

Resonances in Coupled-Channel Scattering from Lattice QCD

David Wilson
for the Hadron Spectrum Collaboration

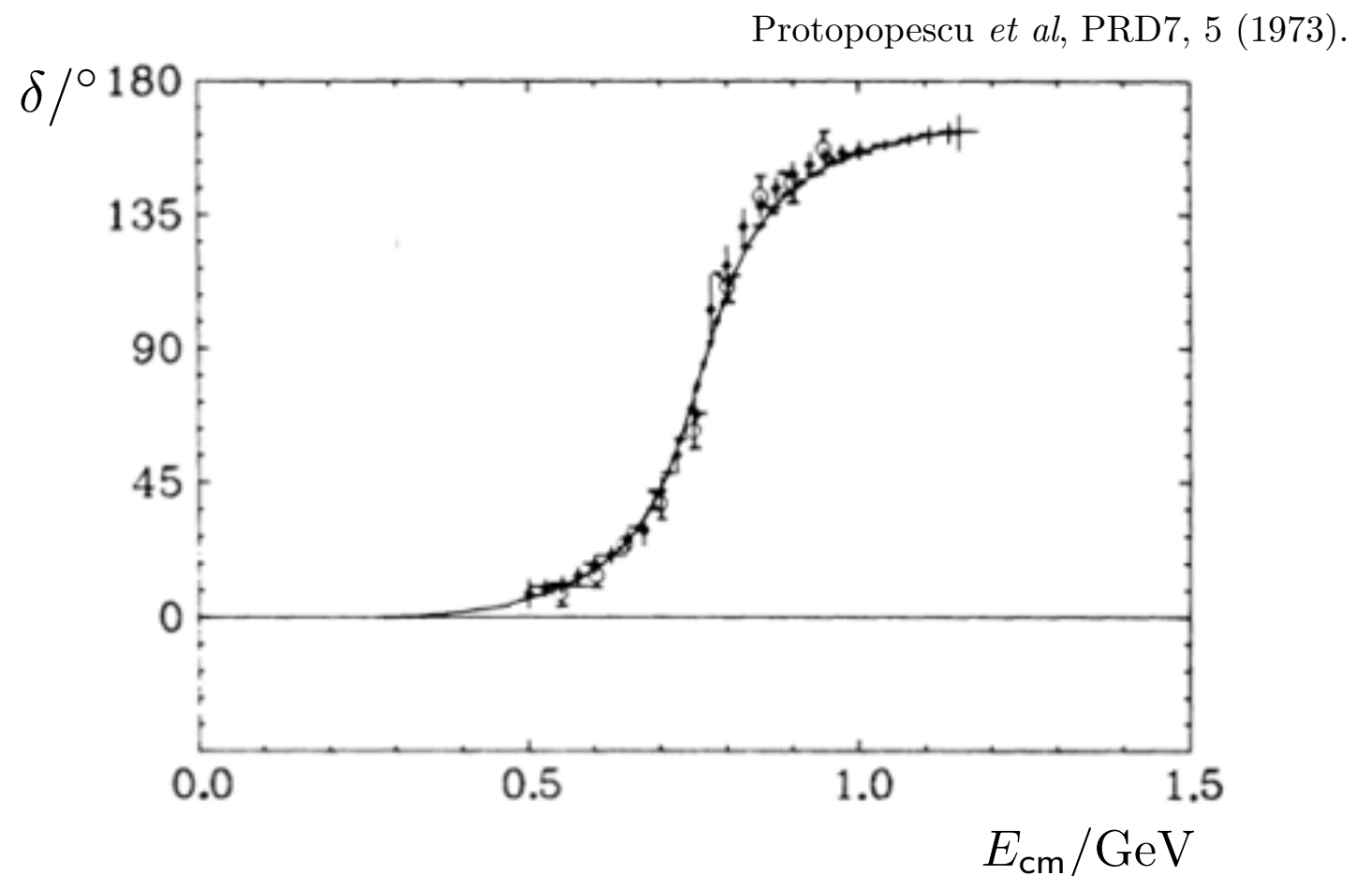
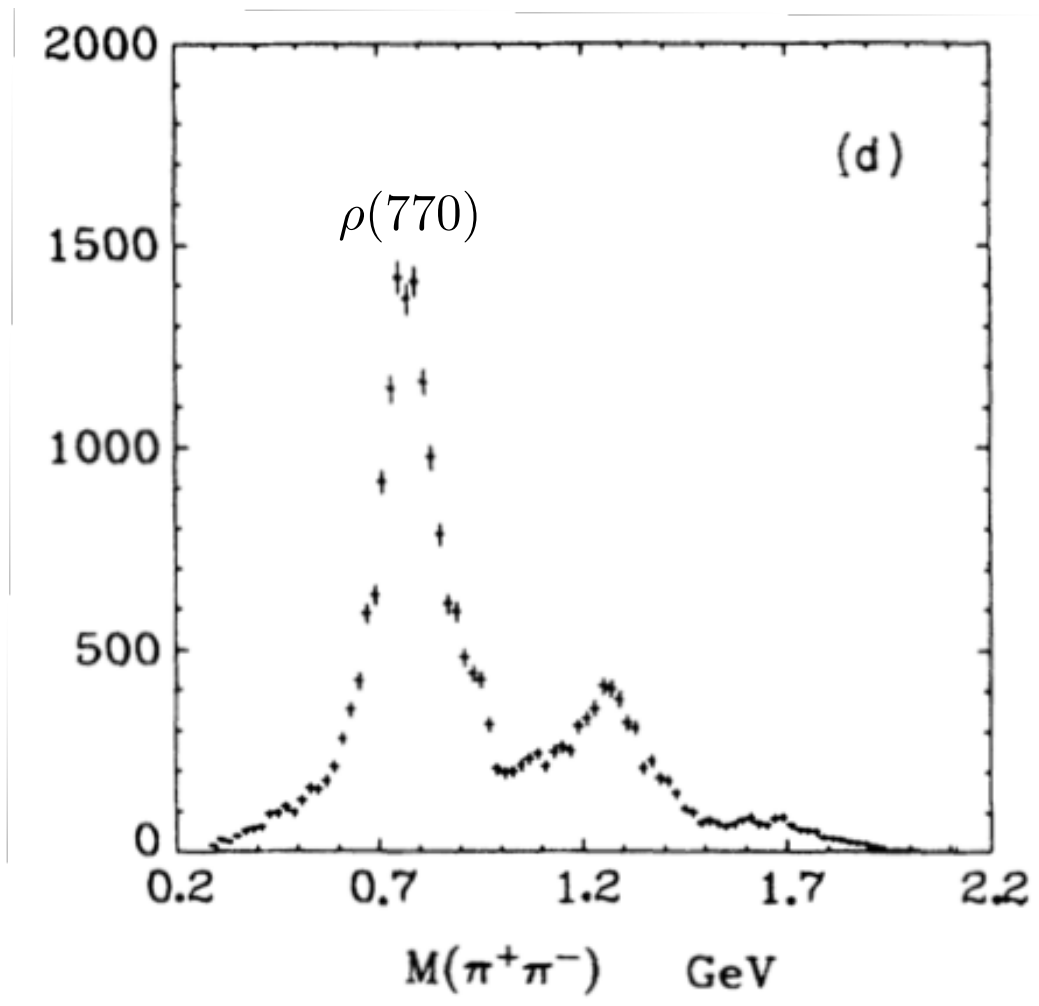
ICHEP 2016 Chicago
3rd-10th August



**UNIVERSITY OF
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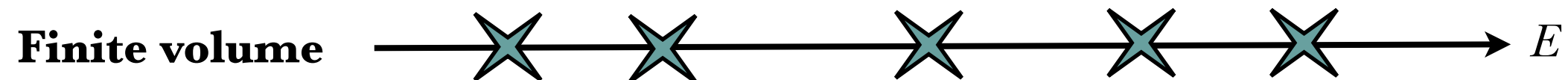
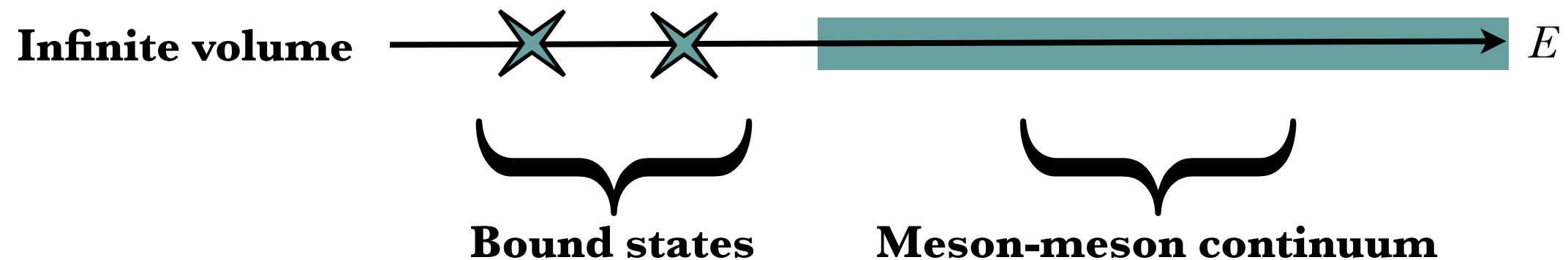
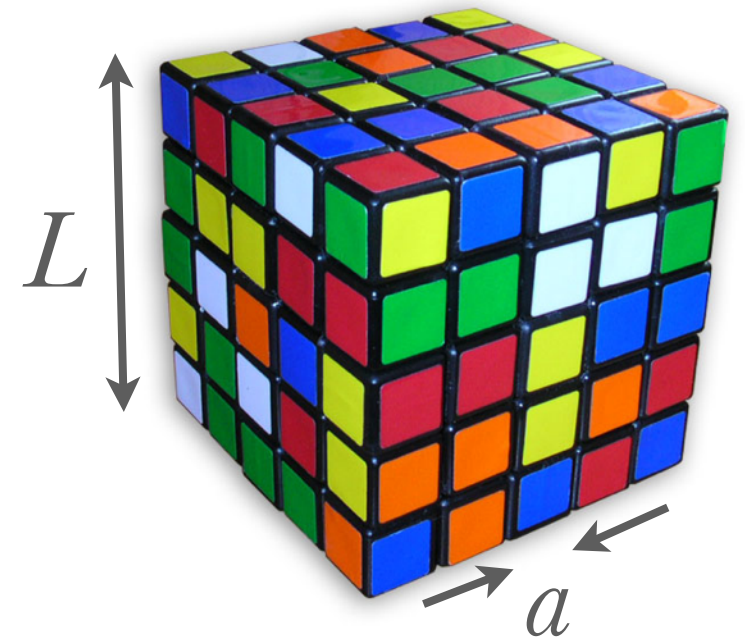
Extracting resonance properties

excited states seen as resonant enhancements
in the scattering of lighter stable particles



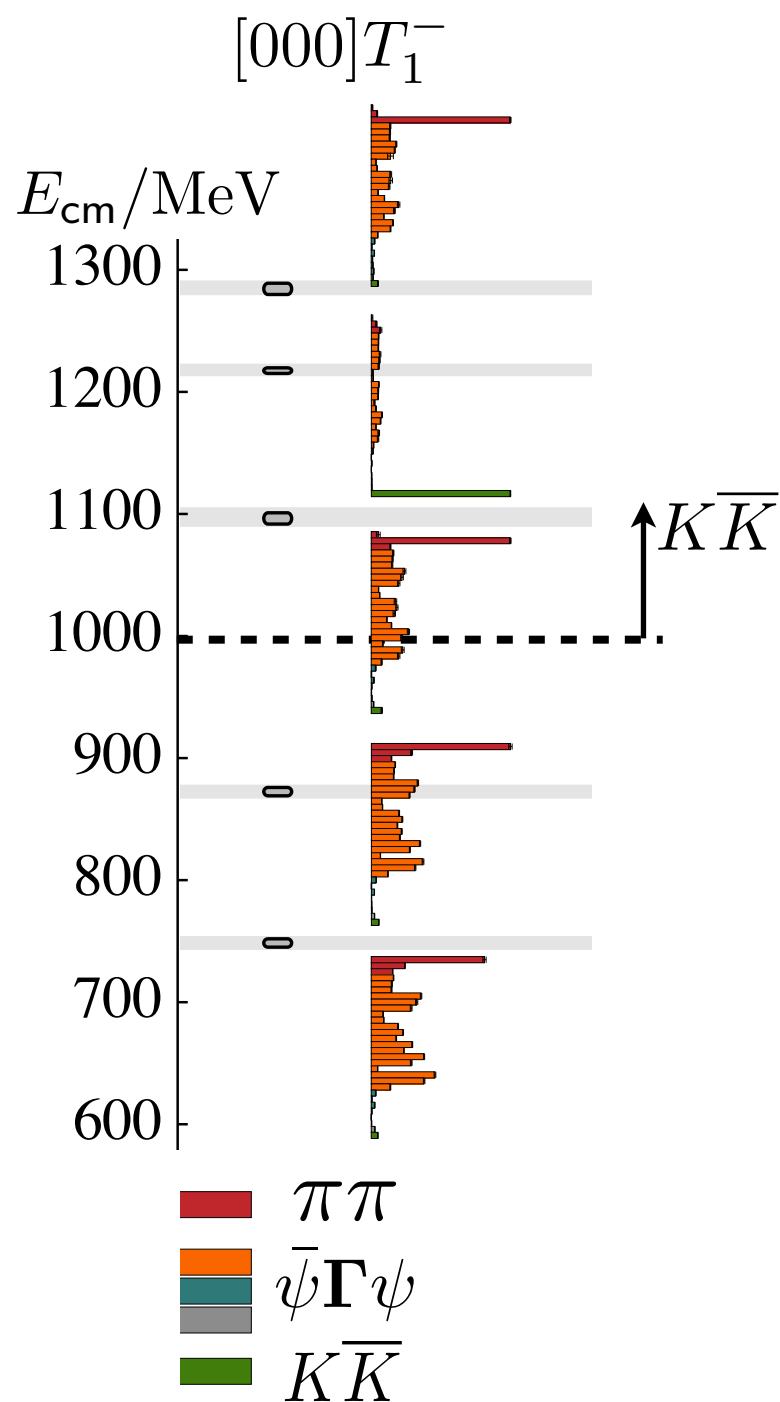
Extracting resonance properties

excited states seen as resonant enhancements
in the scattering of lighter stable particles



ρ resonance

Infinite volume phase shifts from a finite volume



$$\psi(0) = \psi(L), \quad \frac{\partial\psi}{\partial x}\bigg|_{x=0} = \frac{\partial\psi}{\partial x}\bigg|_{x=L}$$

$$\sin\left(\frac{pL}{2} + \delta(p)\right) = 0$$

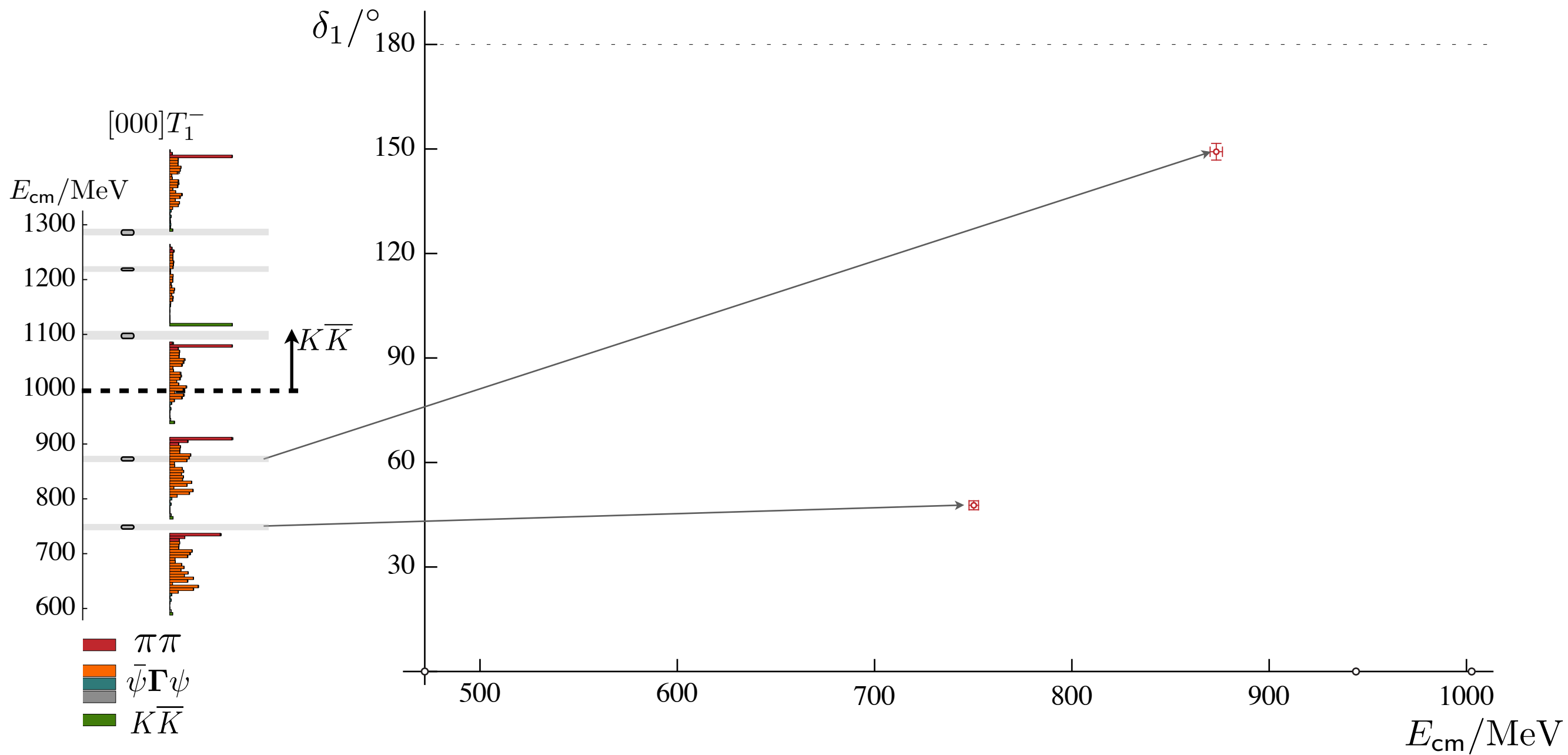
$$p = \frac{2\pi n}{L} - \frac{2}{L}\delta(p)$$

$$m_{\pi} = 236 \text{ MeV}$$

ρ resonance

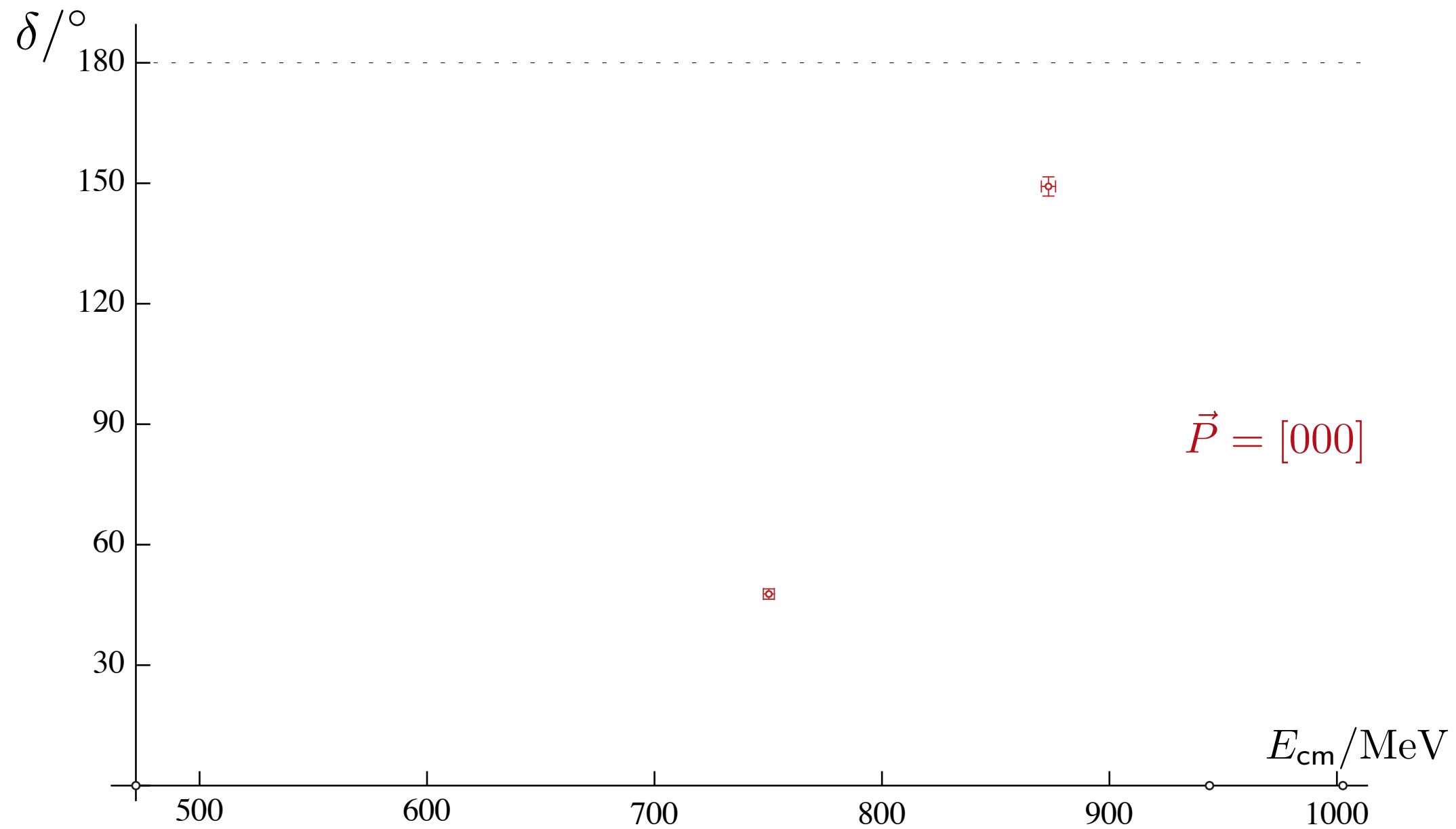
Phase shifts via the Lüscher method: $\tan \delta_1 = \frac{\pi^{3/2} q}{\mathcal{Z}_{00}(1; q^2)}$

$$\mathcal{Z}_{00}(1; q^2) = \sum_{n \in \mathbb{Z}^3} \frac{1}{|\vec{n}|^2 - q^2}$$



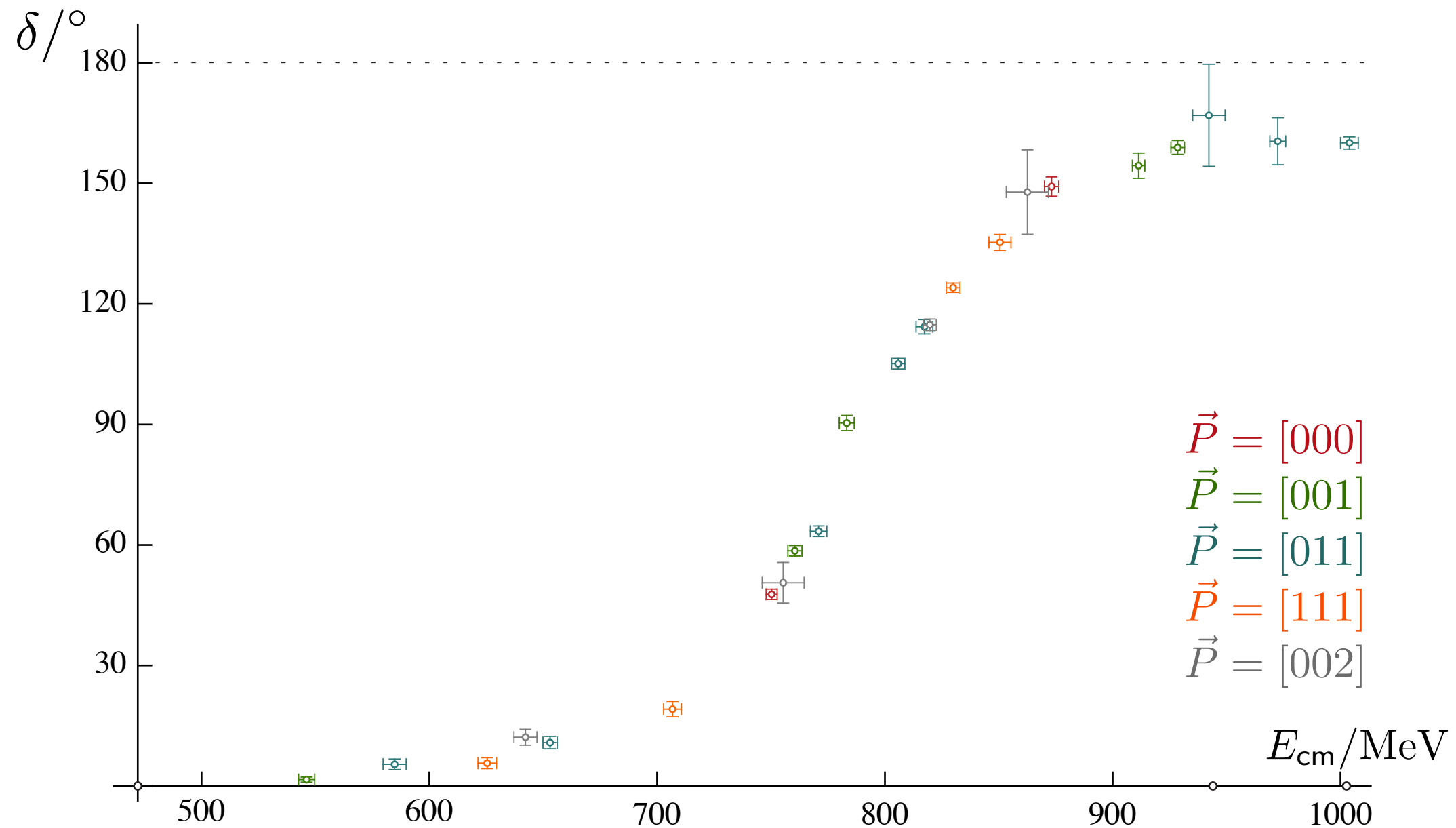
$$m_\pi = 236 \text{ MeV}$$

ρ resonance with moving frames



$$m_\pi = 236 \text{ MeV}$$

ρ resonance with moving frames



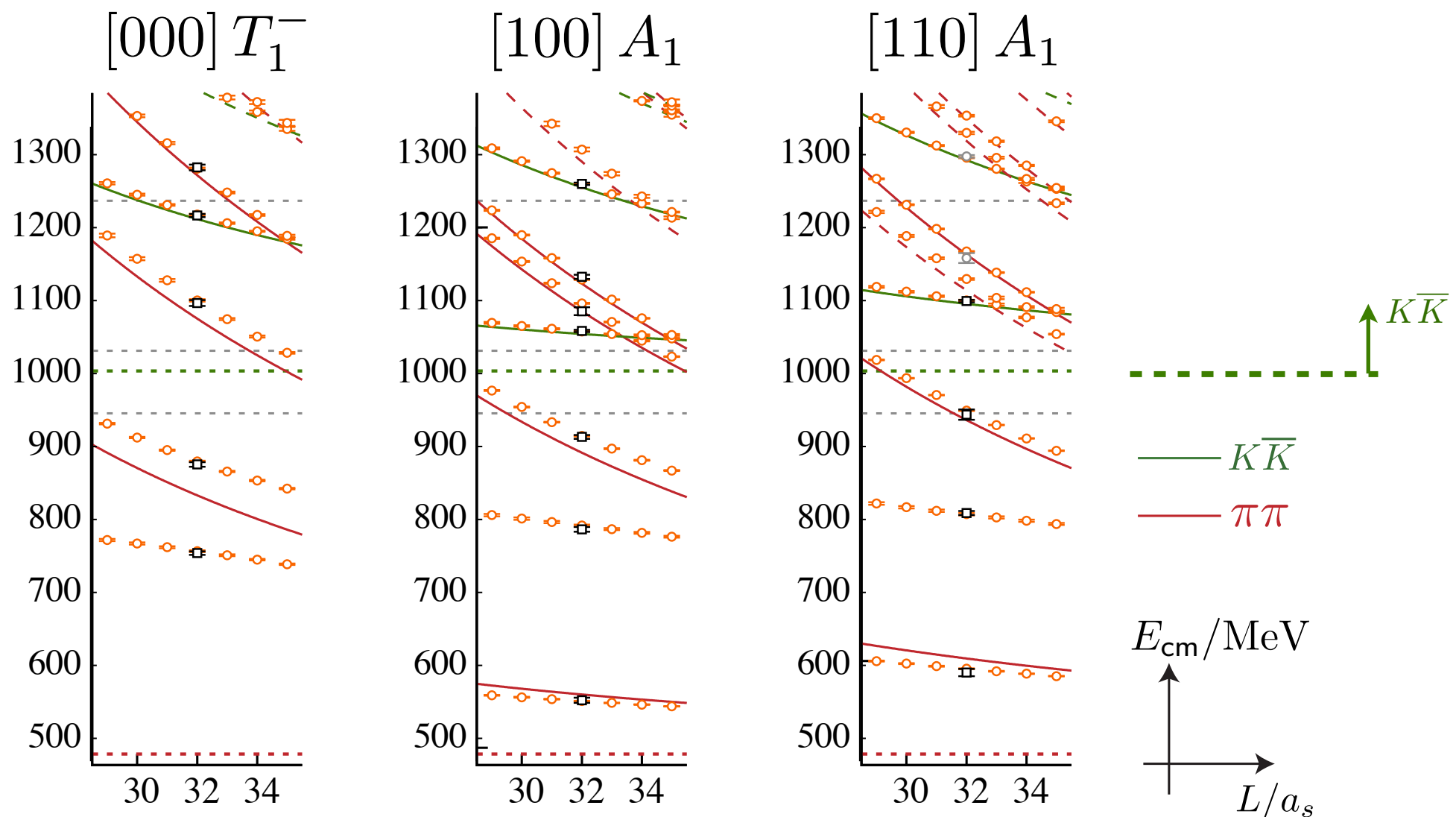
$$m_\pi = 236 \text{ MeV}$$

ρ resonance into the coupled-channel region

$$\mathbf{t} = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$

$$\mathbf{t}^{-1} = \mathbf{K}^{-1} - i\rho$$

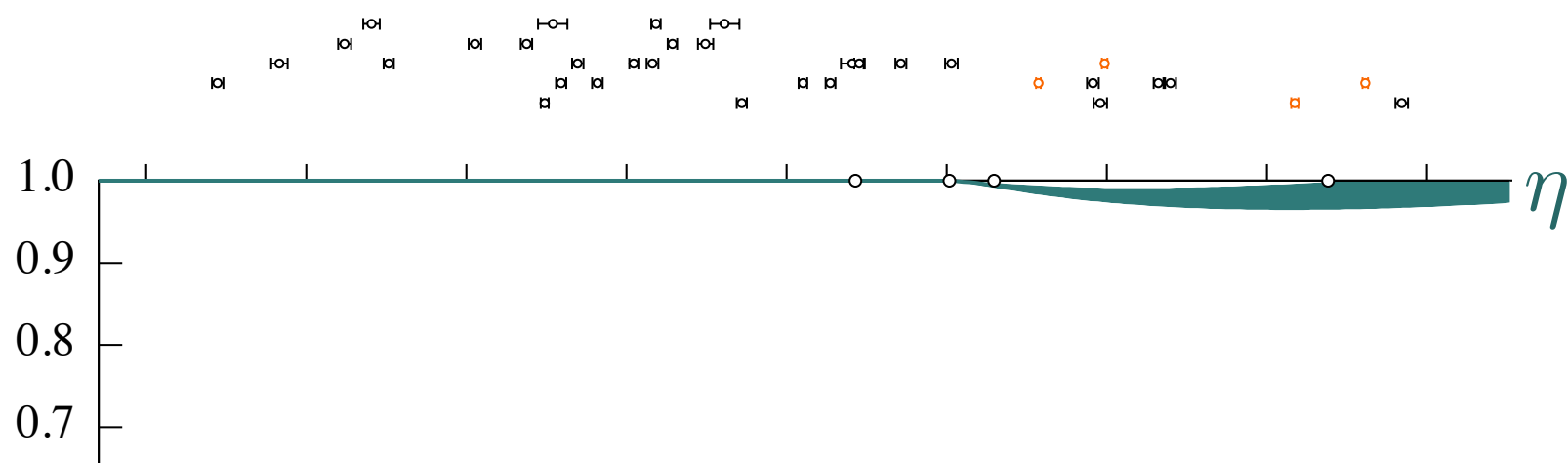
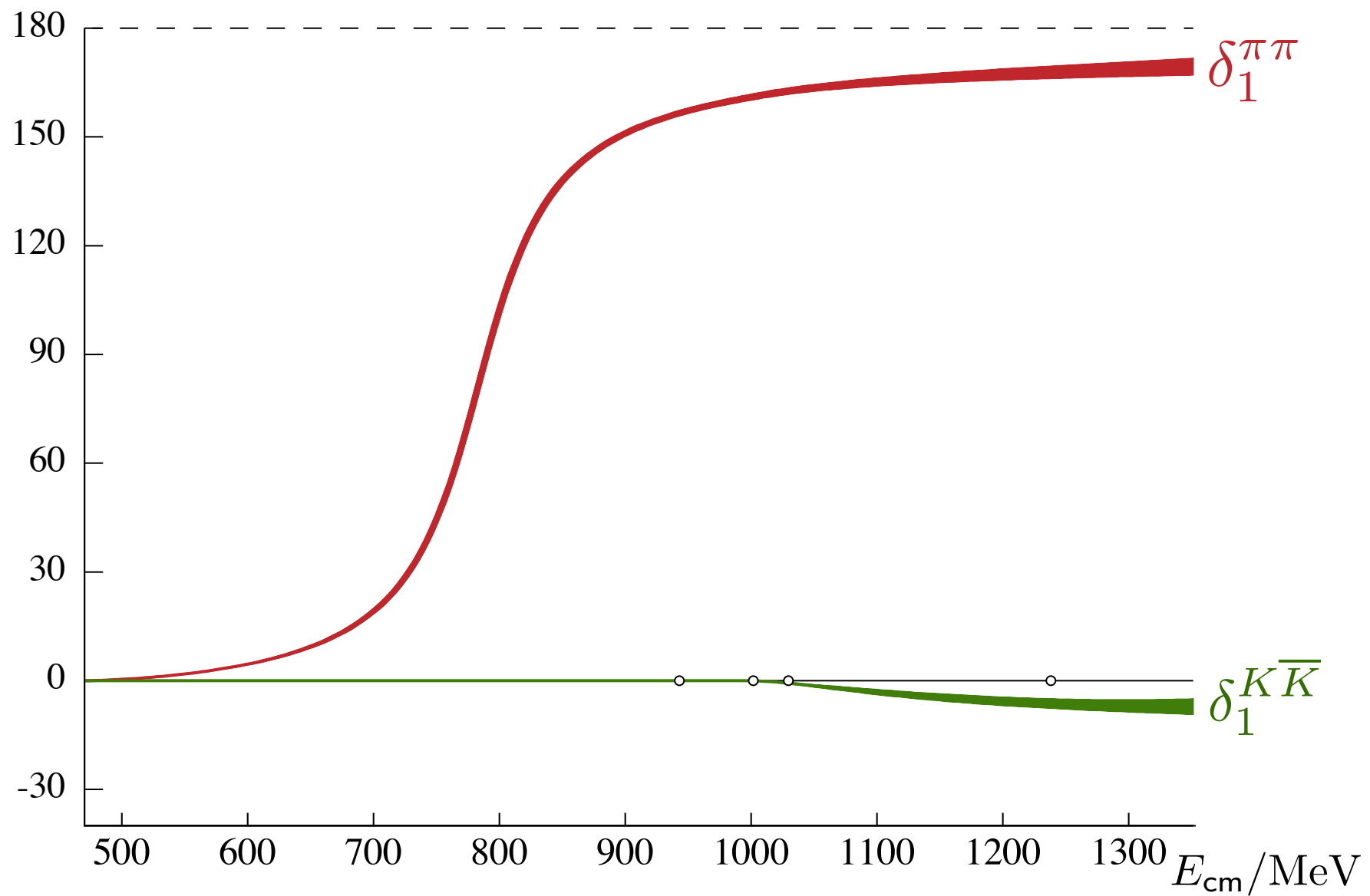
$$\text{e.g.: } K_{ij} = \frac{g_i g_j}{m^2 - s} + \gamma_{ij}$$



$$m_\pi = 236 \text{ MeV}$$

ρ resonance into the coupled-channel region

PRD 92 094502, arXiv:1507.02599



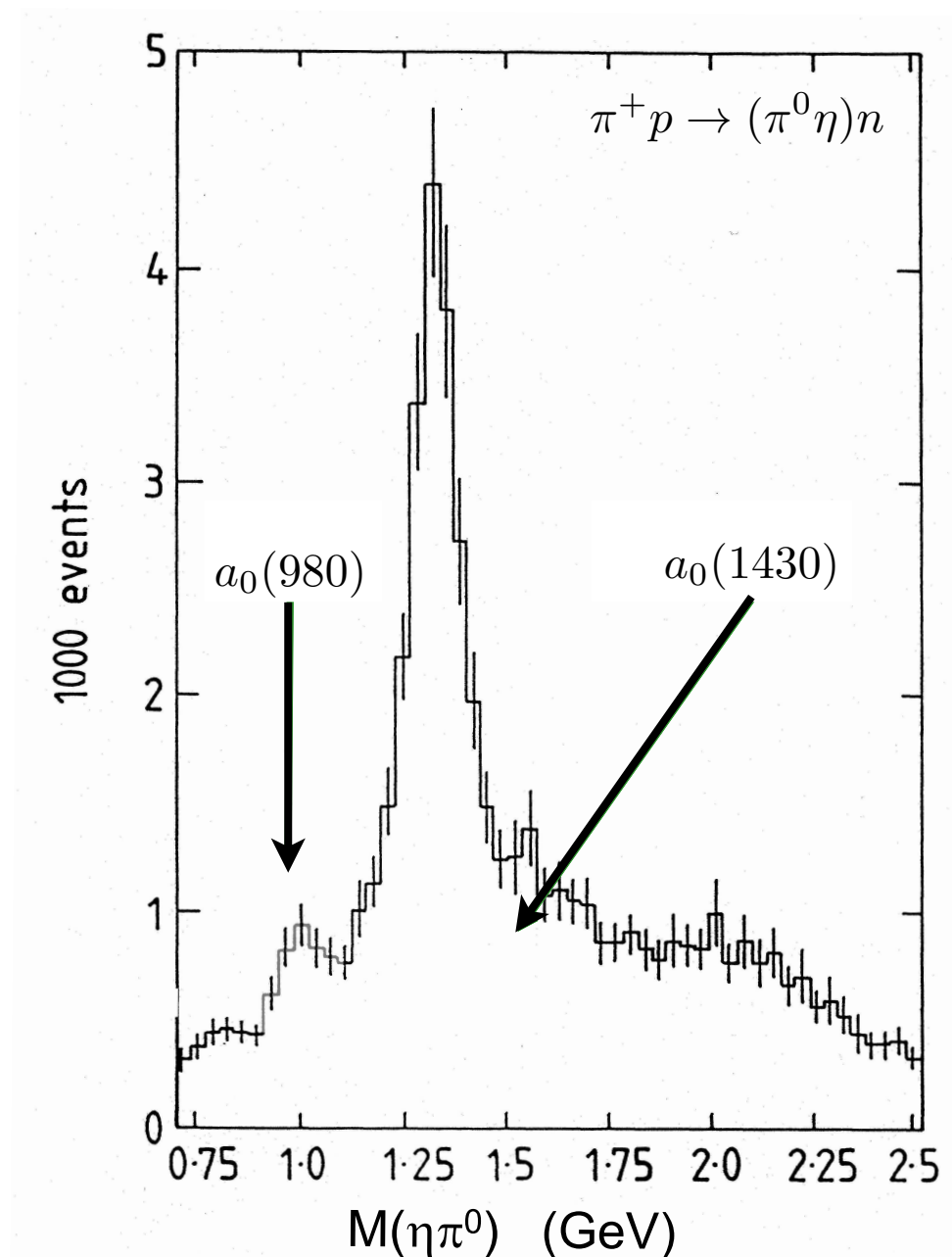
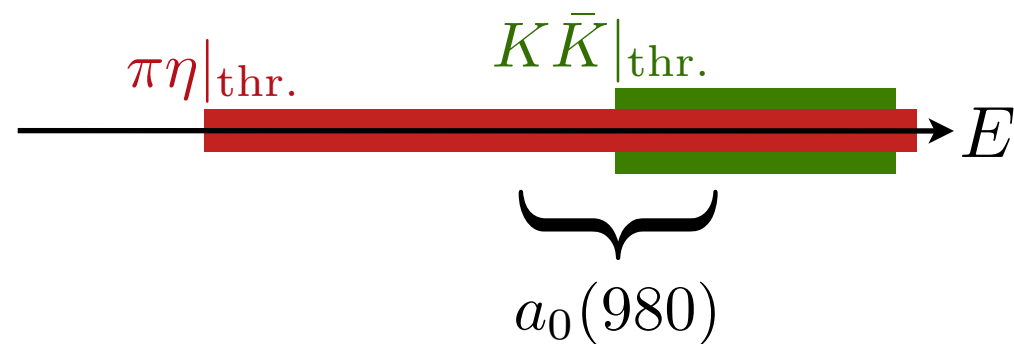
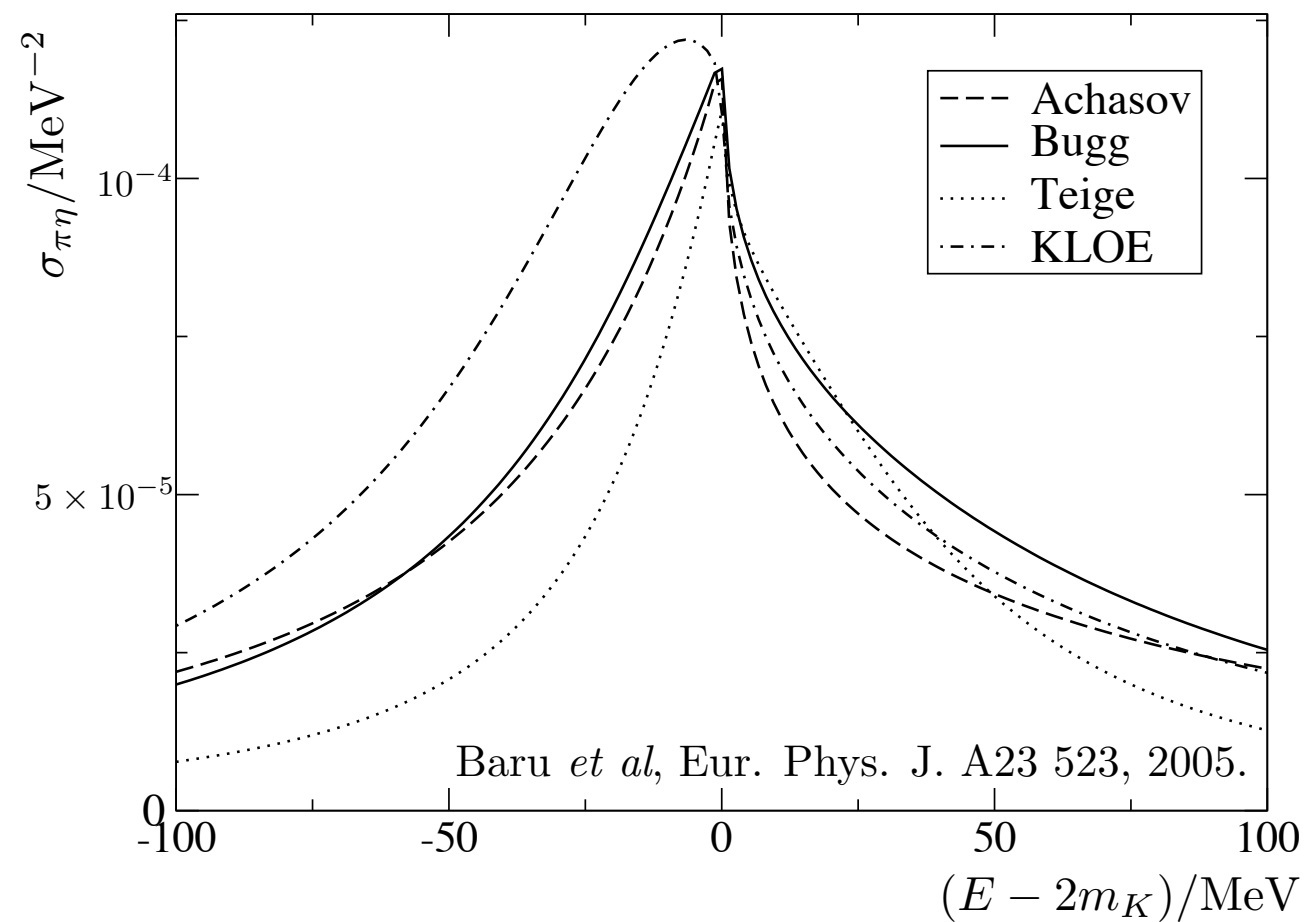
$$m_{\pi} = 236 \text{ MeV}$$

An a_0 resonance

$$\pi\eta - K\bar{K} - \pi\eta'$$

$$I = 1 \quad J = 0$$

PRD 93 094506, arXiv:1602.05122



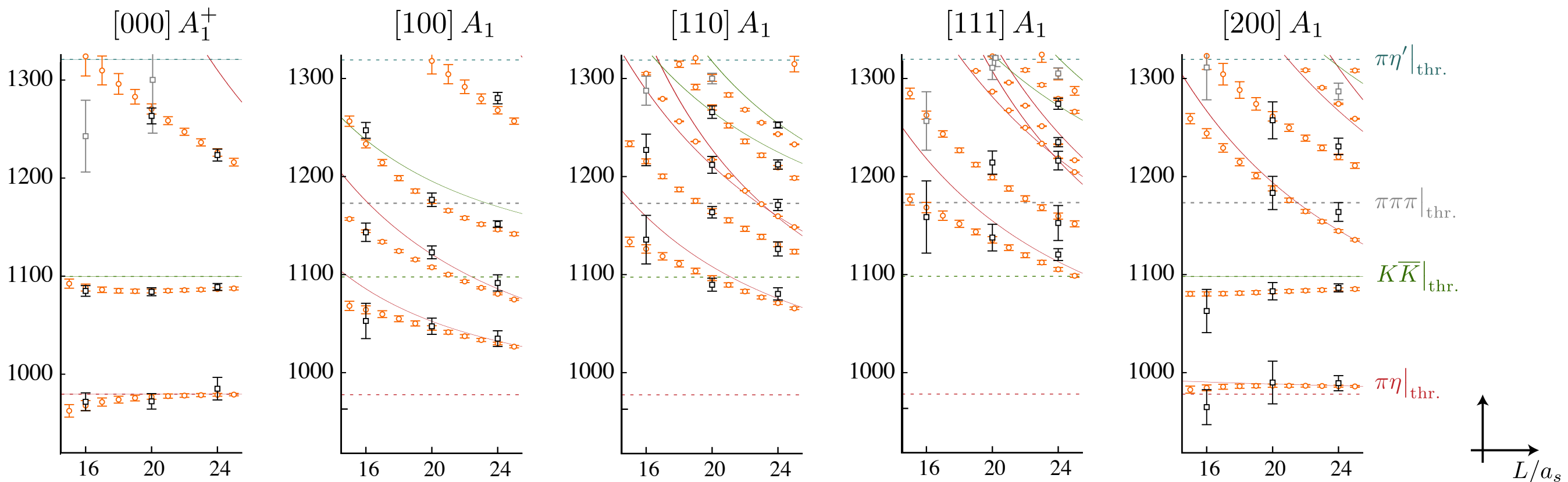
GAMS, Alde *et al* PLB 203 397, 1988.

$$m_\pi = 391 \text{ MeV}$$

a_0 resonance - two channel region

$$\pi\eta\text{-}K\bar{K}$$

$$K_{ij} = \frac{g_i g_j}{m^2 - s} + \gamma_{ij}$$



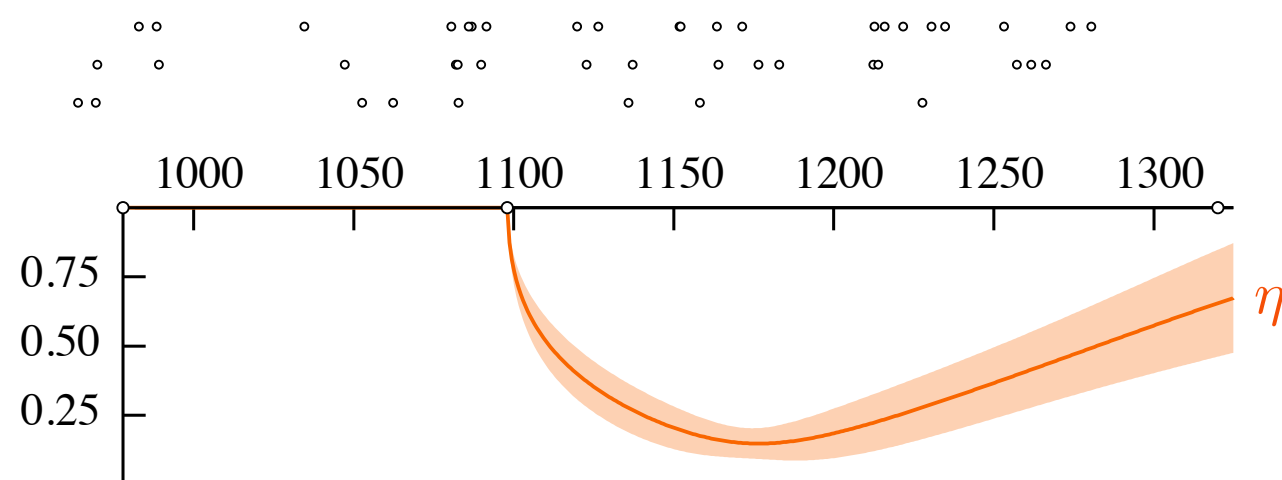
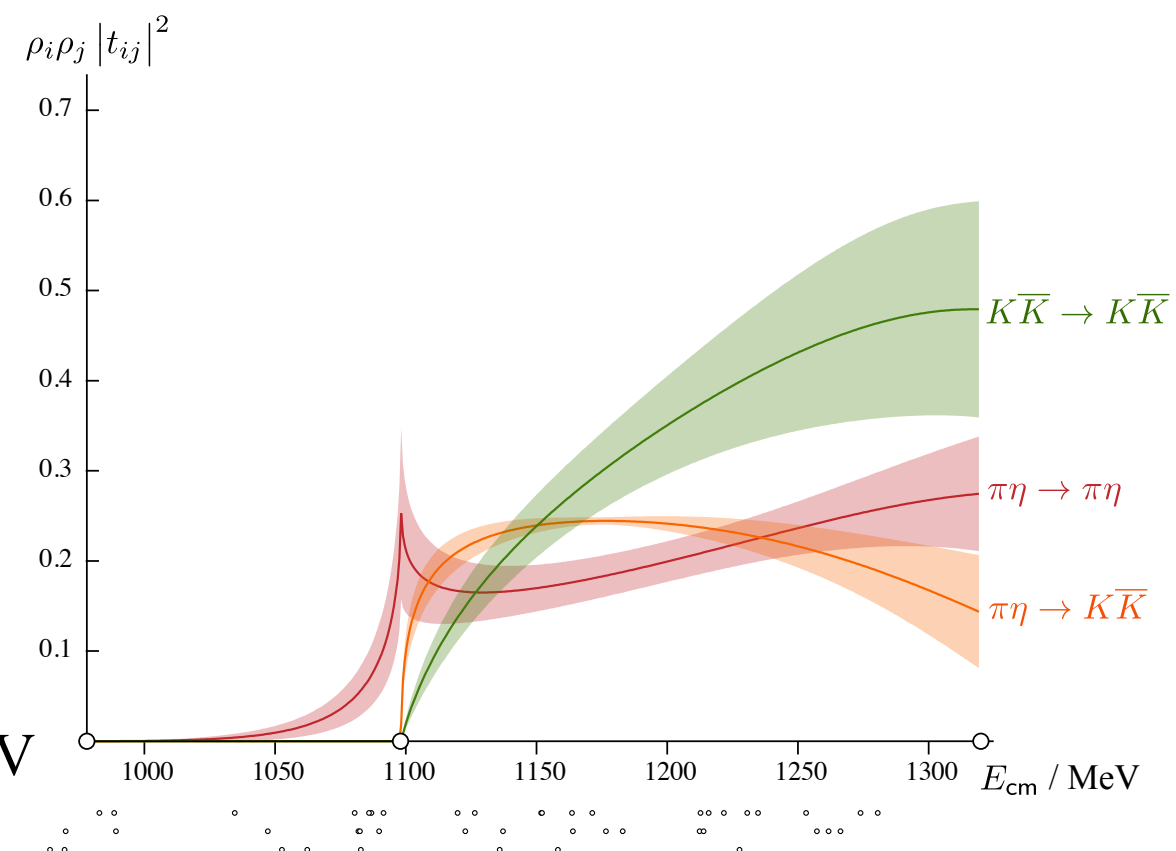
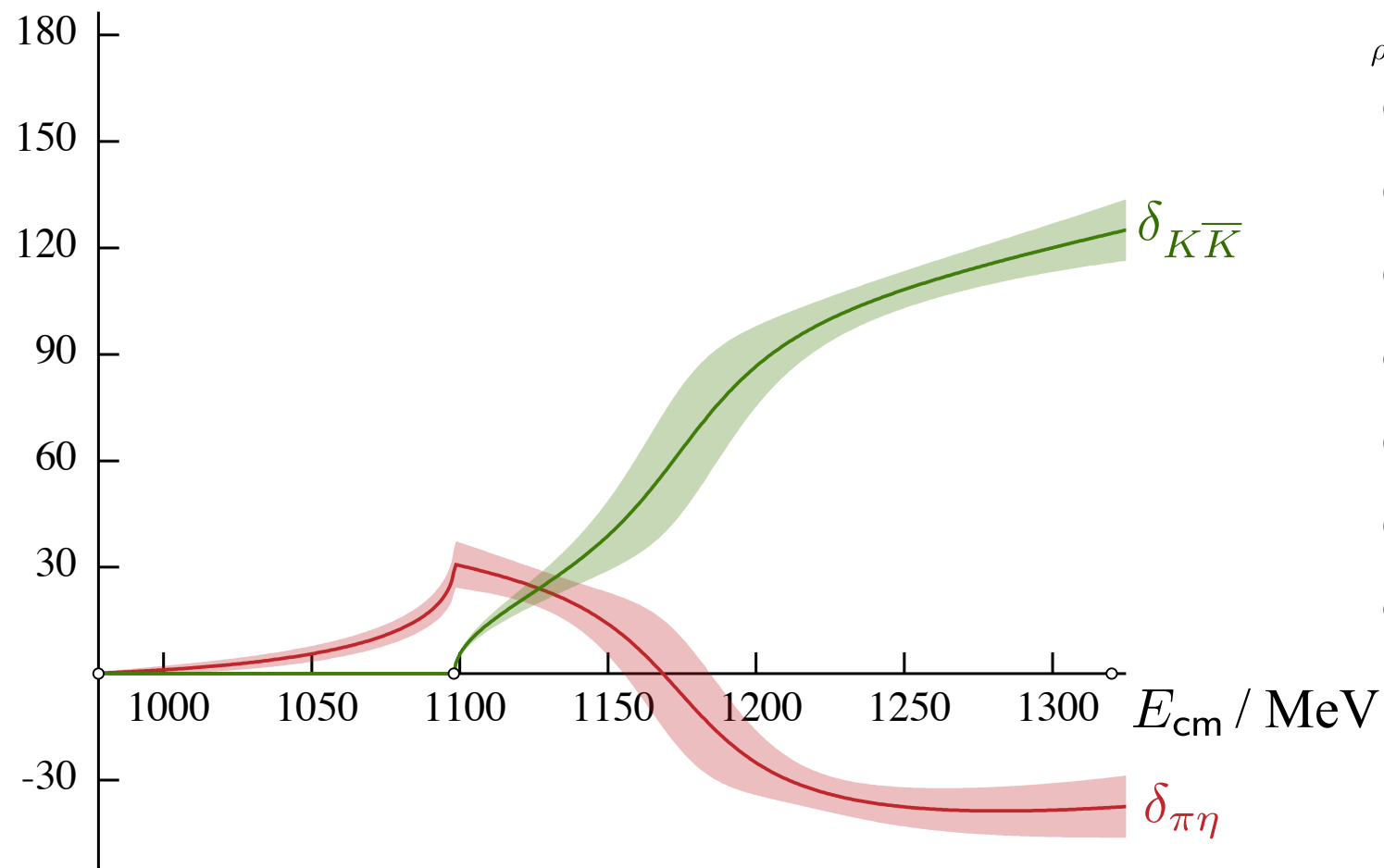
$$\begin{aligned} m &= (1254 \pm 16) \cdot \text{GeV} \\ g_{\pi\eta} &= (515 \pm 16) \cdot \text{GeV} \\ g_{K\bar{K}} &= (-730 \pm 85) \cdot \text{GeV} \\ \gamma_{\pi\eta, \pi\eta} &= -0.16 \pm 0.24 \\ \gamma_{\pi\eta, K\bar{K}} &= -0.56 \pm 0.29 \\ \gamma_{K\bar{K}, K\bar{K}} &= 0.12 \pm 0.38 \end{aligned} \begin{bmatrix} 1 & 0.58 & -0.06 & -0.51 & 0.39 & 0.02 \\ & 1 & -0.63 & -0.87 & 0.84 & -0.49 \\ & & 1 & 0.52 & -0.68 & 0.83 \\ & & & 1 & -0.90 & 0.53 \\ & & & & 1 & -0.78 \\ & & & & & 1 \end{bmatrix}$$

$$\chi^2/N_{\text{dof}} = \frac{58.0}{47-6} = 1.41$$

$$m_\pi = 391 \text{ MeV}$$

a_0 resonance - two channel region

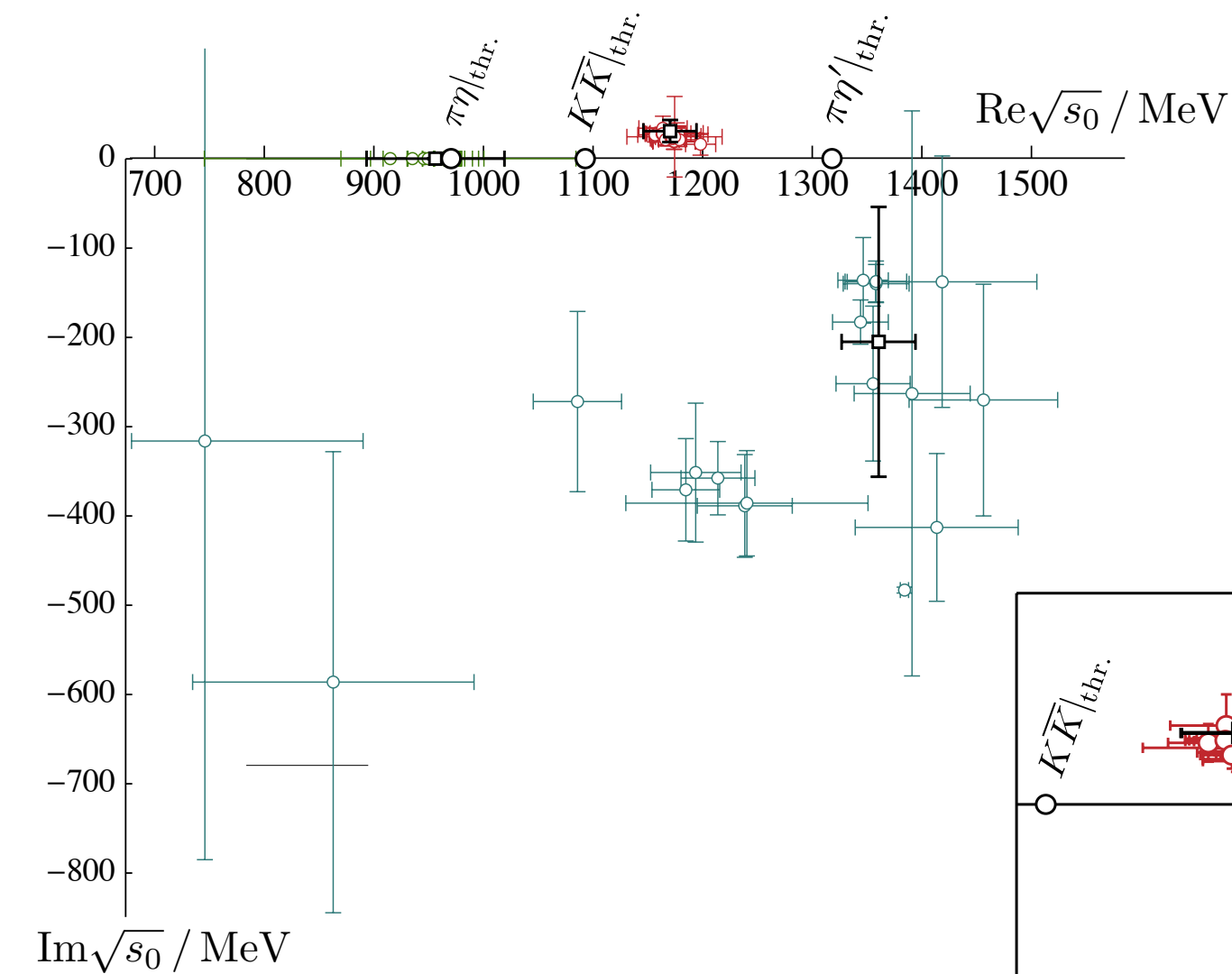
S-wave $\pi\eta$ - $K\bar{K}$



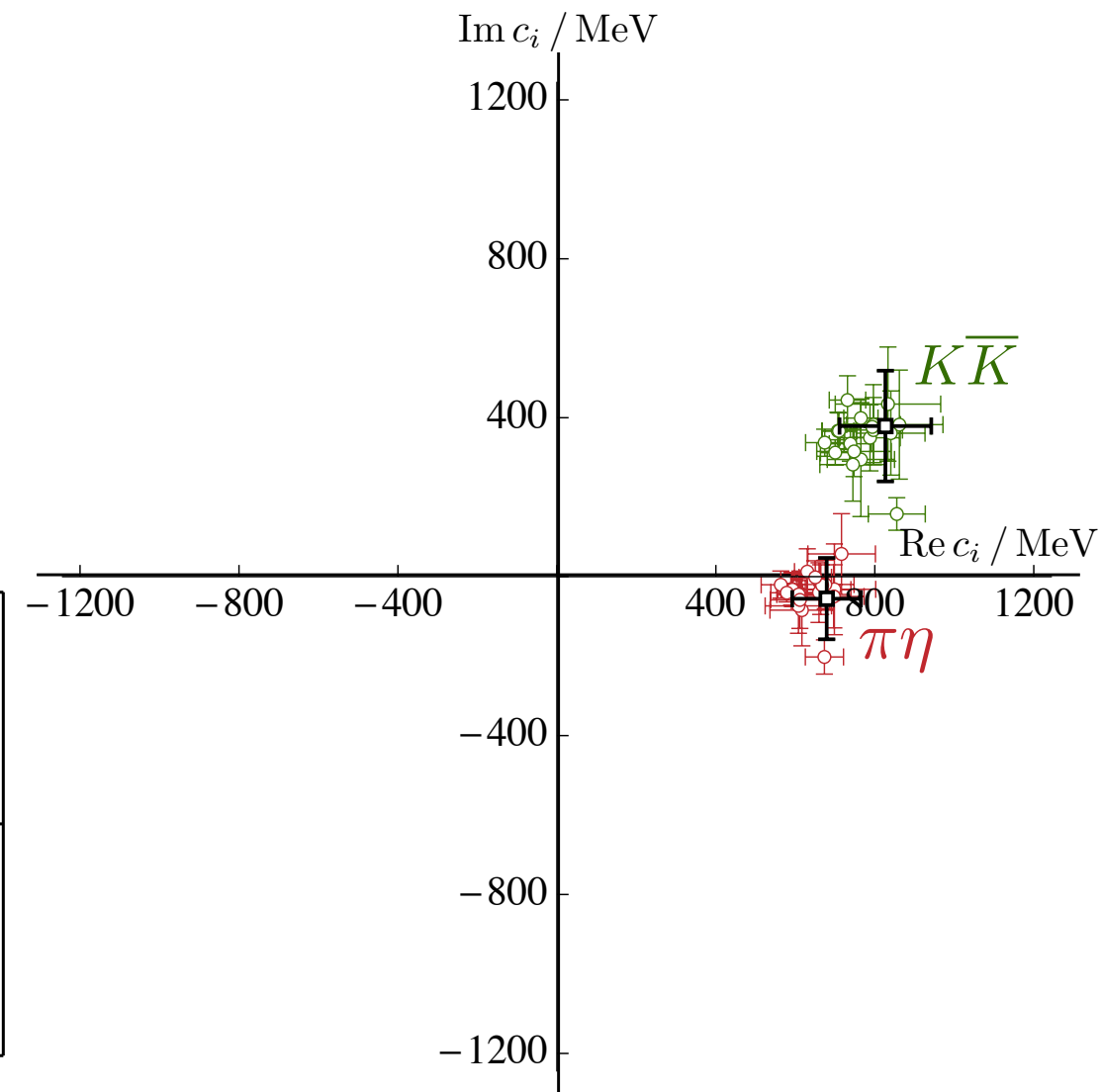
$$m_{\pi} = 391 \text{ MeV}$$

a_0 resonance pole

$$t_{ij} \sim \frac{c_i c_j}{s_0 - s}$$



$$\sqrt{s_0} = \left((1177 \pm 27) + \frac{i}{2}(49 \pm 33) \right) \text{ MeV}$$



$$|c_{\pi\eta}| = 652(130) \text{ MeV}$$

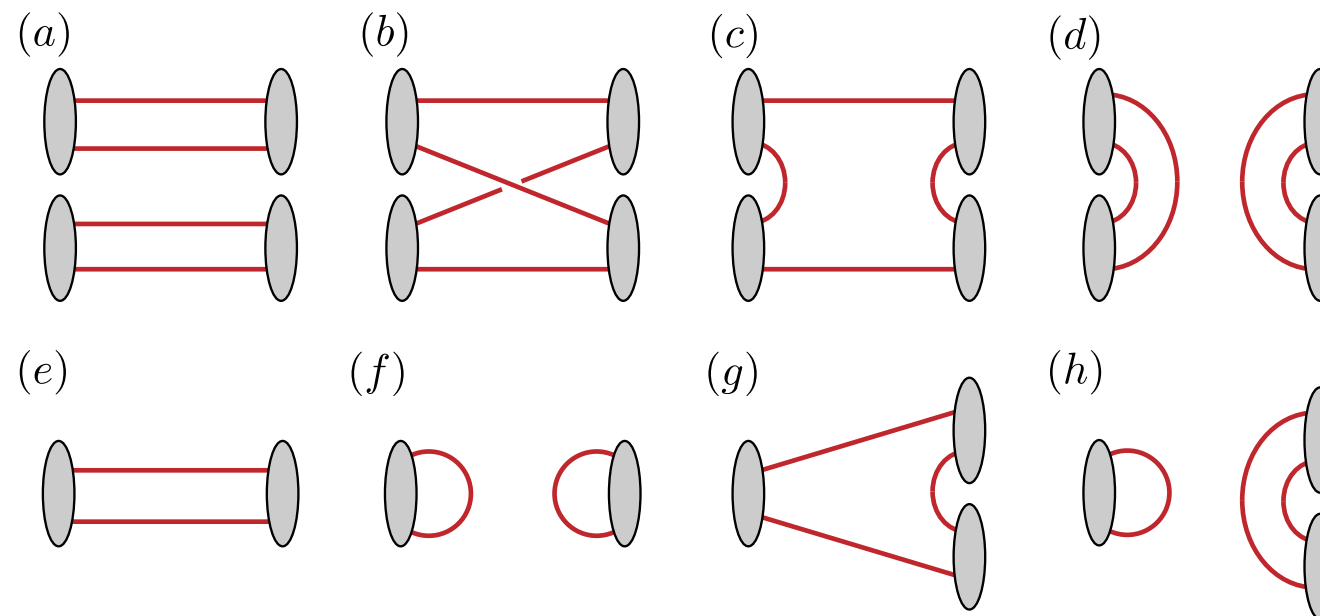
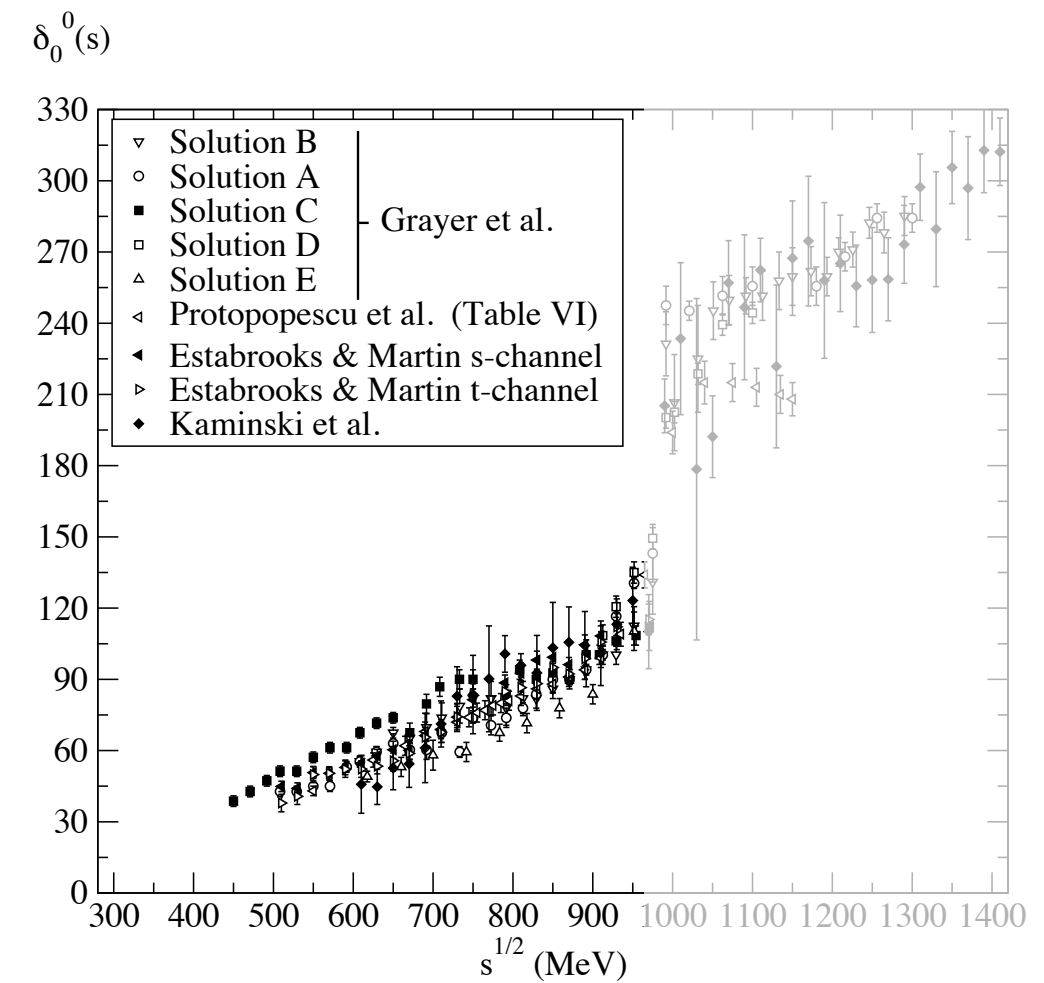
$$|c_{K\bar{K}}| = 844(170) \text{ MeV}$$

$$m_\pi = 391 \text{ MeV}$$

The $f_0(500)/\sigma$ resonance

arXiv:1607.05900

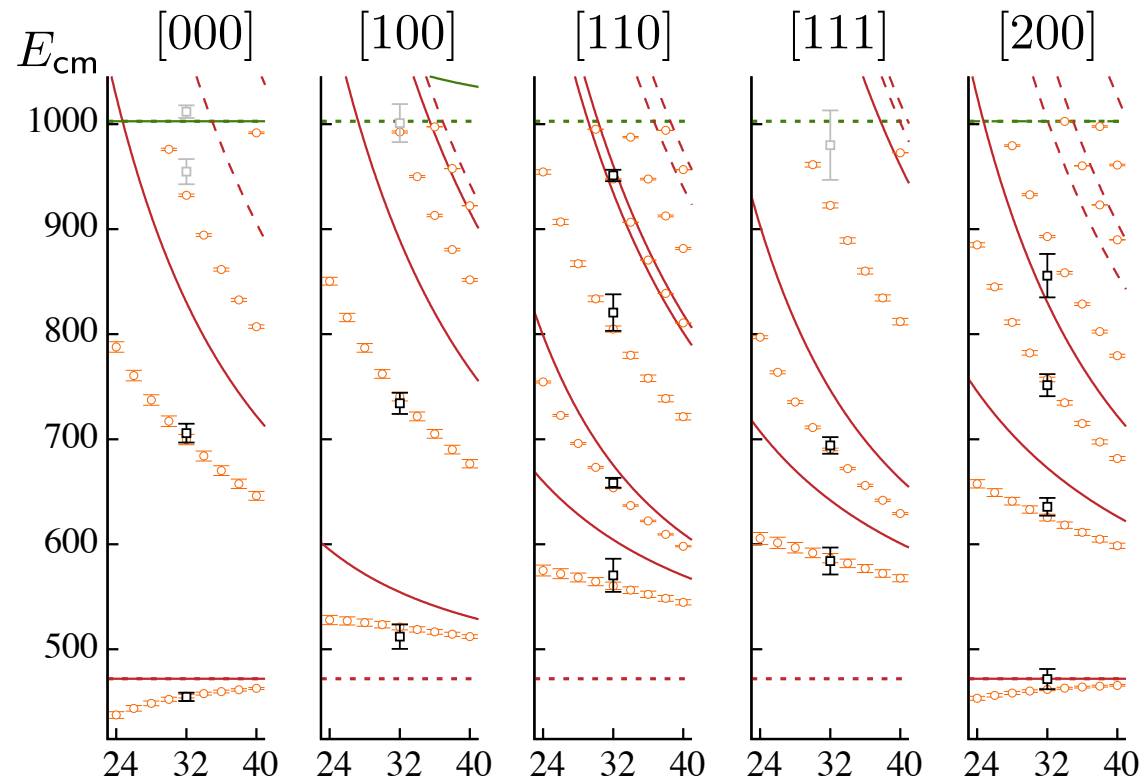
elastic scattering with
vacuum quantum numbers
 $\pi\pi$ in $I = 0, J = 0$



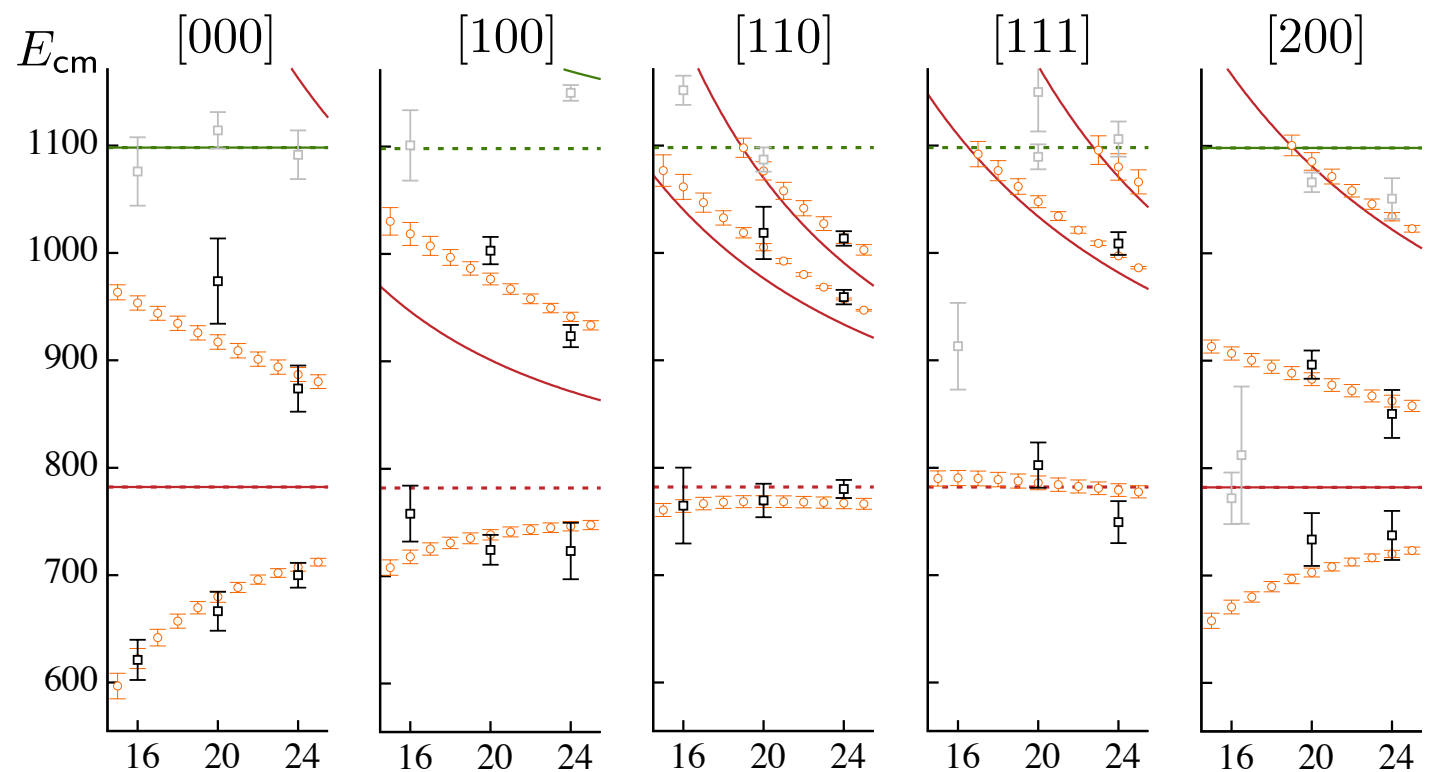
The $f_0(500)/\sigma$ resonance

elastic scattering with
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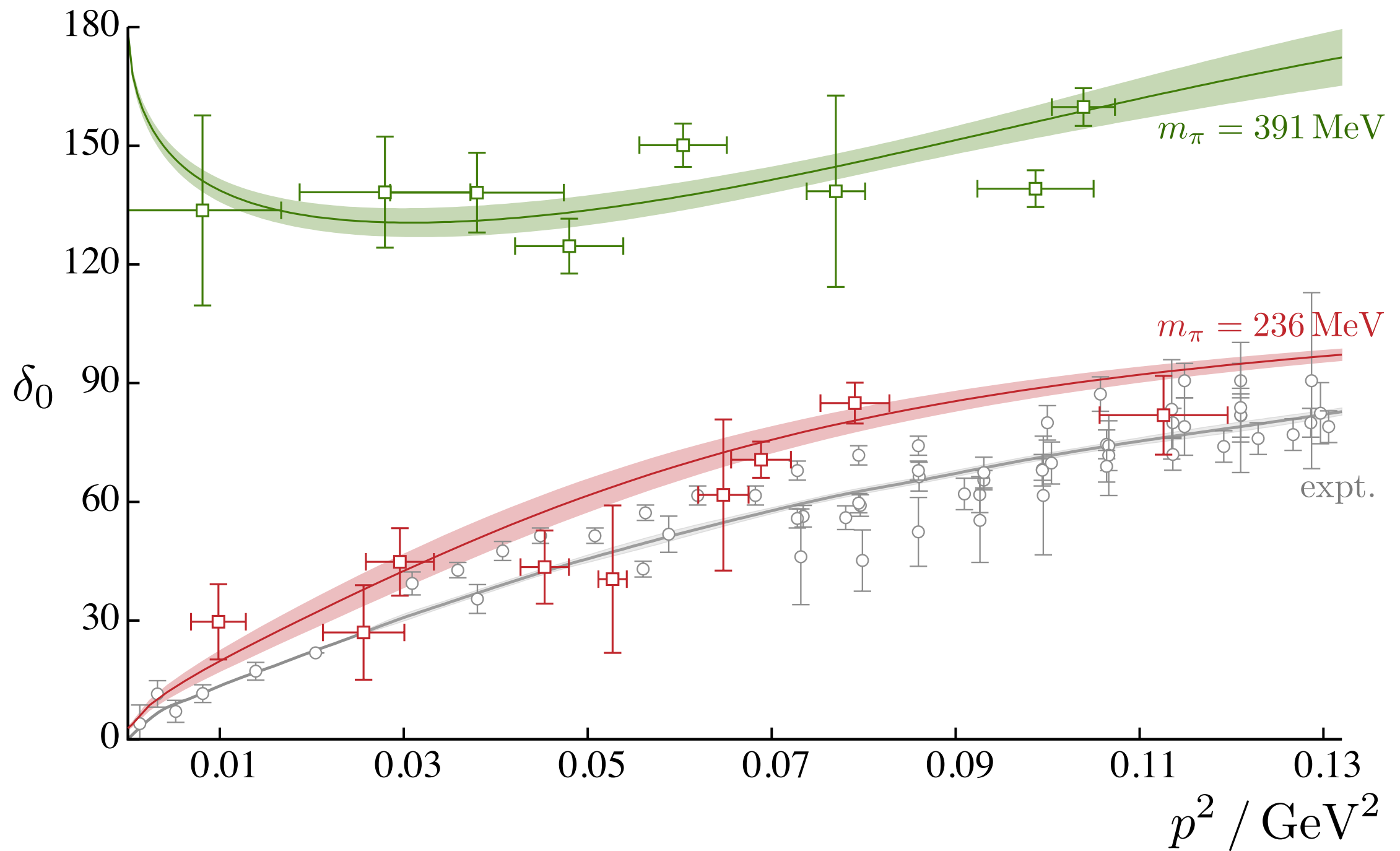
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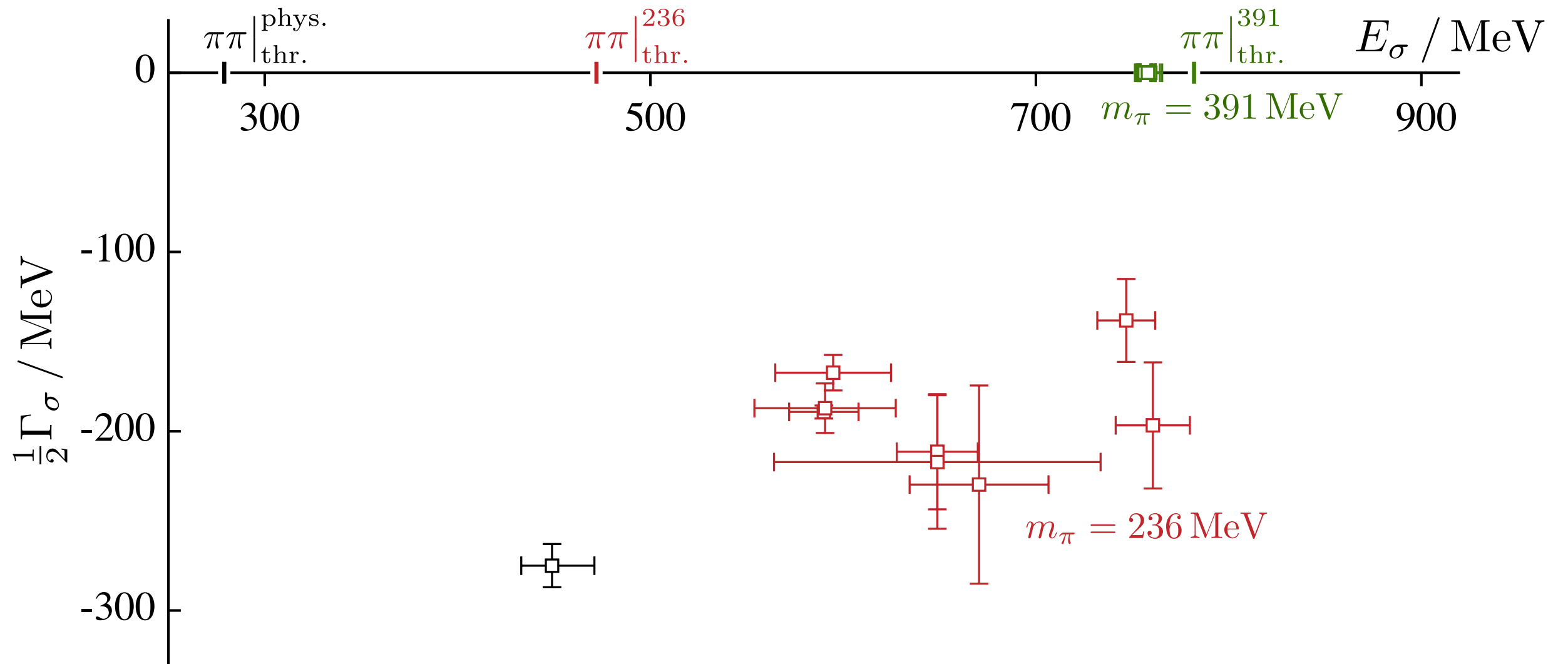


The $f_0(500)/\sigma$ resonance



The $f_0(500)/\sigma$ resonance

arXiv:1607.05900



Future directions

two-body coupled-channel

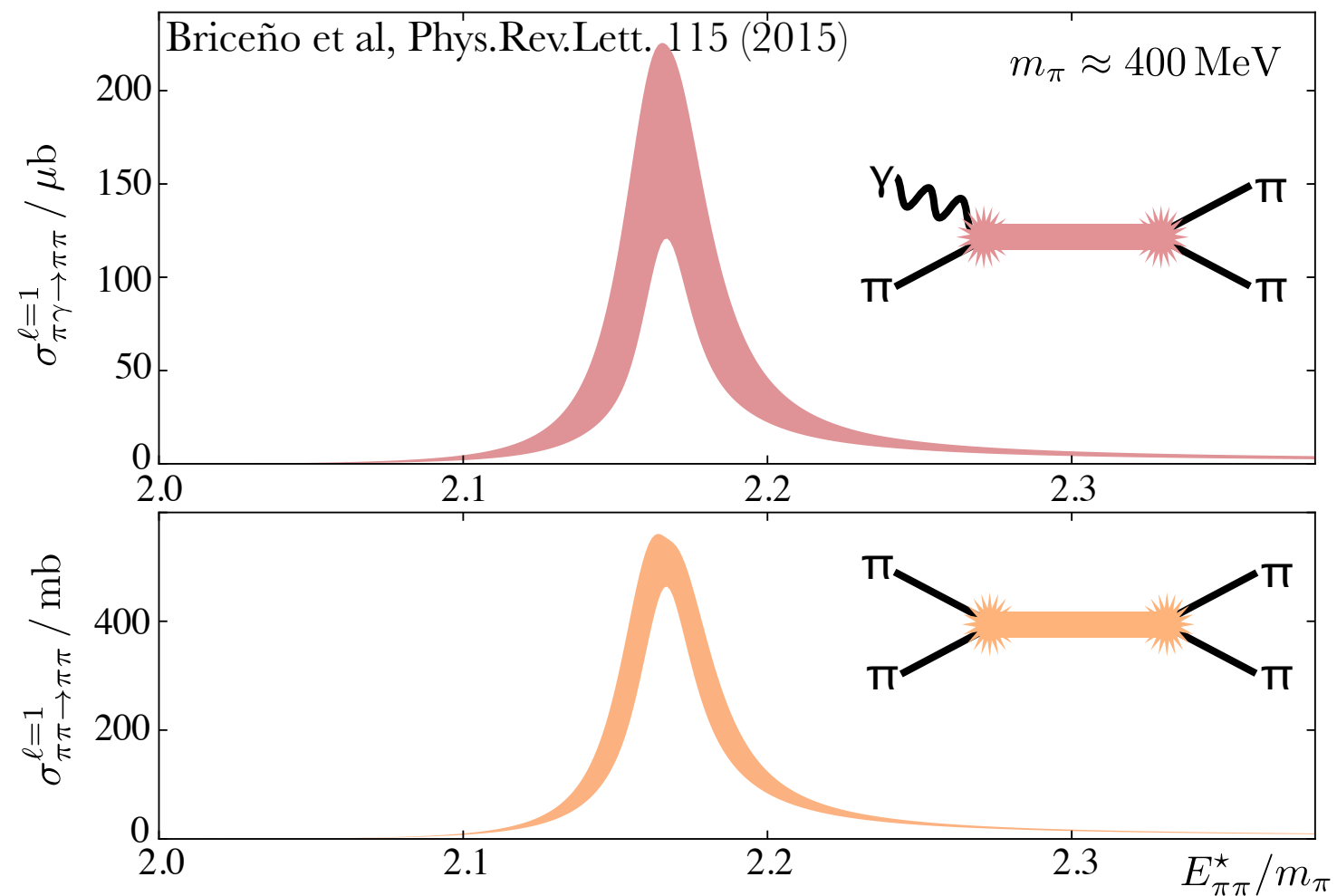
$f_0(980)$

$D\bar{D}$

$D\bar{D}^*$

$N\pi$

$\gamma a \rightarrow bc$



further operator structures - glueball, tetraquark, ...

formalism for three-body and beyond

- needed for higher energies
- needed to get closer to the physical mass

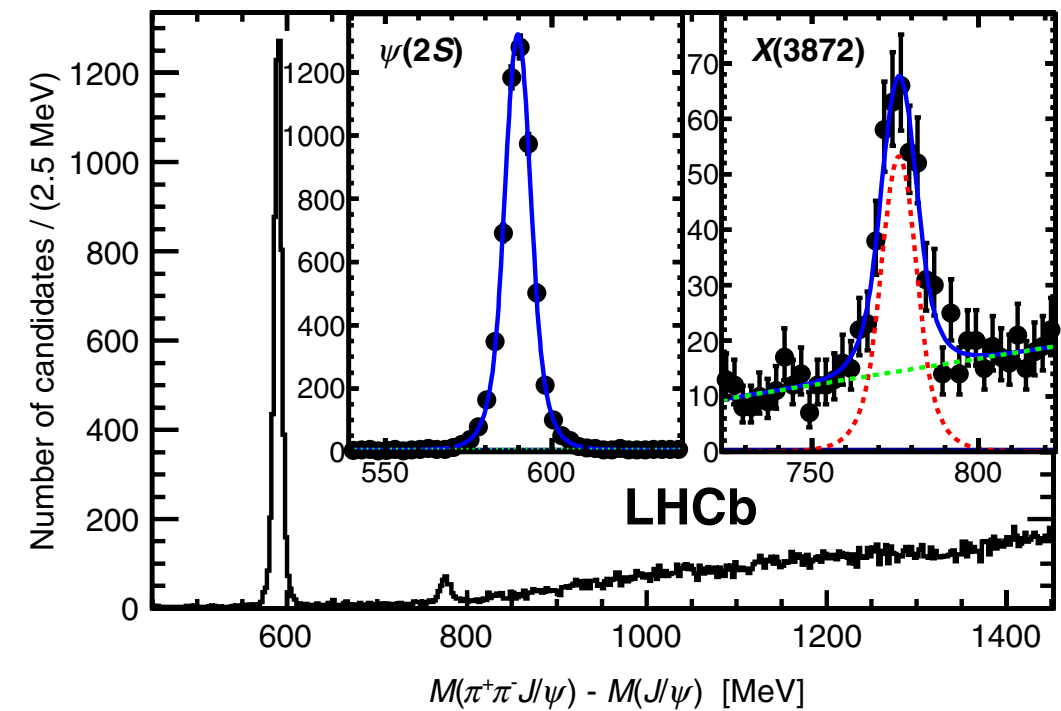
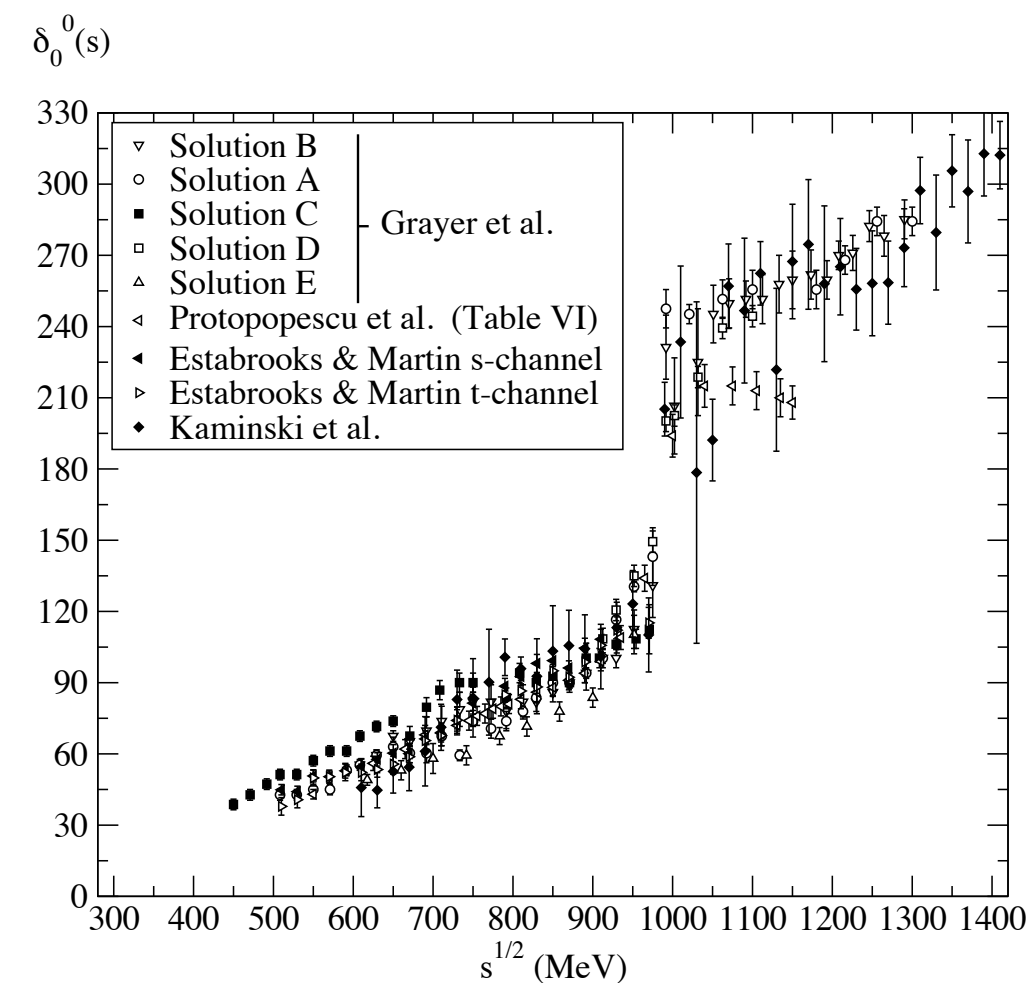
Backup

$$m_{\pi} = 391 \text{ MeV}$$

Coupled-channel scattering

$a_0(980)$, $f_0(980)$
 $a_1(1260)$
 $X(3872)$, and other XYZ states
 $N^*(1440)$, $\Lambda(1405)$, ...

all decay into multiple final states
 all are resonant enhancements in multiple channels
 to understand these rigorously, we need coupled-channel analyses



Coupled-channel scattering

Direct extension of the elastic quantization condition derived by Lüscher

$$\det [\mathbf{1} + i\boldsymbol{\rho}(E) \cdot \boldsymbol{t}(E) \cdot (\mathbf{1} + i\mathcal{M}(E, L))] = 0$$

phase space

infinite volume scattering
 t -matrix

known finite-volume
functions

Many derivations, **all in agreement**:

He, Feng, Liu 2005 - two channel QM, strong coupling

Hansen & Sharpe 2012 - field theory, multiple two-body channels

Briceño & Davoudi 2012 - strongly-coupled Bethe-Salpeter amplitudes

Guo et al 2012 - Hamiltonian & Lippmann-Schwinger

Also derivations in specific channels, or for a specific parameterization of the interactions like NREFT, chiral PT, Finite Volume Hamiltonian, etc.

Briceño 2014 - Generalised to scattering of particles with non-zero spin, and spin- $1/2$.

Significant steps towards a general 3-body quantization condition have been made

Amplitude parameterization

$$\mathbf{t} = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$

$$\det [\mathbf{1} + i\boldsymbol{\rho}(E) \cdot \mathbf{t}(E) \cdot (\mathbf{1} + i\boldsymbol{\mathcal{M}}(E, L))] = 0$$

determinant condition:

- several unknowns at each value of energy
- energy levels typically do not coincide
- underconstrained problem for a single energy

one solution: use energy dependent parameterizations

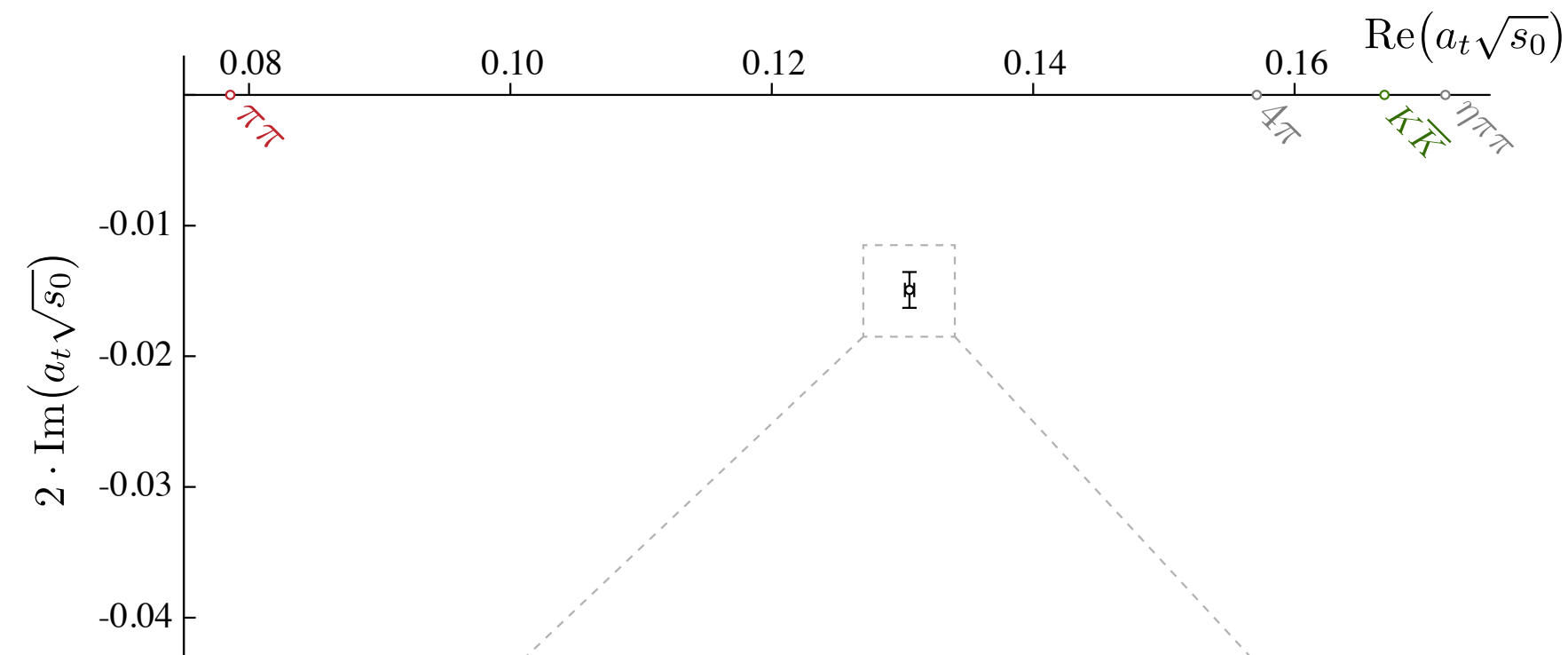
- Constrained problem when $\#(\text{energy levels}) > \#(\text{parameters})$
- Essential amplitudes respect unitarity of the S-matrix

$$\mathbf{S}^\dagger \mathbf{S} = \mathbf{1} \quad \rightarrow \quad \text{Im } \mathbf{t}^{-1} = -\boldsymbol{\rho} \quad \rho_{ij} = \delta_{ij} \frac{2k_i}{E_{\text{cm}}}$$

K-matrix approach:

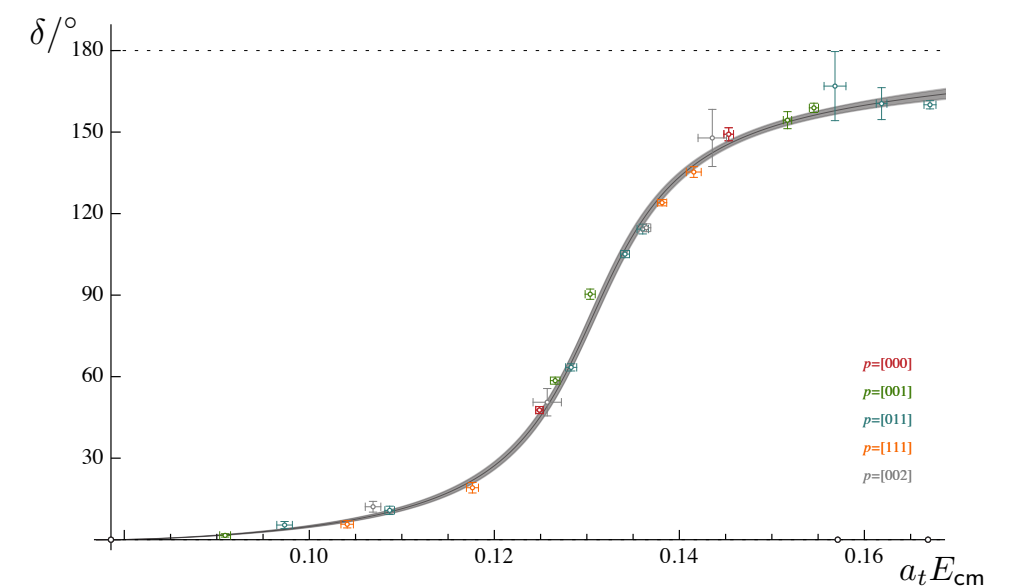
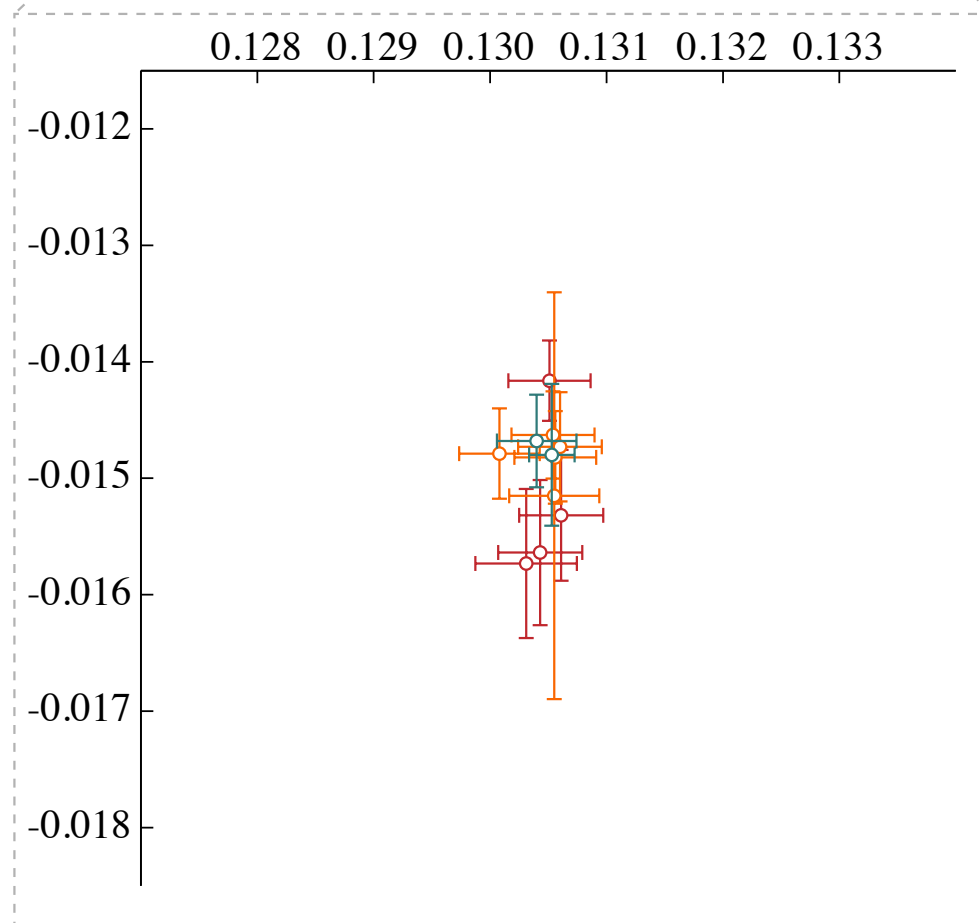
$$\mathbf{t}^{-1} = \mathbf{K}^{-1} - i\boldsymbol{\rho} \quad \text{e.g.: } K_{ij} = \frac{g_i g_j}{m^2 - s} + \gamma_{ij}$$

ρ resonance pole



near a pole:

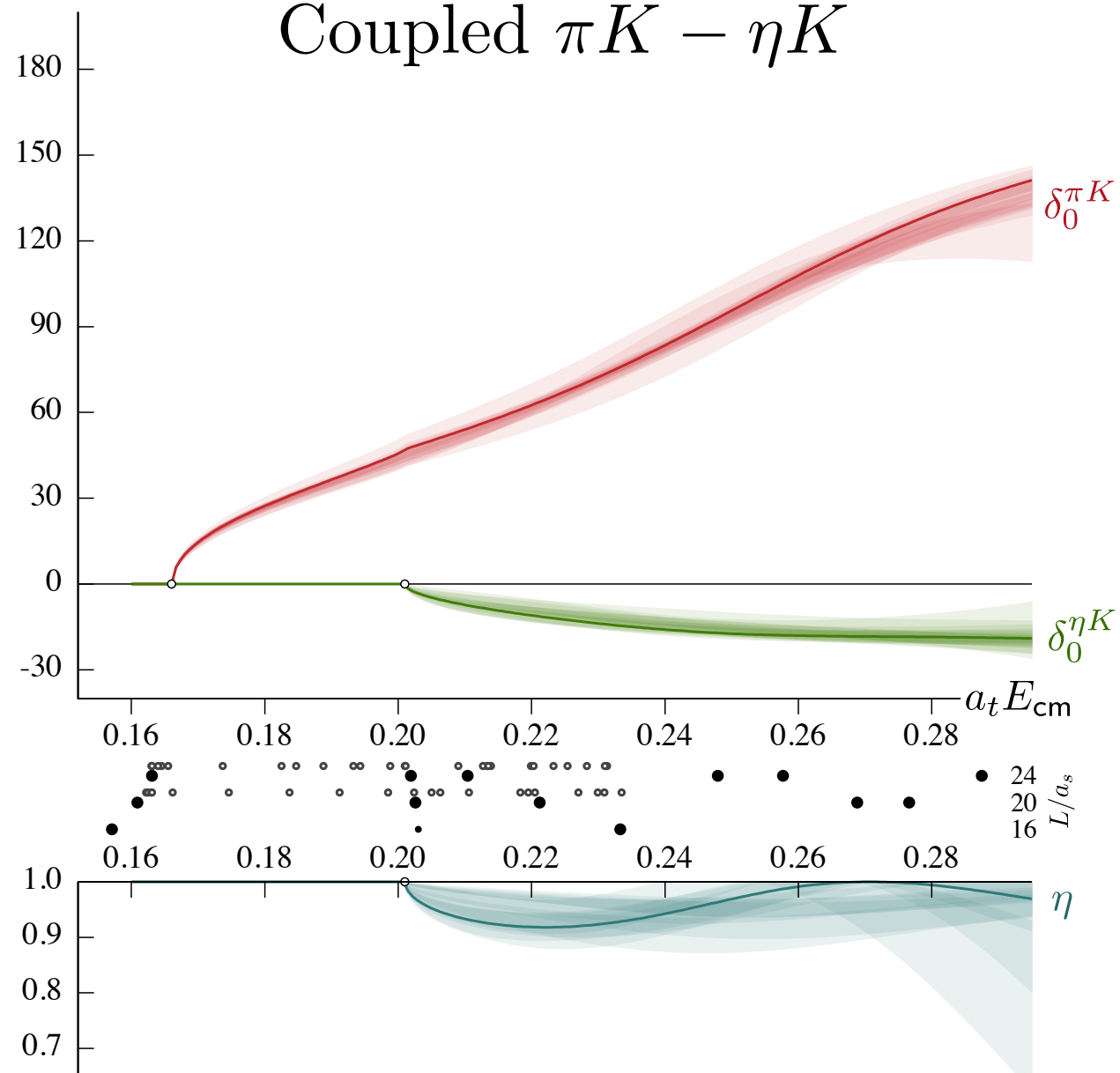
$$t_{ij} \sim \frac{c_i c_j}{s_0 - s}$$



$$m_\pi = 236 \text{ MeV}$$

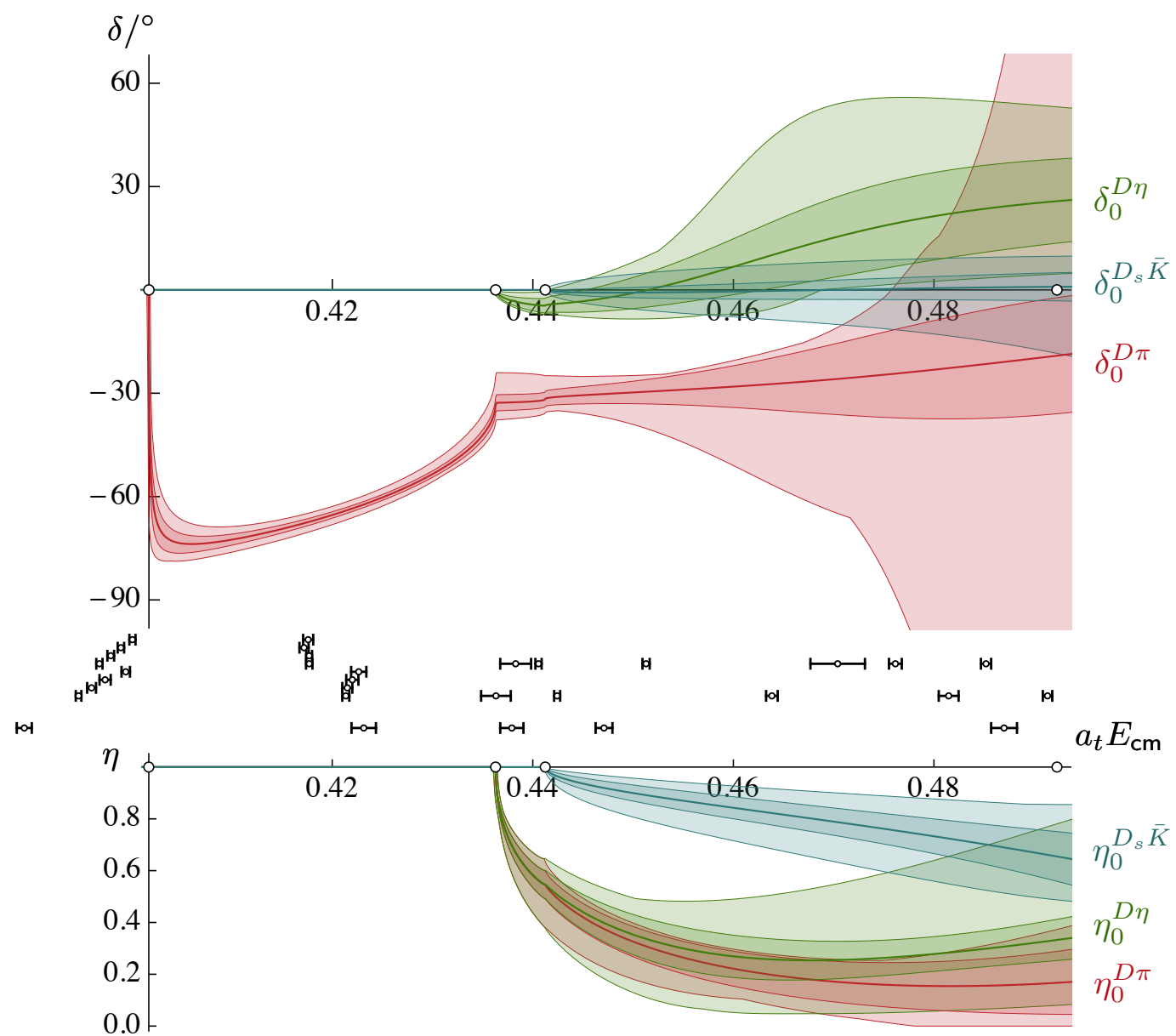
Other calculations

Coupled $\pi K - \eta K$



Combined S & P-wave analysis
80 energy levels from 3 volumes
arXiv:1406.4158, PRL 113 (2014) no.18, 182001

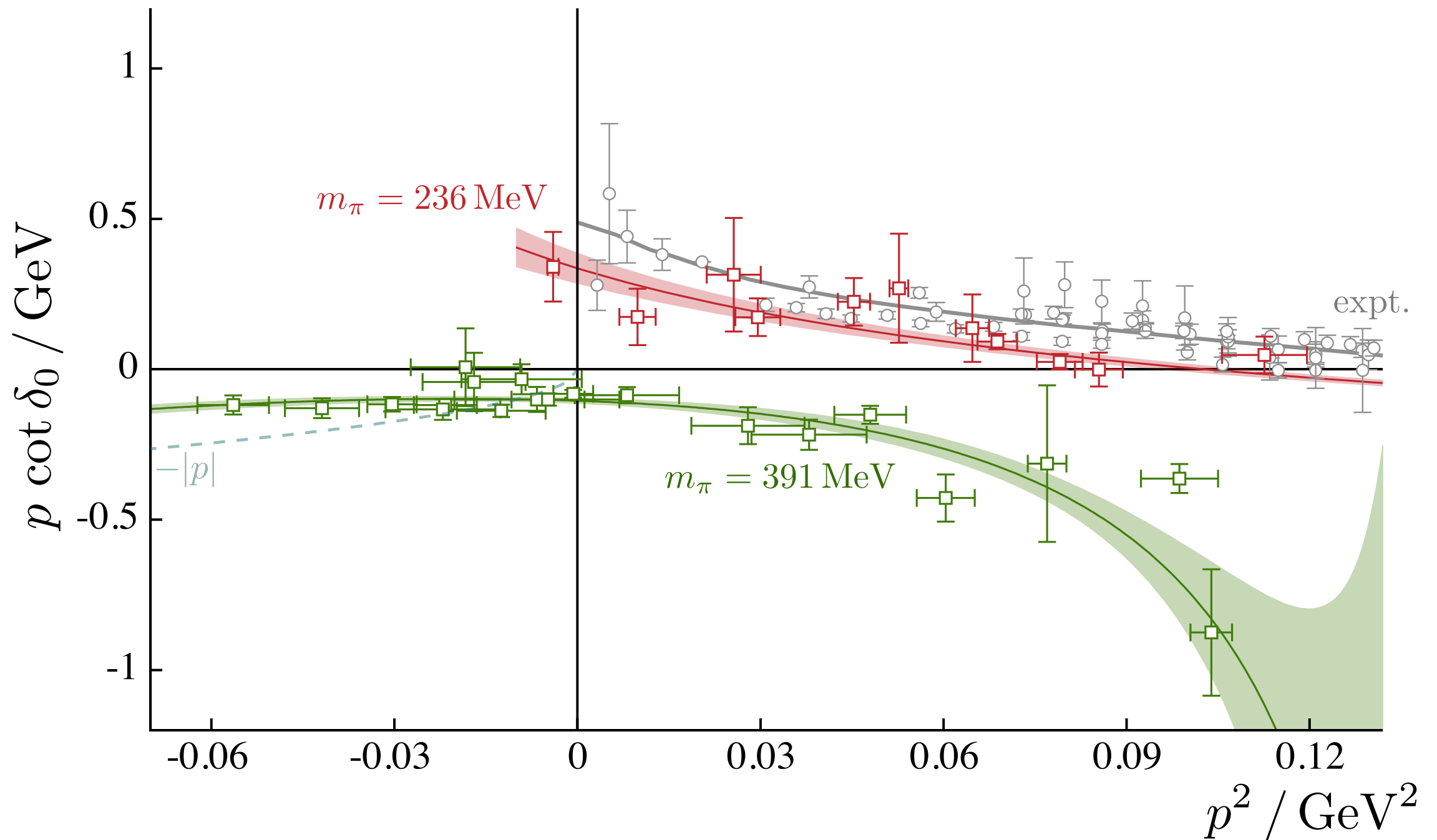
Coupled $D\pi - D\eta - D_s \bar{K}$



Combined S & P-wave analysis
3 coupled channels in S-wave
47 energy levels from 3 volumes
arXiv:1607.07093

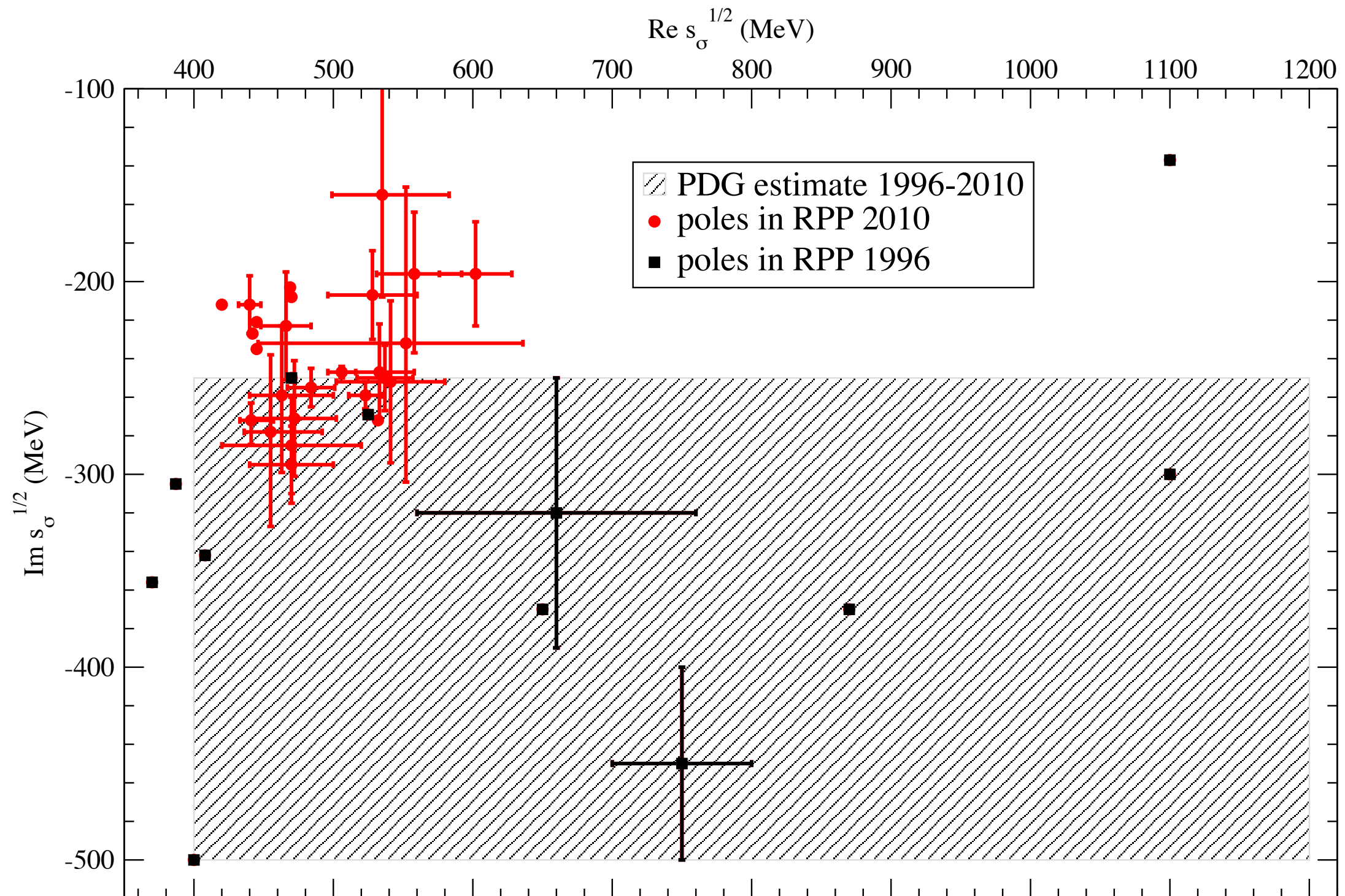
$$m_\pi = 391 \text{ MeV}$$

The $f_0(500)/\sigma$ resonance



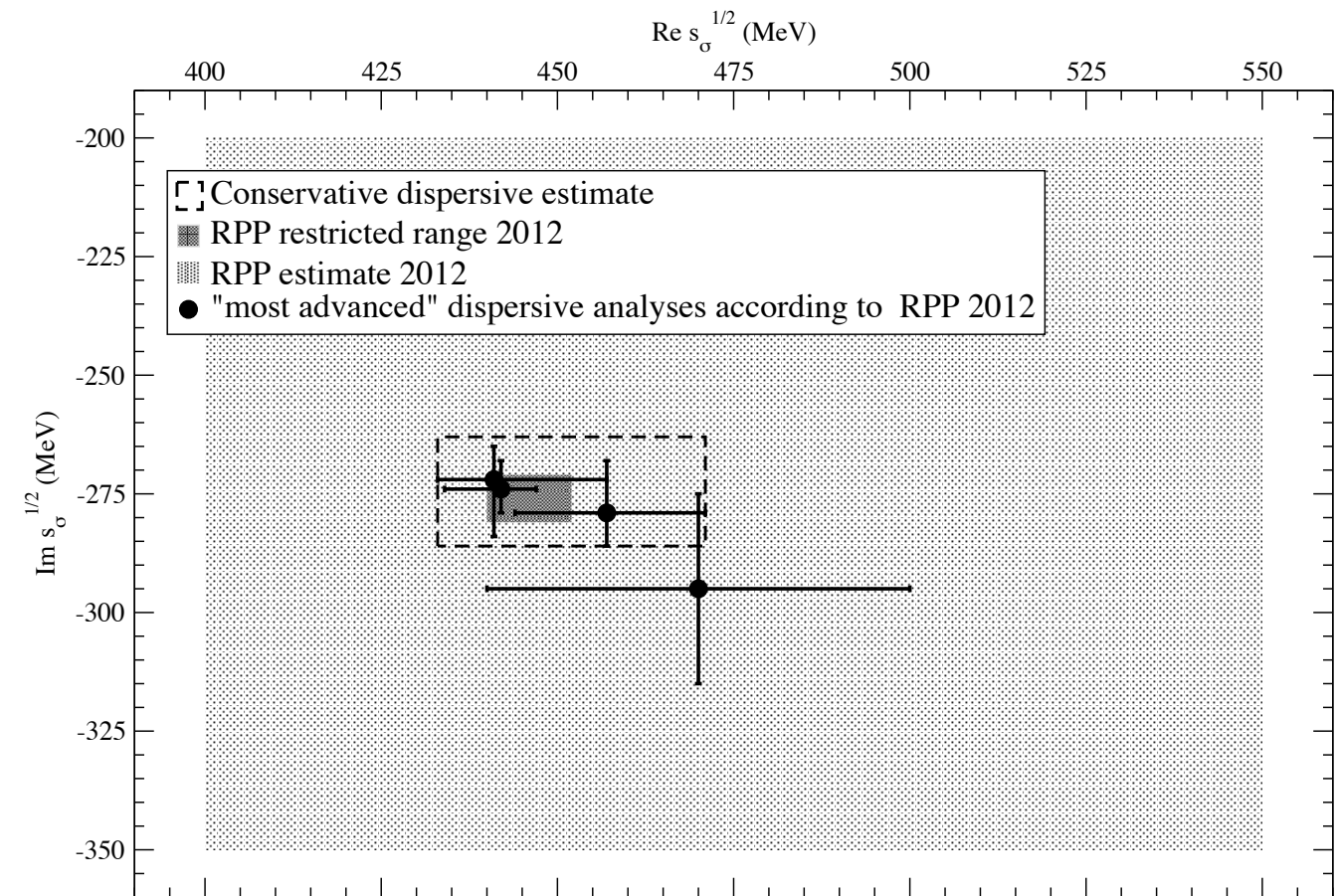
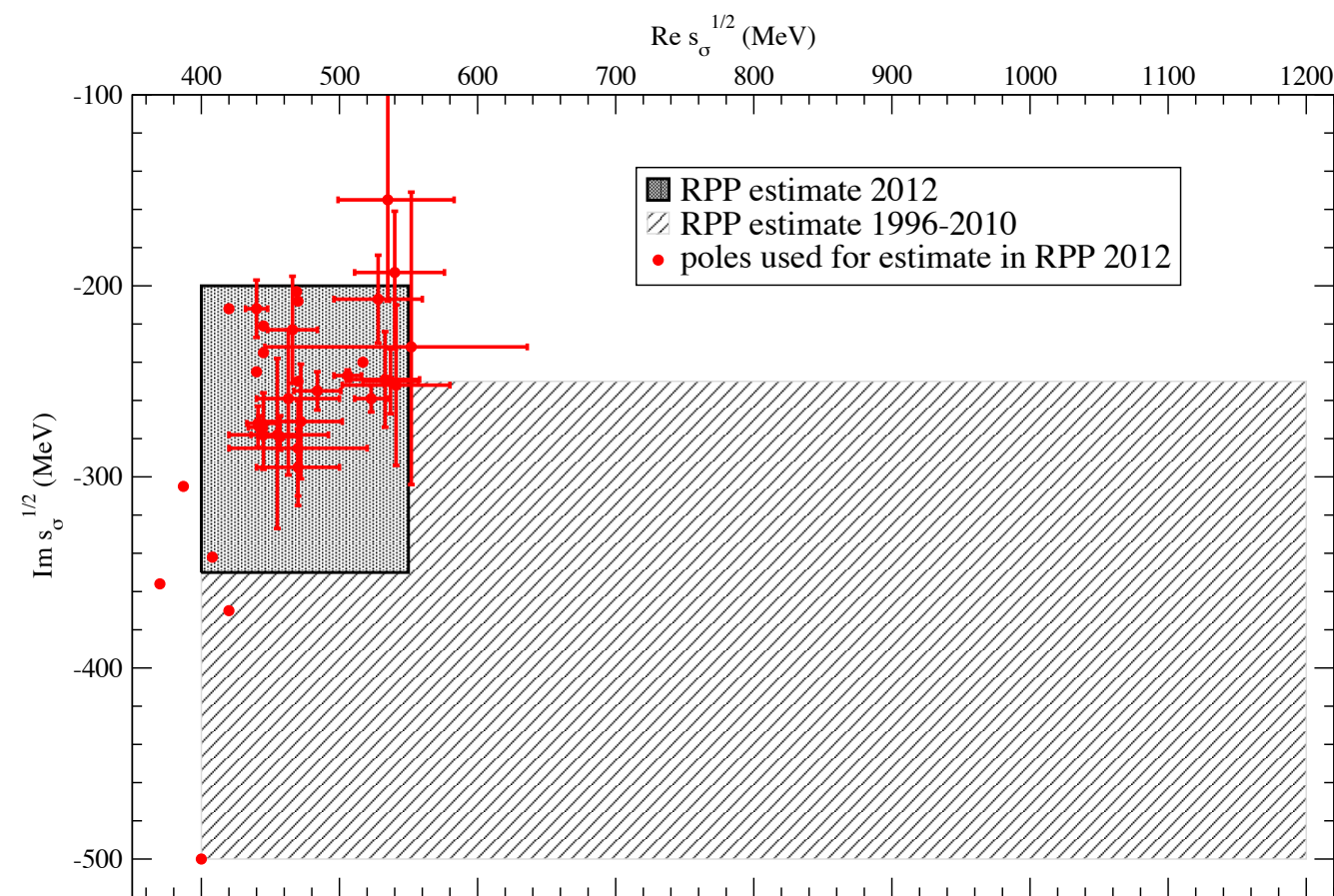
The $f_0(500)/\sigma$ resonance

from J. R. Pelaez, arXiv:1510.00653



The $f_0(500)/\sigma$ resonance

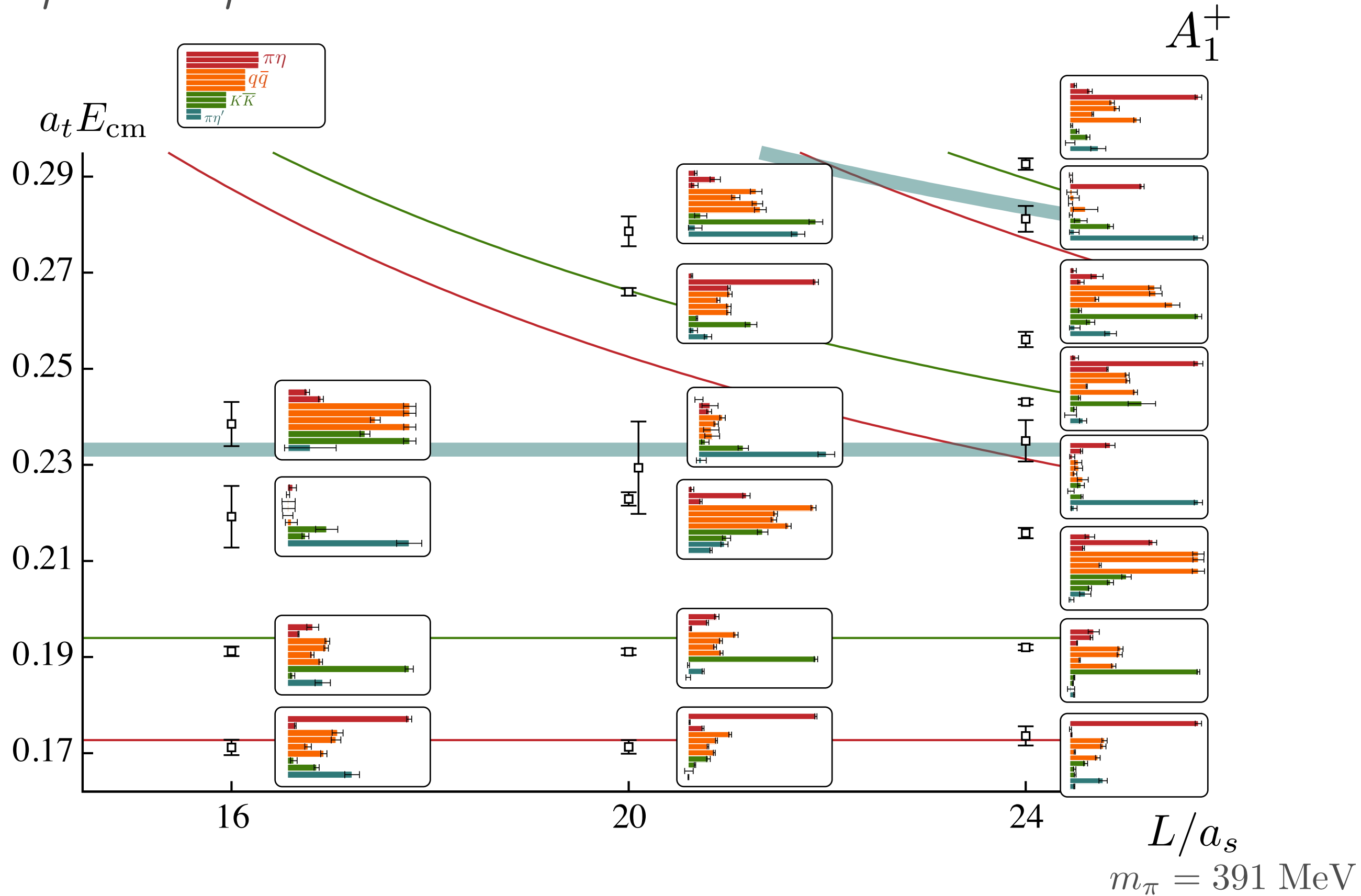
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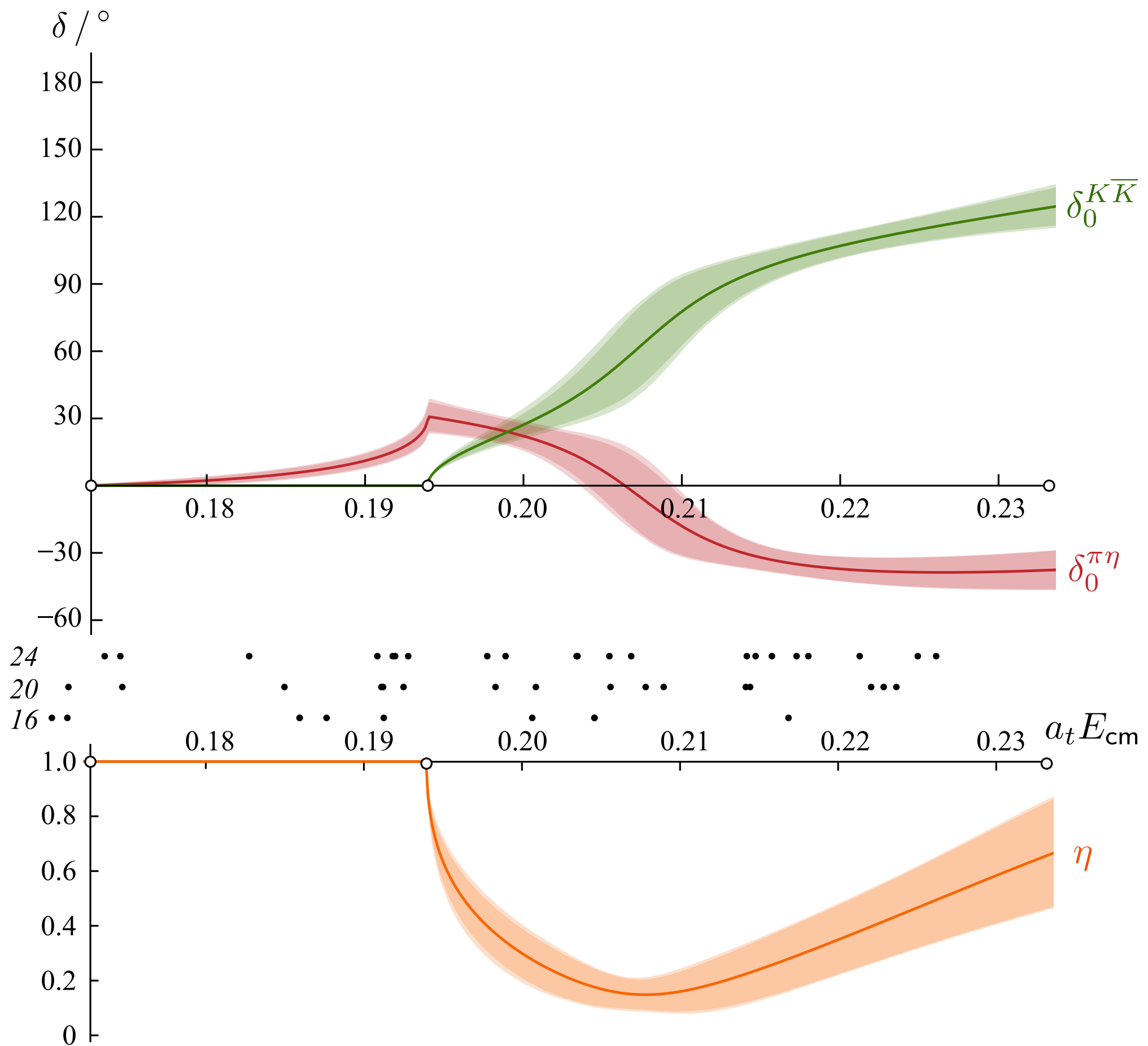


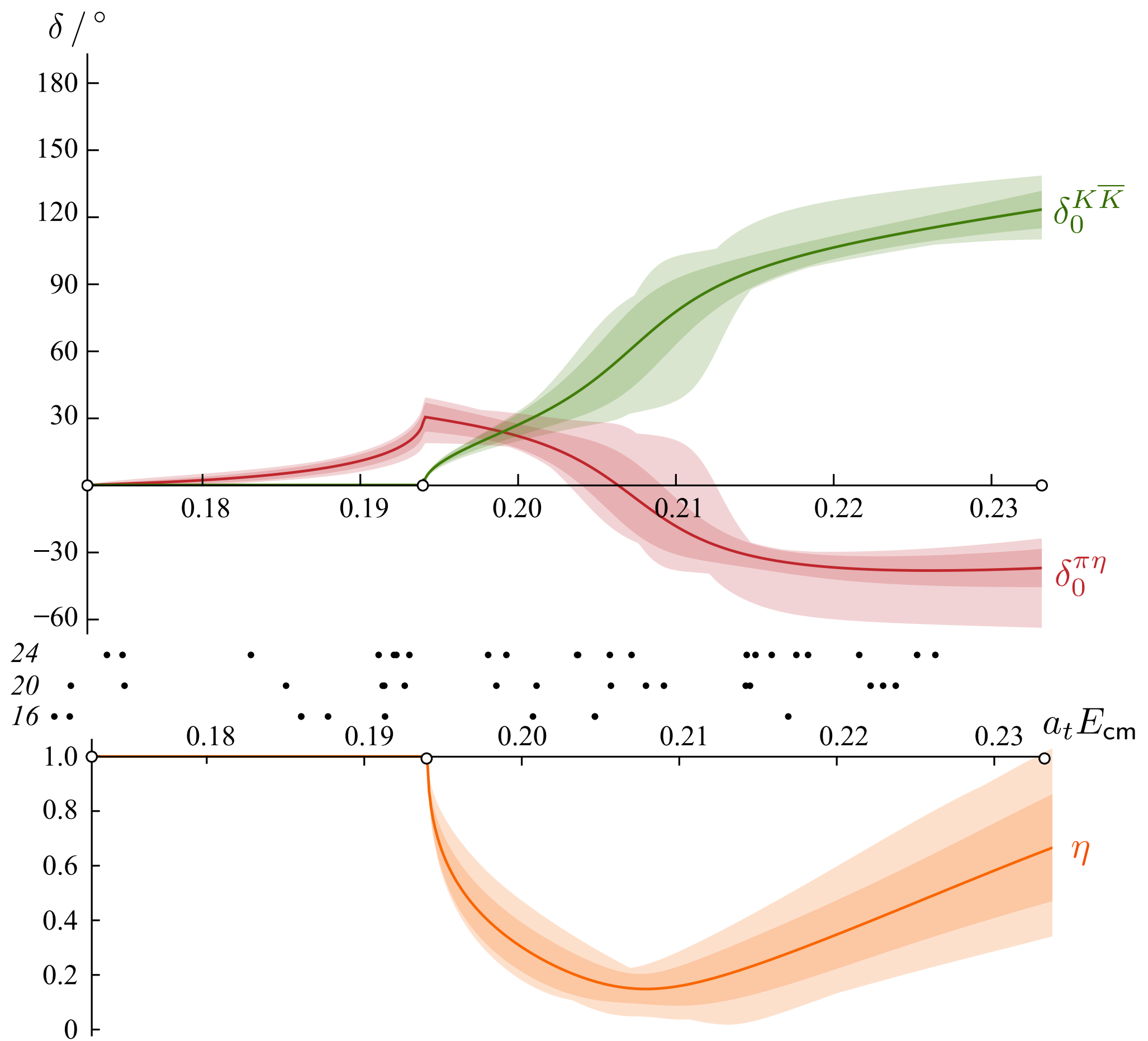
An a_0 resonance

- for more see Jozef Dudek, Tuesday 26 Jul 2016 at 16:50

$$\pi\eta-K\bar{K}-\pi\eta'$$

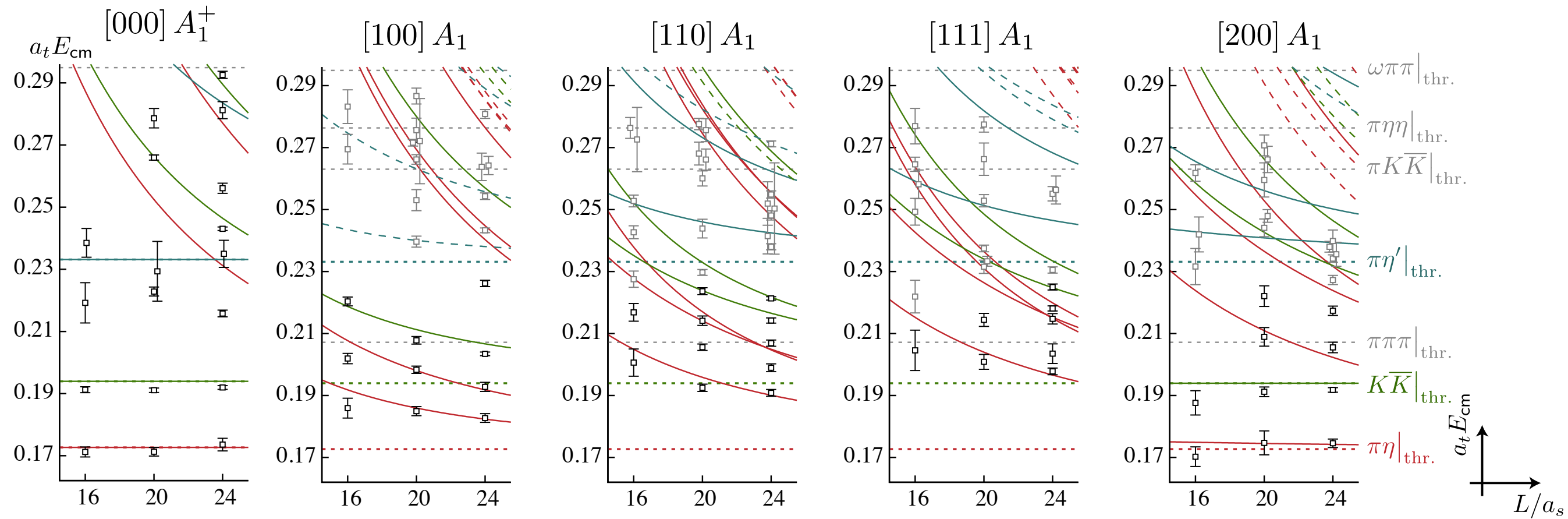






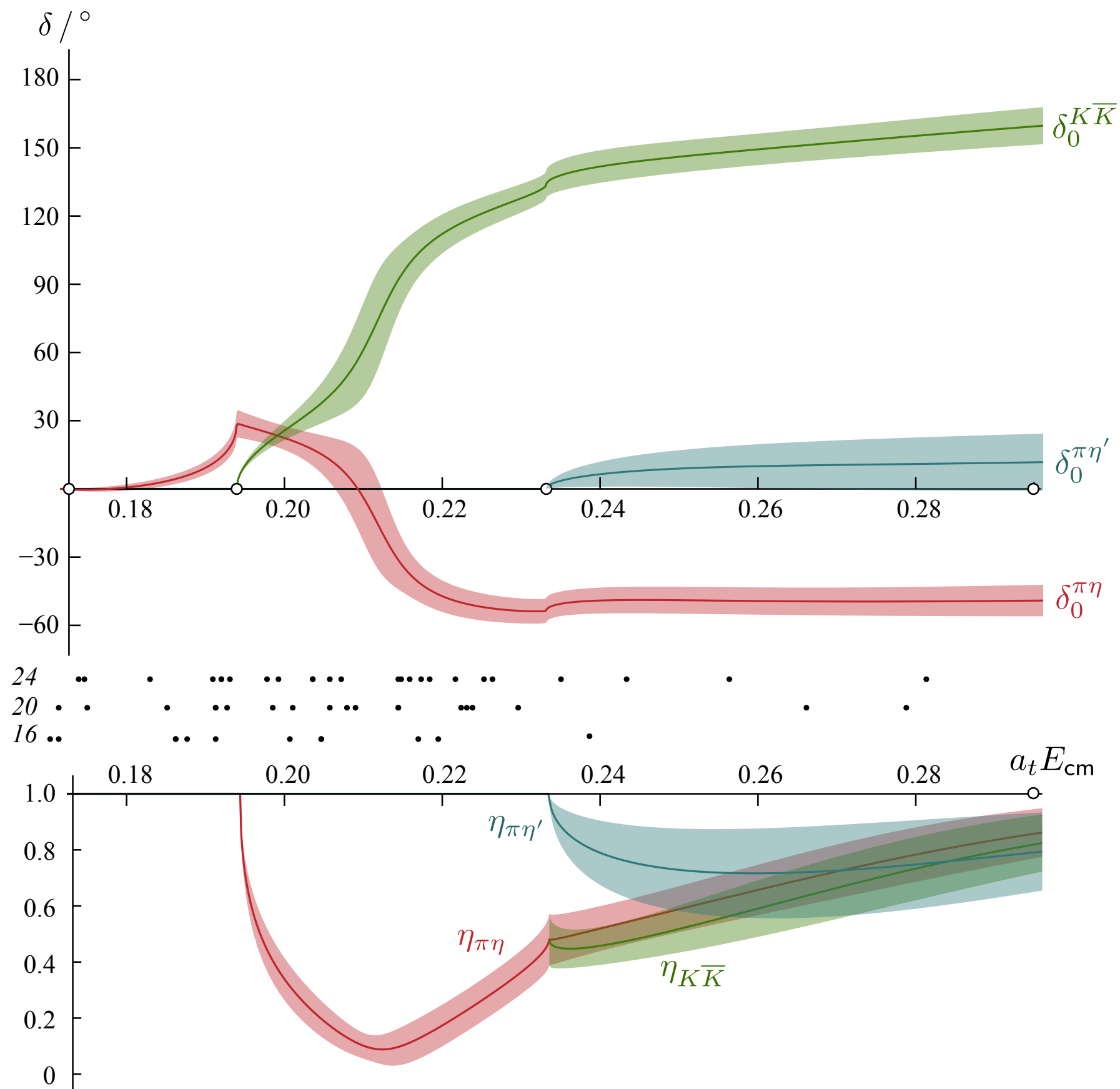
An a_0 resonance

$$\pi\eta-K\bar{K}-\pi\eta'$$



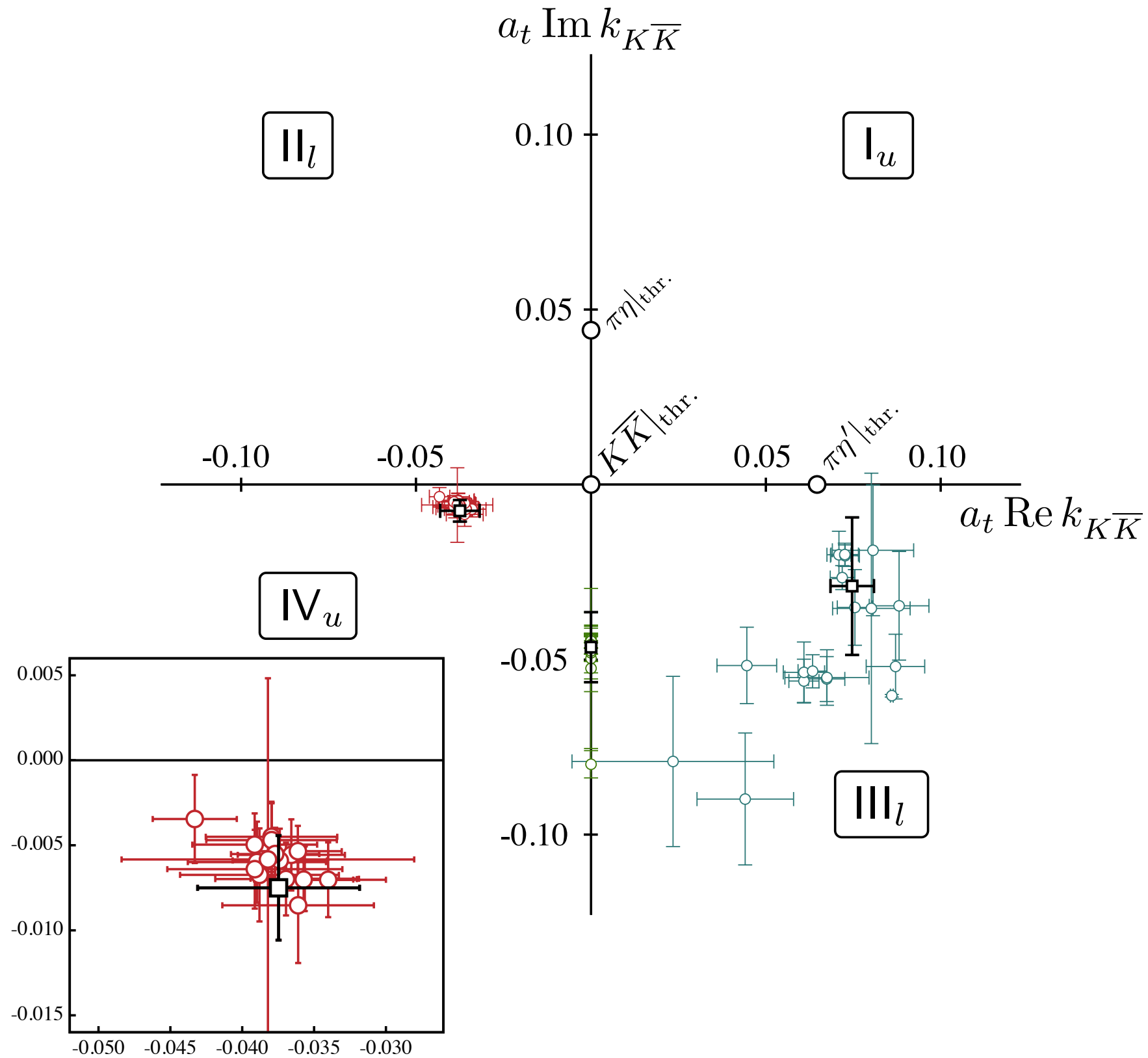
$$m_\pi = 391 \text{ MeV}$$

An a_0 resonance - three channel region



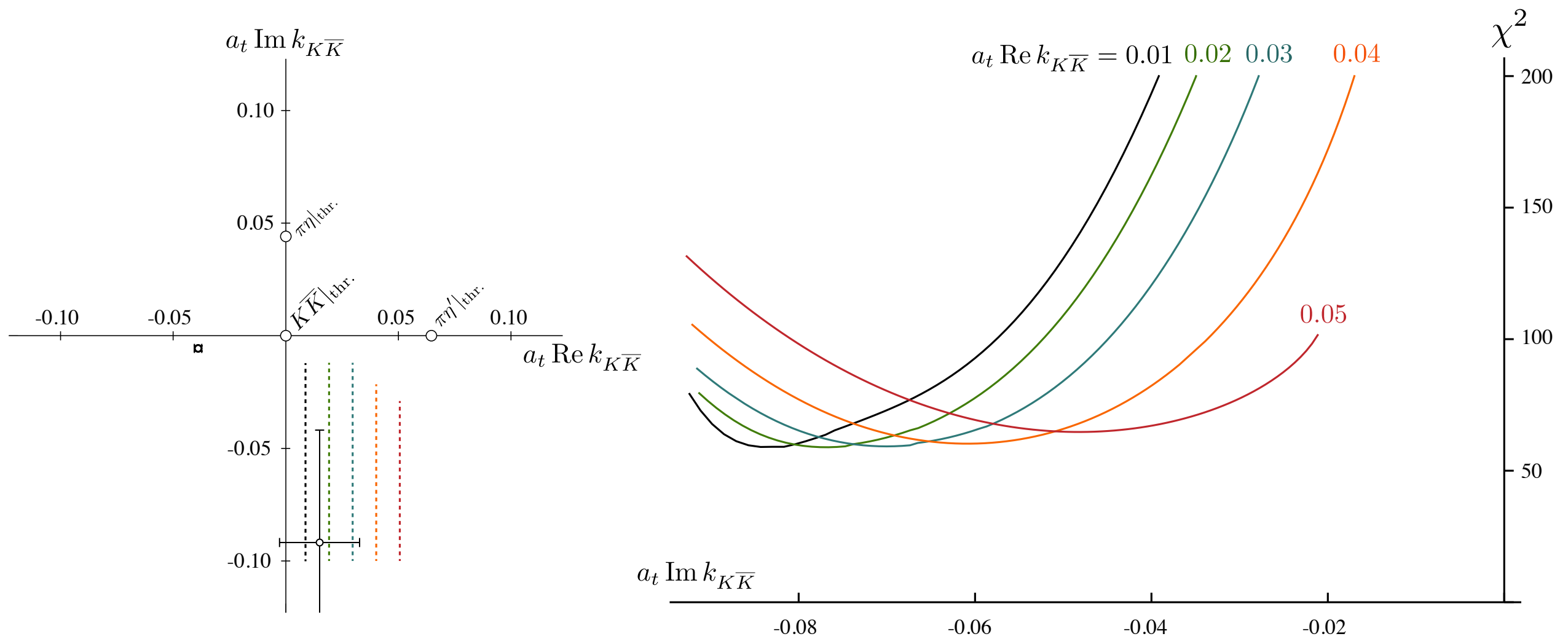
$$m_\pi = 391 \text{ MeV}$$

Poles



$$m_\pi = 391 \text{ MeV}$$

Poles



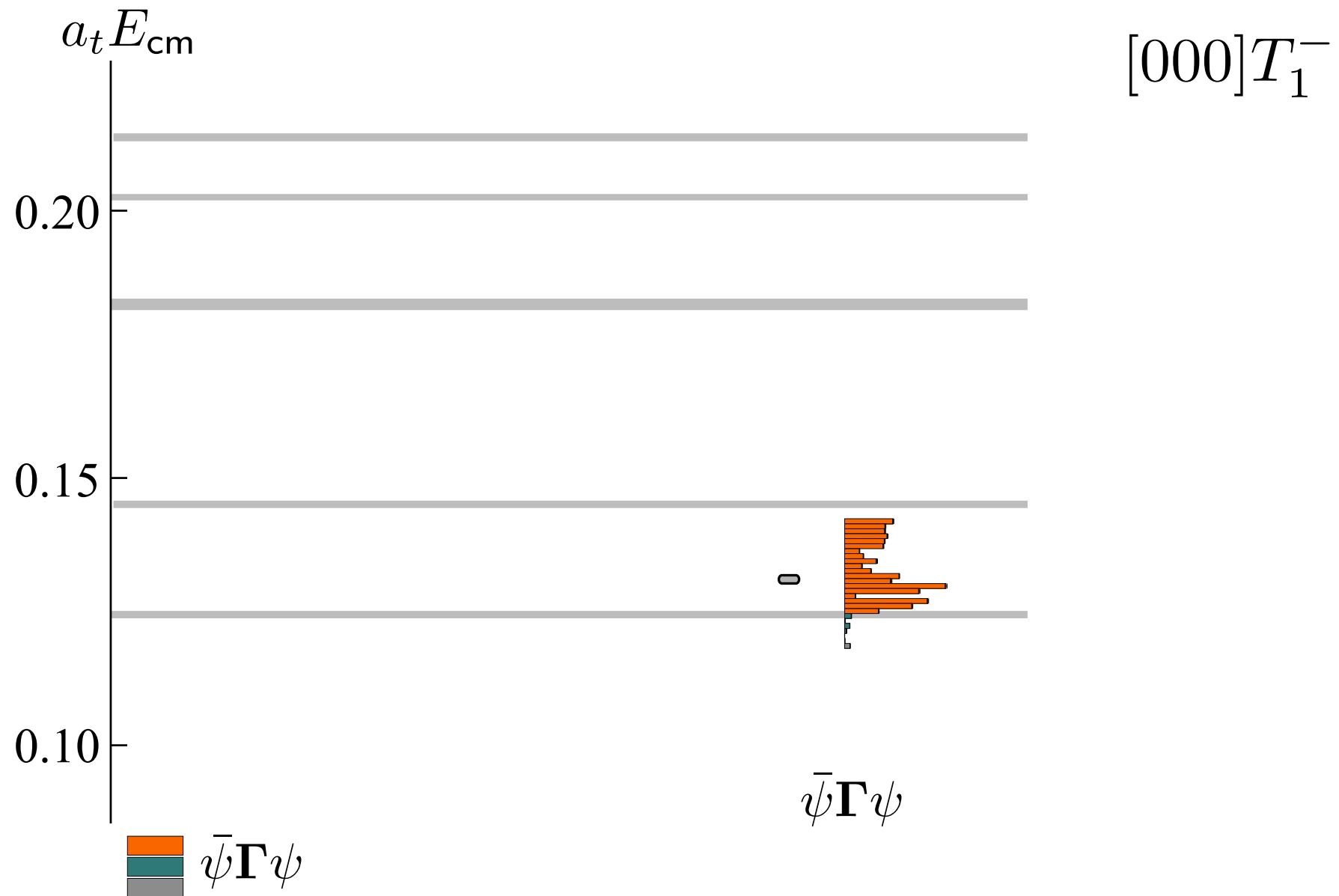
$$m_\pi = 391 \text{ MeV}$$

Extracting resonance properties

build a large basis of operators: $\mathcal{O}^\dagger \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$

compute large correlation matrices: $C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$

solve GEVP: $C_{ij}(t) v_j^n = \lambda_n(t, t_0) C_{ij}(t_0) v_j^n$

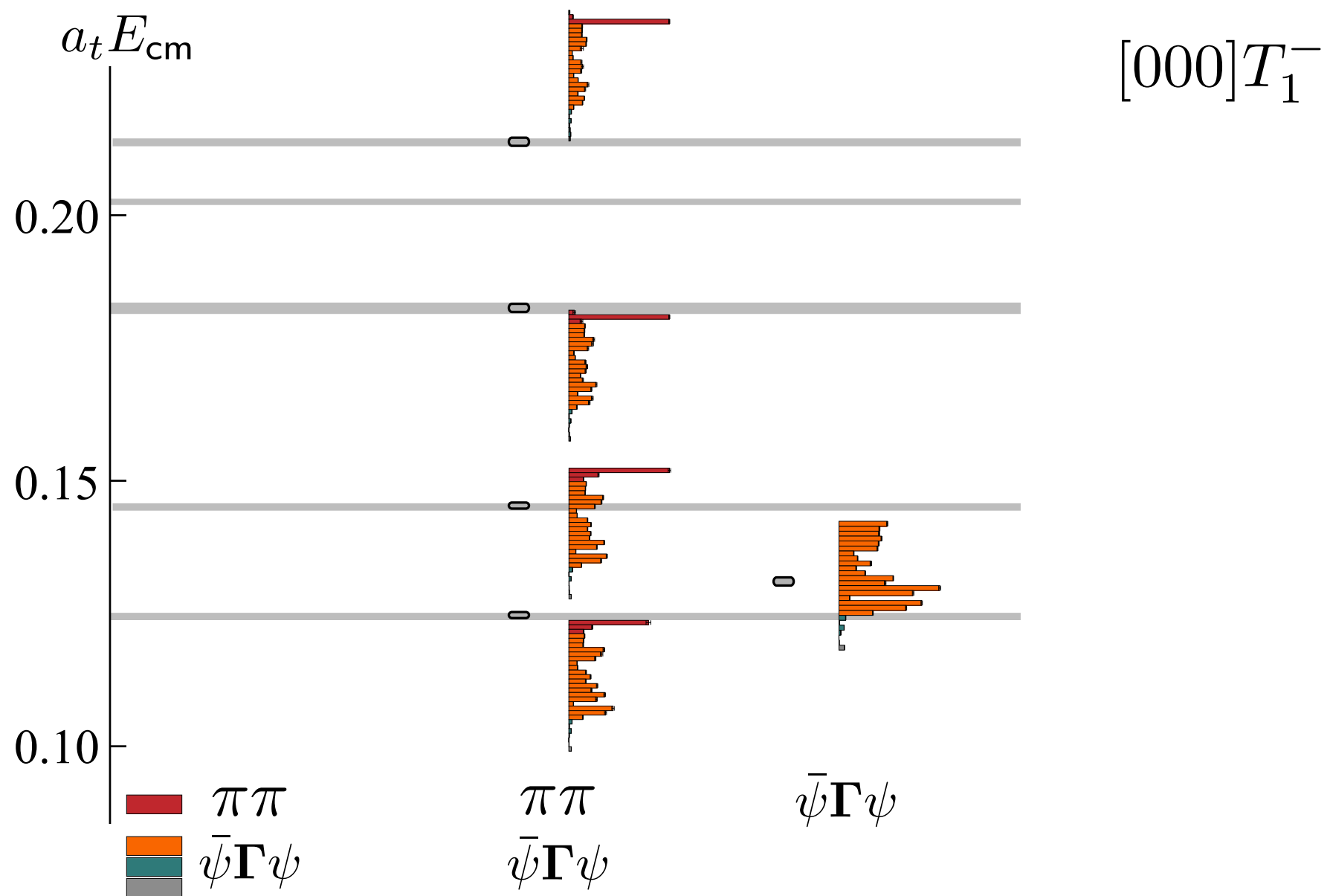


$$m_\pi = 236 \text{ MeV}$$

Extracting resonance properties

add in $\pi\pi$ operators using a variationally optimal pion $\pi^\dagger = \sum_i v_i^\pi \mathcal{O}_i^\dagger$

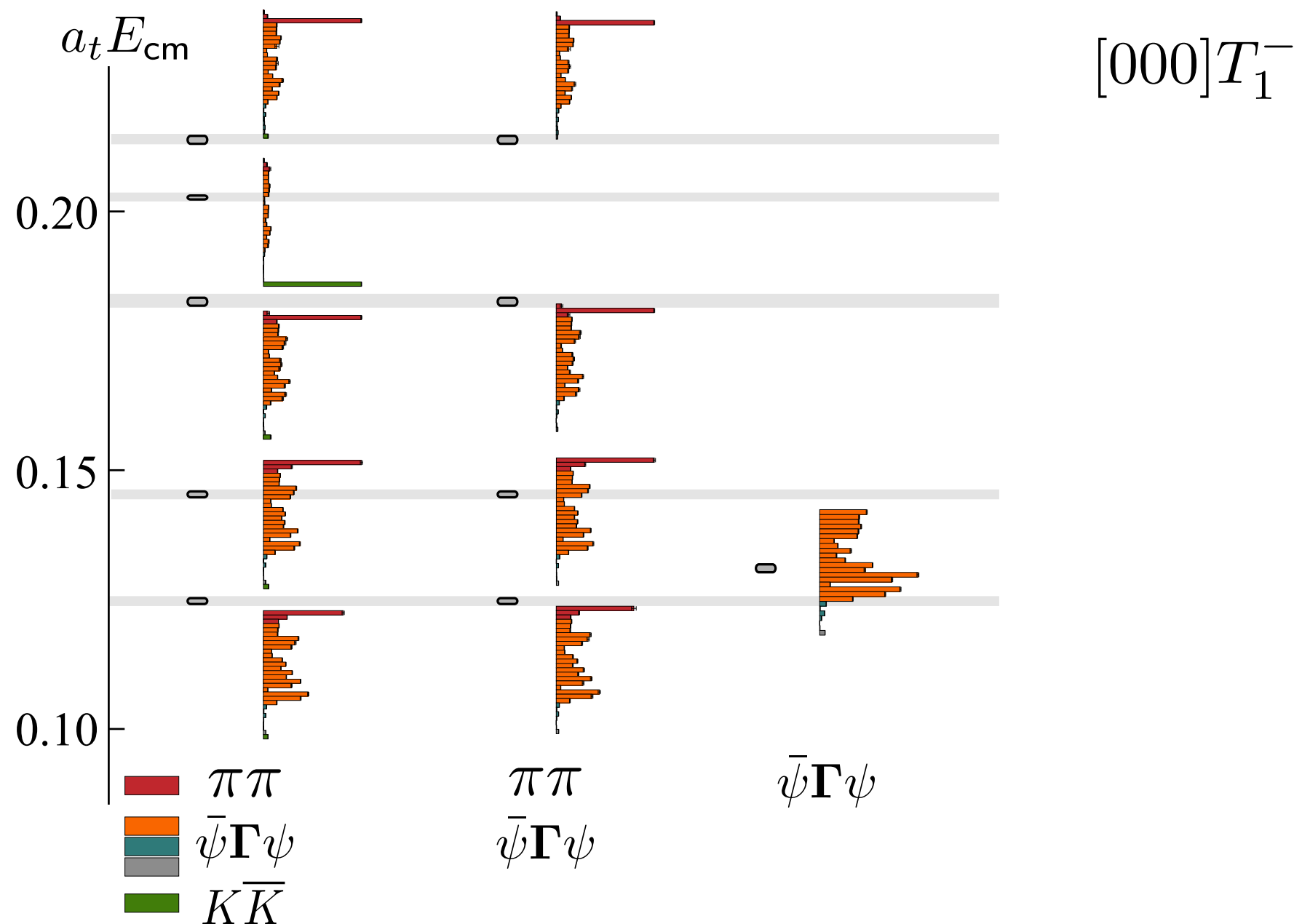
combine in pairs $(\pi\pi)^\dagger = \sum_{\vec{p}_1 + \vec{p}_2 = \vec{P}} \mathcal{C}(\vec{p}_1, \vec{p}_2) \pi^\dagger(\vec{p}_1) \pi^\dagger(\vec{p}_2)$



$$m_\pi = 236 \text{ MeV}$$

Extracting resonance properties

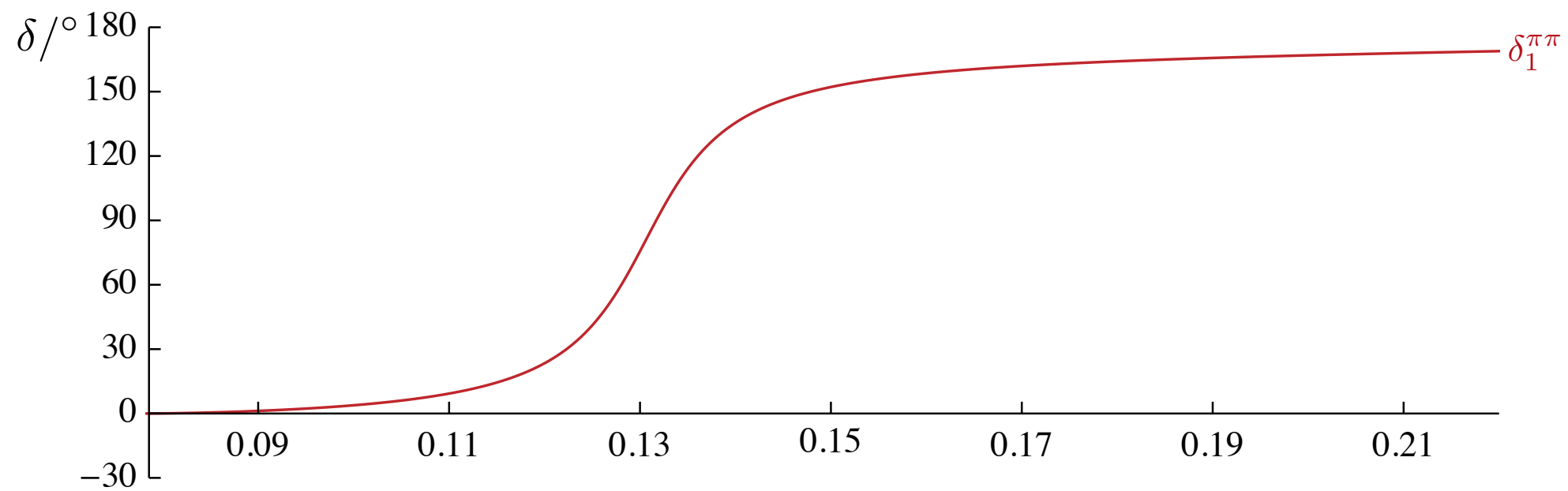
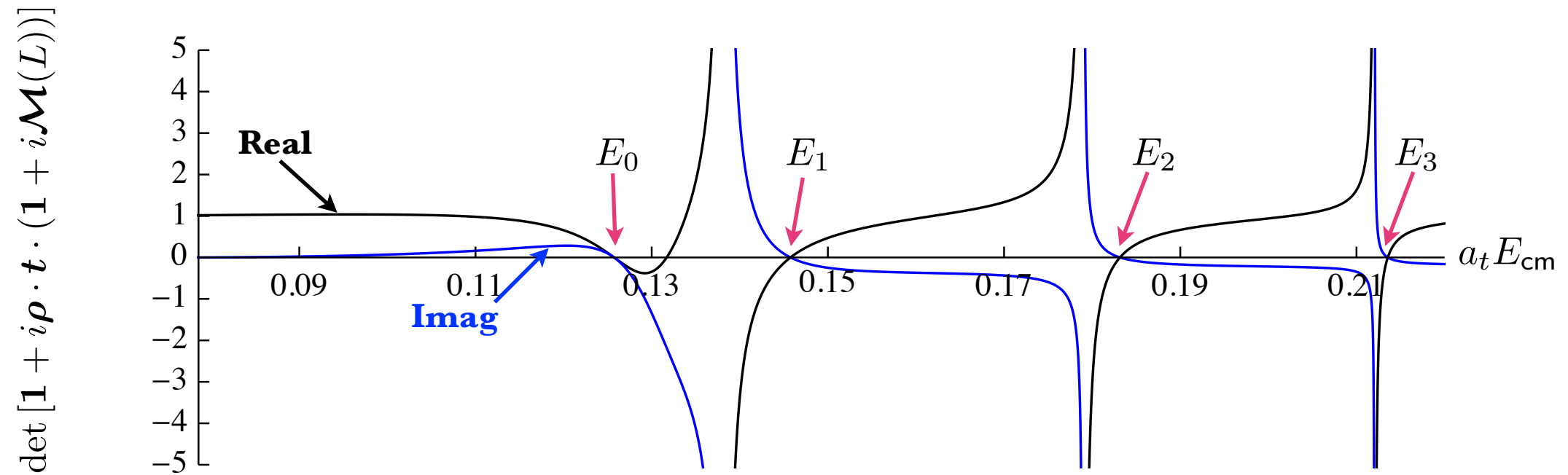
essential to have operators that overlap onto “meson” and “meson-meson” contributions to the physical spectrum



$$m_\pi = 236 \text{ MeV}$$

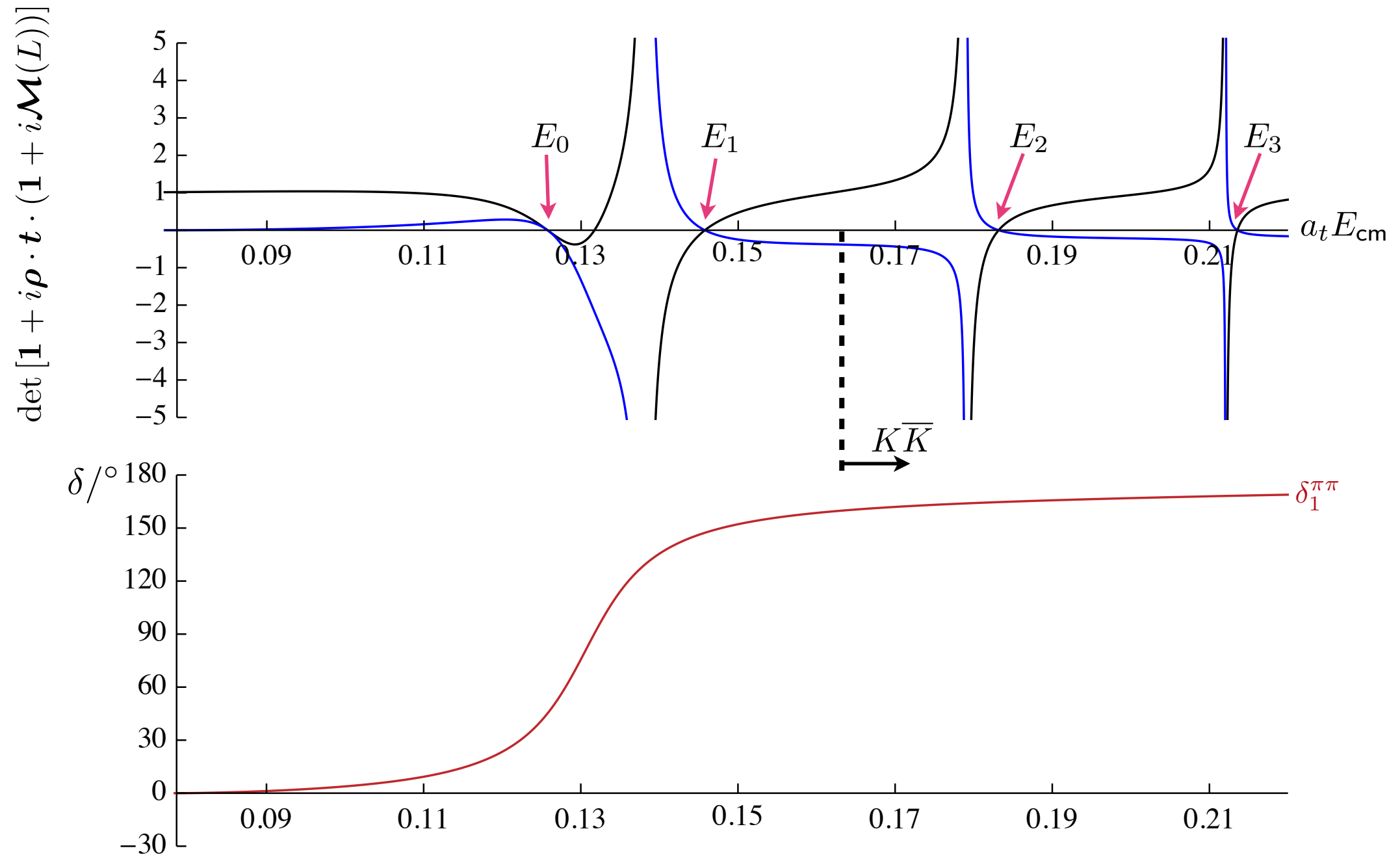
Determinant

$$t = (\pi\pi \rightarrow \pi\pi)$$



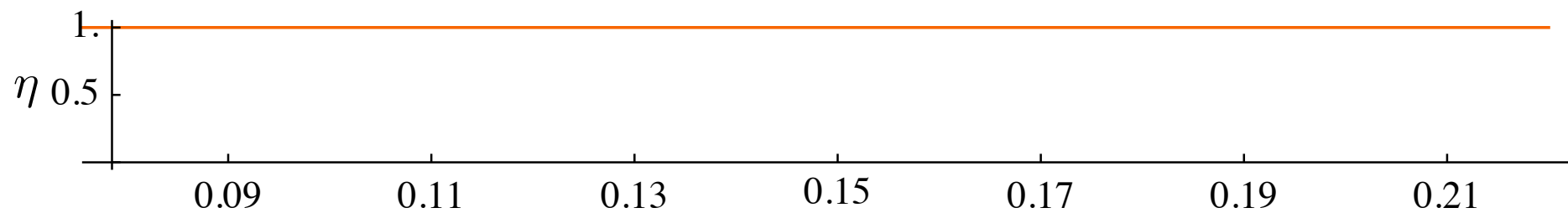
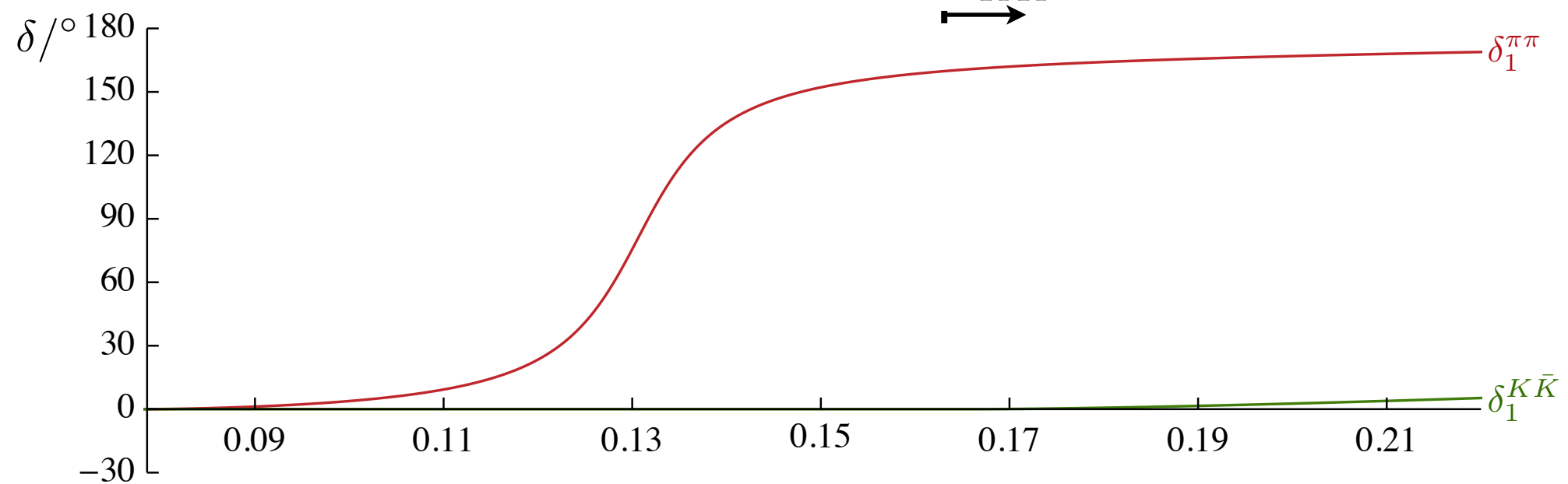
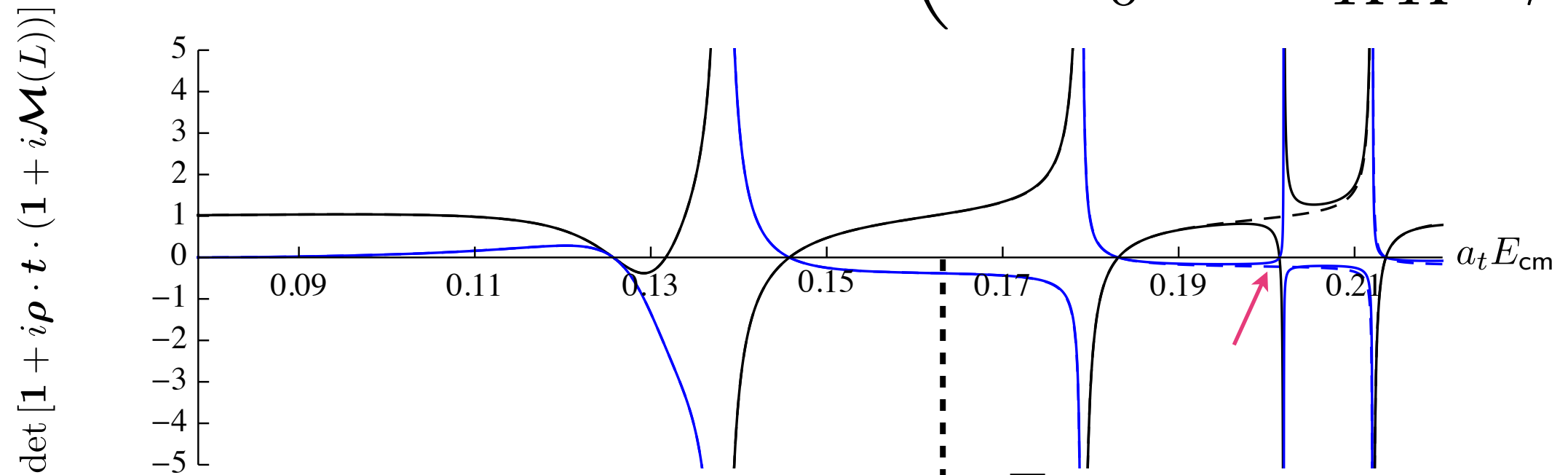
Determinant

$$t = (\pi\pi \rightarrow \pi\pi)$$



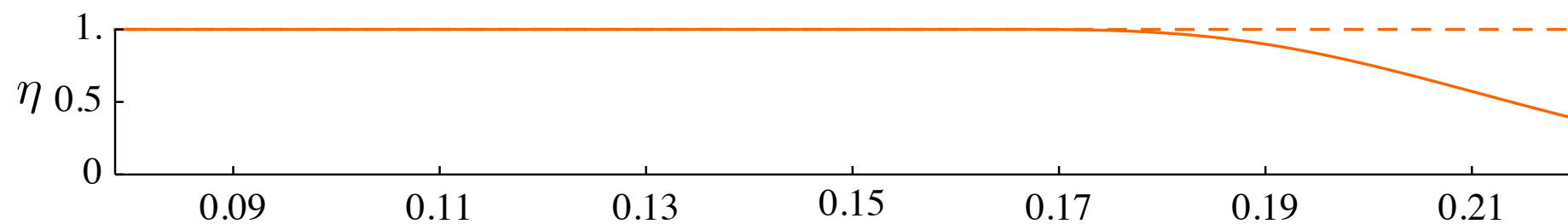
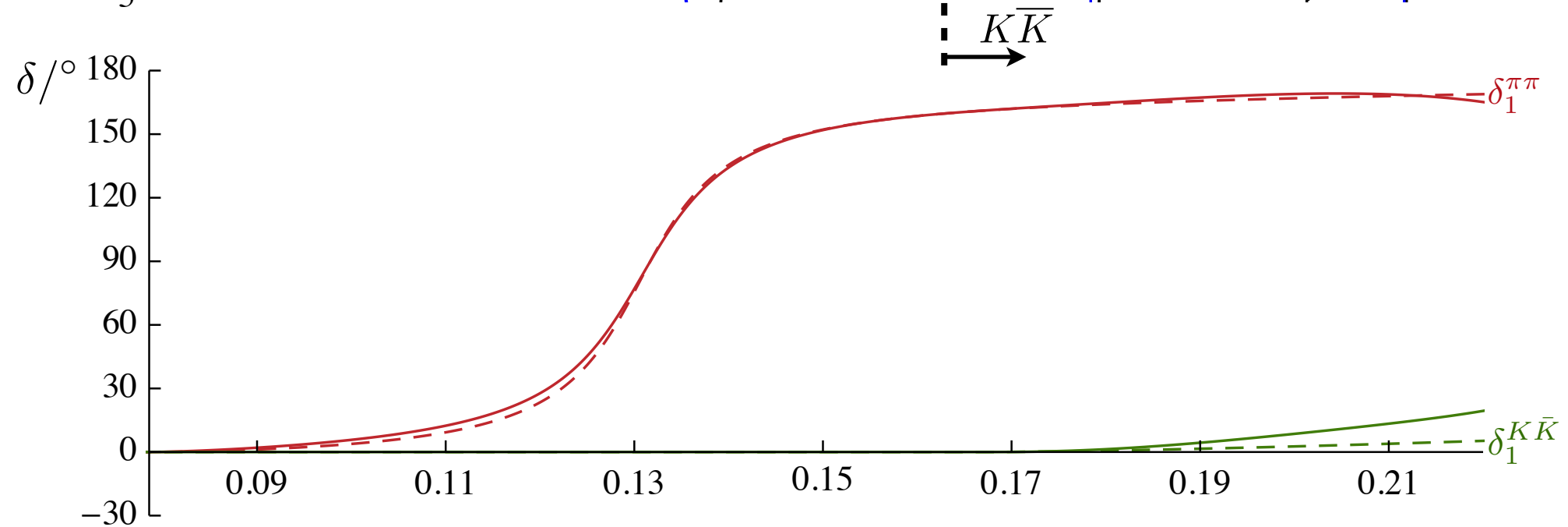
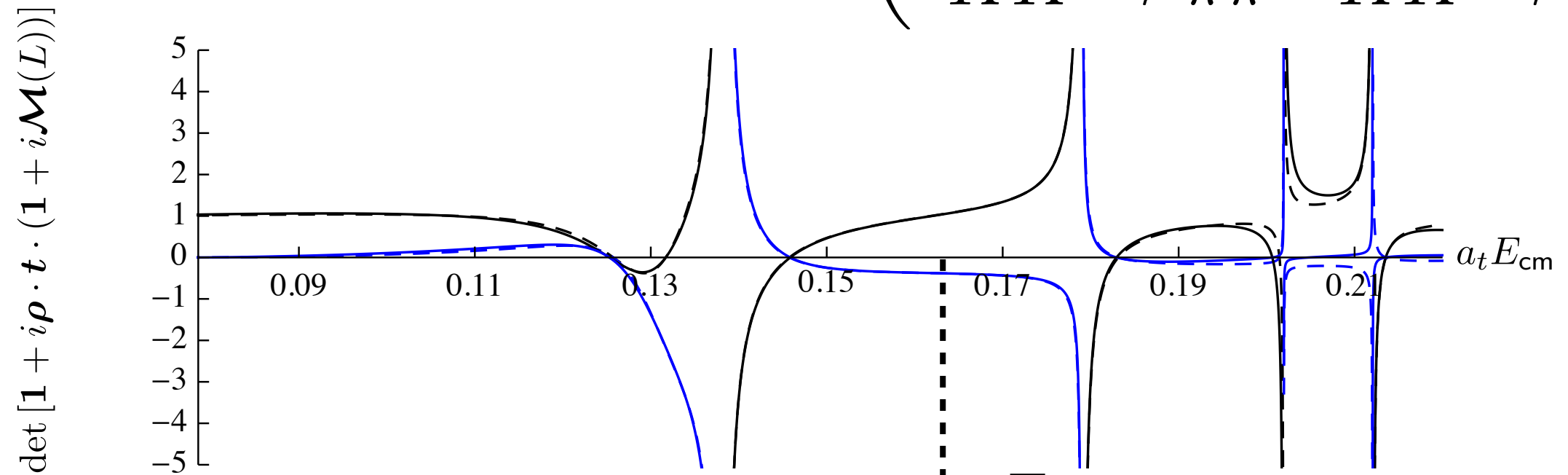
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & 0 \\ 0 & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



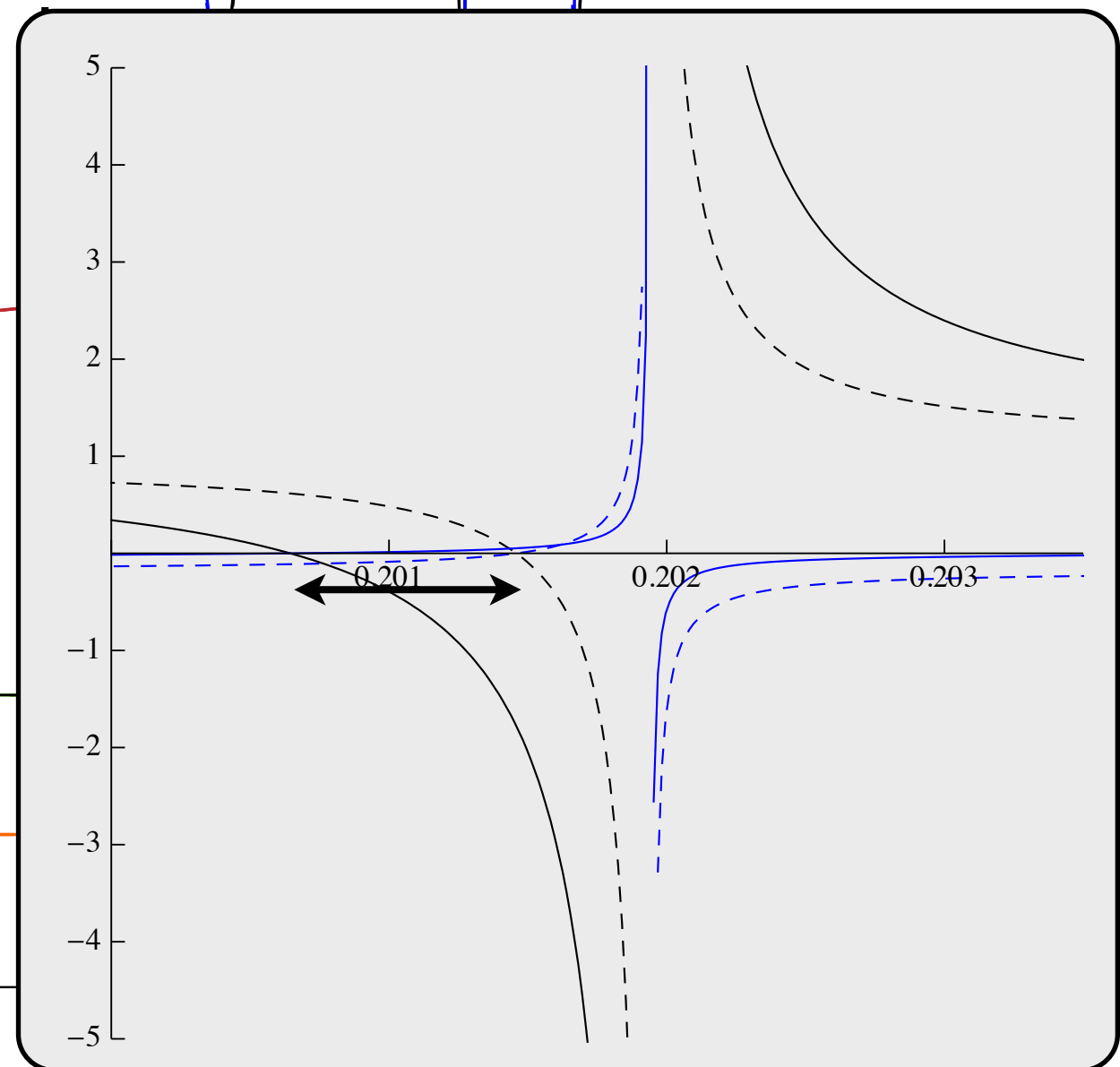
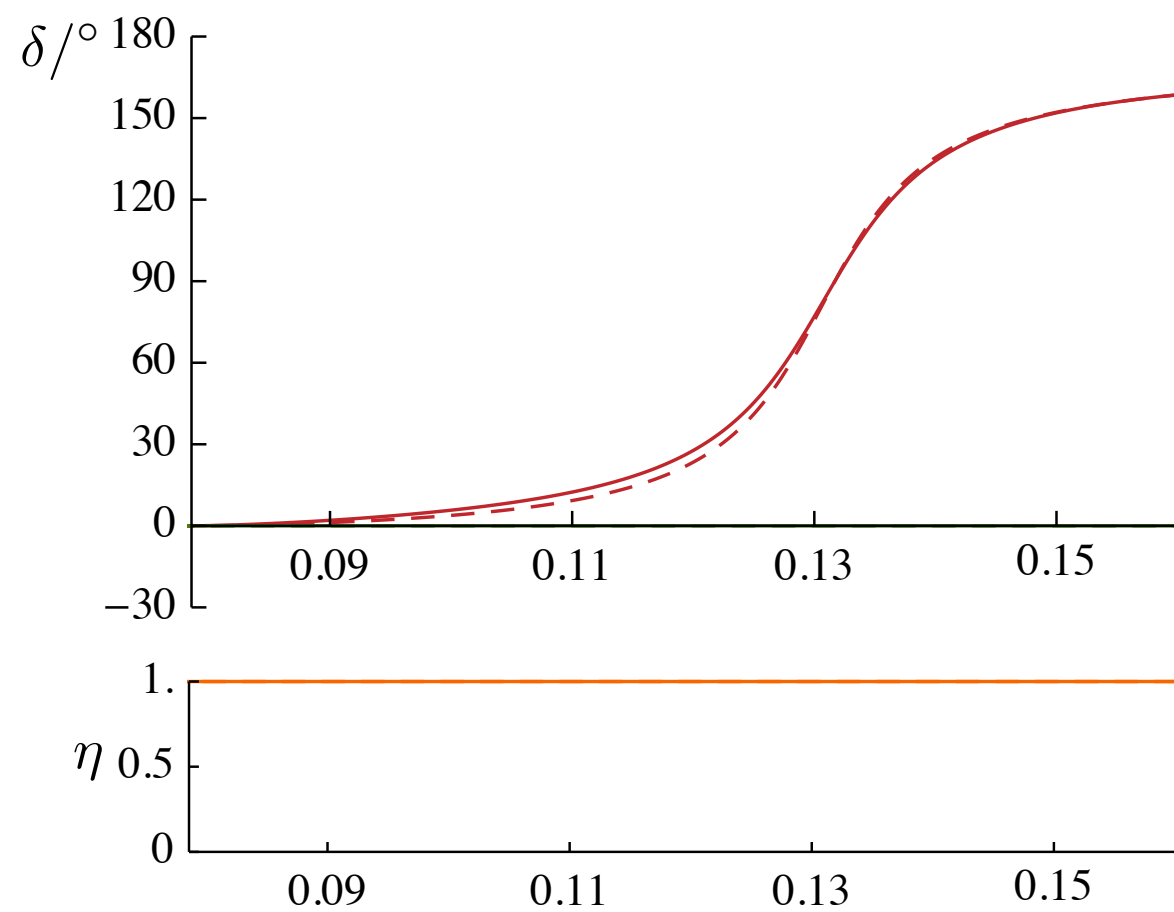
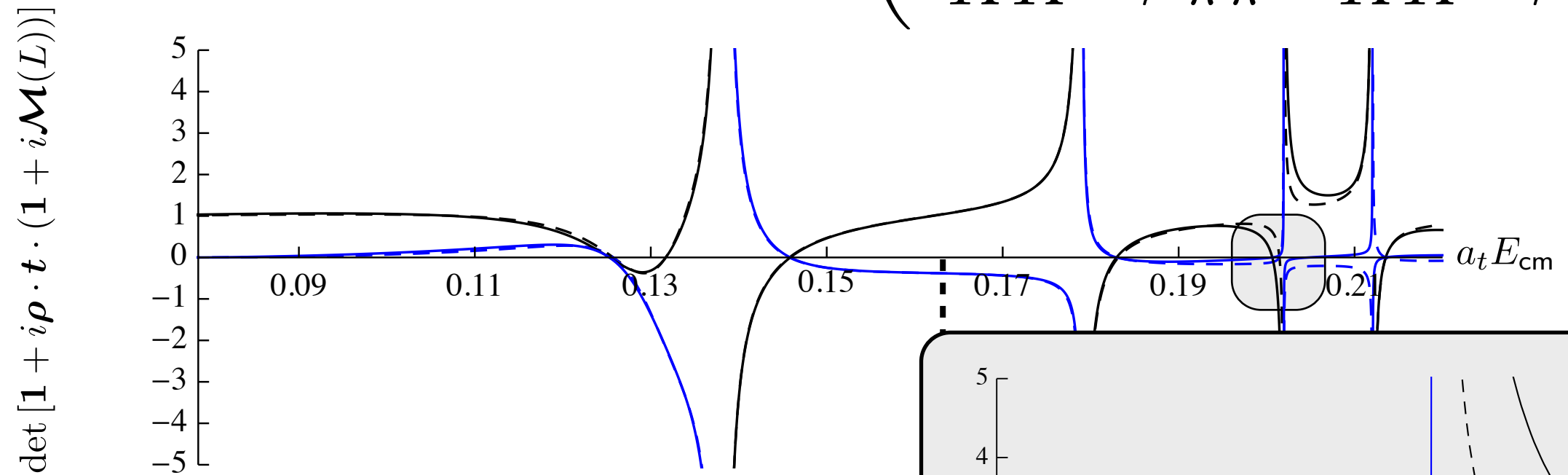
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



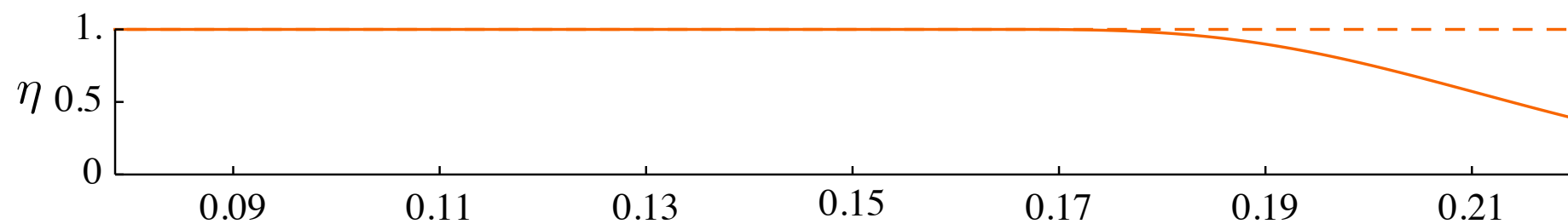
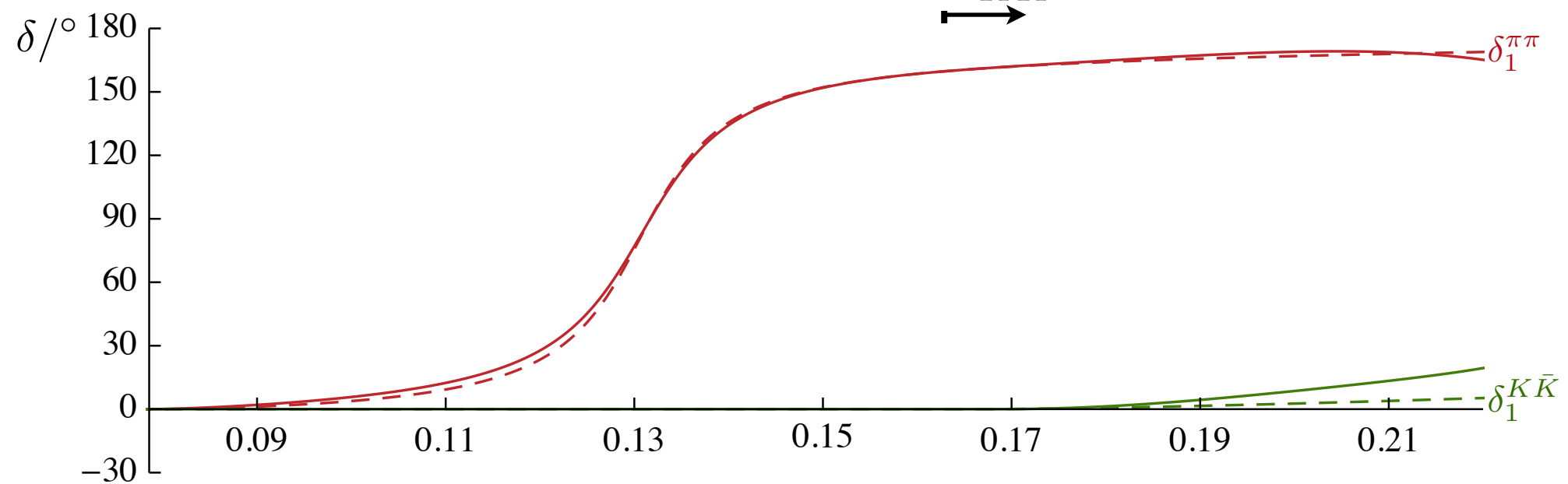
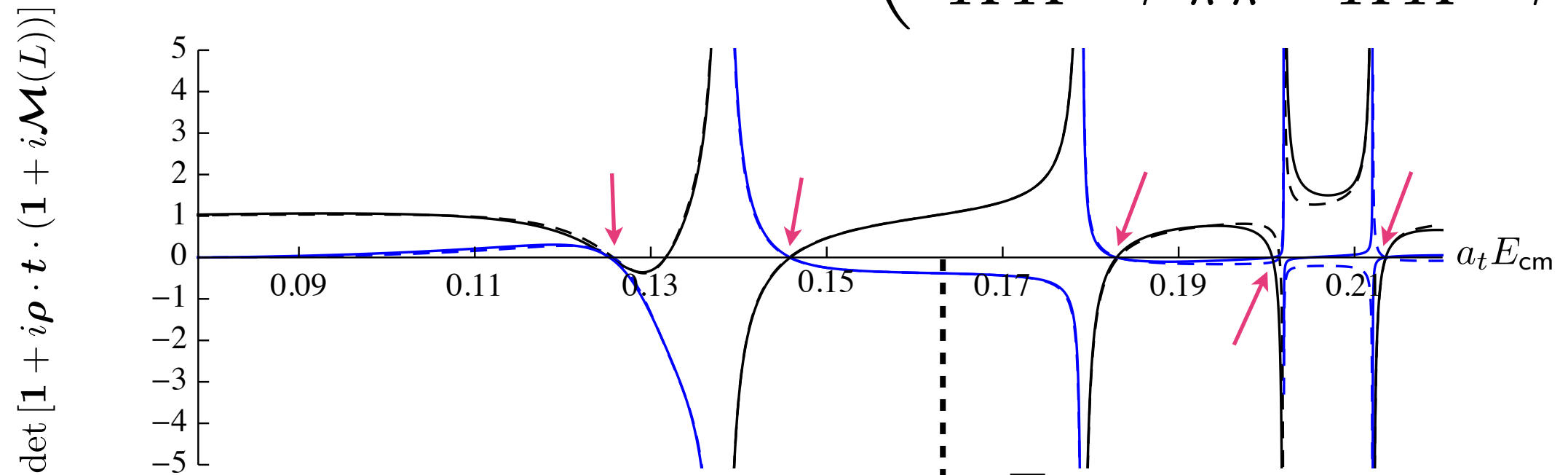
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



Chiral perturbation theory ρ extrapolation

