



38th INTERNATIONAL CONFERENCE
ON HIGH ENERGY PHYSICS

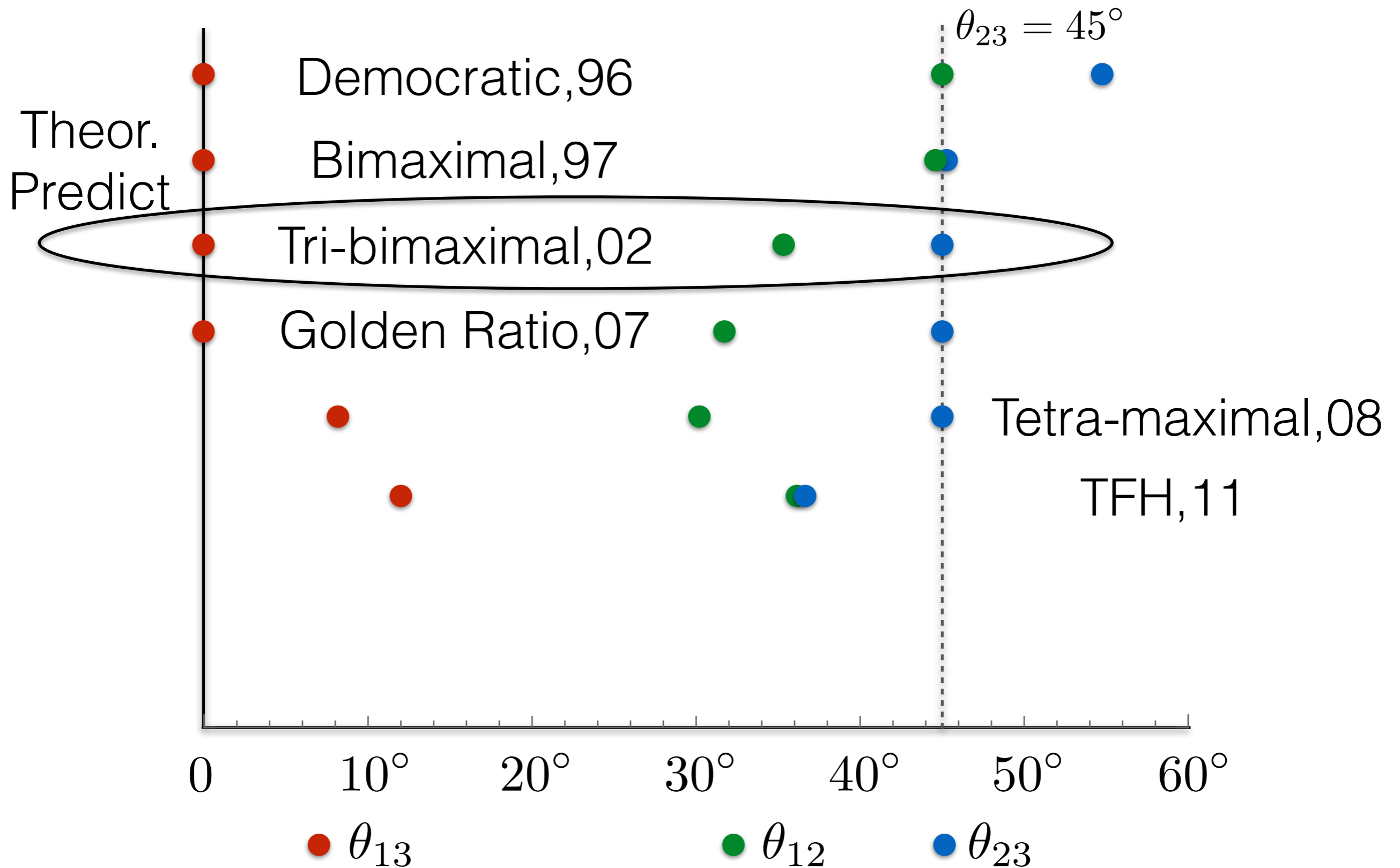
AUGUST 3 - 10, 2016
CHICAGO

Theoretical motivations for precision measurements of oscillation parameters

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Mixing angles: theoretical predictions



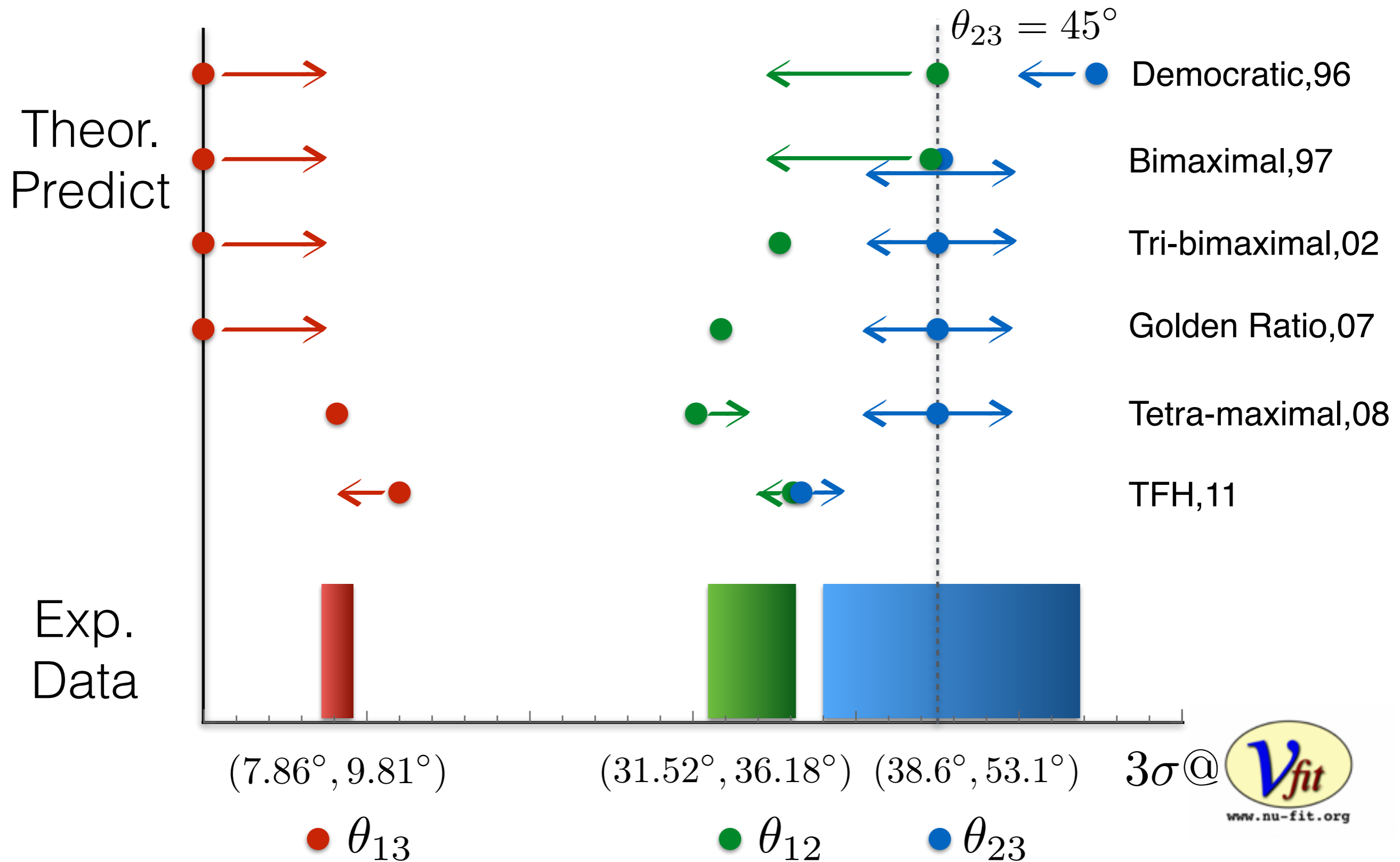
Neutrino oscillation experiments

$$\Delta m_{ji}^2 \equiv m_j^2 - m_i^2$$

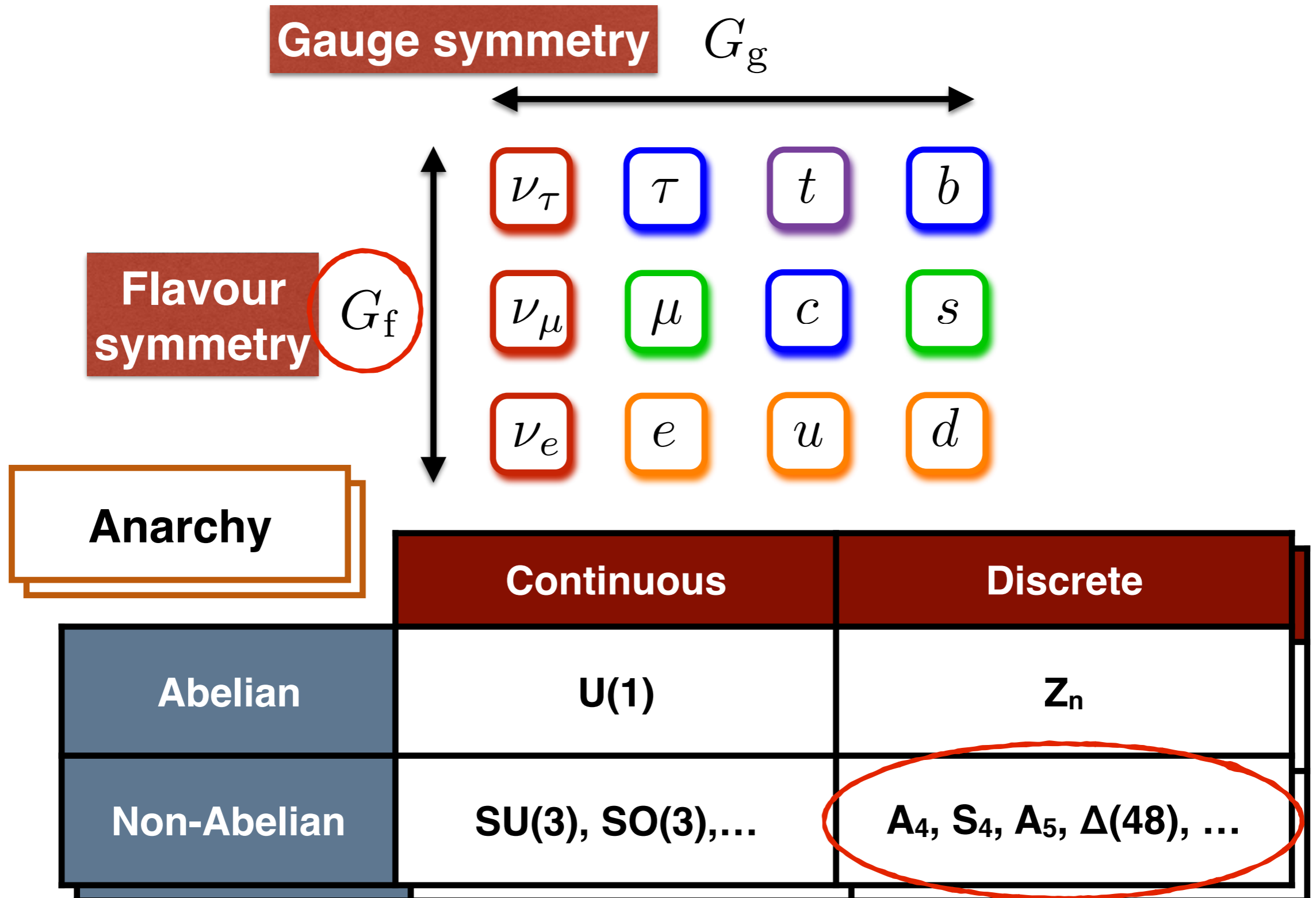
Type	Probability	Existing Experiments	$\Delta_{ji} \equiv \Delta m_{ji}^2 L / (4E)$
Solar	P_{ee}	Super-K, SNO, Borexino, ...	
	$P_{ee} \approx \sin^2 \theta_{12}$	$P_{ee} \approx 1 - \frac{1}{2} \sin^2 2\theta_{12}$ (MSW effect)	
Atmospheric	$P_{\mu\mu}, P_{\mu\tau}$	Super-K, ...	
	$P_{\mu\tau} \approx 1 - P_{\mu\mu} \approx \sin^2 2\theta_{23} \sin^2 \Delta_{31}$		
Reactor	$P_{\bar{e}\bar{e}}$	KamLAND, Daya Bay, Reno, Double Chooz, ...	
	$P_{\bar{e}\bar{e}} \approx 1 - \sin^2 2\theta_{12} \sin^2 \Delta_{21}$	$P_{\bar{e}\bar{e}} \approx 1 - \sin^2 2\theta_{13} \sin^2 \Delta_{31}$	
Accelerator	$P_{\mu\mu}, P_{\mu e}$	K2K, T2K, MINOS, NOvA, ...	
	$P_{\mu\tau} \approx 1 - P_{\mu\mu} \approx \sin^2 2\theta_{23} \sin^2 \Delta_{31}$	OPERA, $P_{\mu\tau}$	
	$P_{\mu e} \approx \frac{s_{23}^2 \sin^2 2\theta_{13}}{(1 - \hat{A})^2} \sin^2[(1 - \hat{A})\Delta_{31}] + \frac{\alpha \tilde{J} \sin(\Delta_{31} + \delta)}{\hat{A}(1 - \hat{A})} \sin(\hat{A}\Delta_{31}) \sin[(1 - \hat{A})\Delta_{31}],$		
	$\tilde{J} = \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23}, \quad \alpha = \Delta m_{21}^2 / \Delta m_{31}^2, \quad \hat{A} = A_{CC} / \Delta m_{31}^2$		

Next generation: DUNE, T2HK, JUNO, PINGU,.....

Mixing angles: theories vs experiments



Flavour symmetry

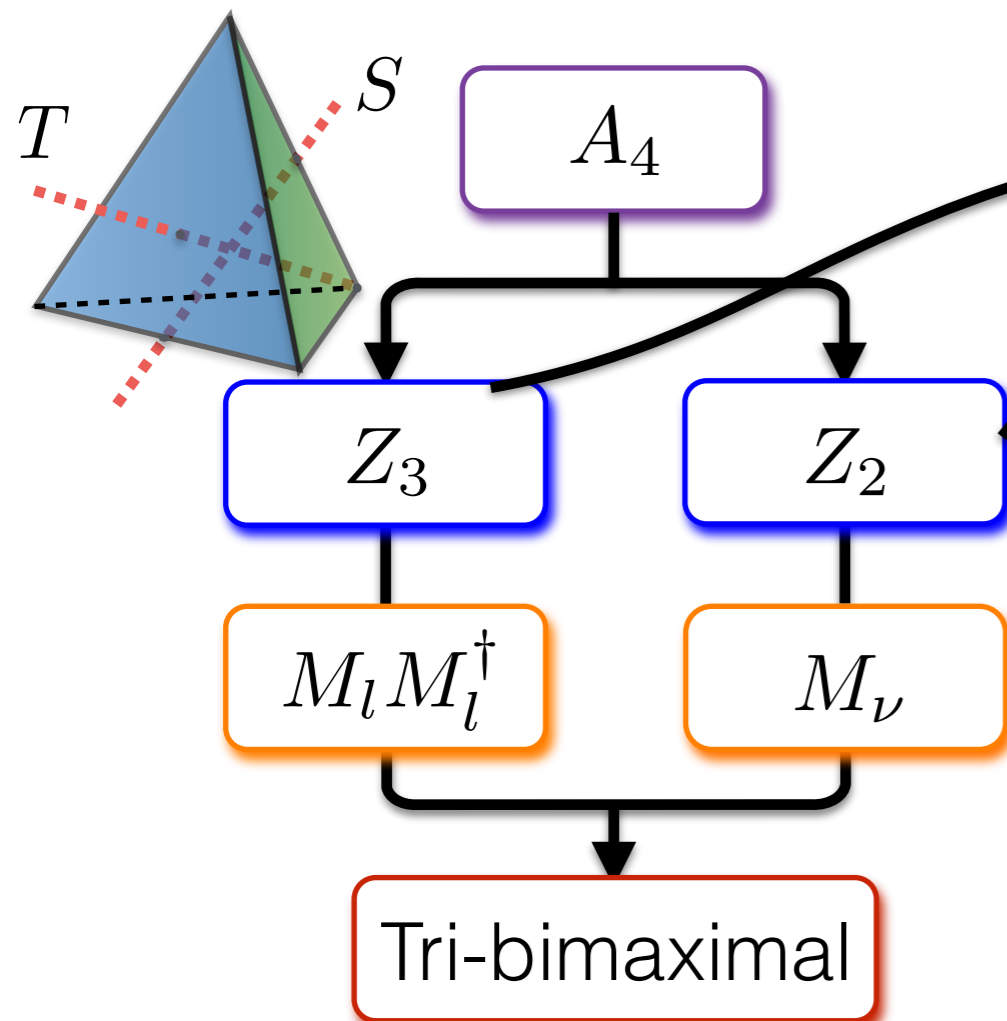


Constant mixing patterns

- Before 2012, Tri-bimaximal (TBM)

$$|U| = \begin{pmatrix} \frac{2}{\sqrt{6}} & \frac{1}{\sqrt{3}} & 0 \\ \frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} \quad \begin{aligned} \sin^2 \theta_{13} &= 0 \\ \sin^2 \theta_{12} &= 1/3 \\ \sin^2 \theta_{23} &= 1/2 \end{aligned}$$

Harrison, Perkins, Scott, 02; Xing, 02



Altarelli, Feruglio, 05; 06.

- After 2012, special corrections are needed:

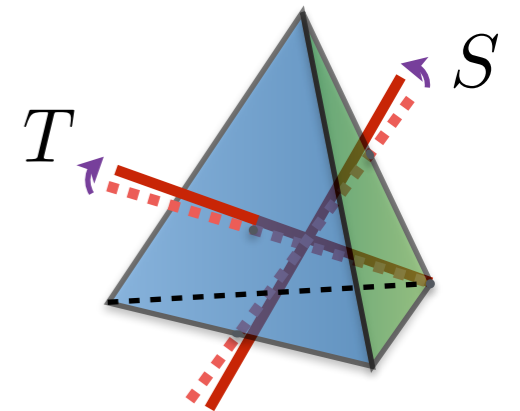
$$U = U_{\text{TBM}} + \delta U$$

Interference between Z_3 and Z_2

Break Z_3 and Z_2

Modify mass structures

Deviation of TBM
sizable θ_{13} and δ



Mixing induced by flavon cross couplings

● An economical model

Pascoli, YLZ, JHEP 1606, 073 (2016)

$$V_{\text{flavon}} = V(\varphi) + V(\chi) + V(\varphi, \chi) + \dots$$

\Downarrow
 \Downarrow
Self Couplings

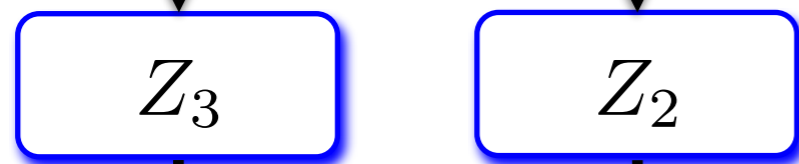
\Downarrow
Cross Couplings

$$V(\varphi) = \mu_\varphi^2 I_1 + g_1 I_1^2 + g_2 I_2$$

$$I_1 = \varphi_1^2 + \varphi_2^2 + \varphi_3^2$$

$$I_2 = \varphi_1^2 \varphi_2^2 + \varphi_2^2 \varphi_3^2 + \varphi_3^2 \varphi_1^2$$

$g_2 < 0$ $g_2 > 0$



Tri-bimaximal

$$\sin \theta_{12} = \frac{1}{\sqrt{3}} (1 - 2|\epsilon_\varphi| \cos \theta_\varphi + 2\epsilon_\chi)$$

$$\sin \theta_{23} = \frac{1}{\sqrt{2}} (1 + |\epsilon_\varphi| \cos \theta_\varphi)$$

$$\sin \theta_{13} = \sqrt{2} |\epsilon_\varphi| \sin \theta_\varphi$$

$$\delta = \begin{cases} 270^\circ - 2|\epsilon_\varphi| \sin \theta_\varphi, & \theta_\varphi > 0, \\ 90^\circ - 2|\epsilon_\varphi| \sin \theta_\varphi, & \theta_\varphi < 0, \end{cases}$$

$$\epsilon_\varphi = |\epsilon_\varphi| e^{i\theta_\varphi}$$

Sum rule $\delta \approx \begin{cases} 270^\circ - \sqrt{2}\theta_{13}, & \theta_\varphi > 0, \\ 90^\circ + \sqrt{2}\theta_{13}, & \theta_\varphi < 0. \end{cases}$

Origin of CP violations

- Explicit CP violations

Complex coefficients

$$\mathcal{L} = |\lambda|e^{i\alpha}\phi^4 + |\lambda|e^{-i\alpha}\phi^{*4} + \dots$$

$$\xrightarrow{CP} \mathcal{L}' = |\lambda|e^{i\alpha}\phi^{*4} + |\lambda|e^{-i\alpha}\phi^4 + \dots \neq \mathcal{L}$$

- Spontaneous CP violations

Real coefficients

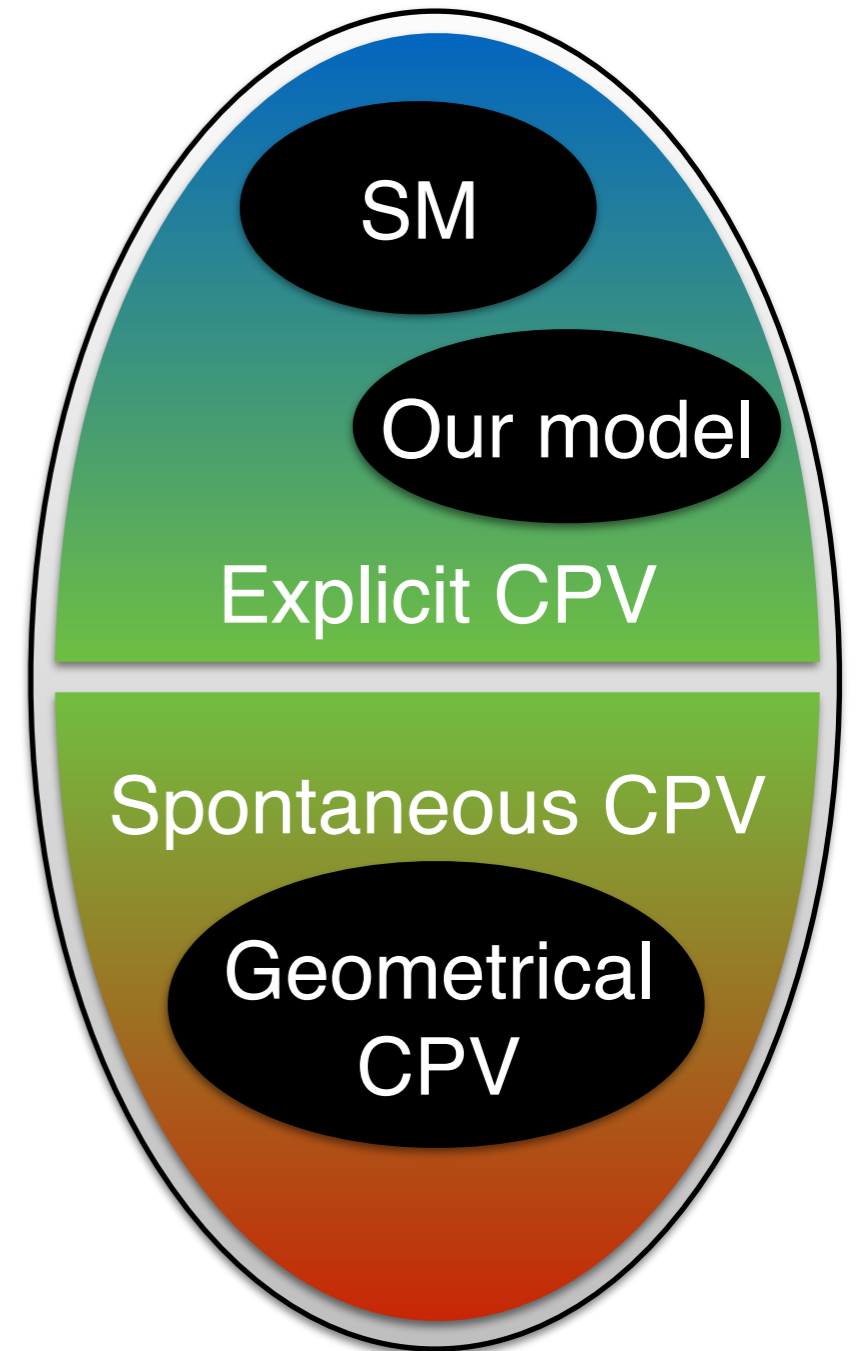
No CP violations at some high scale.

CP violation arises from complex vacuum expectation values (VEV) of some scalars.

$$\langle \phi \rangle = (v_1 e^{i\alpha_1}, v_2 e^{i\alpha_2}, v_3 e^{i\alpha_3})^T$$

Geometrical CP Violation: analytical calculable CP-violating phases

Branco, Gerard, Grimus, 84; de Medeiros Varzielas, Emmanuel-Costa, 11



Generalised CP symmetry

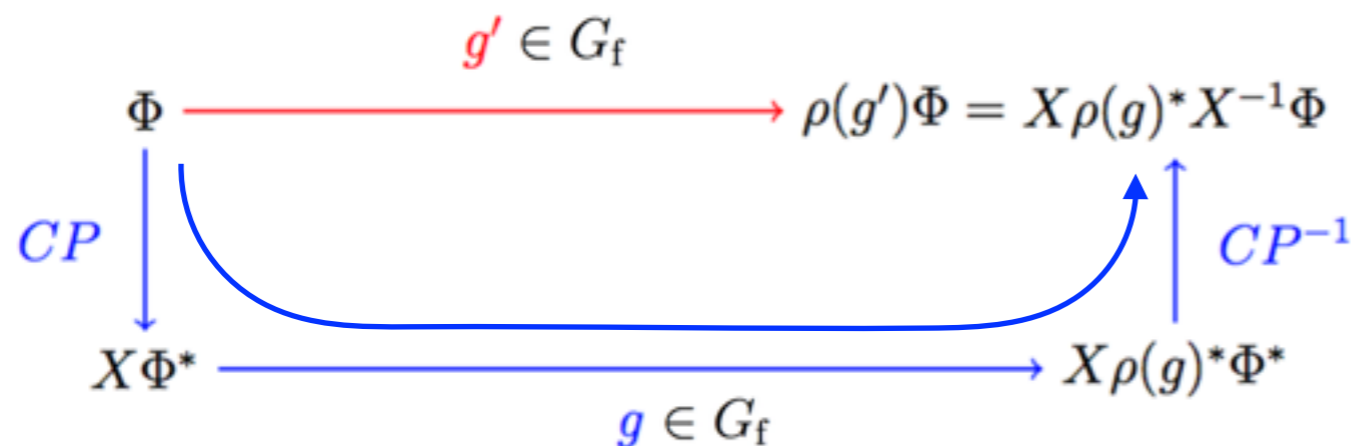
- Generalised CP symmetry (GCP)

$$\begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix} \rightarrow X \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}^* \quad \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \rightarrow X \begin{pmatrix} \nu_e^c \\ \nu_\mu^c \\ \nu_\tau^c \end{pmatrix} \quad \text{with} \quad (x^0, \vec{x}) \rightarrow (x^0, -\vec{x})$$

X is a unitary matrix (may not be an identity matrix).

- Consistency condition

Holthausen, Lindner, Schmidt, 1211.6953



Once the flavour symmetry is fixed, all GCP candidates are determined.

- Powerful prediction

Feruglio, Hagedorn, Ziegler, 1211.5560

Row permutation

Rotation

Column permutation

$$U_{\text{PMNS}} = P_{\alpha\beta\gamma} V R(\theta) K_\nu P_{ijk}$$

Constant matrix

$\text{diag}\{1, \pm 1, \pm 1\}$

Predictions of generalised CP symmetry

- CP conserving

$$X = \mathbf{1}_{3 \times 3} \quad \delta = 0, 180^\circ.$$

- CP violations from generalised CP symmetry

μ - τ reflection symmetry \Rightarrow Maximal CP violations

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \rightarrow \begin{pmatrix} \nu_e^c \\ \nu_\tau^c \\ \nu_\mu^c \end{pmatrix} \quad M_\nu = \begin{pmatrix} A & B & B^* \\ B & C & D \\ B^* & D & C^* \end{pmatrix} \quad \begin{array}{l} \text{Babu, Ma, Valle, 02;} \\ \text{Harrison, Scott, 02.} \end{array}$$

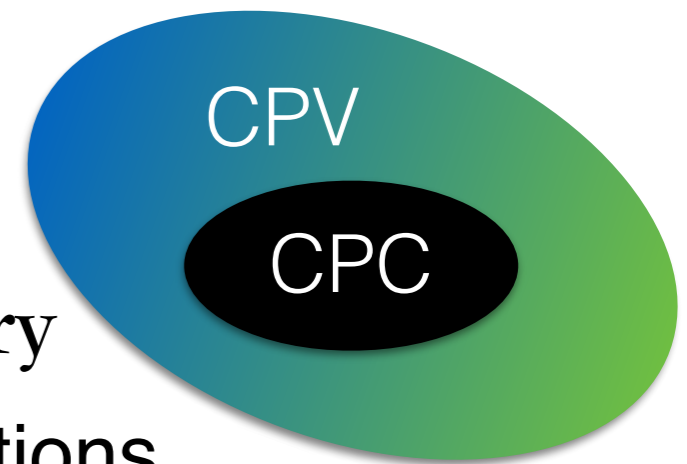
$$\Rightarrow \theta_{23} = 45^\circ, \delta = 90^\circ, 270^\circ.$$

μ - τ reflection symmetry is predicted by A_4 , and any groups containing A_4 , e.g. S_4 , A_5 , ...

Feruglio, Hagedorn, Ziegler, 1211.5560; 1303.7178; Girardi, Meroni, Petcov, Spinrath, 1312.1966; Li, Ding, 1312.4401; Ballett, Pascoli, Turner, 1503.07543..... For a review, see Xing, Zhao, 1512.04207.

Generalised CP symmetries also provide other possibilities that δ does not take specific values, e.g. $\Delta(48)$, $\Delta(96)$, and some other groups in the $\Delta(3n^2)$, $\Delta(6n^2)$ series.

Ding, YLZ, 1304.2645; 1312.5222; Ding, King, 1403.5846; Hagedorn, Meroni and Molinaro, 1408.7118; Ding, King and Neder, 1409.8005; Ding, King, 1510.03188.



Lepton asymmetry in the early Universe

- Baryon asymmetry in our observed Universe

$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} \approx 6.2 \times 10^{-10}$$

- Baryogenesis via classical leptogenesis (in type I seesaw)

N are Majorana fermions

Complex Yukawa couplings

Decay of lightest N

Lepton asymmetry

Sphaleron $(\Delta B = \Delta L)$

Baryon asymmetry

- Lepton asymmetry $\Delta f_{l_\alpha} \equiv f_{l_\alpha} - f_{\bar{l}_\alpha}$ in classical leptogenesis

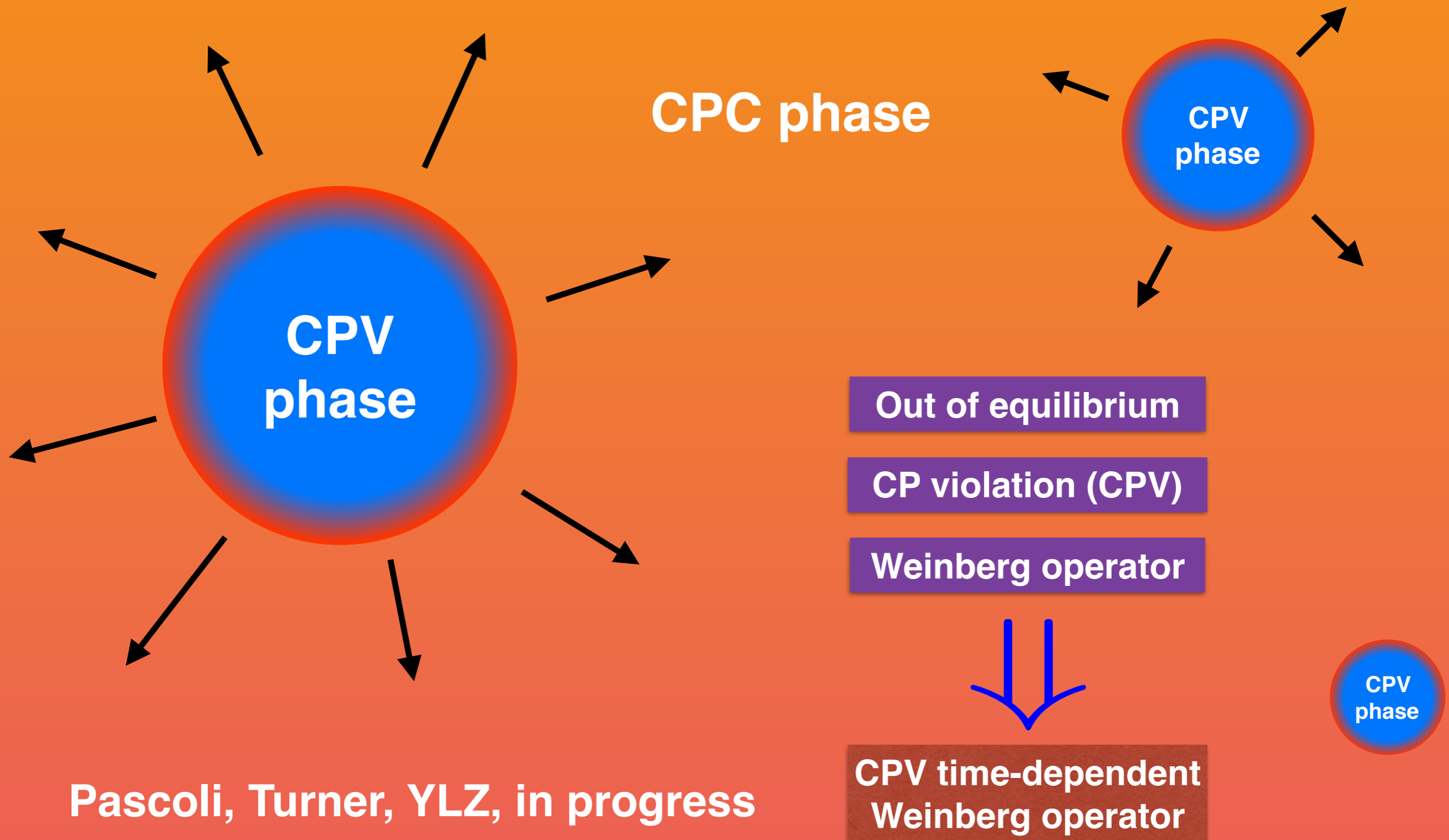
$$\Delta f_{l_\alpha} \propto \text{Im} \left\{ \begin{array}{c} \text{Diagram 1: } N_1 \text{ decaying to } l_\alpha \text{ and } H \\ \text{Diagram 2: } N_1 \text{ decaying to } l_\alpha \text{ and } N_j \text{, then } N_j \text{ decaying to } l_\beta \text{ and } H \\ \text{Diagram 3: } N_1 \text{ decaying to } N_j \text{ and } H \text{, then } N_j \text{ decaying to } l_\beta \text{ and } l_\alpha \end{array} \right\}$$

$$\propto \text{Im} \{ Y_{\nu\alpha 1}^* (Y_\nu^\dagger Y_\nu)_{1j} Y_{\nu\alpha j} \} \quad [\text{Fukugita, Yanagida, 1986}]$$

Yukawa couplings between ℓ and N Y_ν cannot be real or proportional to a unitary matrix (ignoring flavour effect).

Alternative way for lepton asymmetry: CPPT

CP phase transition (CPPT) below the seesaw scale



Alternative way for lepton asymmetry: CPPT

- Time-dependent coupling in the Weinberg operator

$$\mathcal{L}_{d=5} = \frac{M_{\nu\alpha\beta}^*(t)}{v_H^2} \ell_{\alpha L}^c \ell_{\beta L}^c H^{*2} + \frac{M_{\nu\alpha\beta}(t)}{v_H^2} \ell_{\alpha L} \ell_{\beta L} H^2$$

- Lepton asymmetry

$$\Delta f_{l_\alpha} \propto \text{Im} \left\{ \begin{array}{c} \begin{array}{c} \text{Diagram 1: } l_\alpha^+ \text{ (red arrow left), } l_\beta^- \text{ (purple arrow right), } M_{\nu\alpha\beta}^*(t_1) \text{ (green dashed arrow down), } H \text{ (blue dashed arrow up)} \end{array} \\ \times \\ \begin{array}{c} \text{Diagram 2: } l_\alpha^- \text{ (red arrow right), } l_\beta^+ \text{ (purple arrow left), } M_{\nu\alpha\beta}(t_2) \text{ (green dashed arrow up), } H \text{ (blue dashed arrow down)} \end{array} \end{array} \right\}$$

- Closed Time Path (CTP) formalism

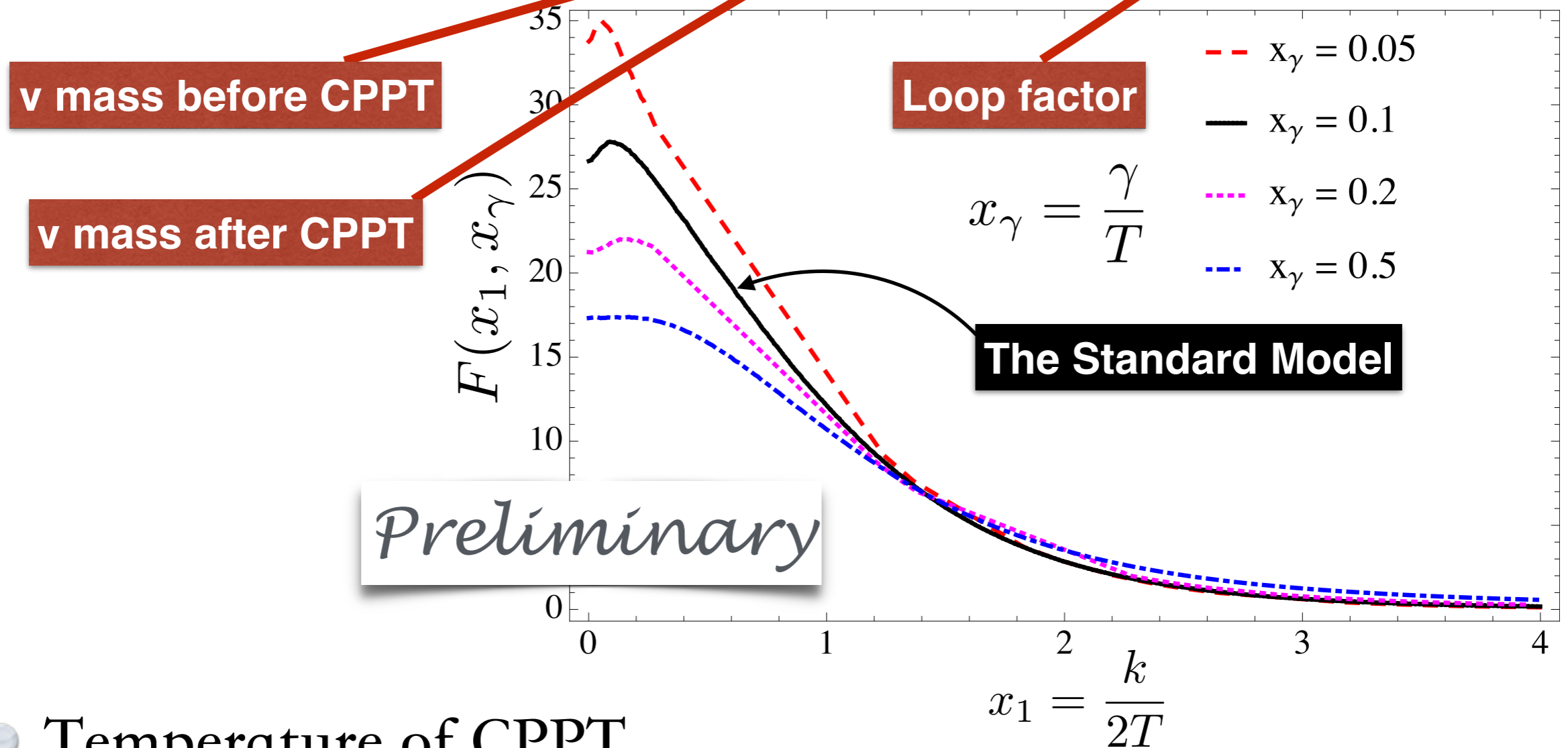
$$\Delta f_{l_\alpha} \propto \int_{t_i}^{t_f} dt_1 \int_{t_i}^{t_f} dt_2 \begin{array}{c} \text{Diagram: Closed Time Path (CTP) loop with } l_\alpha \text{ (red), } l_\beta \text{ (purple), } M_{\nu\alpha\beta}^*(t_1) \text{ (green), } M_{\nu\alpha\beta}(t_2) \text{ (green), and } H \text{ (blue dashed) paths.} \end{array}$$

Alternative way for lepton asymmetry: CPPT

- Lepton asymmetry

Pascoli, Turner, YLZ, in progress

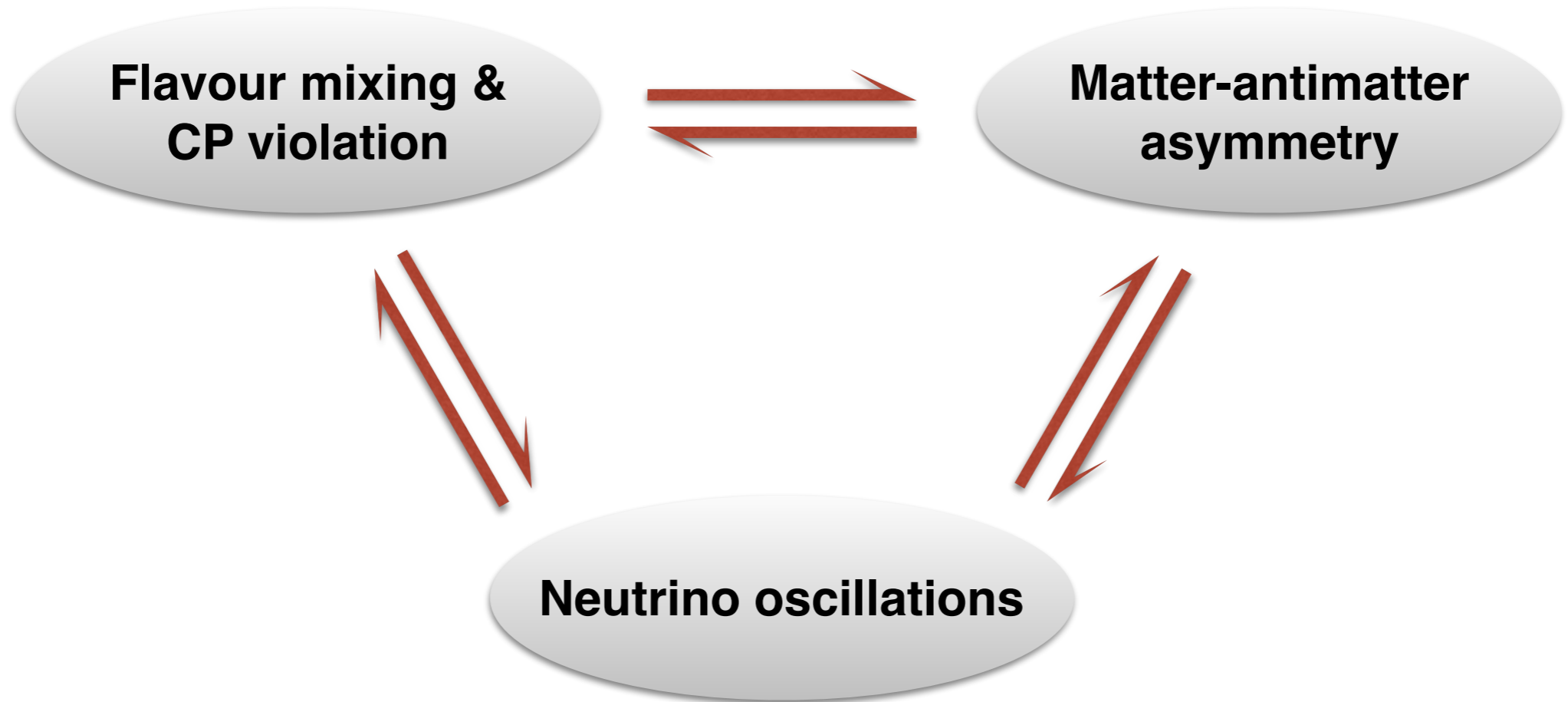
$$\Delta f_\ell \equiv \sum_{\alpha=e,\mu,\tau} \Delta f_{\ell_\alpha}(k) = \text{Im}\{\text{tr}[M_\nu^0 M_\nu^*]\} \frac{3T^2}{(2\pi)^4 v_H^4} F\left(\frac{k}{2T}, \frac{\gamma}{T}\right)$$



- Temperature of CPPT

$$T \sim 10 \sqrt{\Delta f_\ell} \frac{v_H^2}{m_\nu}$$

$$\Delta f_\ell > \eta_B \quad m_\nu = 0.1 \text{eV} \quad \Rightarrow \quad T > 10^{11} \text{ GeV}$$



Thank you very much