



# Neutrino Physics: Present and Future

Boris Kayser

CERN

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A deep field image of the universe, showing a vast field of galaxies and stars. The background is a dark, starry sky filled with numerous galaxies of various shapes and colors, including spiral, elliptical, and irregular forms. The galaxies are scattered across the field, with some appearing as bright, distinct objects and others as faint, diffuse clouds. The overall scene is a rich, multi-colored tapestry of cosmic structures.

# The Open Questions

- What is the absolute scale of neutrino mass?
- Are neutrinos their own antiparticles?
- Are there “sterile” neutrinos?

**We must be alert to surprises!**

- What is the pattern of mixing among the different types of neutrinos?

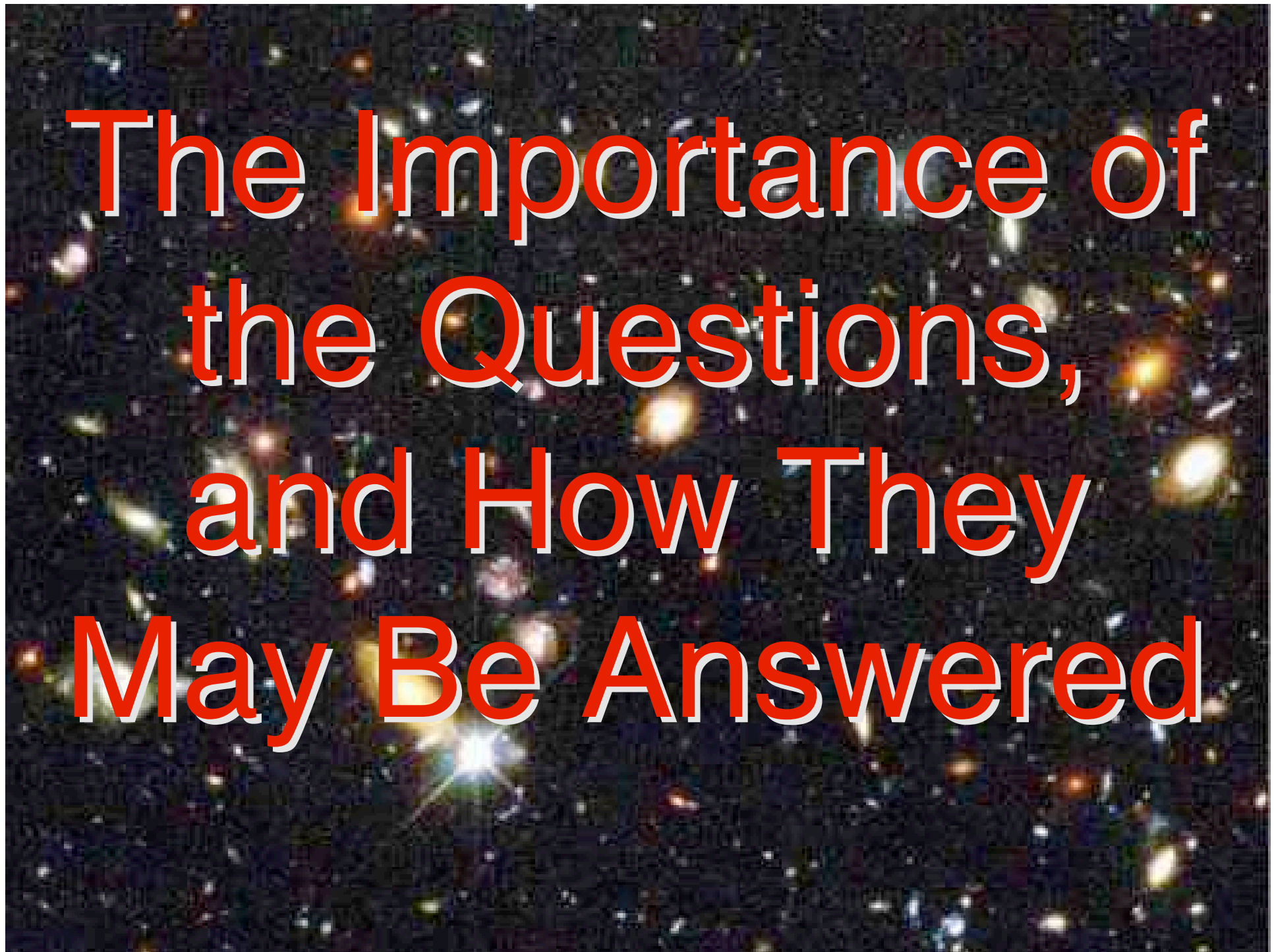
What is  $\theta_{13}$ ? Is  $\theta_{23}$  maximal?

- Is the spectrum like  $\underline{=}$  or  $\underline{=}$  ?

- Do neutrinos violate the symmetry CP?

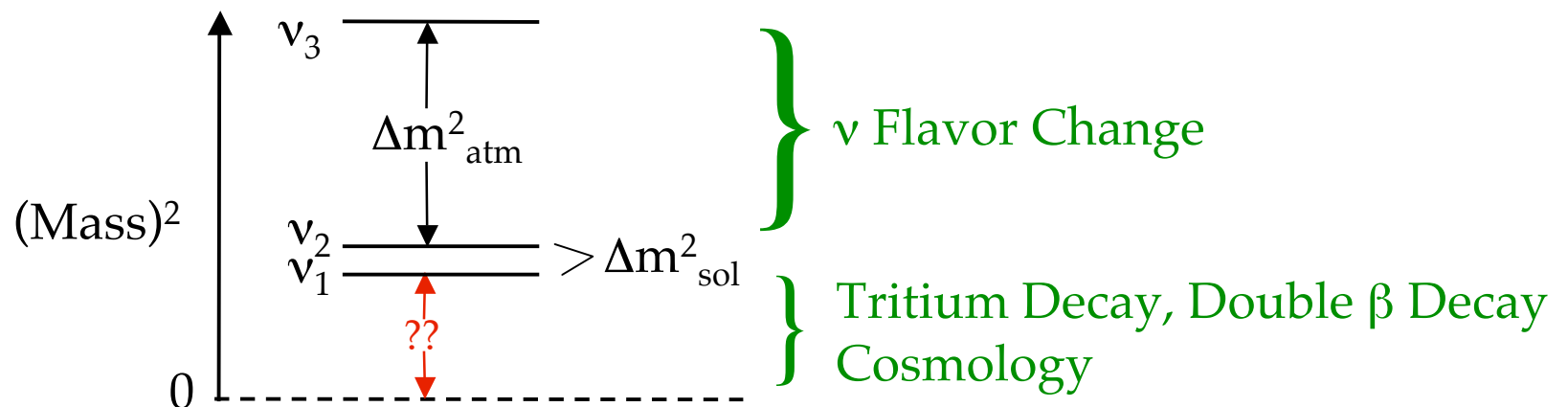
Is  $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\nu_\alpha \rightarrow \nu_\beta)$  ?

- What can neutrinos and the universe tell us about one another?
- Is CP violation by neutrinos the key to understanding the matter – antimatter asymmetry of the universe?
- What physics is behind neutrino mass?



# The Importance of the Questions, and How They May Be Answered

# What Is the Absolute Scale of Neutrino Mass?



How far above zero  
is the whole pattern?

# A Cosmic Connection

Oscillation Data  $\Rightarrow \sqrt{\Delta m^2_{\text{atm}}} < \text{Mass}[\text{Heaviest } \nu_i]$

Cosmological Data + **Cosmological Assumptions**  $\Rightarrow$

$$\Sigma m_i < (0.17 - 1.0) \text{ eV} .$$

Mass( $\nu_i$ )  $\uparrow$  ( Seljak, Slosar, McDonald )  
Pastor

If there are only **3** neutrinos,

$$0.04 \text{ eV} \lesssim \text{Mass}[\text{Heaviest } \nu_i] < (0.07 - 0.4) \text{ eV}$$

$\sqrt{\Delta m^2_{\text{atm}}}$   $\uparrow$

Cosmology  $\uparrow$



Are Neutrinos Their  
Own Antiparticles?

# What Is the Question?

For each *mass eigenstate*  $\nu_i$ , does —

- $\bar{\nu}_i = \nu_i$  (Majorana neutrinos)

or

- $\bar{\nu}_i \neq \nu_i$  (Dirac neutrinos) ?

Equivalently, is the **Lepton Number**  $L$  defined by—

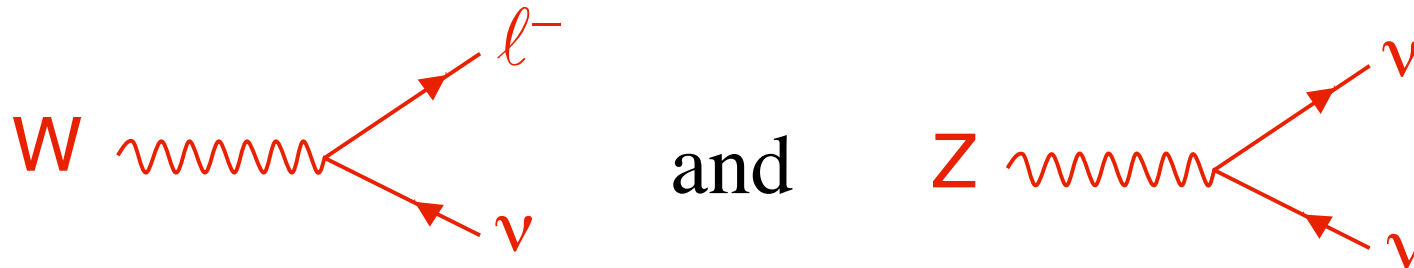
$$L(\nu) = L(\ell^-) = -L(\bar{\nu}) = -L(\ell^+) = 1 \text{ conserved?}$$

If not, then nothing distinguishes  $\bar{\nu}_i$  from  $\nu_i$ . We then have Majorana neutrinos.



**How Can the Standard  
Model be Modified to  
Include Neutrino Masses?**

# The S(tandard) M(odel)



couplings conserve the **Lepton Number L** defined by—

$$L(\nu) = L(l^-) = -L(\bar{\nu}) = -L(l^+) = 1.$$

So do the Dirac charged-lepton mass terms

$$m_\ell \bar{l}_L l_R \quad \longrightarrow \quad \begin{array}{c} \xrightarrow{l^{(\mp)}} \quad \times \quad \xrightarrow{l^{(\mp)}} \\ m_\ell \end{array}$$

- Original SM:  $m_\nu = 0$ .
- Why not add a **Dirac** mass term,

$$m_D \bar{\nu}_L \nu_R$$


Then everything conserves  $L$ , so for each mass eigenstate  $\nu_i$ ,

$$\bar{\nu}_i \neq \nu_i \quad (\text{Dirac neutrinos})$$


$$[L(\bar{\nu}_i) = -L(\nu_i)]$$

- The SM contains no  $\nu_R$  field, only  $\nu_L$ .

To add the Dirac mass term, we had to add  $\nu_R$  to the SM.

Unlike  $\nu_L$ ,  $\nu_R$  carries no Electroweak Isospin.

Thus, no SM principle prevents the occurrence of the **Majorana** mass term

$$m_R \overline{\nu_R^c} \nu_R$$


But this does not conserve L, so now

$$\bar{\nu}_i = \nu_i \quad (\text{Majorana neutrinos})$$

[No conserved L to distinguish  $\bar{\nu}_i$  from  $\nu_i$ ]

We note that  $\bar{\nu}_i = \nu_i$  means —

$$\bar{\nu}_i(\mathbf{h}) = \nu_i(\mathbf{h})$$



The objects  $\nu_R$  and  $\overline{\nu_R^c}$  in  $m_R \overline{\nu_R^c} \nu_R$  are not the mass eigenstates, but just the neutrinos in terms of which the model is constructed.

$m_R \overline{\nu_R^c} \nu_R$  induces  $\nu_R \leftrightarrow \overline{\nu_R^c}$  mixing.

As a result of  $K^0 \leftrightarrow \overline{K^0}$  mixing, the neutral K mass eigenstates are —

$$K_{S,L} \cong (K^0 \pm \overline{K^0})/\sqrt{2} .$$

As a result of  $\nu_R \leftrightarrow \overline{\nu_R^c}$  mixing, the neutrino mass eigenstate is —

$$\nu_i = \nu_R + \overline{\nu_R^c} = \text{“} \nu + \overline{\nu} \text{”} .$$

# Many Theorists Expect Majorana Masses

The Standard Model (SM) is defined by the fields it contains, its **symmetries** (notably Electroweak Isospin Invariance), and its renormalizability.

Leaving neutrino masses aside, anything allowed by the SM symmetries occurs in nature.

If this is also true for neutrino masses, then neutrinos have Majorana masses.



- The presence of Majorana masses
- $\bar{\nu}_i = \nu_i$  (Majorana neutrinos)
- L not conserved

— are all equivalent

**Any one implies the other two.**

(Recent work: Hirsch, Kovalenko, Schmidt)

To Determine If  
Neutrinos Are  
Majorana Particles

# The Promising Approach — Neutrinoless Double Beta Decay [ $0\nu\beta\beta$ ]

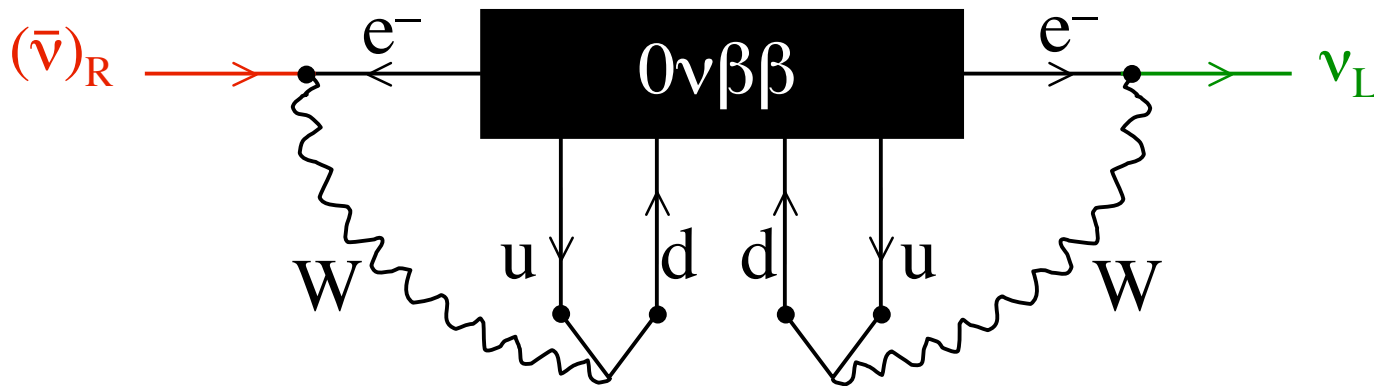


If we start with *a lot* of parent nuclei (say, one ton of them), we can cope with the smallness of  $\mathcal{X}$ .

Observation would imply  $\mathcal{X}$  and  $\bar{\nu}_i = \nu_i$ .

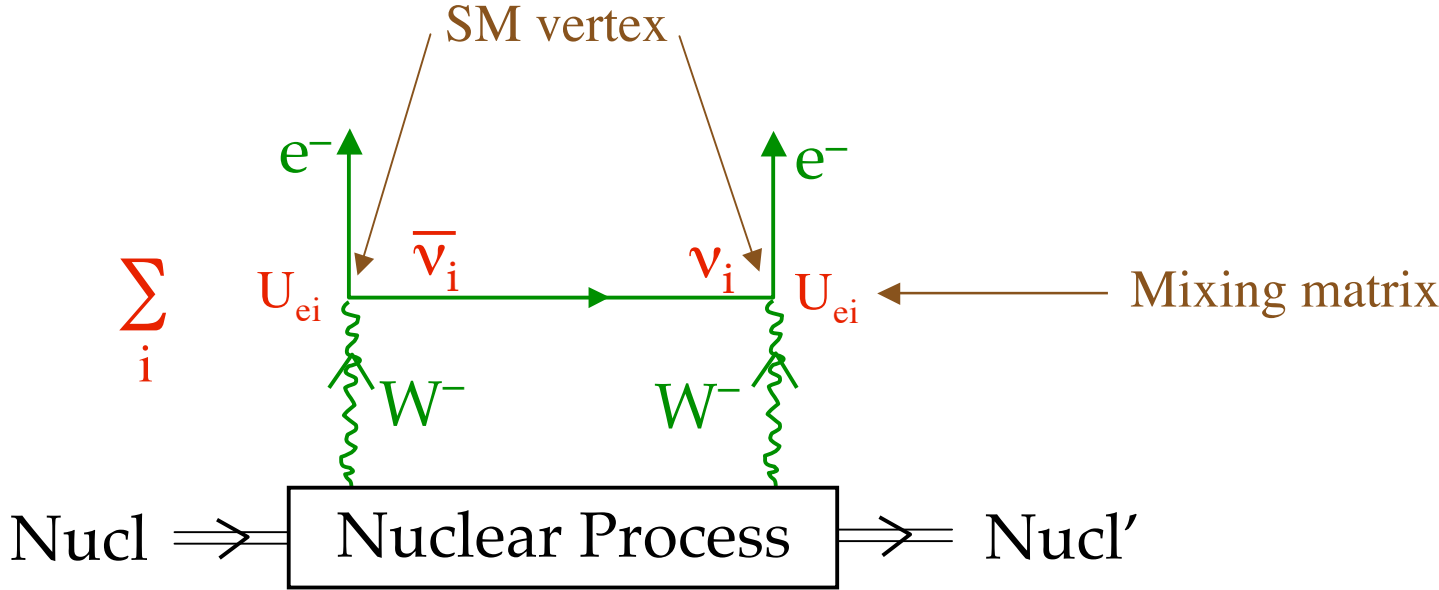
Whatever diagrams cause  $0\nu\beta\beta$ , its observation would imply the existence of a Majorana mass term:

Schechter and Valle

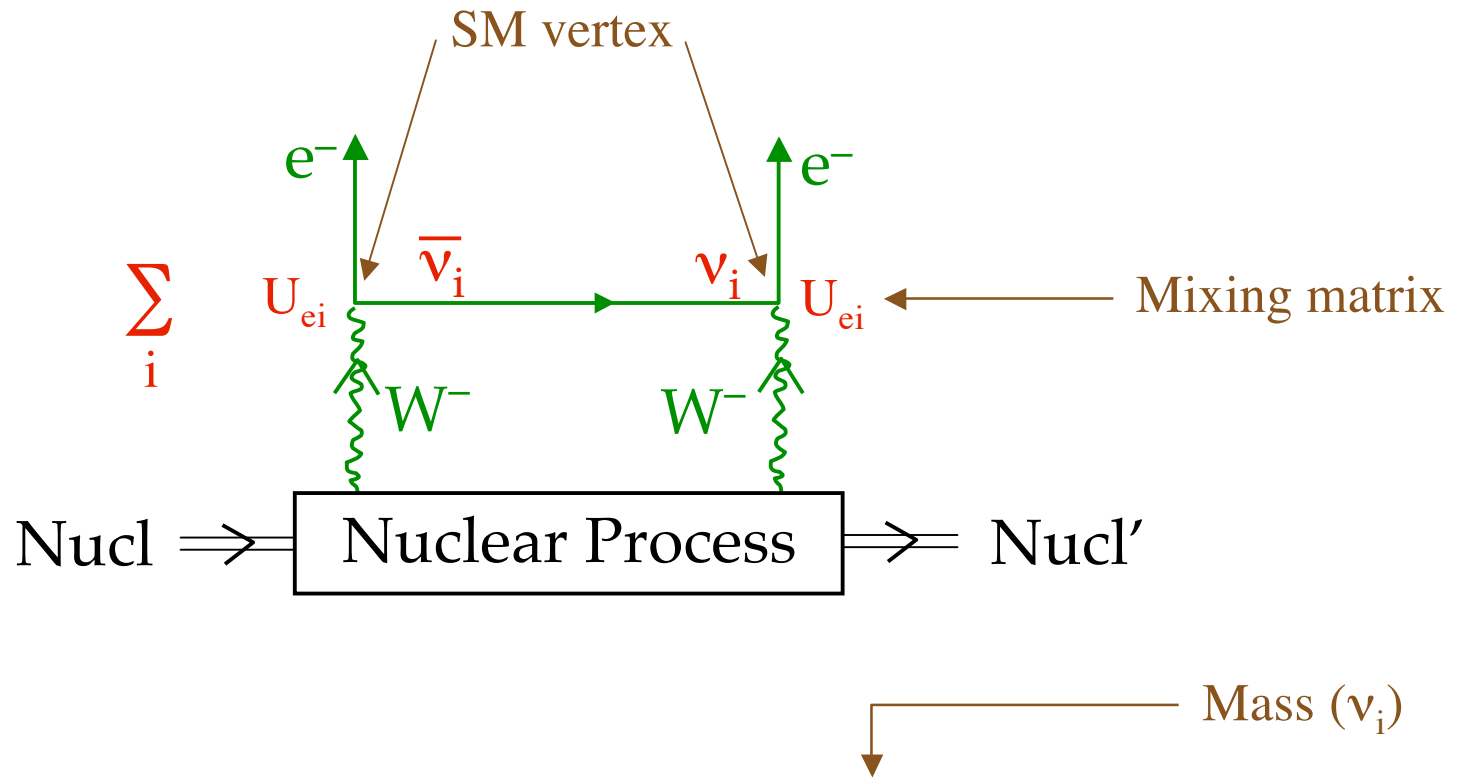


$(\bar{\nu})_R \rightarrow \nu_L$  : A Majorana mass term

We anticipate that  $0\nu\beta\beta$  is dominated by a diagram with Standard Model vertices:



In —



the  $\bar{\nu}_i$  is emitted [RH + O{ $m_i/E$ }LH].

Thus, Amp [ $\nu_i$  contribution]  $\propto m_i$

$$\text{Amp}[0\nu\beta\beta] \propto \left| \sum_i m_i U_{ei}^2 \right| \equiv m_{\beta\beta}$$

The proportionality of  $0\nu\beta\beta$  to mass is no surprise.

$0\nu\beta\beta$  violates L. But the SM interactions conserve L.

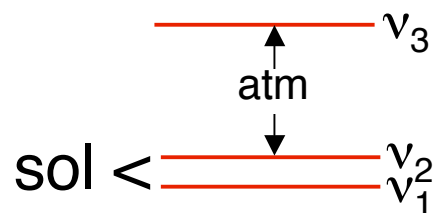
The L – violation in  $0\nu\beta\beta$  comes from underlying  
**Majorana mass terms.**

# How Large is $m_{\beta\beta}$ ?

How sensitive need an experiment be?

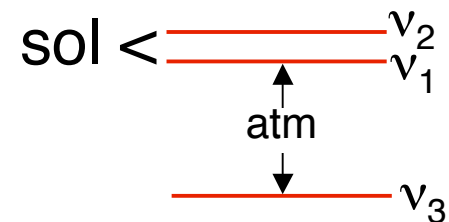
Suppose there are only 3 neutrino mass eigenstates.  
(More might help.)

Then the spectrum looks like —



Normal hierarchy

or

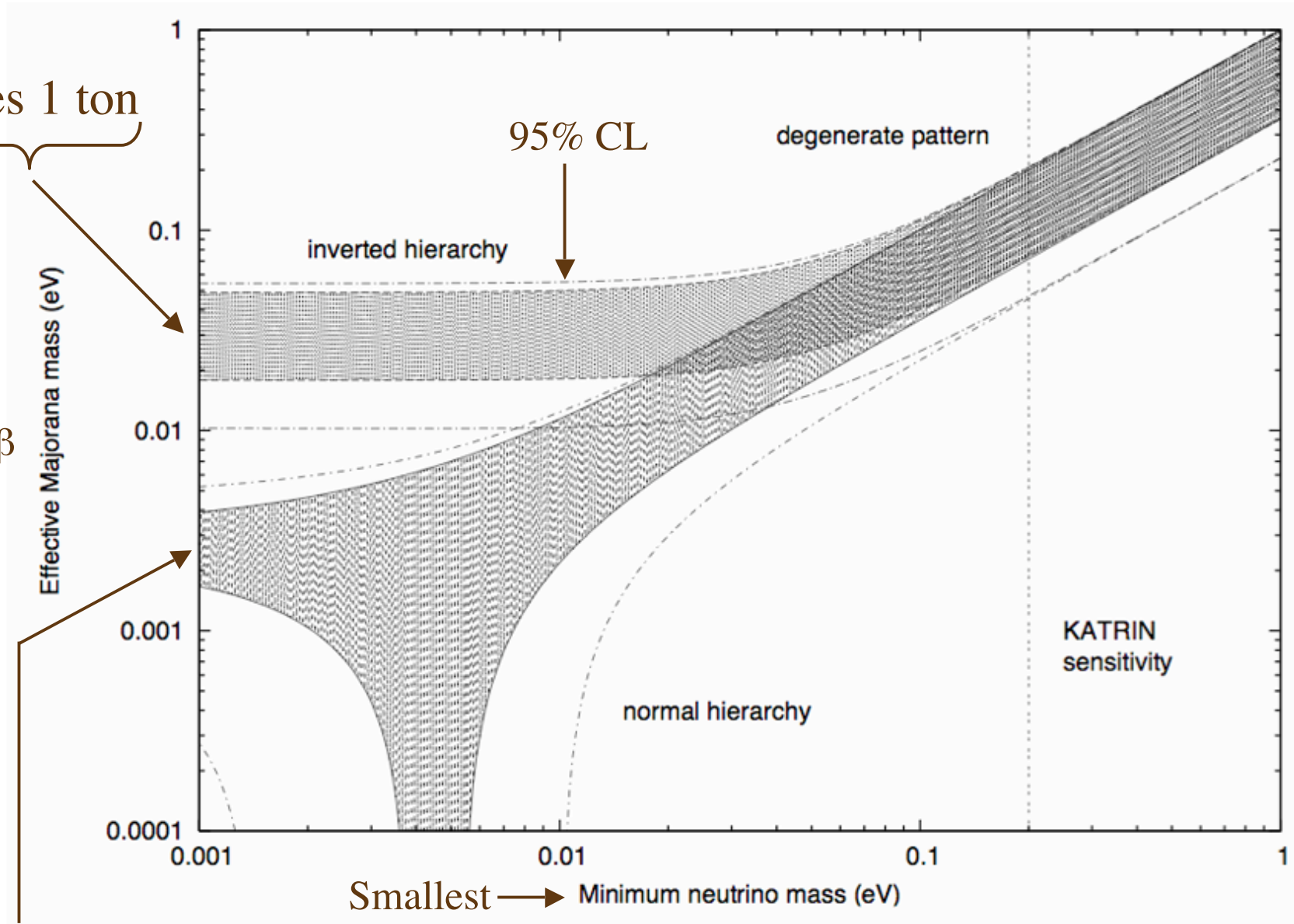


Inverted hierarchy



Takes 1 ton

$m_{\beta\beta}$

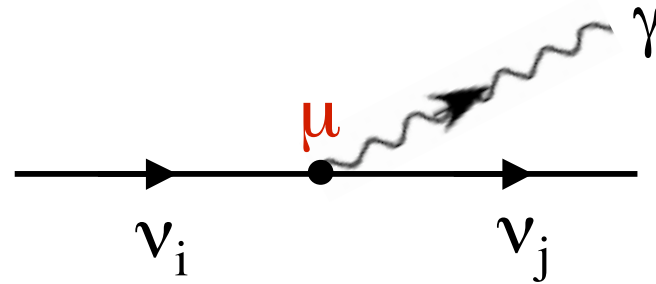


Takes  
100 tons

$m_{\beta\beta}$  For Each Hierarchy

# Possible Information From Neutrino Magnetic Moments

Both Majorana and Dirac neutrinos can have *transition* magnetic dipole moments  $\mu$ :



For *Dirac* neutrinos,  $\mu < 10^{-15} \mu_{\text{Bohr}}$

For *Majorana* neutrinos,  $\mu < \text{Present bound}$

$$\text{Present bound} = \begin{cases} 7 \times 10^{-11} \mu_{\text{Bohr}} ; \text{Wong et al. (Reactor)} \\ 3 \times 10^{-12} \mu_{\text{Bohr}} ; \text{Raffelt (Stellar E loss)} \end{cases}$$

*An observed  $\mu$  below the present bound but well above  $10^{-15} \mu_{Bohr}$  would imply that neutrinos are *Majorana* particles.*

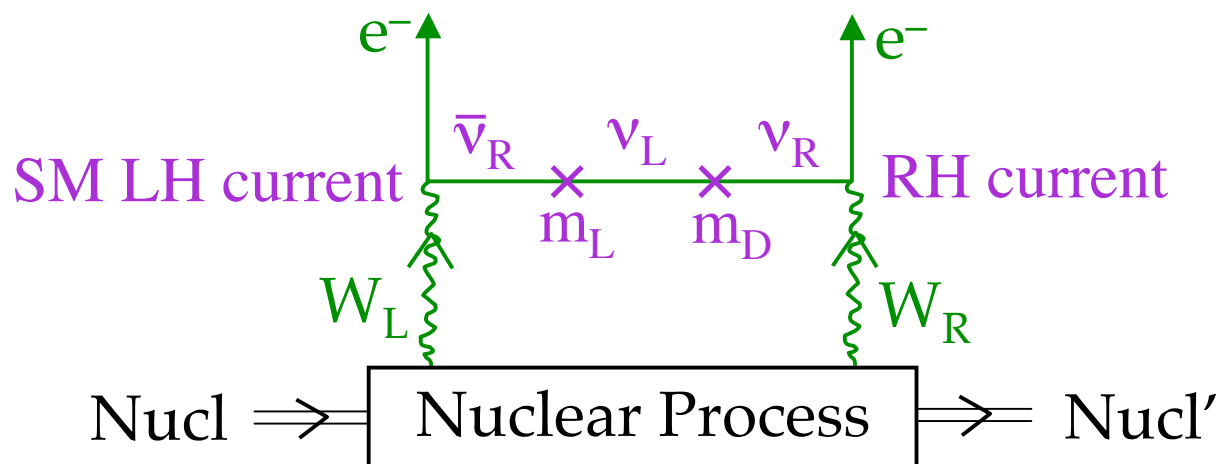
However, a dipole moment that large requires L-violating new physics below 100 TeV.

( Bell, Cirigliano, Davidson, Gorbahn, Gorchtein, Ramsey-Musolf, Santamaria, Vogel, Wise, Wang )

Neutrinoless double beta decay at the planned level of sensitivity only requires this new physics at  $\sim 10^{15}$  GeV, near the Grand Unification scale.

# Backup Slides

Wouldn't the dependence on neutrino mass be eliminated by a Right-Handed Current?



The SM LH current does not violate L.

An identical current, but of opposite handedness, wouldn't violate L either.

We still need the L-violating **Majorana neutrino mass** to make this process occur.

With a RH current at one vertex,

$$\text{Amp}[0\nu\beta\beta] \propto (\nu \text{ mass})^2 .$$

Contributions with a RH current at one vertex  
are not likely to be significant.

{ BK, Petcov, Rosen  
Enqvist, Maalampi, Mursula }