

Computation of QCD cross sections of large multiplicity final states at NLO

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 - need for NLO corrections of high multiplicity final state processes at the LHC
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- **Latter**
 - is good for making improvements
 - could lead to unnecessary repetitions (waste)

Workshops could be useful to reduce waste
(at the end efficient and correct codes are needed)

Cross sections at NLO

$$\sigma = \sigma^{\text{LO}} + \sigma^{\text{NLO}} + \dots$$

Born:

$$\sigma^{\text{LO}} = \int_m d\sigma_m^{\text{B}} = \int d\Phi_m \left| M_m^{(0)} \right|^2 J_m$$

NLO:

$$\begin{aligned} \sigma^{\text{NLO}} &= \int_m d\sigma_m^{\text{V}} + \int_{m+1} d\sigma_{m+1}^{\text{R}} \\ &= \int d\Phi_m 2 \operatorname{Re} \left(M_m^{(0)*} M_m^{(0)} \right) J_m + \int d\Phi_{m+1} \left| M_{m+1}^{(0)} \right|^2 J_{m+1} \end{aligned}$$

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- Most important messages for tree level:
 - solutions to efficient PS generation (PHEGAS, MADEVENT)
 - recursive approach to amplitudes:
BG is most efficient for large m
 - MC summation over helicity is useful
 - MC summation over colour is the ‘only practical’ way for ‘colourful’ processes

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 - generalized unitarity
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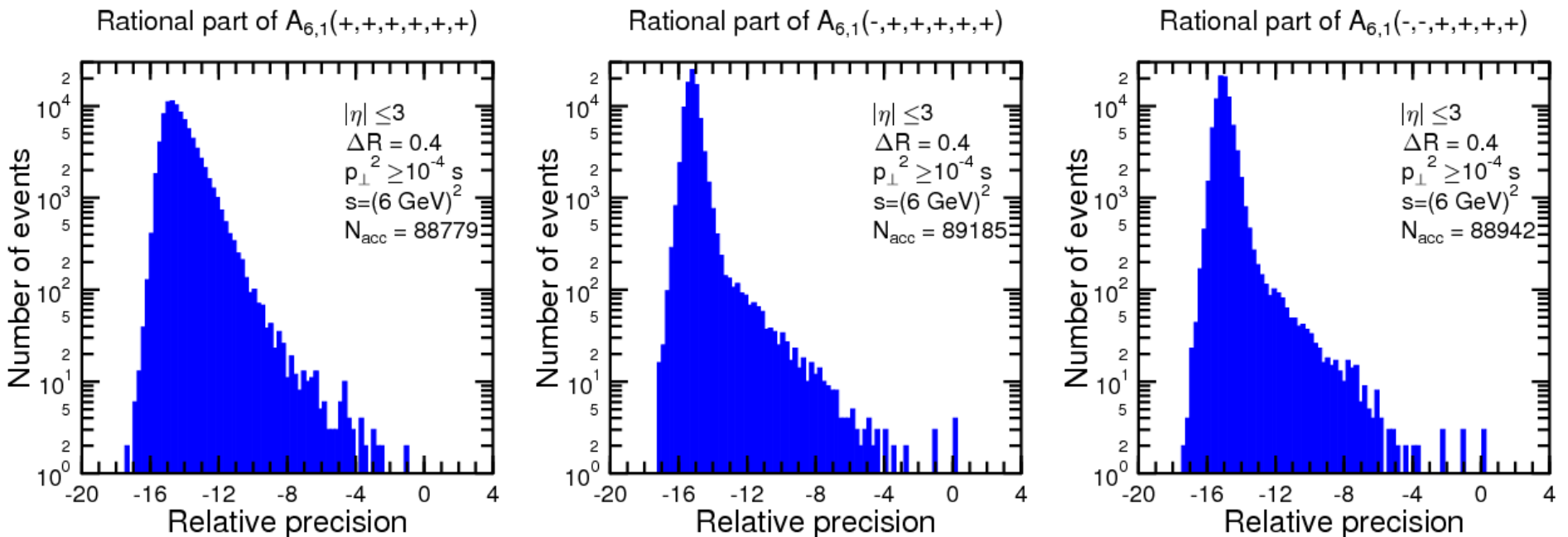
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- Essentially
 - same for cut-constructible (**C**) parts
 - different in computing the rational (**R**) parts

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Precision plots for six-gluon amplitudes

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 - dipole (CS+DT)
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- Majority of calculations used dipole subtractions in the last decade (except MC@NLO), but it
 - cannot directly be generalized to NNLO (⇒ antennae?)
 - inconvenient for matching with PS
 - becomes inefficient for multileg processes due to proliferation of dipoles (⇒ FKS?)

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- Easy separation of soft & collinear singularities (to all orders)
- General subtraction scheme at NNLO
- New scheme at NLO that solves previous objections simultaneously
 - minimum number of subtraction terms at NLO (as FKS)
 - uses exact PS factorisations/convolutions (similarly to CS)
 - permutation symmetry of the SME (if any) is respected (as FKS)
 - works at any order in PT