Lattice QCD Properties and Interactions of Light Nuclei



William Detmold, MIT

ICHEP 2016, Chicago,, August 6th 2016

The intensity frontier

- Seek new physics through quantum effects
- Precise experiments
 - Sensitivity to probe the rarest interactions of the SM
 - Look for effects where there is no SM contribution
 - Important focus of HEP(NP) experimental program
 - Dark matter direct detection
 - Neutrino physics
 - Charged lepton flavour violation, EDMs...
- Major component is nuclear targets





The intensity frontier

- Dark matter direct detection: nuclear recoils in large bucket of nuclei as signal
 - Detection rate/bounds depends on dark matter properties/dynamics and x-sec on nucleus
- Positive signals would be unambiguous
- Post-detection: precise nuclear x-sec (with quantified uncertainties) to discern underlying dynamics
- Potentially understand seemingly conflicting positive and negative signals

og₁₀(S2_b/S1) x,y,:

Inform experimental design and backgrounds



http://www.hep.ucl.ac.uk/darkMatter/



Nuclear uncertainties

How well do we know nuclear matrix elements?

- Stark example of problems: Gamow-Teller transitions in nuclei
 - Well measured for large range of nuclei (30<A<60)
 - Many nuclear structure calcs (QRPA, shell-model,...) – spectrum well described
 - Matrix elements systematically off by 20–30%
 - "Correct" by "quenching" axial charge in nuclei ...



$$T(GT) \sim \sqrt{\sum_{f} \langle \boldsymbol{\sigma} \cdot \boldsymbol{\tau} \rangle_{i \to f}}$$

$$\langle \boldsymbol{\sigma} \boldsymbol{\tau}
angle = rac{\langle f || \sum_k \boldsymbol{\sigma}^k \boldsymbol{t}_{\pm}^k || i
angle}{\sqrt{2J_i + 1}}$$

Nuclear theory at the intensity frontier

- Coming need for precision determinations of nuclear matrix elements
 - Must be based on the Standard Model (no hand-waving)
 - Must have fully quantified uncertainties
 - Timeframe and precision goals set by experiment
- Current state is far from this
- Need to develop appropriate tools



Precision Nuclear Physics

Goal: Predictive capability for properties of nuclei

- Exploit effective degrees of freedom
- Establish quantitative control through linkages between different methods
 - QCD forms a foundation determines few body interactions & matrix elements
 - Match existing EFT and many body techniques onto QCD



Quantum Chromodynamics

- Lattice QCD: tool to deal with quarks and gluons
 - Formulate problem as functional integral over quark and gluon d.o.f. on R₄

$$\langle \mathcal{O} \rangle = \int dA_{\mu} dq d\bar{q} \, \mathcal{O}[q, \bar{q}, A] e^{-S_{QCD}[q, \bar{q}, A]}$$

- Discretise and compactify system
 - Finite but large number of d.o.f $(\sim 10^{10})$
- Integrate via importance sampling (average over important configurations)
- Undo the harm done in previous steps



External currents and nuclei

- Xe in LQCD not likely any time soon
- Nuclear effective field theory:
 - I-body currents are dominant
 - 2-body currents are sub-leading but non-negligible
- LQCD: determine one body current from single nucleon
- LQCD: determine few-body contributions from A=2,3,4...
- Match EFT and many body methods to LQCD to make predictions for larger nuclei





QCD for Nuclear Physics

Nuclei in LQCD are a hard



QCD for Nuclear Physics

- Nuclei in LQCD are a hard
- Physics at multiple scales



QCD for Nuclear Physics

- Nuclei in LQCD are a hard
- Physics at multiple scales
- Two exponentially difficult challenges for LQCD
 - Contraction complexity grows factorially
 - Probabilistic method statistical uncertainty grows exponentially with A (naively)



Unphysical nuclei

NPLQCD collaboration

Case study QCD with $m_u = m_d = m_s^{\text{phys}}$



I. Spectrum of light nuclei (A < 5); NN interactions

[PRD 87 (2013), 034506, PRC 88 (2013), 024003]]

- 2. Nuclear structure: magnetic moments, polarisabilities (A<5) [PRL **II3**, 252001 (2014)]
- 3. Nuclear reactions: np \rightarrow dy [PRL 115, 132001 (2015)]



Brian Tiburzi

CCNY/RBC

Martin Savage

U. Washington

Assumpta Parreno

Barcelona







Emmanuel Chan U. Washington







MIT

Will Detmold MIT



Kostas Orginos William & Mary





Phiala Shanahan MIT

Mike Wagman **U**.Washington



Light nuclei

• Light hypernuclear binding energies @ m_{π} =800 MeV



More states bound; deeper bindings;

Magnetic moments of nuclei

- Hadron/nuclear energies are modified by presence of fixed external fields
- Eg: fixed B field

$$E_{h;j_z}(\mathbf{B}) = \sqrt{M_h^2 + (2n+1)|Q_h eB|} - \mu_h \cdot \mathbf{B} - 2\pi \beta_h^{(M0)} |\mathbf{B}|^2 - 2\pi \beta_h^{(M2)} \langle \hat{T}_{ij} B_i B_j \rangle + ..$$

- QCD calculations with multiple fields enable extraction of coefficients of response
 - Eg: magnetic moments, polarisabilities, ...
 - Not restricted to simple EM fields (axial, twist-2,...)





Magnetic field in z-direction (strength quantised by lattice periodicity)

Magnetic moments from spin splittings

$$\delta E^{(B)} \equiv E^{(B)}_{+j} - E^{(B)}_{-j} = -2\mu |\mathbf{B}| + \gamma |\mathbf{B}|^3 + \dots$$

 Extract splittings from ratios of correlation functions

$$R(B) = \frac{C_j^{(B)}(t) \ C_{-j}^{(0)}(t)}{C_{-j}^{(B)}(t) \ C_j^{(0)}(t)} \xrightarrow{t \to \infty} Z e^{-\delta E^{(B)}t}$$

 Careful to be in single exponential region of each correlator



[NPLQCD PRL **II3**, 252001 (2014)]

Magnetic moments of nuclei



Thermal Neutron Capture Cross-Section

[NPLQCD PRL 115, 132001 (2015)]

- Thermal neutron capture cross-section: $np \rightarrow d\gamma$
 - Critical process in Big Bang Nucleosynthesis
 - Historically important: nucleus is not just nucleons

 $d = np ({}^{3}S_{1})$

First QCD nuclear reaction!

np ($|S_0)$

np→dγ

Presence of magnetic field mixes $I_z=J_z=0$ ³S₁ and ¹S₀ *np* systems



- Calculate energies in presence of B fields
- Shift of eigenvalues determined by transition amplitude [WD, & M Savage 2004]

$$\Delta E_{3S_1,1S_0} = \mp \left(\kappa_1 + \overline{L}_1\right) \frac{eB}{M} + \dots$$



Lattice correlator with ${}^{3}S_{1}$ source and ${}^{1}S_{0}$ sink

$$Iz=Jz=0 \text{ correlation matrix}$$
$$C(t; \mathbf{B}) = \begin{pmatrix} C_{3S_{1},3S_{1}}(t; \mathbf{B}) & C_{3S_{1},1S_{0}}(t; \mathbf{B}) \\ C_{1S_{0},3S_{1}}(t; \mathbf{B}) & C_{1S_{0},1S_{0}}(t; \mathbf{B}) \end{pmatrix}$$

np→dγ

Generalised eigenvalue problem

$$[\mathbf{C}(t_0;\mathbf{B})]^{-1/2}\mathbf{C}(t;\mathbf{B})[\mathbf{C}(t_0;\mathbf{B})]^{-1/2}v = \lambda(t;\mathbf{B})v$$

Ratio of correlator ratios to extract 2-body

$$R_{{}^{3}\!S_{1},{}^{1}\!S_{0}}(t;\mathbf{B}) = \frac{\lambda_{+}(t;\mathbf{B})}{\lambda_{-}(t;\mathbf{B})} \xrightarrow{t \to \infty} \hat{Z} \exp\left[2 \ \Delta E_{{}^{3}\!S_{1},{}^{1}\!S_{0}}t\right]$$

$$\delta R_{3S_{1},1S_{0}}(t;\mathbf{B}) = \frac{R_{3S_{1},1S_{0}}(t;\mathbf{B})}{\Delta R_{p}(t;\mathbf{B})/\Delta R_{n}(t;\mathbf{B})} \to A \ e^{-\delta E_{3}} S_{1,1S_{0}}(\mathbf{B})t$$

$$\delta E_{{}^{3}S_{1},{}^{1}S_{0}} \equiv \Delta E_{{}^{3}S_{1},{}^{1}S_{0}} - [E_{p,\uparrow} - E_{p,\downarrow}] + [E_{n,\uparrow} - E_{n,\downarrow}]$$

$$\rightarrow 2\overline{L}_{1}|e\mathbf{B}|/M + \mathcal{O}(\mathbf{B}^{2})$$

[NPLQCD PRL **115**, 132001 (2015)]

np→dγ

Correlator ratios 0.8 1.2 m_{π} =450 MeV m_{π} =800 MeV 1.0 0.6 $\delta R^{3}S_{1,1}S_{0}(t;B)$ F 0.8 $\begin{array}{c} \delta R_{3S_1,1S_0}(t;B) \\ 0.0 \\$ 0.2 0.2 0.0 0.0 10 15 10 15 0 5 20 0 5 20 t/a t/a Field strength & mass dependence 0.4 0.00 -0.02 0.3 $a \, \delta E_{^3S_{1,}{}^1S_0}(\tilde{n})$ [<u>I</u>] [<u>N</u>] -0.04-0.060.1 -0.08 Slopes give \overline{L}_l 0.0L -0.10 L 0.2 0.6 0.4 0.8 1.0 0.0 10 2 8 12 0 4 ۲ $m_\pi^2 \, [\text{GeV}^2]$ **|B**|

[NPLQCD PRL **II5**, 132001 (2015)]

np→dγ

- Key point: extract short-distance contribution at physical mass $\overline{L}_1^{lqcd} = 0.285(^{+63}_{-60})$ NM
- Use EFT to combine with phenomenological nucleon magnetic moment, scattering parameters

$$\sigma^{\text{lqcd}}(np \to d\gamma) = 307.8(1 + 0.273 \ \overline{L}_1^{\text{lqcd}}) \ \text{mb}$$

$$\sigma^{\text{lqcd}}(np \to d\gamma) = 332.4(^{+5.4}_{-4.7}) \text{ mb}$$

c.f. phenomenological value

$$\sigma^{\text{expt}}(np \to d\gamma) = 334.2(0.5) \text{ mb}$$

■ NB: at m_{π} =800 MeV, use LQCD for all inputs (ab initio)

$$\sigma^{800 \text{ MeV}}(np \to d\gamma) \sim 10 \text{ mb}$$

[NPLQCD PRL **115**, 132001 (2015)]

QCD for nuclei

- Nuclei are under serious study directly from QCD
 - Spectroscopy of light nuclei and exotic nuclei
 - Structure: magnetic moments and polarisabilities
 - Electromagnetic interactions: thermal capture cross-section
 - Weak interactions: M Savage, Neutrinos 09:52 am Sat
- Prospect of a quantitative connection to QCD makes this a very exciting time for nuclear physics
 - Same techniques work for IF relevant matrix elements
 - Critical role in current and upcoming intensity frontier experimental program



