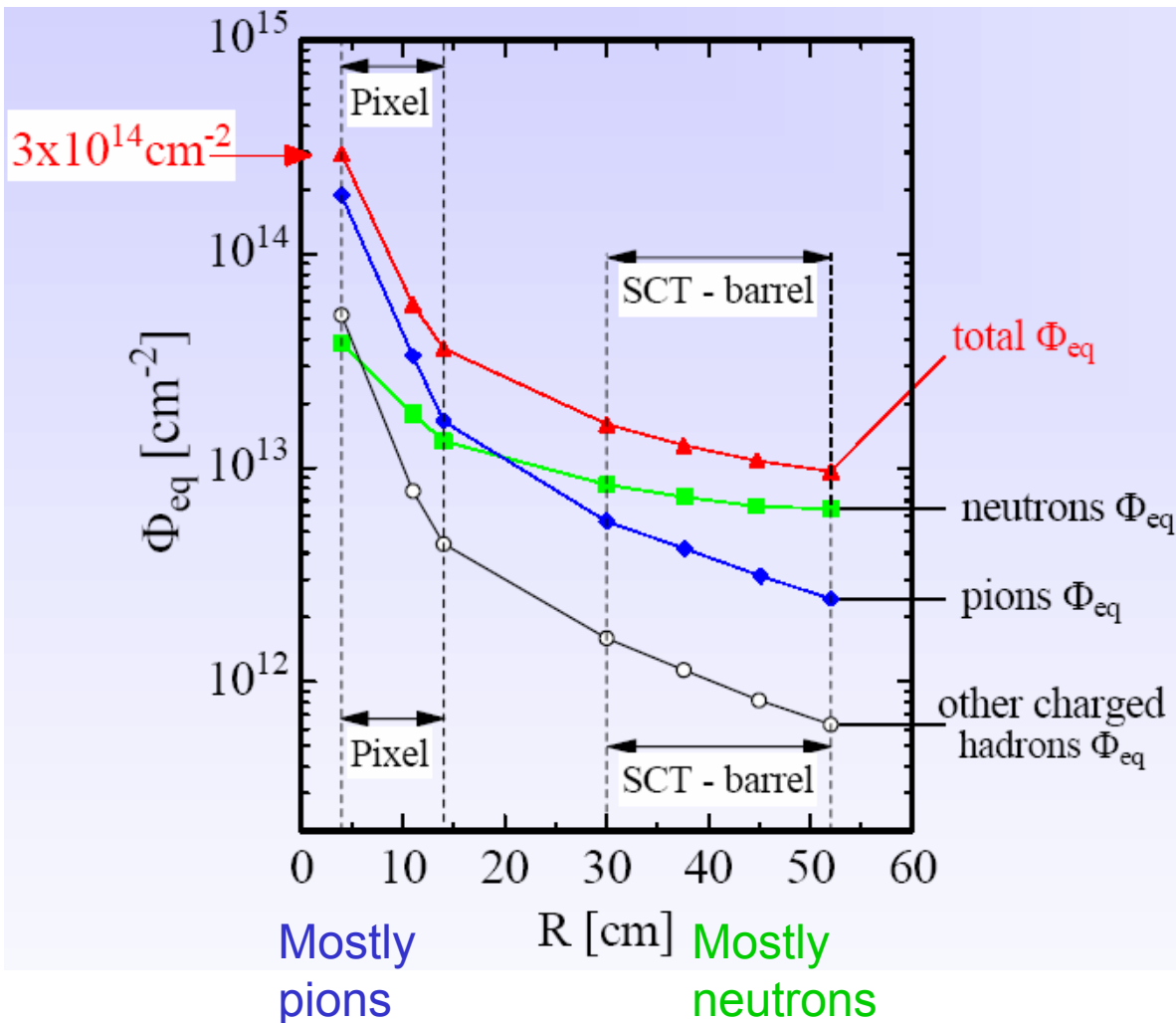


$$\tau_{eff} (10^{15}) = 2ns \quad w = v_{sat} \tau_{eff} = (10^7 \text{ cm/s}) \times 2 \times 10^{-9} \text{ s} = 200 \mu m$$

$$\tau_{eff} (10^{16}) = 0.2ns \quad w = v_{sat} \tau_{eff} = (10^7 \text{ cm/s}) \times 2 \times 10^{-10} \text{ s} = 20 \mu m$$

Huge
Decrease
in CCE?

Current LHC trackers



LHC Fluence per year at full luminosity

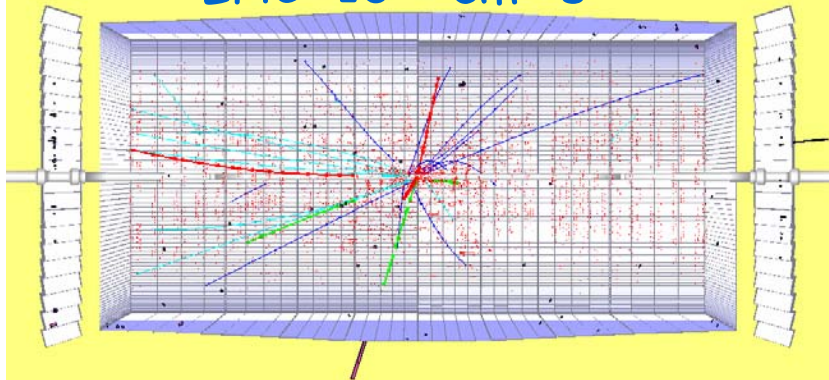
CMS

- Three regions to match radiation damage and occupancy
 - n-on-n pixels ($r < 20 \text{ cm}$)
 - $\Phi = 3.0 \times 10^{14} \text{ cm}^{-2}/\text{year}$, 270 μm thick sensors Low resistivity (1.5-2 $\text{K}\Omega\cdot\text{cm}$) oxygenated for the barrel.
 - p-on-n strips
 - Inner region ($20 \text{ cm} < r < 50 \text{ cm}$)
 - $\Phi = 1.6 \times 10^{14} \text{ cm}^{-2}$, 320 μm thick, Low resistivity (1.5-2 $\text{K}\Omega\cdot\text{cm}$), pitch $\sim 80 \mu\text{m}$
 - Outer region ($r > 50 \text{ cm}$)
 - $\Phi = 3.5 \times 10^{13} \text{ cm}^{-2}$, 500 μm thick, High resistivity (3.5-7.5 $\text{K}\Omega\cdot\text{cm}$), pitch $\sim 200 \mu\text{m}$

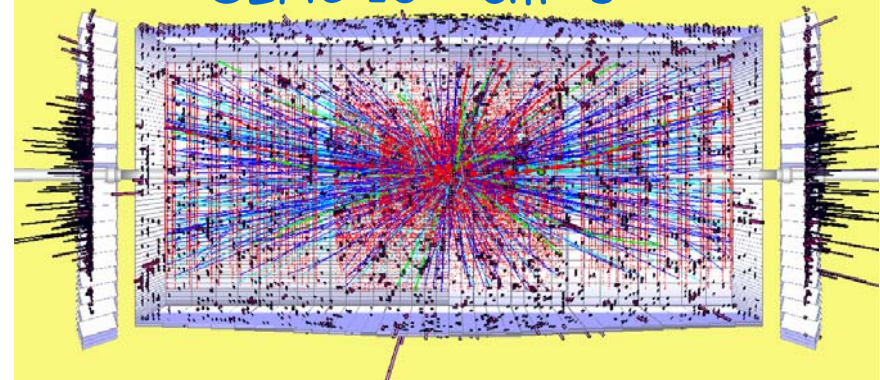
SLHC and tracking

$H \rightarrow ZZ \rightarrow ee\mu\mu$ $m(\text{higgs})=300 \text{ GeV}$ all tracks with $p_T < 1 \text{ GeV}$ removed

LHC $10^{33} \text{ cm}^{-2}\text{s}^{-1}$



SLHC $10^{35} \text{ cm}^{-2}\text{s}^{-1}$



- **Integrated Luminosity** (radiation damage) dictates the detector **technology**
- **Instantaneous rate** (particle flux) dictates the detector **granularity**
 - more granularity if we aim at same performance we expect from the LHC trackers

R (cm)	Φ (p/cm ²)	Technology	S
>50	10^{14}	Present p-in-n (or n-in-p) Limitation I_{leak}	$\sim 20\text{Ke}^-$
20-50	10^{15}	Present n-in-n (or n-in-p) Limitation V_{dep}	$\sim 10\text{Ke}^-$
<20	10^{16}	RD needed Limitation trapping	$\sim 5\text{Ke}^-$

Radiation hard devices for the SLHC (RD50 et al)

■ Silicon Defect Engineering

- Understanding radiation damage
 - Macroscopic effects and Microscopic defects
 - Simulation of defect properties & kinetics
 - Irradiation with different particles & energies
- ★ □ Oxygen rich Silicon
 - DOFZ, Cz, MCZ, EPI
- Oxy. dimer & hydrogen enriched Si
- Pre-irradiated Si
- Influence of processing technology

■ New Materials

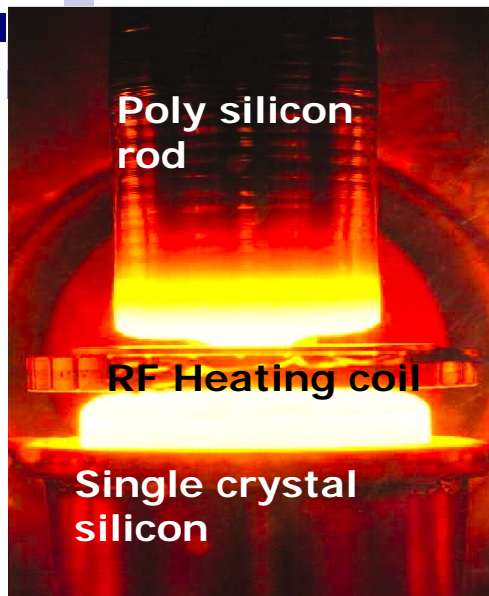
- ★ □ Silicon Carbide (SiC), Gallium Nitride (GaN)
- ★ □ Diamond: CERN RD42 Collaboration
- ★ □ Amorphous silicon (TFA)

• Device Engineering

- p-type silicon detectors (n-in-p) ★
- thin detectors
- 3D ★
- Semi 3D detectors Stripixels ★
- Cost effective detectors
- Simulation of highly irradiated detectors ★
- Monolithic devices

• Change operational conditions

- CERN-RD39
“Cryogenic Tracking Detectors” ★

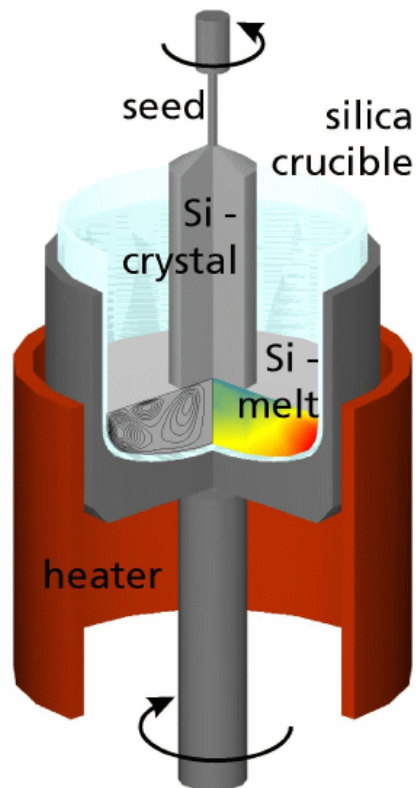


■ Float Zone

- Using a single Si crystal seed, melt the vertically oriented rod onto the seed using RF power and “pull” the **monocrystalline ingot**
- Can be oxygenated by diffusion at high T

■ Czochralski silicon

- Pull Si-crystal from a Si-melt contained in a silica crucible while rotating.
- Silica crucible is dissolving oxygen into the melt ⇒ **high concentration of O in CZ**
- Material used by IC industry (cheap), now available in high purity for use as particle detector (MCz)

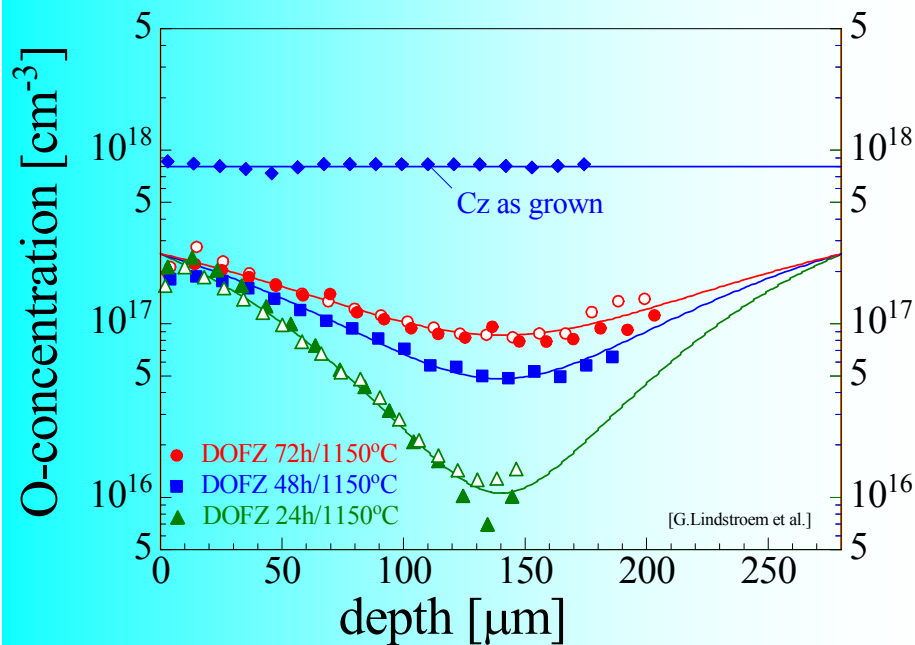


■ Epitaxial silicon

- Chemical-Vapor Deposition (CVD) of Silicon
- CZ silicon substrate used ⇒ **diffusion of oxygen**
- Growth rate about $1\mu\text{m}/\text{min}$
- Excellent homogeneity of resistivity
- $150\mu\text{m}$ thick layers produced (thicker is possible)
- price depending on thickness of epi-layer but not extending ~ 3 x price of FZ wafer

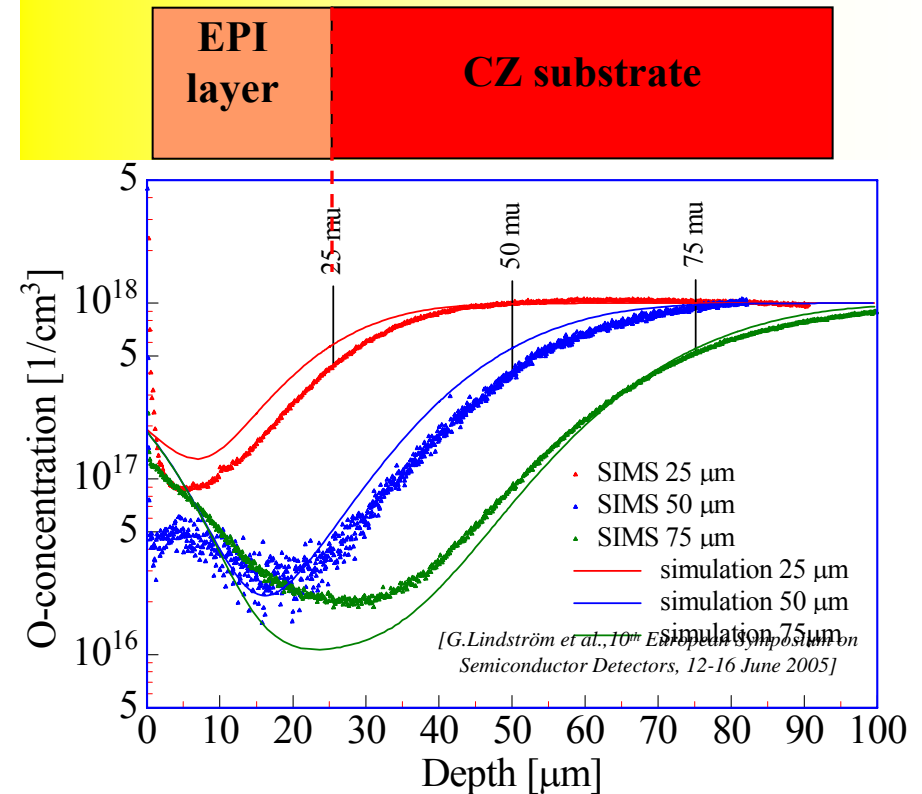
Oxygen concentration in FZ, CZ and EPI

- CZ: high homogeneous concentration and formation of Thermal Donors (reducing acceptors due to radiation)



- DOFZ: inhomogeneous oxygen distribution, increasing with time at high temperature

- Epitaxial silicon



- EPI: inhomogeneous O concentration due to diffusion from substrate into epi-layer during production

Standard FZ, DOFZ, Cz and MCz Silicon

■ Standard FZ silicon

- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- strong N_{eff} increase at high fluence

■ Oxygenated FZ (DOFZ)

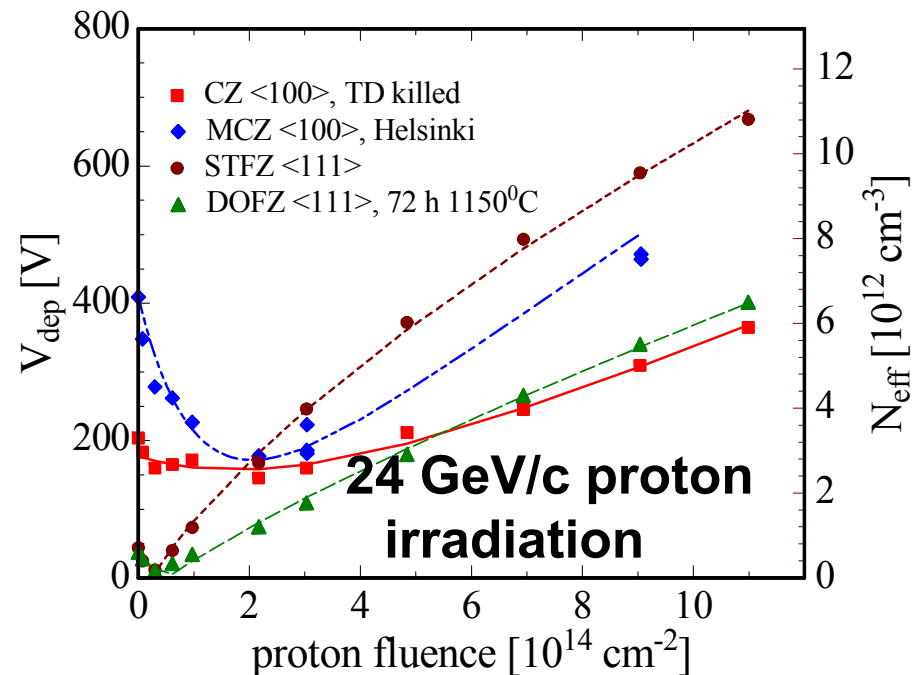
- type inversion at $\sim 2 \times 10^{13}$ p/cm²
- reduced N_{eff} increase at high ϕ

■ CZ silicon and MCZ silicon

- no type inversion in fluence range
- **Verified for CZ and MCz silicon by TCT measurements** \Rightarrow
donor generation overcompensates acceptor generation

■ Common to all materials (after hadron irradiation):

- **reverse current increase**
- **increase of trapping (electrons and holes) within $\sim 20\%$**

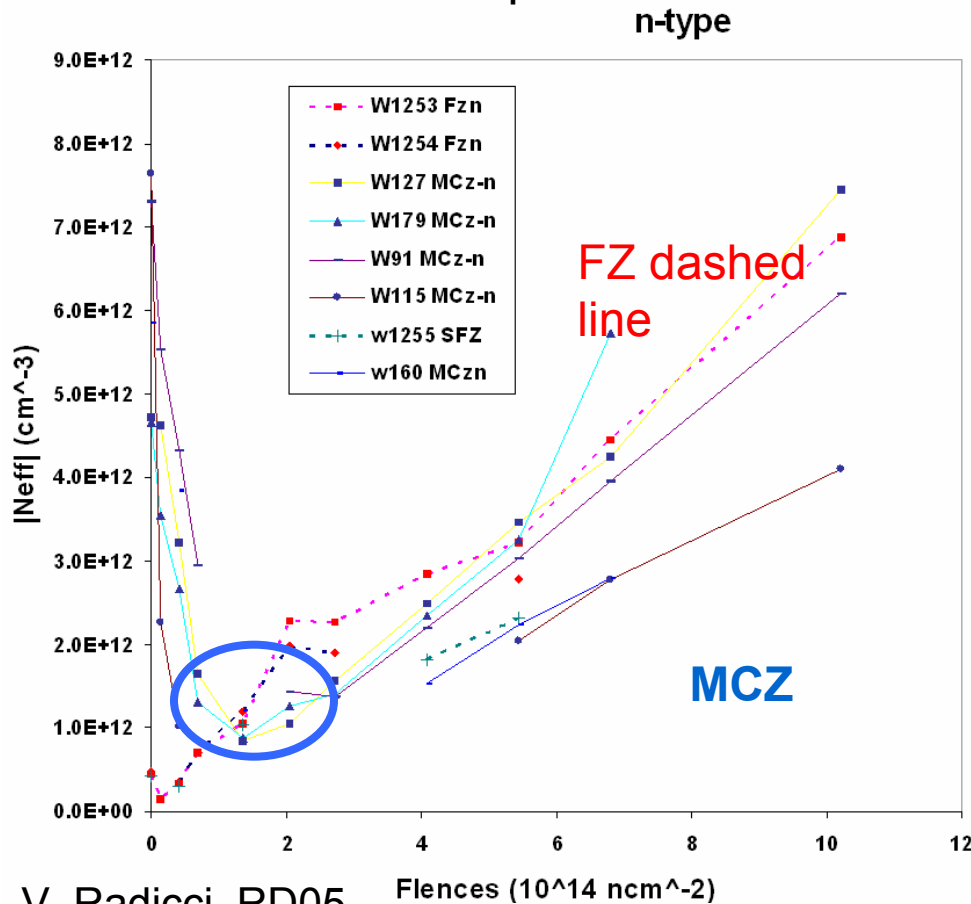


Many groups are studying MCz: INFN, Glasgow, BNL, Helsinki Institute of Physics HIP, Purdue, Liverpool, Rochester etc....

Irradiation studies

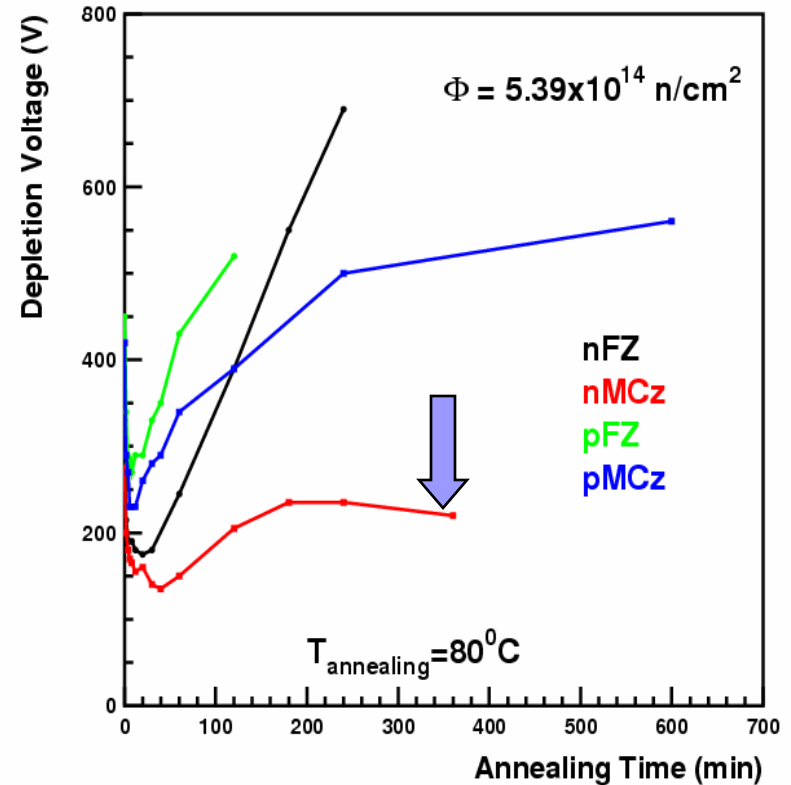
SMART

- Irradiation with 24 GeV/c protons at CERN SPS to 6.0×10^{13} , 3.0×10^{14} , $3.4 \times 10^{15} n_{eq}/cm^2$; Irradiation with 26 MeV protons Karlsruhe in the range: $1.4 \times 10^{13} - 2.0 \times 10^{15} n_{eq}/cm^2$



V. Radicci, RD05

Flences ($10^{14} ncm^{-2}$)

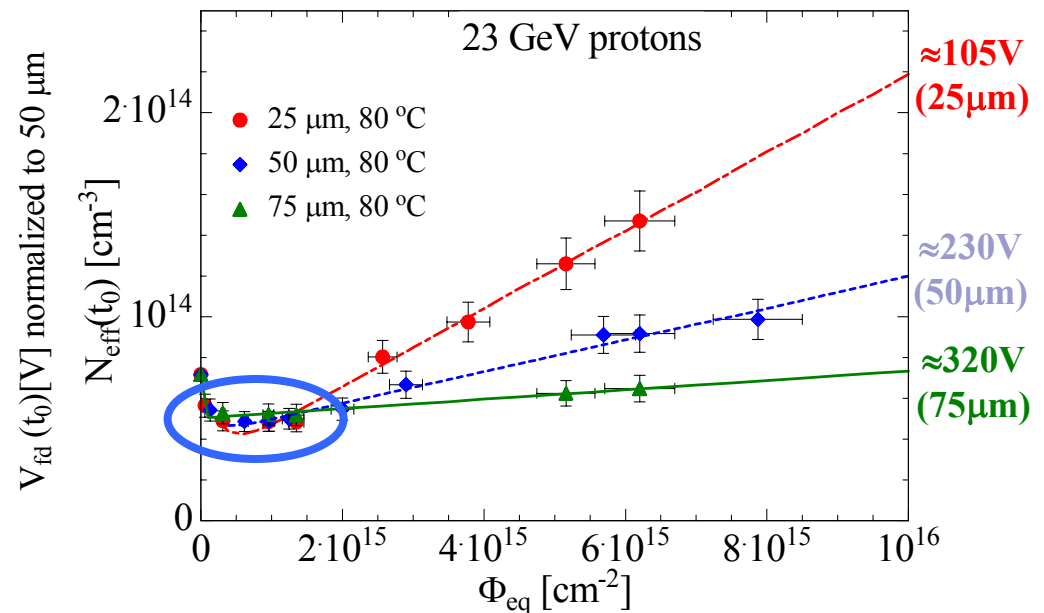
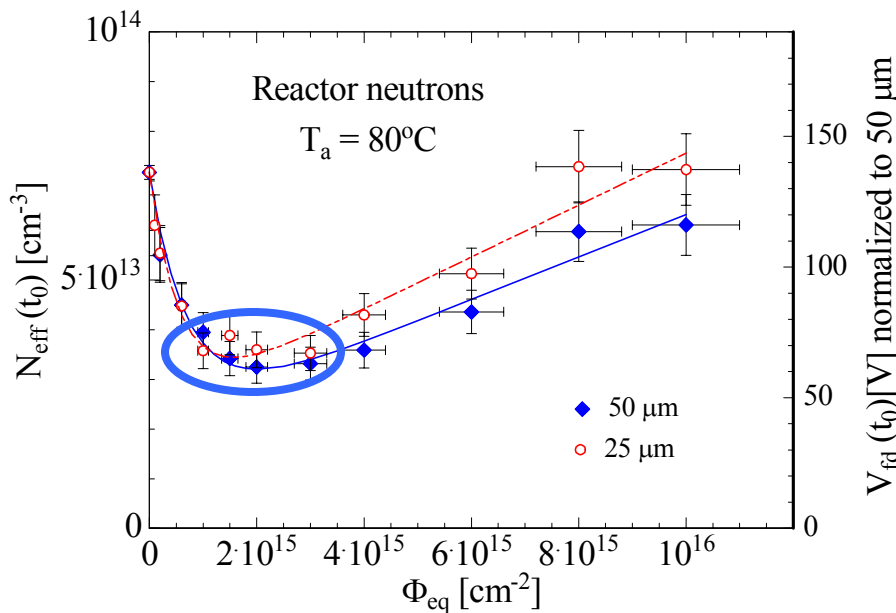


MCz: improved reverse annealing could simplify operational conditions

EPI Irradiation

G.Lindström et al., 10th European Symposium on Semiconductor Detectors, 12-16 June 2005

- Epitaxial silicon grown by ITME
 - Layer thickness: 25, 50, 75 μm ; resistivity: $\sim 50 \Omega\text{cm}$
 - Oxygen: $[\text{O}] \approx 9 \times 10^{16} \text{cm}^{-3}$; Oxygen dimers

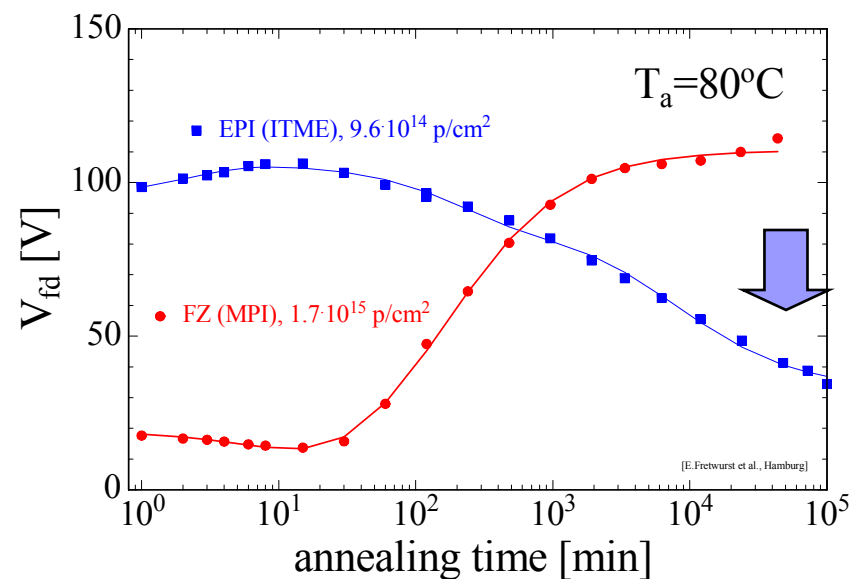
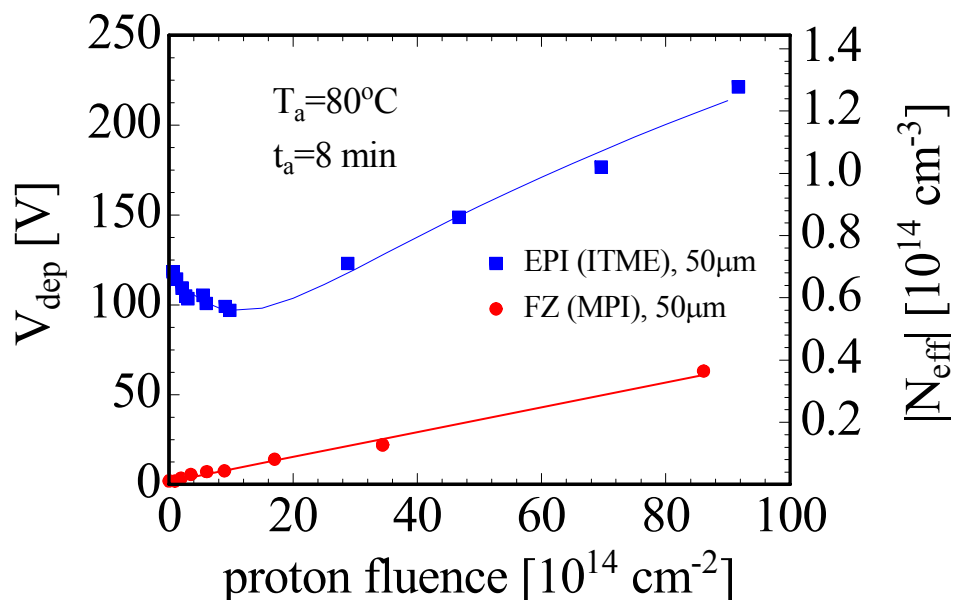


- **No type inversion** in the full range up to $\sim 10^{16} \text{p/cm}^2$ and $\sim 10^{16} \text{n/cm}^2$ (type inversion only observed during long term annealing)
- Proposed explanation: introduction of shallow donors bigger than generation of deep acceptors

EPI Annealing

[E.Fretwurst et al., RESMDD - October 2004]

- 50 μm thick silicon detectors:
 - **Epitaxial silicon** (50 Ωcm on CZ substrate, ITME & CiS)
 - **Thin FZ silicon** (4K Ωcm , MPI Munich, wafer bonding technique)

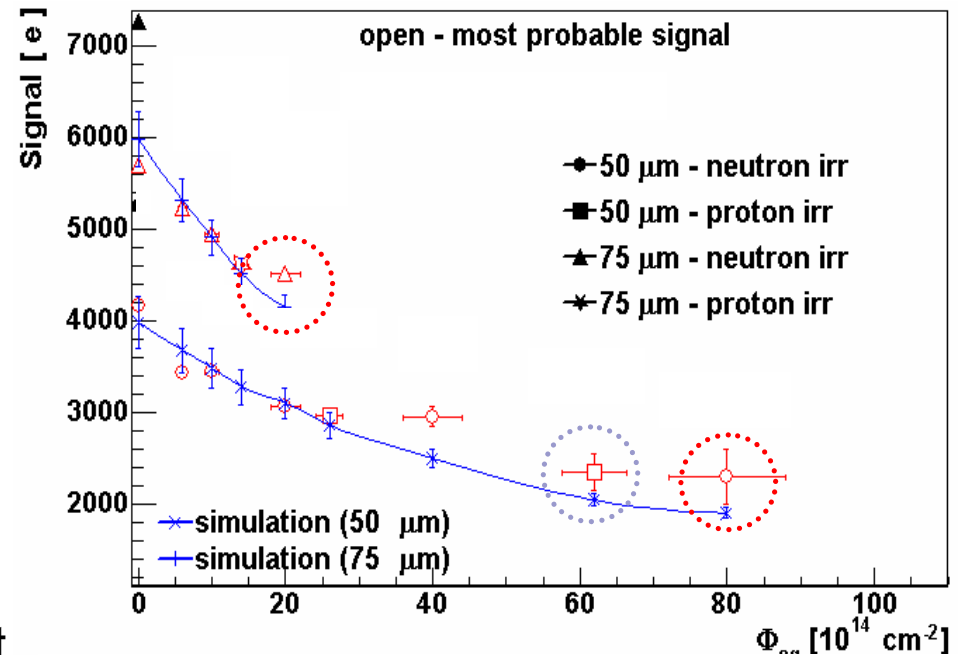
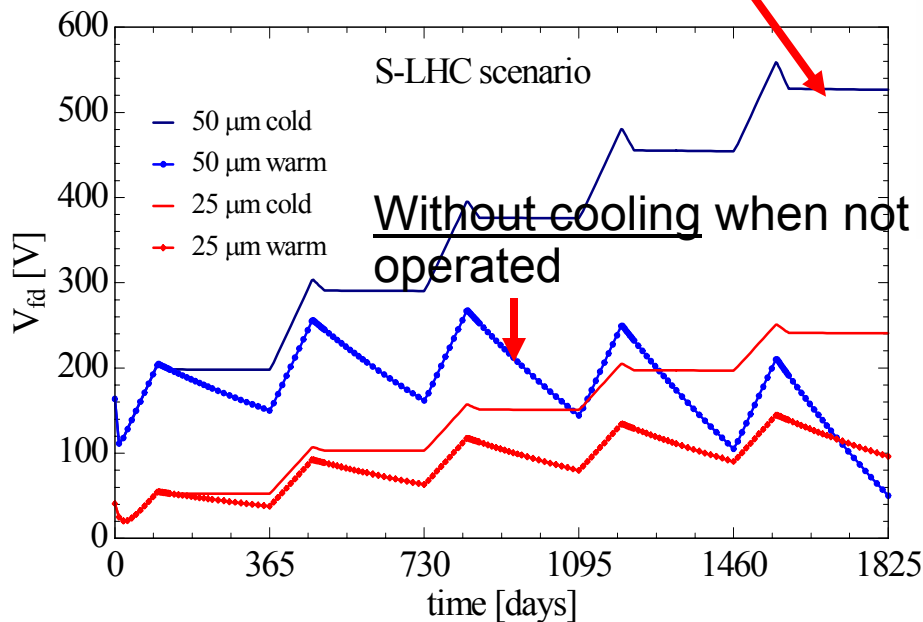


- **Thin FZ silicon:** Type inverted, increase of depletion voltage with time
- **Epitaxial silicon:** No type inversion, decrease of depletion voltage with time \Rightarrow No need for low temperature during maintenance of SLHC detectors!

EPI SLHC

- Radiation @ 4cm: $\Phi_{eq}(\text{year}) = 3.5 \times 10^{15} \text{ cm}^{-2}$
- SLHC-scenario:
 - 1 year = 100 days beam (-7°C)
 - 30 days maintenance (20°C)
 - 235 days no beam (-7°C or 20°C)

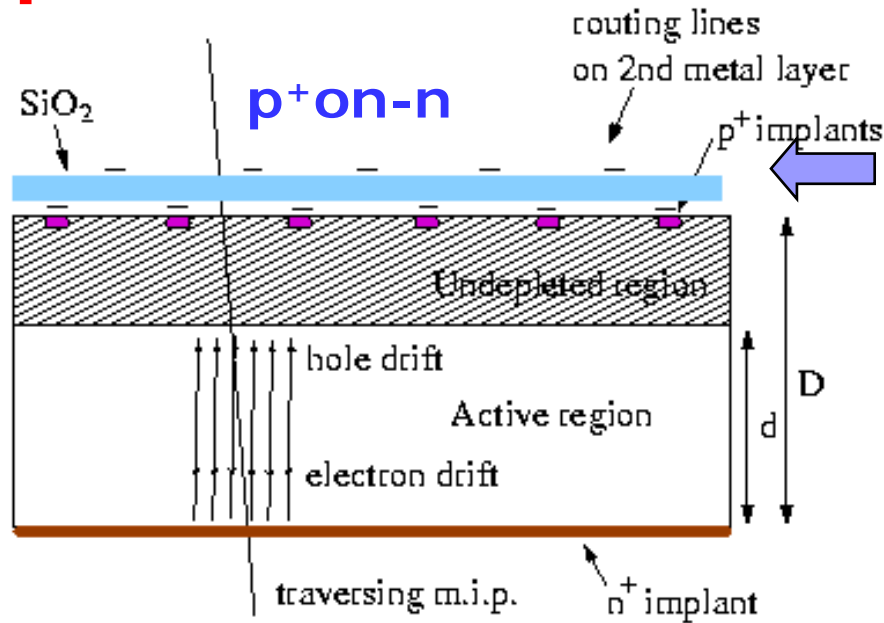
With cooling when not operated



- CCE measured with β from ^{90}Sr
 - 25ns shaping time
 - proton and neutron irradiations of 50 μm and 75 μm epi layers
 - CCE (50 μm) $\Phi_{eq} = 8 \times 10^{15} \text{ n/cm}^{-2}$, 2300 e
 - CCE (75 μm) $\Phi = 2 \times 10^{15} \text{ n/cm}^{-2}$, 4500 e
 - CCE (50 μm): $\Phi = 1 \times 10^{16} \text{ p/cm}^{-2}$ 2400 e

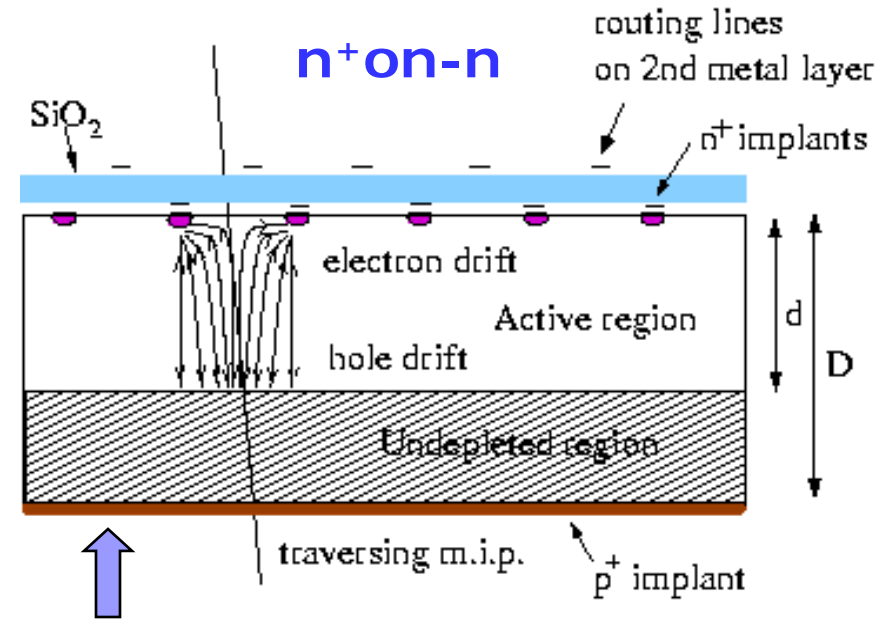
50 μm EPI silicon: a solution for pixels detectors at SLHC?

p-in-n versus n-in-n detectors



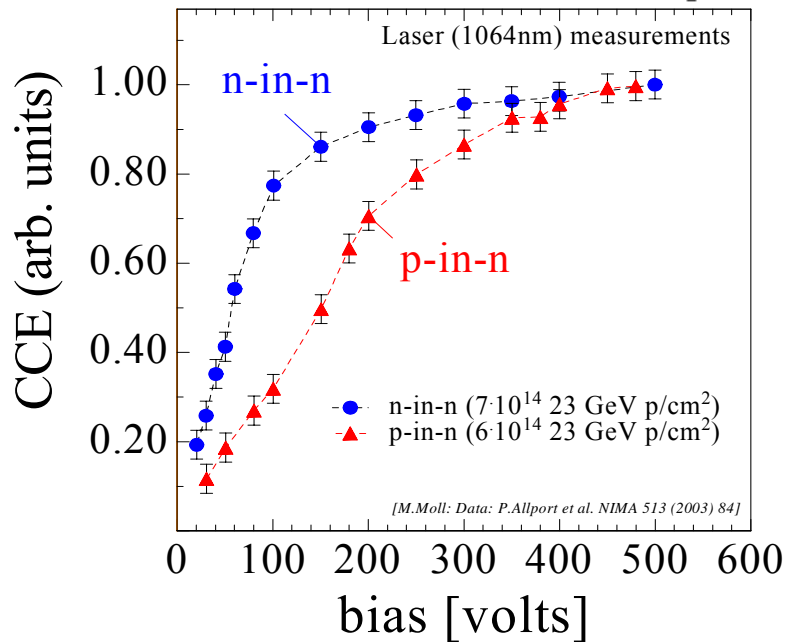
p-on-n silicon, under-depleted:

- Charge spread \Rightarrow degraded resolution
- Charge loss \Rightarrow reduced CCE



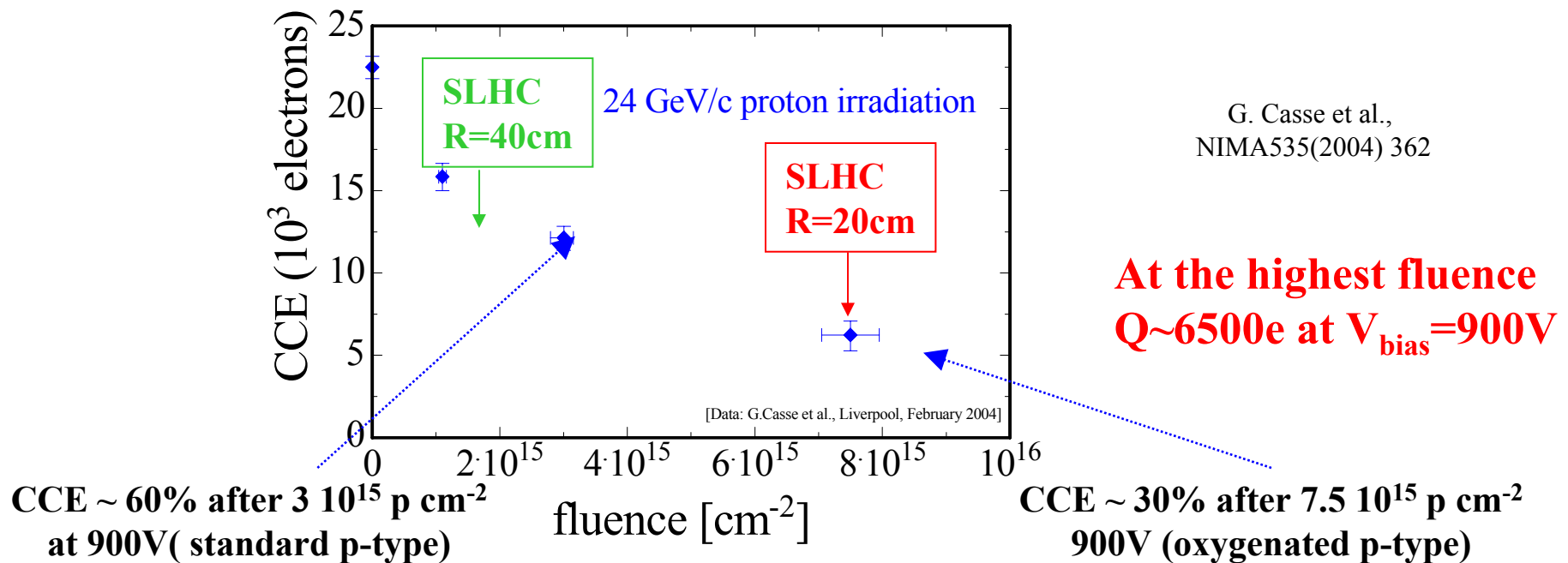
n-on-n silicon, under-depleted:

- Limited loss in CCE
- Less degradation if under-depleted
- Collect electrons (fast)



n-in-p microstrip detectors

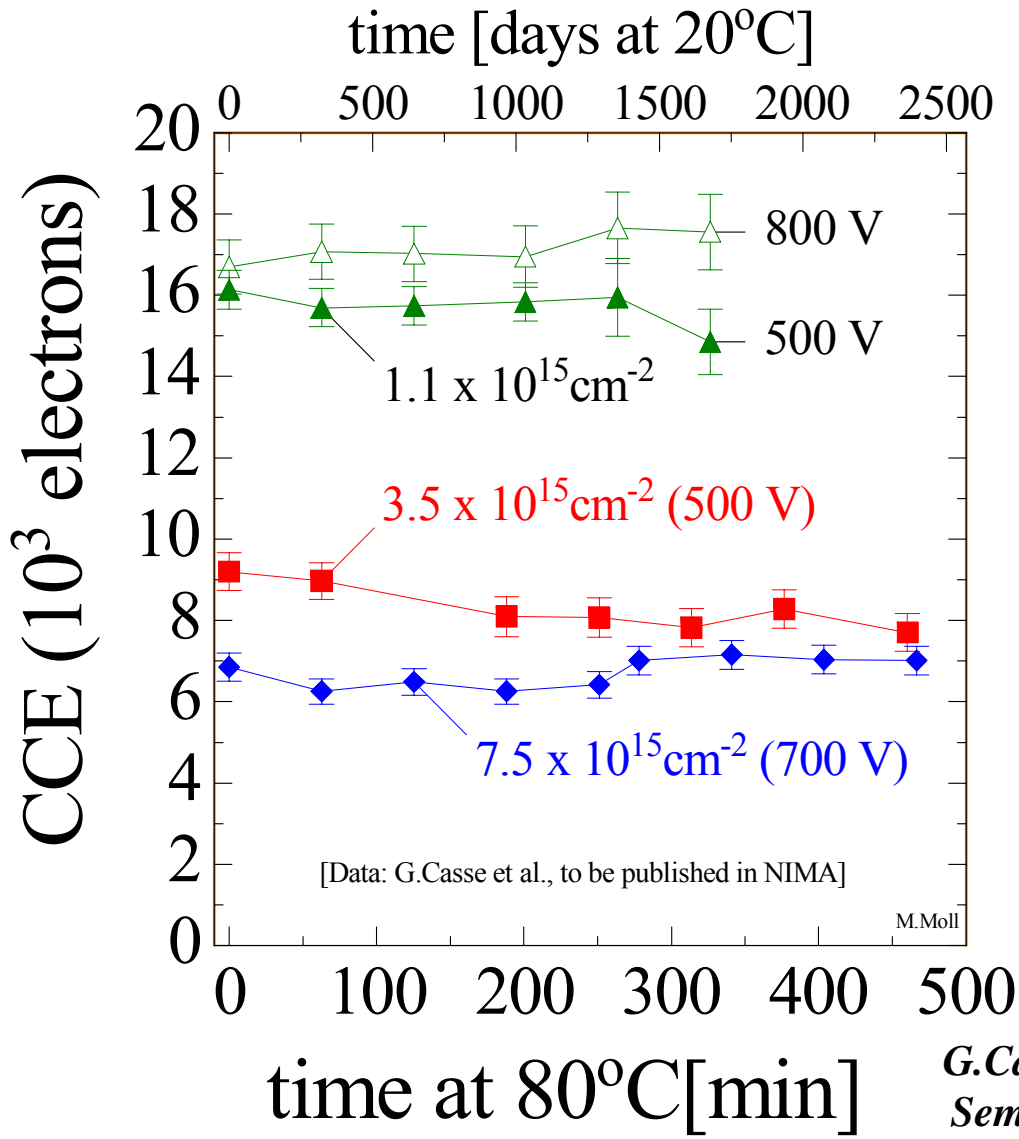
- Miniature n-in-p microstrip detectors (280 μm)
- Detectors read-out with LHC speed (40MHz) chip (SCT128A)
- Material: standard p-type and oxygenated (DOFZ) p-type



Charge collection in planar silicon detectors might be sufficient for all but inner-most Pixel layer!

Benefit: Single sided processing ~50% cheaper than n-in-n

Annealing of p-type sensors



- p-type strip detector (280 μm) irradiated with 23 GeV p ($7.5 \times 10^{15} \text{ p/cm}^2$)
- expected from previous CV measurement of V_{dep} :
 - before reverse annealing: $V_{\text{dep}} \sim 2800\text{V}$
 - after reverse annealing $V_{\text{dep}} > 12000\text{V}$
- no reverse annealing visible in the CCE measurement !

G.Casse et al., 10th European Symposium on Semiconductor Detectors, 12-16 June 2005

Novel Materials

Property	Diamond	GaN	4H SiC	Si
E_g [eV]	5.5	3.39	3.26	1.12
$E_{\text{breakdown}}$ [V/cm]	10^7	$4 \cdot 10^6$	$2.2 \cdot 10^6$	$3 \cdot 10^5$
μ_e [cm^2/Vs]	1800	1000	800	1450
μ_h [cm^2/Vs]	1200	30	115	450
v_{sat} [cm/s]	$2.2 \cdot 10^7$	-	$2 \cdot 10^7$	$0.8 \cdot 10^7$
Z	6	31/7	14/6	14
ϵ_r	5.7	9.6	9.7	11.9
e-h energy [eV]	13	8.9	7.6-8.4	3.6
Density [g/cm^3]	3.515	6.15	3.22	2.33
Displacem. [eV]	43	19.2±2	25	13-20

■ Wide bandgap
diamond=5.5
SiC=3.3eV

< leakage current
than silicon

■ Signal:

Diamond	36 e/ μm
SiC	51 e/ μm
Si	89 e/ μm

> charge than
diamond

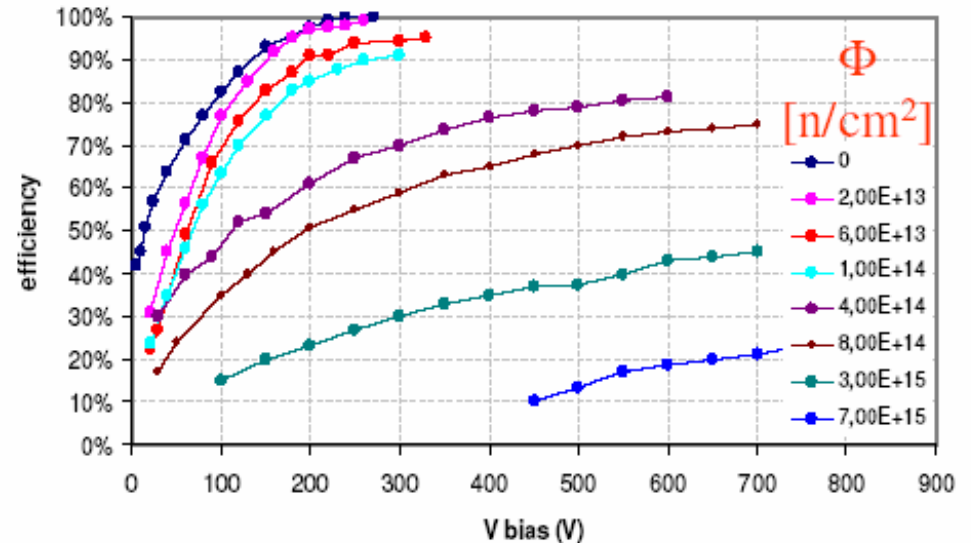
■ > displacement
threshold than silicon
⇒ radiation harder than
silicon (?)

■Diamond:

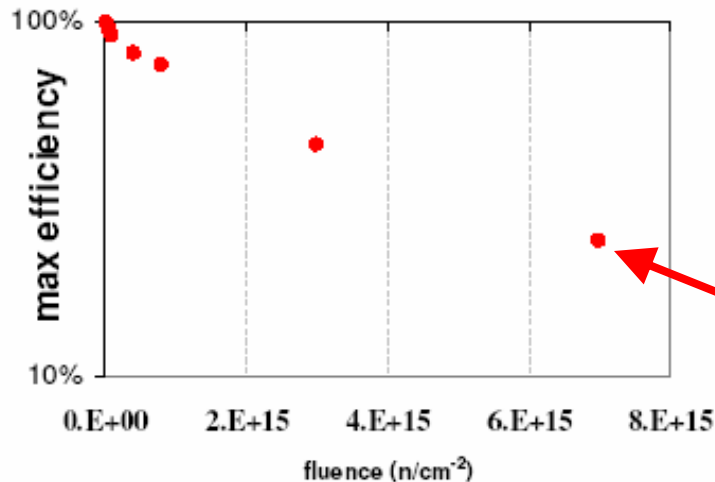
- Dielectric constant ($2.1 \times$ lower than Si) → low capacitance
- Higher Electron and hole mobility → fast collection times

SiC: CCE after irradiation

- Material: epitaxial layers by CREE Res. Inc. and IKZ (Institut für Kristallzüchtung, Berlin)
- Devices: Schottky diodes, Alenia Marconi Systems (Rome)
- Depletion depth: 20-40 μm
- Effective doping: $5.3 \times 10^{14} \text{ cm}^{-2}$
- Irradiated with protons at CERN PS to $1.6 \times 10^{16}/\text{cm}^2$ and neutrons at Ljubjana to $7 \times 10^{15}/\text{cm}^2$



S.Sciortino et al., presented on the RESMDD 04 conference, in press with NIMA



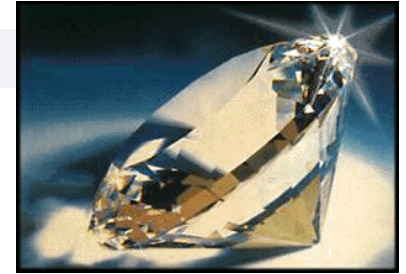
CCE before irradiation

- 1100 e⁻ @400 V with α particles
- 1400 e⁻ @200 V with MIPS (100% CCE)

CCE after irradiation

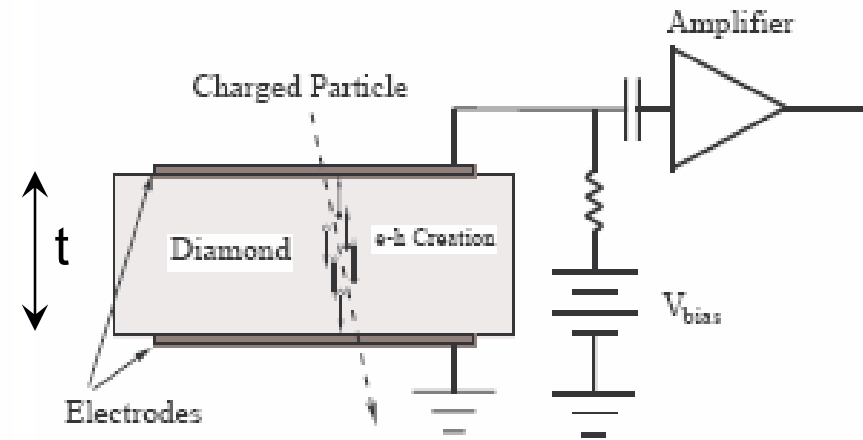
- 20% CCE (α) after $7 \times 10^{15} \text{ n/cm}^2$!
- 35% CCE (β) ($\sim 300 \text{ e}^-$) after $1.4 \times 10^{16} \text{ p/cm}^2$ much less than in silicon

Diamond

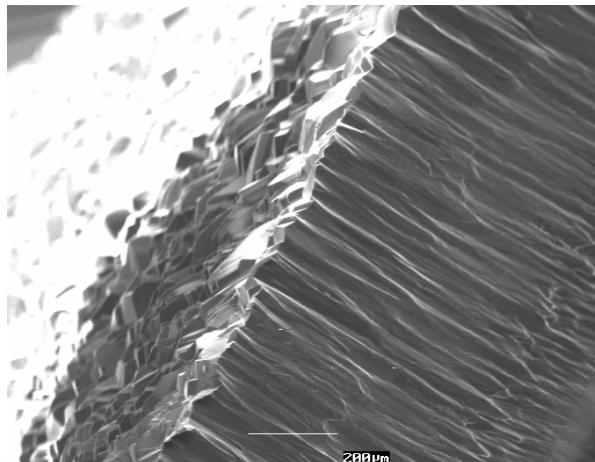


- Ionization energy is high: $MIP \approx 2x$ less signal for same X_0 (w.r.t. SI)
 - Diamond: $\sim 13.9ke^-$ in $361 \mu m$
 - SI: $\sim 26.800 ke^-$ in $282 \mu m$
- In Polycrystalline Diamond grain-boundaries, dislocations, and defects:
 - limits carrier lifetime, mobility and charge collection distance and position resolution

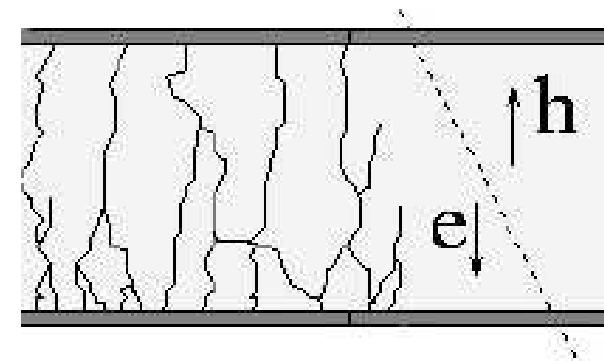
Signal formation



$Q = Q_0 d / t$ where $d =$ collection distance = distance e-h pair move apart

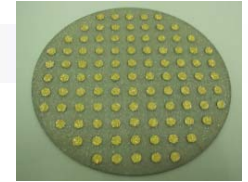


Polycrystalline Diamonds traditionally grown by CVD



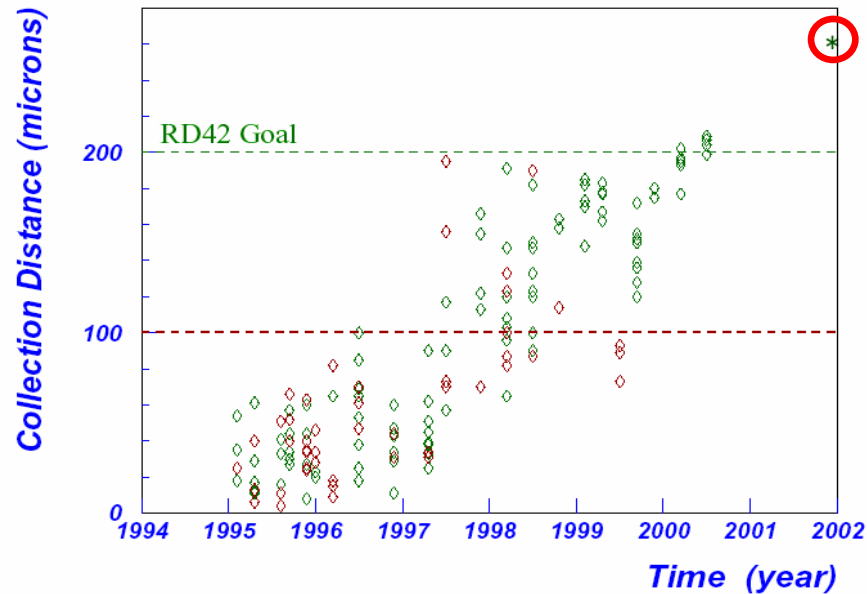
Daniela Bortoletto

Polycrystalline Diamonds



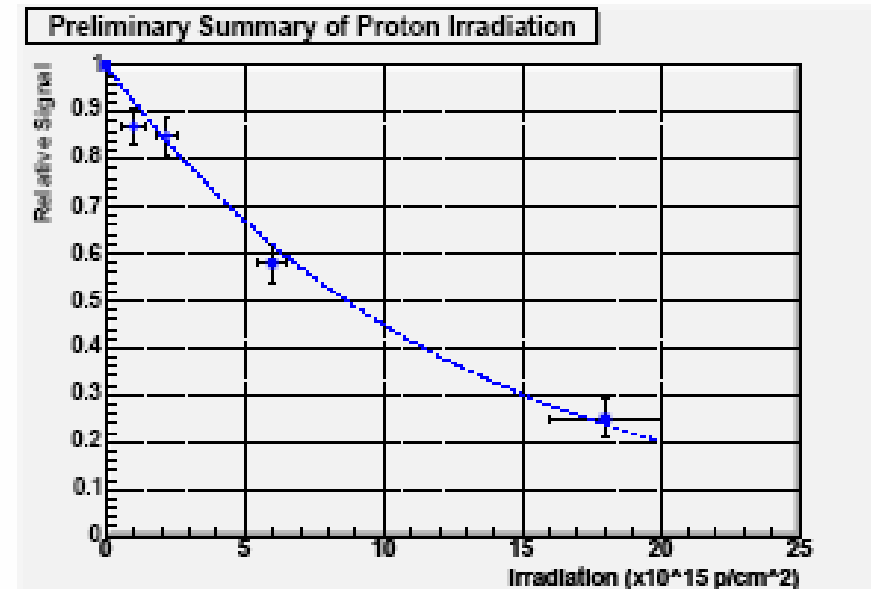
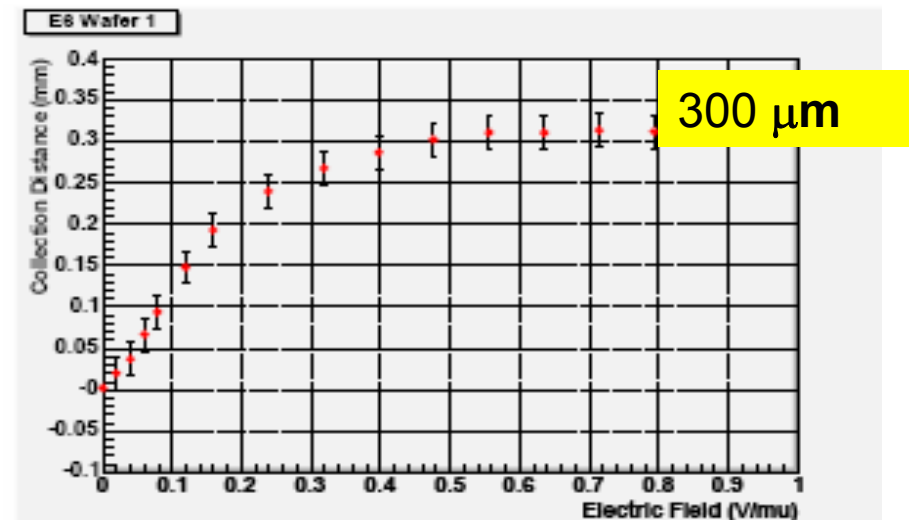
- RD42 in collaboration with vendors have achieved collection distance > 300 μm

Charge Collection in DeBeers CVD Diamond



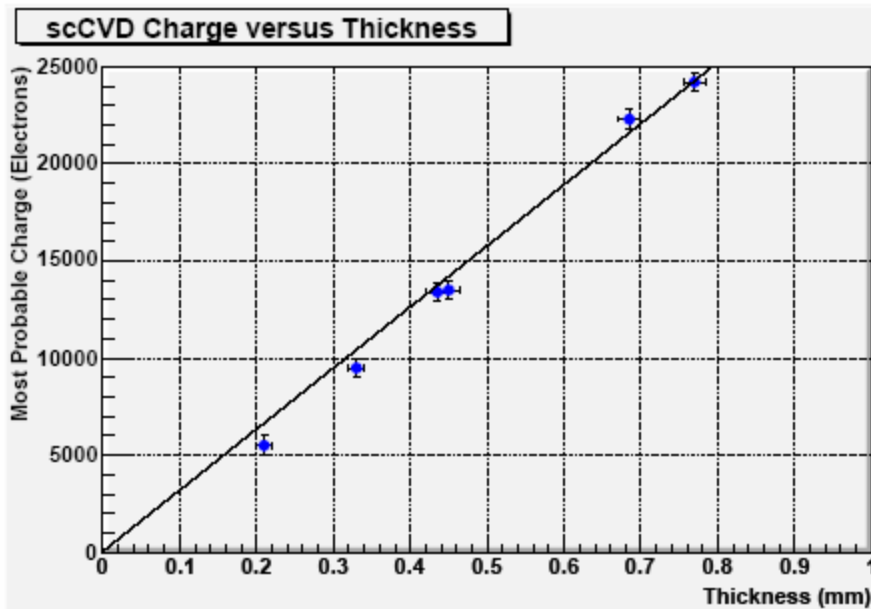
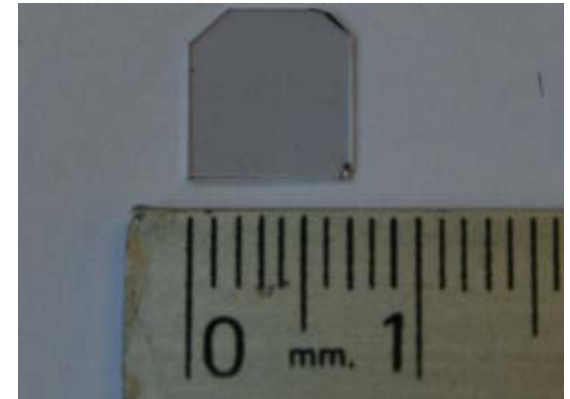
–Wafers diameter >12 cm

- Excellent radiation hardness
 - 60% CCE at $2.9 \times 10^{15} \text{ p/cm}^2$
 - 23% improvement in resolution
 - 25% CCE at $1.8 \times 10^{16} \text{ p/cm}^2$
- Used in successfully for radiation monitoring for BaBar, Belle, CDF, CMS



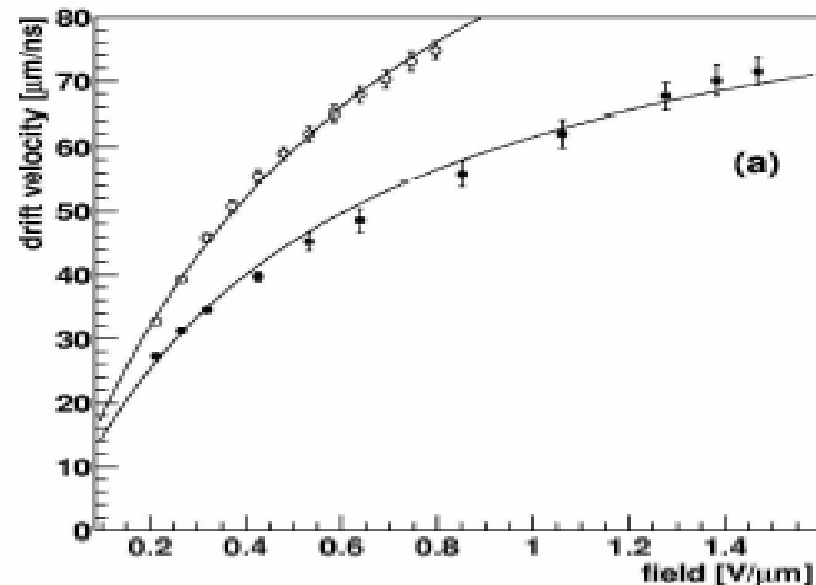
Single Crystal Diamond

- Single crystal diamond has been fabricated with Element six $\approx 10 \text{ mm} \times 10 \text{ mm}$, $>1 \text{ mm}$ thickness.
- Largest scCVD diamond $\approx 14 \text{ mm} \times 14 \text{ mm}$.



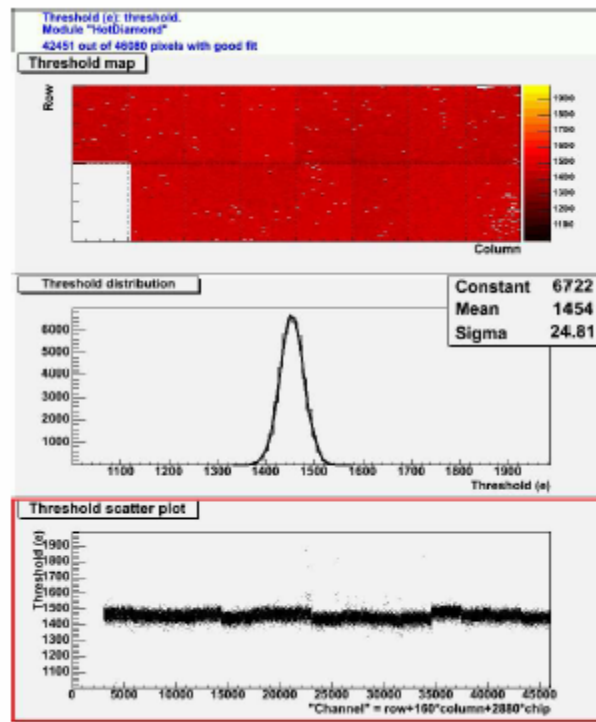
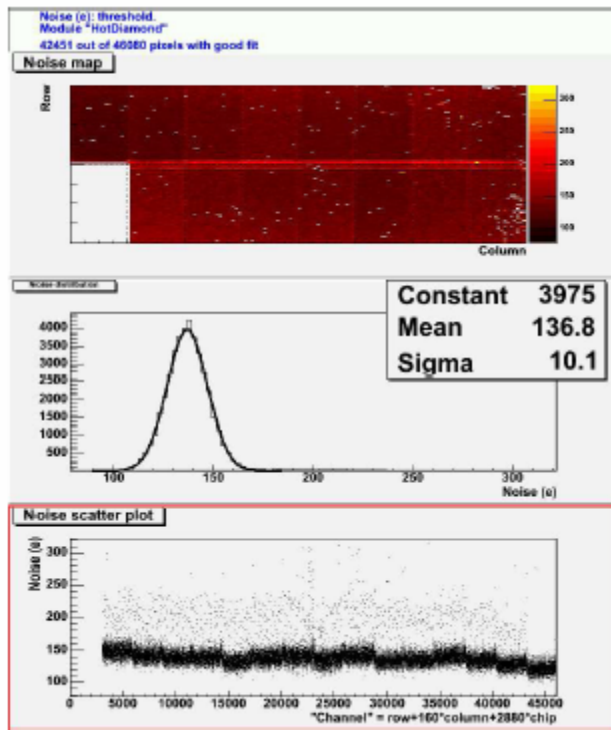
- High quality scCVD diamond can collect full charge
- Width of Landau distribution is $\approx 1/2$ that of silicon, $\approx 1/3$ that of pCVD diamond

- Excellent mobility. For this sample:
 - $\mu_{oh} = 1714 \text{ cm}^2/\text{Vs}$, $\mu_{oe} = 2064 \text{ cm}^2/\text{Vs}$
- High drift velocity \Rightarrow better lifetimes \Rightarrow charge trapping might not be an issue

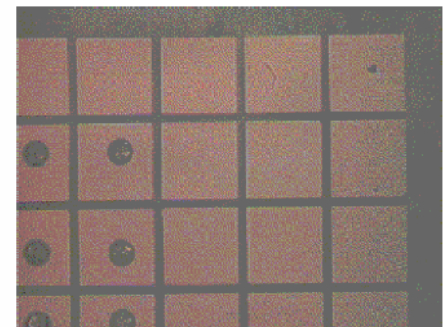


Diamond Atlas module

- Beam test at DESY with 4-6 GeV electrons
- Results: Noise $\sim 137e$, Mean Threshold $1454e$, Threshold Spread $\sim 25e$.



Harris Kagan
Vertex 2005
Nikko, Japan



CMS diamond pixel

- Preliminary efficiency $>97.5\%$
 - still need to correct for dead or missing channels.

Device Engineering: 3D detectors

(Introduced by S.I. Parker et al., NIMA 395 (1997) 328)

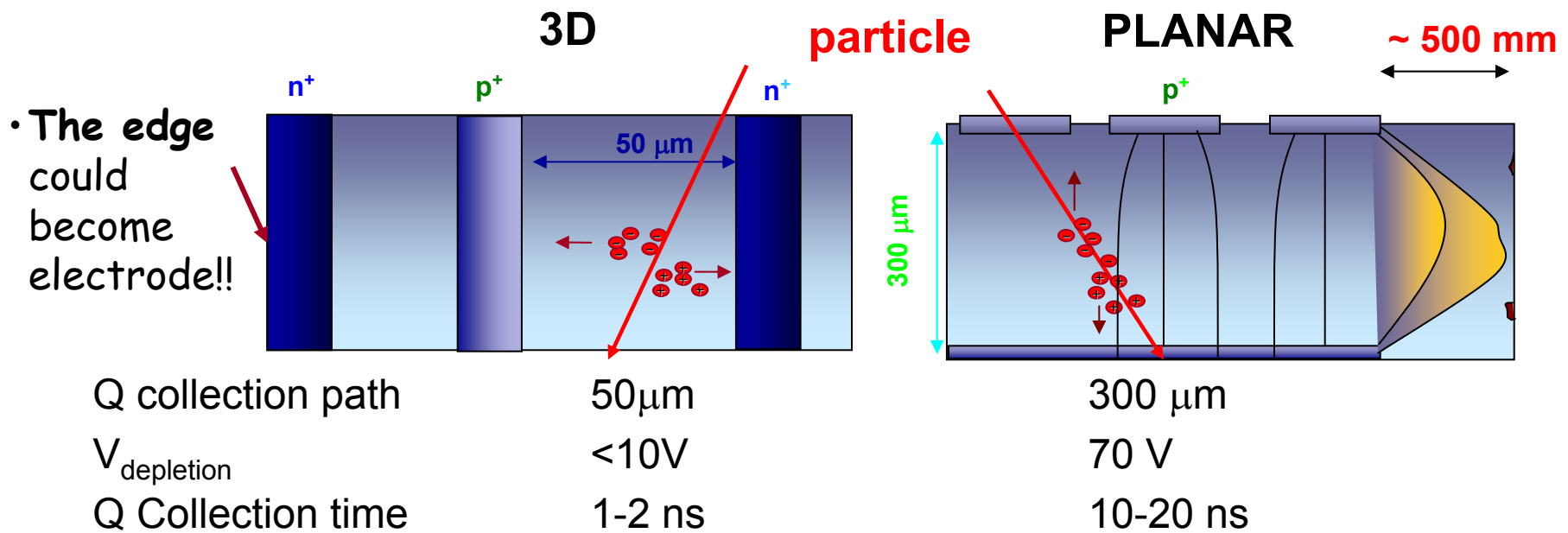
■ Combine **VLSI** and **MEMS** (Micro Electro Mechanical Systems)

■ **Electrodes:**

- Narrow columns processed inside the bulk instead then implanted on surface: 3D
- Diameter: $10\mu\text{m}$; Distance: $50\text{-}100\mu\text{m}$

■ **Lateral depletion:**

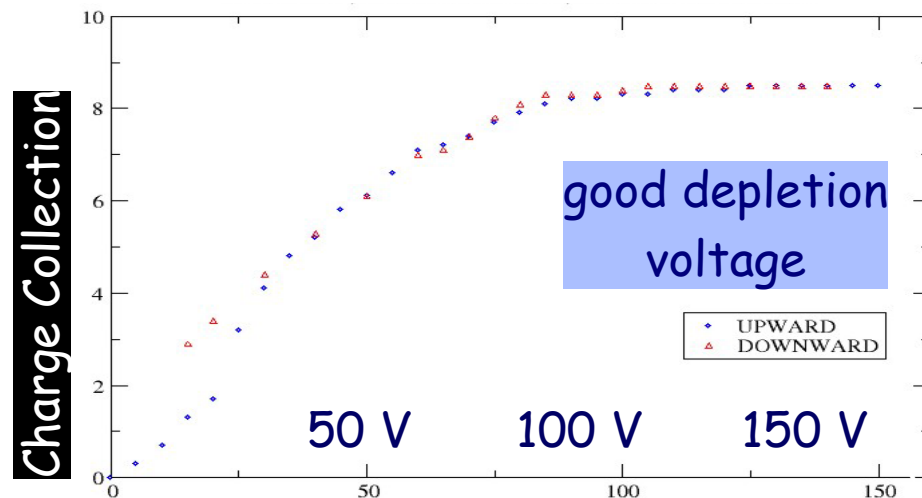
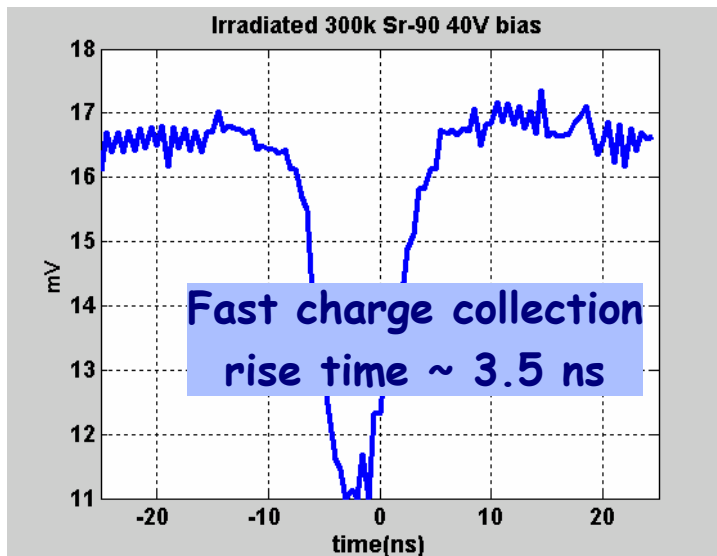
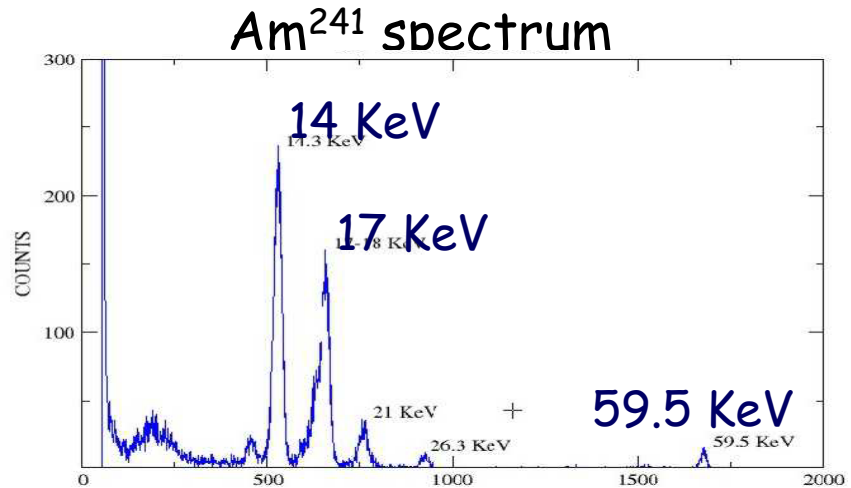
- Lower depletion voltage
- Short collection distance \Rightarrow fast signal
- More rad hard



3D detectors: characteristics

- Low leakage currents
- Low depletion voltages
- Gaussian X ray lines
- Fast charge collection

Performance after irradiation $\sim 10^{15}$ p/cm²



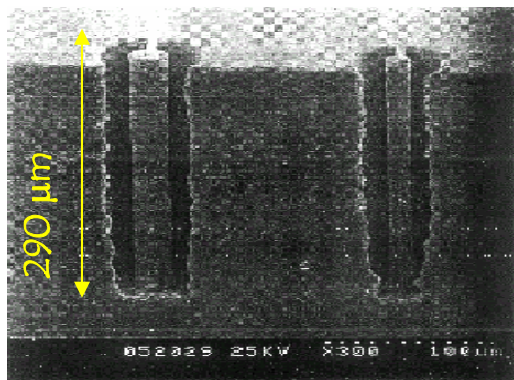
3D DETECTOR FABRICATION

Non Standard Processing: Wafer bonding, Deep reactive ion etching, Low pressure chemical vapor deposition, Metal deposition \Rightarrow Mass production expensive

1) ETCHING THE ELECTRODES



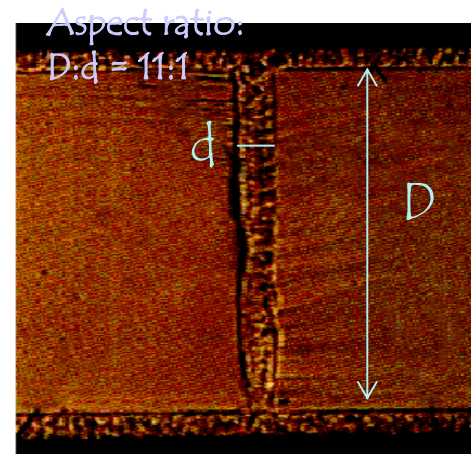
WAFER BONDING
(mechanical stability)
 $\text{Si-OH} + \text{HO-Si} \rightarrow \text{Si-O-Si} + \text{H}_2\text{O}$



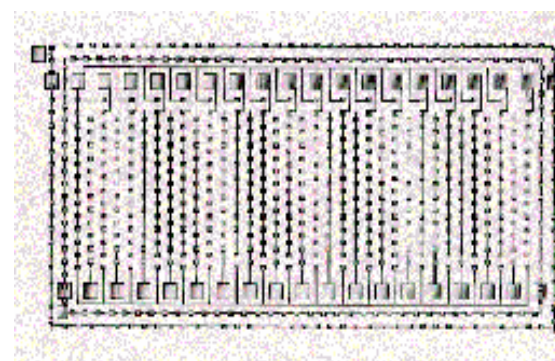
DEEP REACTIVE ION ETCHING
(electrodes definition)
Bosh process
 SiF_4 (gas) + C_4F_8 (teflon)

C shaped test structure
 $\sim 1 \mu\text{m}$ difference between top and bottom

2) FILLING THE ELECTRODES



LOW PRESSURE CHEMICAL VAPOR DEPOSITION
(Electrodes filling with conformal doped polysilicon)
 $2\text{P}_2\text{O}_5 + 5\text{Si} \rightarrow 4\text{P} + 5\text{SiO}_2$
 $2\text{B}_2\text{O}_3 + 3\text{Si} \rightarrow 4\text{B} + 3\text{SiO}_2$



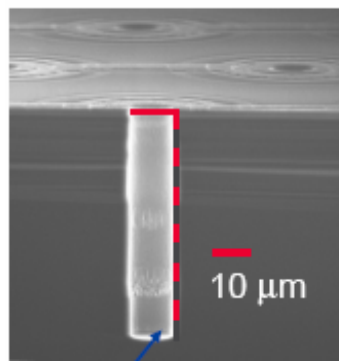
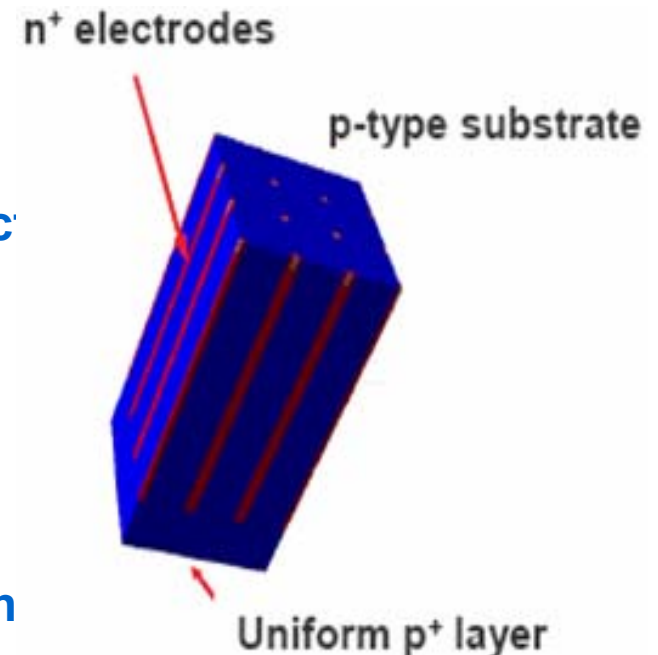
METAL DEPOSITION
Shorting electrodes of the same type with Al for strip electronics readout or deposit metal for bump-bonding

3D detectors

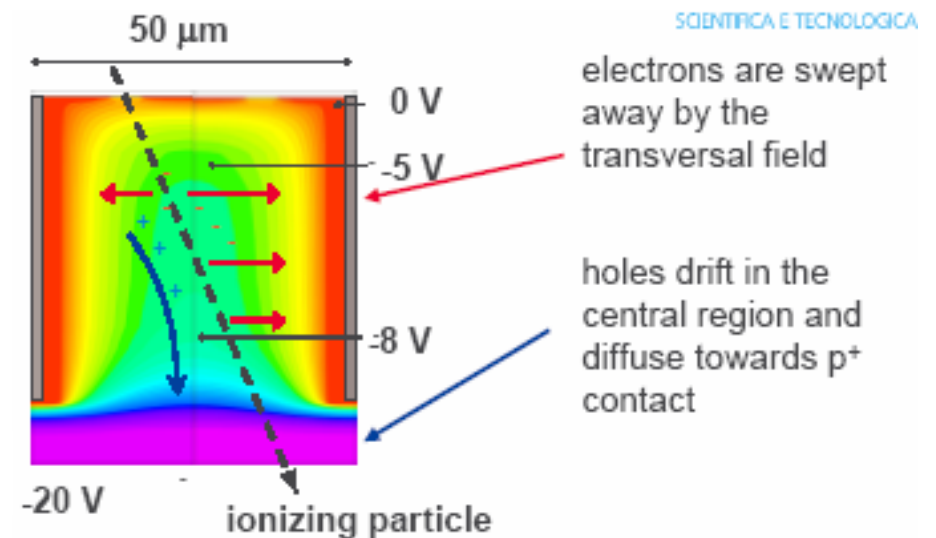
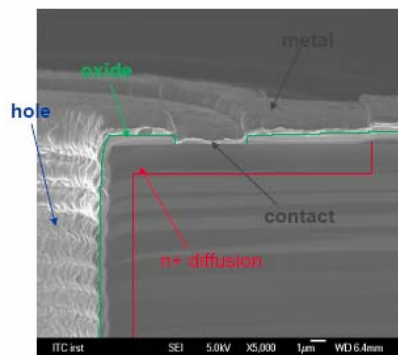
Claudio Piemonte (*ITC-irst*)

- **3D Single Type Column (3D-STC)** aiming at process simplification

- n⁺ columns in p-type substrate
- Bulk contact provided by a uniform p⁺ contact on backside
- Holes not etched through the wafer
- No hole filling (holes are doped but not filled with polysilicon)
- CNM: Hole etching (DRIE); IRST: other processing (contacts or polysilicon deposition etc.)



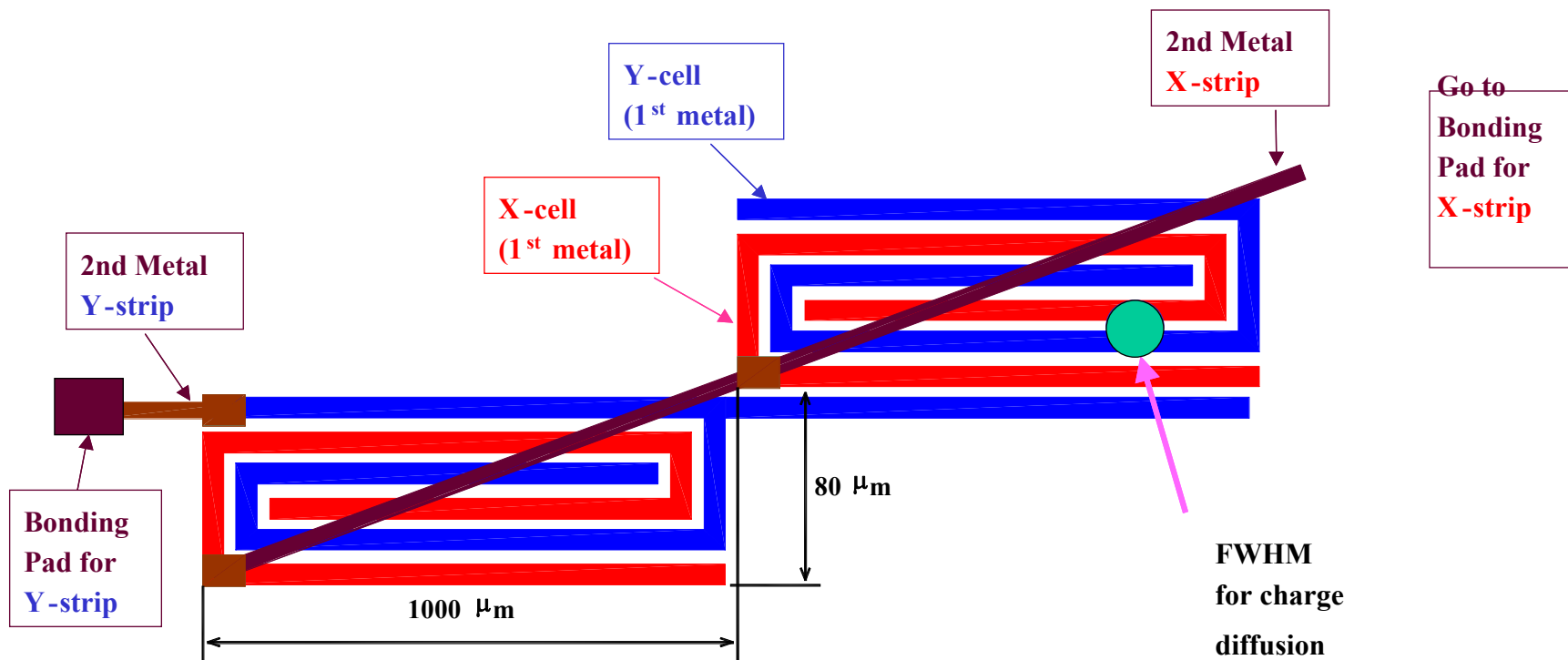
Hole depth: 120 μm



IRST-Trento and CNM Barcelona

Other new structures: Stripixel

- Several concepts for new (**planar and mixed planar & 3D**) detector structures aiming for improved radiation tolerance or less costly detectors (see e.g. Li - 6th RD50 workshop, or Bortoletto- 5th RD50 Workshop)
- Example: Stripixel concept or semi 3D:



Z. Li, D. Lissauer, D. Lynn, P. O'Connor, V. Radeka

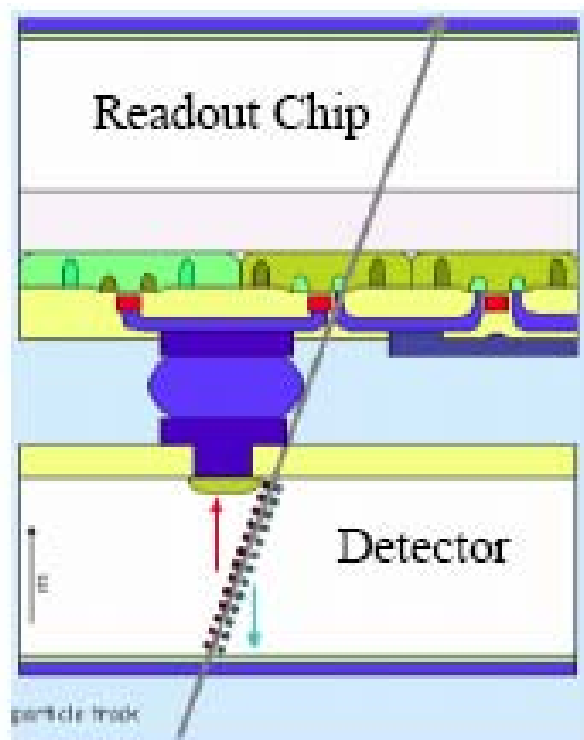
Monolithic Active Pixel Sensors (MAPS)

■ Hybrid Pixel sensors have achieved a level of maturity in HEP. Future problems are cost, mass, and cooling of detectors under high radiation.

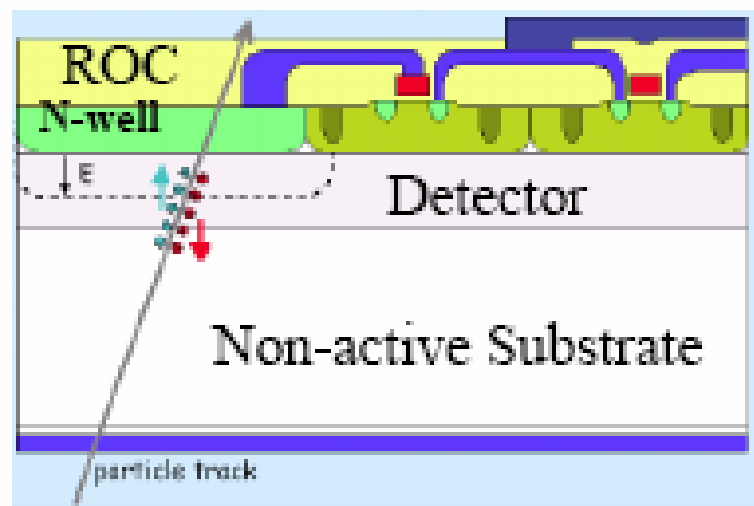
■ Much work is being done on MAPS to reduce mass (ILC)
■ A MAPS is a silicon structure where the detector and the primary readout electronics are processed on the same substrate.

- Only the top few microns of an IC contain active circuitry.
- The rest is merely a support structure.

Bum
bond



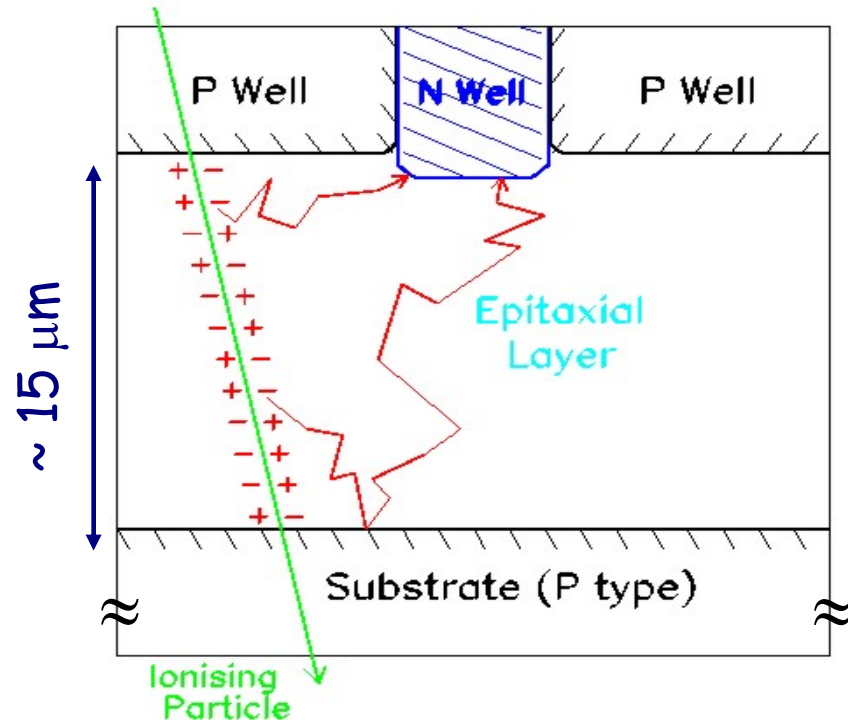
Hybrid Pixel Principle



MAPS Principle

Kucewicz, Krakow

MAPS with Standard CMOS processes



- Many groups studying this concept: RAL, IReS, Hawaii, INFN SLIM etc.

- Principle of operation:

- Signal charge created in the epitaxial layer $Q=80e\text{-h}/\mu\text{m}$
- The charge is $\approx 1000 e\text{-}$ and it can be measured because of the small capacitance of the electrode
- Used in CMOS cameras

- Advantages

- signal processing integrated on sensor substrate
- Sensors may be thinned down to $<20 \mu\text{m}$
- Standard processing \Rightarrow chip and fast turn around

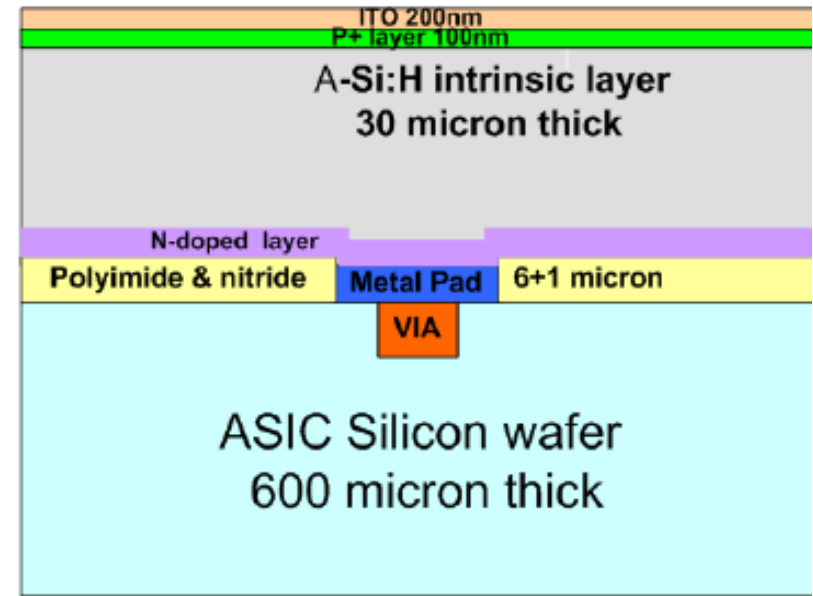
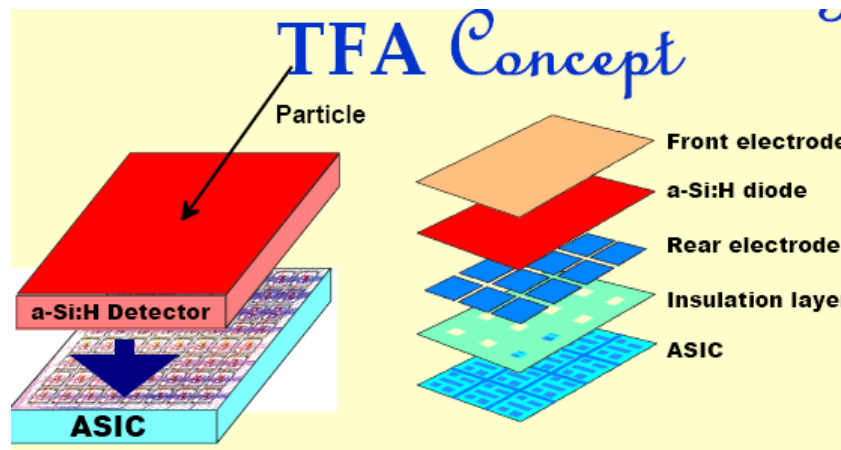
- Challenges:

- Transistors options are limited
- Newer processes have thinner or no-epi \Rightarrow vary small signals
- Sparsification difficult to implement
- Limited readout speed
- Triple well processes in $0.13 \mu\text{m}$ are promising

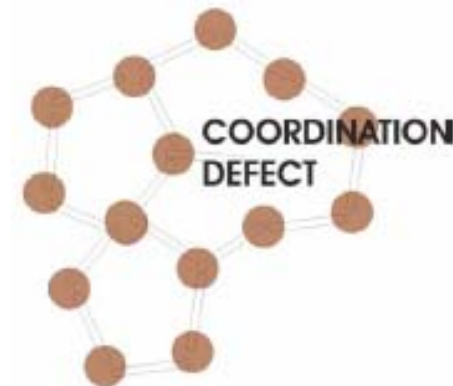
Thin Film on Asic (TFA)

Development centered at CERN, Pierre Jarron's talk at Vertex 04

- TFA is an emerging pixel sensor technology
- Deposition of a-Si:H layer above readout ASICs



A-Si:H



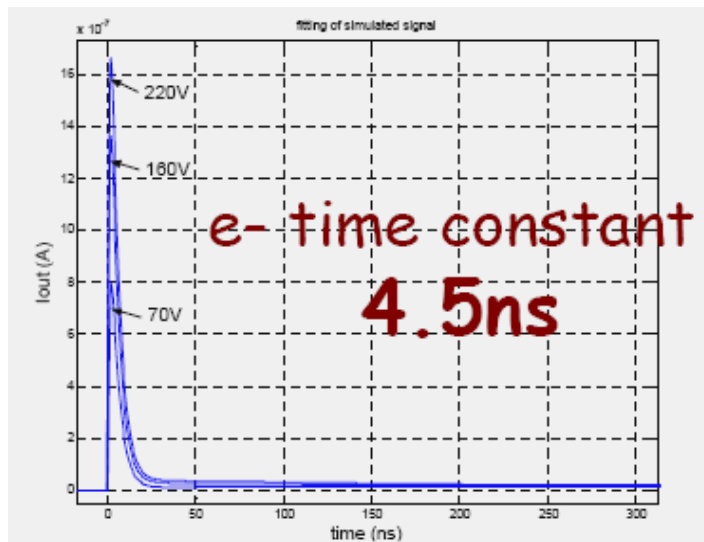
Amorphous structure

- Dangling bonds compensated by H
- H compensates impurities or radiation-induced defects
- Short time annealing
- Band tail formation due to bonding disorder

A:Si-H properties and rad hardness

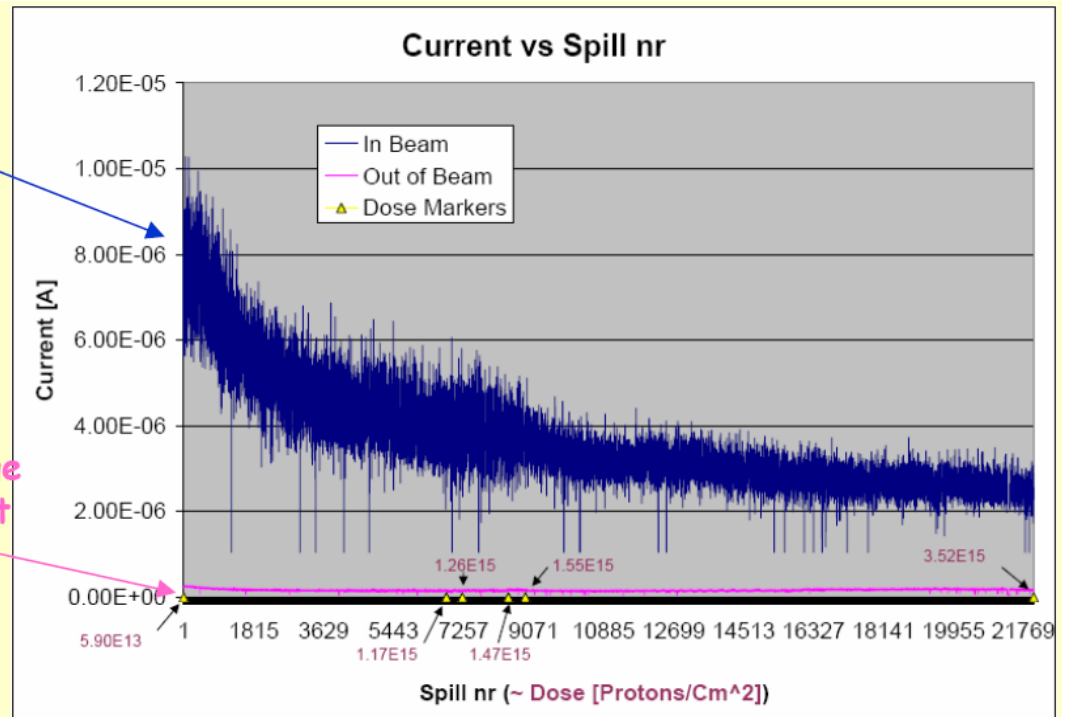
■ a-Si:H film ultra rad hard:

- No I_{leak} increase after a fluence of $1.8 \times 10^{16} p/cm^2$.
- Slight degradation of CCE and Photo-Conductivity after $3.5 \times 10^{15} p/cm^2$ CCE. It can be recovered after 1 hour annealing at $150^\circ C$



Proton signal

Leakage current

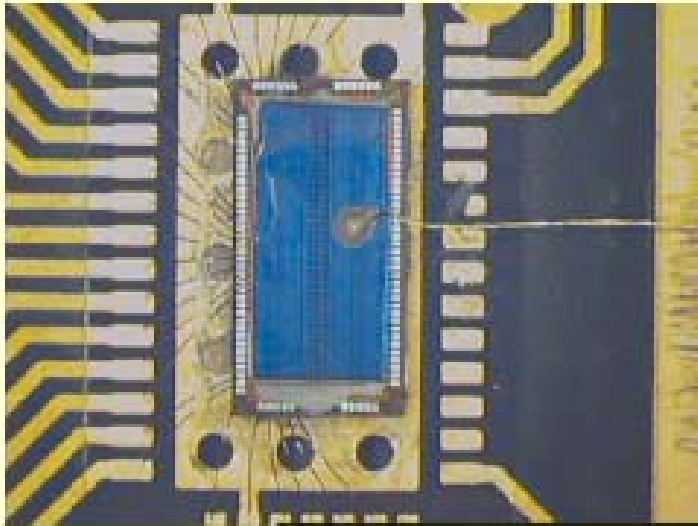


a-Si:H detector film signal

- fast component from e⁻ drift

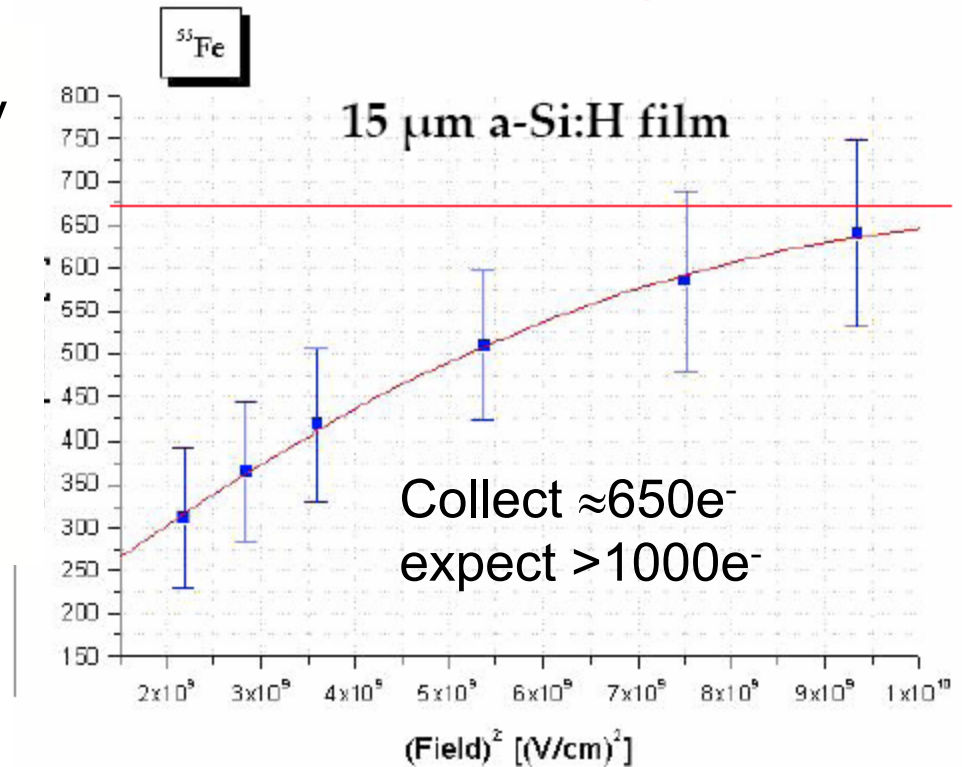
Advantages/Technical challenges

- CCE Not yet understood:
 - e-h creation energy 3.6-5 eV
- Develop thicker film
- Low noise readout to handle small signal
- Non-commercial process, Deposition with plasma reactor at IMT, Neuchatel



30 μm thick sample

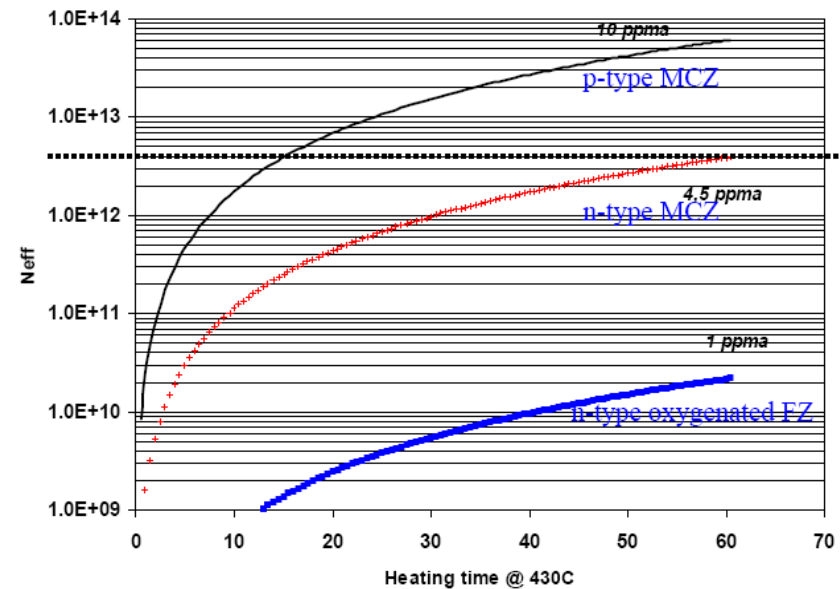
Planarity \Rightarrow edge effects



- Ultra-rad hard sensors
- Low cost pixel detector technology
 - Film deposition cost estimated $\approx 30\%$ of deep sub-micron wafer.
 - For 8", about \$2000 + \$600 for 100cm² (yield 75%)

Change of operation conditions

- Device Recovery/Improvement Via Elevated temperature annealing (DRIVE)
- Thermal annealing of radiation damage in FZ or DOFZ:
 - If $T > 450$ °C will destroy the detector and/or electronics
 - If $T < 450$ °C, reverse annealing (generation of more negative space charges) makes detectors worse
- For high resistivity MC_Z Si thermal annealing with $T = 200$ °C - 450 °C will generate thermal donor (TD, positive space charge) due to high oxygen concentration [O]
- IDEA: Adjust annealing T and time, and [O] so that TD creation rate cancels compensates the original negative space charges due to irradiation \Rightarrow manageable V_{dep}



DRIVE first results

Z. Lee BNL

■ Detectors:

- p⁺-n-n⁺ pad structures with multi-GRs n-type MCz Si $\rho=1$ kOhm·cm,
- p⁺-p-n⁺ pad structures with multi-GRs p-type MCZ Si $\rho=3$ kOhm·cm

■ Irradiation 24 GeV and 20 GeV protons

■ Measurements:

□ After each annealing step (at BNL):

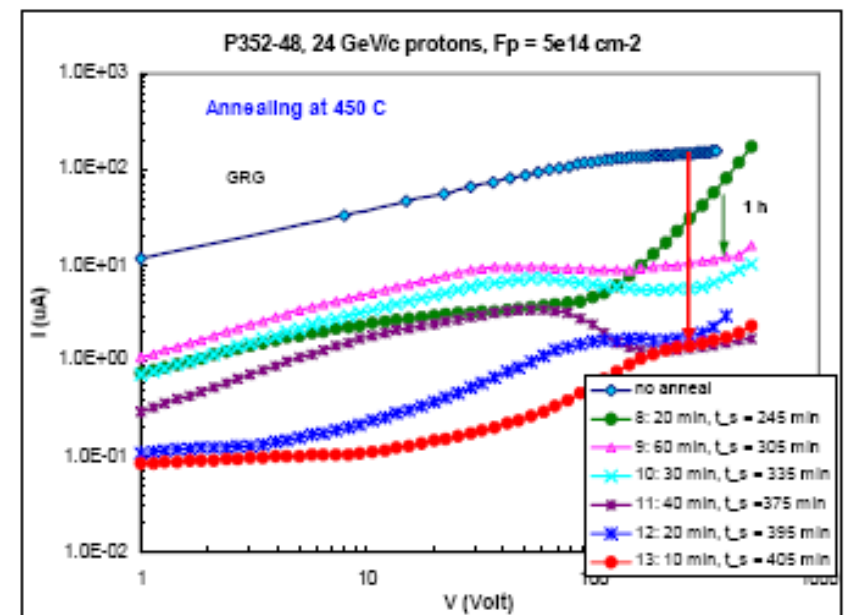
- I-V and C-V dependences
- a current pulse response using TCT with a laser pulse generation of non-equilibrium carriers

□ After final annealing (at Ioffe Inst.):

- Spectra of deep levels (C-DLTS)

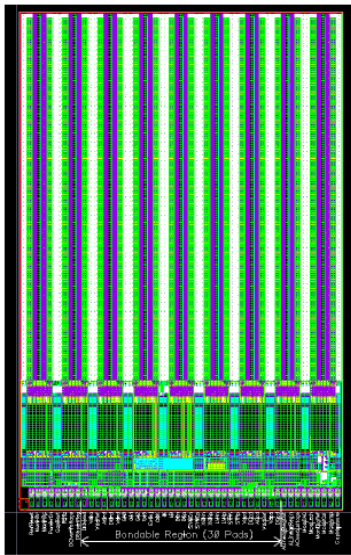
• Interesting results. For ex. :

- I_{leak} decreases after 1st anneal, but increases at high biases
- After $t_{ann} = 305$ min: I_{leak} becomes saturated
- After final anneal, I_{leak} decreased by more than 2 orders of magnitude

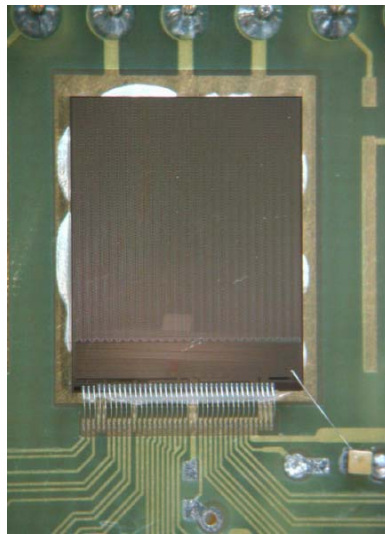


Sub-micron technology

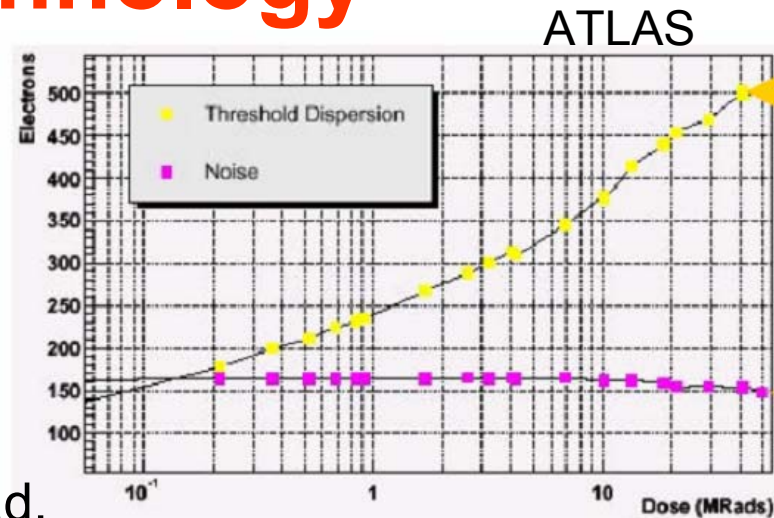
- LHC pixel development:
 - 5K channels per cm^2
⇒ unprecedented level of integration
 - 40 MHz operation at power density below $0.5\text{W}/\text{cm}^2$
 - radiation hardness of 50MRad, and high SEU-tolerance.



Atlas Chip



CMS Chip



Increase in threshold dispersion (re-tuning necessary)

No noise change

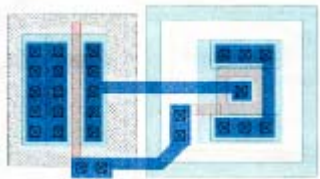
- Technology was initially unavailable. Adequate performance achieved in rad-soft $0.8\mu\text{m}$ CMOS, but transfer to rad-hard failed to meet requirements.
- The arrival of $0.25\mu\text{m}$ CMOS with rad-tolerant layout rules was key to development of the current ROCS.
- Working with mainstream commercial vendor (IBM) ⇒ average yield of $\sim 80\%$.

Sub-micron technology

- Both ATLAS and CMS are investigating IBM 0.13 μ m CMOS8RF electronics:
 - First measurements show that 0.13 μ m CMOS appears to be more radiation hard than 0.25 μ m CMOS.
 - VT shifts at about 70MRad are reduced by factors of 5 or more (X-ray studies).
 - Cost of runs is high: 0.25 μ m engineering run is \approx 150K\$, 0.13 μ m is \approx 600K\$
- CMS/ATLAS common chip development??



Standard



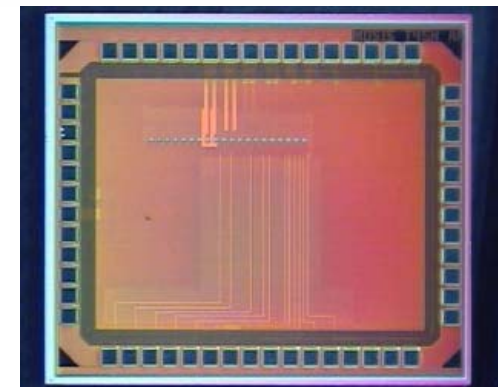
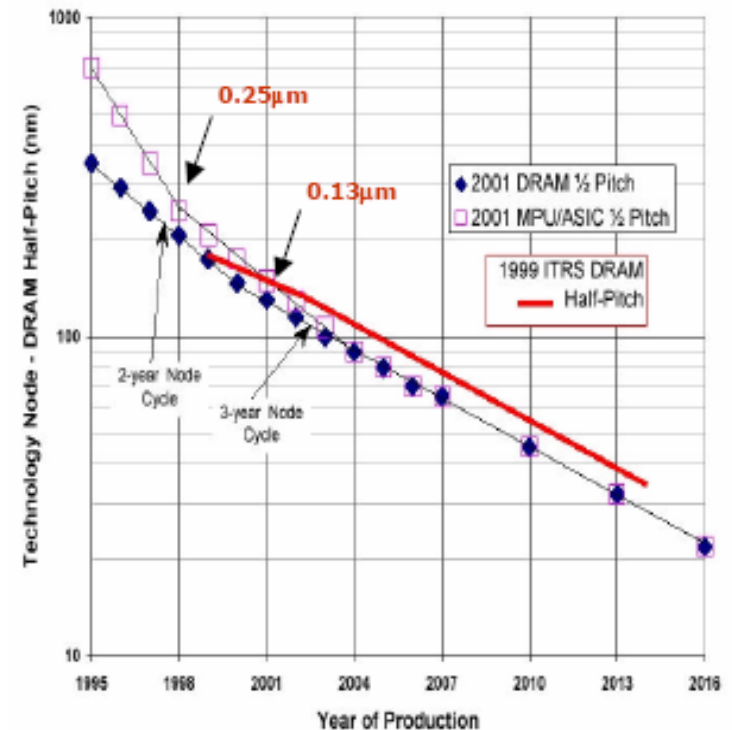
ELT

LHC luminosity upgrade

- Can we fully use gain of 0.13 μ m?
Can we waive some of the special design rules ?
- In 0.25 μ m DSM size penalty due to Enclosed Layout Transistor (ELT). Size penalty, reduces sensitivity to SEU.
- ATLAS Designed test chip in IBM 0.13 μ m CMOS technology to addresses performance of SEU-tolerant storage schemes.

Daniela Bortoletto

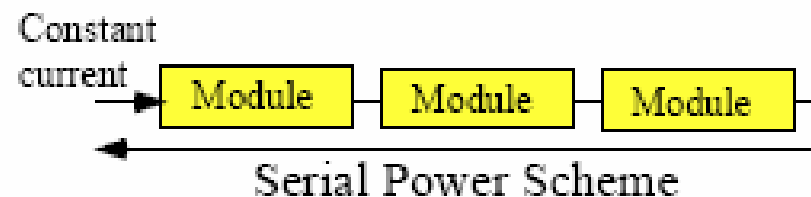
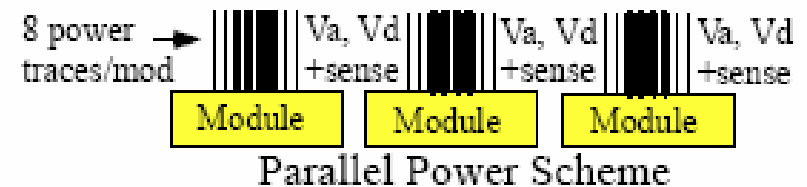
ITRS Roadmap Acceleration Continues...Half Pitch



Other major issues

- Uncertainty in the crossing frequency: (25ns, 15 ns, 10 ns, and 12.5ns) \Rightarrow Critical for ROC
 - CMOS logic power is \propto to CV^2f
 - The power supply voltage decreases from 2.5 to 1.5 V going from 0.25 μm to 0.13 μm \Rightarrow power reduction of 2.77 at fixed f
 - The gate capacitance/unit area goes up, the gate area goes down \Rightarrow decrease in C \approx 1.73 for a given complexity
 - The frequency might increase limiting the power savings.
 - Power in analog section is $\propto I_{\text{rms}} \times V$
 - Decrease in V reduces the power by 1.66 at constant I
 - The current in the analog section may actually be increased to compensate for lower dynamic range limiting power saving

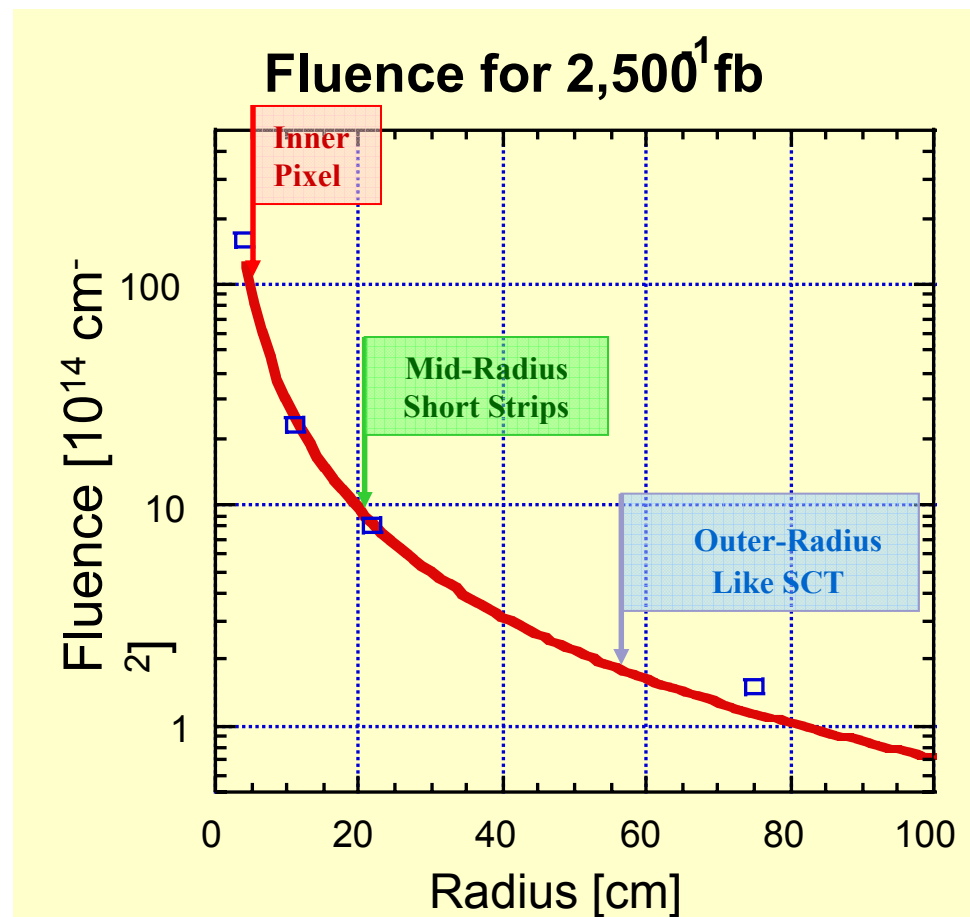
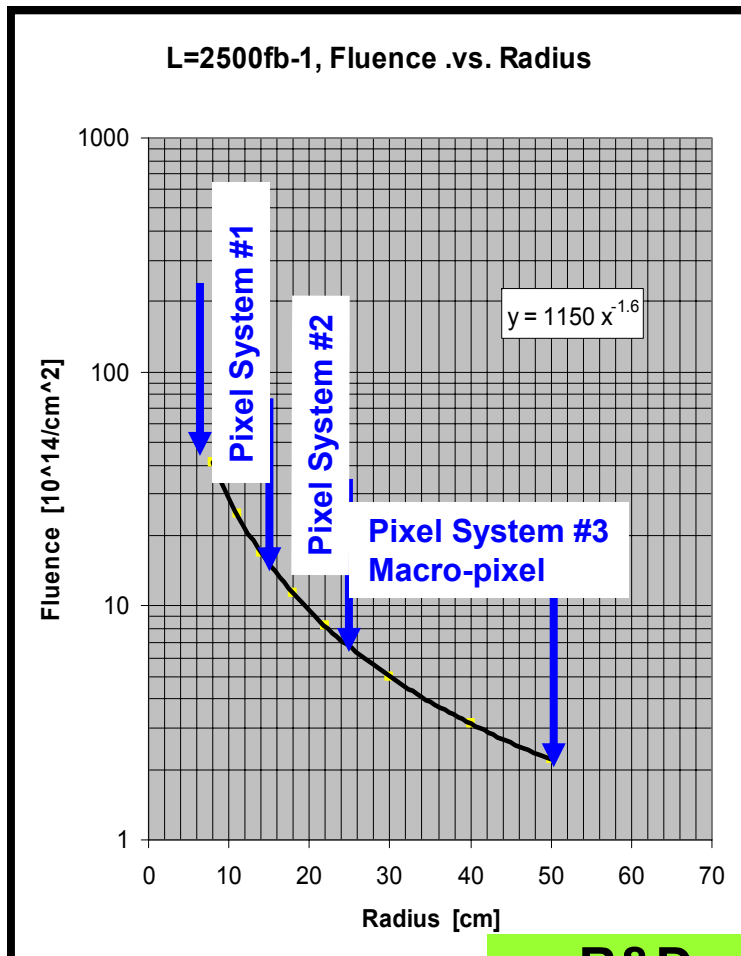
- The requirement of more granularity \Rightarrow increase in channel counts \Rightarrow cooling \Rightarrow more mass
- Serial powering of high density detectors can offer: reduced number of power cables, reduce mass, reduced power dissipation



- Large systems are hard to build
 - Qualification must be taken seriously
 - Many complex system issues

Detector specific developments

- CMS and Atlas are starting look at detector configurations:



CMS-Horisberger

LHC luminosity upgrade

- R&D needed below <20 cm
- Cost issues everywhere

Summary

- At fluences up to 10^{15}cm^{-2} (Outer layers of a SLHC detector) the change of depletion voltage and the large area is the major problem:
 - **CZ silicon detectors** could be a cost-effective radiation hard solution (no type inversion, use p-in-n technology)
 - **p-type silicon** microstrip detectors show very encouraging results: $\text{CCE} \approx 6500 \text{ e}$; $\Phi_{\text{eq}} = 4 \times 10^{15} \text{ cm}^{-2}$, $300\mu\text{m}$, collection of electrons no reverse annealing observed in CCE measurement!
- At the fluence of 10^{16}cm^{-2} (Innermost layer of a SLHC detector) the active thickness is significantly reduced due to trapping. Options are:
 - **Thin/EPI detectors**
 - **3D detectors, TFA**
 - **Diamond**
- Performance of 0.13 DSM electronics is promising
- Integration issue (power, cooling and mass) are complex. LHC experience will be very useful to avoid mistakes



**Radiation hard
electronics for
low signals**





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