

# A Realistic Unified Gauge Coupling from the Micro-Landscape of Orbifold GUTs

Arthur Hebecker (Heidelberg)

Based on work with Christian Gross (0812.4267 [hep-ph])

and on earlier idea with Gero von Gersdorff [hep-th/0504002]

## General Setting

- Orbifold GUTs (coming e.g. from anisotropic heterotic orbifolds)
- Size of 5th Dimension is GUT- rather than string-scale
- This size sets the 4d unified gauge coupling
- **Question: What determines this ‘largish’ size of the 5th dimension?**

- 5d to 4d compactification:

$$g_4^2 = \frac{g_5^2}{2\pi R}.$$

- 1- and 2-loop Casimir energy generate a radion effective potential

$$V(R) \sim \frac{1}{R^4} + \frac{g_5^2}{R^5}.$$

(one may think of this as being due to Kähler corrections to non-scale Kähler potential for  $T = R + \dots$ , cf. Luty, Okada, '02).

- Brane-localized operators lead to log-enhancement:

$$V(R) \sim \frac{1}{R^4} + \frac{g_5^2}{R^5} \ln(g_5^2/R).$$

- It only takes a small numerical accident to make  $R_{\min}$  largish, so that  $\alpha_{\text{GUT}} \simeq 1/25$ .
- The freedom of ‘distributing matter between bulk and branes’ (i.e. untwisted and twisted matter) easily allows for this ‘accident’.
- **Indeed, 12 of the 256 models of our ‘micro landscape’ have an  $\alpha_{\text{GUT}}$  between 1/20 and 1/30.**

**Final Comment:**

**We also have an interesting proposal for ‘uplifting’ our AdS vacuum using small 5d warping (cf. Bagger,Belyaev,’02; Bagger,Redi,’03; Falkowski,’05).**

# Phenomenology of Supersymmetric Gauge-Higgs Unification

Based on work with

Felix Brümmer, Sylvain Fichtel, Sabine Kraml (0906.2957 [hep-ph])

and on earlier work with

John March-Russell and Robert Ziegler (0801.4101 [hep-ph])

## General Setting

- As before: Orbifold GUT, SUSY is broken by  $F_T$  and  $F_\phi$
- Further assumption: Gauge-Higgs unification based on SU(6)
- Crucial observation:  $35 = 24 + 5 + \bar{5} + 1$   
(cf. Burdman, Nomura, 2002)

# Known Facts

(Choi/Haba/Jeong/Okumura/Shimizu/Yamaguchi 2004;

cf. also Lopez-Cardoso et al., '94; Antoniadis et al., '94; Brignole et al., '95)

- **5d action in terms of  $\mathcal{N} = 1$  superfields, coupled to supergravity à la Marti/Pomarol, contains terms**

$$\int d^2\theta T \operatorname{tr} W^2 \quad , \quad \int d^4\theta \bar{\varphi}\varphi \frac{\operatorname{tr}(\Phi + \bar{\Phi})^2}{T + \bar{T}} ,$$

**which implies:**

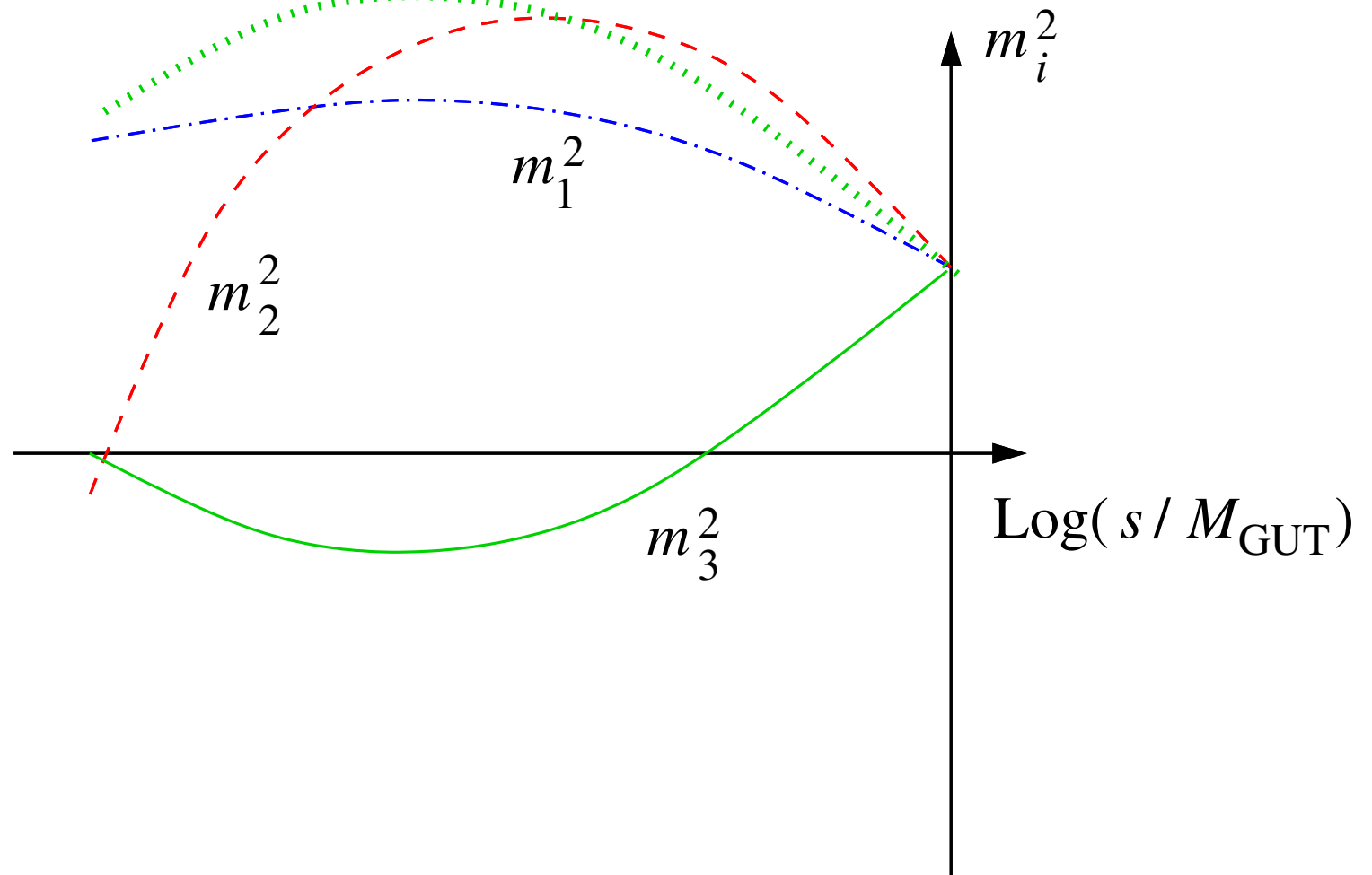
$$M_{1/2} = \frac{\bar{F}_T}{2R} \quad , \quad \mu = \bar{F}_\varphi - \frac{\bar{F}_T}{2R} \quad , \quad m_i^2 = |F_\varphi|^2 - \frac{F_\varphi \bar{F}_T + \text{h.c.}}{2R}$$

**(our conventions:  $m_{1,2}^2 \equiv |\mu|^2 + m_{H_{d,u}}^2$  and  $m_3^2 \equiv B\mu$ ).**

**This is inconsistent with realistic low-energy phenomenology!**

**(without severe fine-tuning)**

Running of  $m_{1,2}^2 = |\mu|^2 + m_{H_{d,u}^2}^2$  and  $m_3^2 = B\mu$



# Our Suggestion

- The SUSY CS-term (**generically present!**) corrects

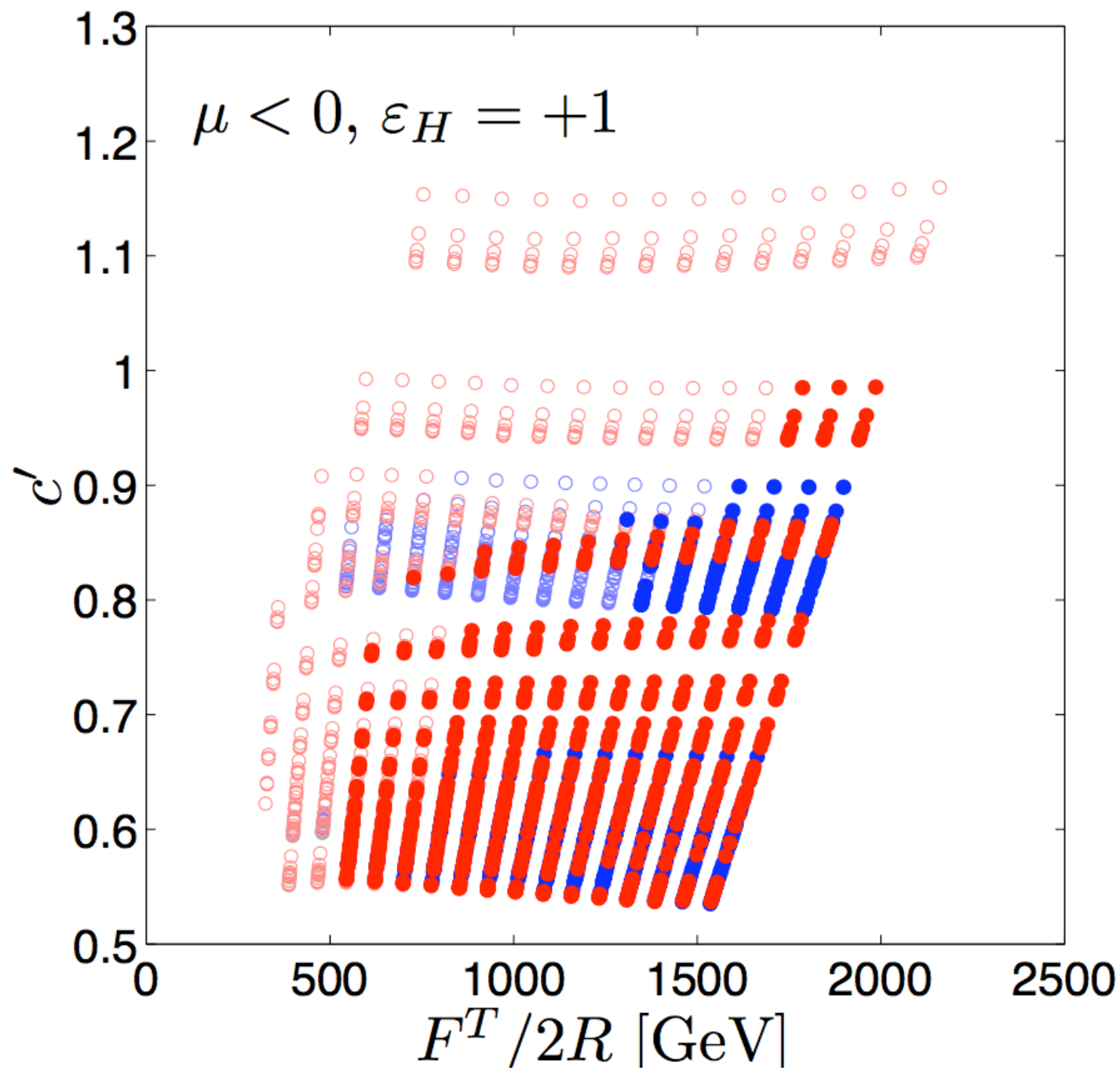
$$\int d^2\theta T \operatorname{tr} W^2 \quad , \quad \int d^4\theta \bar{\varphi}\varphi \frac{\operatorname{tr}(\Phi + \bar{\Phi})^2}{T + \bar{T}}$$

by

$$\int d^2\theta \operatorname{tr} \Phi W^2 \quad , \quad \int d^4\theta \bar{\varphi}\varphi \frac{\operatorname{tr}(\Phi + \bar{\Phi})^3}{(T + \bar{T})^2} .$$

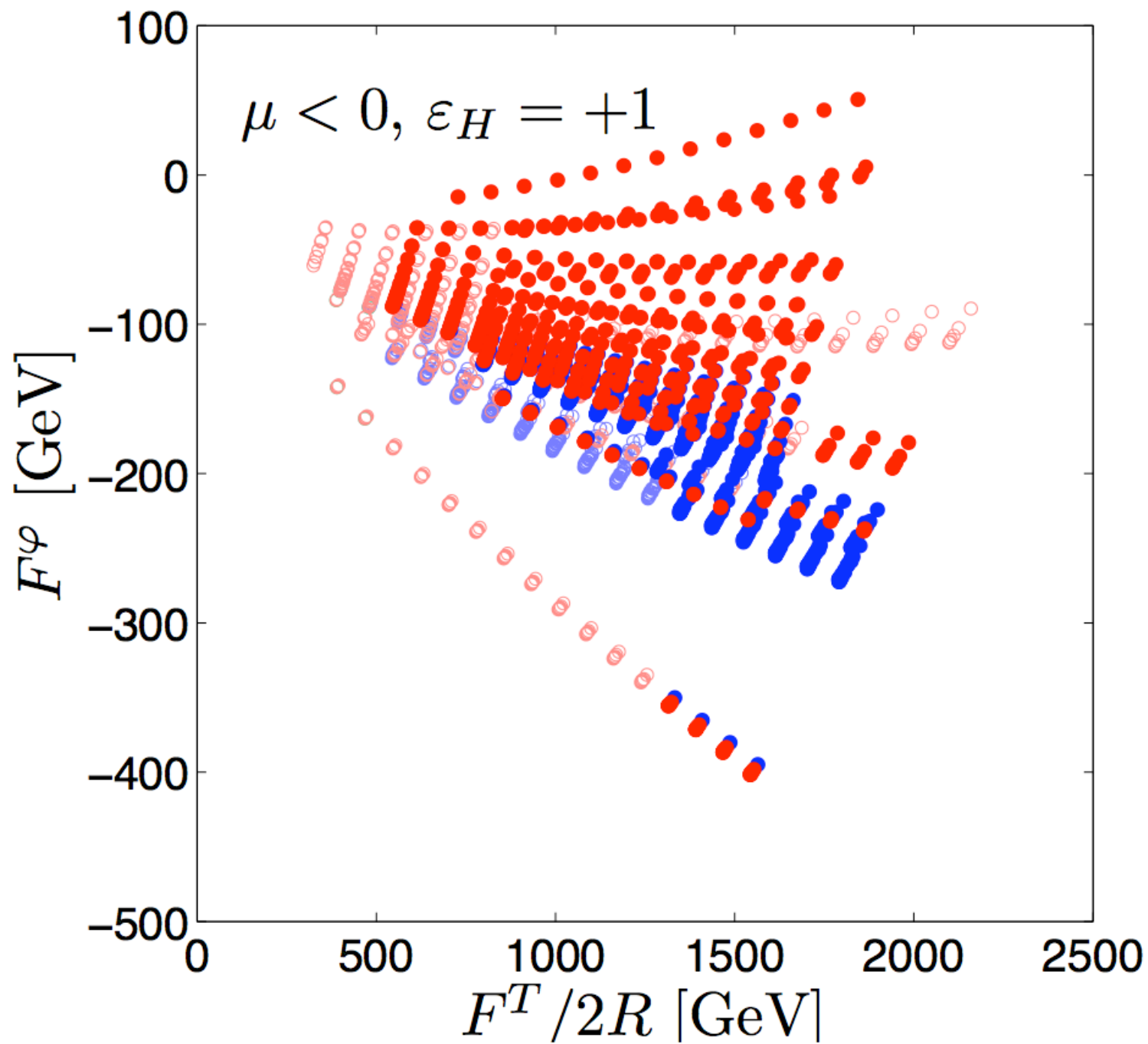
- Due to the extra parameter ( $c' \equiv$  **coefficient of the CS term**) one finds large regions with excellent phenomenology.
- Similar to HENS; cf. Evans, Morrissey, Wells, '06

# Importance of Chern-Simons term





# $F$ term ratios



# Neutralino and slepton masses

