

# ***CP violation in the three-body $B^\pm$ phase-space***

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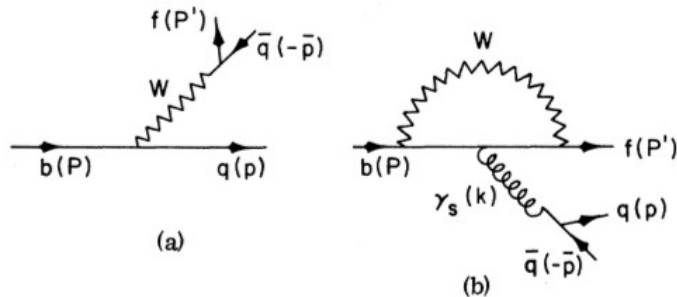
# CP noninvariance in charged decays

- CPV: Different decay rates for particle and antiparticle
- Charged systems: First mechanism to CPV by BSS
  - No CPV at tree level – Need to go beyond first Born approximation
  - Absorption from loop-diagrams is needed
  - QCD: “penguin” diagrams with timelike gluon transition
- Writing the decay amplitude in a general form:  $\mathcal{M} = \mathcal{D} - (i/2)\mathcal{A}$

$$a_{\text{CP}} = \Gamma(i \rightarrow f) - \Gamma(\bar{i} \rightarrow \bar{f}) \propto \text{Im}(\mathcal{D}^* \mathcal{A})$$

M. Bander, D. Silverman and A. Soni,

Phys. Rev. Lett. 43, 242 (1979)



# Basic ingredients to CPV at quark level

- Schematically we can write as an example the  $b \rightarrow s$  amplitude as:

$$\mathcal{M}_{b \rightarrow s} = v_u A_u + v_c A_c + v_t A_t \equiv v_u \Delta_{ut} + v_c \Delta_{ct}$$

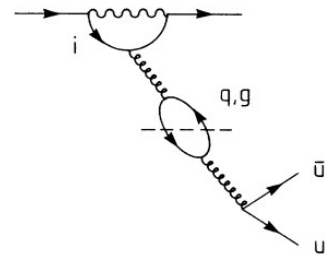
- The asymmetry for this amplitude is:

$$a_{\text{CP}} \propto \text{Im}(v_u^* v_c) \text{Im}(\Delta_{ut}^* \Delta_{ct})$$

- This equation shows that we need basically two ingredients to have CPV:

- 1) different weak CP violating phases and
- 2) absorptive part coming from a loop diagram (hard FSI's);

- Additional terms with respect to CPT;



J.M. Gérard and W.S. Hou, Phys. Rev. Lett. 62, 855 (1989)

FIG. 2. The extra  $O(\alpha_s^2)$  graph needed to satisfy the  $CPT$  theorem at the  $O(\hbar^2)$ ,  $O(\alpha_s^2)$  level. Only the absorptive part, i.e., the cut graph, of the vacuum bubble is needed.

# Soft final state interactions

- Develop a formalism based on CPT invariance and Unitarity;
    - Hadronic FSI replace the absorptive part of the penguin graph;
    - Soft final state interactions do not disappear for large  $m_b$ ;
    - Soft FSI should have an important effect for CPV;
  - Meson-meson inelastic scattering at 5 GeV:
    - Important strong phase effect in B decays;
    - Depends on how probable the state is as a final state;
- 
- L. Wolfenstein, Phys. Rev. D 43, 151 (1991)
  - J. F. Donoghue, E. Golowich, A. A. Petrov, and J. M. Soares, Phys. Rev. Lett. 77, 11 (1996))
    - M. Suzuki and L. Wolfenstein, Phys. Rev. D 60, 074019 (1999)

# Experimental facts

- Significant part of the observed CPV distribution is located in a region where hadronic channels are strongly coupled;

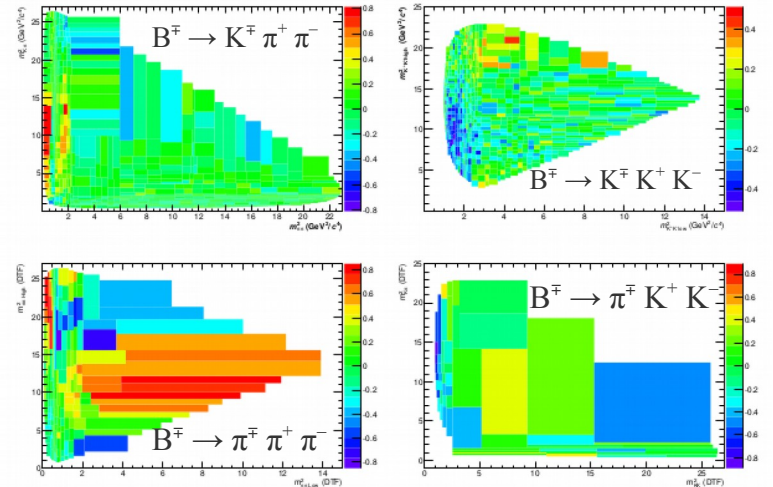
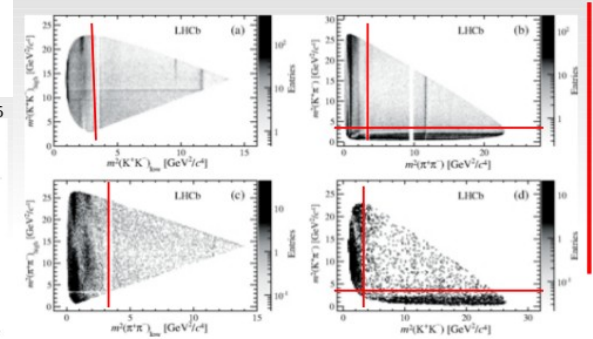
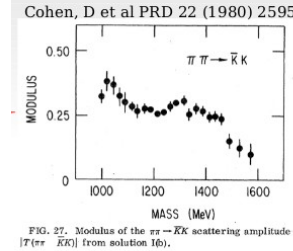
- Events concentrated at low invariant mass;

- Huge phase-space available;

- Reasonable assumptions:

- 1) Final states factorize in a two body interacting system plus a bachelor;
- 2) Rescattering effects in three-body B decays happens in  $3 \rightarrow 3$  channels;

- D. Cohen, D. S. Ayres, R. Diebold, S. L. Kramer, A. J. Pawlicki, and A. B. Wicklund, Phys. Rev. D 22, 2595(1980).
- R. Aaiji et al. (LHCb Collaboration), Phys. Rev. D 90, 112004 (2014).



# Model asymmetry

- Split the S-matrix in two parts

$$S = S_0 + S_1,$$

where  $S_0$  is the diagonal part with the elastic scattering and  $S_1$  is perturbatively treated as the first order inelastic transition;

- Factorizes the short-range and long-range contributions in the decay amplitude:

$$\mathcal{A}_{\lambda LO}^{\pm} = A_{0\lambda} + e^{\pm i\gamma} B_{0\lambda} + i \sum_{\lambda'} t_{\lambda\lambda'} (A_{0\lambda'} + e^{\pm i\gamma} B_{0\lambda'})$$

$$\Delta\Gamma_{\lambda} = |A_{\lambda}^{-}|^2 - |A_{\lambda}^{+}|^2 = \Gamma(h \rightarrow \lambda) - \Gamma(\bar{h} \rightarrow \bar{\lambda})$$

$$\Delta\Gamma_{\lambda} = 4 (\sin \gamma) \text{Im}[B_{0\lambda}^{*} A_{0\lambda} + i \sum_{\lambda'} (B_{0\lambda}^{*} t_{\lambda\lambda'} A_{0\lambda'} - B_{0\lambda'}^{*} t_{\lambda\lambda'}^{*} A_{0\lambda})]$$

- Channels coupled by strong interactions implies  $\Delta\Gamma_{\alpha} = -\Delta\Gamma_{\beta}$ 
  - I. Bediaga, T. Frederico, and O. Lourenço, Phys. Rev. D 89, 094013 (2014).

# S and P-wave resonant interferences

- Two resonant intermediate states, with different weak phases, share the same kinematical region and hadronic final state;
- CP asymmetry from S and P wave interference in same hadronic final state.
  - Resonant interferences ( $\rho(770)$  and  $f_0(980)$  resonances) + hadronic inelastic rescattering

$$\mathcal{A}_{0\lambda}^{\pm} = A_0^{\rho} F_{\rho}^{\text{BW}} k(s) \cos \theta + A_0^{f_0} F_{f_0}^{\text{BW}} + A_0^{\text{NR}},$$

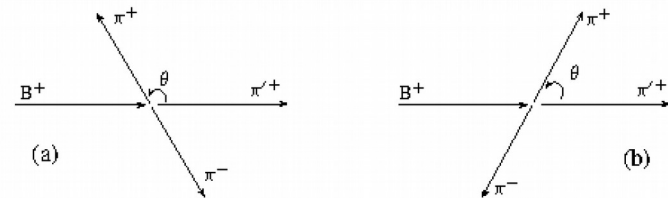
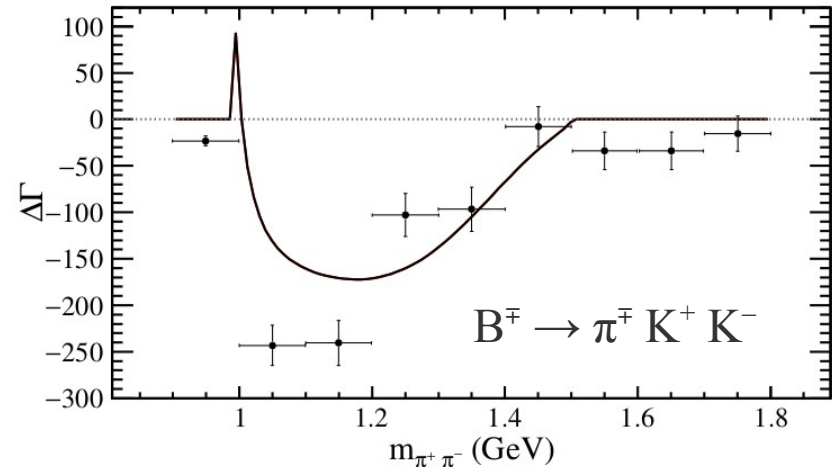
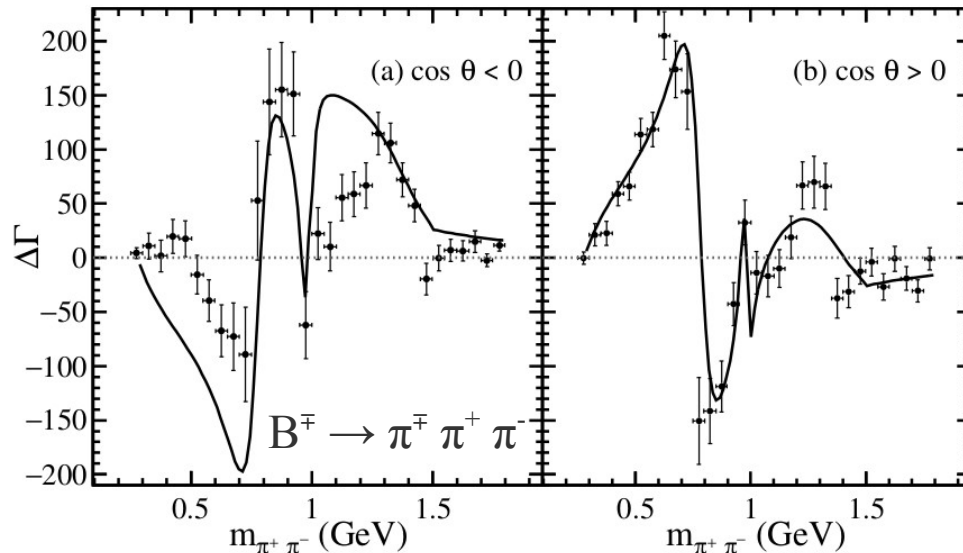
$$\text{where } A_0^i = a_0^i + b_0^i e^{\pm i\gamma} \quad (i = \rho, f_0, \text{NR})$$

$$\text{S-matrix: } S_{\lambda,\lambda'} = \delta_{\lambda,\lambda'} + i t_{\lambda,\lambda'}, \quad t_{\pi\pi \rightarrow \text{KK}} = \sqrt{1-\eta^2}(s) e^{i(\delta_{\text{KK}} + \delta_{\pi\pi})} \quad (\text{using } \delta_{\text{KK}} \approx \delta_{\pi\pi})$$

$$\begin{aligned} \Delta\Gamma_{\lambda} = & \mathcal{B} \sqrt{1-\eta^2}(s) \cos[2\delta_{\pi\pi}(s)] + |F_{\rho}^{\text{BW}}(s)|^2 k(s) \cos \theta \{ \mathcal{D}(m_{\rho}^2 - s) + \mathcal{E} m_{\rho} \Gamma_{\rho}(s) \} + \\ & \mathcal{G} |F_{\rho}^{\text{BW}}(s)|^2 |F_{f_0}^{\text{BW}}(s)|^2 k(s) \cos \theta [(m_{\rho}^2 - s) m_{f_0} \Gamma_{f_0}(s) + m_{\rho} \Gamma_{\rho}(s) (m_{f_0}^2 - s)] \end{aligned}$$

- I. Bediaga, T. Frederico, and O. Lourenço, Phys. Rev. D 89, 094013 (2014).
- J.H.A.N., I. Bediaga, A.B.R. Cavalcante, T. Frederico, O. Lourenço. Phys.Rev. D 92, 054010 (2015)
- I.B., I.I. Bigi, A. Gomes, G. Guerrer, J. Miranda and A.C. Dos Reis-Phys. Rev. D80, 096006 (2009)
  - J.R. Pelaez and F.J. Ynduráin, Phys. Rev. D 71, 074016 (2005)

# Fitting the data



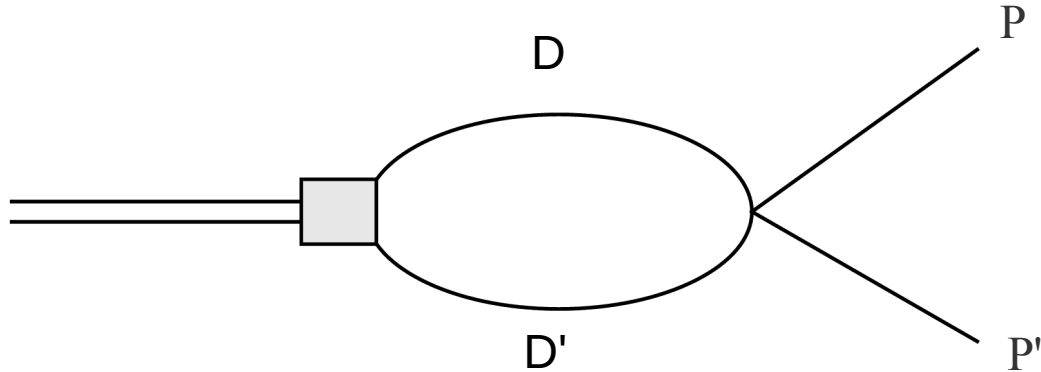
➤ Only region  $\cos \theta > 0$  of  $B^- \rightarrow \pi^- \pi^+ \pi^-$  is fitted

➤ J.H.A.N., I. Bediaga, A.B.R. Cavalcante, T. Frederico, O. Lourenço, Phys.: Conf. Ser. 706 042010 (2016)

➤ Same study was done for  $B^- \rightarrow K^+ \pi^+ \pi^-$  (obtaining  $B^- \rightarrow K^+ K^+ K^-$  as output) in J.H.A.N., I. Bediaga, A.B.R. Cavalcante, T. Frederico, O. Lourenço. Phys.Rev. D 92, 054010 (2015)

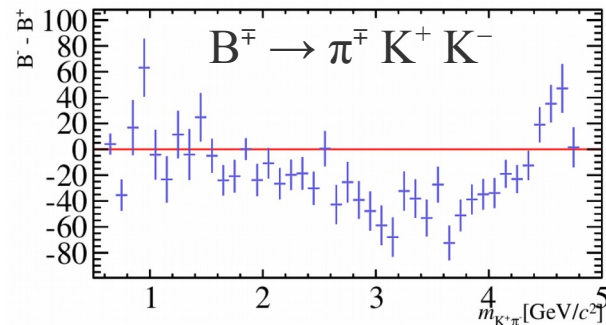
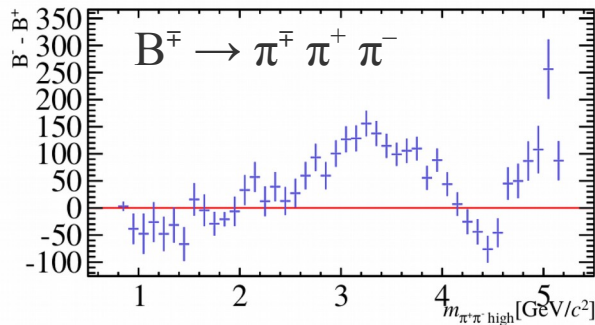
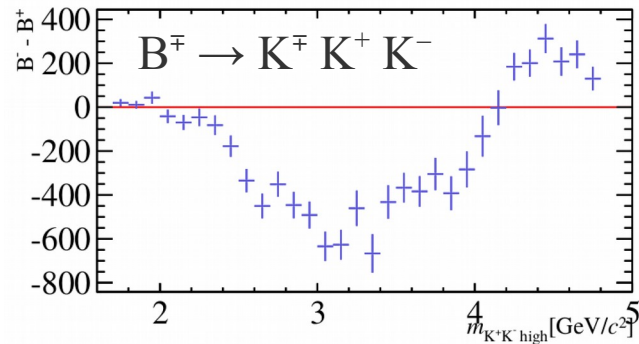
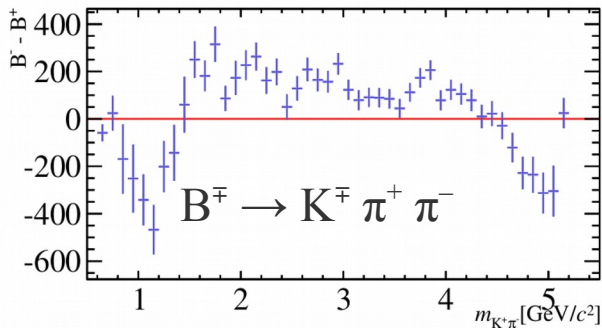


# High mass CP distribution



- Involving two-body coupling between light mesons ( $PP'$ ) and heavy mesons (double charm  $DD'$ );
  - $P$  and  $P'$  can be kaon or pion;

# CP asymmetry in the high mass region



- R. Aaiji et al. (LHCb Collaboration), Phys. Rev. D 90, 112004 (2014).  
New LHCb results: available at <https://cds.cern.ch/record/1751517?ln=en>

# CP asymmetry in the high mass region

- Three light pseudoscalar mesons can couple with channels like  $D\bar{D}h$ ,  $h = K, \pi$ ;
  - Rescattering can contribute to CPV in regions of large two-body invariant mass above the  $D\bar{D}$  threshold;
- Final state interactions in hadronic B decays involving transitions from heavy to light mesons in the literature;
- $D\bar{D}_s \rightarrow K\pi$  rescattering computed within chiral Lagrangian based models;
- Or parametrized using an effective expansion for the S-matrix;
- Branching ratio to double charm channels is bigger than to light mesons channels;
  - Even if the coupling is small, the BR is a factor of 10-100 bigger;

➤ H.-Y. Cheng, C.-K. Chua, and A. Soni, Phys. Rev. D 71, 014030 (2005).

➤ J.A. Oller, E. Oset, Phys. Rev. D 60, 074023 (1999).

➤ D. Gamermann, E. Oset, D. Strottman and M. J. Vicente Vacas, Phys. Rev. D 76, 074016 (2007).

# CP asymmetry in the high mass region

- Preliminary result (using  $t_{K\pi \rightarrow DD_s}$  from chiral lagrangian model):
- Two-channel unitary S-matrix parametrization:

- $|S_{LL \rightarrow HH}|^2 = \sqrt{1 - \eta^2(s)} = \Lambda (s/s_{th} - 1)^{1/2} / (s/s_{th})^{2.5}$

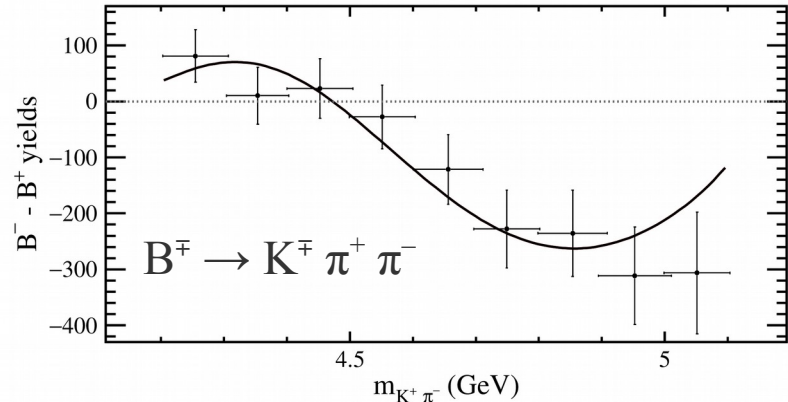
$$S_\lambda = (k \cot \delta_\lambda + i k_\lambda) / (k \cot \delta_\lambda - i k_\lambda), \quad k_\lambda = (s - s_{th\lambda})^{1/2} / 2$$

- $\lambda=LL$  channel with a pole above the HH threshold:

$$k \cot \delta_{LL} = -c / (1 - k_{LL} / k_{0LL}), \quad k_{0LL} = (s_0 - s_{thLL})^{1/2} / 2$$

- $\lambda=HH$  channel with a virtual bound state:  $k \cot \delta_{HH} = -1/a$

Our parametrization agrees with this result.



[S] parametrization similar to the one in J.R. Pelaez and F.J. Ynduráin, Phys. Rev. D 71, 074016 (2005)

# Summary and Outlook

- CPT must be a practical constraint in the analysis of the three-body B meson decays;
- Formalism is based on CPT invariance and Unitarity;
- Soft FSI is explicitly factorized and shows to be essential in our study;
- The B meson phase space presents resonant interferences with particular features;
- The decay channels  $B^{\mp} \rightarrow \pi^{\mp} \pi^+ \pi^-$  and  $B^{\mp} \rightarrow \pi^{\mp} K^+ K^-$  (also  $B^{\mp} \rightarrow K^{\mp} \pi^+ \pi^-$  and  $B^{\mp} \rightarrow K^{\mp} K^+ K^-$ ) presents asymmetries that seems to be related by the CPT constraint;
- $PP \rightarrow HH$  ( $P = \pi, K$  and  $H = D, D_s$ ) scattering seems to be related with the high mass CP distribution;
- Understand hadron-hadron scattering below 5 GeV using CPV data;
- Next steps:
  - Include the three-body rescattering effects with the spectator meson;
  - Address the quark-level processes as source amplitudes for the hadronic rescattering;



**Thank you!**