

Automatic Loop Calculations in the MSSM with Non-minimal Flavor Violation

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CERN, January 2005

on collaboration with

S. Heinemeyer and *T. Hahn*

- FeynArts, FormCalc, LoopTools with **NMFV**
- FeynHiggs2.2 with **NMFV**
- E.g.: Corrections to M_W , $\sin^2 \theta_{\text{eff}}$, m_h
- Conclusions

NMFV \Rightarrow Connection to Flavour Physics

FEYNARTS

FORMCALC

and

LOOPTOOLS

FeynArts

T. Hahn, Comput. Phys. Commun. 140 (2001) 418, hep-ph/0012260

The MSSM model files are documented in:

Comput. Phys. Commun. 143 (2002) 54, hep-ph/0105349

FormCalc, LoopTools

T. Hahn, M. Perez-Victoria,

Comput. Phys. Commun. 118 (1999) 153, hep-ph/9807565

Latest version:

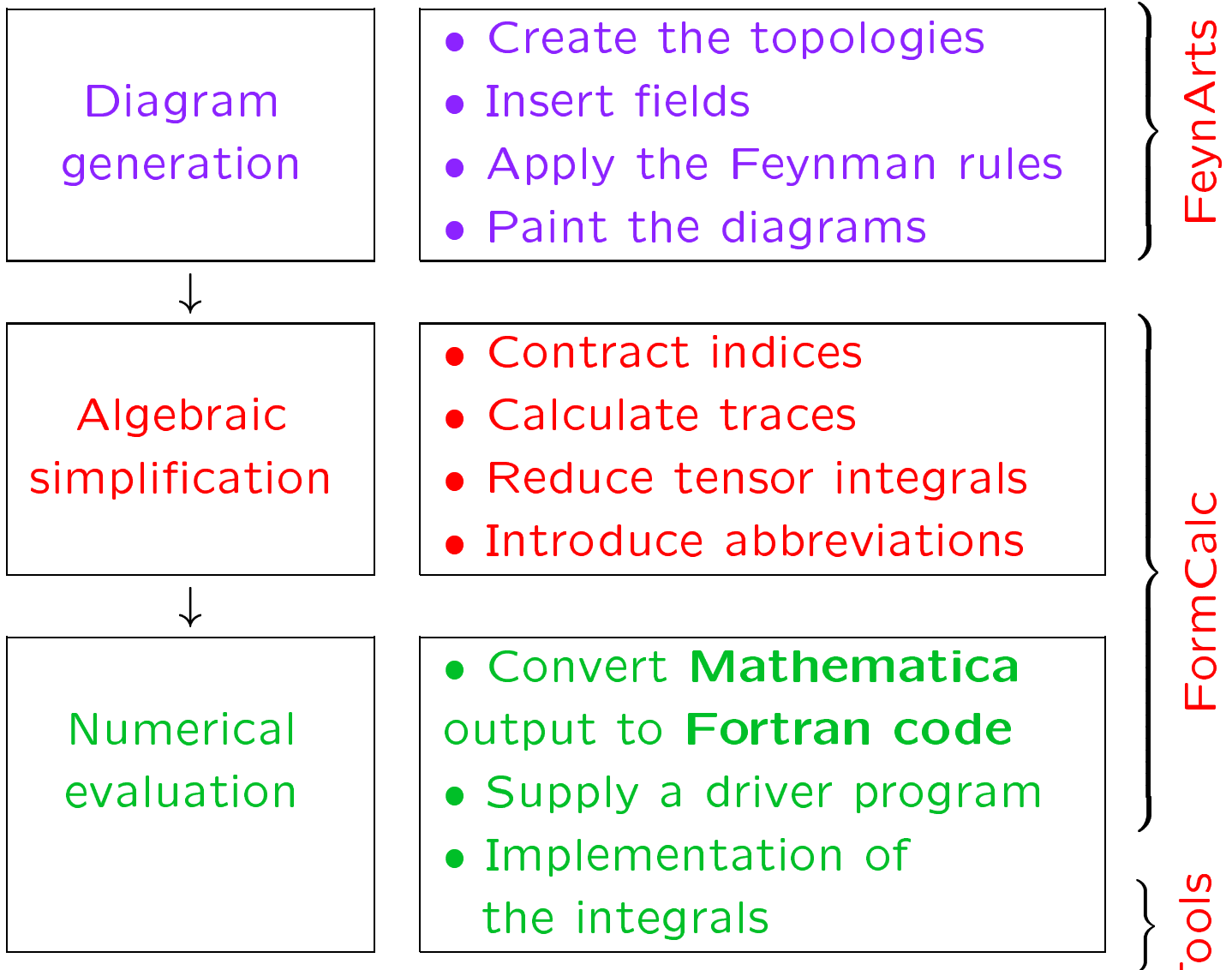
<http://www.feynarts.de> - (10 Jan 05)

<http://www.feynarts.de/formcalc> - (20 Jan 05)

<http://www.feynarts.de/looptools> - (7 Dec 04)

Automatic Installation with FeynInstall

Evaluating Feynman Diagrams



The three programs can be used for one-loop calculations and they work together smoothly

FeynArts is a Mathematica package

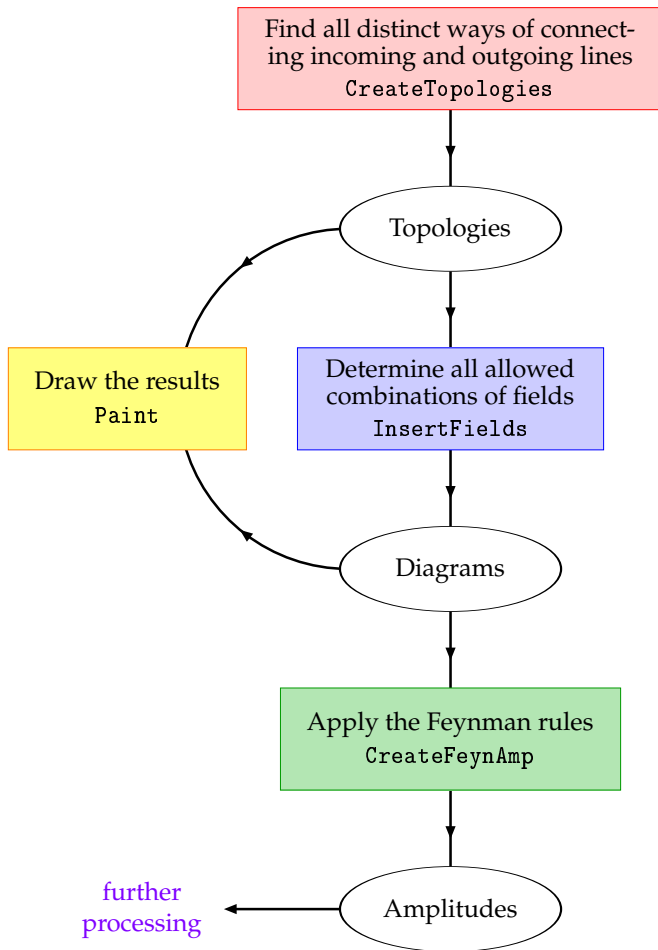
FormCalc is a Mathematica package, uses FORM

J. Vermaseren, math-ph/0010025, <http://www.nikhef.nl/~form/>

LoopTools implementation in Fortran, can be used in Fortran/Mathematica/C++

How FeynArts works ?

FeynArts



EXAMPLE: generating the Higgs self-energy

```
top = CreateTopologies[ 1, 1 -> 1 ]
```

one loop
one incoming particle
one outgoing particle

```
Paint[top]
```

```
ins = InsertFields[ top, S[1] -> S[1],  
Model -> SM ]
```

use the Standard Model

the name of the
Higgs boson in the
"SM" model file

```
Paint[ins]
```

```
amp = CreateFeynAmp[ins]
```

```
amp >> HiggsSelfEnergy.amp
```

Generating diagrams in just a few lines - aren't there any strings attached ?

Yes, one has to set up, once and for all, a MODEL FILE containing the particles and the allowed couplings .

★ FeynArts Model files can be supplied/modified by the user

Model Files Currently included:

- **Standard Model (SM) [w/QCD]**
All one-loop counter terms included.
- **Minimal Supersymmetric Standard Model (MSSM) [w/QCD]**
Counter terms by T. Fritzsche.
- **Two-Higgs-Doublet Model (THDM)**
Counter terms not included yet.

In preparation: **FVMSSM**

by T.Hahn and J.I.Illana

FVMSSM.mod

Model file for the flavour-violating MSSM

[FeynArts-3.2/Models/FVMSSM.mod](#) -

last modified 7/10/04

- Feynman rules with FV are included by modifying an existing model file **MSSMQCD.mod**
→ **FVMSSM.mod** is technically not a full model file, but loads **MSSM.mod** and applies changes
- **FormCalc** files must be modified as well
→ Set up the **model.h** and **mssm-ini.F** files which declares and computes the squarks mixing matrices

FVMSSM.mod

Feynman rules included:

[FFS] 2 Leptons - Higgs

2 Quarks - Higgs

Chargino - Lepton - Slepton

Chargino - Quark - Squark

Gluino - Quark - Squark

Lepton - Neutralino - Slepton

Neutralino - Quark - Squark

[FFV] 2 Leptons - Gauge boson

2 Quarks - Gauge boson

2 Quarks - Gluon

[SSS] Higgs - 2 Sleptons

Higgs - 2 Squarks

[SSSS] 2 Higgs - 2 Sletons

2 Higgs - 2 Squarks

2 Sleptons - 2 Squarks

4 Sleptons

4 Squarks

[SSV] 2 Sleptons - Gauge bosons

2 Squarks - Gauge bosons

2 Squarks - Gluon

[SSVV] 2 Sleptons - 2 Gauge bosons

2 Squarks - 2 Gauge bosons

2 Squarks - 2 Gluons

2 Squarks - Gauge boson - Gluons

E.g. the **SSS** coupling is declared by:

- in the Classes model file (here for $h_0 \tilde{u}_g^s \bar{\tilde{u}}_g^s$)

$$\begin{aligned}
& C[S[1], S[13, s1, j1, o1], -S[13, s2, j2, o2]] == \\
& ((-I/6) * EL * IndexDelta[j1, j2] * IndexDelta[o1, o2] * \\
& (Conjugate[USf[3, j1][s1, 1]] * \\
& ((MW * MZ * SAB * SB * (-3 + 4 * SW^2) + \\
& 6 * CA * CW * Mass[F[3, j1]]^2) * USf[3, j1][s2, 1] + \\
& 3 * CW * (CA * Af[3, j1] + SA * Conjugate[MUE]) * \\
& Mass[F[3, j1]] * USf[3, j1][s2, 2]) + \\
& Conjugate[USf[3, j1][s1, 2]] * \\
& (3 * CW * (MUE * SA + CA * Conjugate[Af[3, j1]]) * \\
& Mass[F[3, j1]] * USf[3, j1][s2, 1] - \\
& 4 * MW * MZ * SAB * SB * SW^2 * USf[3, j1][s2, 2] + \\
& 6 * CA * CW * Mass[F[3, j1]]^2 * USf[3, j1][s2, 2])) \\
& / (CW * MW * SB * SW)
\end{aligned}$$

USf - sfermion mixing matrices (6x6)

Mass[F] - sfermion masses

NMFV in the MSSM

→ Mixing of scalar quark families (beyond CKM)

E.g.: General case of mixing between the third and second generation of squarks ($\tilde{t}/\tilde{c}, \tilde{b}/\tilde{s}$ sectors)

Squark Generation Mixing via Soft Breaking

Parametrization of non-diagonal squark mass matrices

$$M_{\tilde{u}}^2 = \begin{pmatrix} M_{\tilde{L}_c}^2 & \Delta_{LL}^t & m_c X_c & \Delta_{LR}^t \\ \Delta_{LL}^t & M_{\tilde{L}_t}^2 & \Delta_{RL}^t & m_t X_t \\ m_c X_c & \Delta_{RL}^t & M_{\tilde{R}_c}^2 & \Delta_{RR}^t \\ \Delta_{LR}^t & m_t X_t & \Delta_{RR}^t & M_{\tilde{R}_t}^2 \end{pmatrix}$$

$$M_{\tilde{d}}^2 = \begin{pmatrix} M_{\tilde{L}_s}^2 & \Delta_{LL}^b & m_s X_s & \Delta_{LR}^b \\ \Delta_{LL}^b & M_{\tilde{L}_b}^2 & \Delta_{RL}^b & m_b X_b \\ m_s X_s & \Delta_{RL}^b & M_{\tilde{R}_s}^2 & \Delta_{RR}^b \\ \Delta_{LR}^b & m_b X_b & \Delta_{RR}^b & M_{\tilde{R}_b}^2 \end{pmatrix}$$

$$M_{\tilde{L}_q}^2 = M_{\tilde{Q}_q}^2 + m_q^2 + \cos 2\beta M_Z^2 (T_3^q - Q_q s_w^2)$$

$$M_{\tilde{R}_q}^2 = M_{\tilde{U}_q}^2 + m_q^2 + \cos 2\beta M_Z^2 Q_q s_w^2 \quad (q = t, c)$$

$$M_{\tilde{R}_q}^2 = M_{\tilde{D}_q}^2 + m_q^2 + \cos 2\beta M_Z^2 Q_q s_w^2 \quad (q = b, s)$$

$$X_q = A_q - \mu (\tan \beta)^{-2T_3^q}$$

E.g. $\Delta_{LL}^t = \lambda^t M_{\tilde{L}_t} M_{\tilde{L}_c}, \Delta_{LL}^b = \lambda^b M_{\tilde{L}_b} M_{\tilde{L}_s}, \dots$

→ λ^t and λ^b correspond to $(\delta_{LL}^u)_{23}$ and $(\delta_{LL}^d)_{23}$

Mass eigenstates : To diagonalize the 4×4 squark mass matrices, two 4×4 rotation matrices, $R_{\tilde{u}}$ and $R_{\tilde{d}}$, are needed

model.h

Common blocks for the model parameters

- SM parameters
- MSSM parameters
- flavour-violating parameters

mssm-ini.F

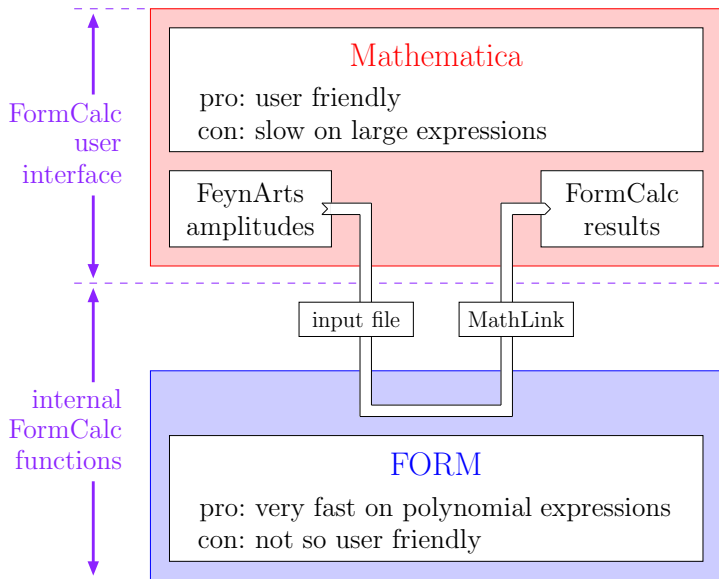
initialize the flavour-violating MSSM
based on FormCalc's mssm-ini.F

- Input parameters
- All parameters appearing in the Model File
- Exclusion limits
- Compute the squark masses:
 - 1) "regular" ones with only left-right mixing (needed for computing the Higgs mass)
 - 2) fully 6x6 generalized mixing matrices

Modified versions are not (yet) public : –(
Checks are being done !

Files available on request !!

How FormCalc works ?



Example: Calculating Higgs self-energy

```
In[1]:= << FormCalc'
```

```
FormCalc 4
by Thomas Hahn
last revised 11 Nov 04
```

```
In[2]:= CalcFeynAmp[<<HiggsSelfEnergy.amp]
```

```
preparing FORM code in /tmp/m6.frm
running FORM... ok
```

```
Out[2]= Amp[{MH} -> {MH}] [
```

$$\frac{\text{Alfa} \left(-\frac{MW^2}{CW^2} + MZ^2\right)}{4 CW^2 \text{Pi} SW^2} -$$

$$\frac{3 \text{Alfa} MH^2 A0[MH^2]}{16 MW^2 \text{Pi} SW^2} -$$

$$\frac{\text{Alfa} \left(3 + \frac{MH^2}{2 MW^2}\right) A0[MW^2]}{4 \text{Pi} SW^2} - \frac{\text{Alfa} \left(\frac{3}{CW^2} + \frac{MH^2}{2MW^2}\right) A0[MZ^2]}{8 \text{Pi} SW^2} +$$

$$\frac{9 \text{Alfa} MH^2^2 B0[MH^2, MH^2, MH^2]}{32 MW^2 \text{Pi} SW^2} +$$

$$\frac{\text{Alfa} \left(-\frac{MH^2}{2} + \frac{MH^2^2}{4 MW^2} + 3 MW^2\right) B0[MH^2, MW^2, MW^2]}{4 \text{Pi} SW^2} +$$

$$\frac{\text{Alfa} \left(-\frac{MH^2}{CW^2} + \frac{MH^2^2}{2MW^2} + \frac{7 MW^2}{CW^2^2} - \frac{MZ^2}{MW^2}\right) B0[MH^2, MZ^2, MZ^2]}{16 \text{Pi} SW^2} +$$

$$\frac{\text{Alfa} MH^2 B1[MH^2, MW^2, MW^2]}{4 \text{Pi} SW^2} + \frac{\text{Alfa} MH^2 B1[MH^2, MZ^2, MZ^2]}{8 CW^2 \text{Pi} SW^2},$$

```
+ Fermionic Contributions
```

Implemented constraints

- Exclusion limits:
lower bound on $m_{\tilde{f}}, m_{\tilde{\chi}}, m_{\tilde{g}}, m_h$
- $\Delta\rho \leq 2 \times 10^{-3}$

Example: Calculating Higgs self-energy

```
In[1]:= << FormCalc'
```

```
FormCalc 4
```

```
by Thomas Hahn
```

```
last revised 11 Nov 04
```

```
In[2]:= CalcFeynAmp[<<HiggsSelfEnergy.amp]
```

```
preparing FORM code in /tmp/m6.frm
```

```
running FORM... ok
```

```
Out[2]= Amp[{MH} -> {MH}] [
```

$$\frac{\text{Alfa} \left(-\left(\frac{M_W^2}{C_W^2}\right) + M_Z^2\right)}{4 C_W^2 \text{Pi} S_W^2} - \frac{3 \text{Alfa} M_H^2 A_0[M_H^2]}{16 M_W^2 \text{Pi} S_W^2} -$$
$$\frac{\text{Alfa} \left(3 + \frac{M_H^2}{2 M_W^2}\right) A_0[M_W^2]}{4 \text{Pi} S_W^2} - \frac{\text{Alfa} \left(\frac{3}{C_W^2} + \frac{M_H^2}{2 M_W^2}\right) A_0[M_Z^2]}{8 \text{Pi} S_W^2} +$$
$$\frac{9 \text{Alfa} M_H^2 B_0[M_H^2, M_H^2, M_H^2]}{32 M_W^2 \text{Pi} S_W^2} +$$
$$\frac{\text{Alfa} \left(\frac{-M_H^2}{2} + \frac{M_H^2^2}{4 M_W^2} + 3 M_W^2\right) B_0[M_H^2, M_W^2, M_W^2]}{4 \text{Pi} S_W^2} +$$
$$\frac{\text{Alfa} \left(-\left(\frac{M_H^2}{C_W^2}\right) + \frac{M_H^2^2}{2 M_W^2} + \frac{7 M_W^2}{C_W^2} - \frac{M_Z^2}{M_W^2}\right) B_0[M_H^2, M_Z^2, M_Z^2]}{16 \text{Pi} S_W^2} +$$
$$\frac{\text{Alfa} M_H^2 B_1[M_H^2, M_W^2, M_W^2]}{4 \text{Pi} S_W^2} + \frac{\text{Alfa} M_H^2 B_1[M_H^2, M_Z^2, M_Z^2]}{8 C_W^2 \text{Pi} S_W^2},$$

```
(
```

$$\frac{\text{Alfa} A_0[M_f^2[2, \text{Gen1}]] M_f^2[2, \text{Gen1}]}{2 M_W^2 \text{Pi} S_W^2}$$
$$\frac{\text{Alfa} M_H^2 B_1[M_H^2, M_f^2[2, \text{Gen1}], M_f^2[2, \text{Gen1}]] M_f^2[2, \text{Gen1}]}{4 M_W^2 \text{Pi} S_W^2}$$
$$\frac{\text{Alfa} B_0[M_H^2, M_f^2[2, \text{Gen1}], M_f^2[2, \text{Gen1}]] M_f^2[2, \text{Gen1}]^2}{2 M_W^2 \text{Pi} S_W^2} +$$
$$\frac{3 \text{Alfa} A_0[M_f^2[3, \text{Gen1}]] M_f^2[3, \text{Gen1}]}{2 M_W^2 \text{Pi} S_W^2}$$
$$\frac{3 \text{Alfa} M_H^2 B_1[M_H^2, M_f^2[3, \text{Gen1}], M_f^2[3, \text{Gen1}]] M_f^2[3, \text{Gen1}]}{4 M_W^2 \text{Pi} S_W^2}$$
$$\frac{3 \text{Alfa} B_0[M_H^2, M_f^2[3, \text{Gen1}], M_f^2[3, \text{Gen1}]] M_f^2[3, \text{Gen1}]^2}{2 M_W^2 \text{Pi} S_W^2} +$$
$$\frac{3 \text{Alfa} A_0[M_f^2[4, \text{Gen1}]] M_f^2[4, \text{Gen1}]}{2 M_W^2 \text{Pi} S_W^2}$$
$$\frac{3 \text{Alfa} M_H^2 B_1[M_H^2, M_f^2[4, \text{Gen1}], M_f^2[4, \text{Gen1}]] M_f^2[4, \text{Gen1}]}{4 M_W^2 \text{Pi} S_W^2}$$
$$\frac{3 \text{Alfa} B_0[M_H^2, M_f^2[4, \text{Gen1}], M_f^2[4, \text{Gen1}]] M_f^2[4, \text{Gen1}]^2}{2 M_W^2 \text{Pi} S_W^2}) \text{SumOver}[\text{Gen1}, 3]$$

Results for

M_W , $\sin^2 \theta_{\text{eff}}$ and m_h

S. Heinemeyer, W. Hollik, F. Merz, S.P.

Eur.Phys.J.C37, 481-493, 2004, hep-ph/0403228

Simplest scenario (Numerical result)

Mixing only between the left-handed components of \tilde{t}, \tilde{c} and \tilde{b}, \tilde{s}

$$\begin{aligned}\Delta_{LL}^t &= \lambda^t M_{\tilde{L}_t} M_{\tilde{L}_c}, & \Delta_{LR}^t &= \Delta_{RL}^t = \Delta_{RR}^t = 0, \\ \Delta_{LL}^b &= \lambda^b M_{\tilde{L}_b} M_{\tilde{L}_s}, & \Delta_{LR}^b &= \Delta_{RL}^b = \Delta_{RR}^b = 0.\end{aligned}$$

The same flavor mixing parameter in the \tilde{t}/\tilde{c} and \tilde{b}/\tilde{s} sectors is assumed : $\lambda = \lambda^t = \lambda^b$

→ $\lambda^t = \lambda^b = 0$ corresponds to the MSSM with MFV

→ λ^t and λ^b correspond to $(\delta_{LL}^u)_{23}$ and $(\delta_{LL}^d)_{23}$

Numerical analysis performed in 2 benchmark scenarios, but with a free scale M_{SUSY} :

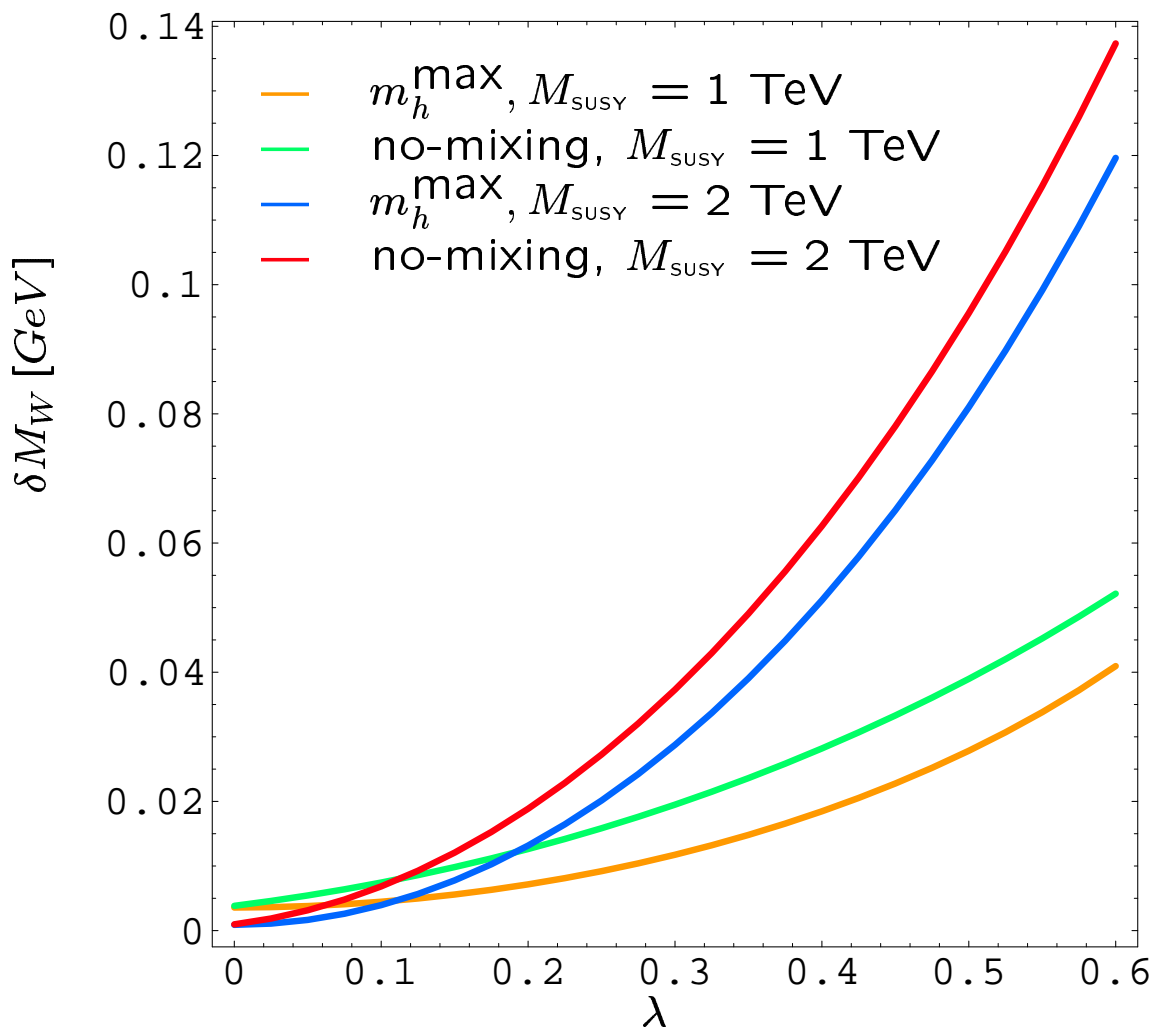
M. Carena et. al, hep-ph/0202167

- m_h^{max} (A_t is not a free parameter, obeying $X_t = 2M_{\text{SUSY}}$, with $X_t = A_t - \mu \cot \beta$)
- **no-mixing** (no mixing in the MFV \tilde{t} sector ($X_t = 0$))

Analytical result:

evaluation with arbitrary NMFV couplings

δM_W as a function of λ :



→ The induced shifts in M_W can become as large as 0.14 GeV for no-mixing, $M_{\text{SUSY}} = 2 \text{ TeV}$, $\lambda = 0.6$.

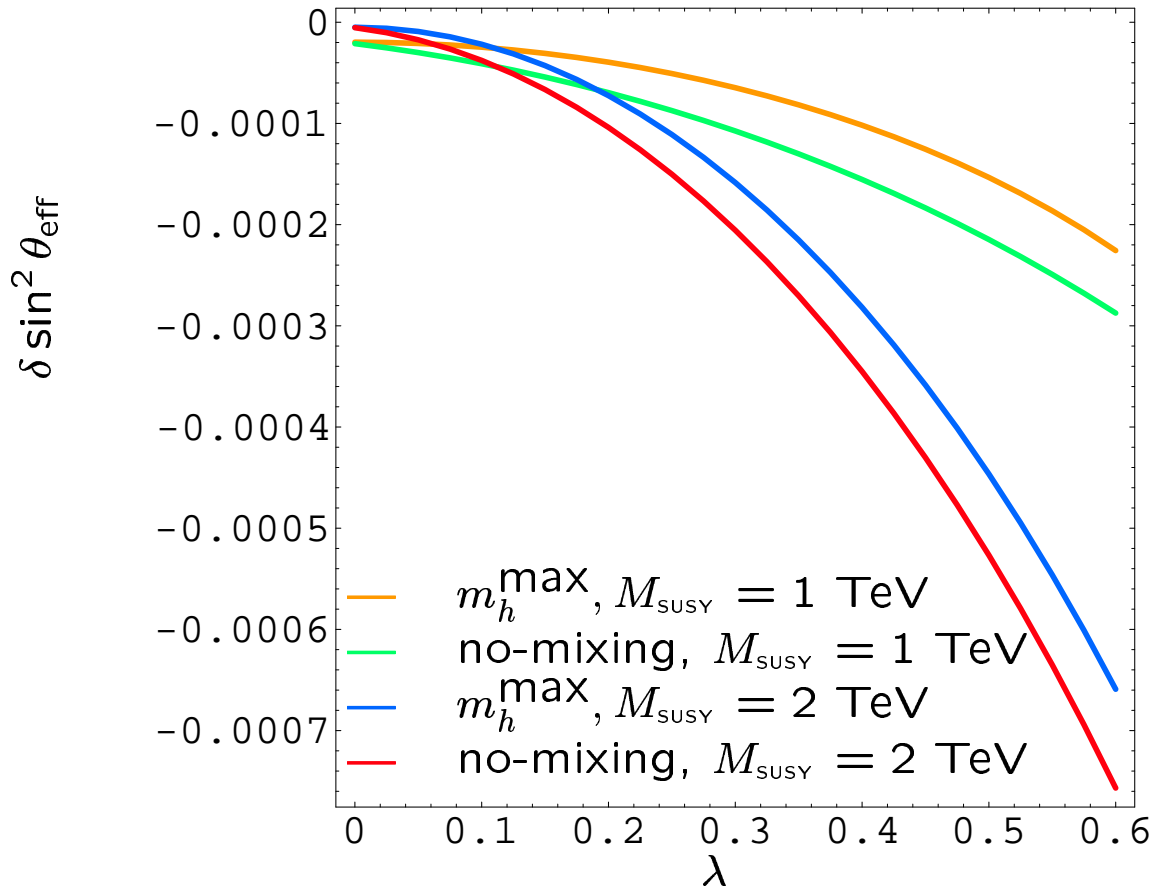
→ $\delta M_W \lesssim 0.05 \text{ GeV}$ in the m_h^{\max} scenario, but still sizeable.

$$\delta M_W^{\text{exp,today}} = 34 \text{ MeV}$$

$$\delta M_W^{\text{exp,future}} = 7 \text{ MeV}$$

⇒ extreme parameter regions already ruled out

$\delta \sin^2 \theta_{\text{eff}}$ as a function of λ :



→ The shifts $\delta \sin^2 \theta_{\text{eff}}$ can reach values up to 7×10^{-4} for no-mixing scenario, $M_{\text{SUSY}} = 2 \text{ TeV}$, $\lambda = 0.6$,

→ smaller, but still sizeable, for the other scenarios.

$$\delta \sin^2 \theta_{\text{eff}}^{\text{exp, today}} = 17 \times 10^{-5}$$

$$\delta \sin^2 \theta_{\text{eff}}^{\text{exp, future}} = 1.3 \times 10^{-5}$$

⇒ extreme parameter regions already ruled out

Higgs mass corrections

Contrary to the SM: m_h is not a free parameter

- Large radiative corrections:

Dominant one-loop corrections: $\sim G_\mu m_t^4 \ln \left(\frac{m_{\tilde{t}_1} m_{\tilde{t}_2}}{m_t^2} \right)$

MSSM with MFV : Dominant one-loop contributions involves third-generation quarks and squarks.

- **MSSM with NMFV**

The squark loops have to be modified by introducing the generation-mixed squarks.

- Measurement of m_h , Higgs couplings

⇒ test of the theory

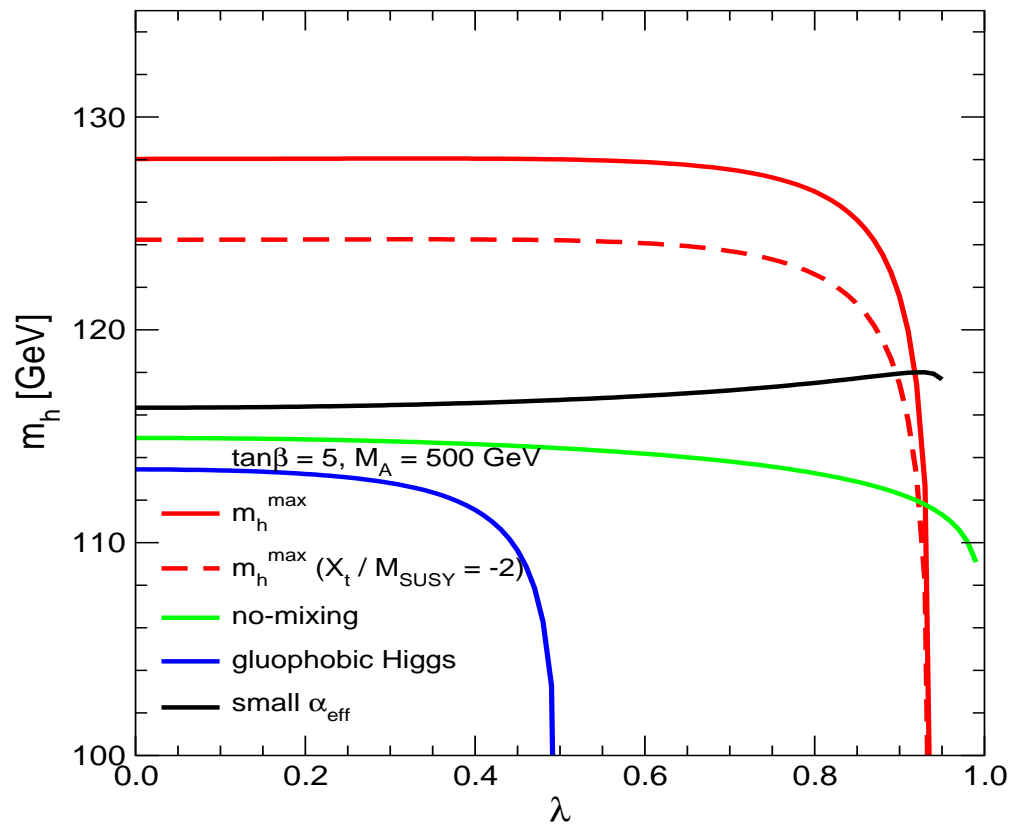
LHC: $\Delta m_h \approx 0.2 \text{ GeV}$

LC: $\Delta m_h \approx 0.05 \text{ GeV}$

Higgs boson sector analysis performed in **5 benchmark scenarios** : *M. Carena et al, hep-ph/0202167*

- m_h^{\max} : to maximize the lightest Higgs boson mass
- **constrained m_h^{\max}** : with $X_t/M_{\text{SUSY}} = -2$ for $b \rightarrow s\gamma$
- **no-mixing** in the MFV \tilde{t} sector
- **gluophobic Higgs** : with reduced ggh coupling
- **small α_{eff}** : with reduced $h\bar{b}b$ and $h\tau^+\tau^-$ coupling

m_h as a function of λ :



$$M_{\text{SUSY}} = M_{\tilde{Q}_q} = M_{\tilde{U}_q} = M_{\tilde{D}_q}, \quad A_s = A_b = A_c = A_t$$

All scenarios show a similar behavior

⇒ small effects for small/moderate λ

⇒ $\delta m_h = \mathcal{O}(5 \text{ GeV})$ only for very large λ

(around 0.5 in the gluophobic Higgs scenario, and around 0.9 in the other four scenarios)

⇒ mostly decreasing m_h , but also increase possible

(in small α_{eff} -scenario it can be enhanced by up to 2 GeV)

⇒ the impact of NMFV on m_h is in general rather small

FEYNHIGGS

FeynHiggs

S. Heinemeyer, W. Hollik, G. Weiglein,
Comput.Phys.Commun. 124 (2000) 76-89, hep-ph/9812320

T. Hahn, S. Heinemeyer, W. Hollik, G. Weiglein,
MPP-2003-147, see: hep-ph/0406152

Latest version:

<http://www.feynhiggs.de> - (Jan 05)

FeynHiggs2.2

- A high precision tool for Higgs physics at colliders
- Provides Higgs boson masses, mixing angles, couplings, branching ratios, ...
in the MSSM with/without complex parameters

AND FOR NMFV

- Additional constraints:
 - $\Delta\rho \leq 2 \times 10^{-3}$, $\Delta\rho$ at $\mathcal{O}(\alpha, \alpha\alpha_s)$
 - $(g_\mu - 2)_{\text{SUSY}}$: Full one-, leading/subleading two-loop SUSY corrections.
 - SUSY Les Houches Accord can also be used
 - FeynHiggs with NMFV MSSM
 - Evaluation of
Higgs masses and the mixing angle α
 - $\Delta\rho$ evaluation is included
- [S. Heinemeyer, W. Hollik, F.Merz, S.P '04]

Subroutines contained in *FeynHiggs2.2*:

- **FHSetFlags**
 - set the flags of the calculation,
- **FHSetPara**
 - set the MSSM input parameters directly, *or*:
FHSetSLHA
 - extract the input parameters from an SLHA data structure,
- **FHGetPara**
 - retrieve (some of) the derived parameters (e.g. chargino masses),
- **FHHiggsCorr**
 - compute the Higgs masses and mixings,
- **FHCouplings**
 - compute the Higgs couplings and BRs,
- **FHConstraints**
 - evaluate $\Delta\rho$ and $(g_\mu - 2)$ constraints.

Conclusions

- FeynHiggs2.2

MSSM with NMFV:

general 4×4 mixing in \tilde{t}/\tilde{c} and \tilde{b}/\tilde{s} sectors

m_h , mixing angle α and $\Delta\rho$ included

- FeynArts, FormCalc and LoopTools

MSSM with NMFV:

6×6 generalized squarks mixing matrices

- Several checks have been done, e.g.:

- Results for M_W , $\sin^2 \theta_{\text{eff}}$ and $\Delta\rho$ reproduced independently

- Results for $BR(H \rightarrow bs, tc)$

- In preparation:

$BR(b \rightarrow s\gamma)$, $B_s \rightarrow \mu^+\mu^-$, ...

[T. Hahn, W. Hollik, J.I. Illana, S.P '05]

[S. Heinemeyer, W. Hollik, J.I. Illana, S.P '05]

Evaluation of M_W , $\sin^2 \theta_{\text{eff}}$, m_h

- **Analytical results:** for arbitrary mixing
Numerical results: only for LL mixing, parametrized with λ ($(\delta_{LL})_{23}$)
- **Large effects** possible for M_W , $\sin^2 \theta_{\text{eff}}$:
 $\lambda \lesssim 0.2 \Rightarrow \delta M_W \lesssim 20 \text{ MeV}$
 $\lambda \lesssim 0.2 \Rightarrow \delta \sin^2 \theta_{\text{eff}} \lesssim 10^{-4}$
→ We have shown that the effects of scalar quark generation mixing enters essentially through $\Delta\rho$
- **Moderate effects** possible for m_h only for large λ
- **Scenarios for B physics can now be tested**, whether they are compatible with electroweak precision observables
(where the effects are large)
and with Higgs physics
(where the effects are small)