

# **Sparticle Masses in Deflected Mirage Mediation**

KC, Jeong, Nakamura, Okumura, Yamaguchi  
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# Flavor and CP conserving mediation of SUSY breaking

## A. Dilaton or volume modulus mediation:

$$\int d^2\theta \frac{1}{4} T W^{a\alpha} W_\alpha^a + \int d^4\theta (T + T^*)^{n_i} Q_i^* Q_i \quad (T = T_0 + F^T \theta^2)$$

$$M_a = \frac{F^T}{T_0 + T_0^*}, \quad m_i^2 = n_i \left| \frac{F^T}{T_0 + T_0^*} \right|^2, \quad A_{ijk} = (n_i + n_j + n_k) \frac{F^T}{T_0 + T_0^*}$$

( $\ni$  Scherk-Schwarz breaking or radion mediation,  
No-scale model or Gaugino mediation:  $n_i = 0$ )

- Flavor conserving as the moduli weights  $n_i$  are flavor-universal rational numbers.
- CP conserving due to the axionic shift symmetry of  $T$ .

In flux compactification, other moduli, e.g. complex structure moduli having flavor-non-universal couplings, can be naturally decoupled from SUSY breaking.

## B. Gauge mediation:

$$\int d^2\theta \lambda X \Phi^c \Phi$$

( $\Phi, \Phi^c$  = Gauge-charged messenger,  $X = X_0 + F^X \theta^2$ )

$$M_a \sim m_i \sim \frac{1}{8\pi^2} \frac{F^X}{X_0}$$

## C. Anomaly mediation:

$C \equiv$  SUGRA Compensator =  $C_0 + F^C$

$$M_a \sim m_i \sim A_{ijk} \sim \frac{1}{8\pi^2} \frac{F^C}{C_0}$$

So far, phenomenological studies of low energy SUSY have focused mostly on the case that one of these mediations dominate others.

However, it is a plausible possibility that all (or two of) three mediations give comparable contributions to the MSSM soft masses, while preserving flavor and CP:

$$\int d^4\theta CC^* \left[ (T + T^*)^{n_c} + (T + T^*)^{n_x} XX^* + (T + T^*)^{n_z} ZZ^* \right] \\ + \int d^2\theta C^3 \left[ A_1 M_{Pl}^3 e^{-8\pi^2 T/N_1} + A_2 \frac{X^{N_2}}{M_{Pl}^{N_2-3}} + \Delta W(Z) \right] \quad (N_2 > 3)$$

( $\Delta W(Z)$  provides SUSY breaking and nearly vanishing C.C.)

- For generic  $A_{1,2}$  of  $\mathcal{O}(1)$ , and  $\Delta W$  giving  $m_{3/2} \ll M_{Pl}$ ,

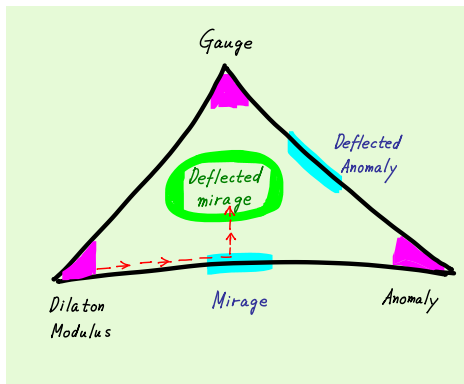
$$8\pi^2 \frac{F^T}{T_0 + T_0^*} \sim \frac{F^X}{X_0} \sim \frac{F^C}{C_0}$$

- The phases of  $F^T$ ,  $F^C$  and  $F^X$  are dynamically aligned as

$$\text{Arg} \left( \frac{F^C}{C_0} \right) = \text{Arg} \left( \frac{F^T}{T_0 + T_0^*} \right) = \text{Arg} \left( \frac{F^X}{X_0} \right) : \text{axion mechanism}$$

# Sparticle Masses in Deflected Mirage Mediation

Deflected mirage mediation, being a general mixed gravity-gauge-anomaly mediation, provides a framework for more general but still theoretically well-motivated pattern of the MSSM soft parameters, which might be useful for the interpretation of experimentally measured sparticle masses.



**Dilaton/Modulus (mSUGRA)  $\longrightarrow$  Mirage  $\longrightarrow$  Deflected Mirage**

# Gaugino masses and light generation sfermion masses

- Gaugino mass ratios are the least sensitive to the other details of the model such as extra matters and/or extra interactions.
- 3rd generation sfermion masses and Higgs masses might severely depend on how  $\mu$  and  $B$  are generated.

## mSUGRA (Dilaton/Modulus Domination)

At  $M_{GUT}$ ,

$$M_a = M_0, \quad m_i^2 = m_0^2 = n_0 M_0^2 \quad \left( M_0 \equiv \frac{F^T}{T_0 + T_0^*} \right)$$

At TeV,

- $M_{\tilde{B}} : M_{\tilde{W}} : M_{\tilde{g}} \simeq 1 : 2 : 6$
- $m_{\tilde{q}_L}^2 : m_{\tilde{u}_R}^2 : m_{\tilde{d}_R}^2 : m_{\tilde{\ell}_L}^2 : m_{\tilde{e}_R}^2$   
 $\simeq (n_0 + 5.0) : (n_0 + 4.6) : (n_0 + 4.5) : (n_0 + 0.5) : (n_0 + 0.15)$

## Mirage $\equiv$ Dilaton/Modulus + Anomaly

KC, Falkowski, Nilles, Olechowski, Pokorski

$$\alpha \equiv \frac{F^C/C_0}{F^T/(T_0 + T_0^*)} \frac{1}{\ln(M_{Pl}/m_{3/2})}$$

$$\left( \alpha \simeq 1 \quad \text{for } W = A_1 e^{-8\pi^2 T/N_1} : \text{KKLT} \right)$$

At TeV,

- $M_{\tilde{B}} : M_{\tilde{W}} : M_{\tilde{g}} \simeq (1 + 0.66\alpha) : (2 + 0.2\alpha) : (6 - 1.8\alpha)$
- $m_{\tilde{q}_L}^2 : m_{\tilde{u}_R}^2 : m_{\tilde{d}_R}^2 : m_{\tilde{\ell}_L}^2 : m_{\tilde{e}_R}^2$   
 $\simeq (n_0 + 5.0 - 3.5\alpha + 0.5\alpha^2) : (n_0 + 4.6 - 3.3\alpha + 0.5\alpha^2) :$   
 $(n_0 + 4.5 - 3.3\alpha + 0.5\alpha^2) : (n_0 + 0.5 - 0.2\alpha - 0.01\alpha^2) :$   
 $(n_0 + 0.15 - 0.05\alpha - 0.01\alpha^2)$

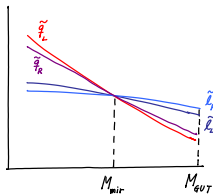
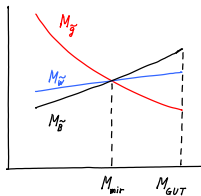
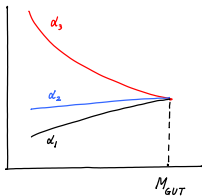
# Mirage unification of soft masses KC, Jeong, Okumura

$$\frac{1}{g_a^2(\mu)} = \frac{1}{g_{GUT}^2} - \frac{b_a}{8\pi^2} \ln\left(\frac{\mu}{M_{GUT}}\right)$$

$$M_a(\mu) = M_0 + \frac{b_a}{8\pi^2} M_0 g_a^2(\mu) \ln\left(\frac{\mu}{M_{mir}}\right)$$

$$m_i^2(\mu) = m_0^2 - \frac{1}{4\pi^2} \gamma_i(\mu) M_0^2 \ln\left(\frac{\mu}{M_{mir}}\right) - \frac{1}{8\pi^2} \dot{\gamma}_i(\mu) M_0^2 \left[ \ln\left(\frac{\mu}{M_{mir}}\right) \right]^2$$

$$\left( M_{mir} = M_{GUT} \left( \frac{m_{3/2}}{M_{Pl}} \right)^{\alpha/2} \right)$$





# Deflection from mirage unification due to gauge mediation

Everett, Kim, Ouyang, Zurek

Add “gauge mediation” with

$$\frac{F^X}{X_0} \sim \frac{F^C}{C_0} \sim \frac{1}{8\pi^2} \frac{F^T}{T_0 + T_0^*}, \quad N_\Phi \equiv \# \text{ of } \Phi(5) + \Phi^c(\bar{5})$$

Effects of gauge mediation on sparticle masses at TeV:

KC, Jeong, Nakamura, Okumura, Yamaguchi

1) Renormalize the mirage parameters  $\alpha$  and  $n_0 = m_i^2(M_{\text{mir}})/M_a^2(M_{\text{mir}})$ :

$$\alpha \rightarrow \alpha_{\text{eff}} \equiv \frac{\alpha}{(1 + \epsilon)}, \quad n_0 \rightarrow n_0^{\text{eff}} \equiv \frac{n_0}{(1 + \epsilon)^2}$$

$$\epsilon = \frac{N_\Phi}{4} \left[ \frac{-F^X/X_0}{8\pi^2 F^T/(T_0 + T_0^*)} - \frac{1}{4\pi^2} \ln \left( \frac{M_{GUT}}{X_0} \right) \right]$$

$$M_{\text{mir}}^{\text{eff}} = M_{GUT} \left( \frac{m_{3/2}}{M_{Pl}} \right)^{\alpha_{\text{eff}}/2}$$

## 2) Deflection of sfermion masses from the mirage pattern:

- $M_{\tilde{B}} : M_{\tilde{W}} : M_{\tilde{g}} \simeq (1 + 0.66\alpha_{\text{eff}}) : (2 + 0.2\alpha_{\text{eff}}) : (6 - 1.8\alpha_{\text{eff}})$
- $m_{\tilde{q}_L}^2 : m_{\tilde{u}_R}^2 : m_{\tilde{d}_R}^2 : m_{\tilde{\ell}_L}^2 : m_{\tilde{e}_R}^2$

$$\simeq (n_0^{\text{eff}} + 5.0 - 3.5\alpha_{\text{eff}} + 0.5\alpha_{\text{eff}}^2 + \delta_{\tilde{q}_L}) :$$

$$(n_0^{\text{eff}} + 4.6 - 3.3\alpha_{\text{eff}} + 0.5\alpha_{\text{eff}}^2 + \delta_{\tilde{u}_R}) :$$

$$(n_0^{\text{eff}} + 4.5 - 3.3\alpha_{\text{eff}} + 0.5\alpha_{\text{eff}}^2 + \delta_{\tilde{d}_R}) :$$

$$(n_0^{\text{eff}} + 0.5 - 0.2\alpha_{\text{eff}} - 0.01\alpha_{\text{eff}}^2 + \delta_{\tilde{\ell}_L}) :$$

$$(n_0^{\text{eff}} + 0.15 - 0.05\alpha_{\text{eff}} - 0.01\alpha_{\text{eff}}^2 + \delta_{\tilde{e}_R})$$

$$\delta_i = \frac{2\epsilon^2}{(1+\epsilon)^2} \sum_a C_2^a(\Phi_i) \left[ \frac{4g_a^4(X_0)}{N_\Phi} - \frac{g_a^2(X_0)}{4\pi^2} \left( \frac{1+2\epsilon}{2\epsilon} + g_a^2(X_0) \right) \ln \left( \frac{M_{GUT}}{X_0} \right) \right]$$

$\delta_i - \delta_j$  represent the true deflection from the mirage pattern, and their relative importance depends on the sign of  $F^X/F^T$ .

For  $F^X/X_0 F^T < 0$ , which is when  $X$  is stabilized by  $X^{N_2}/M_{Pl}^{N_2-3}$ , the deflection is not significant:

$$|\delta_{\tilde{q}_L} - \delta_{\tilde{e}_R}| \lesssim 0.02N_\Phi, \quad |\delta_{\tilde{d}_R} - \delta_{\tilde{\ell}_L}| \lesssim 0.01N_\Phi \quad \text{for } \alpha \simeq 1, 4 \leq N_2 \leq 6$$

## Conclusion

Deflected mirage mediation provides a framework for quite general but still theoretically well-motivated pattern of the MSSM soft parameters, which might be useful for the interpretation of experimentally measured sparticle masses.