

# $B$ -meson semileptonic rare decay phenomenology from lattice QCD

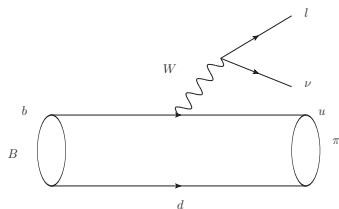
Ran Zhou

Fermilab

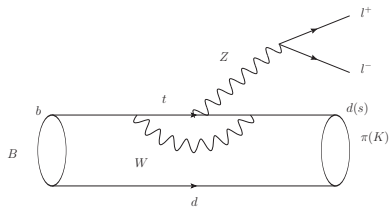
(ICHEP 2016, Chicago, US)

08/06/2016

# Heavy to light semileptonic decays



Tree-level diagram



Loop-level diagram

- $B$ -meson semileptonic decays through tree-level diagram ( $b \rightarrow ul\nu$ ).  
For example,  $B \rightarrow \pi l\nu$ ,  $B_s \rightarrow K l\nu$ ,  $\Lambda_b \rightarrow p l\nu$
- $B$ -meson semileptonic decays through loop-level diagram  
( $b \rightarrow s(d)ll$ ) For example,  $B \rightarrow K(\pi)l^+l^-$ ,  $B \rightarrow K(\pi)\nu\bar{\nu}$ ,  $\Lambda_b \rightarrow \Lambda ll$

# Standard Model prediction

The Standard Model prediction can be written in a generic form:

$$\text{Theo. pred.} = (\text{prefactors}) \times (\text{CKMfactor}) \times \langle f | \hat{O} | i \rangle$$

- Prefactors contain the Wilson coefficients (short distance physics).
- CKM factor depends on the processes.
- Lattice QCD calculates  $\langle f | \hat{O} | i \rangle$  non-perturbatively from first principle.  
(long distance physics)

# Hadronic matrix elements and form factors

- Matrix elements in  $B \rightarrow K(\pi)ll$  and  $B \rightarrow \pi l\nu$  processes:

$$\langle B(p) | \bar{b} \gamma^\mu s | K(k) \rangle, \langle B(p) | \bar{s} \sigma^{\mu\nu} b | K(k) \rangle$$

$$\begin{aligned} \langle B(p) | \bar{b} \gamma^\mu s | K(k) \rangle &= f_+(p^\mu + k^\mu - \frac{m_B^2 - m_K^2}{q^2} q^\mu) + f_0 \frac{m_B^2 - m_K^2}{q^2} q^\mu \\ &= \sqrt{2m_B} \left[ f_{\parallel} \frac{p^\mu}{m_B} + f_{\perp} k_{\perp}^\mu \right] \end{aligned}$$

$$\begin{cases} f_{\parallel}(E_K) = \frac{\langle B(p) | \bar{b} \gamma^0 s | K(k) \rangle}{\sqrt{2m_B}} \\ f_{\perp}(E_K) = \frac{\langle B(p) | \bar{b} \gamma^i s | K(k) \rangle}{2\sqrt{m_B}} \frac{1}{p_i} \end{cases}$$

$$\begin{cases} f_0(E_K) = \frac{2m_B}{m_B^2 - m_K^2} [(m_B - E_K) f_{\parallel}(E_K) + (E_K^2 - m_K^2) f_{\perp}(E_K)] \\ f_+(E_K) = \frac{1}{\sqrt{2m_B}} [f_{\parallel}(E_K) + (m_B - E_K) f_{\perp}(E_K)] \end{cases}$$

# Hadronic matrix elements and form factors

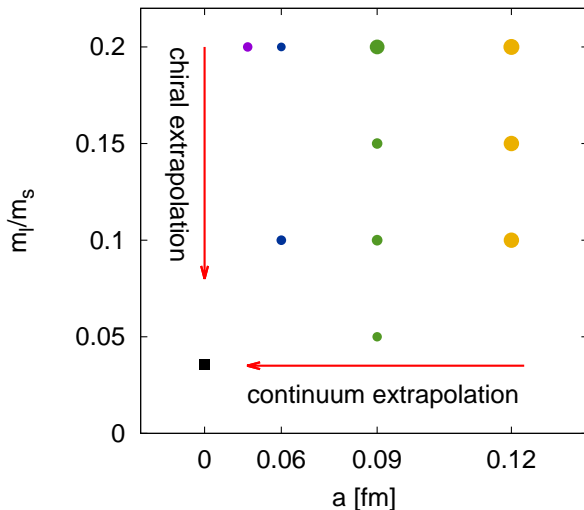
Semileptonic  $B \rightarrow K l l$  transition from tensor current:

$$q_\nu \langle K(k) | \bar{s} \sigma^{\mu\nu} b | B(p) \rangle = \frac{i f_T}{m_B + m_K} [q^2(p^\mu + k^\mu) - (m_B^2 - m_K^2)q^\mu]$$

Solve for  $f_T$ :

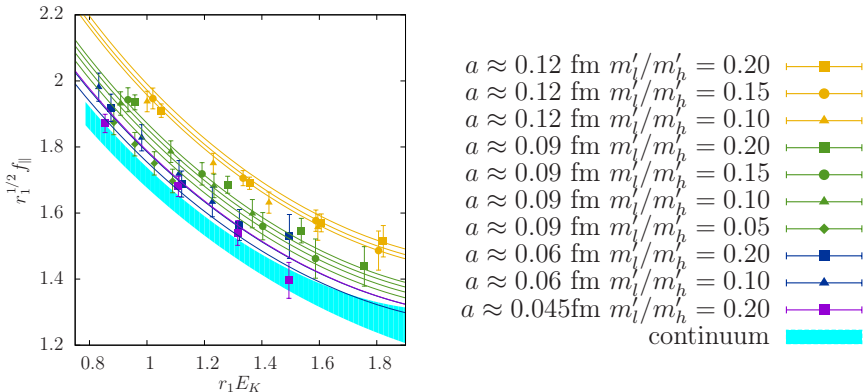
$$f_T = \frac{m_B + m_K}{\sqrt{2m_B}} \frac{\langle K(k) | i b \sigma^{0i} s | B(p) \rangle}{\sqrt{2m_B} k^i}$$

# Lattice ensembles used in $B \rightarrow K(\pi)\ell\ell$ calculations



**Figure :** Ensembles of QCD gauge field configurations used in the simulations.

# $f_{\parallel}$ , $f_{\perp}$ chiral-continuum extrapolations



**Figure :** Chiral-continuum extrapolation on  $f_{\parallel}$  (left).

- Lattice-QCD simulations are in the low  $E_{K(\pi)}$  region.
- Form factors in the full  $q^2$  range are obtained from z-expansion fit.

# Heavy to light semileptonic rare decays (Loop level)

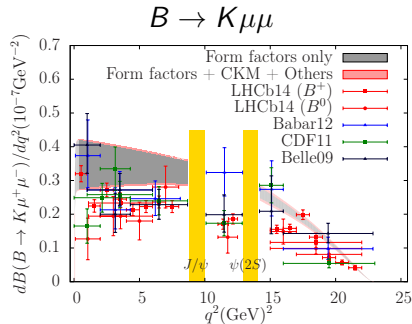
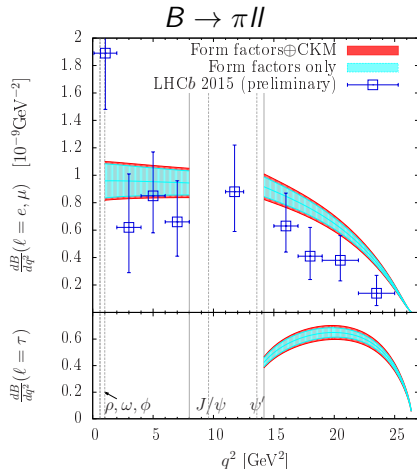
New results on  $B$ -meson rare decays

	Fermilab/MILC	Fermilab/MILC	Detmold and Meinel
process	$B \rightarrow Kll$ ,	$B \rightarrow \pi ll$	$\Lambda_b \rightarrow \Lambda$
kinematics	full $q^2$	full $q^2$	full $q^2$
ensembles	MILC asqtad	MILC asqtad	RBC/UKQCD DWF
$N_f$	2+1	2+1	2+1
$a$	4/0.045-0.12	4/0.045-0.12	2/0.09-0.12
$M_\pi^{\min}$	260	260	227
light quark	asqtad	asqtad	DWF
$b$ quark	Fermilab	Fermilab	RHQ
Ref.	PRD.93.025026	PRL.115.152002	PRD.93.074501

- PRD.93.034005 (Fermilab/MILC,  $B$  rare decay pheno)
- PRD.94.013007 (Meinel and van Dyk,  $\Lambda_b$  rare decay pheno)
- PRD.88.054509, PRL.111.162002 (HPQCD,  $B \rightarrow Kll$  ff and pheno),  
PRD.89.094501, PRL.112.212003 ( $B \rightarrow K^*ll$  ff and pheno)



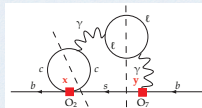
# Standard Model predictions of $B$ rare decays



- Standard-Model predictions of the differential decay rate in  $B \rightarrow \pi l l$  and  $B \rightarrow K l l$  process (PRL.115.152002, PRD.93.034005).

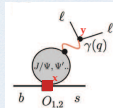
# Resonance states and non-lattice errors

● Note the difference between inclusive and exclusive (high- $q^2$ ) OPE:



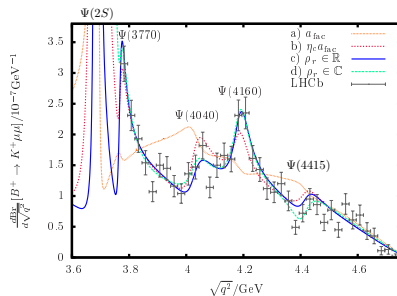
$$(x-y)^2 \sim \frac{1}{(m_b - \sqrt{q^2})^2}$$

The breakdown of the OPE at very large  $q^2$  is independent of the presence of resonant charm loops



$$(x-y)^2 \gg \frac{1}{q^2}$$

The presence of resonant charm loops jeopardize the OPE itself and one has to rely on quark-hadron duality [Beylich, Buchalla, Feldmann]

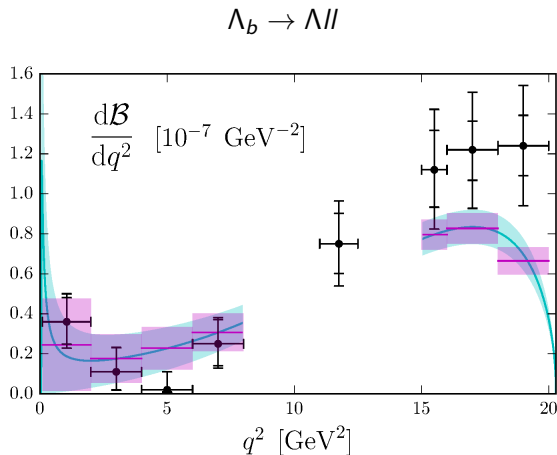


E. Lunghi at KITP 2015

1406.0566

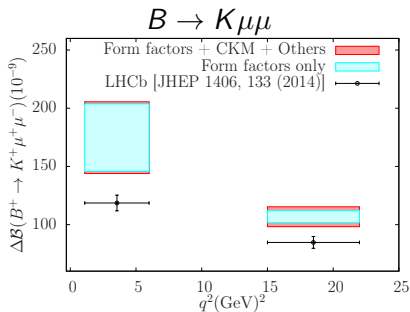
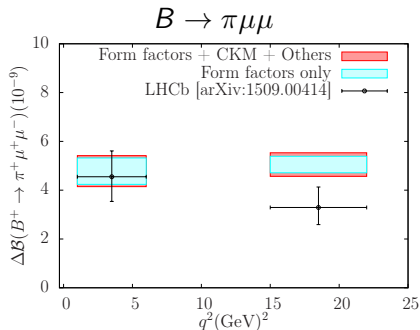
- As the form factor errors become smaller, cautious is needed to treat the non-lattice errors in the Standard-Model predictions.
- The resonance states could introduce the violation of quark-hadron duality in the high  $q^2$  range (1406.0566, PRD.70.114005, EPJC.71.1625).
- For all non-lattice errors, please refer to E. Lunghi's talk at KITP Program "Lattice Gauge Theory for the LHC and Beyond"

# Standard Model predictions of $B$ rare decays



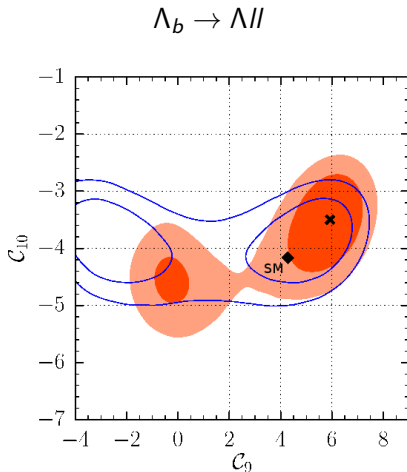
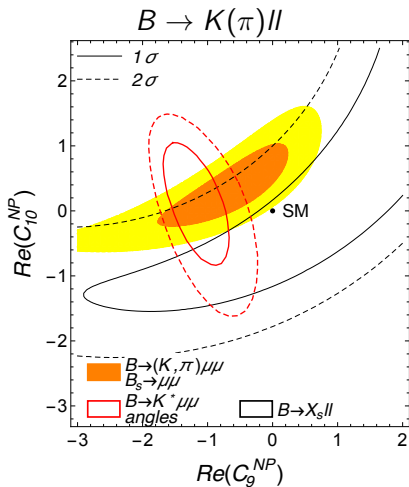
- Standard-Model predictions of the differential decay rate in  $\Lambda_b \rightarrow \Lambda l l$  processes. (PRD.93.074501)

# Standard Model predictions of $B$ rare decays



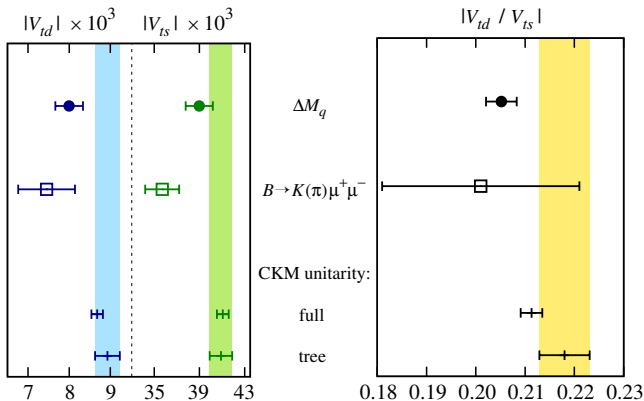
- Four measurements combined disfavor the Standard-Model hypothesis at the  $1.7\sigma$  level (PRD.93.034005).

# Constraints on New Physics.



**Figure :** Constraints to new physics from  $B \rightarrow K(\pi)ll$  plus  $B_s \rightarrow \mu\mu$  decays (PRD.93.034005) and  $\Lambda_b \rightarrow \Lambda ll$  decay (PRD.94.013007).

# Impact from new B-mixing results



- Determinations of  $|V_{td}|$  and  $|V_{ts}|$ , and their ratio from  $B$ -mixing and rare  $B \rightarrow K(\pi)\mu\mu$  decays (PRD.93.113016, PRD.93.034005).
- Tension between FCNC and tree processes: results from  $B$ -mixing (rare B decays) lie below the determinations from CKMfitters full global unitarity triangle fit using only tree-level inputs by  $1.2\text{--}2.1\sigma$ .

## Theoretical studies of the $B \rightarrow K(\pi)\nu\bar{\nu}$

In the Standard Model, the decay rate for  $B \rightarrow K(\pi)\nu\bar{\nu}$  is:  
(arXiv:1409.4557, arXiv:0902.0160)

$$\frac{dB(B \rightarrow K(\pi)\nu\bar{\nu})}{dq^2} = 3\tau_B |N_{K(\pi)}|^2 \frac{X_t^2}{(\sin^2\theta_W)^4} \rho_{K(\pi)}(q^2), \quad (1)$$

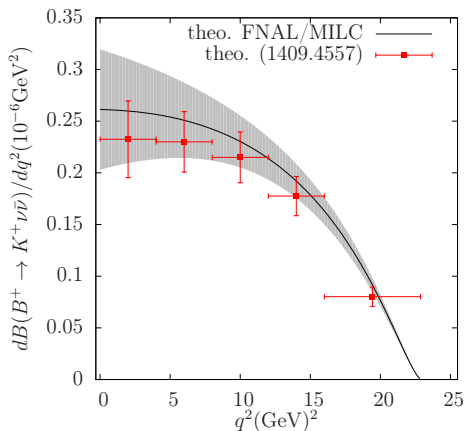
where the numerical coefficient  $N_{K(\pi)}$  depends upon the relevant CKM factors and  $\rho_{K(\pi)}$  is the rescaled hadronic form factor:

$$N_{K(\pi)} = V_{tb} V_{ts(d)}^* \frac{G_F \alpha_{EW}}{16\pi^2} \sqrt{\frac{M_B}{3\pi}}, \quad (2)$$

$$\rho_{K(\pi)}(q^2) = \frac{\lambda^{3/2}(q^2)}{M_B^4} f_+^2(q^2). \quad (3)$$

- The form factor  $f_+$  is the same as in the  $B \rightarrow K(\pi)l^+l^-$  lattice-QCD calculations.

# Theoretical studies of the $B \rightarrow K \nu \bar{\nu}$



- Grey band: theoretical result from FNAL/MILC  $B \rightarrow K$  form factor
- Red points: theoretical result from and lattice-QCD plus LCSR form factor results (arXiv:1409.4557).



# Summary

- There are many progresses in the heavy flavor physics from lattice calculations since 2015.
- Many new  $2\sigma$  hints of NP have been revealed by the improvements of lattice calculations.
- Belle II and LHCb will improve measurements and observe new decays.
- We will continue to sharpen tests of SM and may reveal presence of NP.