

# $|V_{us}|$ FROM HADRONIC $\tau$ DECAYS

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## OUTLINE

- *$V_{us}$  from inclusive flavor-breaking sum rules*
  - *Systematics in the conventional implementation and a resolution of the  $> 3\sigma$  low  $|V_{us}|$  puzzle*
  - *New implementation strategy results + current experimental limitations*
- *A new lattice+inclusive  $us$   $V+A$   $\tau$  data strategy*

## CONTEXT

- $\tau$  vs. non- $\tau$   $|V_{us}|$  determinations: the  $> 3\sigma$  low inclusive FB  $\tau$  FESR  $|V_{us}|$  puzzle

$ V_{us} $	Source
0.2258(9)(?)	3-family unitarity, HT14 $ V_{ud} $
$0.2231(4)_{exp}(7)_{latt}$	$K_{\ell 3}$ , 2+1+1 lattice $f_+(0)$
$0.2250(4)_{exp}(9)_{latt}$	$\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$ , lattice $f_K/f_\pi$
$0.2176(19)_{exp}(10?)_{th}$	Inclusive FB kinematic wt $\tau$ FESR (Passemar CKM14)

- $\tau$  result: from “conventional implementation” of more general inclusive FB FESR framework

## BASICS: HADRONIC $\tau$ DECAYS IN THE SM

- $R_{ij;V/A} \equiv \Gamma[\tau \rightarrow \nu_\tau \text{ hadrons}_{ij;V/A}(\gamma)]/\Gamma[\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e(\gamma)]$
- With  $y_\tau \equiv s/m_\tau^2$ , flavor  $ij$  decays in SM [Tsai PRD4 (1971) 2821]

$$\frac{dR_{ij;V+A}}{ds} = \frac{12\pi^2 |V_{ij}|^2 S_{EW}}{m_\tau^2} [1 - y_\tau]^2 \tilde{\rho}_{ij;V+A}(s)$$

$$\tilde{\rho}_{ij;V+A}(s) \equiv [(1 + 2y_\tau) \rho_{ij;V+A}^{(J=1)}(s) + \rho_{ij;V+A}^{(J=0)}(s)]$$

kinematic weight :  $w_\tau(y) = (1 - y)^2(1 + 2y)$

# THE INCLUSIVE FB $\tau$ $|V_{us}|$ DETERMINATION

- FESRs for  $\Pi_{ud-us;V+A}^{(J=0+1)}(Q^2)$ ,  $\rho_{ud-us;V+A}^{(J=0+1)}(s)$  Cauchy's theorem

$$\int_{s_{th}}^{s_0} \underset{\text{experiment}}{ds} w(s) \rho(s) = \frac{-1}{2\pi i} \oint_{|s|=s_0} \underset{\text{OPE}}{ds} w(s) \Pi(s)$$

- Experiment:  $|V_{ij}|^2 \rho_{ij;V/A}^{(0+1)}(s)$  from  $dR_{ij;V/A}/ds$  mildly model-dependent continuum  $us$  subtraction
- $R_{ij;V/A}^w(s_0)$ : re-weighted  $R_{ij;V/A}$  analogue

$$R_{ij;V/A}^w(s_0) \sim \int_{th}^{s_0} ds \frac{dR_{ij;V/A}}{ds} \frac{w(s/s_0)}{w_\tau(s/m_\tau^2)}$$

- FB differences  $\delta R^w(s_0) \equiv \frac{R_{ud;V+A}^w(s_0)}{|V_{ud}|^2} - \frac{R_{us;V+A}^w(s_0)}{|V_{us}|^2}$

- FESR, OPE for  $\delta R^w(s_0)$ , input  $|V_{ud}| \Rightarrow$

$$|V_{us}| = \sqrt{\frac{R_{us;V+A}^w(s_0)}{\frac{R_{ud,V+A}^w(s_0)}{|V_{ud}|^2} - [\delta R^w(s_0)]^{OPE}}}$$

Self-consistency:  $|V_{us}|$  independent of  $s_0, w$

- **The conventional implementation** [Gamiz et al. JHEP03(2003)060]
  - $s_0 = m_\tau^2$ ,  $w = w_\tau$  only [spectral integrals from inclusive  $ud, us$  BFs, **but no self-consistency tests**]
  - $w_\tau$  degree 3  $\Rightarrow$  OPE to  $D = 8$
  - Strong assumptions re  $D = 6$  (VSA),  $D = 8$  ( $\sim 0$ )

- **Conventional implementation tests** [KM et al. arXiv:1511.08514]

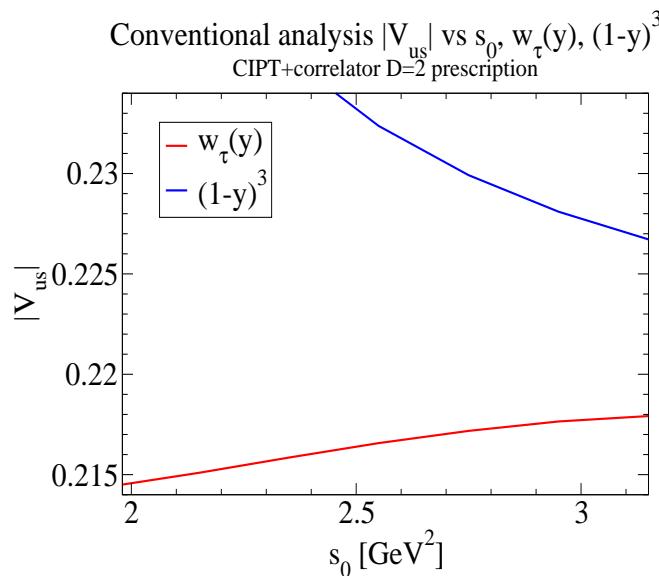
- $|V_{us}|$  stability checks with variable  $s_0 \leq m_\tau^2$

D=8

- Targeted  $D = 6, 8$  assumptions test:  $y = (s/s_0)$ ,

$$w_\tau(y) = 1 - 3y^2 + 2y^3 \text{ c.f. } \hat{w}(y) = 1 - 3y + 3y^2 - y^3$$

D=6



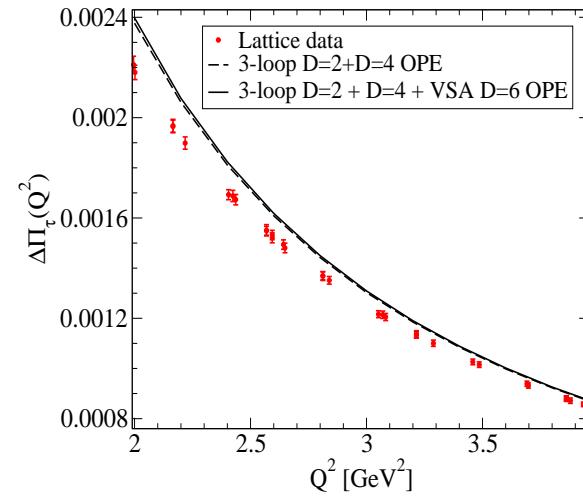
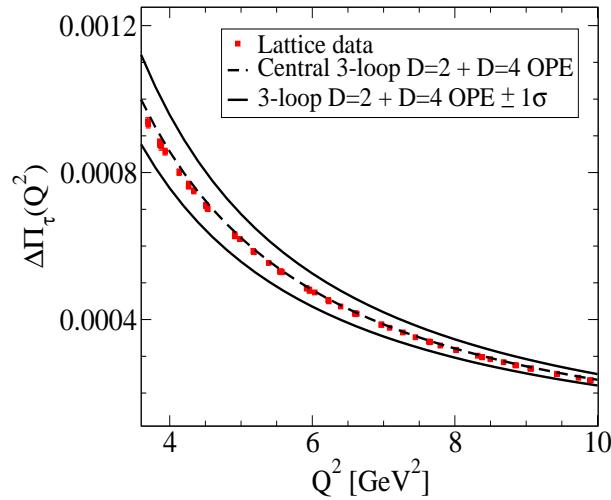
- Slow  $D = 2$  convergence also a potential issue

$$\begin{aligned} [\Delta \Pi_\tau(Q^2)]_{D=2}^{OPE} &= \frac{3}{2\pi^2} \frac{\bar{m}_s^2}{Q^2} \left[ 1 + \frac{7}{3}\bar{a} + 19.933\bar{a}^2 \right. \\ &\quad \left. + 208.746\bar{a}^3 + \dots \right] \end{aligned}$$

$\overline{MS}$  running  $\bar{a} = \frac{\alpha_s(Q^2)}{\pi}$ ,  $\bar{m}_s = m_s(Q^2)$ ,  $\bar{a}(m_\tau^2) > 0.1$

- OPE/lattice  $\Pi_{ud-us:V+A}^{(0+1)}(Q^2)$  comparison
  - $n_f = 2+1$ ,  $m_\pi \sim 300 \text{ MeV}$ ,  $1/a = 2.38 \text{ GeV}$ ,  $m_\pi L \sim 4.1$ ,  $32^3 \times 64$  RBC/UKQCD ensemble
  - Tight cylinder cut for continuum correlator behavior
  - Excellent lattice/ $D = 2+4$  OPE match for fixed scale, 3-loop  $D = 2$ ,  $Q^2 \sim 4 - 10 \text{ GeV}^2$  [FIG]

- Conventional OPE error estimates **VERY** conservative despite slow  $D = 2$  convergence [FIG]
- Confirms non-negligible  $D > 4$ ,  $Q^2 < 4 \text{ GeV}^2$  [FIG]



## AN ALTERNATE FB FESR IMPLEMENTATION

- Theory side
  - No  $D > 4$  assumptions: effective condensates  $C_{D>4}$  from fits to data (requires variable  $s_0$ )
  - 3-loop-truncated FOPT  $D = 2$ , standard  $D = 2 + 4$  error estimates [as per comparison to lattice]
  - $C_{2N+2}$ ,  $|V_{us}|$  from  $w_N(y) = 1 - \frac{y}{N-1} + \frac{y^N}{N-1}$  FESR
  - $|V_{us}|$  from different  $w_N$  as self-consistency check

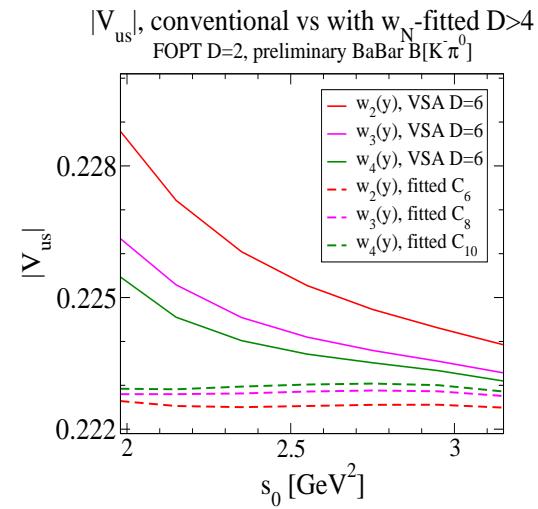
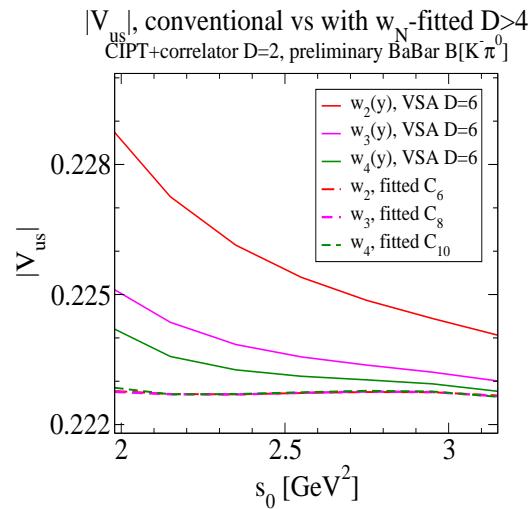
- Experimental input

- Updated/corrected 2013 ALEPH for  $ud$  V+A
- $us$  V+A from sum over exclusive modes
  - \*  $K$  from  $K_{\mu 2}$
  - \*  $K\pi$ ,  $K^-\pi^+\pi^-$ ,  $\bar{K}^0\pi^-\pi^0$ : BaBar, Belle unit-normalized distributions, BFs

[Note: HFAG  $B[K^-\pi^0\nu_\tau] = 0.00433(15)$  (BaBar dominated) c.f. preliminary BaBar (Adametz) thesis  $0.00500(15)$  (recommended by BaBar)]
  - \* Remaining (“residual modes”) from 1999 ALEPH (note:  $\sim 25\%$  errors, some MC)

## RESULTS

- Unphysical  $s_0$ -,  $w(y)$ -dependence problems solved. E.g., for Adametz  $B[K^-\pi^0\nu_\tau]$  input



- $|V_{us}|$  increased by  $\sim 0.0020$  with fitted  $C_{D>4}$

- Significant impact of HFAG 2014 → preliminary BaBar Adametz thesis  $B[K^-\pi^0\nu_\tau]$  (3-weight averages)

$$|V_{us}| = 0.2200(23)_{exp}(5)_{th} \quad (HFAG)$$

$$|V_{us}| = 0.2228(23)_{exp}(5)_{th} \quad (Adametz)$$

- Adametz  $B[K^-\pi^0\nu_\tau]$  input  $w$ -independence example

Weight	$ V_{us} $ CIPT+corr $D=2$	$ V_{us} $ FOPT $D=2$
$w_2$	0.2227(23)	0.2225(23)
$w_3$	0.2227(23)	0.2228(23)
$w_4$	0.2227(23)	0.2230(23)

- New implementation, updated  $B[K^-\pi^0\nu_\tau]$  completely resolves old  $> 3\sigma$  low  $|V_{us}|$  puzzle

- Very favorable ( $\sim 0.0005$ ) theory error situation
- $us$  spectral integral uncertainty dominates current error
- Error budget, 3-weight, Adametz  $B[K^-\pi^0\nu_\tau]$ , 3-loop-truncated FOPT  $D = 2$  fit

Source	$\delta V_{us} $ ( $w_2$ FESR)	$\delta V_{us} $ ( $w_3$ FESR)	$\delta V_{us} $ ( $w_4$ FESR)
$\delta\alpha_s$	0.00001	0.00004	0.00004
$\delta m_s(2 \text{ GeV})$	0.00017	0.00019	0.00019
$\delta\langle m_s\bar{s}s \rangle$	0.00035	0.00035	0.00035
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
$ud$ exp	0.00027	0.00028	0.00028
$us$ exp	0.00226	0.00227	0.00227

- Theory error  $\Rightarrow$  competitive with  $K_{\ell 3}$ ,  $\Gamma[K_{\mu 2}]/\Gamma[\pi_{\mu 2}]$  with sufficient *us* experimental error improvement
- *us* experimental uncertainties currently BF dominated
- BFs more easily improved experimentally than exclusive mode  $dR/ds$  distribution contributions
- Near-term low-multiplicity mode progress likely [combined BaBar, Belle (+Belle II) effort on spectral functions from existing B-factory data under way]
- **However** sub-0.5%  $|V_{us}|$  needs sub-%  $R_{us;V+A}^w$  error

- Exclusive  $us$  mode  $w_N$  spectral integral contributions

Relative exclusive mode  $R_{us:V+A}^w$  contributions

Wt	$s_0$ [GeV $^2$ ]	$K$	$K\pi$	$K\pi\pi$	Other
				(B-factory)	
$w_2$	2.15	0.496	0.426	0.062	0.010
	3.15	0.360	0.414	0.162	0.065
$w_3$	2.15	0.461	0.446	0.073	0.019
	3.15	0.331	0.415	0.182	0.074
$w_4$	2.15	0.441	0.456	0.082	0.021
	3.15	0.314	0.411	0.194	0.081

- “Other”: 1999 ALEPH data/MC,  $\sim 25\%$  error  
 $\Rightarrow$  “sufficient improvement” includes experimentally (much) more challenging higher-multiplicity modes

## A PROMISING $\tau$ -BASED ALTERNATIVE

- Work with J. Hudspith, T. Izubuchi, R. Lewis, H. Ohki, C. Lehner + … (RBC/UKQCD)
- Basic idea: generalized dispersion relations for products of combination  $\tilde{\Pi}$  of  $J = 0, 1$  *us* V+A polarizations with weights having poles at Euclidean  $Q^2$ 
  - $\tilde{\Pi}(Q^2)$ : polarization sum with spectral function  $\tilde{\rho}(s)$  (experimental  $dR_{us;V+A}/ds$ )
  - Theory: Lattice *us* 2-point function data (no OPE)
  - Weights tunable, allow suppression of larger-error, higher-multiplicity *us* spectral contributions

## More on the lattice-inclusive $us \tau$ approach

- $|V_{us}|^2 \tilde{\rho}_{us;V+A}(s)$  from experimental  $dR_{us;V+A}/ds$

$$\tilde{\rho}_{us;V+A}(s) \equiv \left(1 + 2\frac{s}{m_\tau^2}\right) \rho_{us;V+A}^{(J=1)}(s) + \rho_{us;V+A}^{(J=0)}(s)$$

(no continuum  $us J = 0$  subtraction required)

- Associated (kinematic-singularity-free) polarization

$$\tilde{\Pi}_{us;V+A}(Q^2) \equiv \left(1 - 2\frac{Q^2}{m_\tau^2}\right) \Pi_{us;V+A}^{(J=1)}(Q^2) + \Pi_{us;V+A}^{(J=0)}(Q^2)$$

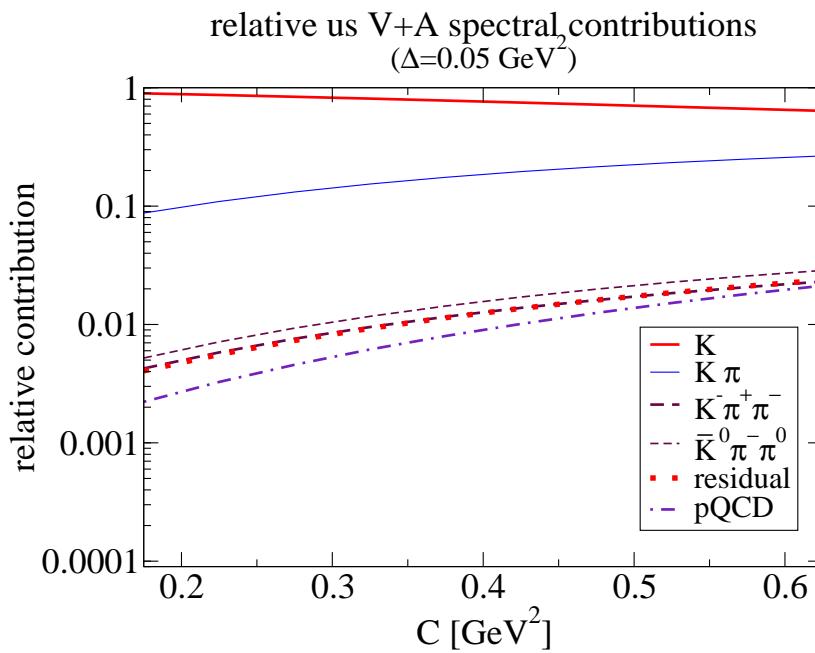
- $\tilde{\rho}_{us;V+A}(s) \sim s$  as  $s \rightarrow \infty$

- For weights  $w_N(s) \equiv \frac{1}{\prod_{k=1}^N (s+Q_k^2)}$ ,  $N \geq 3$ , obtain convergent, unsubtracted 'dispersion relation'

$$\int_{th}^{\infty} ds w_N(s) \tilde{\rho}_{us;V+A}(s) = \sum_{k=1}^N \frac{\tilde{\Pi}_{us;V+A}(Q_k^2)}{\prod_{j \neq k} (Q_j^2 - Q_k^2)}$$

- Lattice data for  $\tilde{\Pi}_{us;V+A}(Q_k^2)$  on RHS
- LHS from experimental  $dR_{us;V+A}/ds$ , up to  $|V_{us}|^2$
- $w_N(s)$ : rapid fall-off if all  $Q_k^2 < 1 \text{ GeV}^2$   
**⇒  $K, K\pi$  dominate LHS, near-endpoint multi-particle,  $s > m_\tau^2$  contributions strongly suppressed**
- Optimization: increasing  $\{Q_k^2\}$  decreases RHS lattice error, increases LHS experimental error

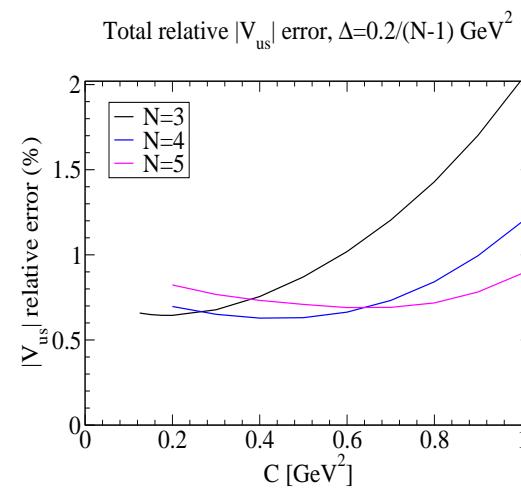
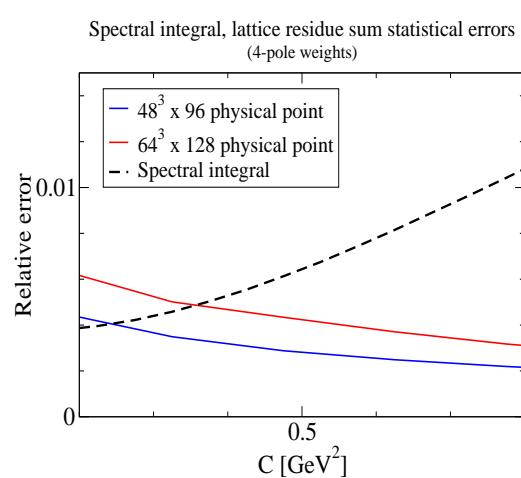
- Below: uniform pole spacing  $\Delta$ , centroid  $C$
- “Tuning” impact example [ $N = 4$ ,  $\Delta = 0.05 \text{ GeV}^2$ ]



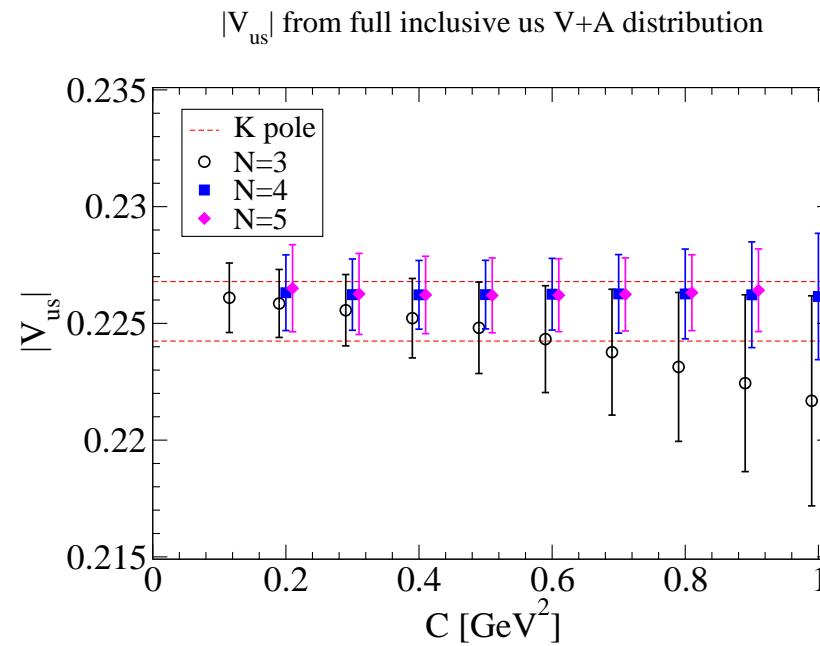
- Sample experimental, statistical, and total errors,  
RBC/UKQCD near-physical-point ensembles

Left: exp't, stat errors vs  $C$ ,  $N = 4$ ,  $\Delta = 0.05 \text{ GeV}^2$

Right: total error vs  $C$ ,  $N = 3, 4, 5$ ,  $\Delta = \frac{0.2}{N-1} \text{ GeV}^2$



- Preliminary  $|V_{us}|$  results



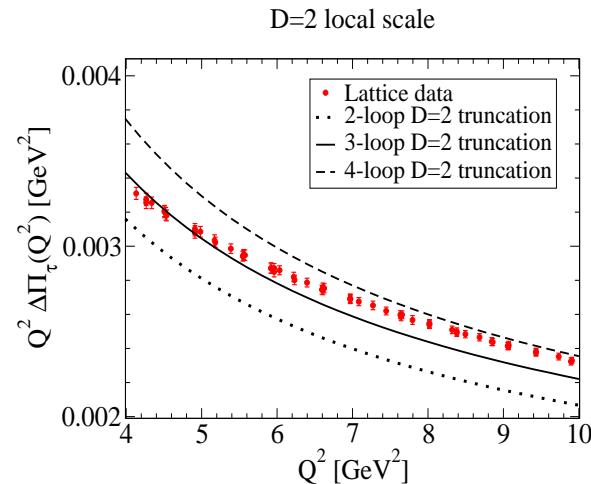
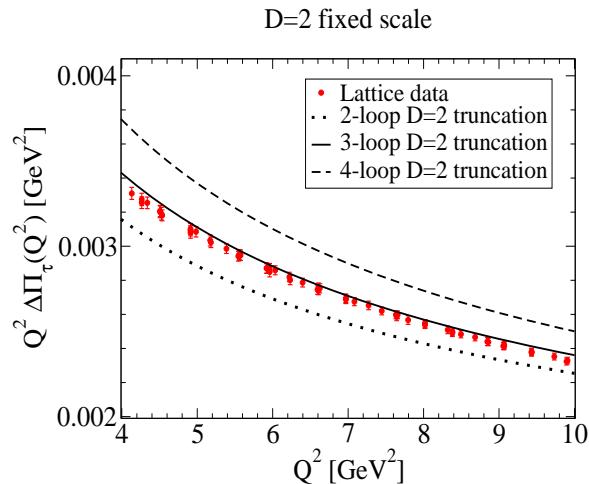
- Full optimization/systematics (FV, continuum limit, IB, small  $m_{\ell,s}$  mistuning corrections) in progress

## SUMMARY

- Old  $3\sigma$  low inclusive FB  $\tau$  FESR  $|V_{us}|$  problem resolved
  - Alternate, no-assumptions implementation:  $|V_{us}|$  higher by  $\sim 0.0020$ , compatible with other determinations
  - Near-term improvements feasible through improvements in  $us$  exclusive mode BFs
  - Highly favorable theoretical error situation
  - However, for competitive  $|V_{us}|$  need improvements to old ALEPH higher-multiplicity, low-statistics data [unlikely in the near-term]

- Advantage of new lattice-inclusive *vs*  $V+A$   $\tau$  approach
  - Theory:
    - \* Lattice in place of OPE; no *us*  $J = 0$  subtraction; improvement through increased statistics
    - \* Easily parasitic on lattice  $a_\mu$  effort (a major effort in the lattice community)
  - Spectral integrals:
    - \* Theory errors still small for weights strongly suppressing higher multiplicity contributions
    - \* Strong  $K, K\pi$  dominance of spectral integral
    - \* Significant experimental improvements possible through just improved  $K\pi$  BFs, distributions

## BACKUP SLIDES



- Higher  $Q^2$ : best (excellent) lattice vs  $D = 2 + 4$  OPE match for 3-loop-truncated, fixed-scale  $D = 2$
- Fixed scale suggests FOPT for FESR  $D = 2$

## OPE, SPECTRAL INPUT

- PDG, FLAG, HPQCD input for  $D = 2, 4$  OPE
- $ud$  V+A spectral data from ALEPH 2013
- $us$  V+A spectral data from sum over exclusive modes [ $> 90\%$  of  $B_{us}^{TOT}$  from  $K_{\ell 2}$ , Belle, BaBar  $K\pi$ ,  $K\pi\pi$ ,  $3K$  results; residual: 1999 ALEPH]
- Favored  $K\pi$  normalization: including preliminary BaBar  $B[\tau \rightarrow K^-\pi^0\nu_\tau]$  update (Adametz thesis)

## MORE ON THE $us$ DATA

- $K$  pole via  $f_K|V_{us}|$  from  $K_{\ell 2}$
- Rather precise unit-normalized  $K^-\pi^0$ ,  $\bar{K}^0\pi^-$ ,  $K^-\pi^+\pi^-$ ,  $\bar{K}^0\pi^-\pi^0$ ,  $3K$  distributions from Belle, BaBar (main uncertainties from BFs)
- $K$ , B-factory modes over 90% of  $B_{us}^{TOT}$
- Residual  $us$  exclusive mode contributions (1999 ALEPH data, covariances) involves significant MC input

## THE EXPERIMENTAL $K\pi$ BF SITUATION

- HFAG 2014  $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0126$
- HFAG 2014  $B[K^-\pi^0\nu_\tau] = 0.00433(15)$  value → preliminary BaBar (Adametz thesis) result  $0.00500(15)$  yields  $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau] = 0.0134$
- Central  $B[(K^-\pi^0 + \bar{K}^0\pi^-)\nu_\tau]$  from  $K_{\ell 3}$ , dispersive analysis expectations [ACLP13]: 0.0133
- 0.07% BF difference “small” but represents  $\sim 2.4\%$  of  $B_{us}^{TOT}$ , hence  $\sim 1.2\%$  increase in  $|V_{us}|$

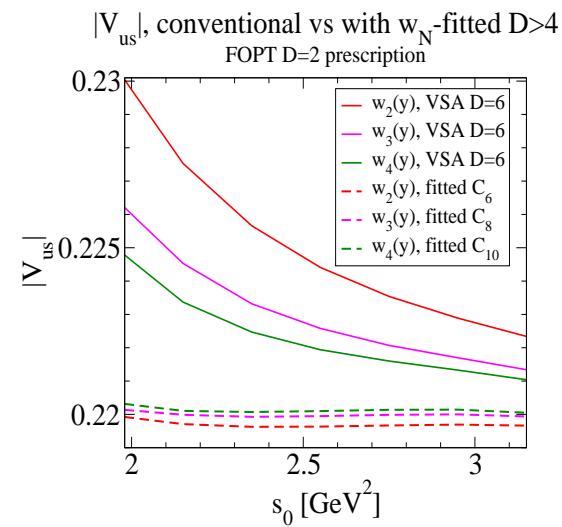
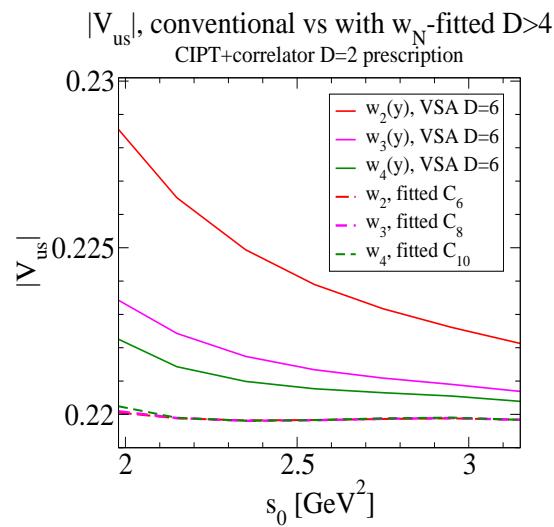
## Results for $|V_{us}|$ for current HFAG 2014 $K\pi$ BFs

Weight	$ V_{us} $ CIPT+corr $D = 2$	$ V_{us} $ FOPT $D = 2$
$w_2$	0.21985(230)	0.21966(230)
$w_3$	0.21985(231)	0.21966(231)
$w_4$	0.21985(231)	0.22009(231)

## Error budget, existing $K\pi$ BFs

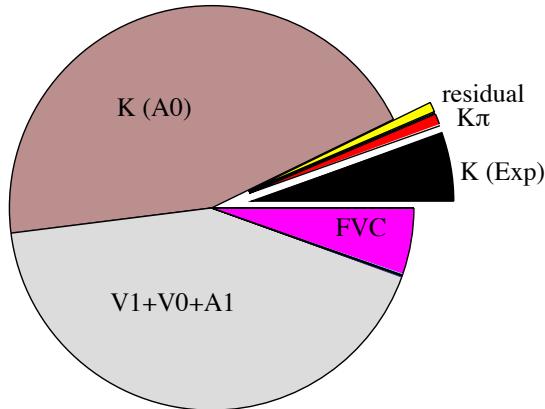
Source	$\delta V_{us} $ ( $w_2$ FESR)	$\delta V_{us} $ ( $w_3$ FESR)	$\delta V_{us} $ ( $w_4$ FESR)
$\delta\alpha_s$	0.00001	0.00003	0.00005
$\delta m_s(2 \text{ GeV})$	0.00017	0.00018	0.00020
$\delta\langle m_s \bar{s}s \rangle$	0.00034	0.00034	0.00034
$\delta(\text{long corr})$	0.00009	0.00009	0.00009
$ud$ exp	0.00027	0.00027	0.00027
$us$ exp	0.00229	0.00229	0.00230

# Stability of $|V_{us}|$ with fitted $C_{2N+2}$ input, existing $K\pi$ BF normalization

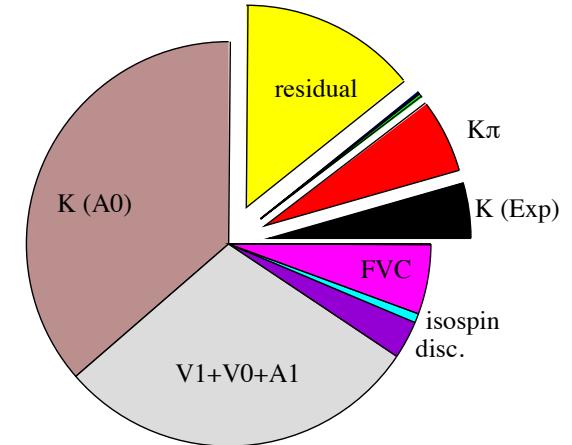


# $|V_{us}|$ error $^2$ budget

error $^2$  (N=4, C=0.2)



error $^2$  (N=4, C=0.5)



$V_{us}$	error $^2$ (%)
<i>Exp.</i>	6
$K(A0)$	45
$V1+V0+A1$	43
<i>disc.</i>	0
<i>isospin</i>	0
<i>FVC</i>	5
<i>OPE</i>	0

$V_{us}$	error $^2$ (%)
<i>Exp.</i>	25
$K(A0)$	36
$V1+V0+A1$	29
<i>disc.</i>	3
<i>isospin</i>	1
<i>FVC</i>	6
<i>OPE</i>	0