

# Standard Model Prediction for Direct CP-violation in K- decays, and Long-Distance Contributions to Kaonic Amplitudes

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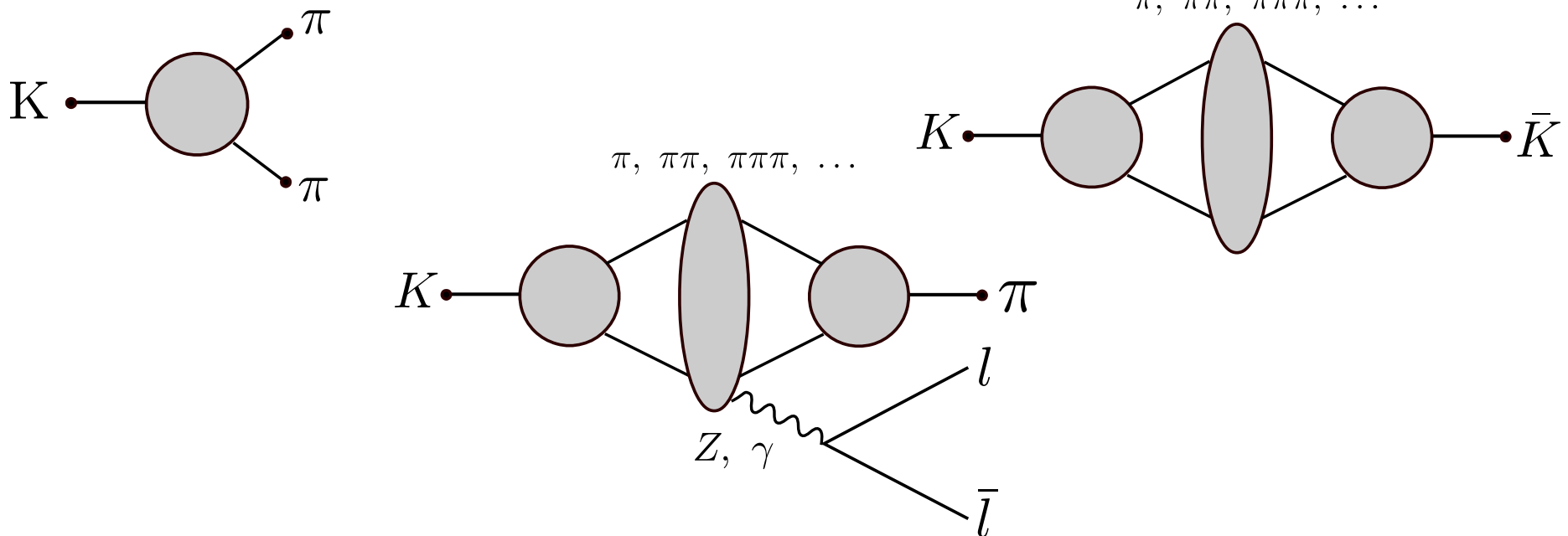
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# Introduction

- Lattice QCD directly simulates low-energy hadronic interactions on supercomputers.
- Only known *ab initio*, **systematically improvable** technique for studying non-perturbative QCD.
- The simulations have now reached levels of precision with which it can make a significant impact on the search for BSM physics.
- In particular, with recent advances in theoretical and computational techniques, we are now able to directly compute matrix elements involving multi-particle states.
- In this talk we focus upon lattice calculations of kaonic matrix elements relating to quantities that are sensitive to BSM physics:



K $\rightarrow$  $\pi\pi$  and  $\varepsilon'/\varepsilon$

# Motivation for studying $K \rightarrow \pi\pi$ Decays

- Likely explanation for matter/antimatter asymmetry in Universe, baryogenesis, requires violation of CP.
- Amount of CPV in Standard Model appears too low to describe measured M/AM asymmetry: tantalizing hint of new physics.
- Direct CPV first observed in late 90s at CERN (NA31/NA48) and Fermilab (KTeV) in  $K^0 \rightarrow \pi\pi$ :

$$\eta_{00} = \frac{A(K_L \rightarrow \pi^0\pi^0)}{A(K_S \rightarrow \pi^0\pi^0)}, \quad \eta_{+-} = \frac{A(K_L \rightarrow \pi^+\pi^-)}{A(K_S \rightarrow \pi^+\pi^-)}.$$

$$\text{Re}(\epsilon'/\epsilon) \approx \frac{1}{6} \left( 1 - \left| \frac{\eta_{00}}{\eta_{\pm}} \right|^2 \right) = 16.6(2.3) \times 10^{-4} \quad (\text{experiment})$$

measure of direct CPV

measure of indirect CPV

- In terms of isospin states:  $\Delta I=3/2$  decay to  $I=2$  final state, amplitude  $A_2$   
 $\Delta I=1/2$  decay to  $I=0$  final state, amplitude  $A_0$

$$A(K^0 \rightarrow \pi^+\pi^-) = \sqrt{\frac{2}{3}}A_0e^{i\delta_0} + \sqrt{\frac{1}{3}}A_2e^{i\delta_2},$$

$$A(K^0 \rightarrow \pi^0\pi^0) = \sqrt{\frac{2}{3}}A_0e^{i\delta_0} - 2\sqrt{\frac{1}{3}}A_2e^{i\delta_2}.$$



$$\epsilon' = \frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}} \left( \frac{\text{Im}A_2}{\text{Re}A_2} - \frac{\text{Im}A_0}{\text{Re}A_0} \right)$$

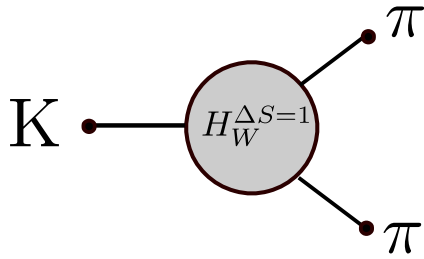
( $\delta_i$  are strong scattering phase shifts.)

$$\omega = \text{Re}A_2/\text{Re}A_0$$

- Small size of  $\epsilon'$  makes it particularly sensitive to new direct-CPV introduced by most BSM models.

# Overview of calculation

- Hadronic energy scale  $\ll M_W$  - use weak effective theory.
- $K \rightarrow \pi\pi$  decays require single insertion of  $\Delta S=1$  Hamiltonian:



10 effective four-quark operators

$$H_W^{\Delta S=1} = \frac{G_F}{\sqrt{2}} V_{ud}^* V_{us} \sum_{j=1}^{10} [z_j(\mu) + \tau y_j(\mu)] Q_j$$

perturbative Wilson coeffs.

$$\tau = -\frac{V_{ts}^* V_{td}}{V_{us}^* V_{ud}} = 0.0014606 + 0.00060408i$$

Imaginary part solely responsible for CPV  
(everything else is pure-real)

LL finite-volume correction

$$A^I = F \frac{G_F}{\sqrt{2}} V_{ud} V_{us} \sum_{i=1}^{10} \sum_{j=1}^7 \left[ (z_i(\mu) + \tau y_i(\mu)) Z_{ij}^{\text{lat} \rightarrow \overline{\text{MS}}} M_j^{I, \text{lat}} \right]$$

renormalization matrix (mixing)

$$M_j^{I, \text{lat}} = \langle (\pi\pi)_I | Q_j | K \rangle \text{ (lattice)}$$

- Operators must be renormalized into same scheme as Wilson coeffs: Use RI-(S)MOM NPR and perturbatively match to  $\overline{\text{MS}}$  at high scale.

# Key challenges of lattice calculation

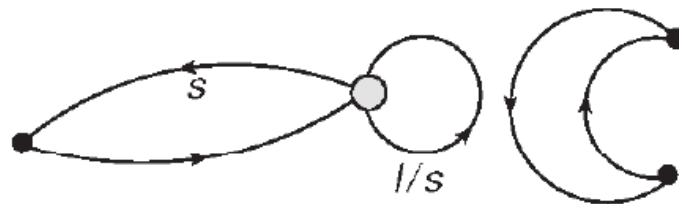
- Primary challenge is to assure physical kinematics. Signal dominated by amplitude with 2 stationary pions in final state.

$$2m_\pi \approx 260 \text{ MeV} \ll m_K \approx 500 \text{ MeV}$$

Solution: Remove stationary pion state from system by manipulating lattice spatial boundary conditions.

- › Antiperiodic BCs on down-quark for  $A_2$
  - › G-parity BCs on both quarks for  $A_0$
- $p_\pi = 0 \rightarrow \pi/L$   
tune L to match  $E_K$  and  $E_{\text{min}}$

- For  $A_0$  serious noise issue due to “disconnected diagrams”



type4

Solution: Use “all-to-all” propagators to tune source to minimize overlap with vacuum and maximize sampling for every configuration.

# Summary of RBC/UKQCD calculations

[Phys.Rev. D91 (2015) no.7, 074502]

- $A_2$  computed on RBC/UKQCD  $64^3 \times 128$  and  $48^3 \times 96$  2+1f Mobius DWF ensembles with the Iwasaki gauge action.
  - $a^{-1} = 2.36$  GeV and  $1.73$  GeV resp - continuum limit taken.
  - 10% and 12% total errors on  $\text{Re}(A_2)$  and  $\text{Im}(A_2)$  resp.
  - Statistical errors sub-percent, dominant systematic errors due to truncation of PT series in computation of renormalization and Wilson coefficients.
- 

[Phys.Rev.Lett. 115 (2015) 21, 212001]

- $A_0$  computed on  $32^3 \times 64$  Mobius DWF ensemble with Iwasaki+DSDR gauge action. G-parity BCs in 3 directions to give physical kinematics.
- Single, coarse lattice with  $a^{-1} = 1.38$  GeV but large physical volume to control FV errors.
- 21% and 65% stat errors on  $\text{Re}(A_0)$  and  $\text{Im}(A_0)$  due to disconn. diagrams and, for  $\text{Im}(A_0)$  a strong cancellation between  $Q_4$  and  $Q_6$ .
- Dominant, 15% systematic error is due again to PT truncation errors exacerbated by low renormalization scale  $1.53$  GeV.



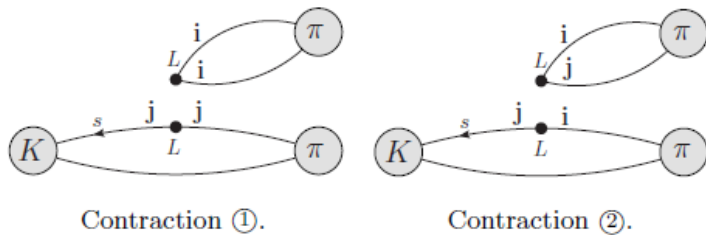
# $\Delta I=1/2$ rule

- In experiment kaons approx 450x (!) more likely to decay into  $I=0$  pi-pi states than  $I=2$ .

$$\frac{\text{Re}A_0}{\text{Re}A_2} \simeq 22.5 \quad (\text{the } \Delta I=1/2 \text{ rule})$$

- Perturbative running to charm scale accounts for about a factor of 2. Is the remaining 10x non-perturbative or New Physics?
- The answer is **low-energy QCD!** RBC/UKQCD [arXiv:1212.1474, arXiv:1502.00263]

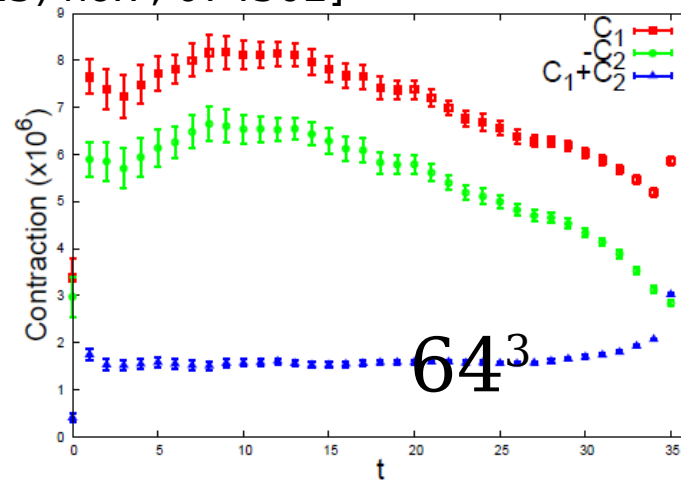
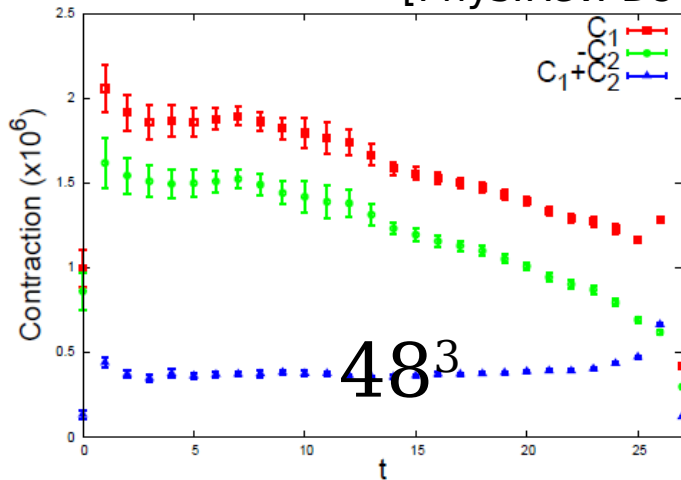
Strong cancellation between the two dominant contractions



$$\text{Re}(A_2) \sim \textcircled{1} + \textcircled{2}$$

find  $\textcircled{2} \approx -0.7\textcircled{1}$  heavily suppressing  $\text{Re}(A_2)$ .

[Phys.Rev. D91 (2015) no.7, 074502]



Pure-lattice calculation

$$\frac{\text{Re}(A_0)}{\text{Re}(A_2)} = 31.1(11.2)$$

[ $\text{Re}(A_0)$  agrees with expt.]

- $\text{Re}(A_0)$  and  $\text{Re}(A_2)$  from expt.
- Lattice values for  $\text{Im}(A_0)$ ,  $\text{Im}(A_2)$  and the phase shifts,

$$\text{Re} \left( \frac{\varepsilon'}{\varepsilon} \right) = \text{Re} \left\{ \frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}\varepsilon} \left[ \frac{\text{Im}A_2}{\text{Re}A_2} - \frac{\text{Im}A_0}{\text{Re}A_0} \right] \right\}$$

$= 1.38(5.15)(4.43) \times 10^{-4},$	(this work)
$16.6(2.3) \times 10^{-4}$	(experiment)

- Total error on  $\text{Re}(\varepsilon'/\varepsilon)$  is  $\sim 3\times$  the experimental error.
- Find reasonable consistency with Standard Model (at  $2.1\sigma$  level).
- Tantalizing hint of discrepancy strong motivation for continued study!

# Outlook

- Error is dominated by those on  $A_0$ .
- Main priority is to increase statistics on  $A_0$  calculation, enabling improved precision and better sys. error estimation. Aim for 4x increase in stats within  $\sim 1$  yr.
- Reduce dominant NPR (15%) and Wilson coefficient (12%) systematic error arising from perturbative  $\overline{\text{MS}}$  matching by increasing renormalization scale. Currently close to completion.
- On longer timescale:
  - Add second lattice spacing and perform continuum limit.
  - Include dynamical charm rather than relying on PT.
  - Investigate (expected %-level) EM and isospin corrections.

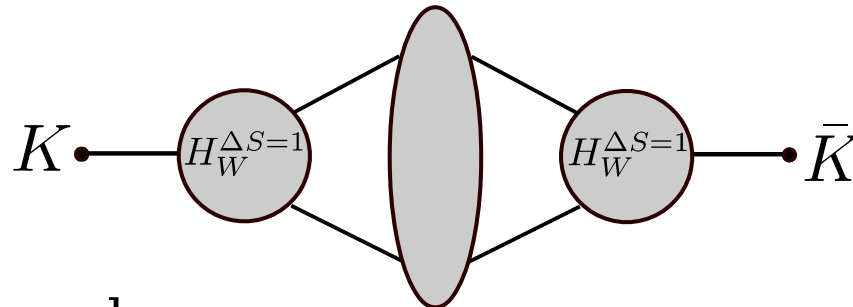
$K_L - K_S$  mass difference

# Introduction

- Neglecting small CP-violation effects, mixing of neutral kaons induced by 2<sup>nd</sup> order weak processes gives rise to mass difference between CP eigenstates  $K_1 \sim K_L$  and  $K_2 \sim K_S$

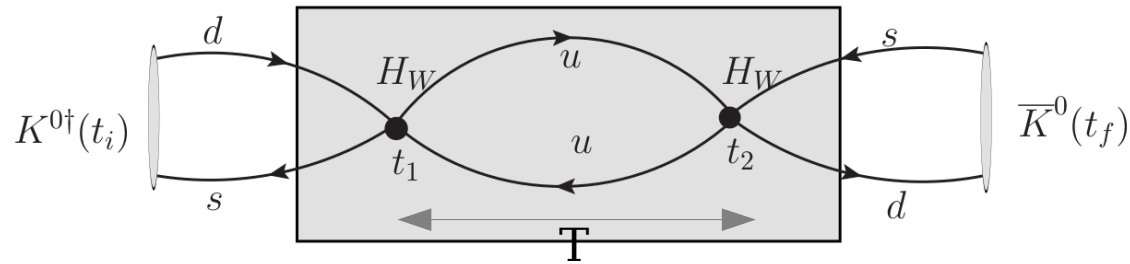
$$\Delta M_K = 2 \sum_n \frac{\langle \bar{K}^0 | H_W | n \rangle \langle n | H_W | K^0 \rangle}{M_K - E_n}$$

- Arises due to FCNC and therefore highly suppressed in SM due to GIM mechanism:  $\Delta m_K = 3.483(6) \times 10^{-12}$  MeV small and highly sensitive to new BSM FCNC.
- PT calc using weak EFT with  $\Delta S=2$  eff. Hamiltonian (charm integrated out) dominated by  $p \sim m_c$ : poor PT convergence at charm scale  $\rightarrow \sim 36\%$  PT sys error.
- PT calc neglects long-distance effects arising when 2 weak operators separated by distance  $\sim 1/\Lambda_{\text{QCD}}$ .
- Use lattice to evaluate matrix element of product of  $H_W^{\Delta S=1, \text{eff}}$  directly:  
 $\pi, \pi\pi, \pi\pi\pi, \dots$



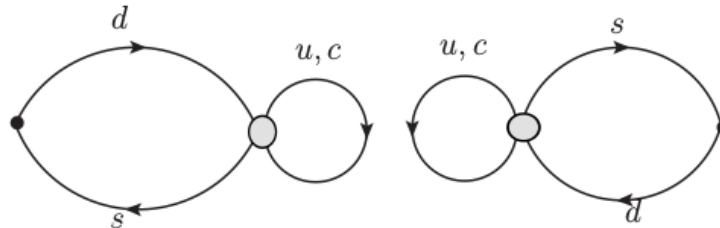
including charm quark.

# Key challenges of lattice calculation



$$\mathcal{A} = N_K^2 e^{-M_K(t_f - t_i)} \sum_n \frac{\langle \bar{K}^0 | H_W | n \rangle \langle n | H_W | K^0 \rangle}{M_K - E_n} \left( -T - \frac{1}{M_K - E_n} + \frac{e^{(M_K - E_n)T}}{M_K - E_n} \right)$$

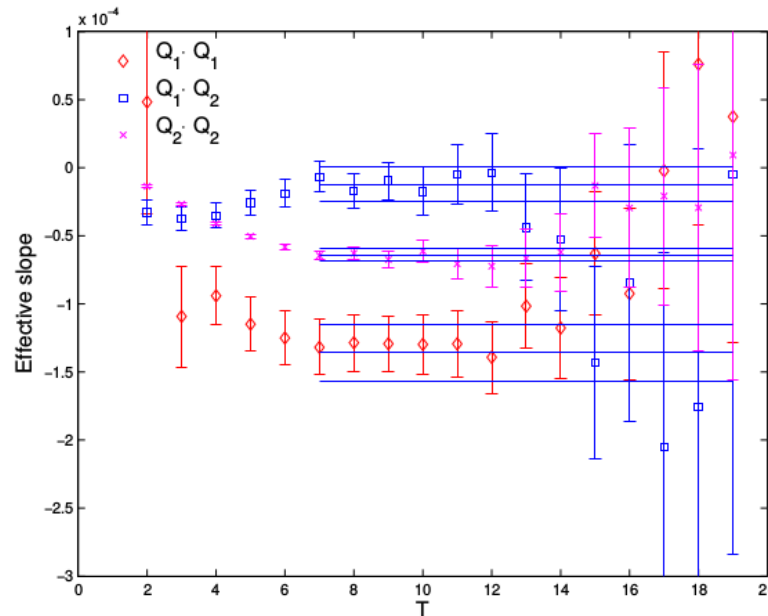
- Vary integration window  $T$  to extract desired matrix element.
- Require explicit subtraction of exponentially-growing term when  $E_n < m_K$
- Disconnected diagrams make the calculation noisy. Requires large statistics.



- Divergence when operators approach removed by GIM - requires (valence) charm on lattice. Need fine lattice to control discretization errors.
- Finite-volume corrections due to multi-particle intermediate states may be large, requiring explicit Lellouch-Lüscher correction.

# Calculation status [Phys.Rev.Lett. 113 (2014) 112003]

- Most recent calculation on one  $24^3 \times 64$  DWF+I lattice with  $a^{-1} = 1.73$  GeV.
- Unphysically heavy pions (330 MeV) - no multi-particle intermediate states.
- Unphysically light charm (949 MeV) for controlled discretization error on this somewhat coarse lattice.
- Large statistics (800 configs) with maximal use of each configuration give good signal for disconnected diagrams.



- Disconnected diagram contribution large, appears to violate OZI “rule”.
- Results in good agreement with expt:

$$\begin{array}{ccc} \text{lattice} & & \text{experiment} \\ \Delta M_K = 3.19(41)(96) \times 10^{-12} \text{ MeV} & & 3.483(6) \times 10^{-12} \text{ MeV} \end{array}$$

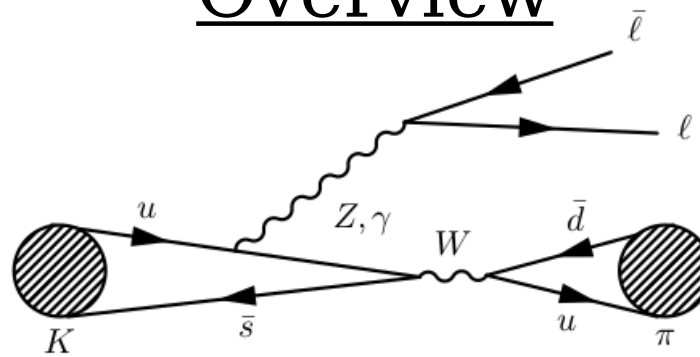
- Dominant error expected to be charm disc. effects  $\sim 30\%$ . Need finer lattices.

Rare kaon decays  $K \rightarrow \pi l \bar{l}$

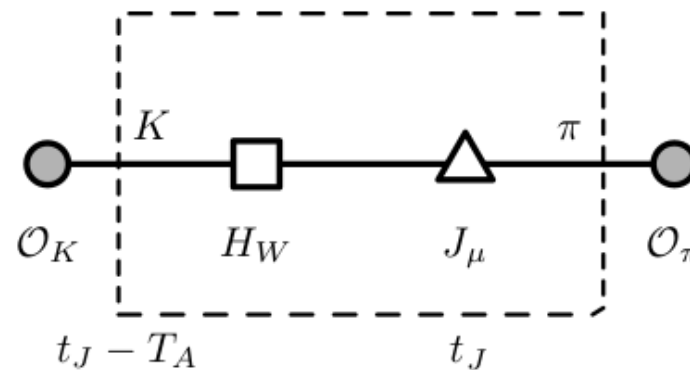


# Overview

[arXiv:1507.03094]



- Additional FCNC processes “rare kaon decays”:  $K^+ \rightarrow \pi^+ l^+ l^-$ ,  $K_S \rightarrow \pi^0 l^+ l^-$
- Amplitude is long-distance dominated: Compute  $K \rightarrow \pi \gamma^*$  on lattice.
- Lattice approach very similar to  $\Delta m_K$  but with EM-current insertion:

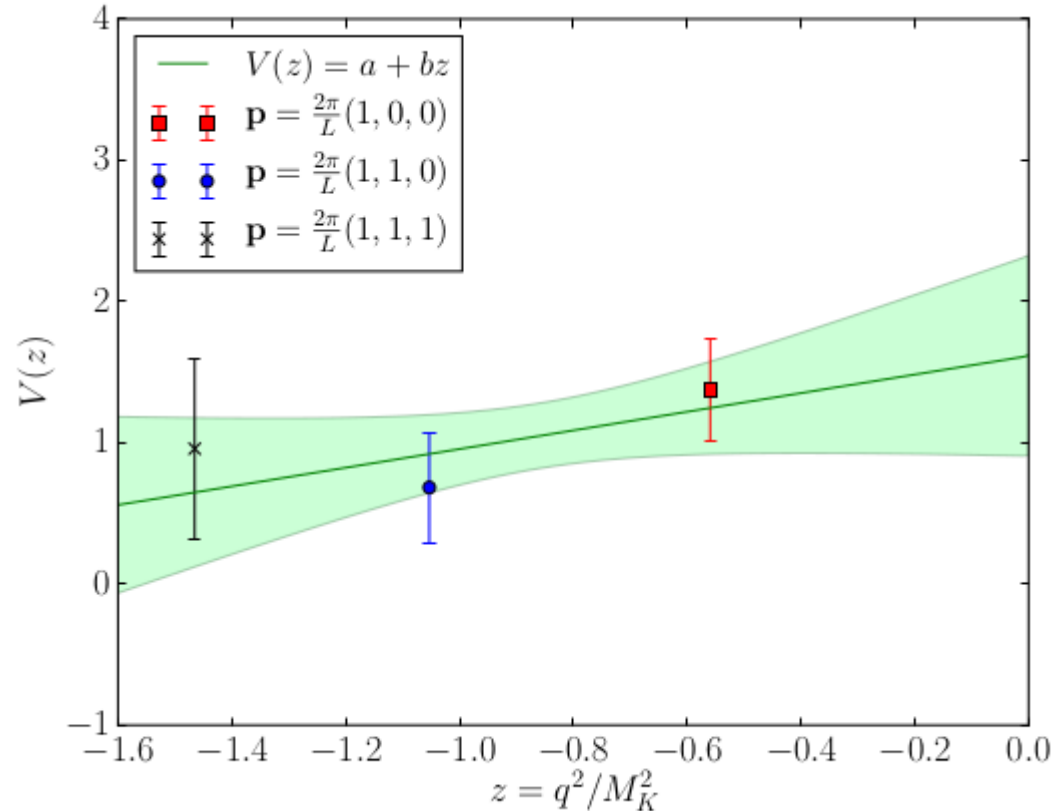


Momentum  
applied to pion to  
study dependence

- Unphysical initial calculation presently underway on  $24^3 \times 64$  ensemble with 430 MeV pions, 620 MeV kaons and  $m_c \sim 533$  MeV.
- Multiple pion momenta allow extraction of form factor.

- Preliminary results from Lattice 2016 talk by A.Lawson (July 28<sup>th</sup>):

$$A_\mu(q^2) \equiv G_F \frac{V(z)}{(4\pi)^2} \left( q^2 (k+p)_\mu - (M_K^2 - M_\pi^2) q_\mu \right), \quad z = q^2/M_K$$

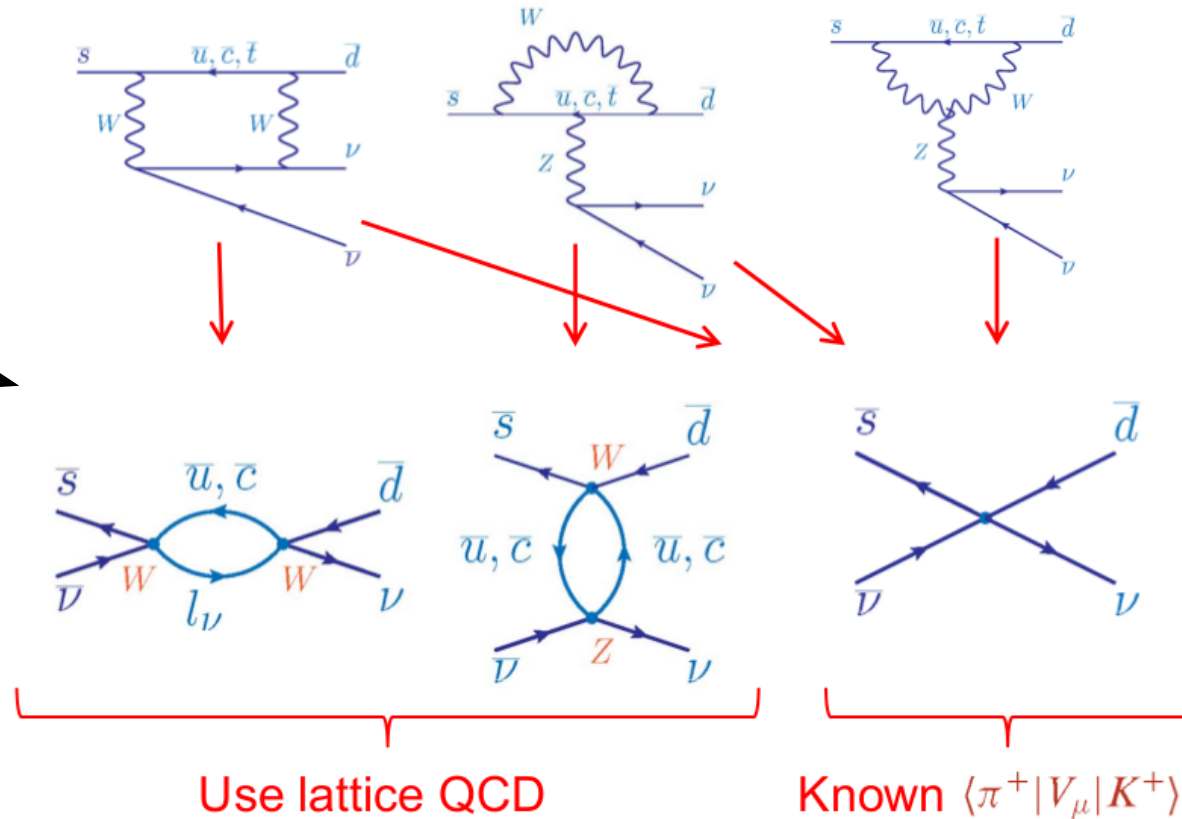


- Future steps are physical pion and kaon masses.
- No statistically significant charm mass dependence observed – is going to physical charm necessary?

Rare kaon decays  $K \rightarrow \pi \nu \bar{\nu}$

- Another FCNC thus far not particularly well experimentally measured.
- NA62 expt expected to provide  $\sim 10\%$  error on 2-3 yr timescale.
- Short-distance dominated but expect  $\sim 5\%$  LD effect: Lattice!
- Again 2x operator insertions:

Effective operators representing both W and Z required



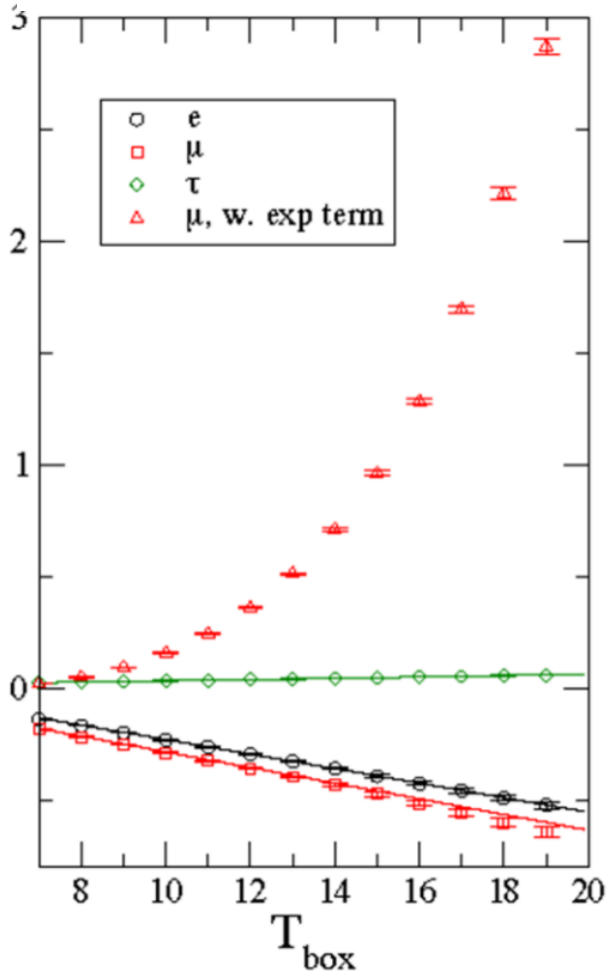
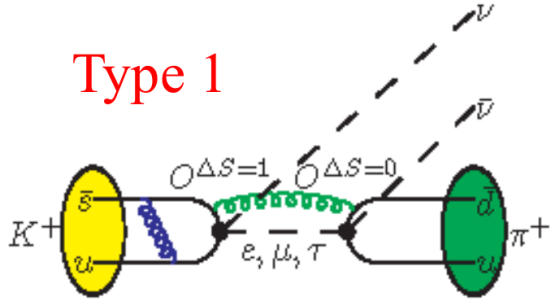
- Short-distance divergence requires NPR "matching" to point operator.

$$\left[ \begin{array}{c} \bar{s} \quad u, c \quad \bar{d} \\ \downarrow \quad \uparrow \quad \downarrow \\ \bar{\nu} \quad W \quad l_\nu \quad W \quad \nu \end{array} + X^{\text{Lat} \rightarrow \text{RI}} \begin{array}{c} \bar{s} \quad \bar{d} \\ \downarrow \quad \uparrow \\ \bar{\nu} \quad \nu \end{array} \right] = 0$$

$p^2 = \mu^2$

- Exploratory calculation on  $16^3 \times 32$   $a^{-1} = 1.73$  GeV ensemble complete.
- Unphysical masses similar to previous slides.
- Preliminary results presented by N. Christ at Lattice 2016 (July 28<sup>th</sup>):

Type 1



- Unphysical kinematics but demonstration of technique:

$$P_C = P_C^{\text{SD}} + \delta P_{cu}$$

total charm contrib.

short-distance part  
from PT: 0.365(12)

bilocal  
contrib.

$$\text{PT } \delta P_{cu} = 0.040(20) \quad [\text{Isodori et al hep-ph/0503107}]$$

$$\Delta P_{cu}(\overline{\text{MS}}, 2 \text{ GeV}) = -0.007(2) \begin{pmatrix} +7 \\ -11 \end{pmatrix}_{\text{RI}} \begin{pmatrix} +5 \\ -21 \end{pmatrix}_{\overline{\text{MS}}}$$

$$WW : -0.032(1) + Z : 0.025(1)$$

- Strong cancellation of WW and Z contribs.
- Wait to observe if situation stays true for more physical kinematics (underway).

# Conclusions

- Lattice techniques now sufficiently advanced to significantly impact search for BSM physics.
- RBC & UKQCD collaboration are leading the way in computing experimentally relevant kaonic amplitudes.
- Complete, physical calculation of  $\varepsilon'/\varepsilon$  performed with hint of tension with experiment.
- Lattice calculation demonstrates  $\Delta I=1/2$  rule arises due to non-perturbative suppression of  $\text{Re}(A_2)$ .
- Matrix elements with 2 operator insertions now possible with controlled systematics, particularly finite-volume errors.
- Early results for FCNC suppressed  $\Delta m_K$  and  $K \rightarrow \pi l \bar{l}$  .
- $K \rightarrow \pi \nu \bar{\nu}$  long-distance contributions computable by similar technique. Preliminary calculation complete.

We hope that these results will spur the experimental community to improve their determinations, particularly for  $\varepsilon'/\varepsilon$ .