

CALORIMETRY

The challenges of future experiments

Richard WIGMANS

CERN, March 15, 2006

Outline:

- A brief history
- Challenges of future experiments (LHC/SLHC, ILC/CLIC)
- Possible solutions / Lessons from the past
- Recent R&D results

A brief history

of design philosophy in particle physics experiments

- $\lesssim 1970$: Bubble Chamber Era : *Target \equiv detector*
- 1965 - 1985: “Electronic Bubble Chambers” (*fixed-target exp.*)
 - Momenta of charged particles: Tracking in magnetic field*
 - Particle ID: TOF, dE/dx , Č counters*
 - Muon ID: Absorber*
 - Photons: “Shower counters”*
- 1980 - now: Calorimeter system cornerstone (*collider exp.*)
 - Strong emphasis on electromagnetic calorimetry (exception: HERA)*
 - Magnetic field optional (UA2, D0)*
- $\gtrsim 2000$: Return to the past? Particle Flow Analysis
 - Use tracking, particle ID and calorimeter information to measure 4-vectors of jets*

Why calorimetry?

- Measure *charged + neutral* particles
- Obtain information on *energy flow*:
Total (missing) transverse energy, jets, *etc.*
- Obtain information *fast*
→ recognize and select interesting events in real time (*trigger*)
- Performance of calorimeters *improves with energy*
($\sim E^{-1/2}$ if statistical processes are the limiting factor)

Important calorimeter features

- Energy resolution
- Position resolution (need 4-vectors for physics)
- Signal speed
- Particle ID capability

The importance of (hadronic) energy resolution

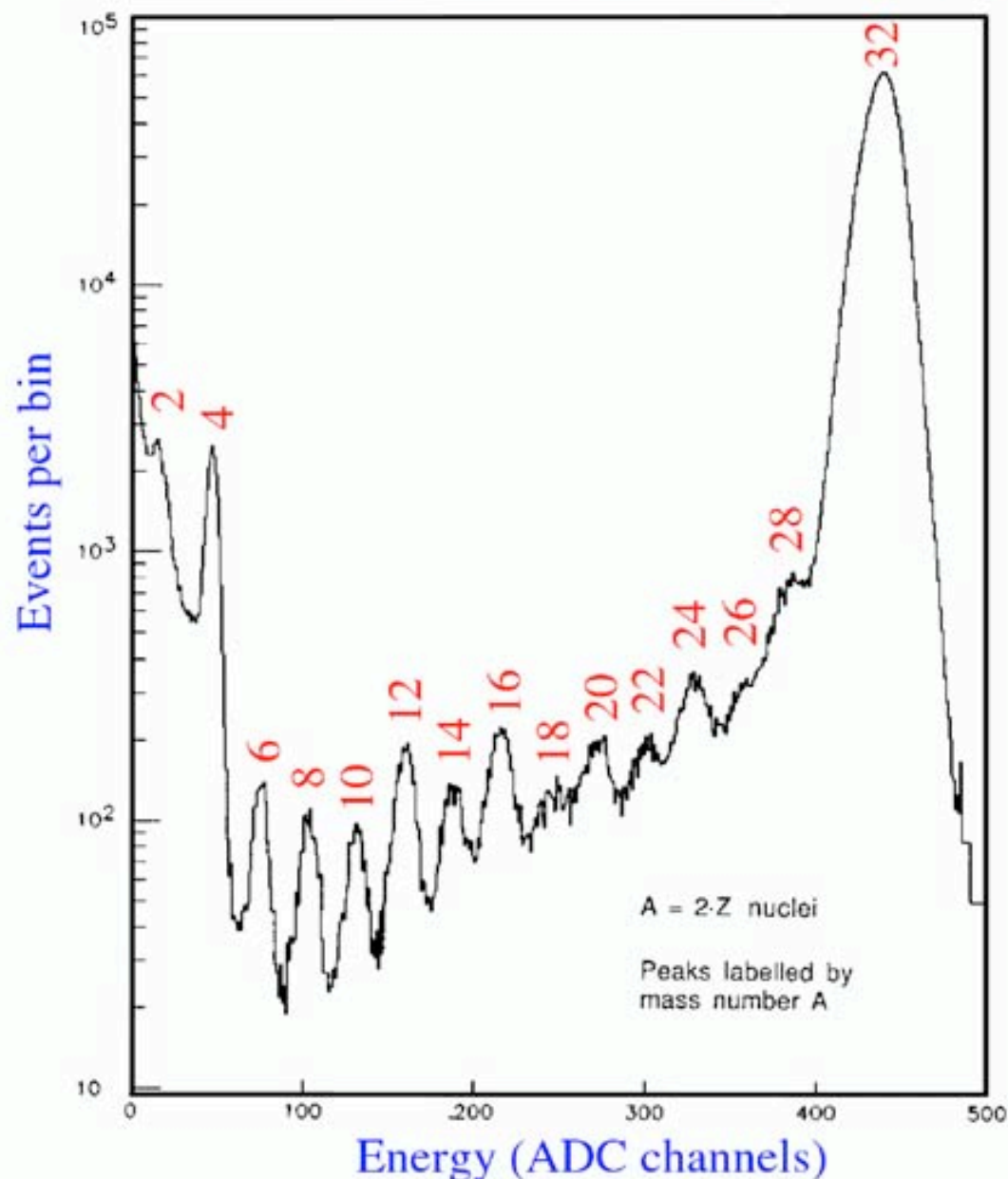


FIG. 7.51. The WA80 calorimeter as a high-resolution spectrometer. Total energy measured with the calorimeter for minimum-bias events revealed the composition of the momentum-selected CERN heavy-ion beam [You 89].

The importance of (hadronic) energy resolution (2)

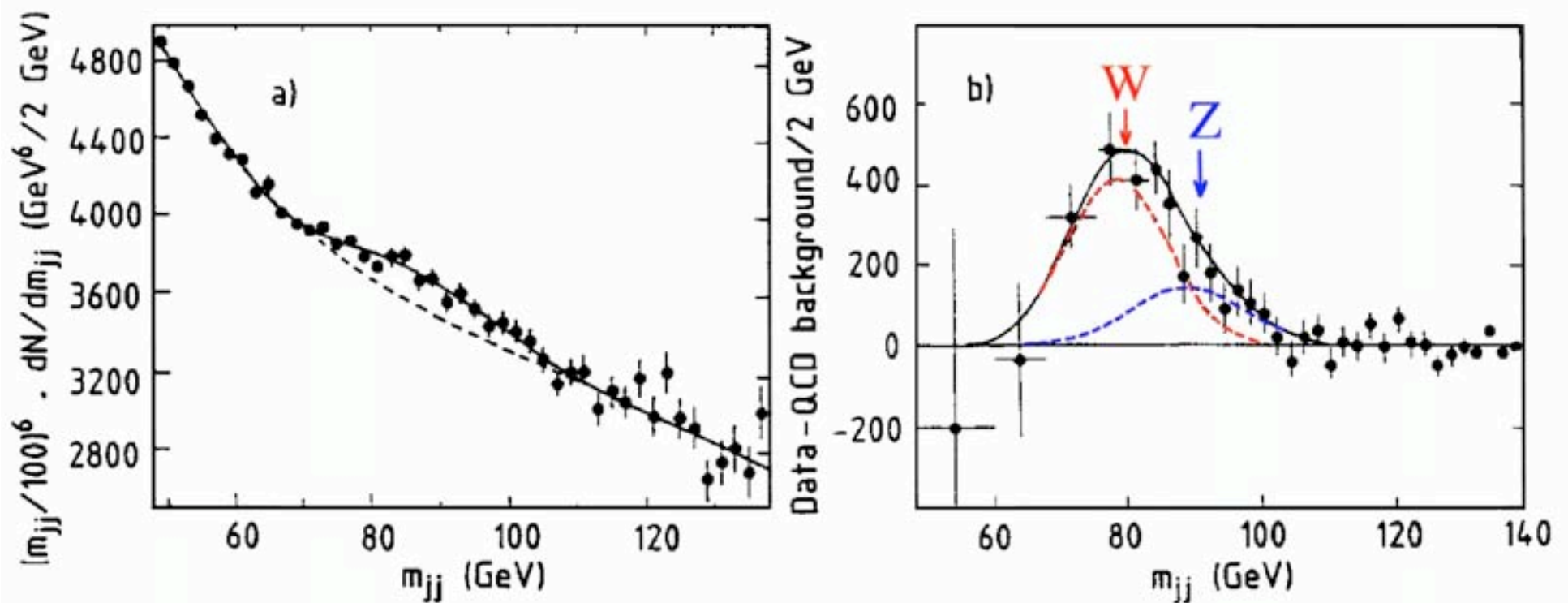


FIG. 7.50. Two-jet invariant mass distributions from the UA2 experiment [Alit 91]. Diagram a) shows the measured data points, together with the results of the best fits to the QCD background alone (*dashed curve*), or including the sum of two Gaussian functions describing $W, Z \rightarrow q\bar{q}$ decays. Diagram b) shows the same data after subtracting the QCD background. The data are compatible with peaks at $m_W = 80$ GeV and $m_Z = 90$ GeV. The measured width of the bump, or rather the standard deviation of the mass distribution, was 8 GeV, of which 5 GeV could be attributed to non-ideal calorimeter performance [Jen 88].

The importance of signal speed

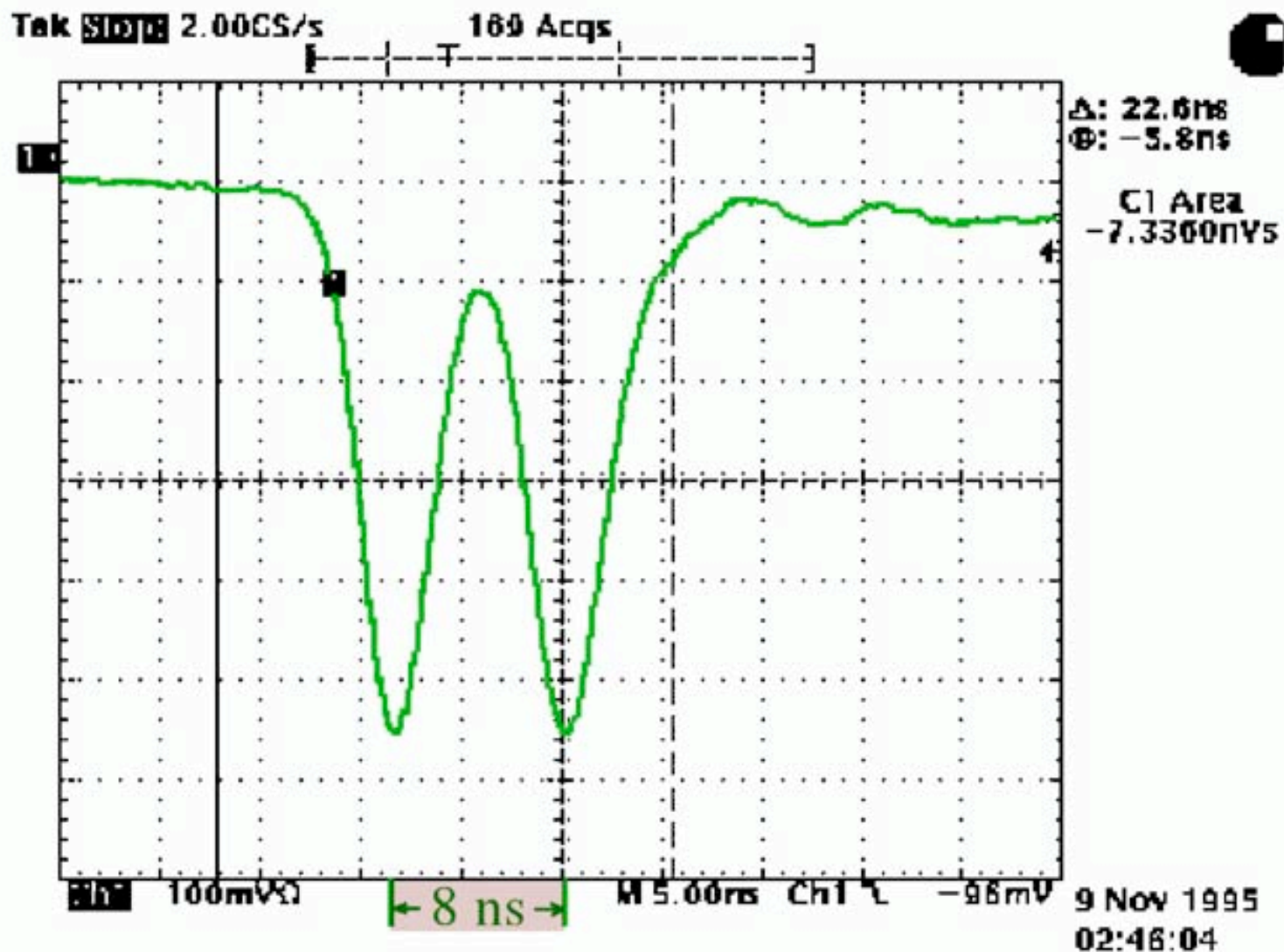


FIG. 7.19. Oscilloscope picture of two events separated by 8 ns in the Zero Degree Quartz Fiber Calorimeter of the NA50 experiment in the CERN heavy-ion beam [Arn 98].

Particle ID does *NOT* require segmentation!

e/π separation using time structure signals

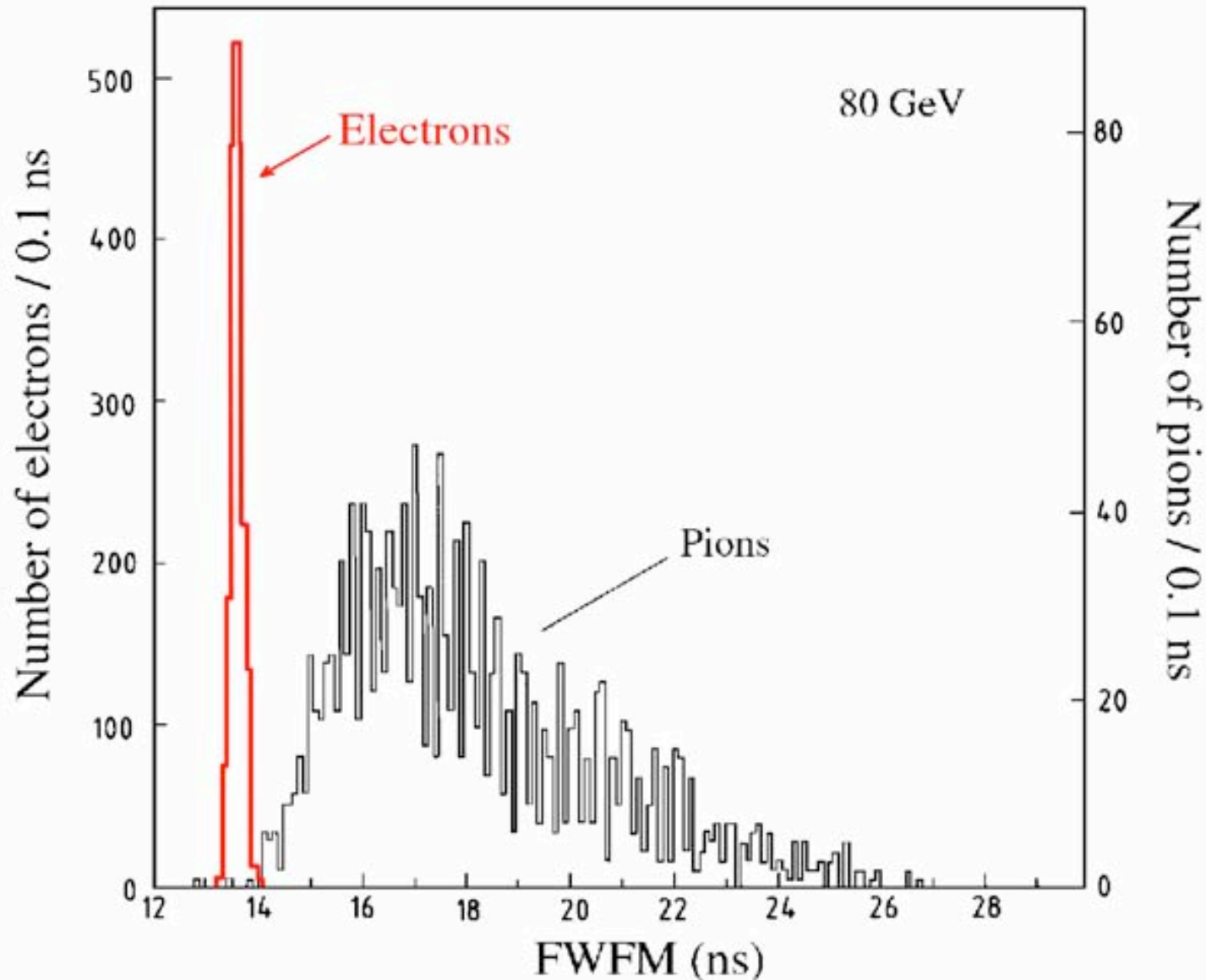


FIG. 7.33. The distribution of the full width at one-fifth maximum (FWFM) for 80 GeV electron and pion signals in SPACAL [Aco 91a].

Important calorimeter features

- Energy resolution
- Position resolution (need 4-vectors for physics)
- Signal speed
- Particle ID capability

but also

- *Gaussian response function* (avoid bias for steeply falling distributions)
- *Signal linearity*, or at least
- Well known relationship between signal & energy (*reliable calibration*)

Most hadron calorimeters fall short in this respect

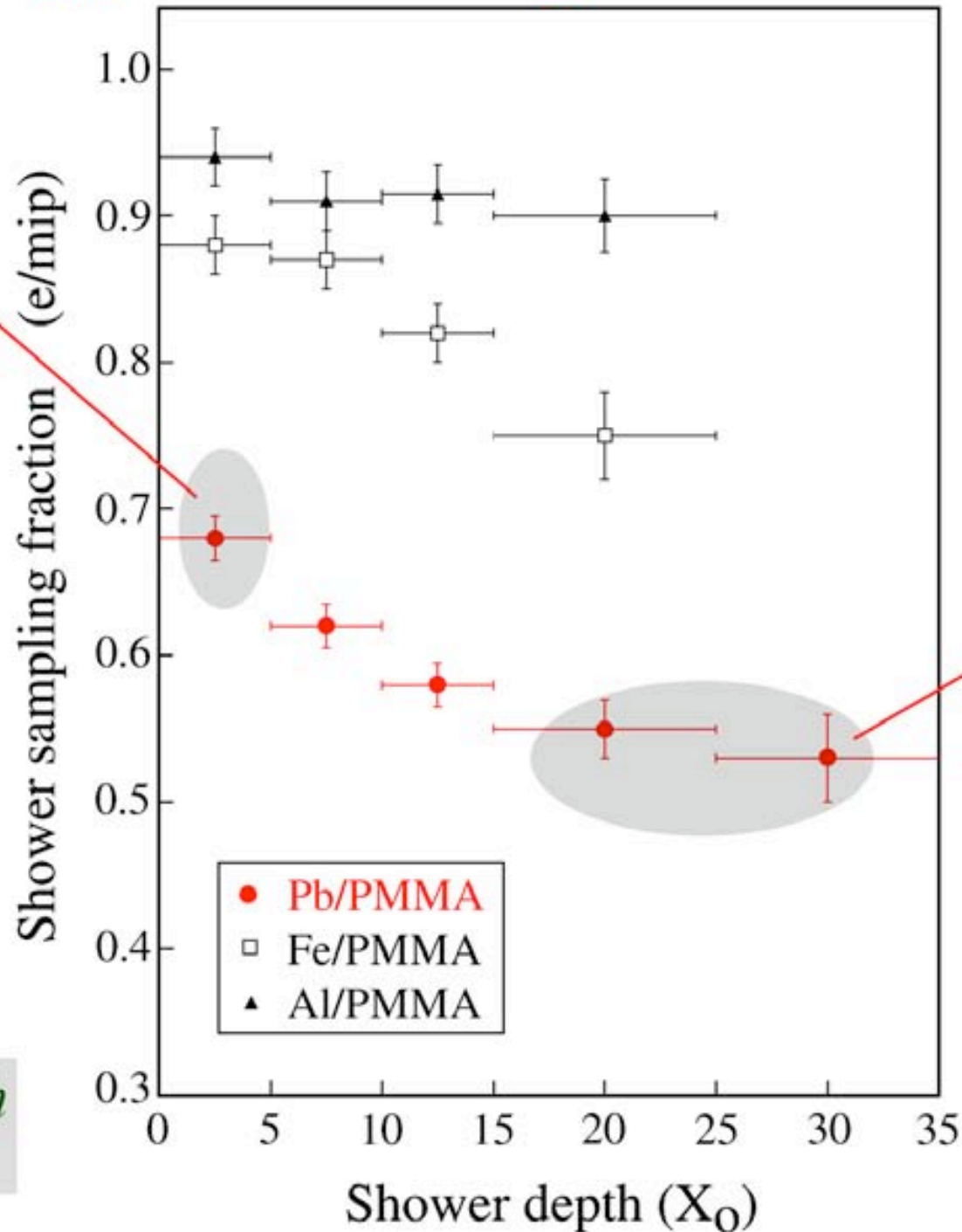
Calibration of calorimeter systems

- Determine relationship between *signal* (pC, p.e.) and *energy* (GeV)
- *Fundamental problem in sampling calorimeters:*
Different shower components are sampled differently
Shower composition changes as shower develops
→ *Sampling fraction changes with the shower age* (also E dependent)

*How to intercalibrate the sections
of a longitudinally segmented calorimeter?*

The sampling fraction changes as shower develops*

shower dominated by mip's

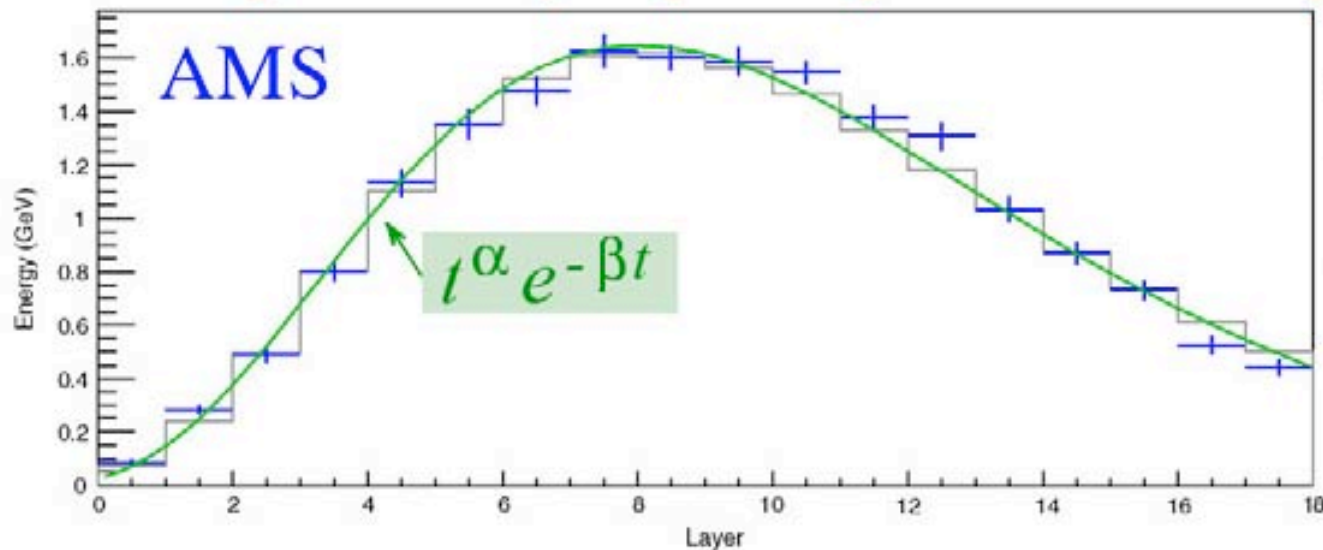


shower dominated by soft γ 's

*By as much as 30%!

CALIBRATION MISERY

Consequences of depth dependence sampling fraction



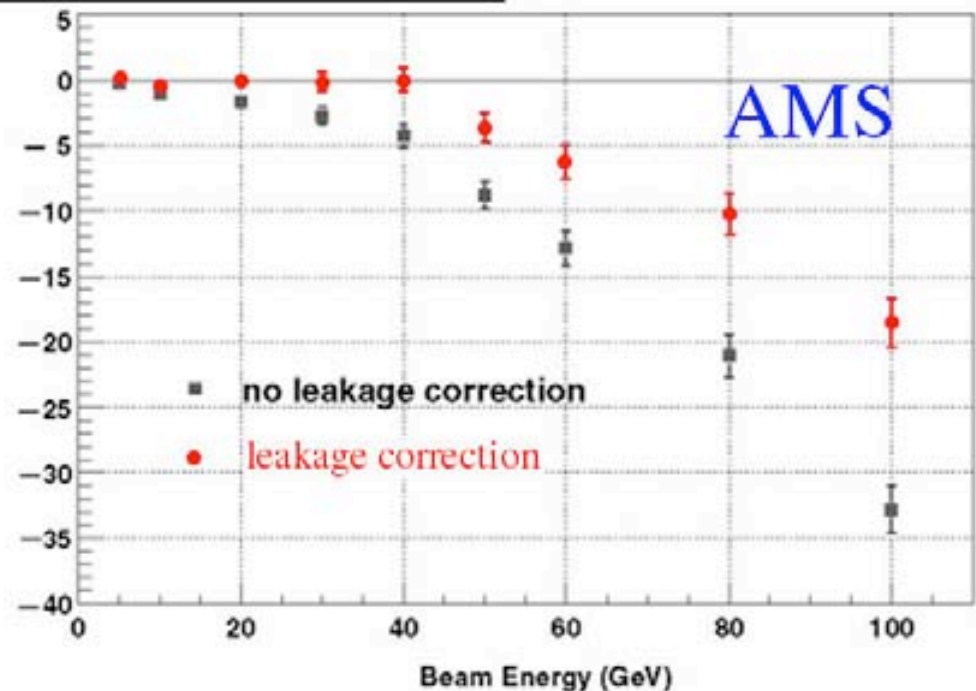
NIM A490, 132
(2002)

Pb/scintillating fiber
18 layers ($17 X_0$)

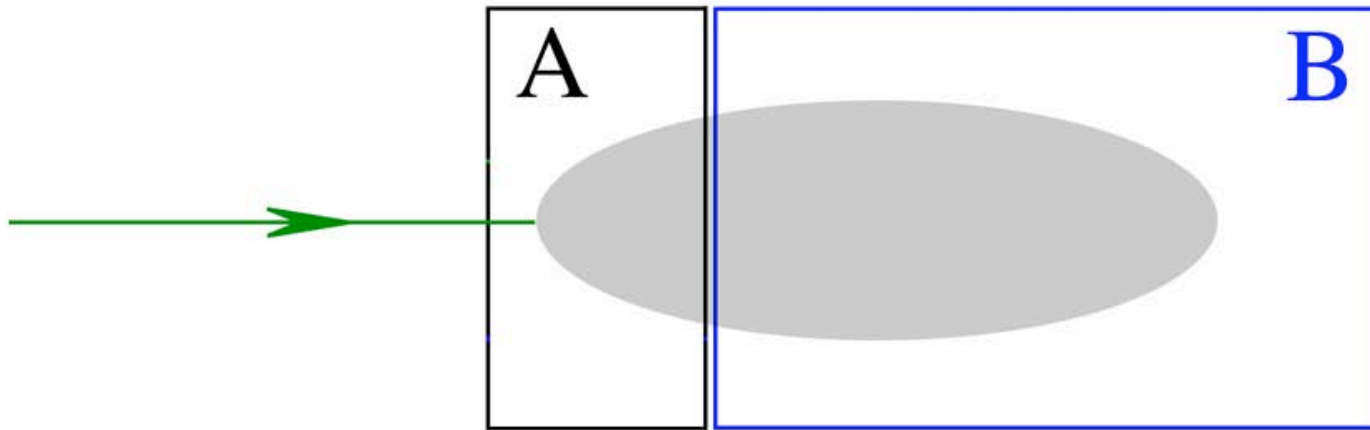
Calibrated with mip's:
11.7 MeV/layer

Shower leakage:
(under)estimated on basis
of fit to longitudinal profile

Measured energy-Beam Energy (%)



A widely used technique for calibrating segmented devices



Minimize
$$Q = \sum_{j=1}^N \left[E - A \sum_{i=1}^n S_{ij}^A - B \sum_{i=1}^n S_{ij}^B \right]^2$$

→ Determine A,B

Calibrating a segmented calorimeter by minimizing total signal width

GIVES WRONG RESULTS!

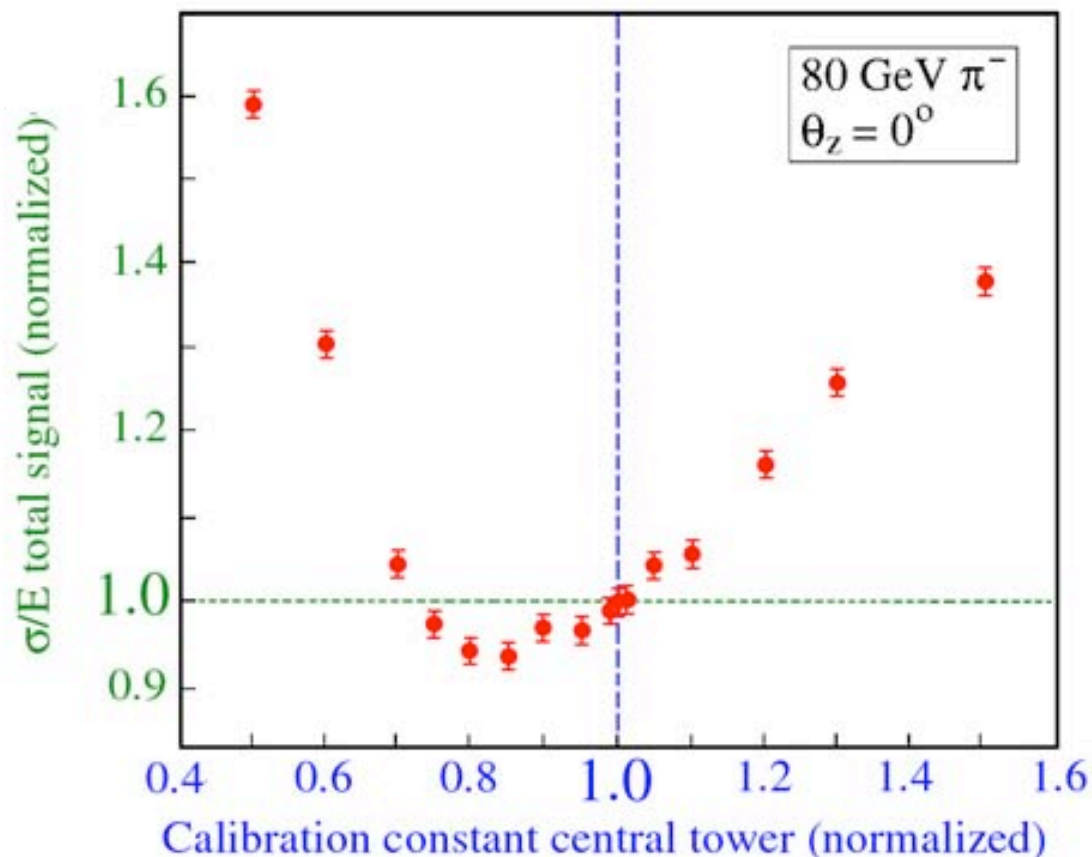


Figure 11: The fractional width, σ/E , of the signal distribution for 80 GeV π^- in the SPACAL detector as a function of the weighting factor applied to signals from the central calorimeter tower into which the pion beam was steered. The calorimeter towers were calibrated with high-energy electrons [7].

Calibration of calorimeter systems

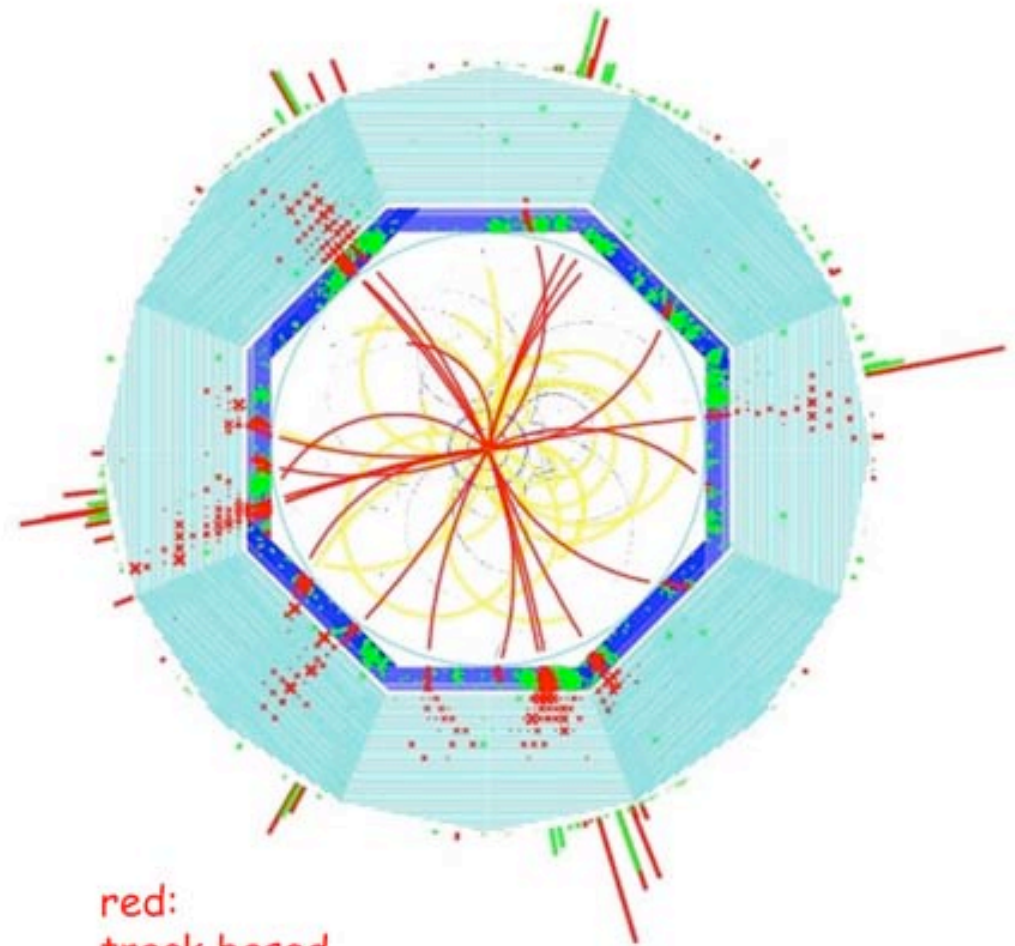
- Determine relationship between *signal* (pC, p.e.) and *energy* (GeV)
- *Fundamental problem in sampling calorimeters:*
Different shower components are sampled differently
Shower composition changes as shower develops
→ *Sampling fraction changes with the shower age* (also E dependent)

*How to intercalibrate the sections
of a longitudinally segmented calorimeter?*

- How do you know if the calibration is correct?
 - Electromagnetic: $Z, J/\psi \rightarrow e^+ e^-$, $\pi^0, \eta \rightarrow \gamma\gamma$
 - Hadronic: ??? $W, Z \rightarrow q\bar{q}$ 'γ-jet balancing' ???

Particle Flow Analysis (Energy Flow Method)

- Use *tracking, particle ID and calorimetry* to measure 4-vectors of jets
- Charged particles represent typically $\sim 65\%$ of the jet energy
However, if only charged jet components are measured:
 $(\sigma/E)_{jet} = 25-30\%$
(independent of jet energy)
 \rightarrow *Calorimetry essential!*
- The problem with this method is *shower overlap*. Need to deconvolute contributions from showering charged particles to avoid *double counting*



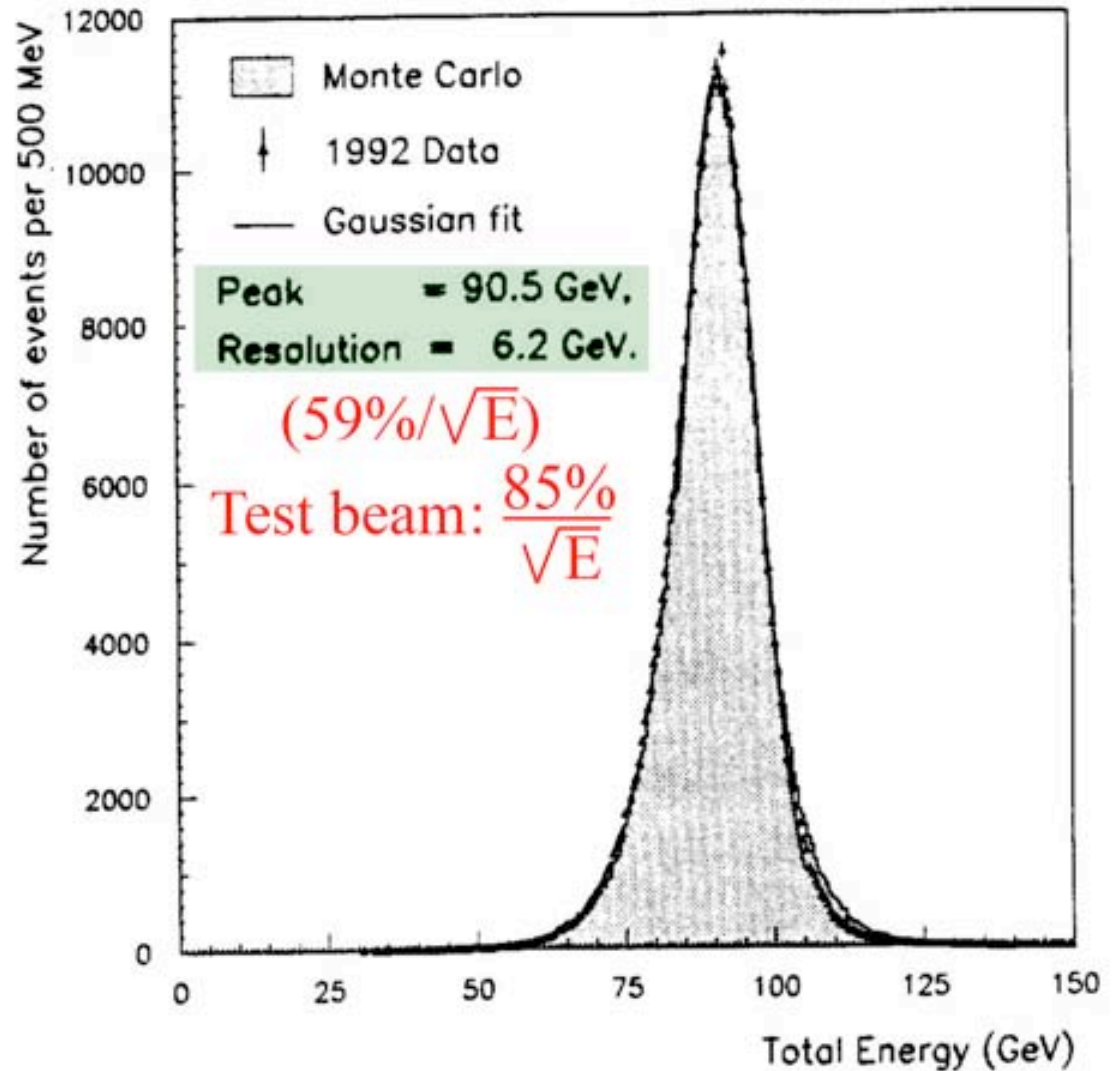
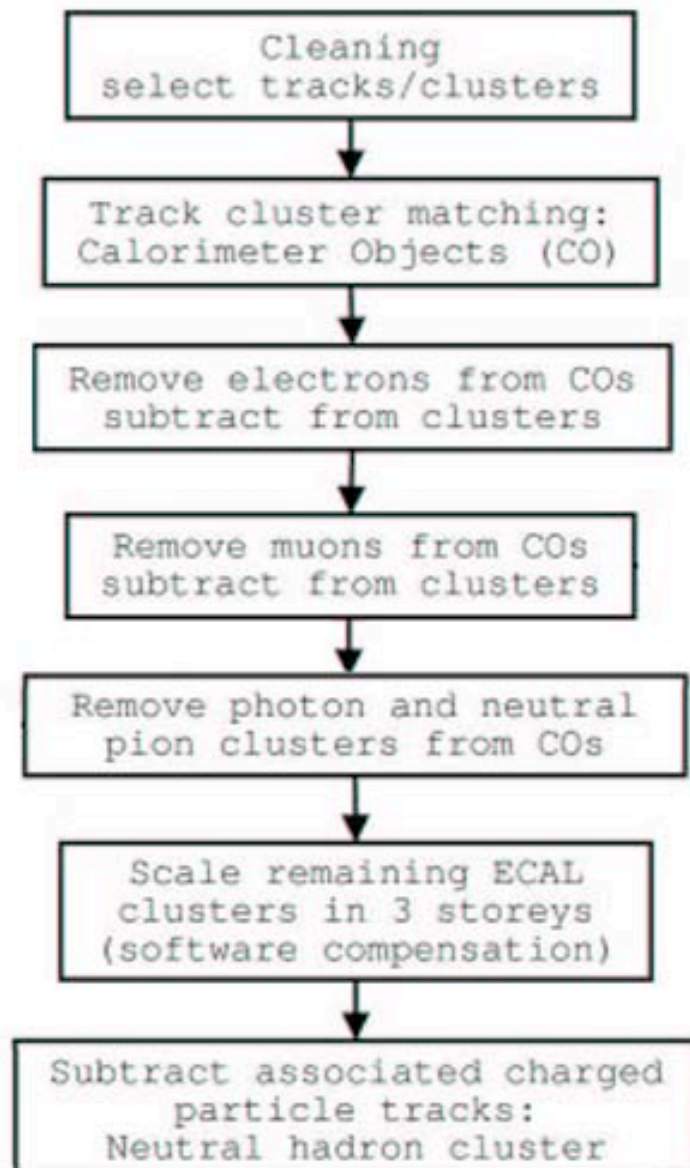
red:
track based

green:
calorimeter based

ZHH \rightarrow qqbbbb

Example of PFA at LEP: ALEPH

Attempt to reconstruct hadronic event structure using particle identification and software compensation



PFA at a Hadron Collider

(no kinematic constraints as at LEP)

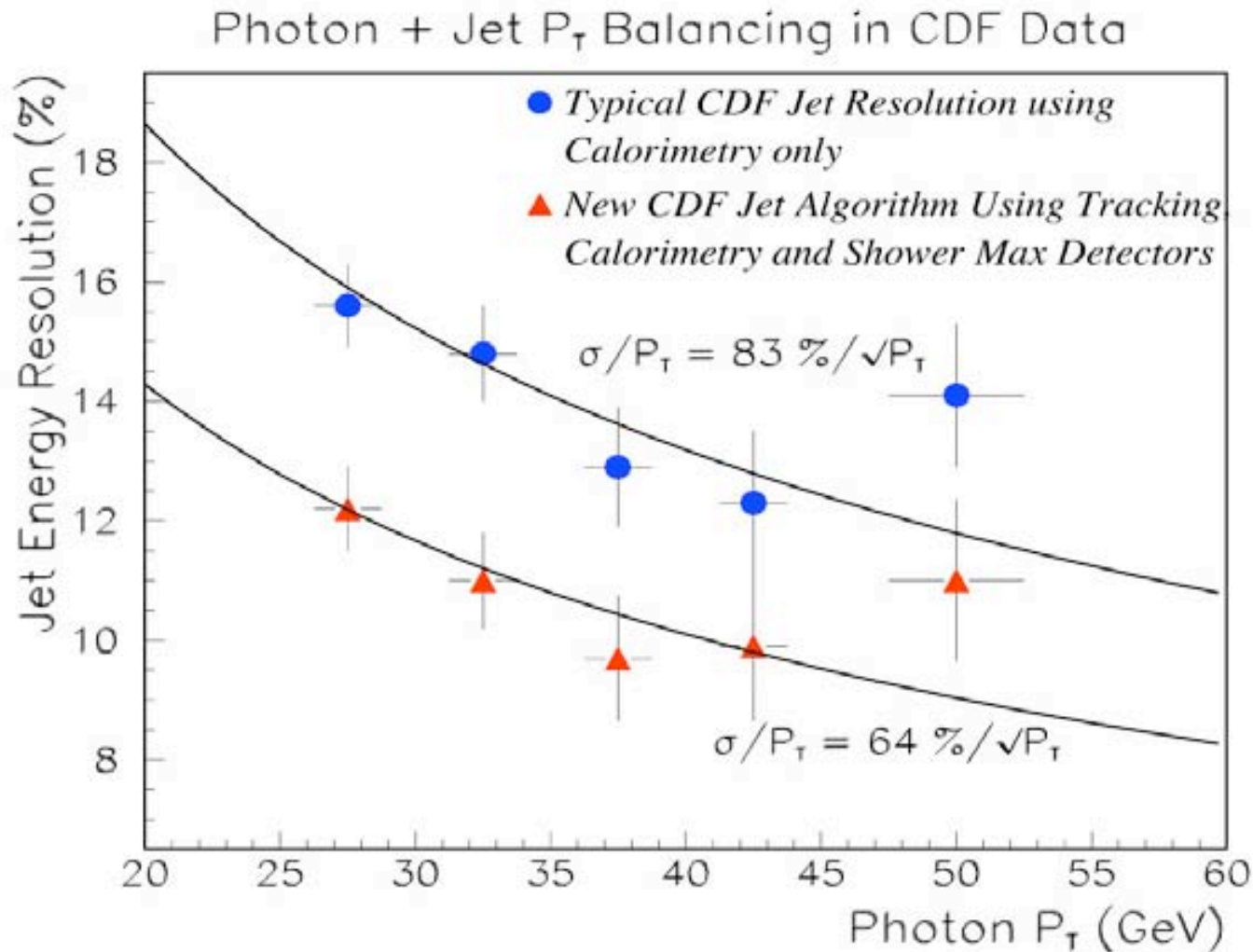


Figure 4: *The central detector resolution σ_D is plotted as a function of P_T^γ for the two methods.*

From: CDF note 5005 (2000)

Improvement of energy resolution expected with EFM (PFA)

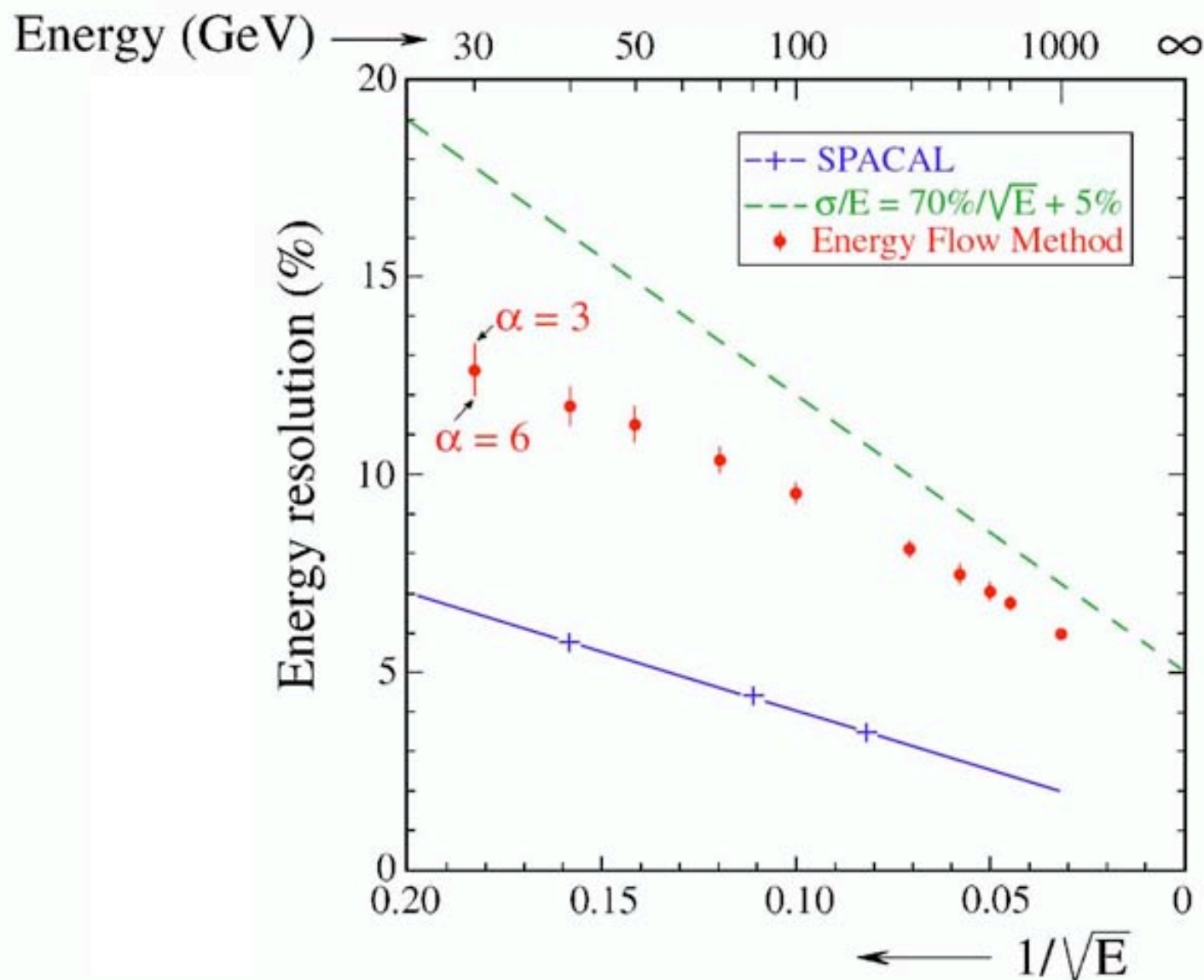


Fig. 11. The jet energy resolution as a function of energy, obtained after applying the Energy Flow Method (the dots), using simulated data from a calorimeter with a jet resolution given by the dashed curve. For comparison, the jet resolution of a compensating calorimeter is given (SPACAL [7]). [From: NIM A495 \(2002\) 107.](#)

Challenges in future experiments (1)

LHC (ATLAS, CMS)

- Em: Calorimeters designed to be sensitive to $H^0 \rightarrow \gamma\gamma$

Challenge: Maintain stability

- Hadrons/jets:

Challenge: The energy scale (segments with different e/h)

Recommendation: Approach calibration in a very systematic way, understood by all users, do not rely on fudge factors. Well documented.

Example: NIM A487 (2002) 381

Challenges in future experiments (2)

SLHC (ATLAS, CMS)

Assuming that bunch spacing does not become smaller:

Effects of increased number of underlying events: Trivial extrapolation

*Main challenge: **The effects of radiation***

- *Radiation damage of scintillators and quartz (CMS) } affects performance*
H₂ build-up in LAr (ATLAS)
- *Activation of detector materials → access/service problematic*
- *Cooling of forward calorimeters.*

The calorie becomes a relevant unit for calorimeters: $1 \text{ cal} = 2.6 \cdot 10^{10} \text{ GeV}$

CMS: 100 GeV per min. bias event deposited in Forward Calorimeters

@ $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$: 12 J/s (W)

@ $10^{35} \text{ cm}^{-2} \text{ s}^{-1}$: 120 W

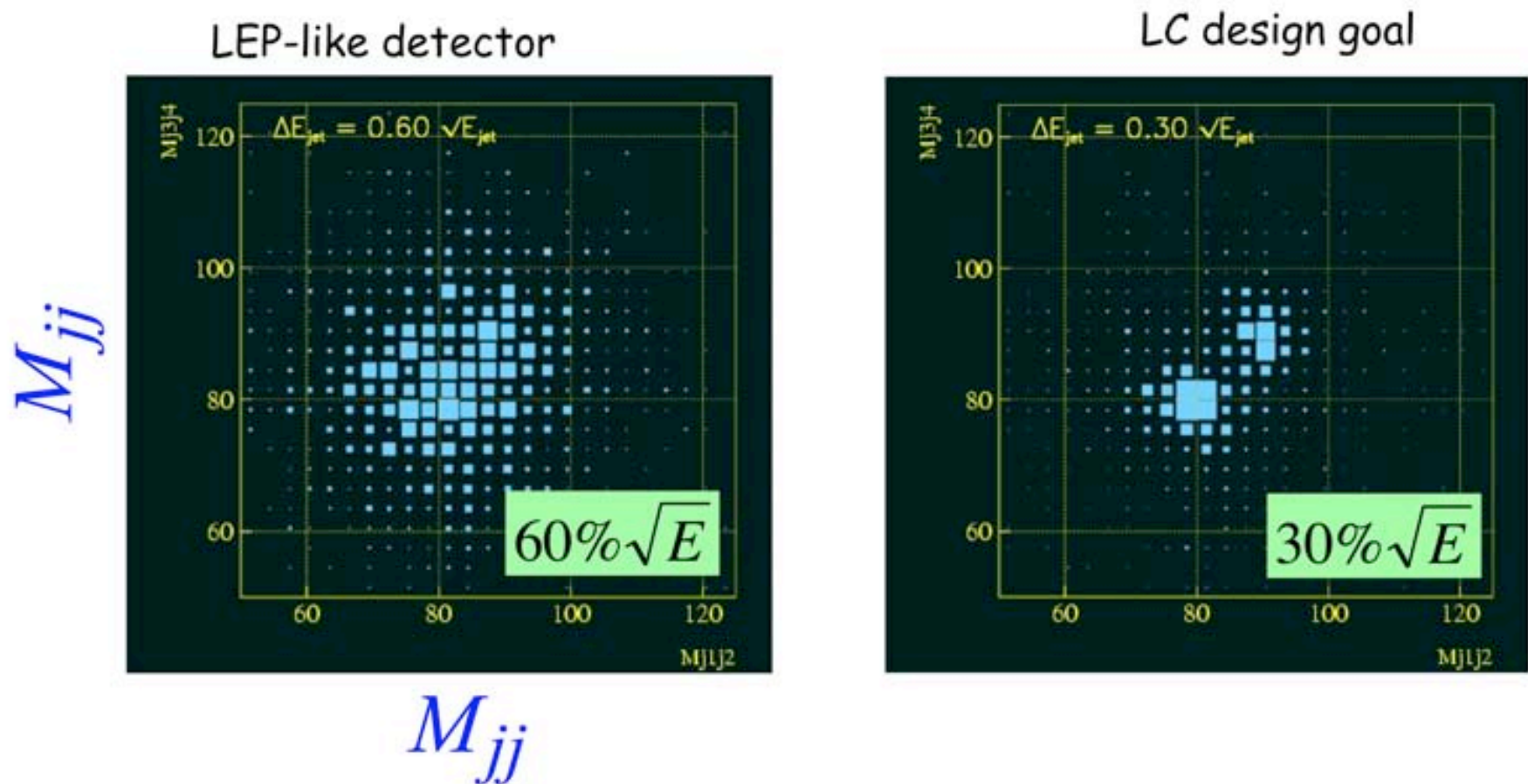
Challenges in future experiments (3)

ILC / CLIC

- *Need to measure 4-vectors of jets with excellent precision.*
Physics program relies heavily on final states with (several) bosons: W, Z, H
Necessary to reconstruct W, Z through their hadronic decay modes.
Hadronic energy resolution very important for this $multi\text{-jet spectroscopy}$.
- *The same argument can also be made for SLHC.*
For example, study of multi-boson couplings is statistics limited if one only would consider leptonically decaying W, Z .
→ SLHC physics program might benefit from improved hadron calorimetry
- *The issue of $H^0 \rightarrow \gamma\gamma$ will presumably be settled during LHC running .*
Therefore, it is conceivable to replace the calorimeter system by one with strongly improved hadronic performance for SLHC era .

Design goal ILC/CLIC: separate $W, Z \rightarrow q\bar{q}$

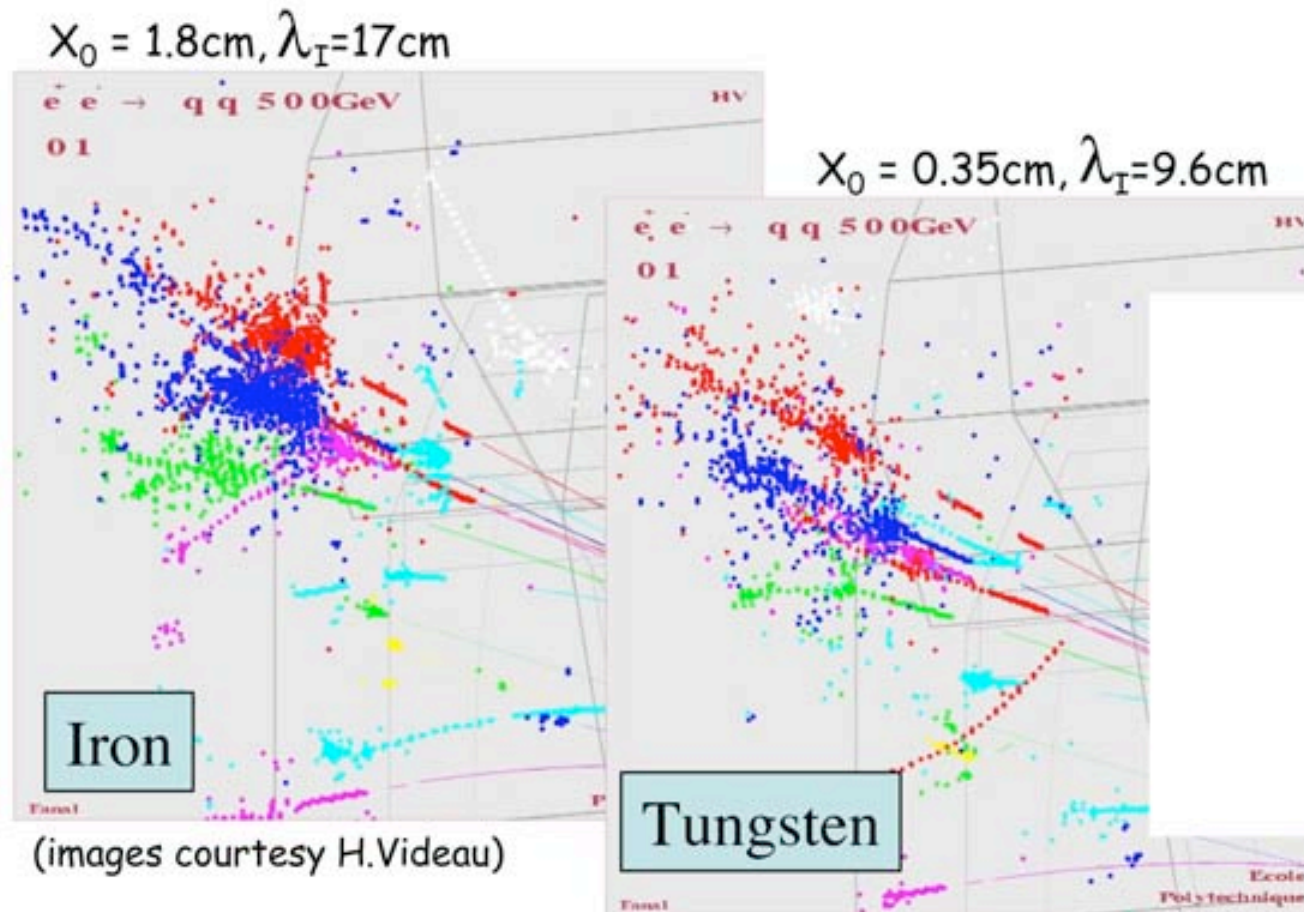
- Hadronic energy resolution very important for this *multi-jet spectroscopy*.



- No kinematic constraints** as in LEP (beamstrahlung)

PFA @ ILC: High Density + Fine Granularity

- In order to reduce the problems of shower overlap, ILC R&D focuses on *reducing the shower dimensions* and *decreasing the calorimeter cell size* (Alternatively, one could increase shower separation through \bar{B} , radius calorimeter)



- *How about calibration??*

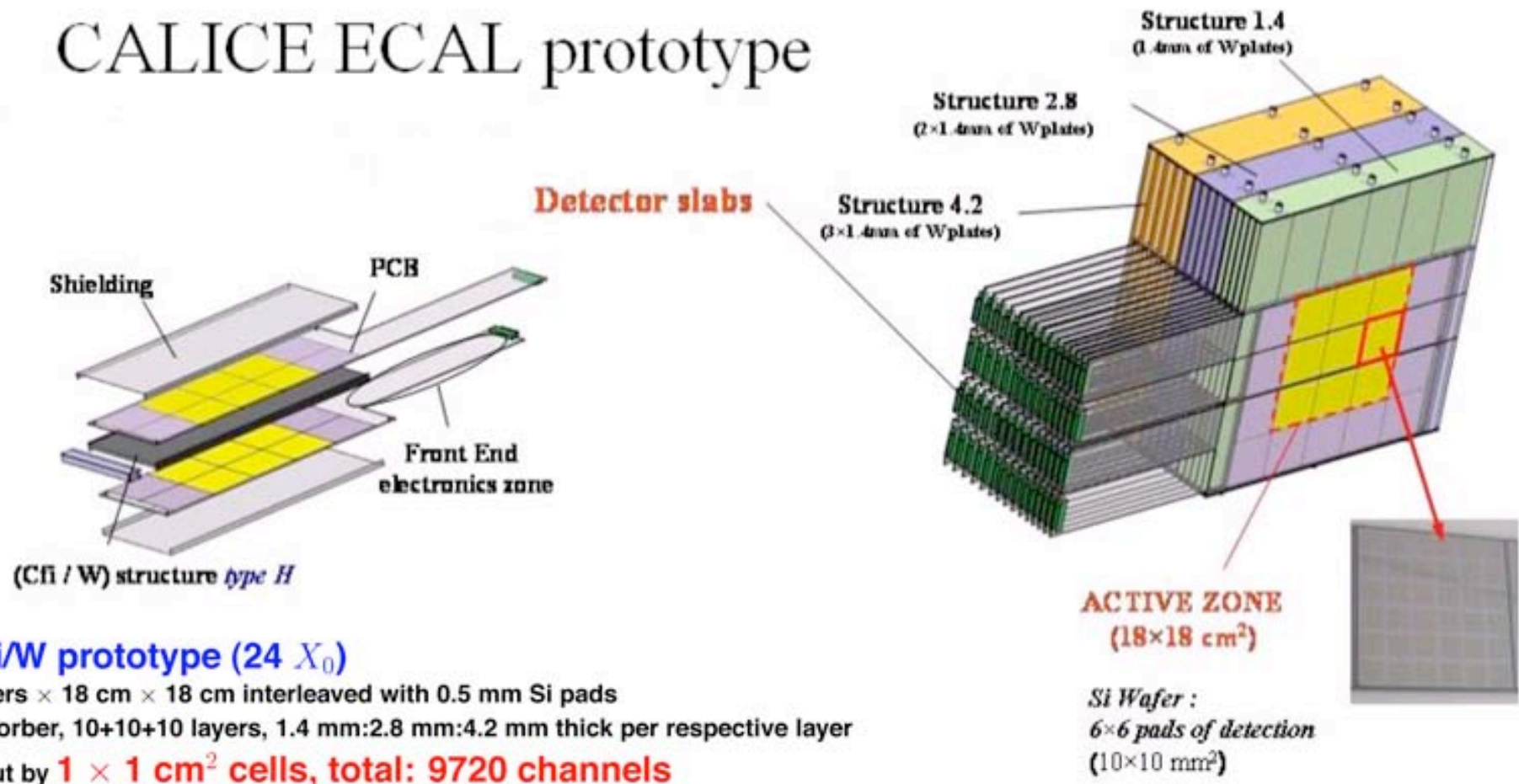
ILC R&D example (PFA): CALICE

164 Physicists, 28 Institutes, 9 Countries: 3 Regions



- ECAL and HCAL together, different options
- electron and hadron beams, start end 2004 at DESY

CALICE ECAL prototype



full Si/W prototype (24 X₀)

- ▷ 30 layers × 18 cm × 18 cm interleaved with 0.5 mm Si pads
- ▷ W absorber, 10+10+10 layers, 1.4 mm:2.8 mm:4.2 mm thick per respective layer
- ▷ readout by **1 × 1 cm² cells, total: 9720 channels**

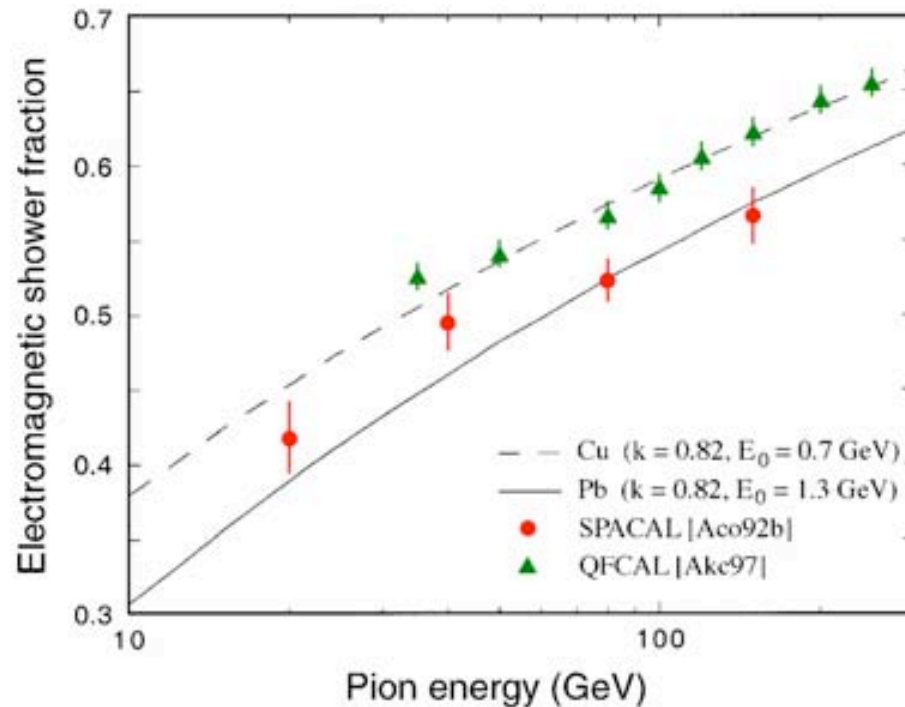
LESSONS FROM 25 YEARS OF R&D

- *LESSON 1:* Energy resolution is determined by *fluctuations*, **not** by average values
- *LESSON 2:* Digital calorimetry has been tried *and abandoned*, **for good reasons**
- *LESSON 3a:* A narrow signal distribution is useless if the mean value is incorrect
Correct energy scale is at least as important as good resolution
- *LESSON 3b:* *Longitudinal segmentation means asking for trouble*
- *LESSON 4:* GEANT based MC simulations of hadronic shower development are *fundamentally flawed* → **useless as design tool**
- *LESSON 5:* If you want to improve hadronic calorimeter performance
→ *reduce/eliminate the (effects of) fluctuations that dominate the performance:*
 - i)* Fluctuations in the em shower fraction, f_{em}
 - ii)* Fluctuations in visible energy (nuclear binding energy losses)

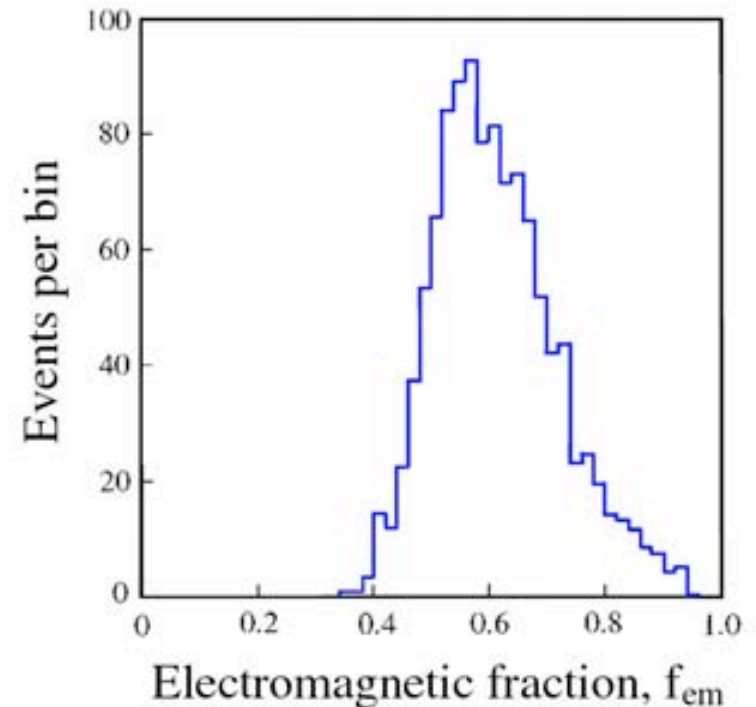
Fluctuations in the em shower component (f_{em})

- *Why are these important ?*
 - Electromagnetic calorimeter response \neq non-em response ($e/h \neq 1$)
 - Event-to-event fluctuations are large and *non-Gaussian*
 - $\langle f_{em} \rangle$ *depends on* shower *energy* and *age*
- *Cause of all common problems in hadron calorimeters*
 - *Energy scale* different from electrons, in energy-dependent way
 - Hadronic *non-linearity*
 - *Non-Gaussian* response function
 - Poor energy *resolution*
 - *Calibration* of the sections of a longitudinally segmented detector

(Fluctuations in) the electromagnetic shower fraction, f_{em}



The em fraction is, on average, *large and energy dependent*



Fluctuations in f_{em} are *large and non-Poissonian*

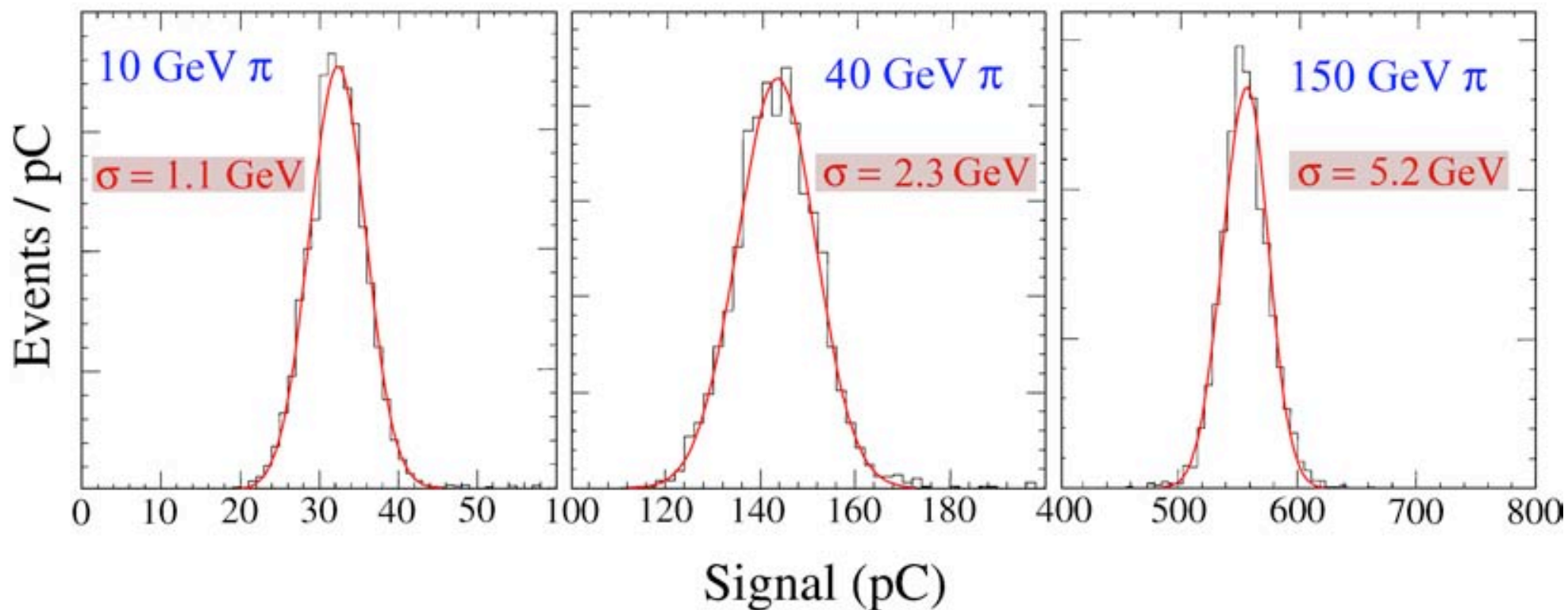
LESSONS FROM 25 YEARS OF R&D

- *LESSON 1:* Energy resolution is determined by *fluctuations*, **not** by average values
- *LESSON 2:* Digital calorimetry has been tried *and abandoned*, **for good reasons**
- *LESSON 3a:* A narrow signal distribution is useless if the mean value is incorrect
Correct energy scale is at least as important as good resolution
- *LESSON 3b:* *Longitudinal segmentation means asking for trouble*
- *LESSON 4:* GEANT based MC simulations of hadronic shower development are *fundamentally flawed* → **useless as design tool**
- *LESSON 5:* If you want to improve hadronic calorimeter performance
→ *reduce/eliminate the (effects of) fluctuations that dominate the performance :*
 - i) Fluctuations in the em shower graction, f_{em}*
 - ii) Fluctuations in visible energy (nuclear binding energy losses)*

This can be done.

ILC requirements were already met 15 years ago

Hadronic signal distributions in a compensating calorimeter



from: NIM A308 (1991) 481

Pros & Cons of Compensating Calorimeters

Pros

- Same *energy scale* for electrons, hadrons and jets. No ifs, ands or buts.
- *Calibrate* with electrons and you are done.
- Excellent hadronic *energy resolution* (SPACAL: $30\%/ \sqrt{E}$).
- *Linearity*, Gaussian *response function* and all that good stuff.
- Compensation fully understood.
We know how to build these things, even though GEANT doesn't

Cons

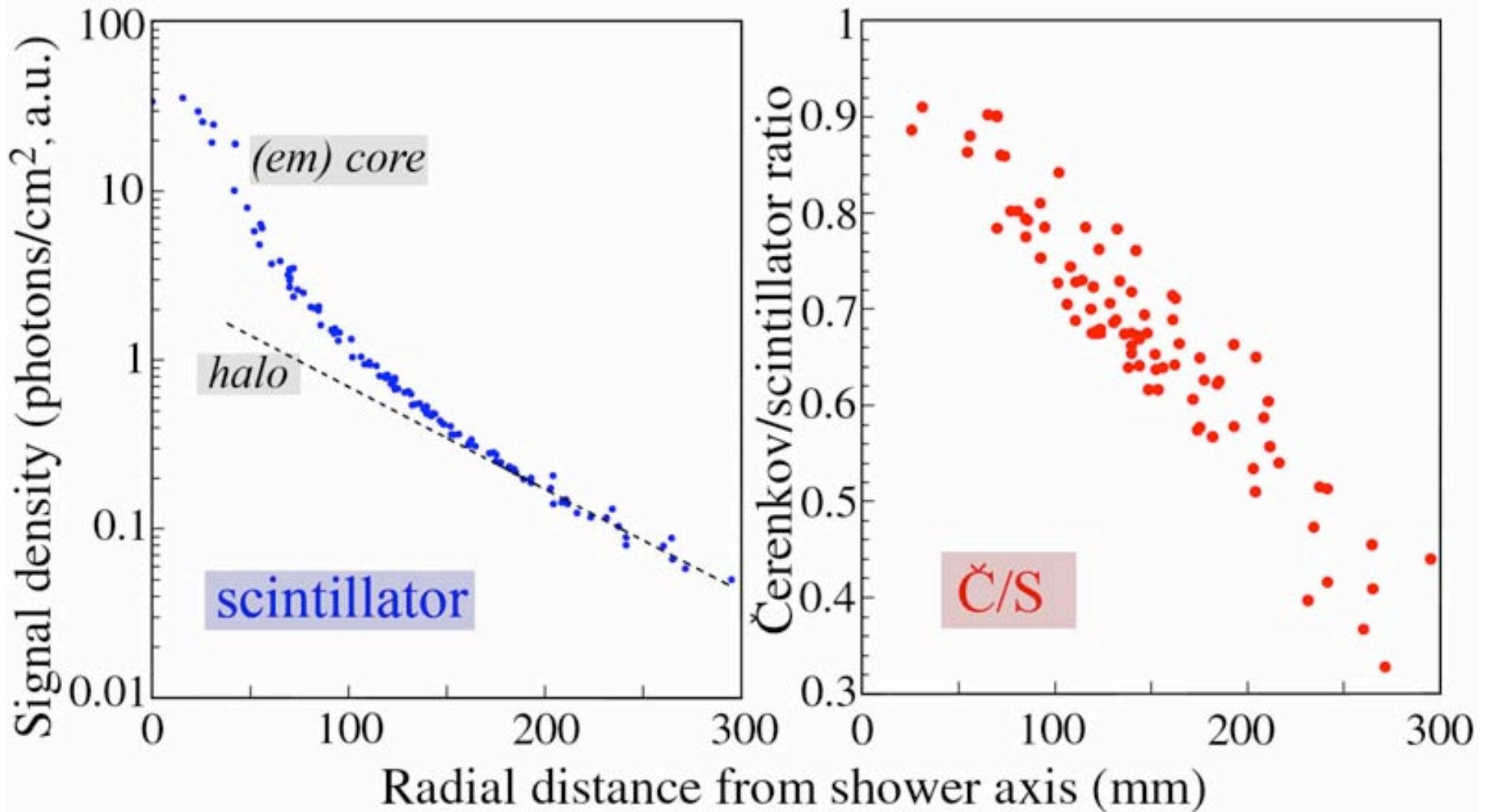
- Small sampling fraction (2.4% in Pb/plastic)
→ *em energy resolution limited* to $10\text{-}15\%/ \sqrt{E}$
- Compensation relies on detecting neutrons
→ Large *integration volume*
→ Long *integration time* (~ 50 ns)

The DREAM principle

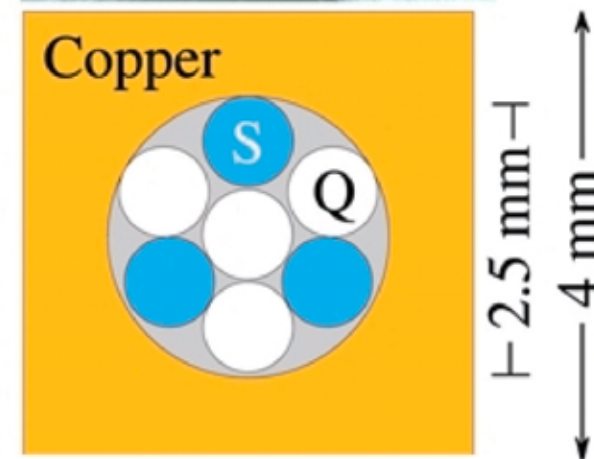
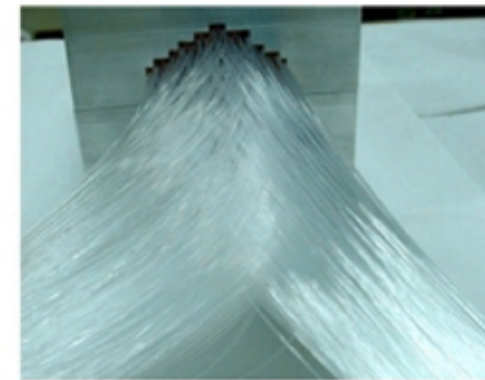
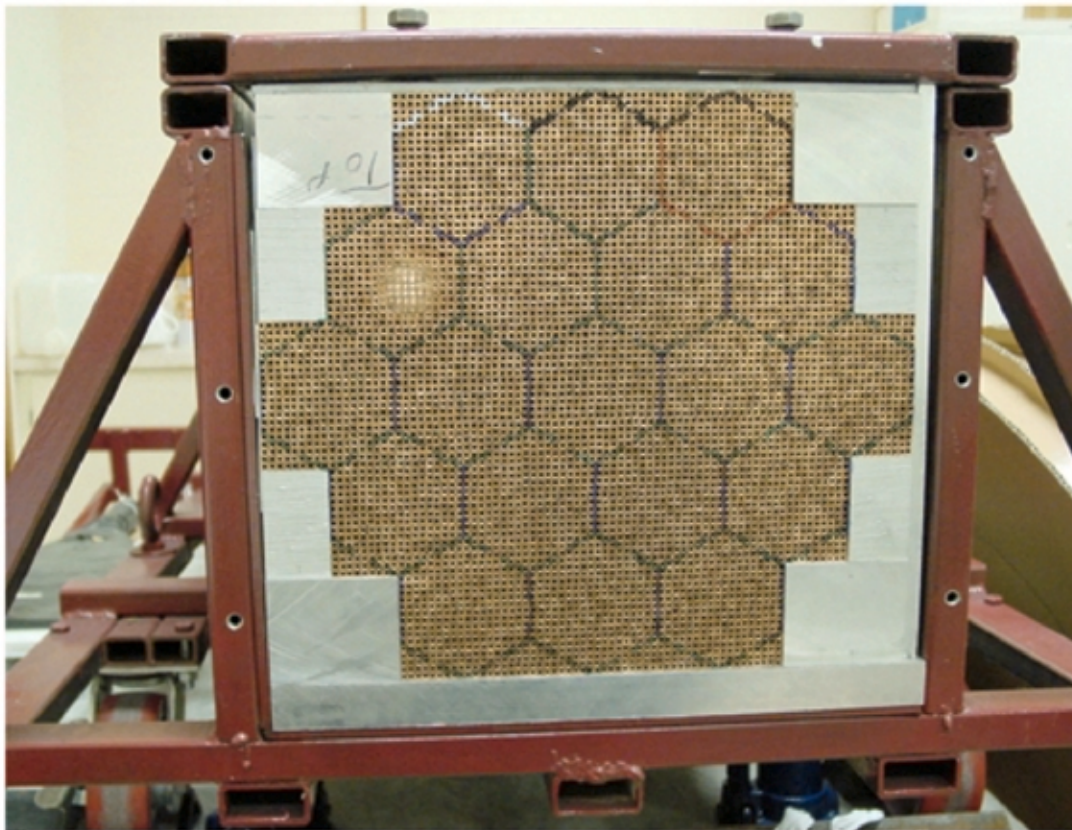
- *Quartz fibers are only sensitive to em shower component!*
 - CMS prototype: $e/h \sim 5$ NIM A399 (1997) 202
 - ➔ Use dual-readout system:
 - Regular readout (scintillator, LAr,...) measures *visible energy*
 - Quartz fibers measure **em shower component** E_{em}
 - Combining both results makes it possible to determine f_{em} and the energy E of the showering hadron
 - *Eliminate dominant source of fluctuations*

DREAM = Dual READout Module

Radial hadron shower profiles (DREAM)



DREAM: Structure

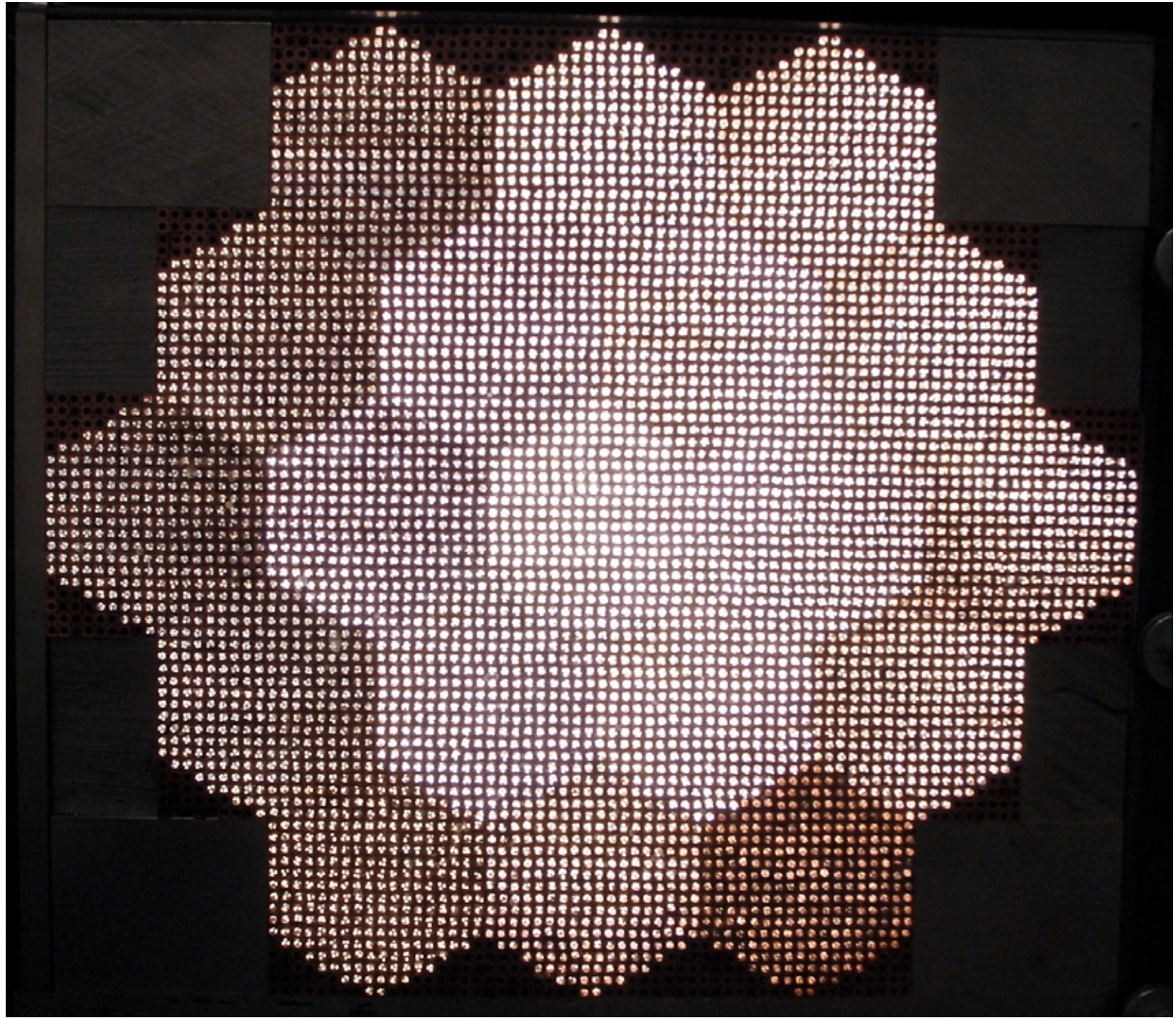


- *Some characteristics of the DREAM detector*

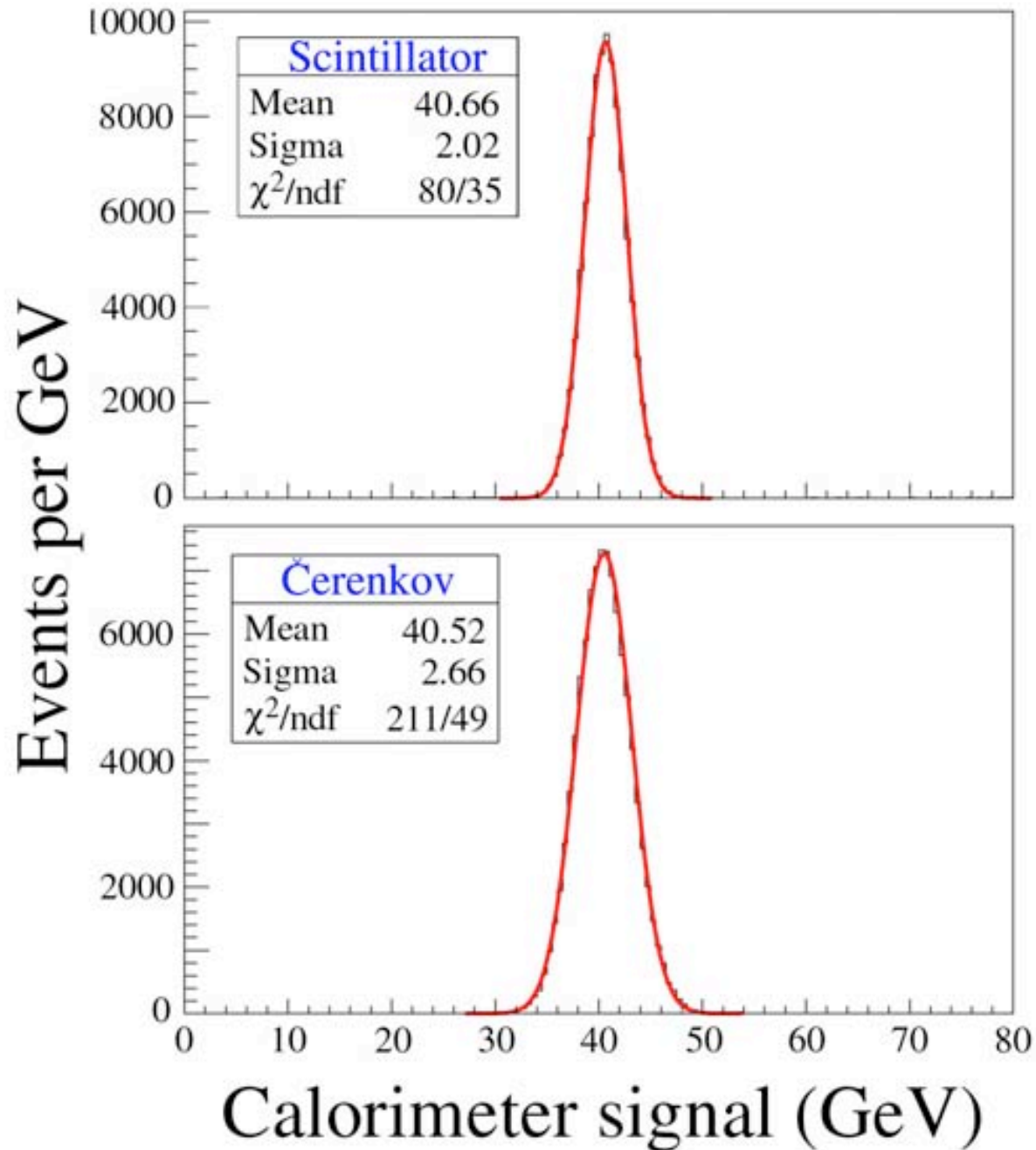
- **Depth** 200 cm ($10.0 \lambda_{\text{int}}$)
- Effective **radius** 16.2 cm ($0.81 \lambda_{\text{int}}$, $8.0 \rho_M$)
- **Mass** instrumented volume 1030 kg
- Number of **fibers** 35910, diameter 0.8 mm, total length ≈ 90 km
- Hexagonal **towers** (19), each read out by 2 PMTs

DREAM readout



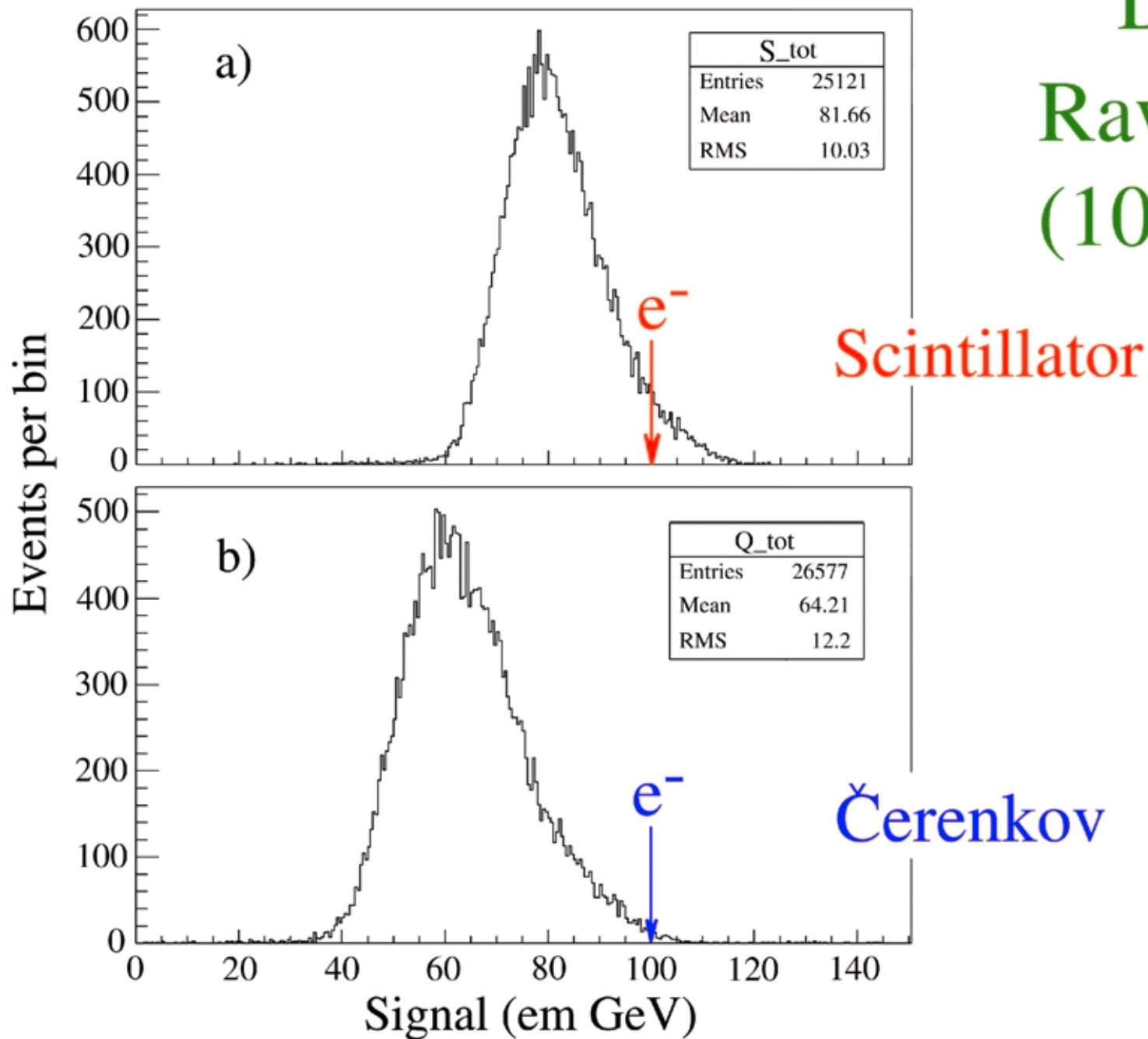


Calibration with 40 GeV electrons (tilt 2°)

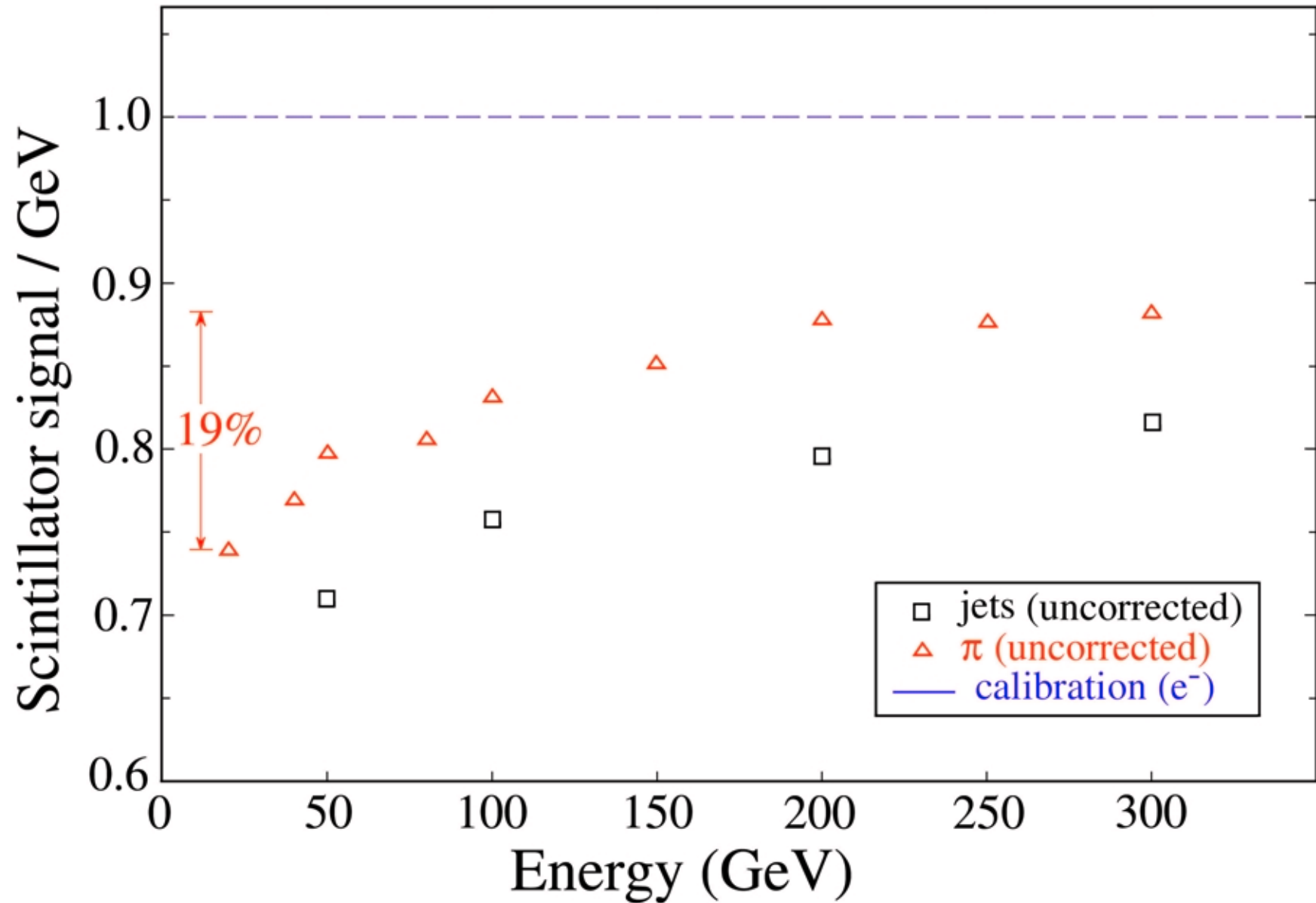


DREAM

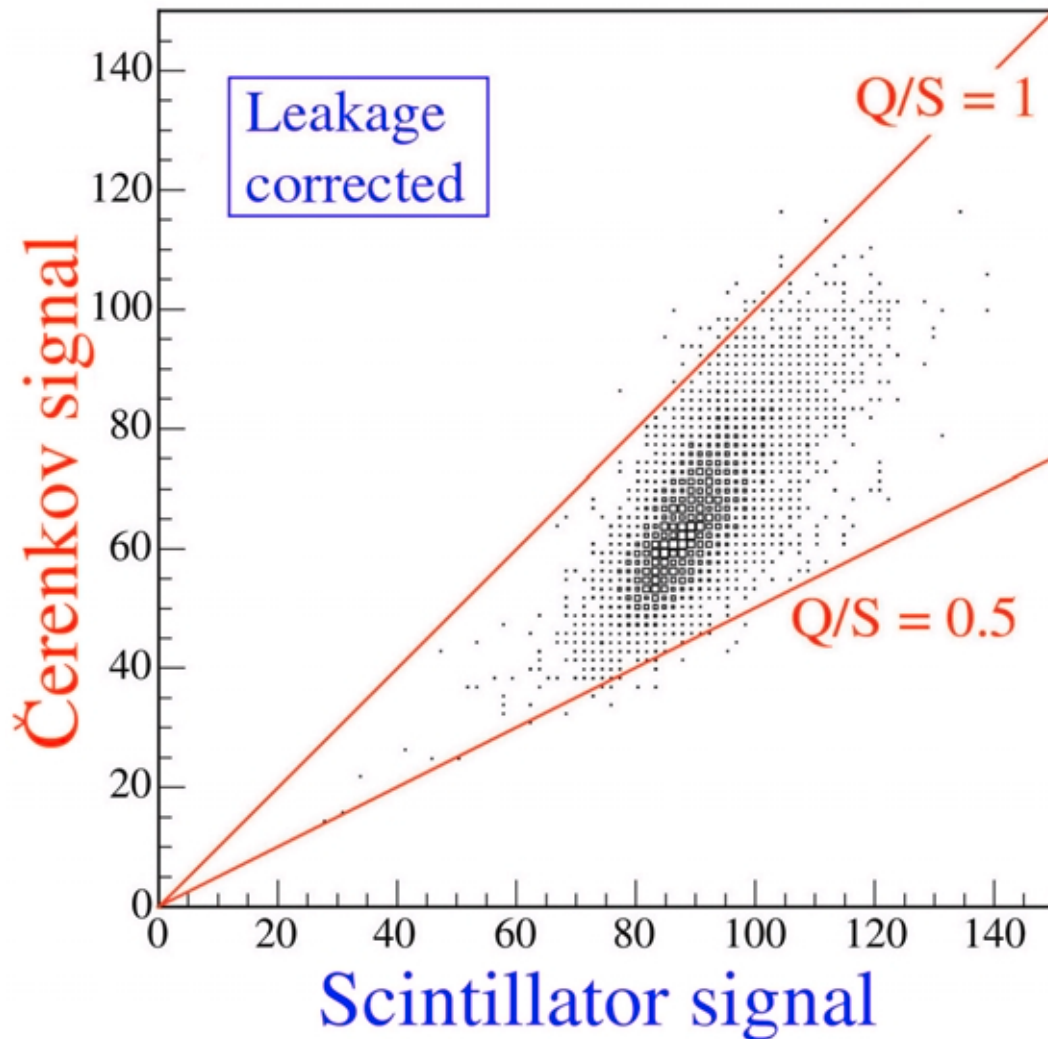
Raw signals (100 GeV π^-)



DREAM: Hadronic response (non-linearity)



DREAM: The (energy-independent) Q/S method



$$S = E \left[f_{\text{em}} + \frac{1}{(e/h)_S} (1 - f_{\text{em}}) \right]$$

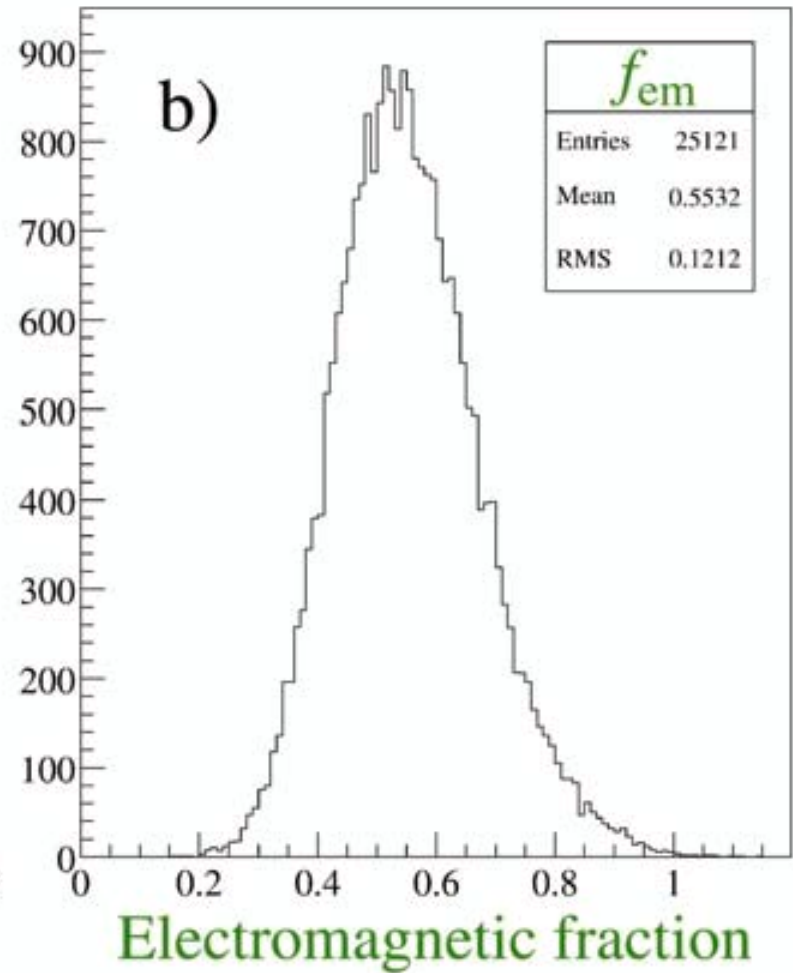
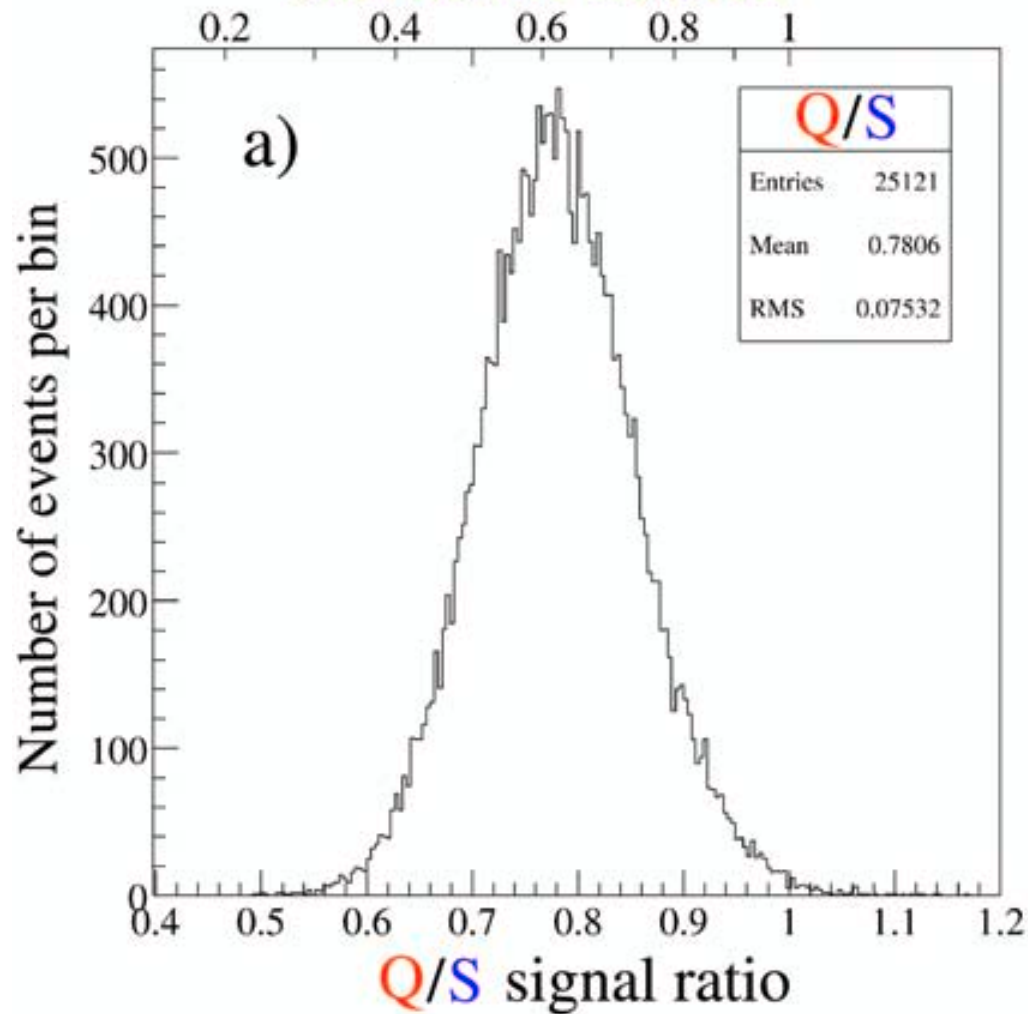
$$Q = E \left[f_{\text{em}} + \frac{1}{(e/h)_Q} (1 - f_{\text{em}}) \right]$$

$$e/h = 1.3 \text{ (S)}, \quad 5 \text{ (Q)}$$

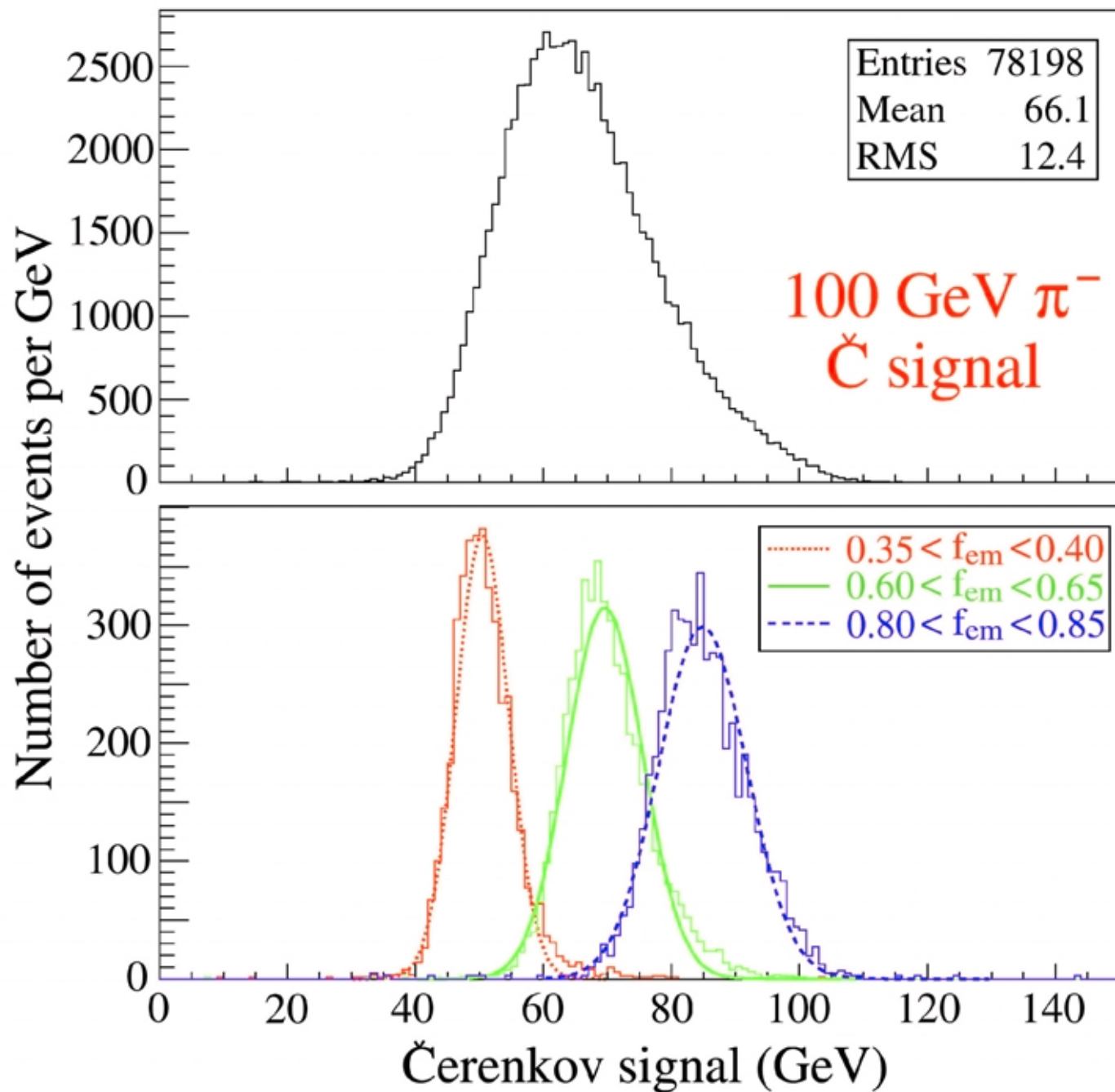
$$\frac{Q}{S} = \frac{f_{\text{em}} + 0.20 (1 - f_{\text{em}})}{f_{\text{em}} + 0.77 (1 - f_{\text{em}})}$$

DREAM: relationship between Q/S ratio and f_{em}

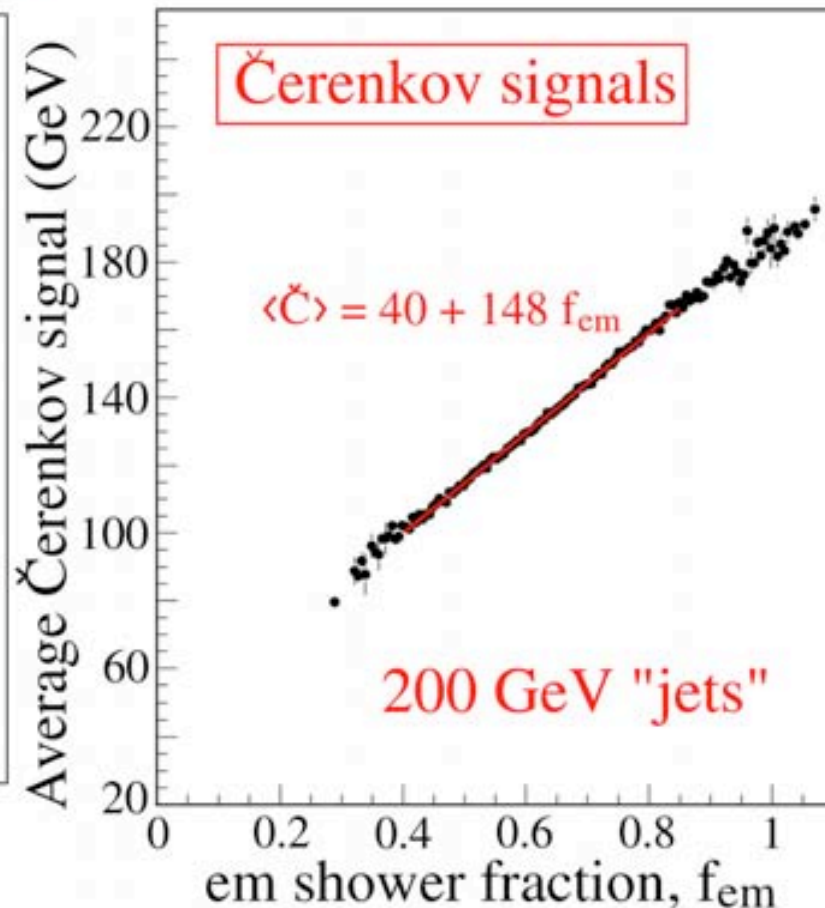
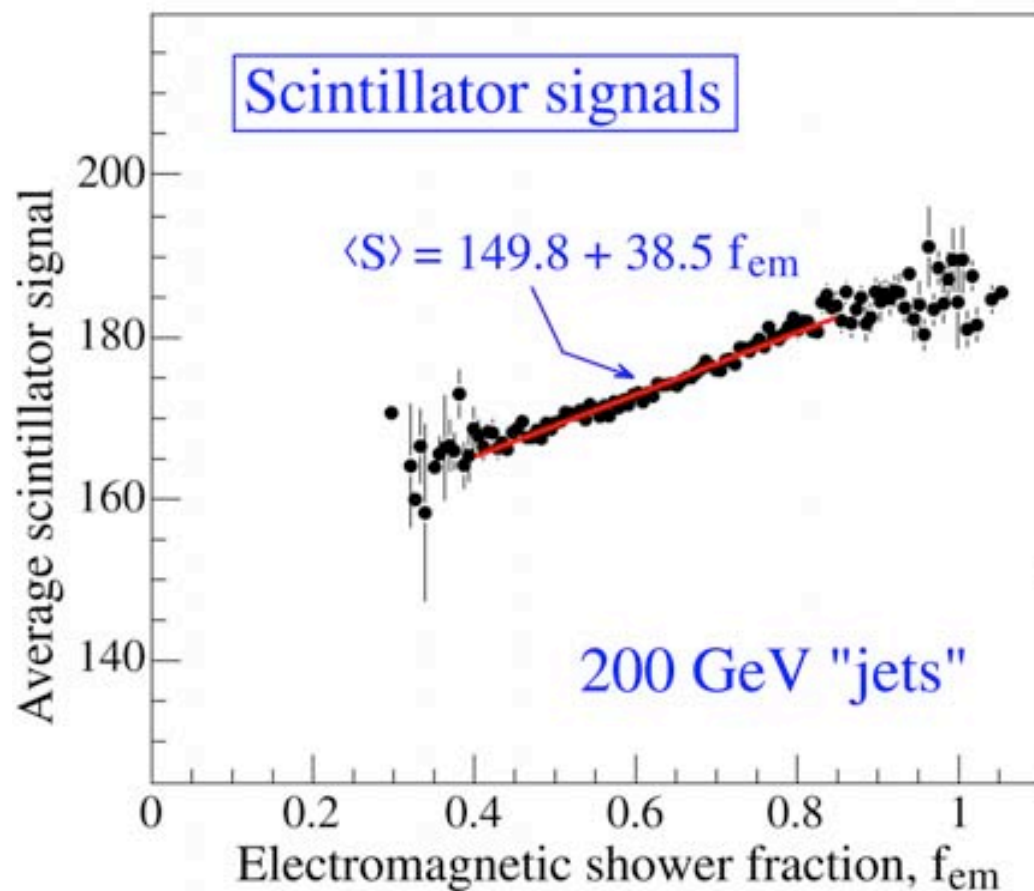
em shower fraction



DREAM: Effect of event selection based on f_{em}



DREAM: Signal dependence on f_{em}



$$R(f_{em}) = p_0 + p_1 f_{em}$$

with

$$\frac{p_1}{p_0} = e/h - 1$$

Cu/scintillator $e/h = 1.3$

Cu/quartz $e/h = 4.7$

Dual-Readout Calorimetry in Practice

The (energy-independent) Q/S method

- Hadronic response (normalized to electrons):

$$R(f_{\text{em}}) = f_{\text{em}} + \frac{1}{e/h} \left[1 - f_{\text{em}} \right], \quad e/h = 1.3 \text{ (S)}, \quad 5 \text{ (Č)}$$

- Q/S response ratio related to f_{em} value \rightarrow find f_{em} from Q/S:

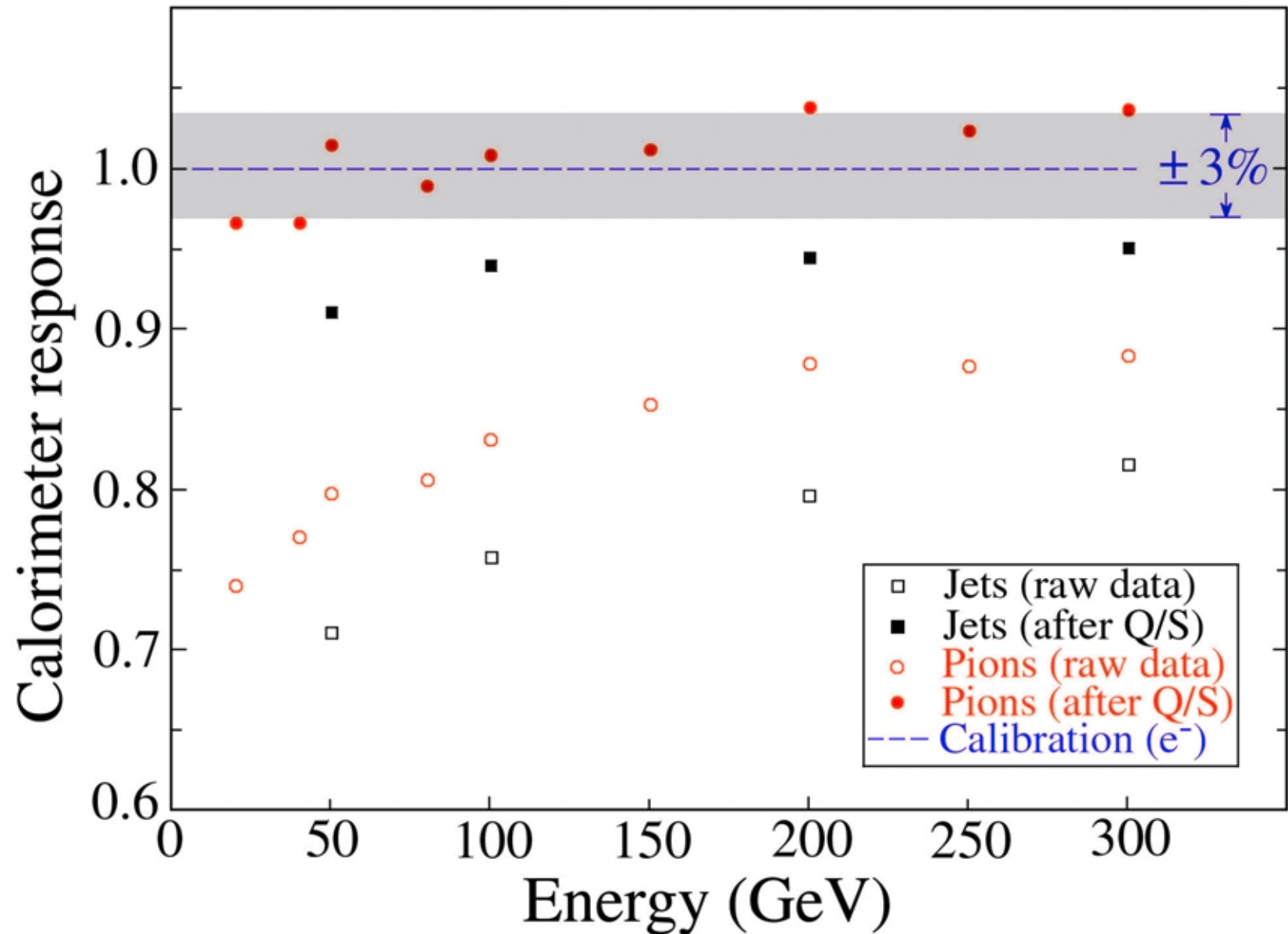
$$\frac{Q}{S} = \frac{R_Q}{R_S} = \frac{f_{\text{em}} + 0.20 (1 - f_{\text{em}})}{f_{\text{em}} + 0.77 (1 - f_{\text{em}})}$$

- Correction to measured signals (regardless of energy):

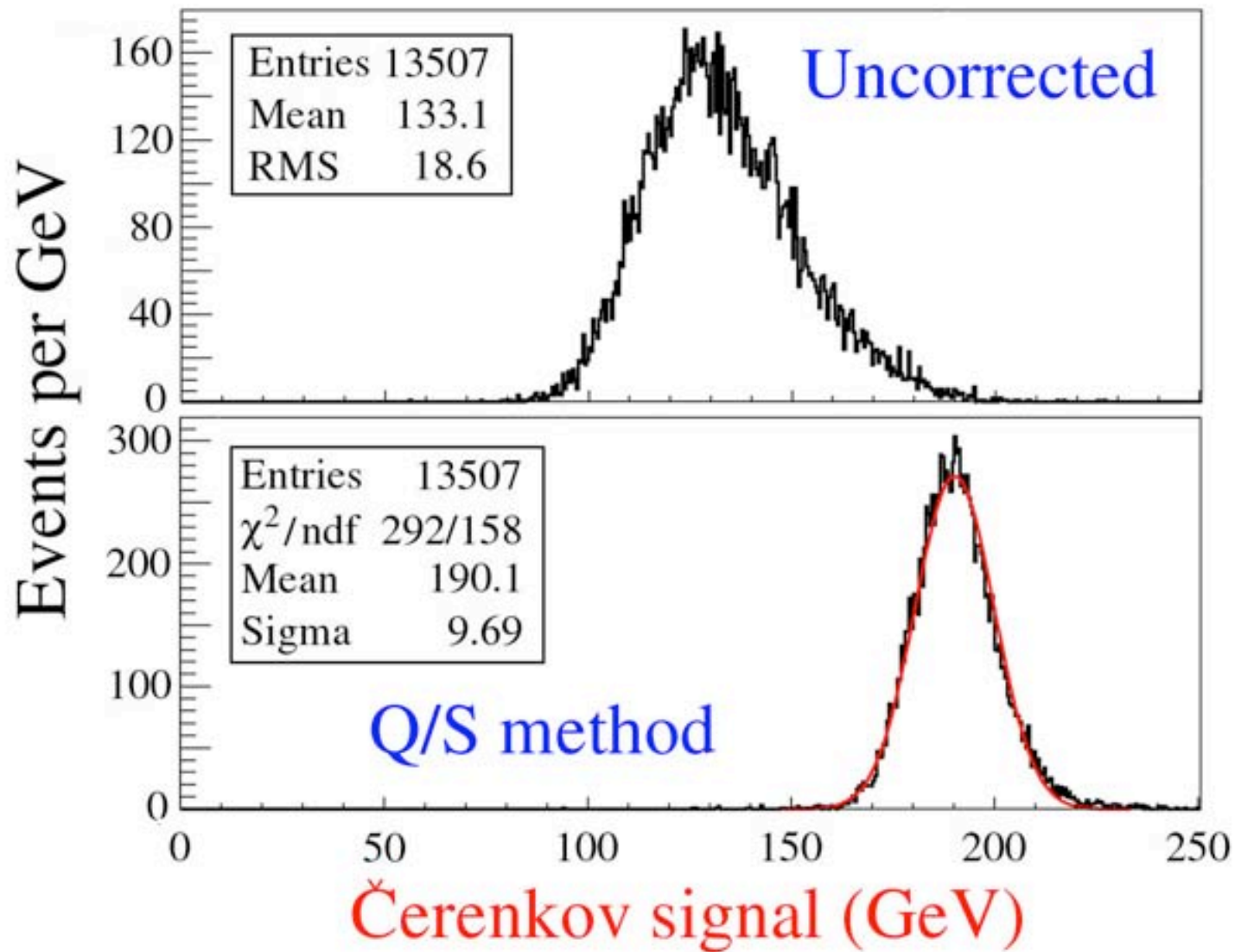
$$S_{\text{corr}} = S_{\text{meas}} \left[\frac{1 + p_1/p_0}{1 + f_{\text{em}} \cdot p_1/p_0} \right], \quad \text{with} \quad \frac{p_1}{p_0} = (e/h)_S - 1$$

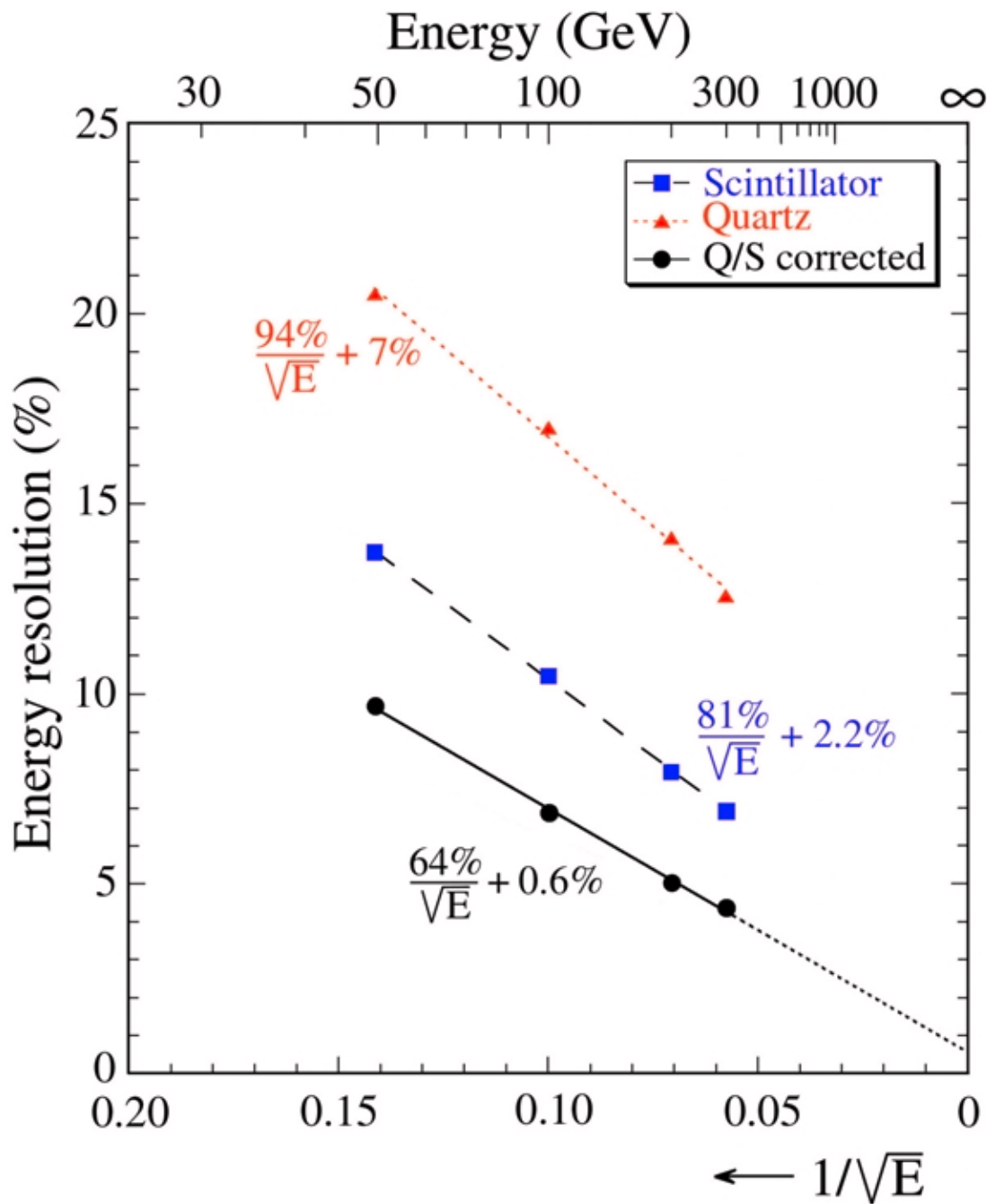
$$Q_{\text{corr}} = Q_{\text{meas}} \left[\frac{1 + p_1/p_0}{1 + f_{\text{em}} \cdot p_1/p_0} \right], \quad \text{with} \quad \frac{p_1}{p_0} = (e/h)_{\check{C}} - 1$$

Hadronic response: Effect Q/S correction



DREAM: Effect of corrections (200 GeV "jets")





DREAM

Energy resolution
"jets"

CONCLUSIONS

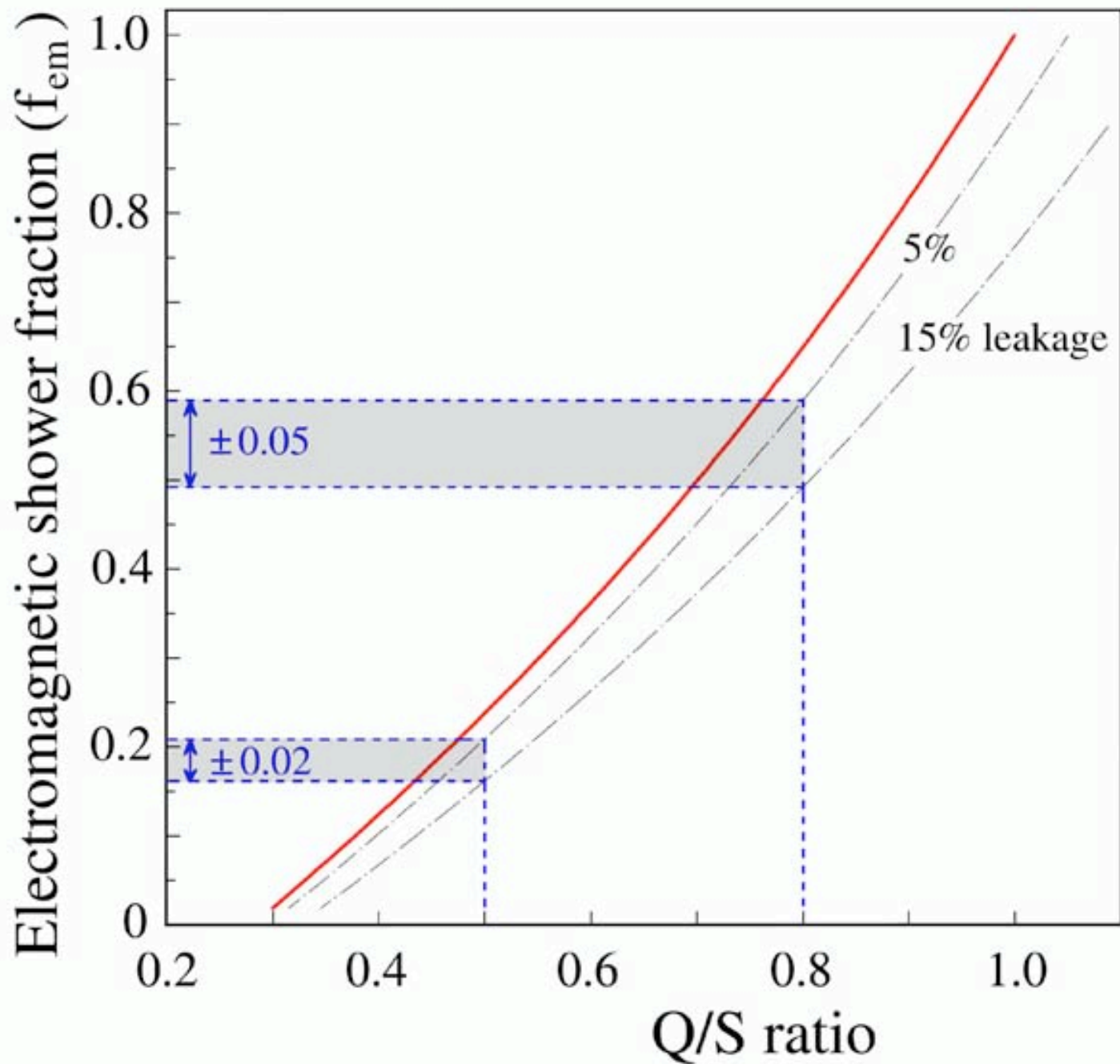
from tests

- **DREAM** offers a powerful technique to *improve* hadronic calorimeter performance:
 - **Correct hadronic energy** reconstruction, *in an instrument calibrated with electrons!*
 - **Linearity** for hadrons and jets
 - **Gaussian** response functions
 - Energy **resolution scales** with $1/\sqrt{E}$
 - $\sigma/E < 5\%$ for high-energy "jets", in a detector with a **mass of only 1 ton!**
dominated by fluctuations in shower leakage
- These, and many other, experimental results are described in 3 papers:
 - Hadrons & jets:** Nucl. Instr. & Meth. A537 (2005) 537
 - Electrons:** Nucl. Instr. & Meth. A536 (2005) 29
 - Muons:** Nucl. Instr. & Meth. A533 (2004) 305

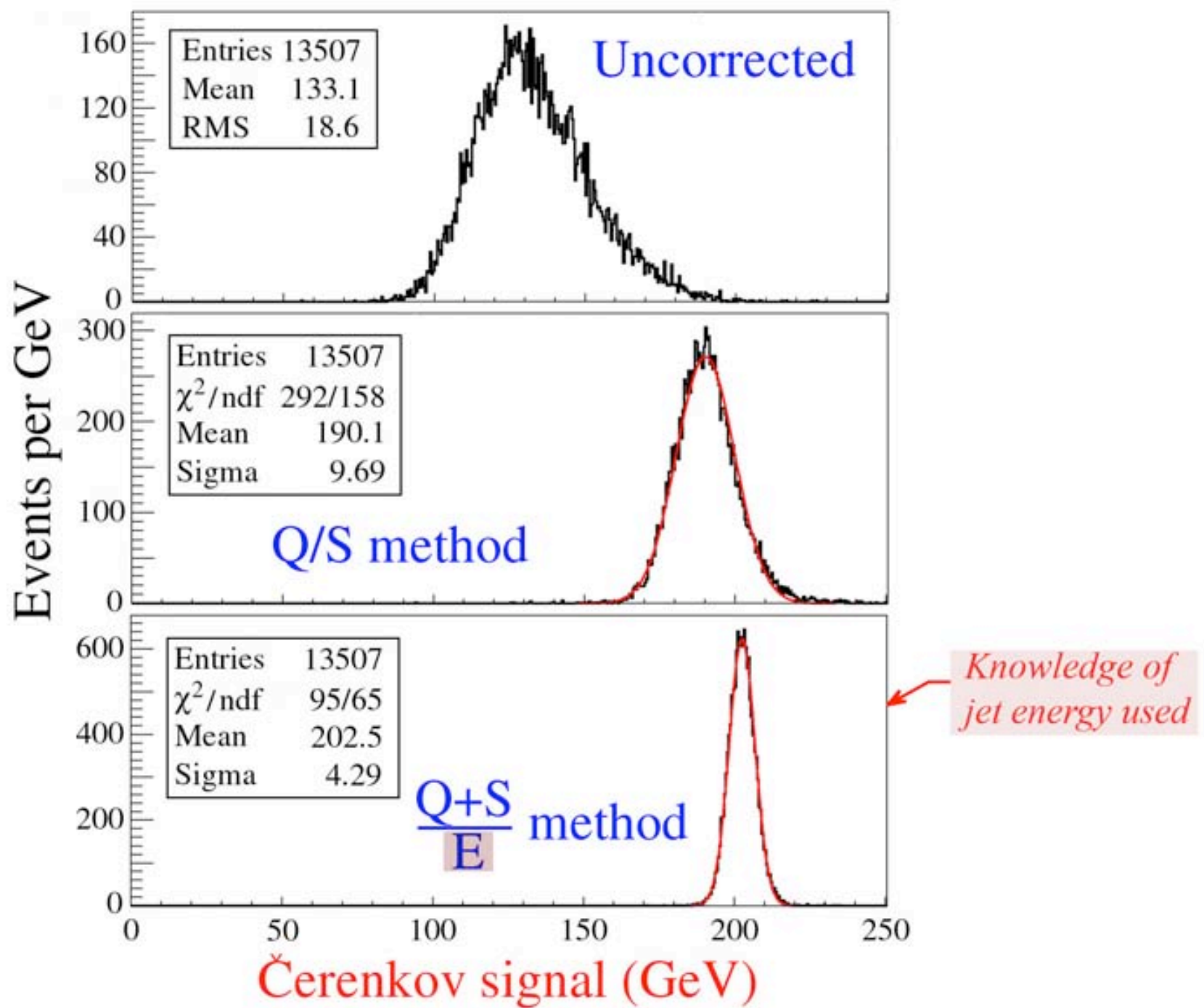
How to improve DREAM performance

- Build a larger detector → *reduce effects side leakage*

DREAM: The importance of leakage and its fluctuations



DREAM: Effect of corrections (200 GeV "jets")

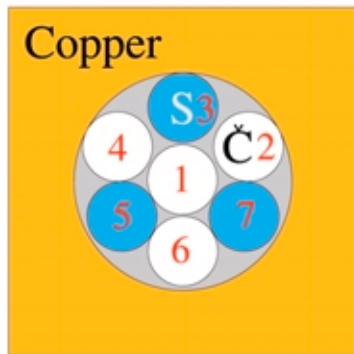


How to improve DREAM performance

- Build a larger detector \longrightarrow *reduce effects side leakage*
- *Increase Čerenkov light yield*
DREAM: 8 p.e./GeV \longrightarrow fluctuations contribute 35%/ \sqrt{E}
No reason why DREAM principle is limited to fiber calorimeters
Homogeneous detector ?!
 \longrightarrow *Need to separate the light into its Č, S components*

DREAM 2

- To what extent can **light** from an optical calorimeter be **separated into** its **scintillation** and **Čerenkov** components?
- *Modified* the DREAM calorimeter

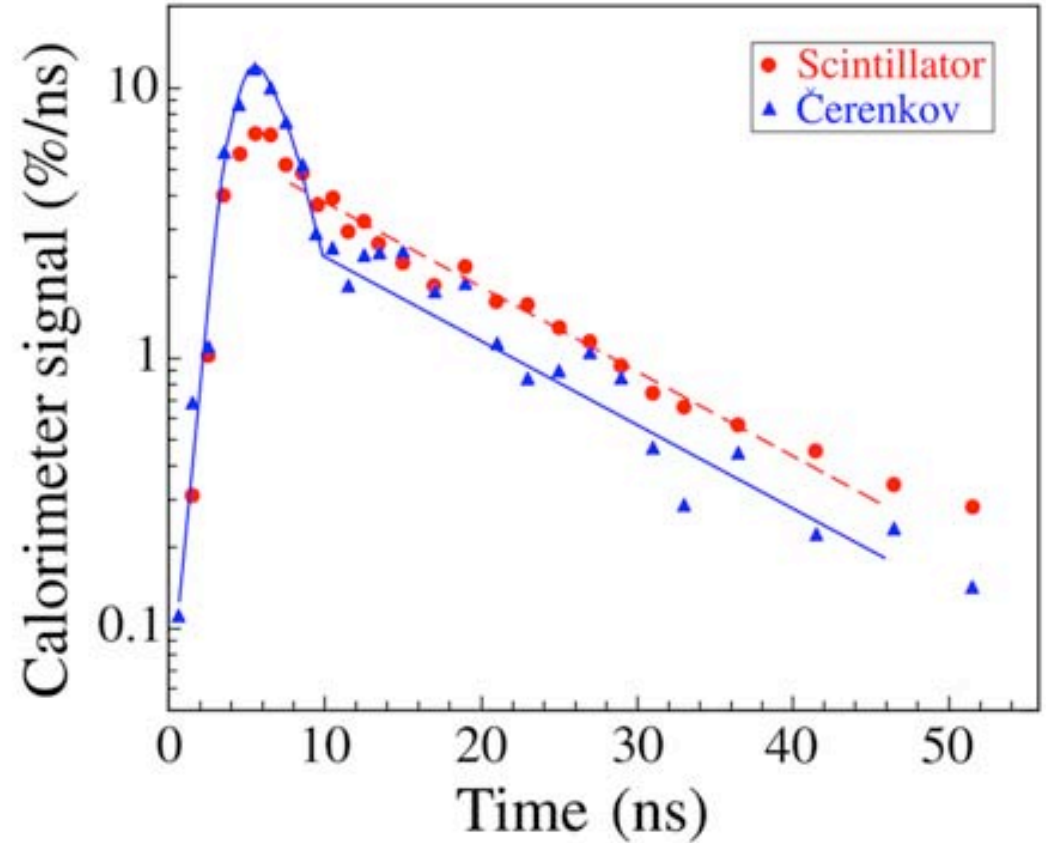
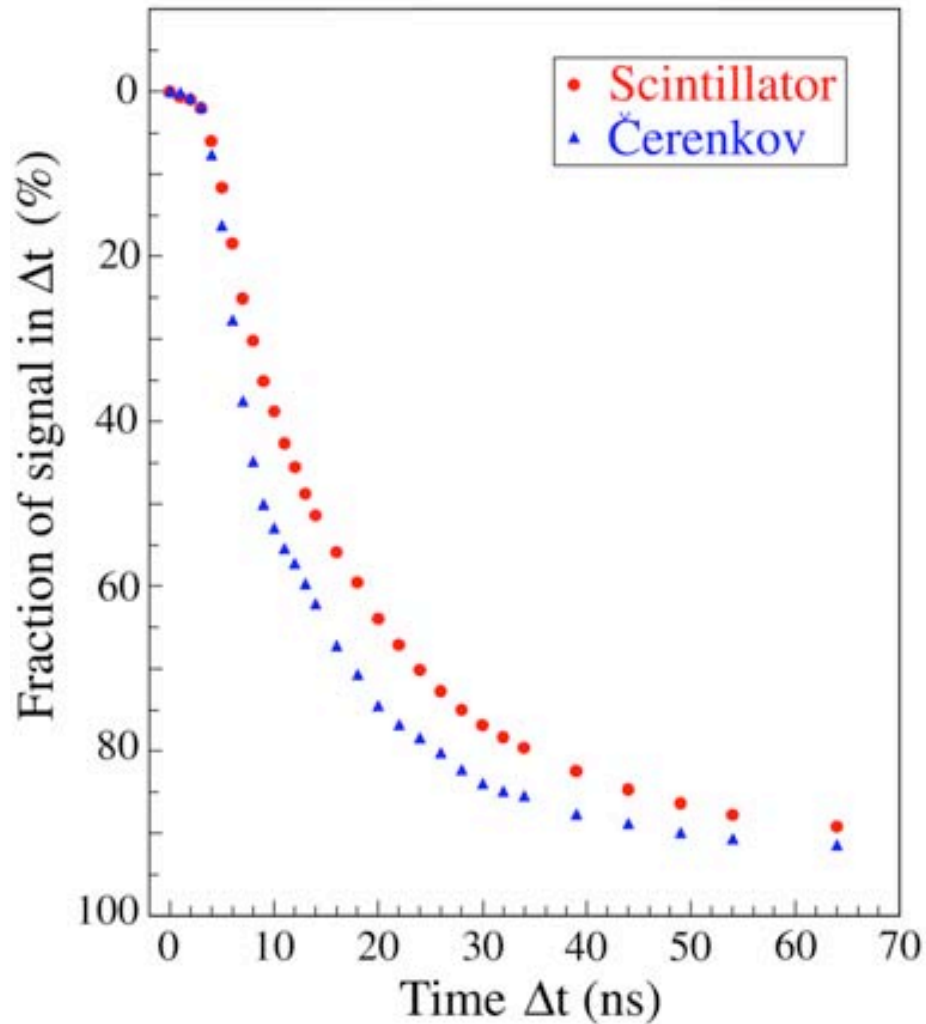


5+7	PMT1 (S)
2+4	PMT2 (Č)
1+3+6	PMT3 (mixed)

Also:
Fibers read out from both ends

- Separation methods based on differences in:
 - Time structure of signals
 - Light directionality
 - Optical spectra
 - Polarization

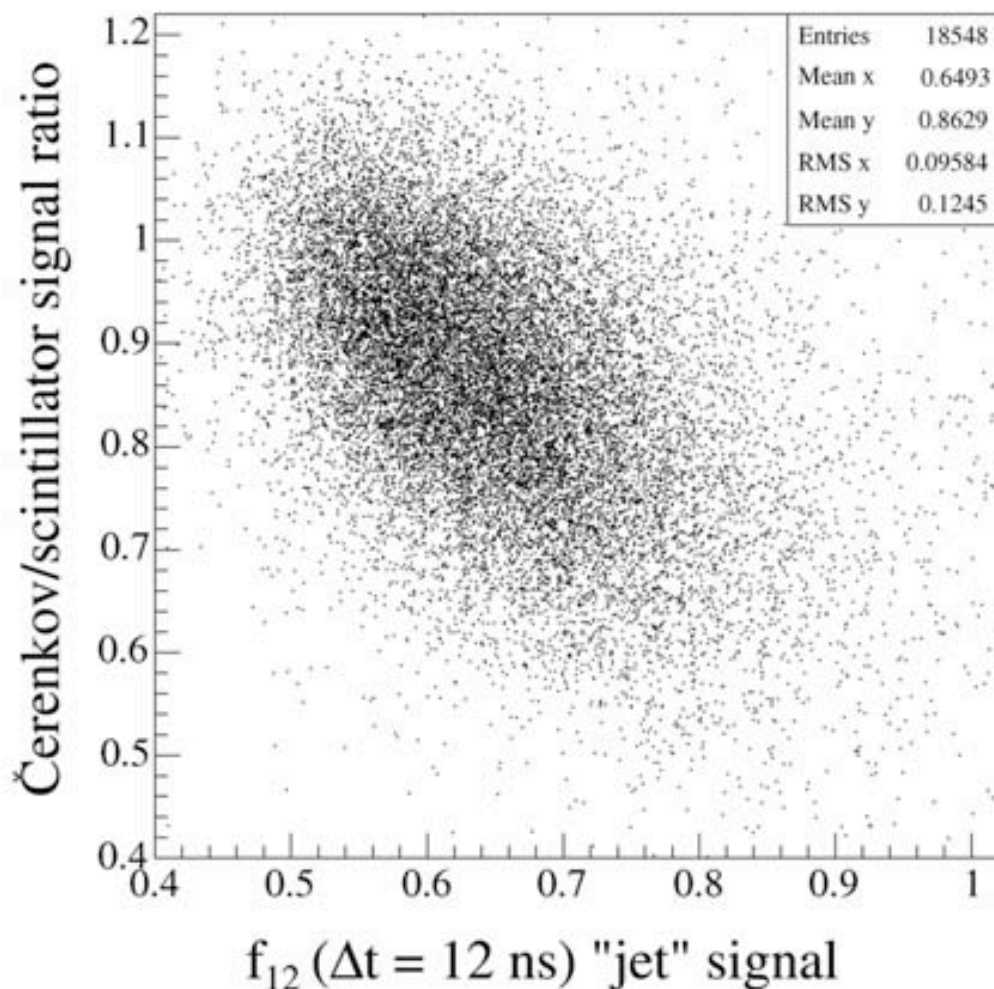
DREAM - Differences in signal time structure



NIM A550 (2005) 185

DREAM: Separation of mixed signal into Č/S components

Use **time structure** event by event



Fraction of signal observed in first 12 ns is measure for Čerenkov/Scintillator ratio

*Limiting factor:
Čerenkov light yield
(8 p.e./GeV)*

NIM A 550 (2005) 185

How to improve DREAM performance

- Build a larger detector \longrightarrow *reduce effects side leakage*
- *Increase Čerenkov light yield*
DREAM: 8 p.e./GeV \longrightarrow fluctuations contribute 35%/ \sqrt{E}
No reason why DREAM principle is limited to fiber calorimeters
Homogeneous detector ?!
 \longrightarrow *Need to separate the light into its Č, S components*
- For ultimate hadron calorimetry (15%/ \sqrt{E}): *Measure E_{kin} (neutrons)*
Is correlated to nuclear binding energy loss (invisible energy)
Can be measured with third type of fiber: **TREAM**

Summary

- New developments in calorimetry are primarily driven by ILC/CLIC. Some of the R&D results may be useful for SLHC.
- Calorimeter design of ATLAS/CMS was driven by $H^0 \rightarrow \gamma\gamma$.
May not be the same choice for the SLHC era.
Redesign calorimeter system with focus on hadronic performance?
- Increased luminosity will have major consequences for calorimeters, in terms of radiation damage and operating conditions.
- The DREAM approach seems to combine the advantages of compensation with a reasonable amount of design flexibility.

(pet pief) Longitudinal segmentation is asking for trouble,
GEANT is (unfortunately) useless for hadron calorimetry

Backup slides

- 1) Digital calorimetry
- 2) Fluctuations in hadron showers
- 3) Monte Carlo benchmark p/π
- 4) Monte Carlo benchmark n
- 5) Calibration CDF

Saturation in "digital" calorimeters (wire chamber readout)

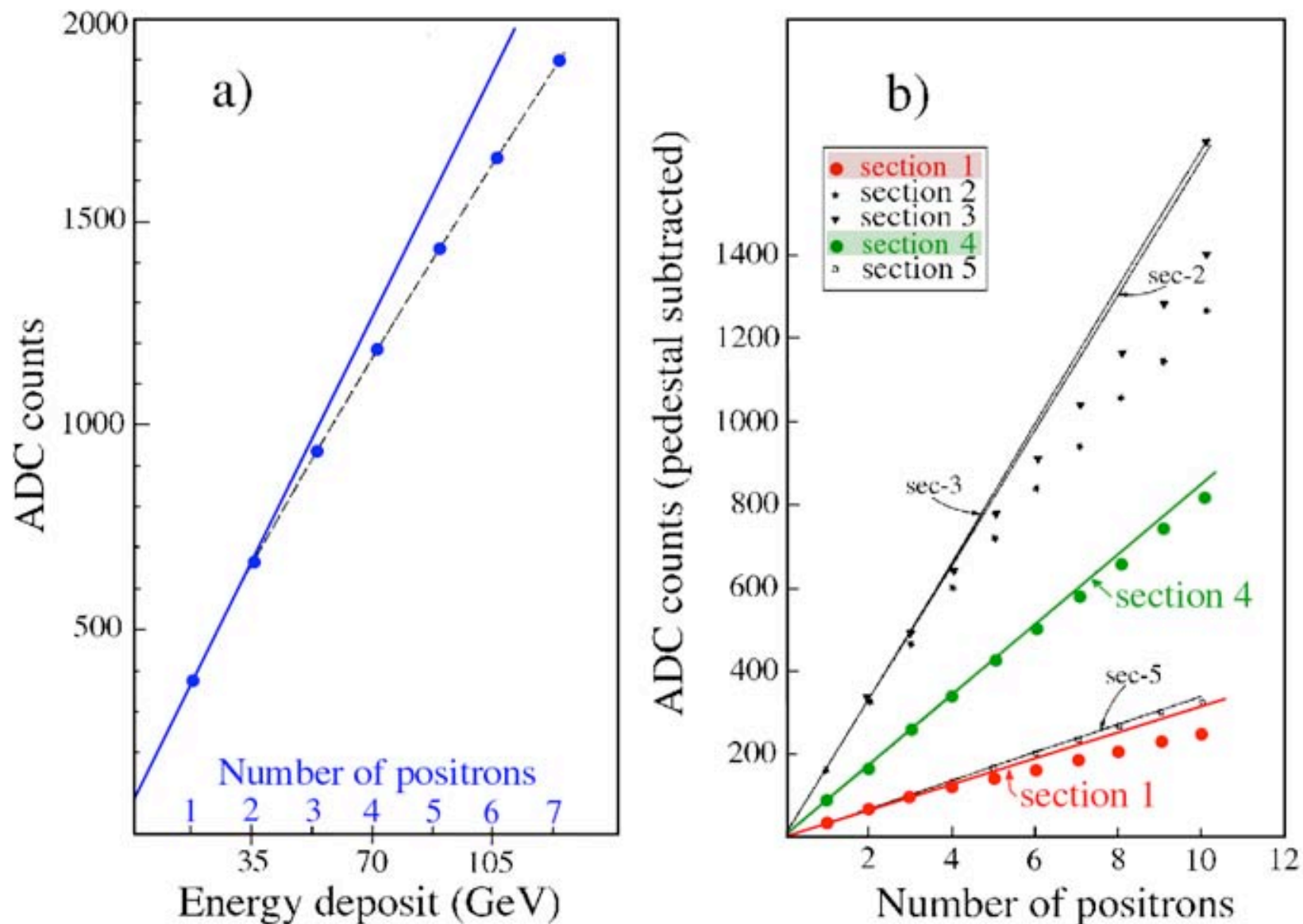


FIG. 3.2. Average em shower signal from a calorimeter read out with gas chambers operating in a "saturated avalanche" mode, as a function of energy. [From: NIM 205 \(1983\) 113.](#)

Hadronic shower profiles: Fluctuations!

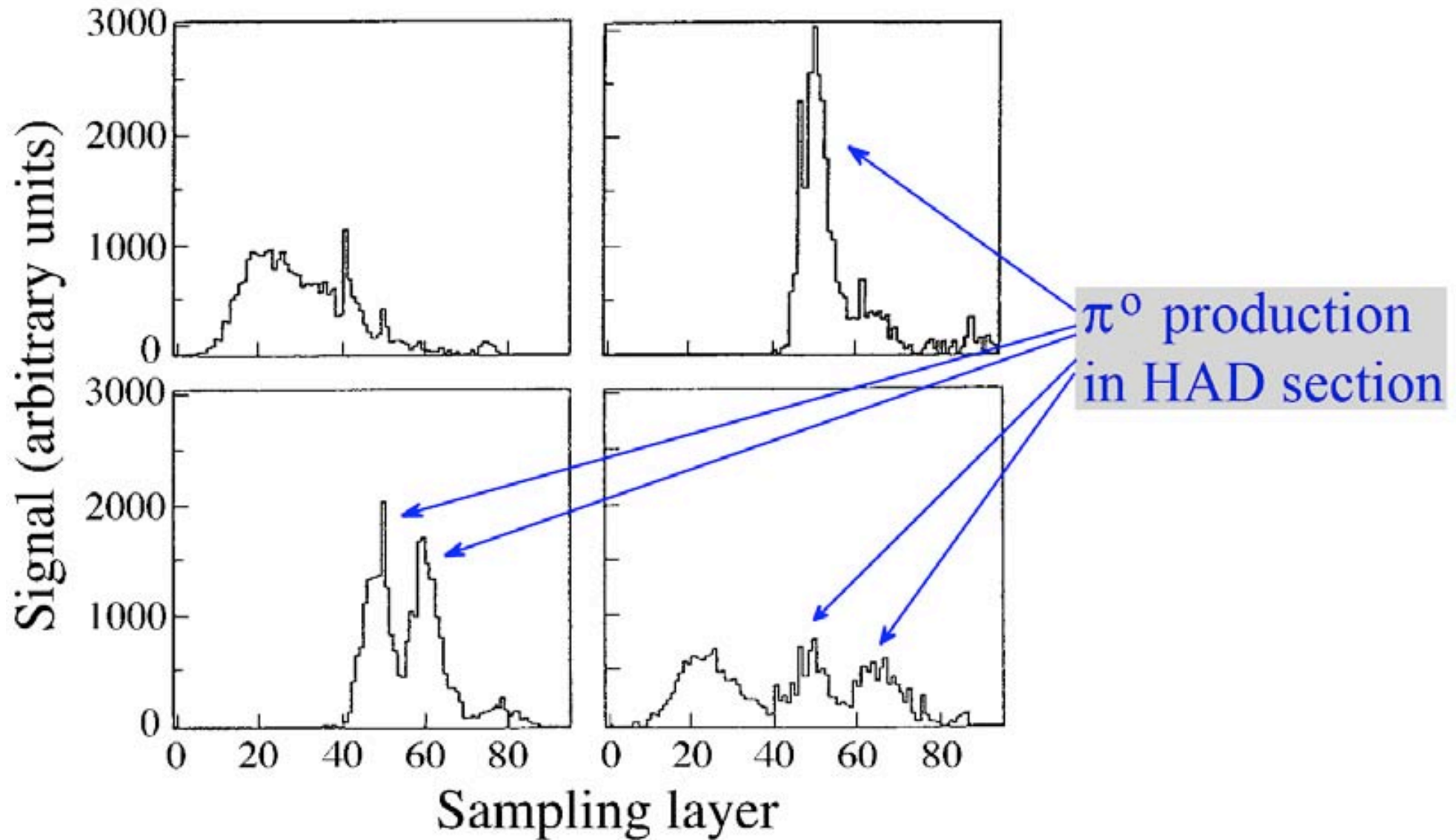


FIG. 2.35. Longitudinal profiles for 4 different showers induced by 270 GeV pions in a lead/iron/plastic-scintillator calorimeter. Data from [Gre 94].

Benchmark data for hadronic Monte Carlo

Test of π^0 production modelling

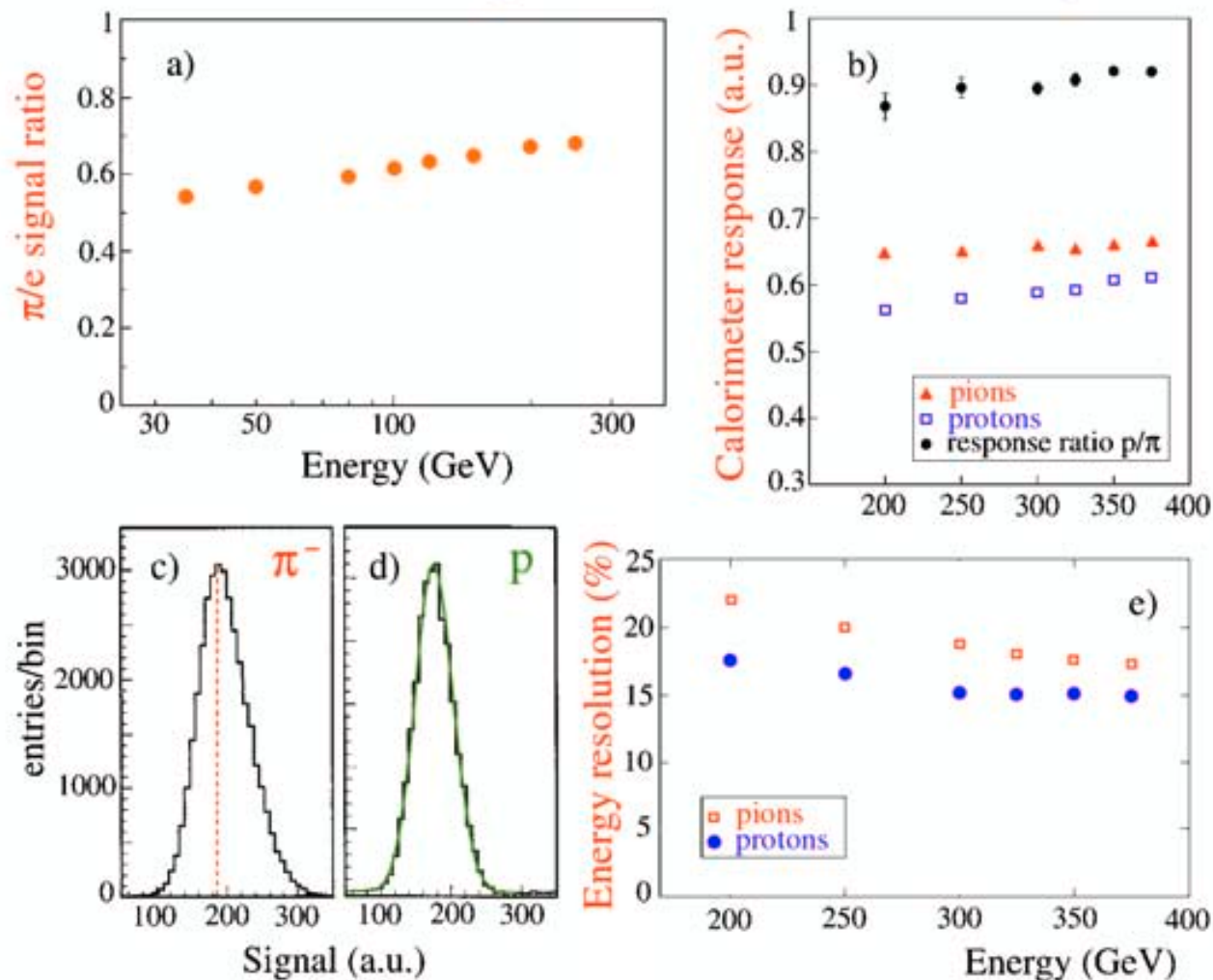


FIG. 8.27. Calorimeter benchmark data for testing the correct implementation of π^0 production in Monte Carlo simulations of hadronic shower development. Experimental data from a copper/quartz-fiber calorimeter, showing the π/e signal ratio as a function of energy (a), the response to protons and pions, as well as the ratio of these responses, as a function of energy (b), the response functions to 300 GeV pions (c) and protons (d), and the energy resolutions for pions and protons as a function of energy (e) [Akc 97].

Benchmark data for hadronic Monte Carlo

Test of description neutron effects

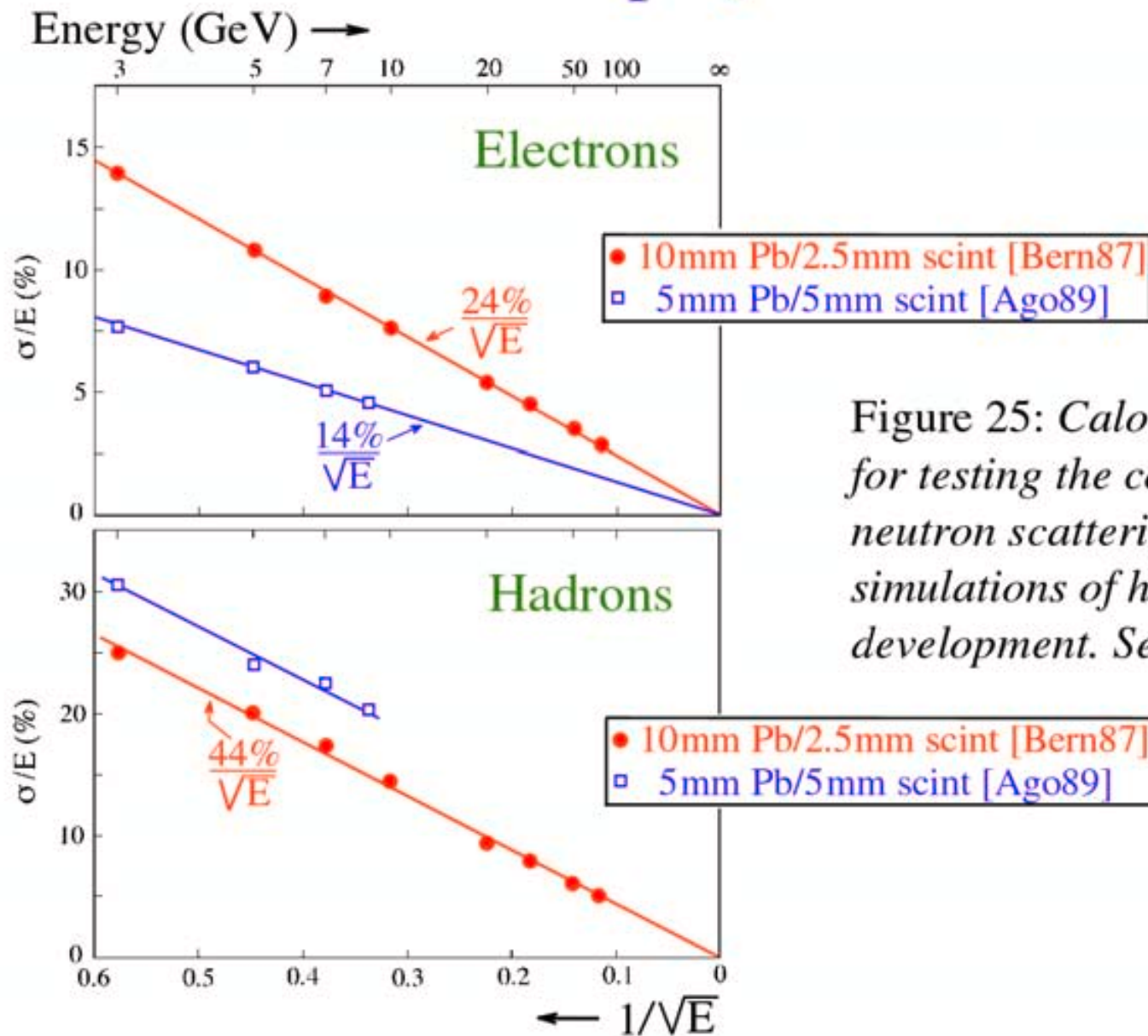


Figure 25: Calorimeter benchmark data for testing the correct implementation of neutron scattering data in Monte Carlo simulations of hadronic shower development. See text for details.

Effects of miscalibration: Mass reconstruction

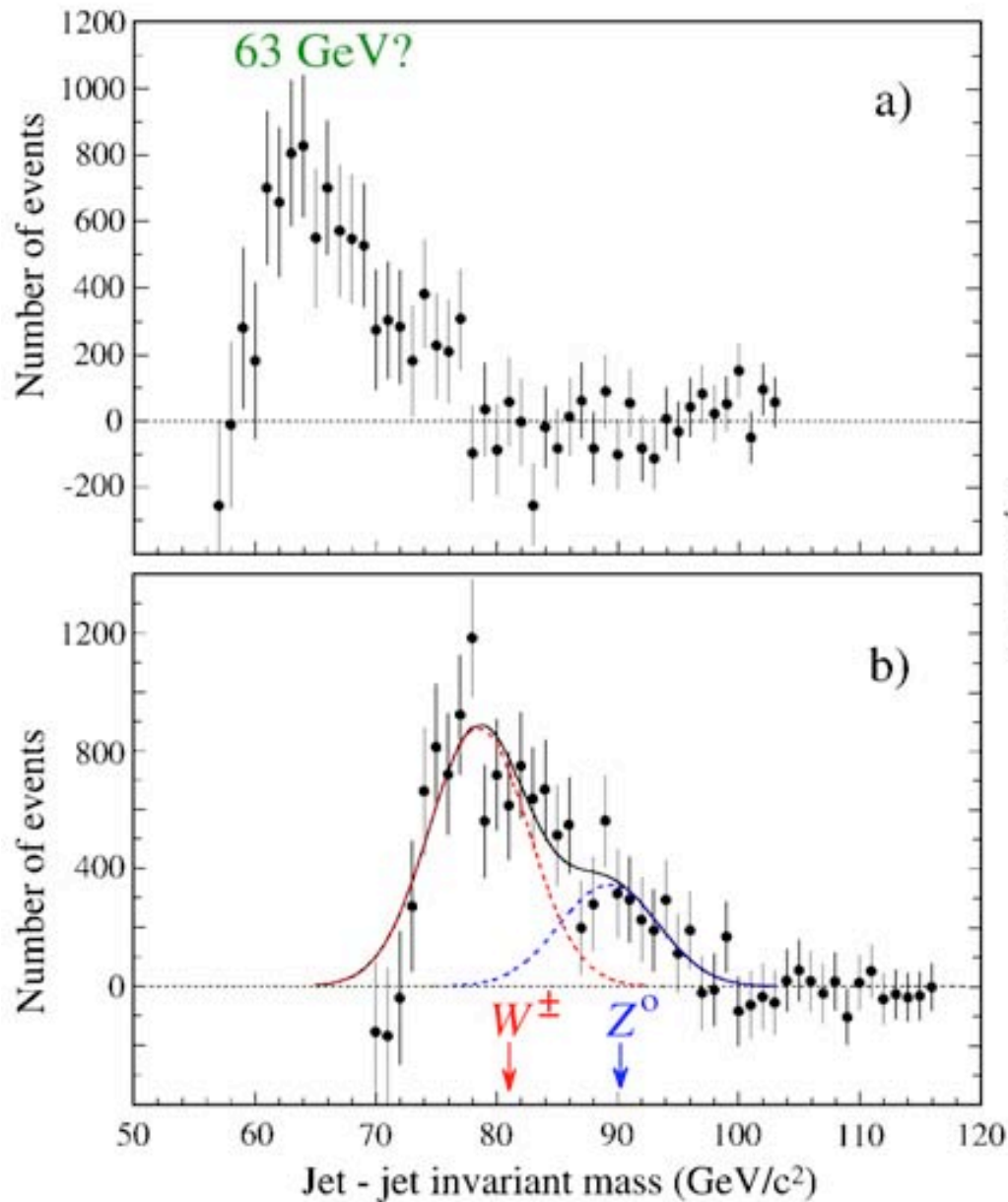


Figure 9: *The jet-jet invariant mass distribution for CDF Run 1c. The background has been subtracted. Results are shown for the standard CDF calibration constants (a) and for our alternative calibration scheme (b).*

From CDF note 6530 (2002)