

# Probing the $H^\pm$ with the $\mu_x$ boosted-bottom-jet tag



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## Two-Higgs doublet models (2HDM)

Two-Higgs doublet models (2HDM) are common features in extensions to the SM, especially in SUSY (which needs at least two Higgs doublets). Symmetry breaking produces three scalar Higgs ( $h, H, H^\pm$ ) and a pseudo-scalar ( $A$ ).

$\tan(\beta)$ : the ratio of the two VEVs  $\beta - \alpha$ : the doublet mixing angle

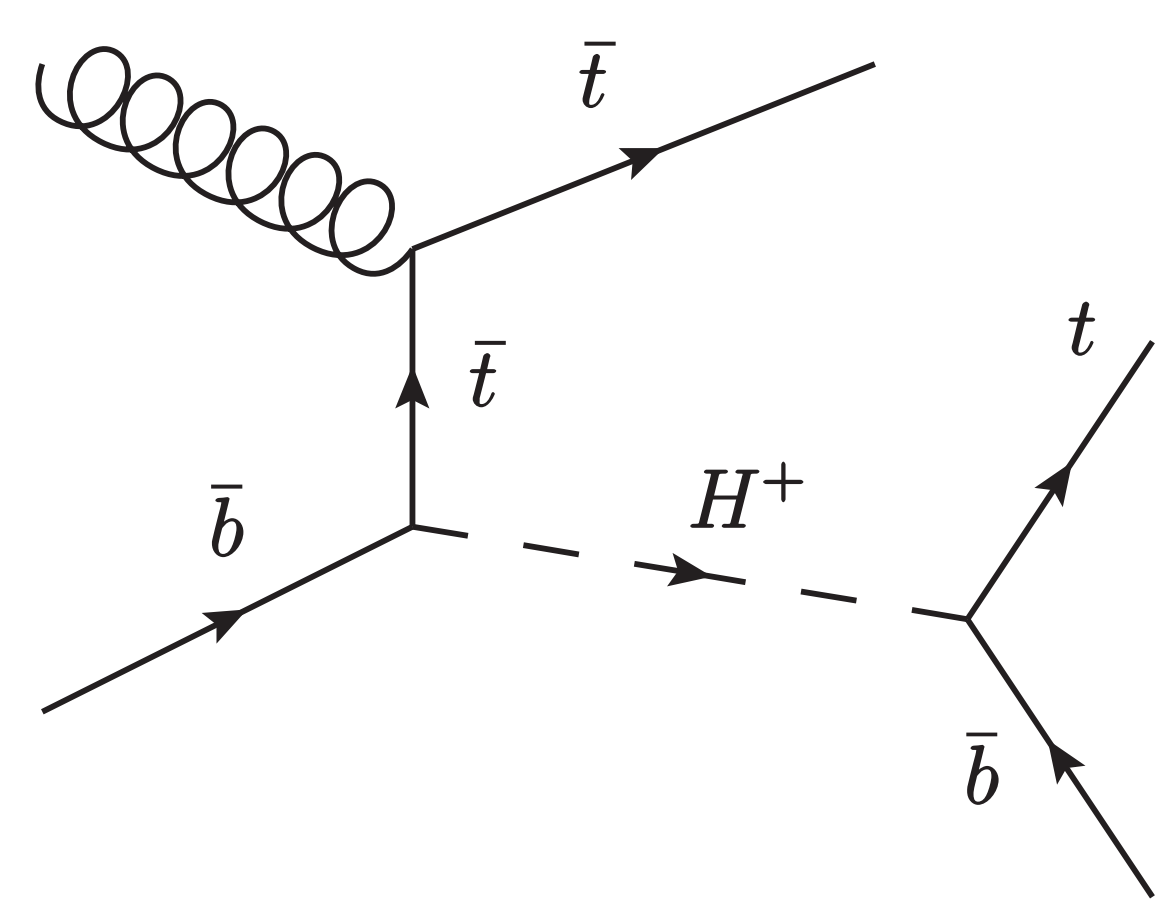
### Realistic 2HDM

- **FCNC** are absent at tree level in **type-II** 2HDM; these also respect SUSY's requirement that  $u_R$  and  $d_R$  couple to opposite doublets.
- There's a **126 GeV boson** that **couples** like a Standard Model Higgs! 2HDM must exist close to the “alignment” limit ( $\cos(\beta - \alpha) \rightarrow 0$ ), so that  $h$  is the SM Higgs.

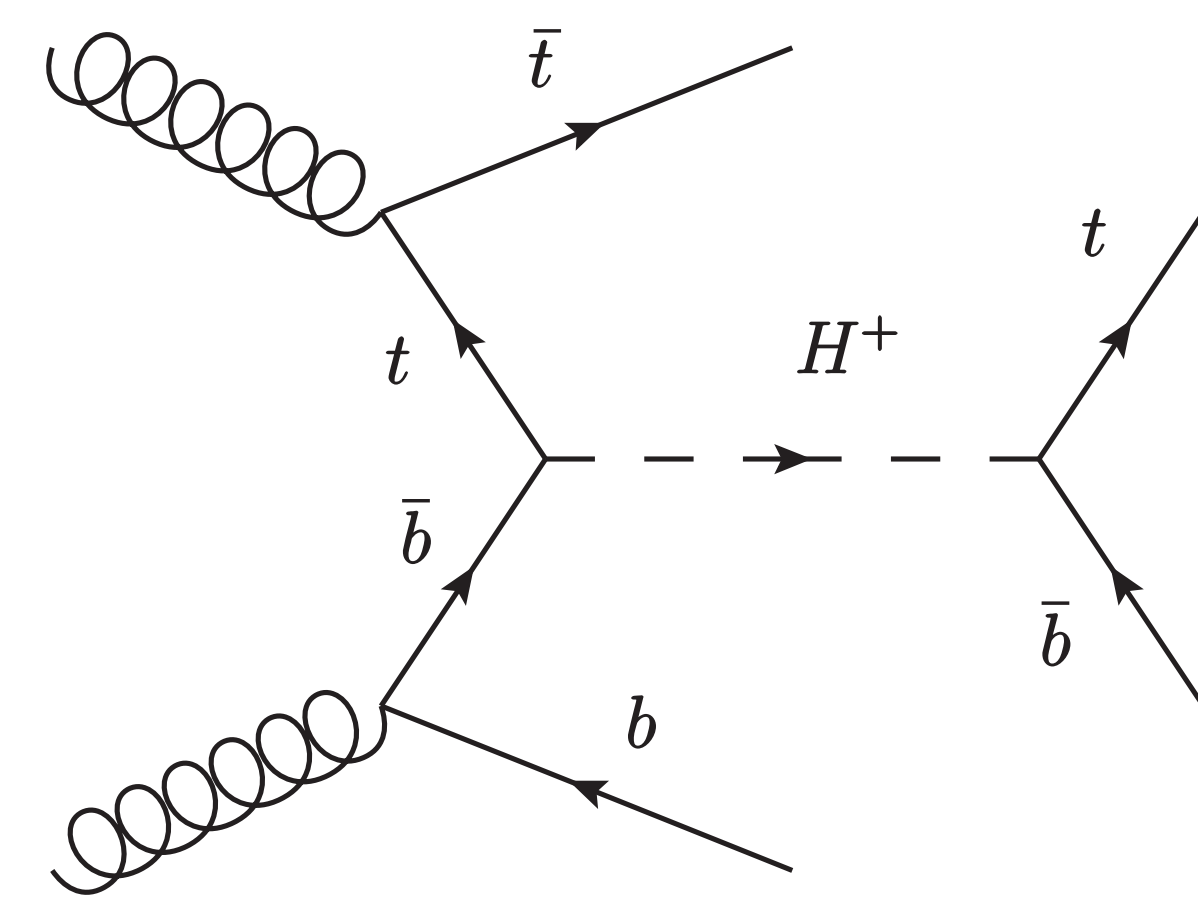
### Phenomenological considerations

- If  $H, A$  and  $H^\pm$  have similar masses (a consequence of SUSY in the alignment limit), they won't decay to each other, but to the SM particles with the largest Yukawa couplings ... **top** and **bottom** quarks (there is no  $ZW^\pm H^\pm$  vertex in 2HDM).
- Interference with SM background makes  $H$  and  $A$  difficult to detect and model. Associated  $H^\pm$  production avoids this, and adds more heavy-flavored jets to tag.

### 3b final state (inclusive cuts)

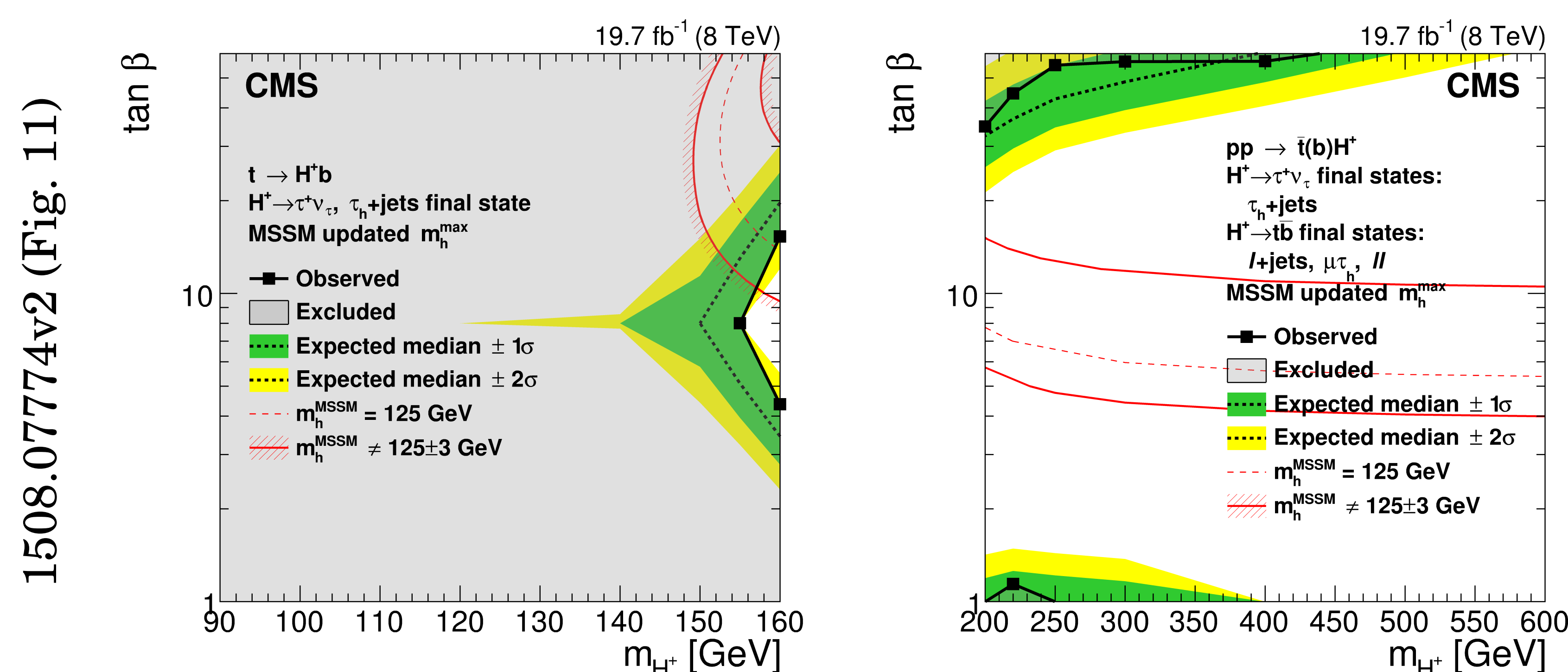


### 4b final state (exclusive cuts)



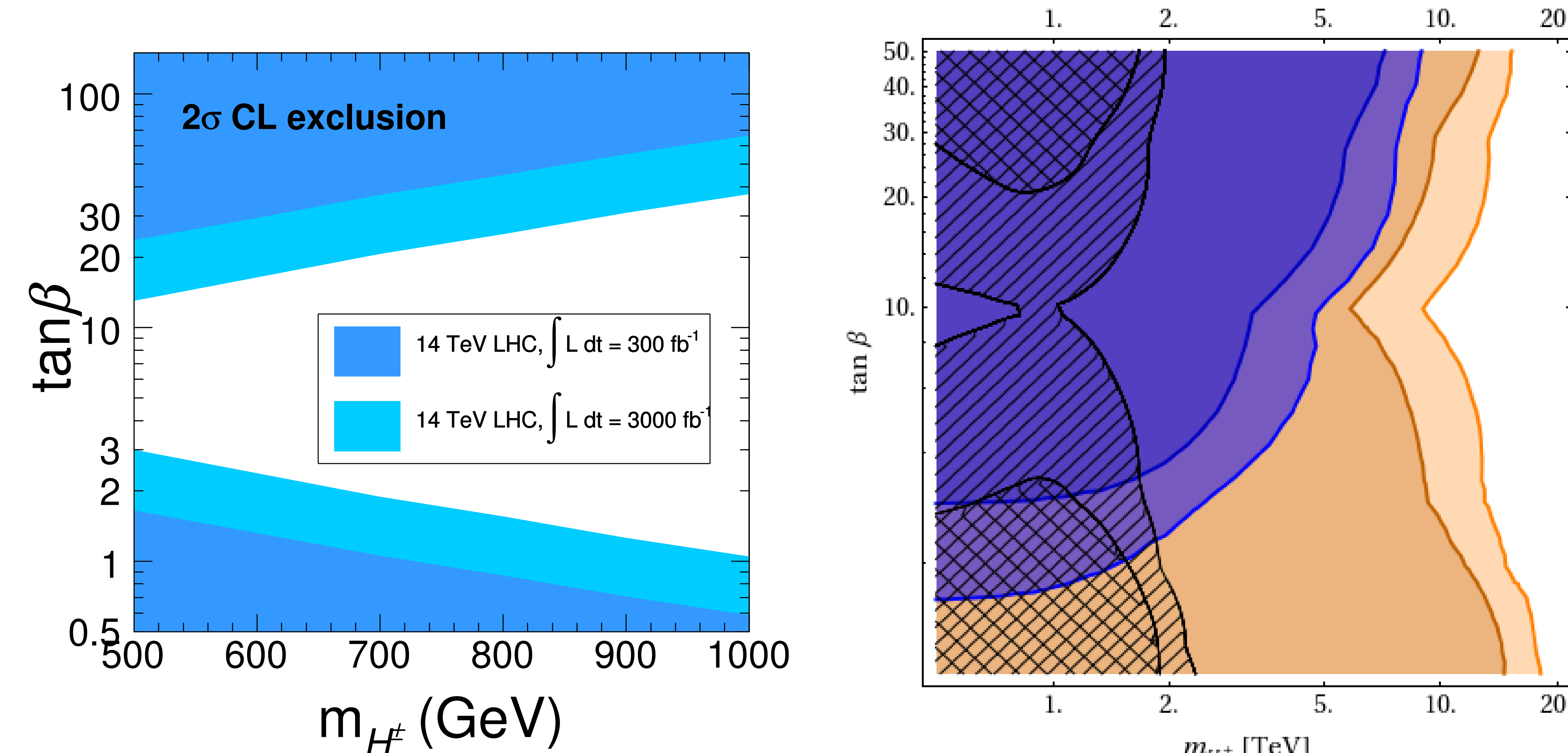
## Experimental challenges

**1. The wedge:** A “natural” 2HDM should have  $\tan(\beta) = O(1-10)$ . This range covers a “wedge” of low production cross section, where the Yukawa coupling to  $H^\pm$  transitions from top-dominated ( $\tan(\beta) < 1$ ) to bottom-dominated ( $\tan(\beta) > 1$ ).



**2. The boosted  $b$ :** As  $H^\pm$  searches push towards TeV masses, the  $H^\pm$  bottom jet becomes increasingly boosted. Unfortunately, standard **track-vertex** bottom jet tags (which use charged tracks to find evidence of a decay vertex slightly displaced from the beam) degrade for boosted  $b$  jets. As jet  $p_T$  increases, the probability to tag a  $b$  jet *decreases*, while the probability to mis-tag a light jet *increases*.

Several groups have made predictions for excluding an  $H^\pm$  heavier than 500 GeV. These predictions rely on simulating track-vertex  $b$  tags inside very boosted jets. What happens when we use a more robust  $b$  tag?



Craig et al. JHEP06(15)137[1504.04630]

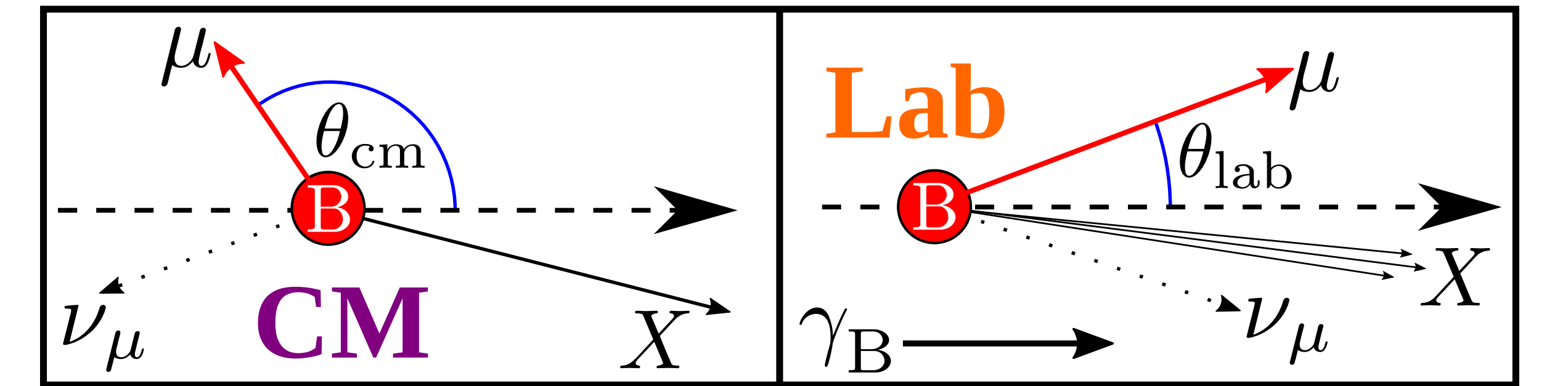
Hajer et al. JHEP11(15)124[1504.07617]  
Checked: 95% @ 14 TeV (300/3000 fb<sup>-1</sup>)  
Tan/orange: 95% @ 100 TeV (3/30 ab<sup>-1</sup>)

## The $\mu_x$ boosted-bottom-jet tag<sup>†</sup>

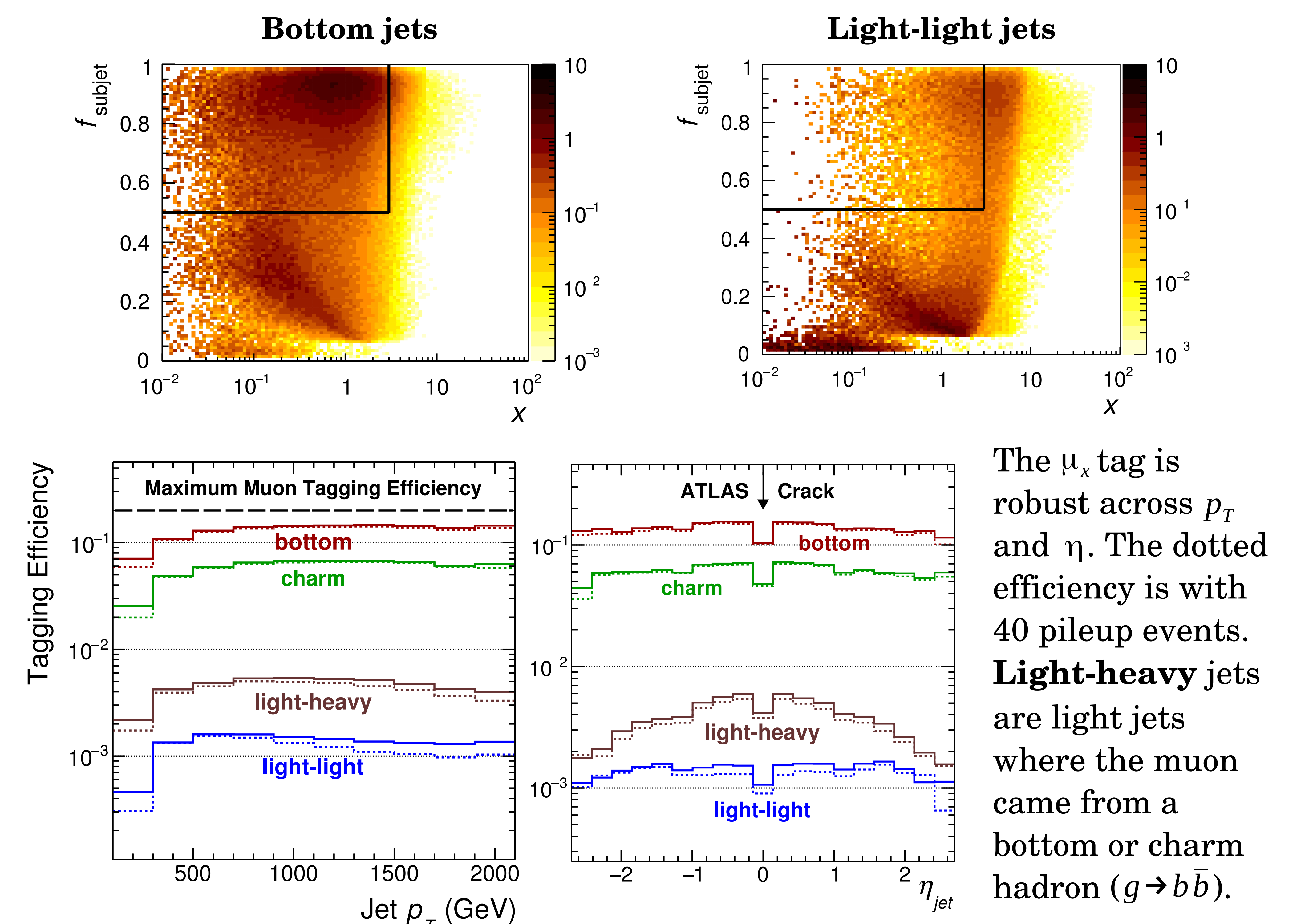
Consider a jet containing a B meson decaying semi-leptonically. In the center-of-momentum (CM) frame, a muon is emitted at some angle  $\theta_{cm}$  w.r.t. the boost axis. In the lab frame, the boost  $\gamma_B$  compresses the B meson products into a subjet. Reconstructing this subjet allows the measurement of a lab frame observable.

$$x \equiv \gamma_B \tan(\theta_{lab}) \approx \frac{\sin(\theta_{cm})}{1 + \cos(\theta_{cm})}$$

$$\frac{dN}{N dx} \approx \frac{2x}{(x^2 + 1)^2}$$



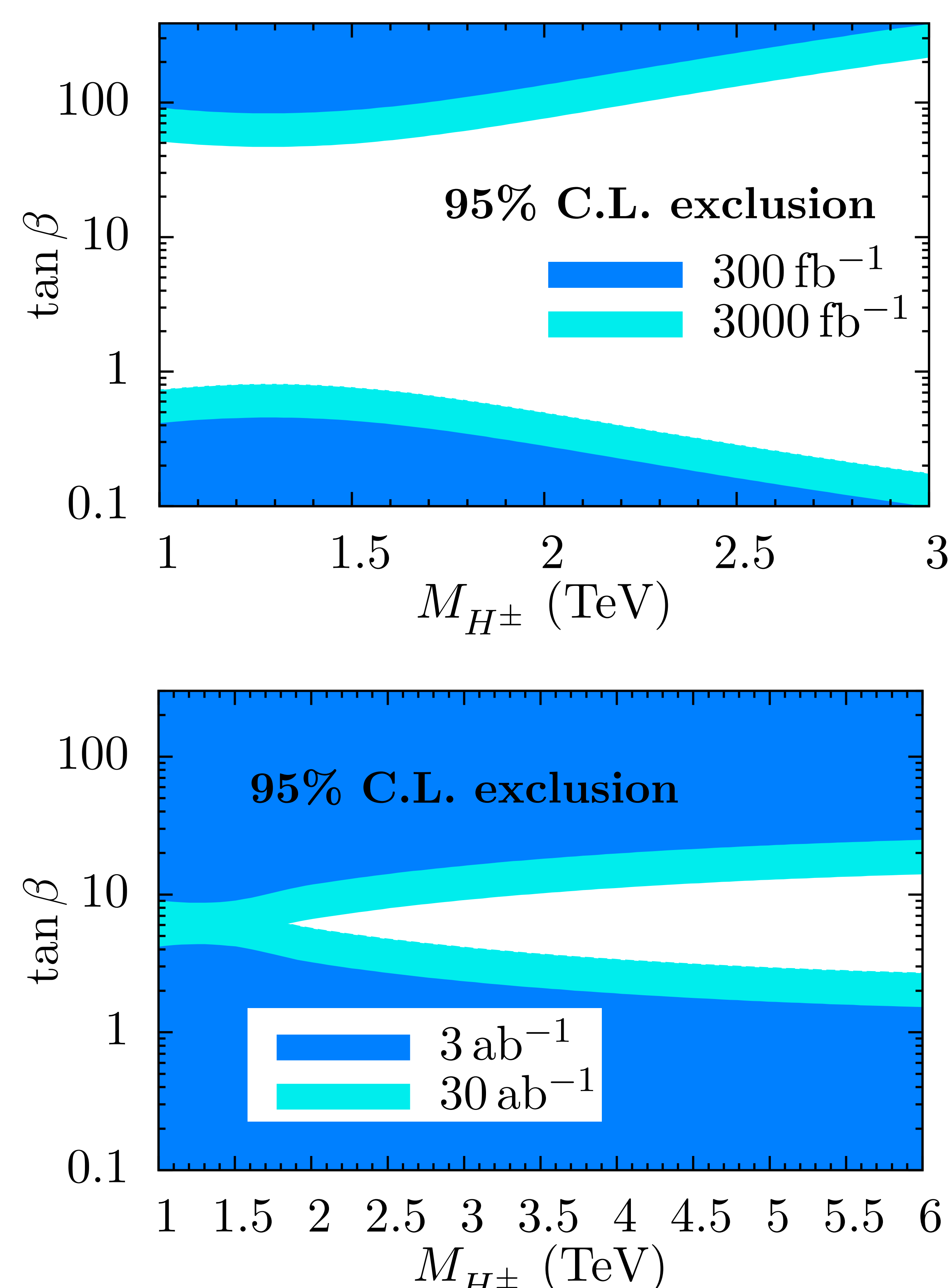
At least 90% of muons arrive in a cone defined by  $x \leq 3$ . In addition, the boosted subjet should carry a large fraction of its jet's momentum  $f_{subjet} = p_{T,subjet}/p_{T,jet} \geq 0.5$ .



The  $\mu_x$  tag is robust across  $p_T$  and  $\eta$ . The dotted efficiency is with 40 pileup events. **Light-heavy** jets are light jets where the muon came from a bottom or charm hadron ( $g \rightarrow b\bar{b}$ ).

## Exclusion predictions at 14 and 100 TeV

We analyzed  $H^\pm$  at LO using MadGraph5, Pythia8, and Delphes3. For the inclusive final state, the  $H^\pm$  is reconstructed from (i) a “narrow” jet (anti-kt,  $R=0.4$ ) tagged with  $\mu_x$  and (ii) a “fat” jet (CA,  $R=0.8$ ) tagged by a boosted hadronic top tag. The associated top is reconstructed as a resolved, leptonic top. The dominant background is  $ttj+X$ , with  $tt(bb/cc)$  about five times smaller. Since the  $3b$  final state accounts for at least 60% of the cross section, we determine that **exclusive** cuts (requiring an associated  $b$ ) do not improve the reach of a search.



We find that at **14 TeV**, the search is entirely limited by the  $H^\pm$  cross section, due to the  $O(10\%)$  branching ratio of the semi-leptonic  $tt$  channel and the  $O(1\%)$  acceptance of the inclusive cuts. However, the signal is quite pure, with signal-over-background ranging from 1/3 to above 1.

At **100 TeV**, the search is background limited, with  $S/B$  ranging from 1% to 5%.

We conclude that our 14 TeV results are comparable to Craig et al., and that the search for  $H^\pm$  in the TeV regime is a strong motivation for a 100 TeV collider.

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<sup>†</sup> K. Pedersen and Z. Sullivan PRD93(16)014014[1511.05990]